

2445-07

Advanced Workshop on Nanomechanics

9 - 13 September 2013

**Nanomechanics:
a brief overview**

Florian Marquardt
Erlangen (Germany)

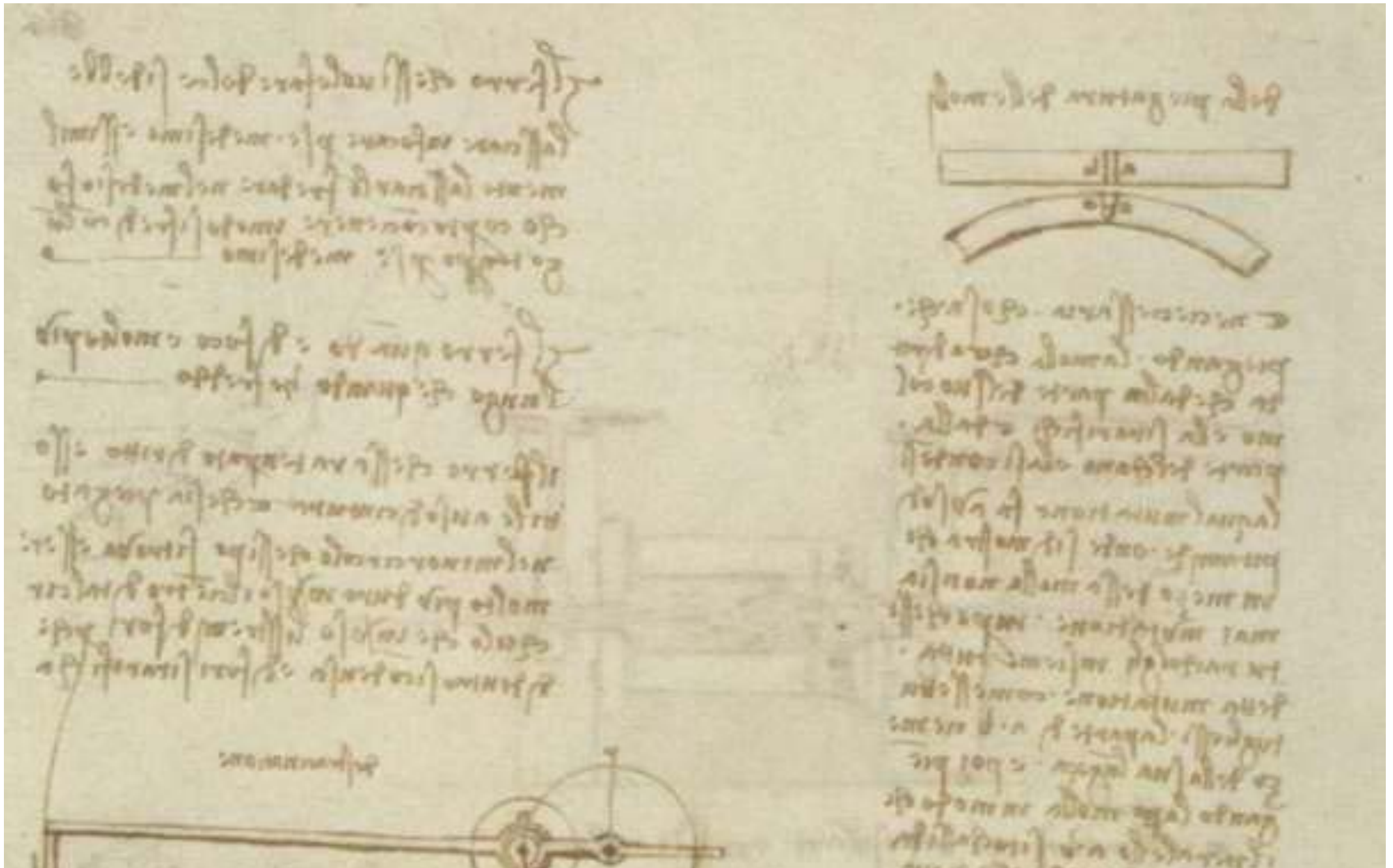
Nanomechanics: a brief overview

Frontiers of Nanomechanics / Trieste 2013

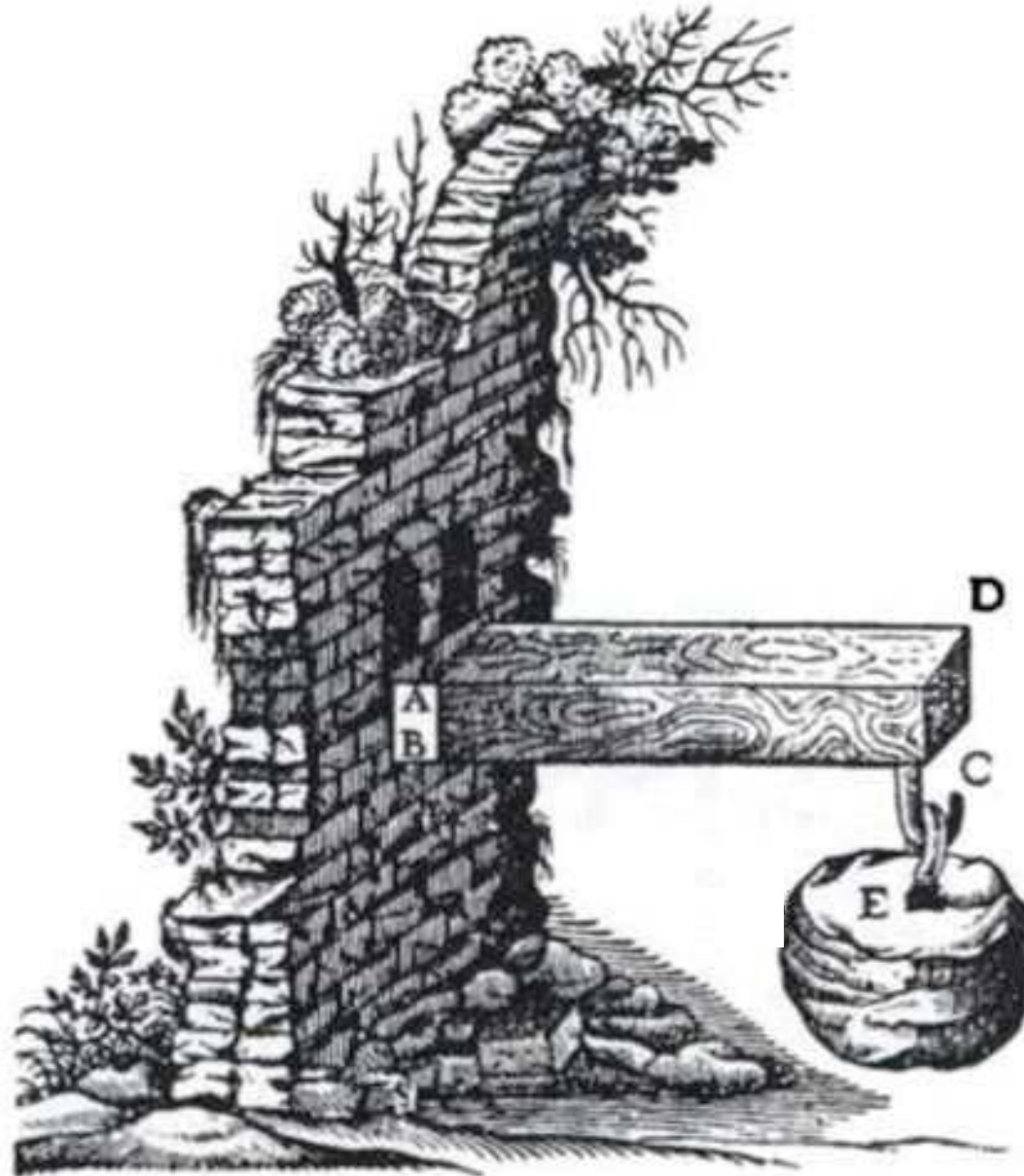
**Florian Marquardt,
Erlangen (Germany)**

**Of bending
(nano-)beams**

Leonardo da Vinci 1493



Galileo Galilei 1638

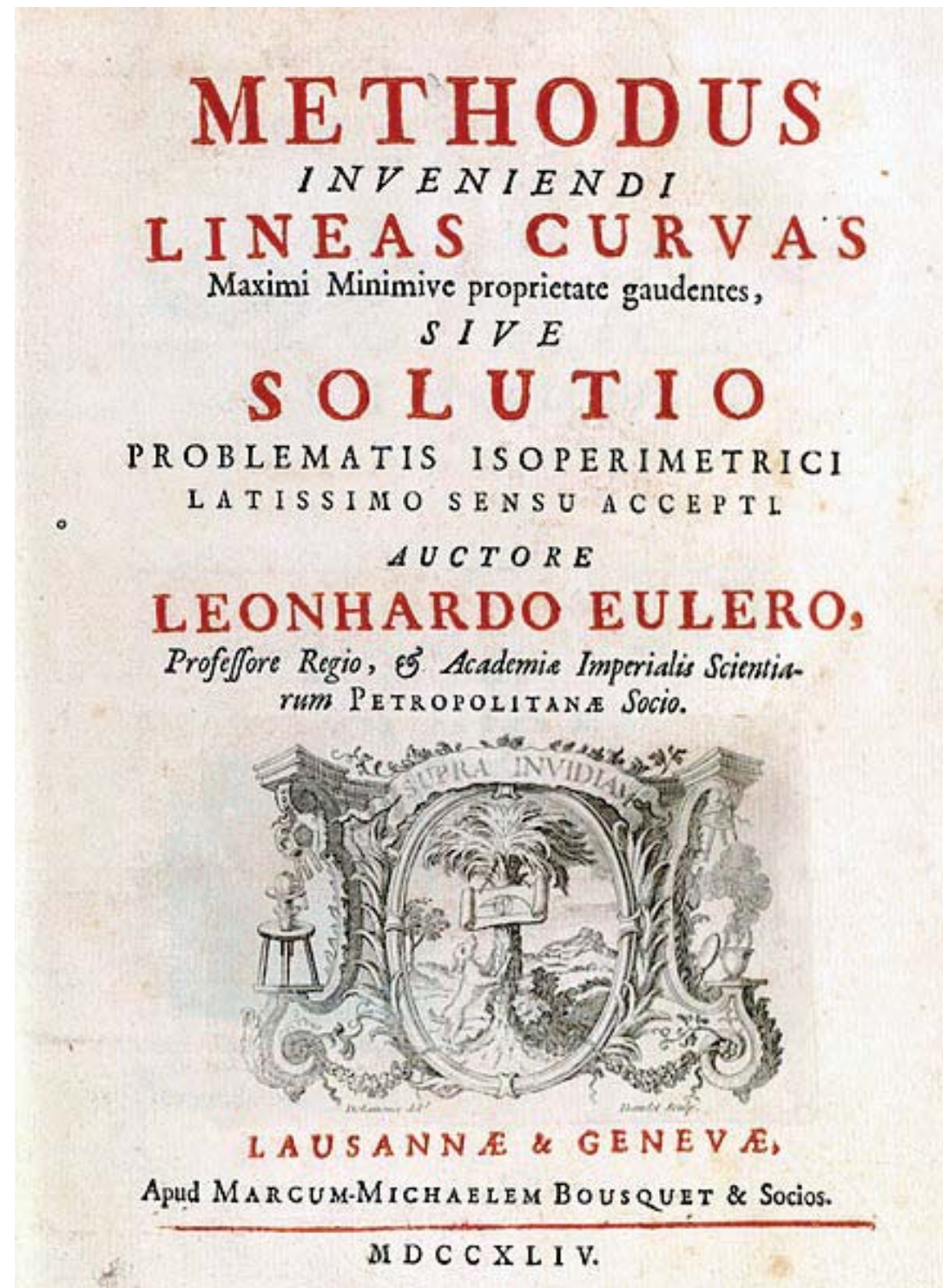


Daniel Bernoulli & Leonhard Euler 1744

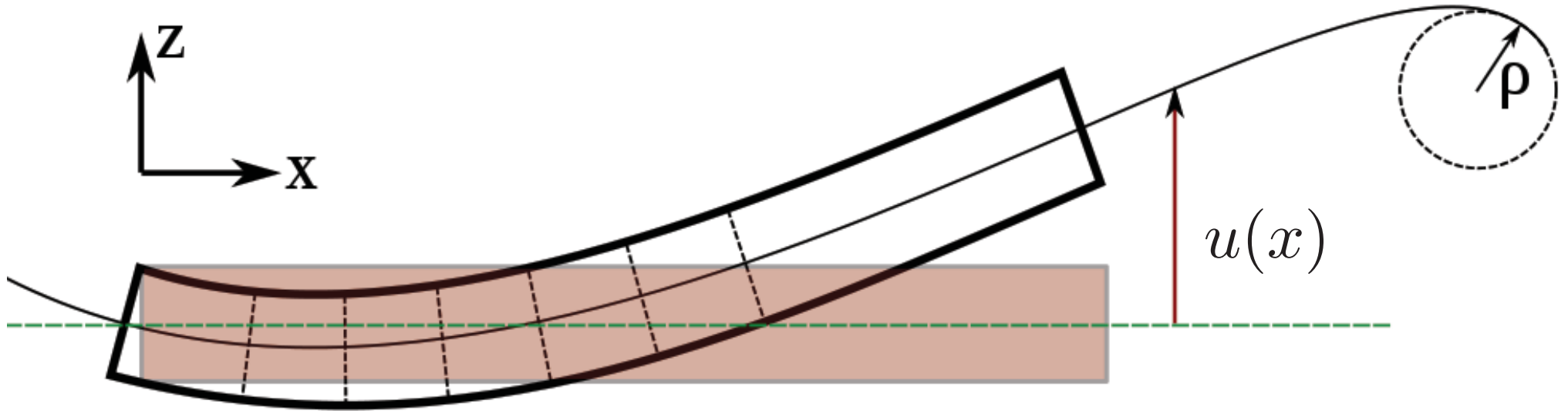


Euler

Elasticity theory & energy approach



Euler-Bernoulli theory for beam bending



energy: bending energy stretching energy

$$U \approx \int_0^L dx \left[\frac{EI}{2} (u''(x))^2 + \frac{F}{2} (u'(x))^2 \right]$$

E : elastic modulus

I : moment of inertia

F : applied force

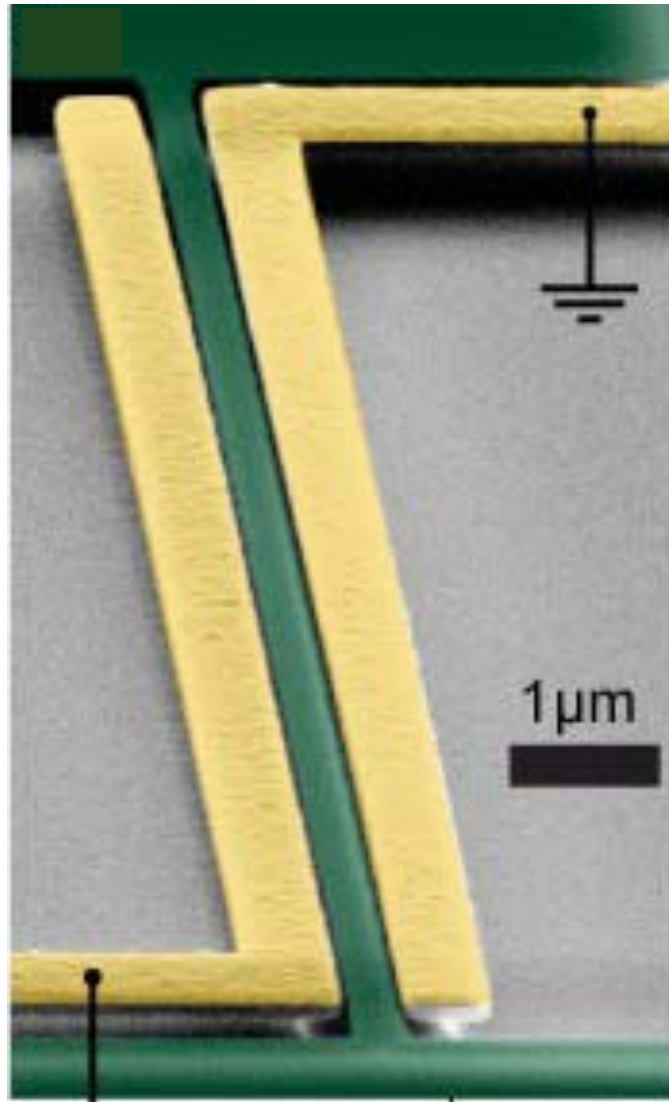
(doubly clamped beam)

$$I = \int dz dy z^2$$



Elasticity theory still works well on the nanometer scale!

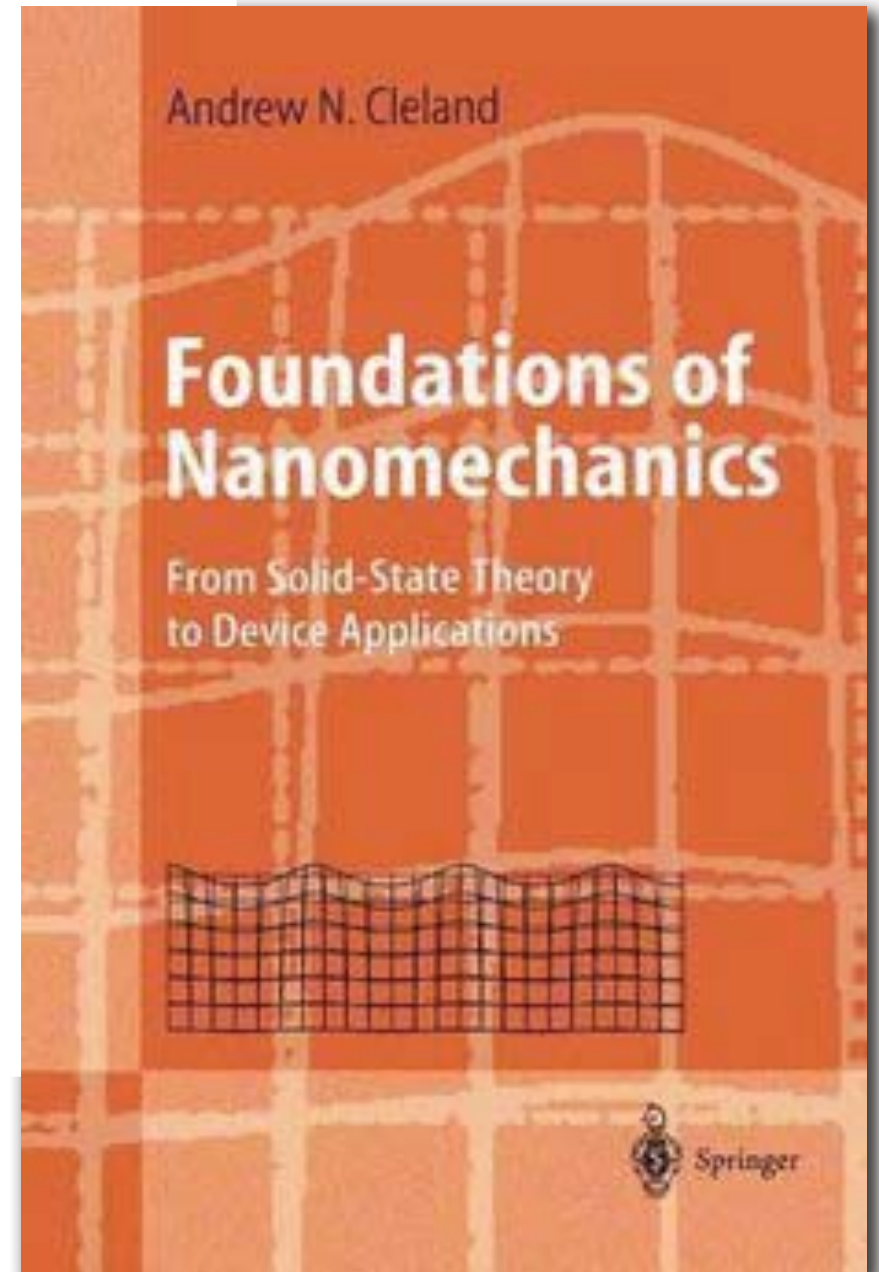
1,000,000 times
smaller!



(Weig)

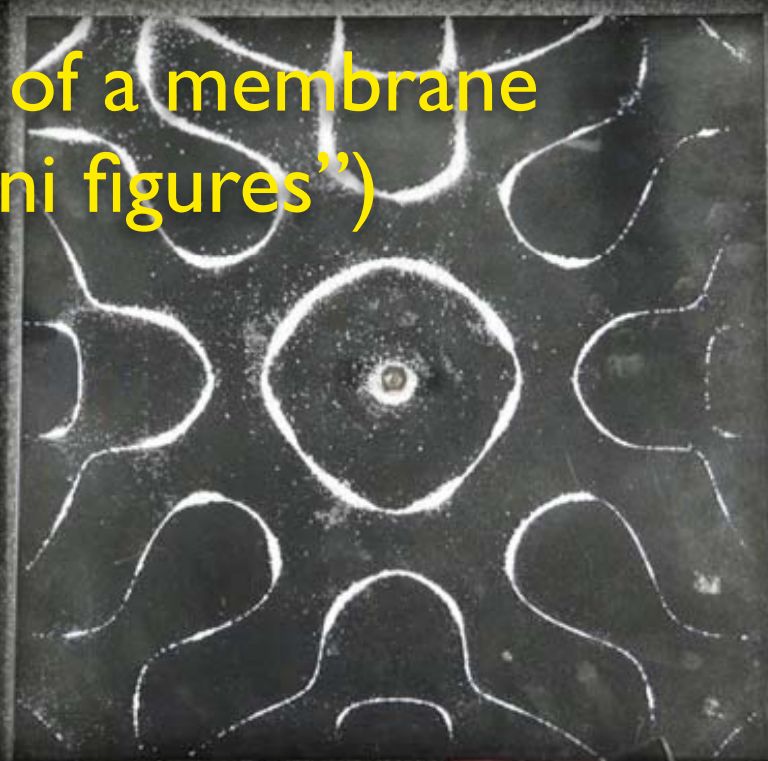
Elasticity theory still works well on the nanometer scale!

1,000,000 times
smaller!



Mechanical vibrations

Eigenmodes of a membrane ("Chladni figures")



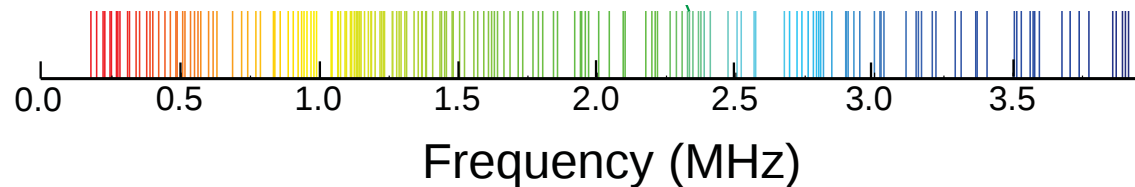
Small vibrations of any mechanical structure described by:

Eigenfrequencies Eigenmodes

$$\Omega_n, \vec{u}_n(\vec{r})$$

deflection from equilibrium

(Yamaguchi)

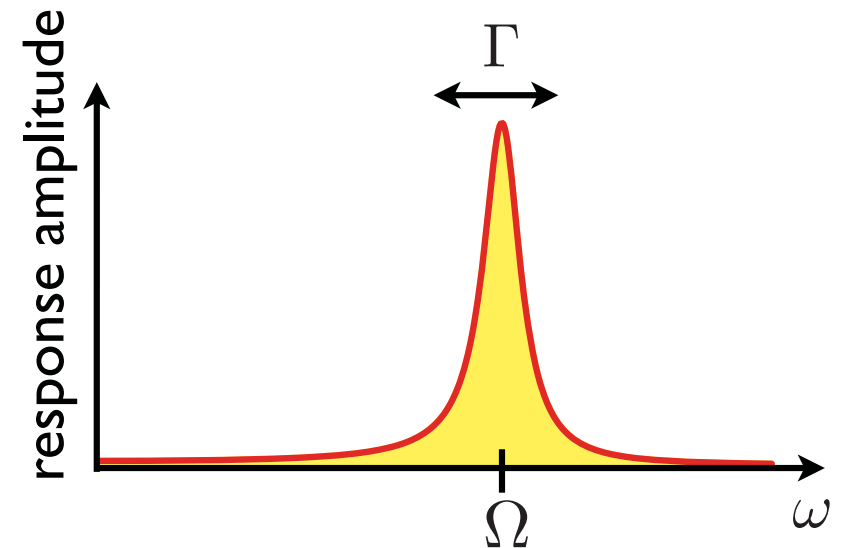


Linear superposition of vibrations

$$\text{displacement field } \vec{u}(\vec{r}, t) = \sum_n x_n(t) \vec{u}_n(\vec{r})$$

Each eigenmode is a harmonic oscillator

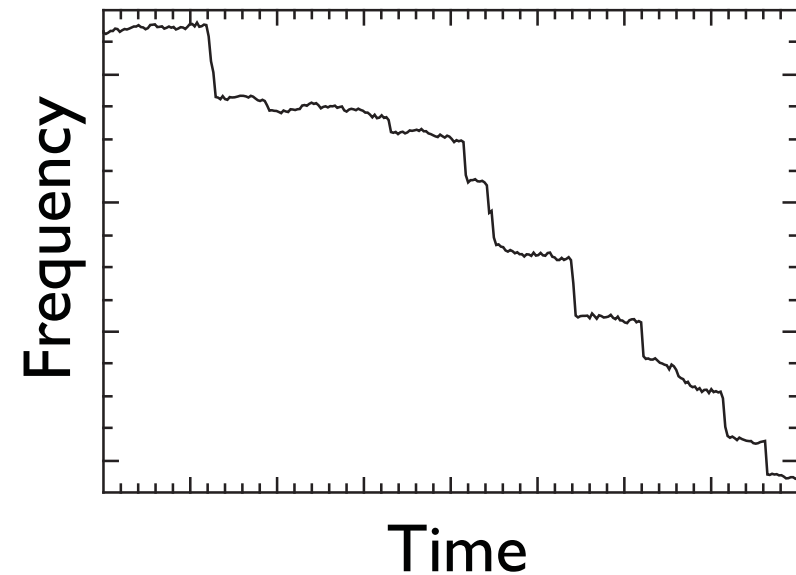
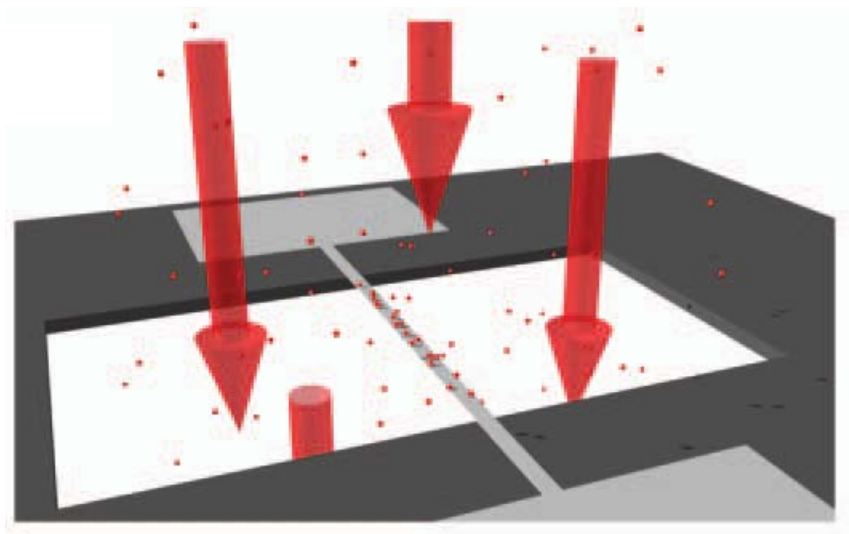
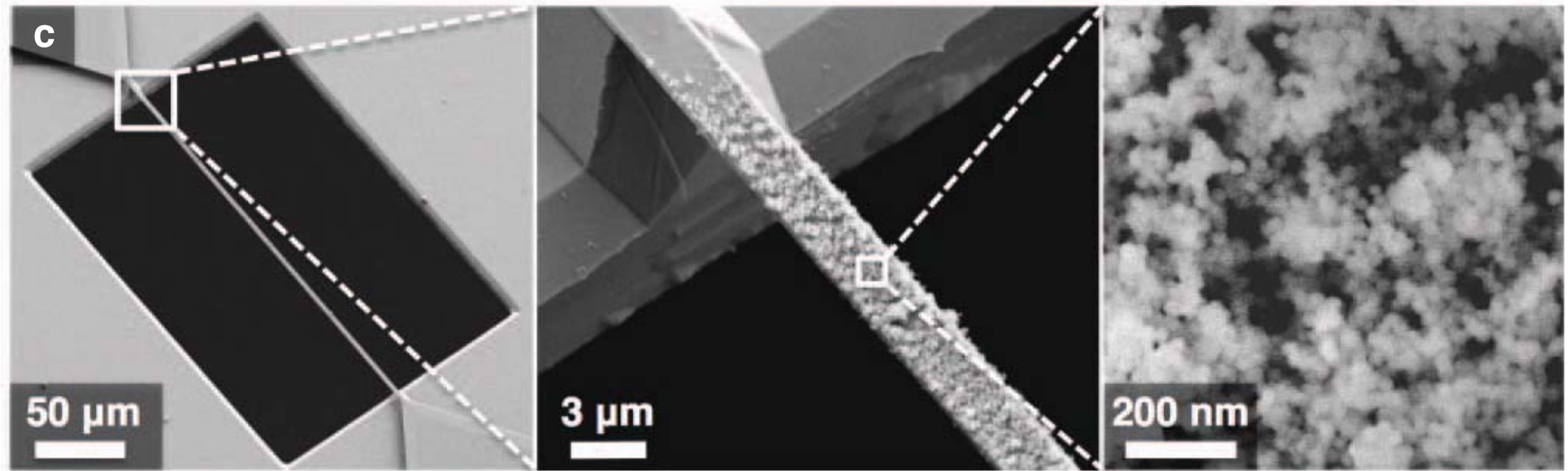
$$m \frac{d^2 x}{dt^2} = \underbrace{-m\Omega^2 x}_{\text{restoring force (linear)}} - \underbrace{m\Gamma \frac{dx}{dt}}_{\text{damping}} + \underbrace{F(t)}_{\text{external force}}$$



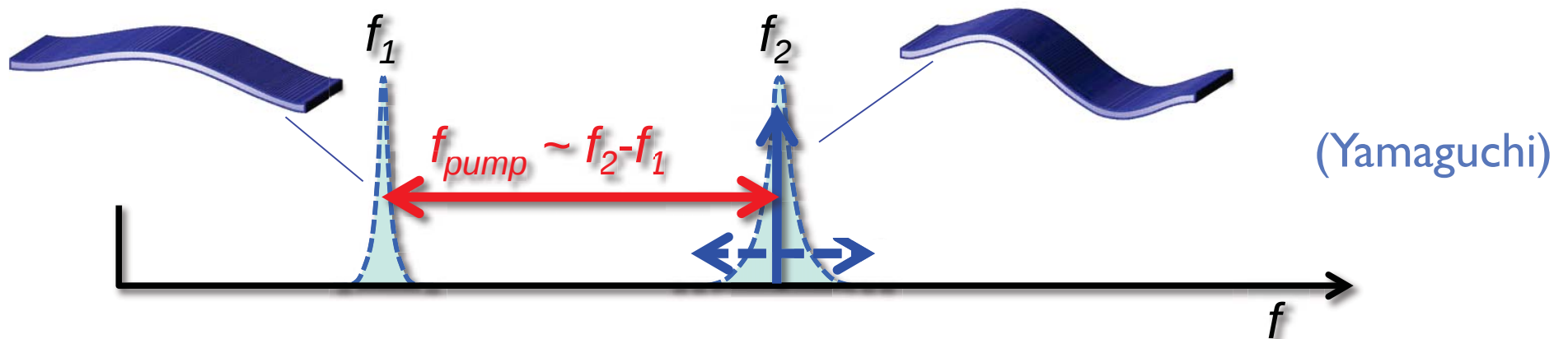
Mass sensing via a shift of the eigenfrequency

$$\Omega = \sqrt{\frac{k}{m}}$$

Silvan Schmid (Friday)

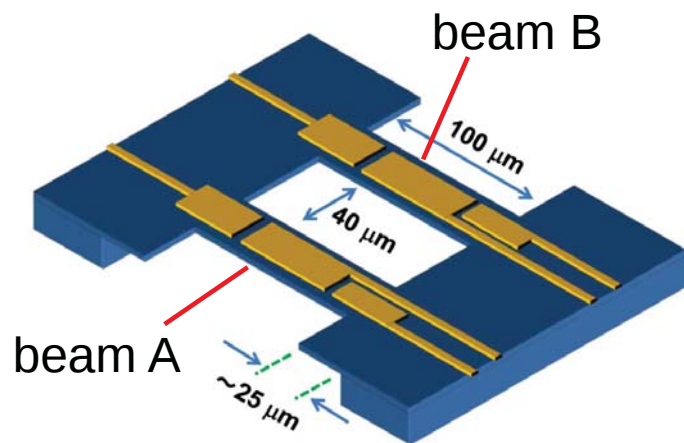


Usually focus on **one** mechanical mode
...but interesting effects for multiple coupled modes!



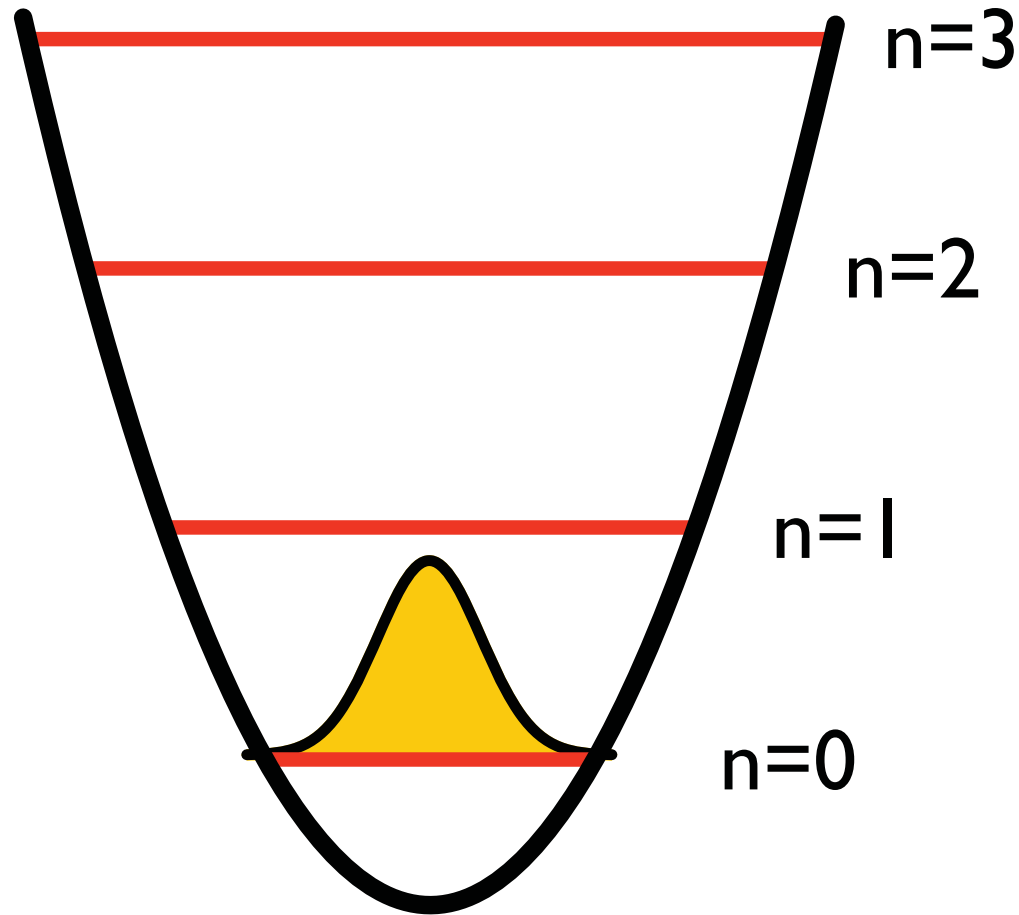
Parametric drive of coupling: $F_1 = K \cos(2\pi f_{\text{pump}} t) x_2 + \dots$

...leads to “Rabi oscillations”
of mechanical energy between
the two modes



Hiroshi Yamaguchi (Thursday)

Quantum-mechanical mechanical harmonic oscillator



$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\Omega^2}{2}\hat{x}^2$$

$$E_n = \hbar\Omega\left(n + \frac{1}{2}\right)$$

phonon number

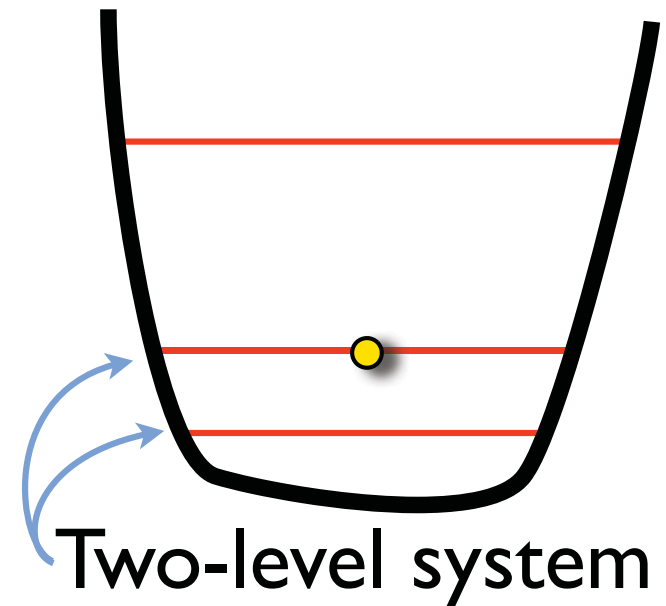
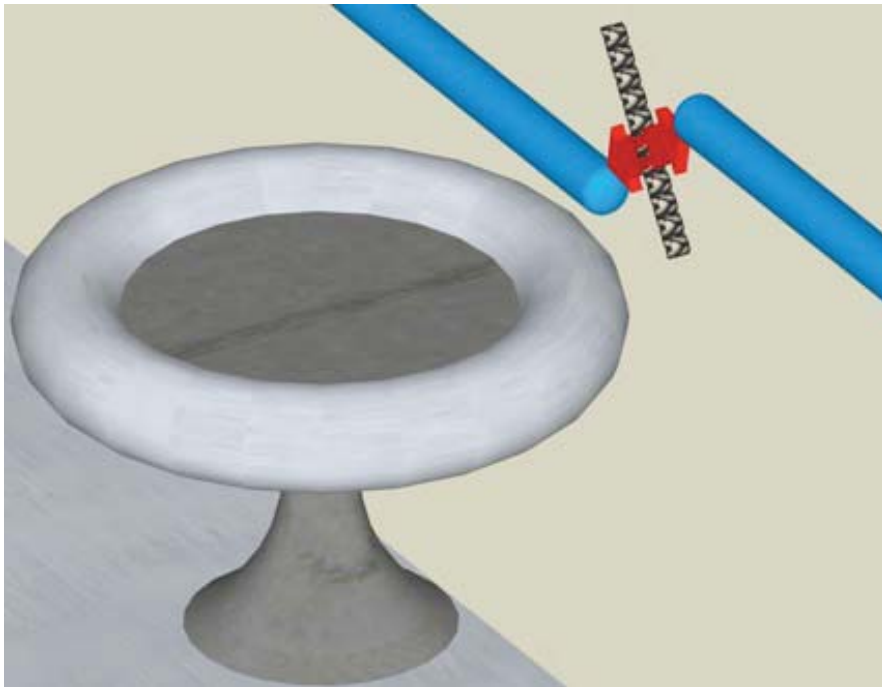
$$\hat{H} = \hbar\Omega\left(\hat{b}^\dagger\hat{b} + \frac{1}{2}\right)$$

$$\hat{x} = x_{\text{ZPF}}(\hat{b} + \hat{b}^\dagger)$$

mechanical zero-point fluctuations (ground state width) $x_{\text{ZPF}} = \sqrt{\hbar/2m\Omega}$

Usually: mechanical modes are **harmonic** oscillators
(typically very good approximation for small vibrations,
e.g. near the single-phonon level)

But: Potential use as qubits if anharmonicity (nonlinearity)
can be made strong enough!



Mechanical damping

Mechanical damping

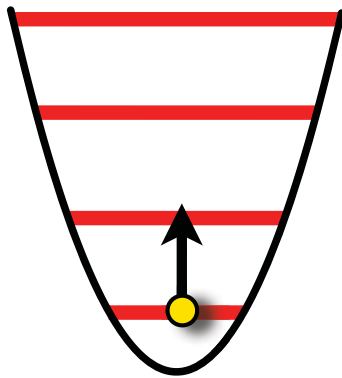
Damping rate

Γ
(different for
each mode)

Quality factor

$$Q = \frac{\Omega}{\Gamma}$$

~ number of oscillations
during damping time
e.g. 10^5



Excitation of ground state
(due to thermal fluctuations)

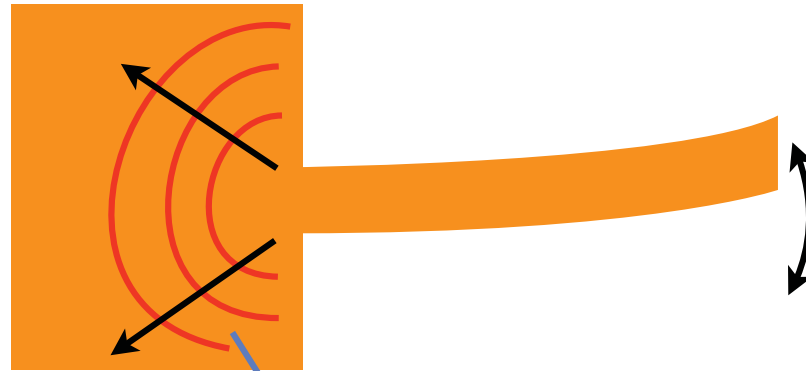
$$\Gamma_{0 \rightarrow 1} = \Gamma \bar{n}_{\text{th}} \approx \frac{k_B T}{\hbar Q}$$

thermal occupation

...sets limits on quantum coherence!

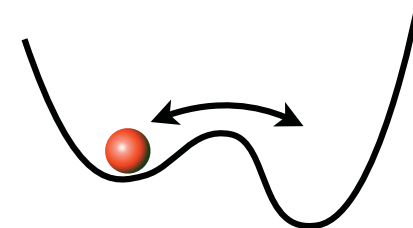
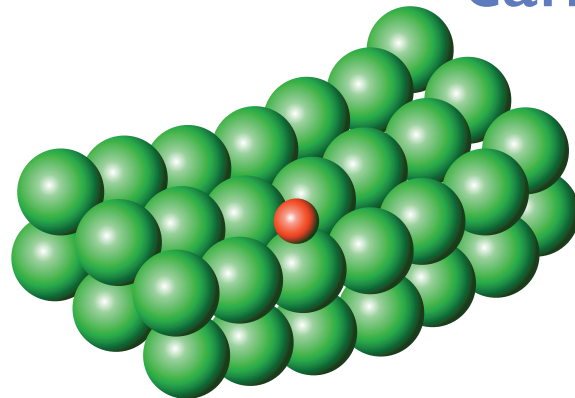
Common sources of mechanical damping

“Clamping losses”: Beam attached to structure



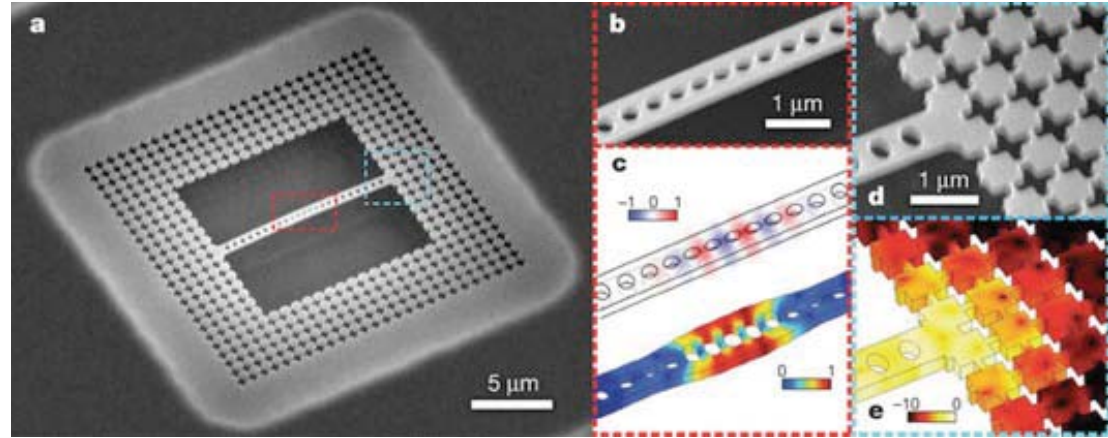
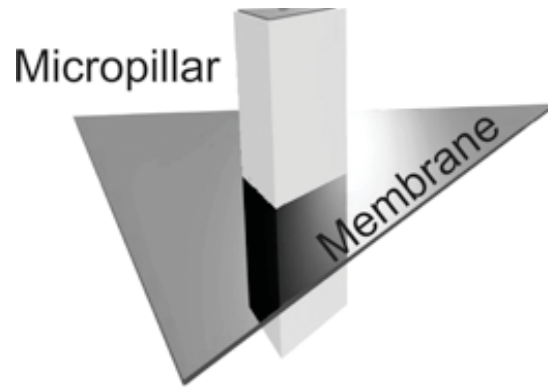
Sound waves radiated into structure

Structural losses; e.g. due to two-level fluctuators
can be excited by vibrations



How to prevent...

“Clamping losses”: Engineer mode shape or surroundings



Antisymmetric mode (LKB group)

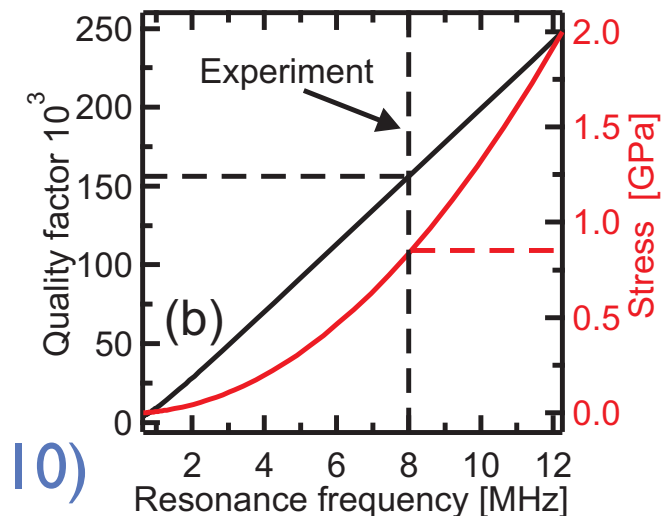
Samuel Deleglise (Thursday)

“Phonon shield” (Painter group)

Amir Safavi-Naeini (Wednesday)

Structural losses: increase tension (oscillation energy)

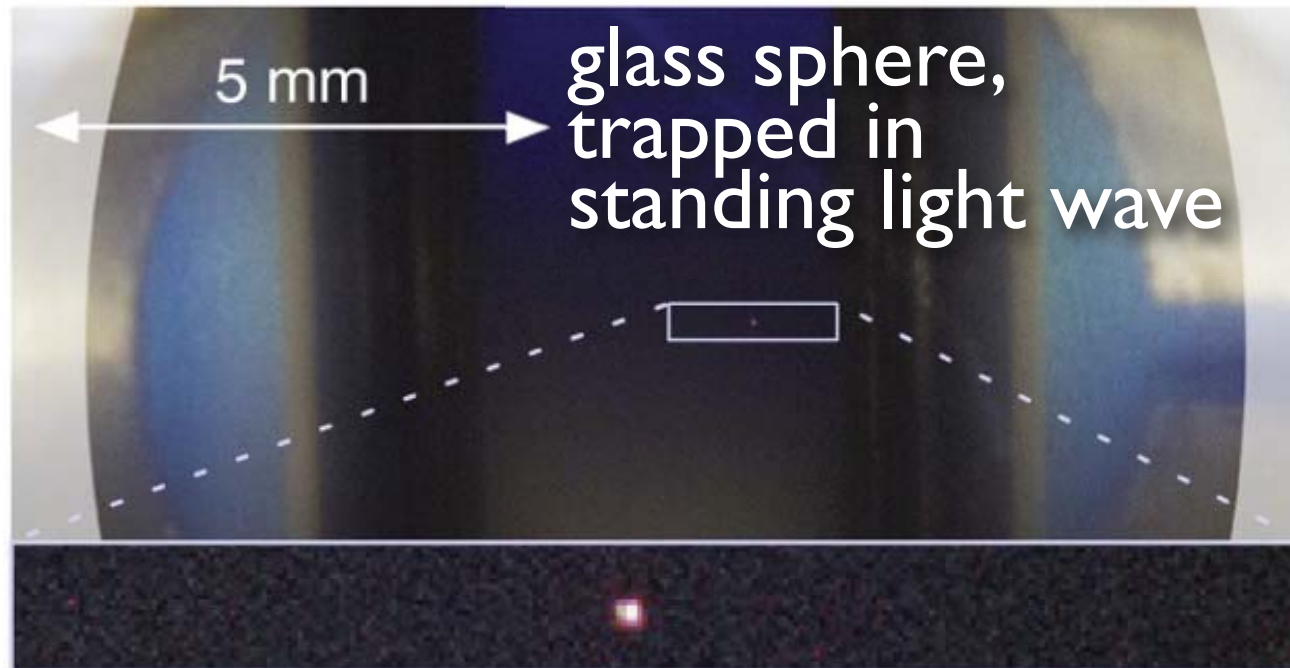
(Unterreithmeier, Faust, Kotthaus, 2010)



How to prevent...

“Clamping losses”: levitate mechanical object!

Nicolai Kiesel (Thursday)



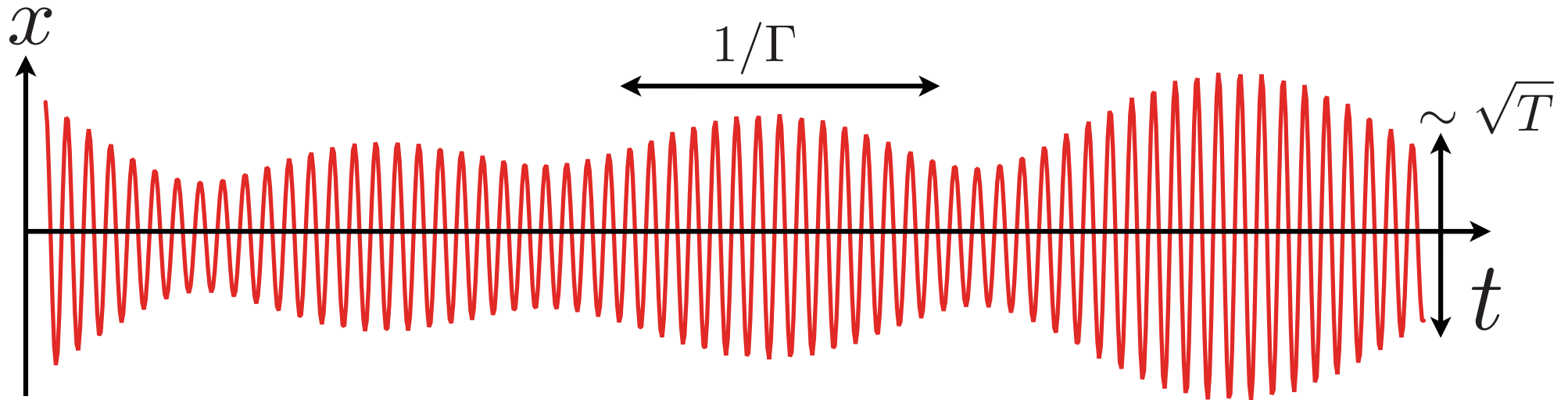
Levitate drop of
superfluid helium
(surface waves!)



Jack Harris (Tuesday)

The mechanical fluctuation spectrum

Thermal fluctuations of a harmonic oscillator



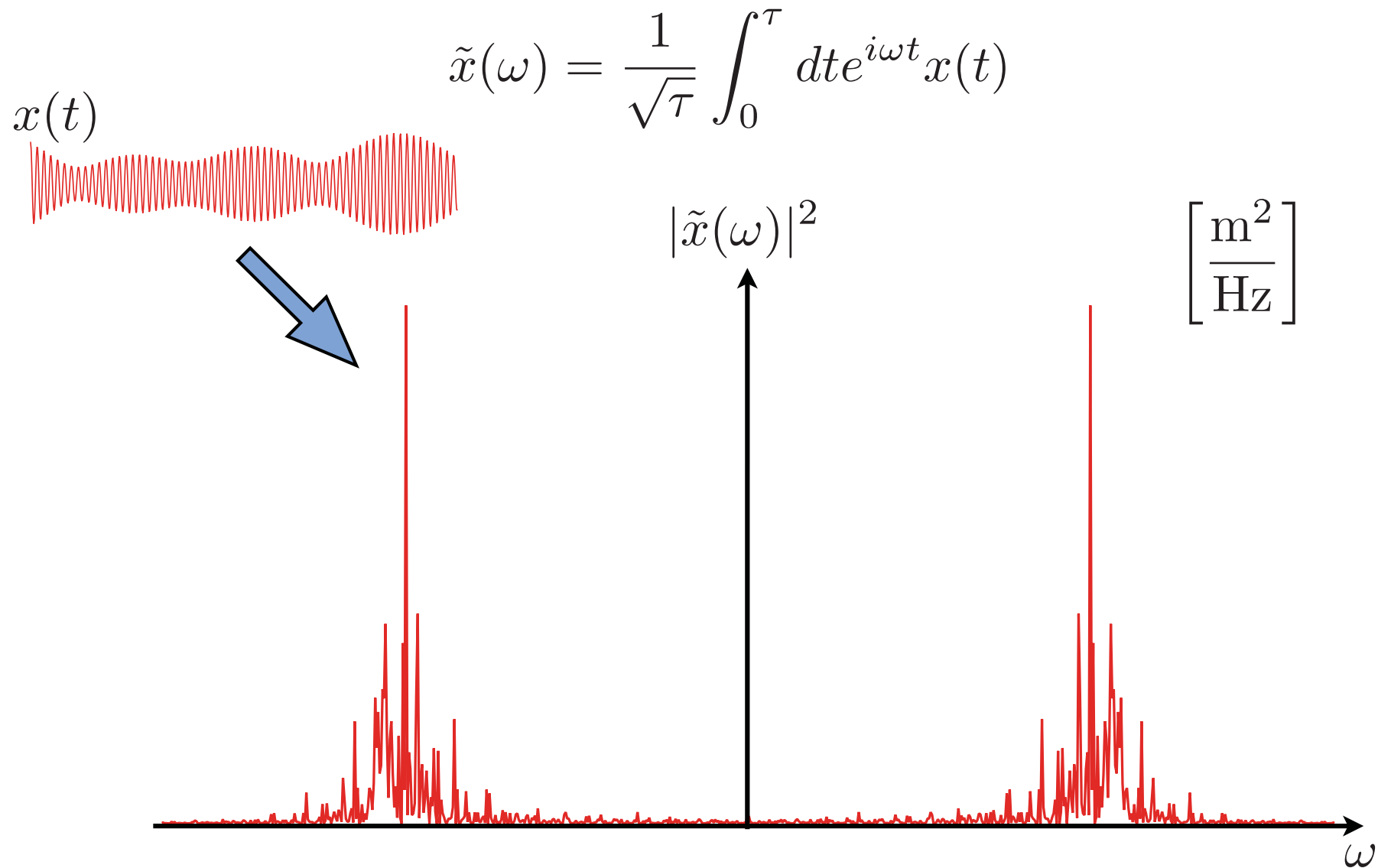
Classical equipartition theorem:

$$\frac{m\omega_M^2}{2} \langle x^2 \rangle = \frac{k_B T}{2} \Rightarrow \langle x^2 \rangle = \frac{k_B T}{m\omega_M^2} \text{ extract temperature!}$$

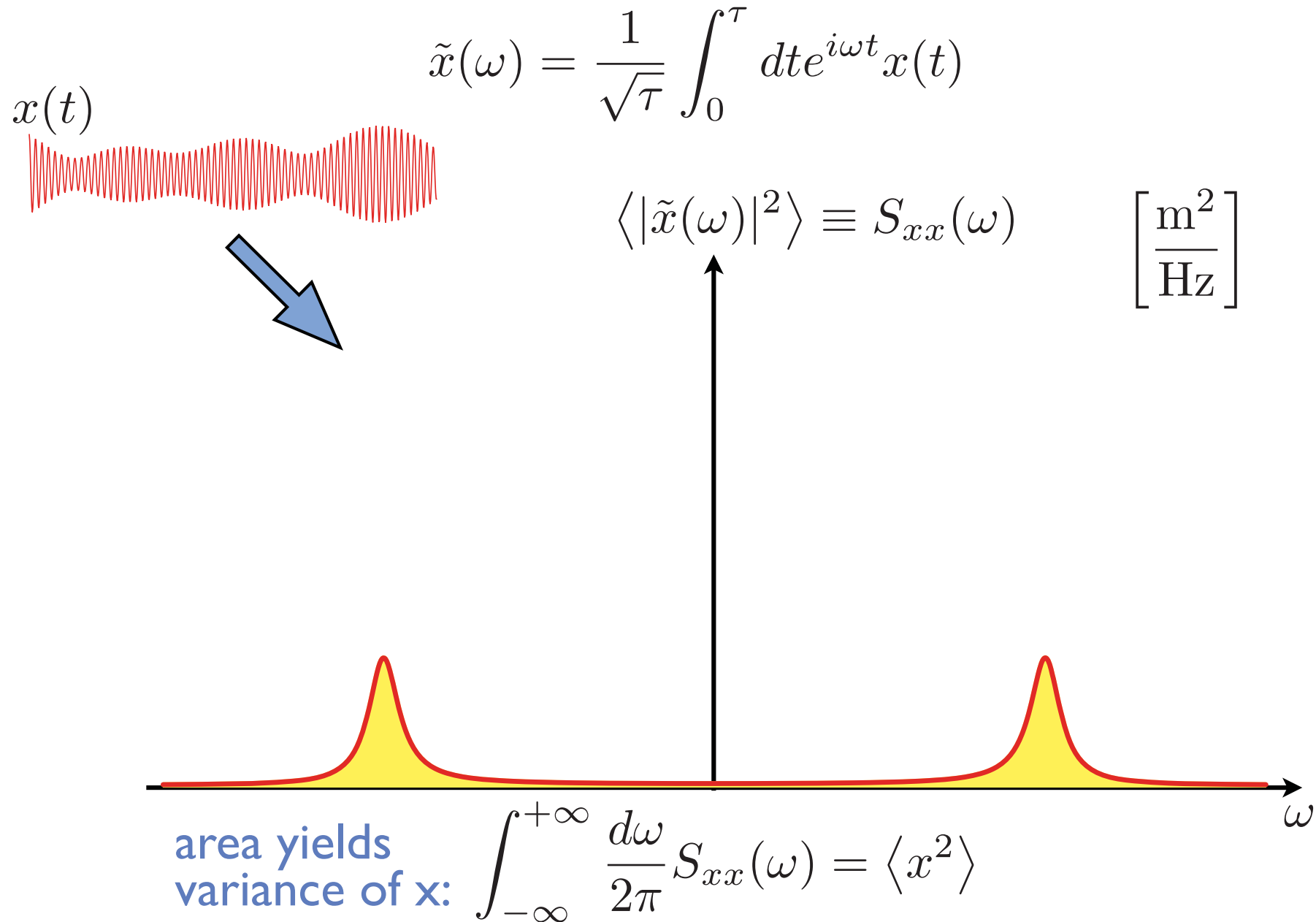
Possibilities:

- Direct time-resolved detection
- Analyze **fluctuation spectrum of x**

The fluctuation spectrum



The fluctuation spectrum



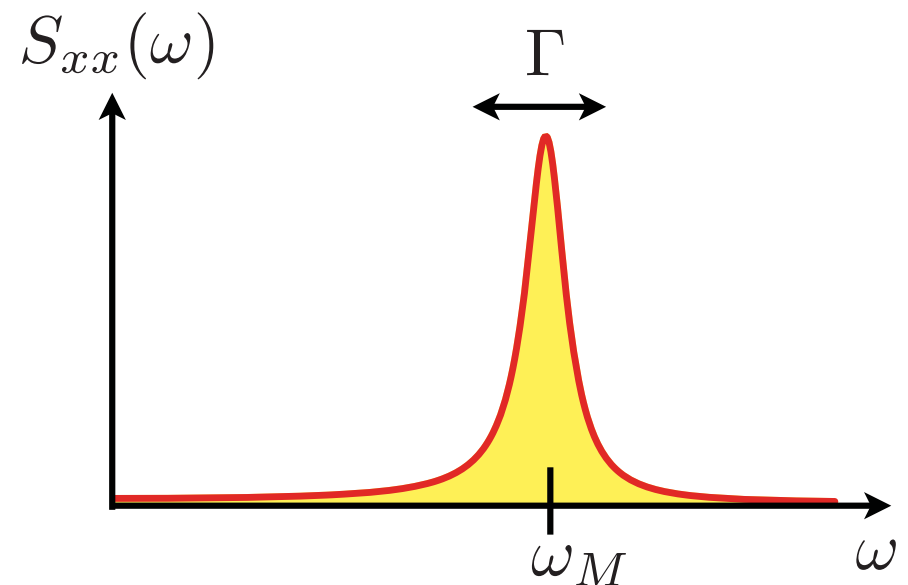
Fluctuation spectrum from the susceptibility: Fluctuation-dissipation theorem

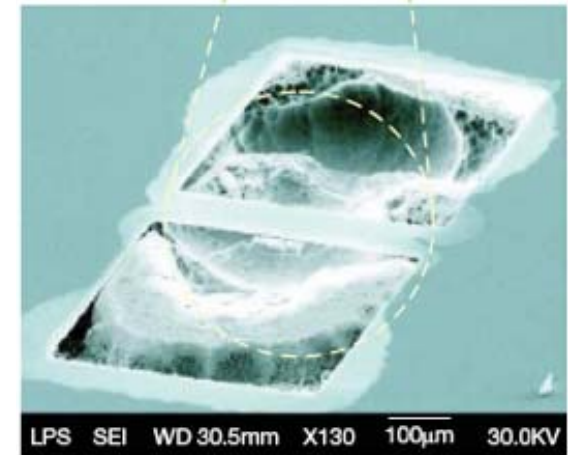
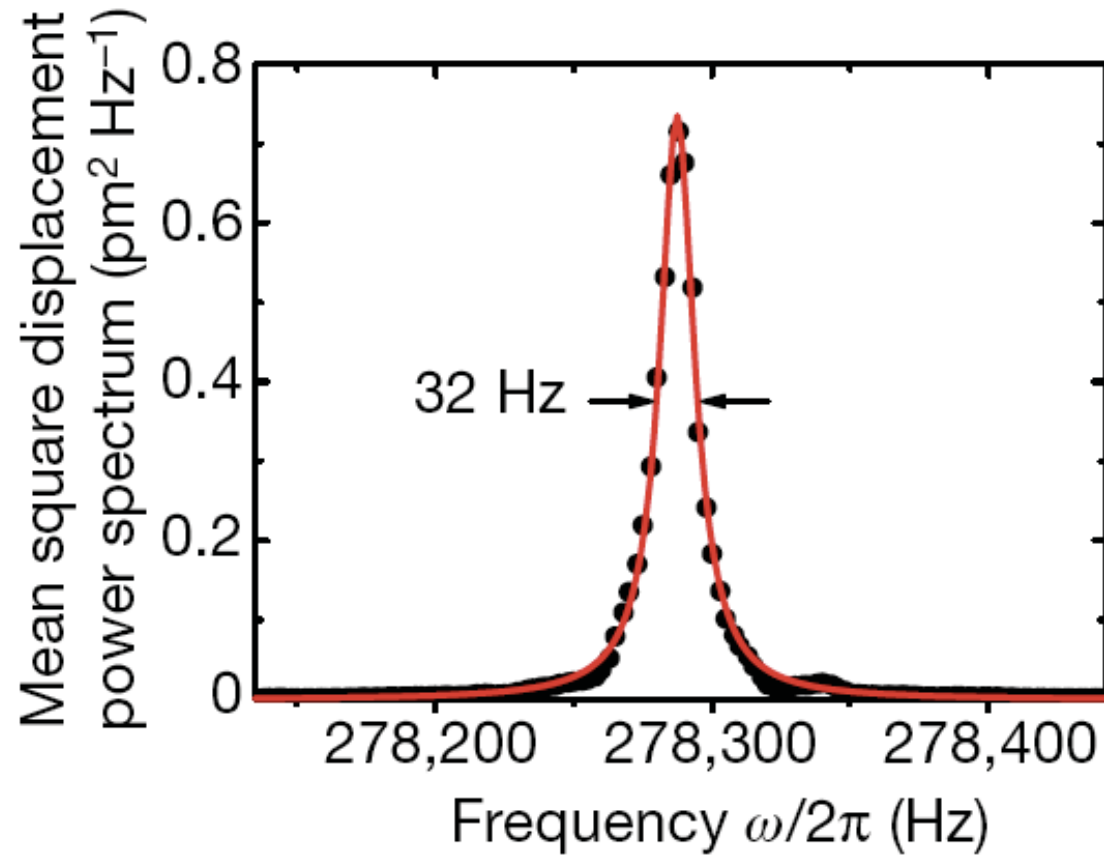
$$\begin{array}{ccc} \text{response} & & \text{force} \\ \langle \delta x \rangle(\omega) = \chi_{xx}(\omega) F(\omega) & & \\ & \text{susceptibility} & \end{array}$$

$$S_{xx}(\omega) = \frac{2k_B T}{\omega} \text{Im} \chi_{xx}(\omega) \quad (\text{classical limit})$$

for the damped oscillator:

$$m\ddot{x} + m\omega_M^2 x + m\Gamma\dot{x} = F$$
$$x(\omega) = \frac{1}{\underbrace{m(\omega_M^2 - \omega^2) - im\Gamma\omega}_{\chi_{xx}(\omega)}} F(\omega)$$

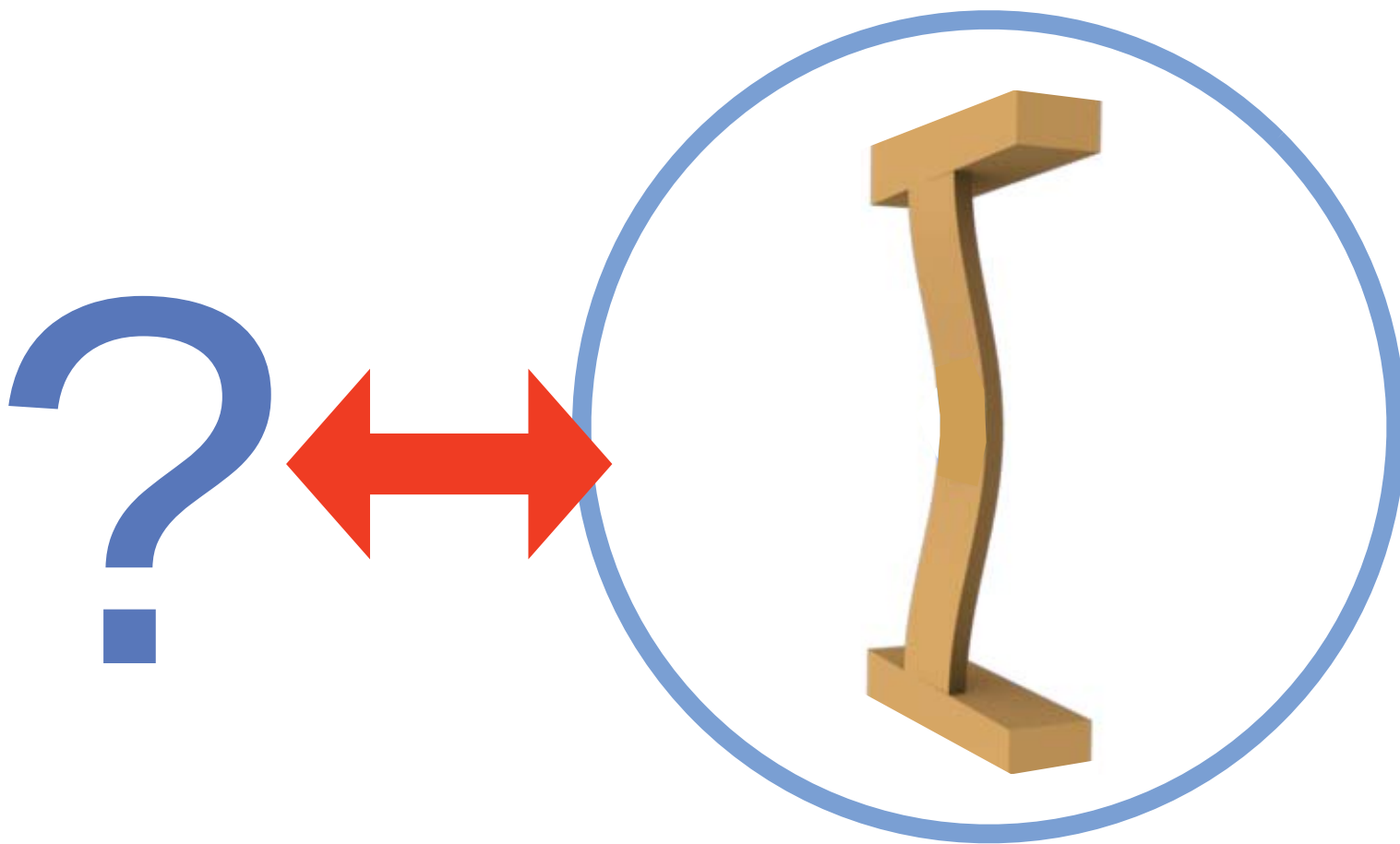




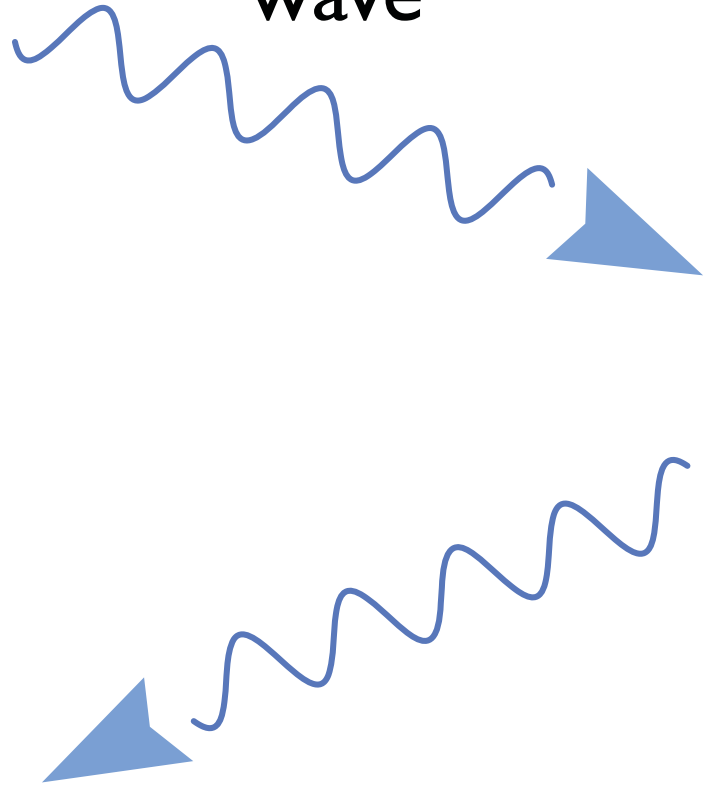
$T=300 \text{ K}$

Experimental curve:
Gigan et al., Nature 2006

Coupling radiation to a mechanical resonator



electromagnetic
wave



radio-frequency (kHz-MHz)

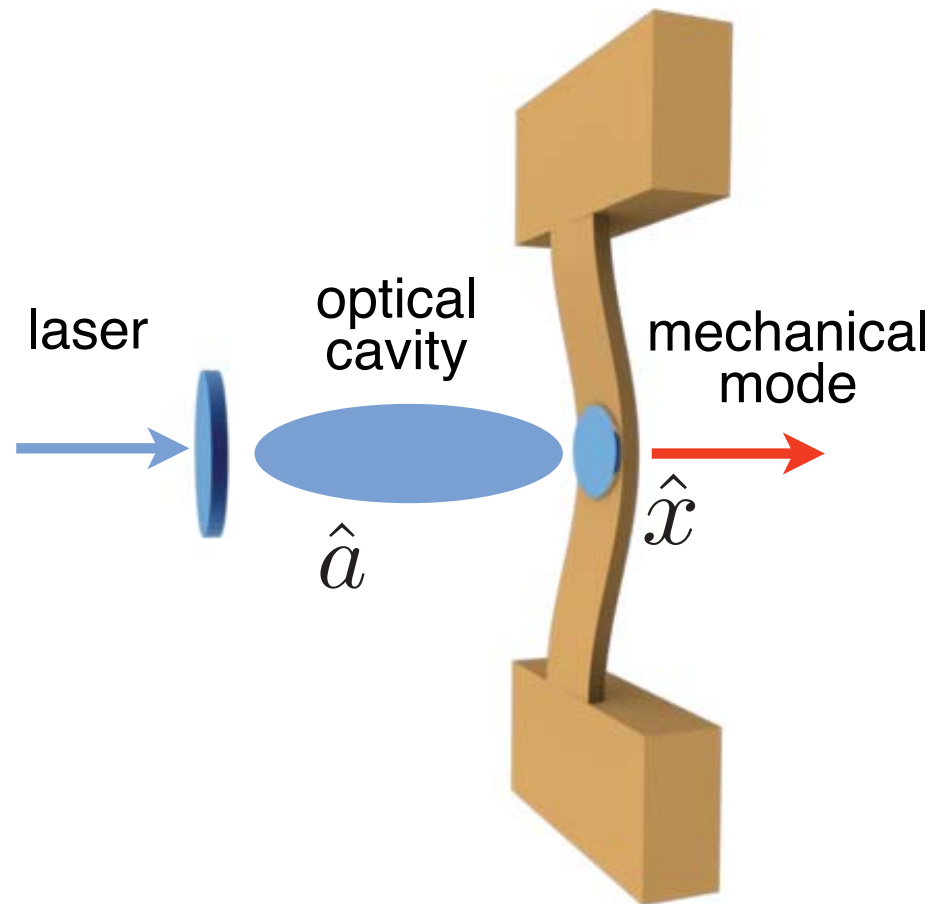
microwaves (GHz)

optical (THz)

optomechanical
coupling
force $\sim E^2(t)$

resonant
coupling
force $\sim E(t)$

The standard optomechanical setup

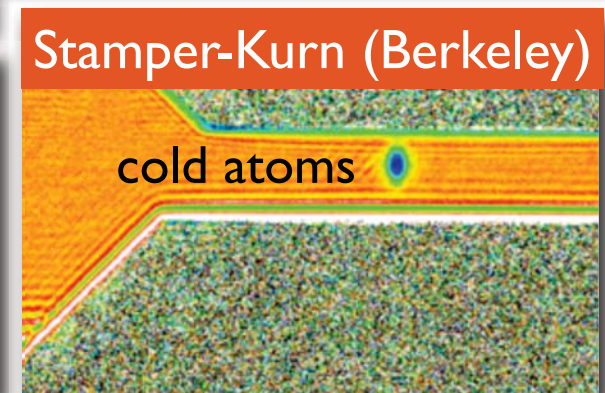
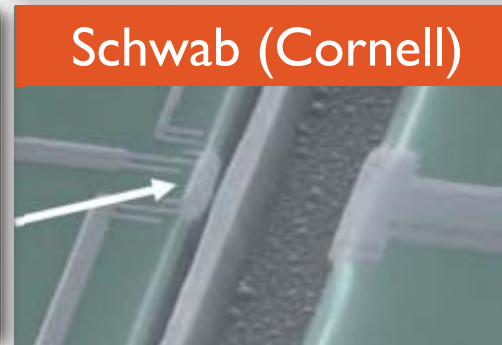
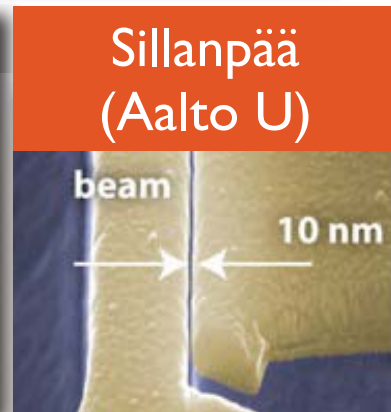
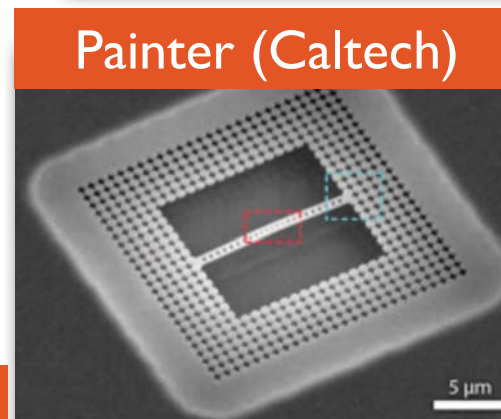
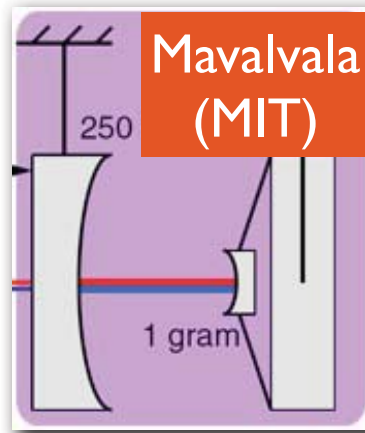
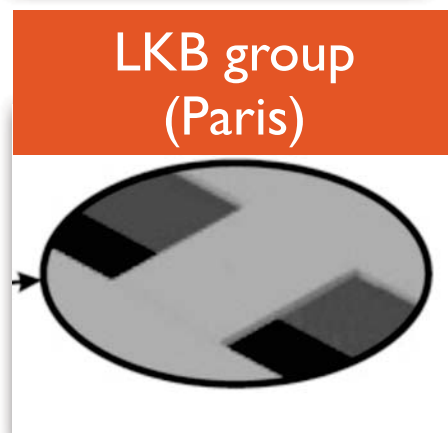
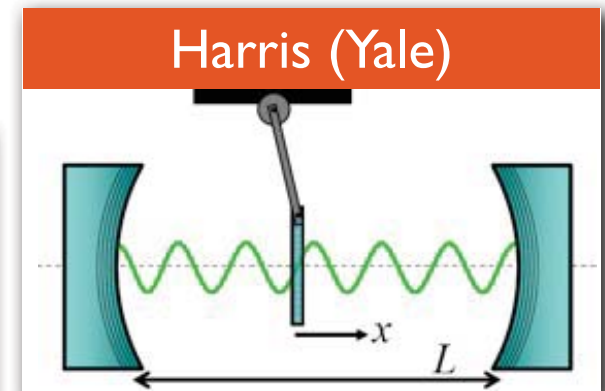
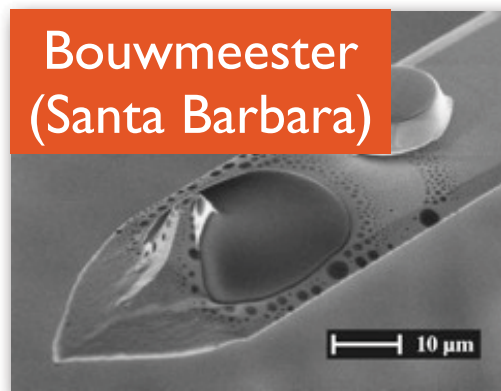
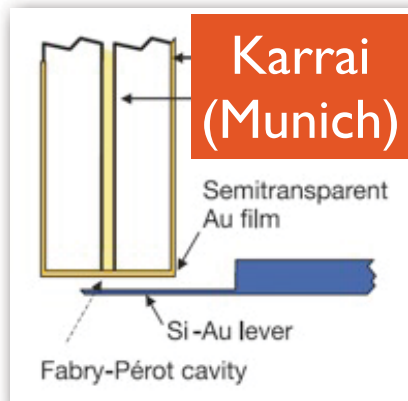


$$\hat{H} = \hbar\omega_{\text{cav}} \cdot (1 - \hat{x}/L)\hat{a}^\dagger\hat{a} + \hbar\Omega\hat{b}^\dagger\hat{b} + \dots$$

Recent Review “Cavity Optomechanics”:
M. Aspelmeyer, T. Kippenberg, FM; arXiv 2013

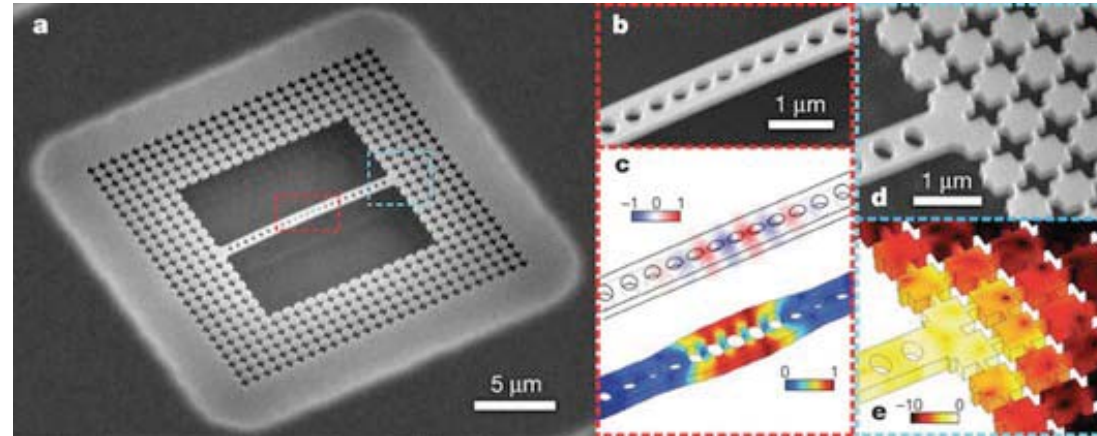
$$\hat{x} = x_{\text{ZPF}}(\hat{b} + \hat{b}^\dagger)$$
$$x_{\text{ZPF}} = \sqrt{\hbar/2m\Omega}$$

Optomechanical experiments (selection)

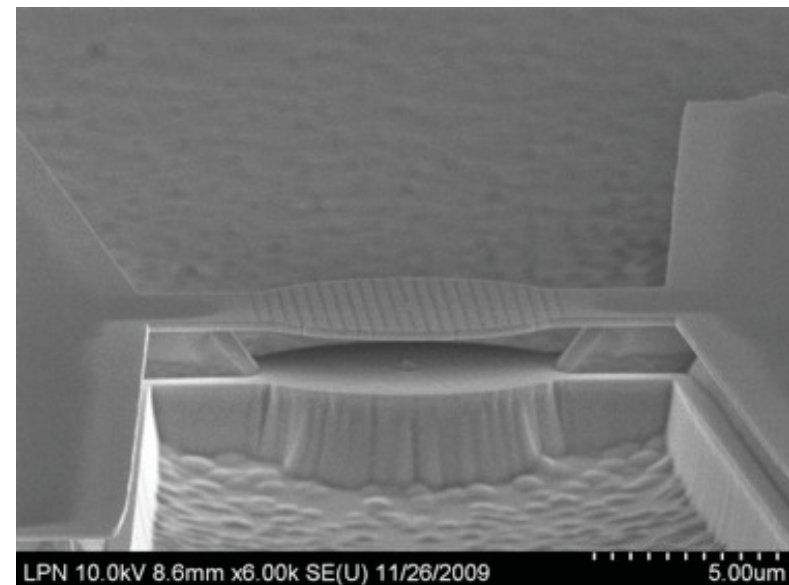
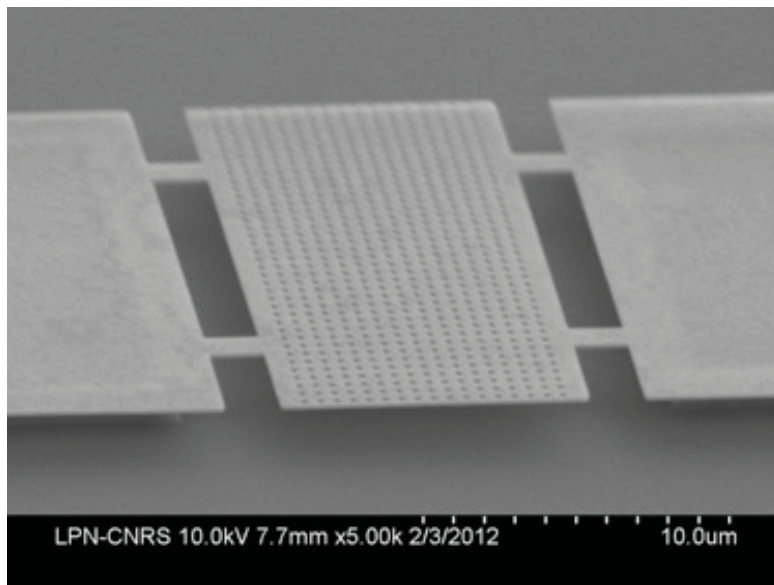


Photonic crystals: Very strong coupling between localized vibrational and optical modes

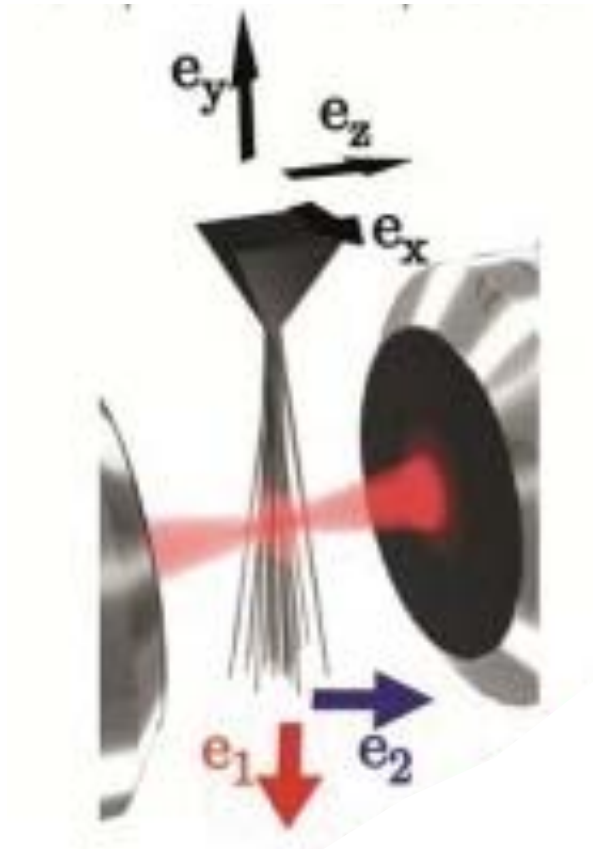
Amir Safavi-Naeini (Wednesday)



Isabelle Robert (Thursday)



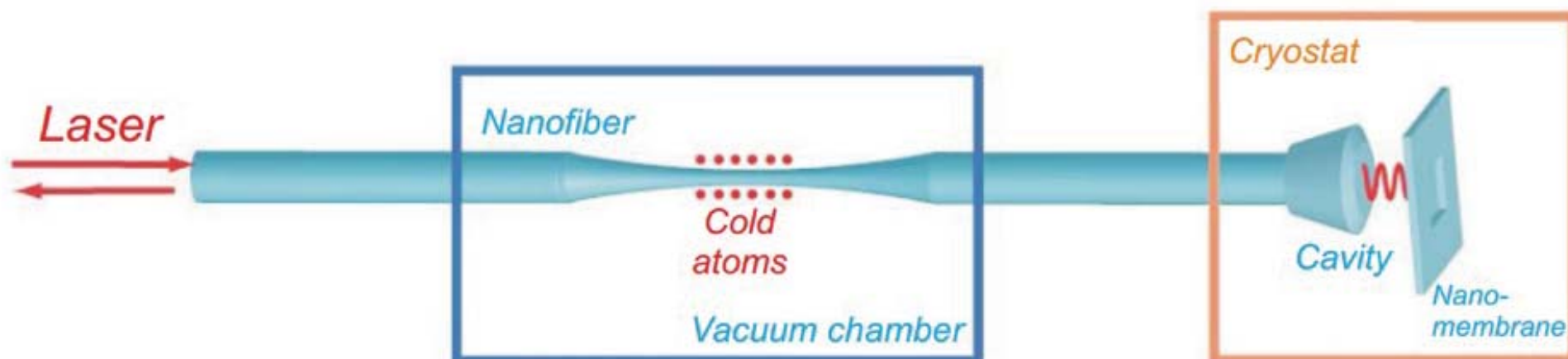
Nano-Optomechanics: Nanowire in a light field



- Ultra-sensitive nano-optomechanical detection of a bi-dimensional nanomechanical degree of freedom
- Topological structure of the radiation force in a focused laser beam

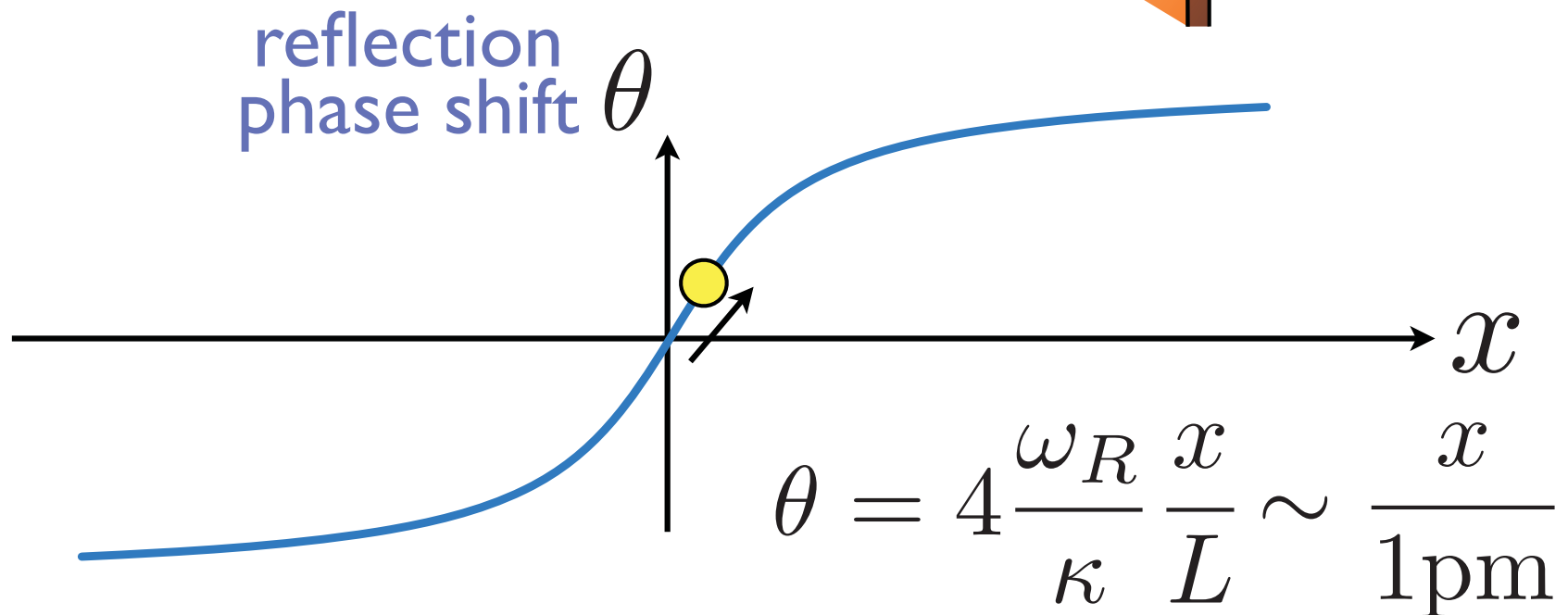
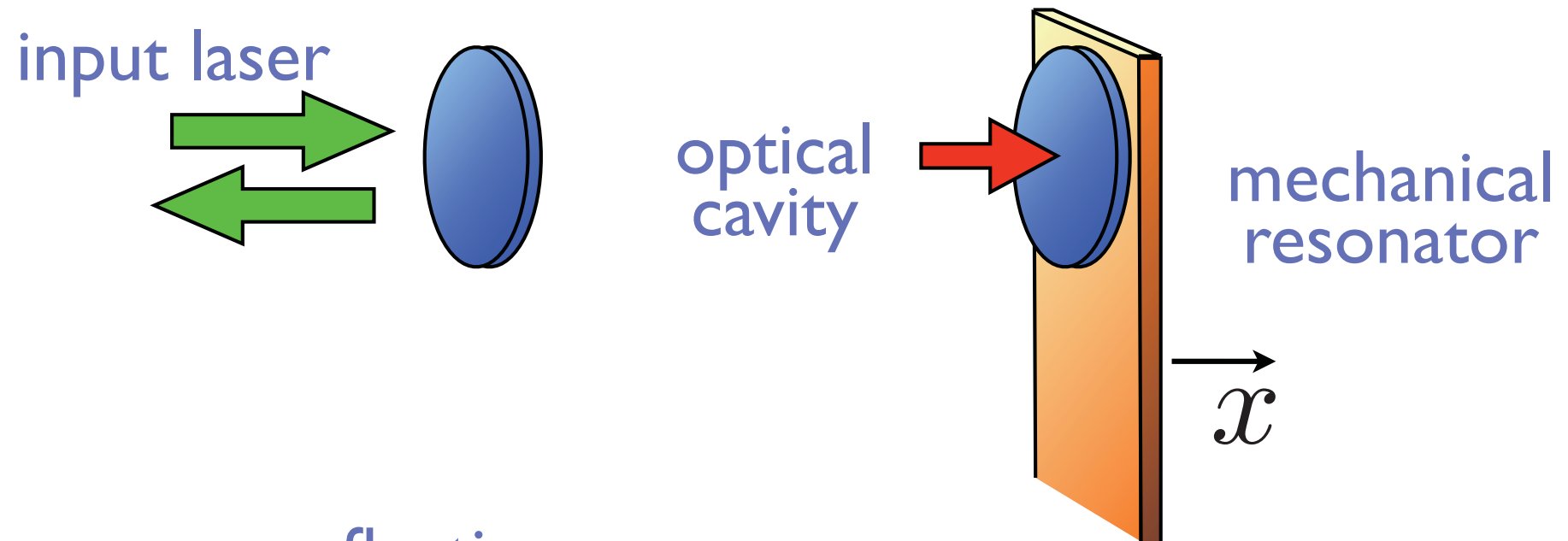
Pierre Verlot (Friday)

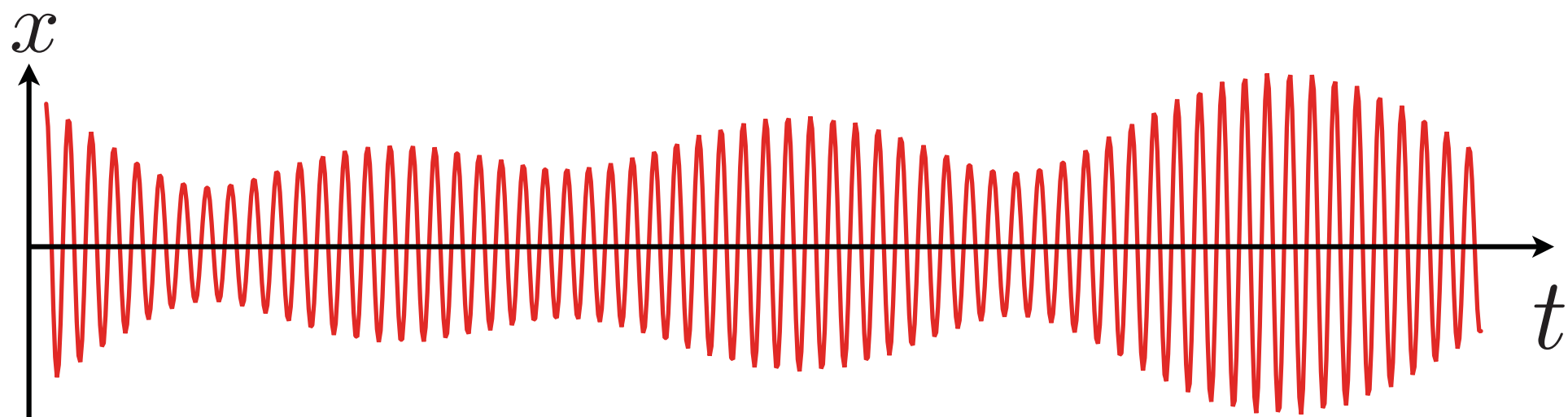
Coupling to atoms

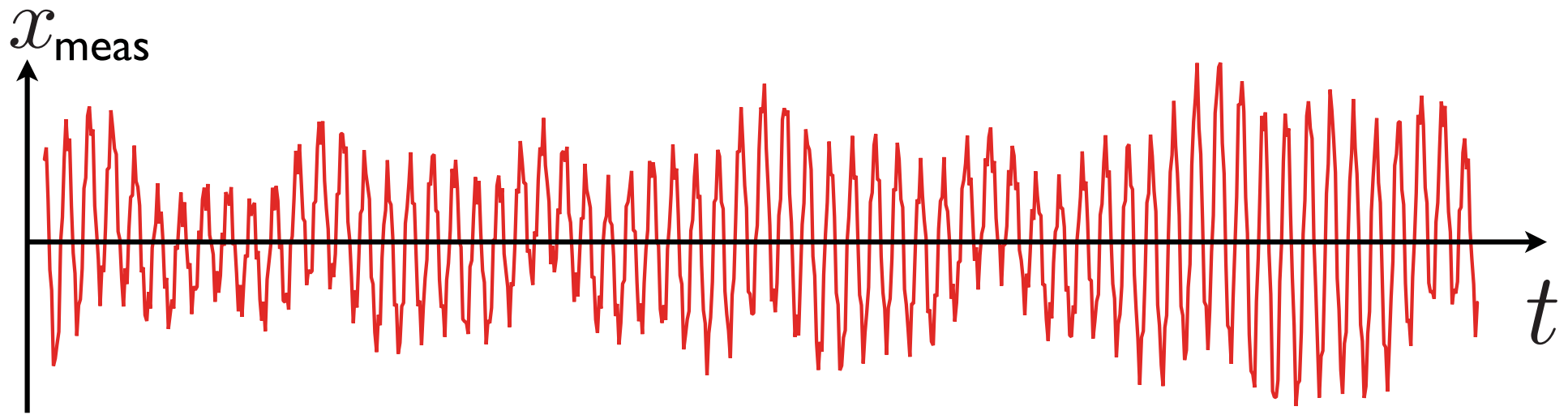


Measuring mechanical motion

Optical detection of mechanical motion







$$x_{\text{meas}}(t) = x(t) + x_{\text{noise}}(t)$$

Two contributions to $x_{\text{noise}}(t)$

1. measurement imprecision

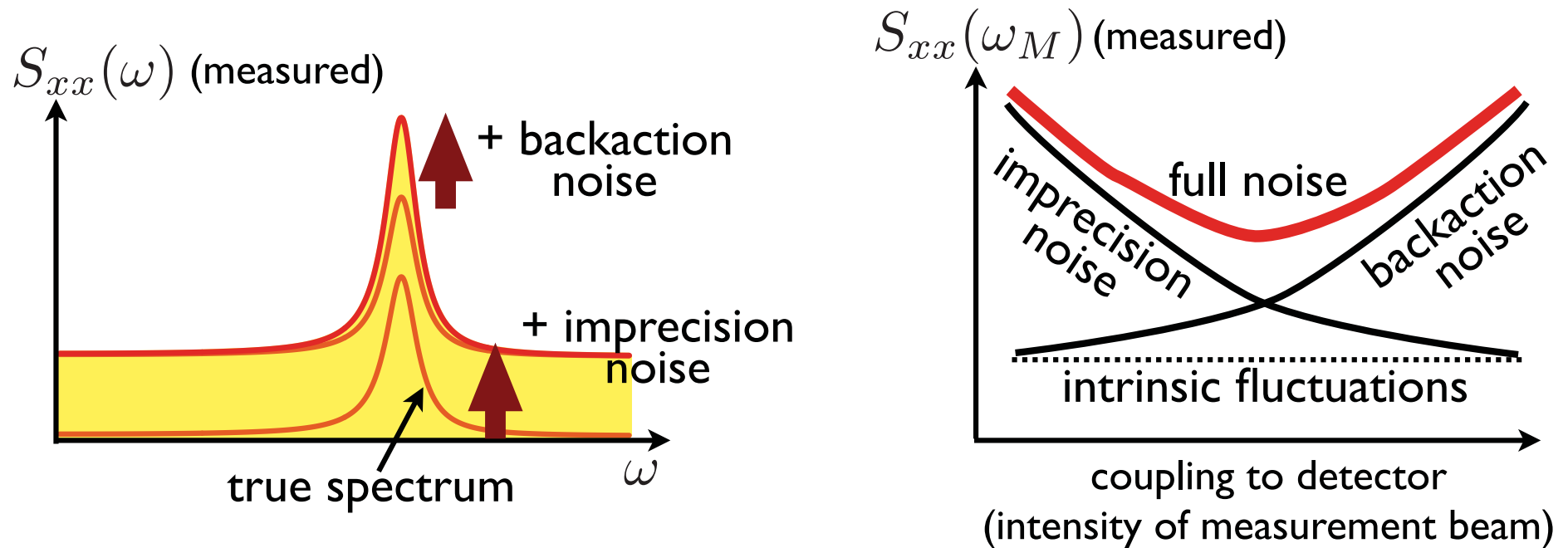
2. measurement back-action:

fluctuating force on system

phase noise of
laser beam (shot
noise limit!)

noisy radiation
pressure force

“Standard quantum limit” of displacement detection



Best case allowed by quantum mechanics:

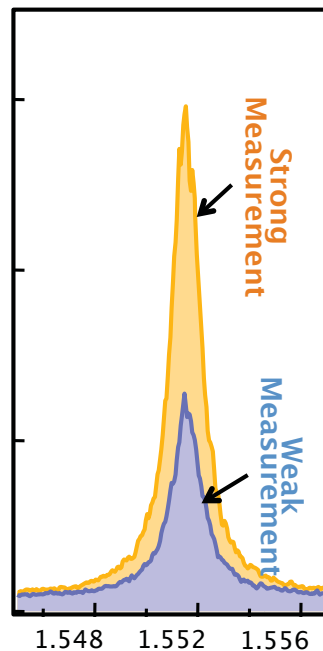
$$S_{xx}^{(\text{meas})}(\omega) \geq 2 \cdot S_{xx}^{T=0}(\omega)$$

“Standard quantum limit (SQL) of displacement detection”

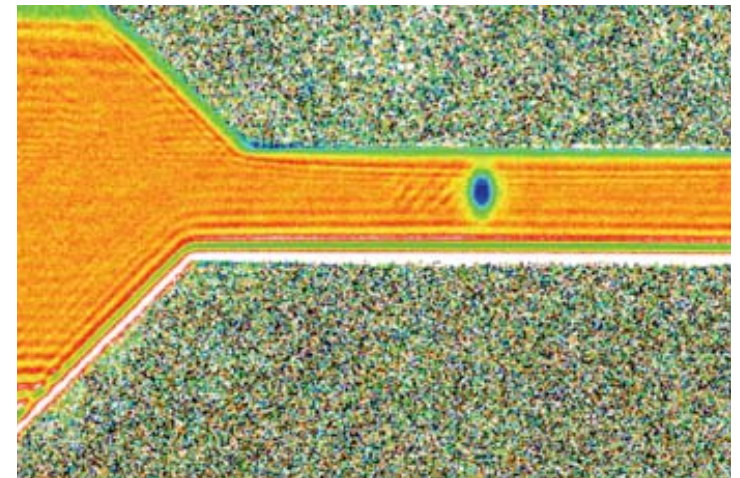
Challenge: Reach optimal regime (where backaction becomes important)

Recent experimental results:

Solid state:
Membrane
resonator



Cold atoms



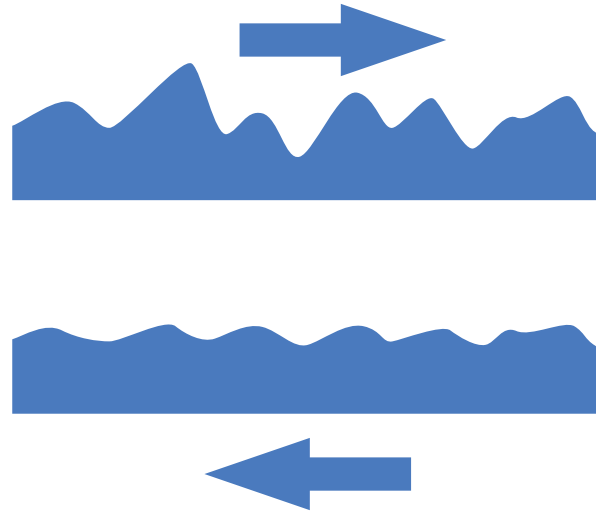
(Berkeley group)

Thomas Purdy (Monday)

Sydney Schreppler (Tuesday)

Strong backaction induces squeezing of radiation field!

Input: laser field with fluctuating intensity



Output: reduced (“squeezed”) noise

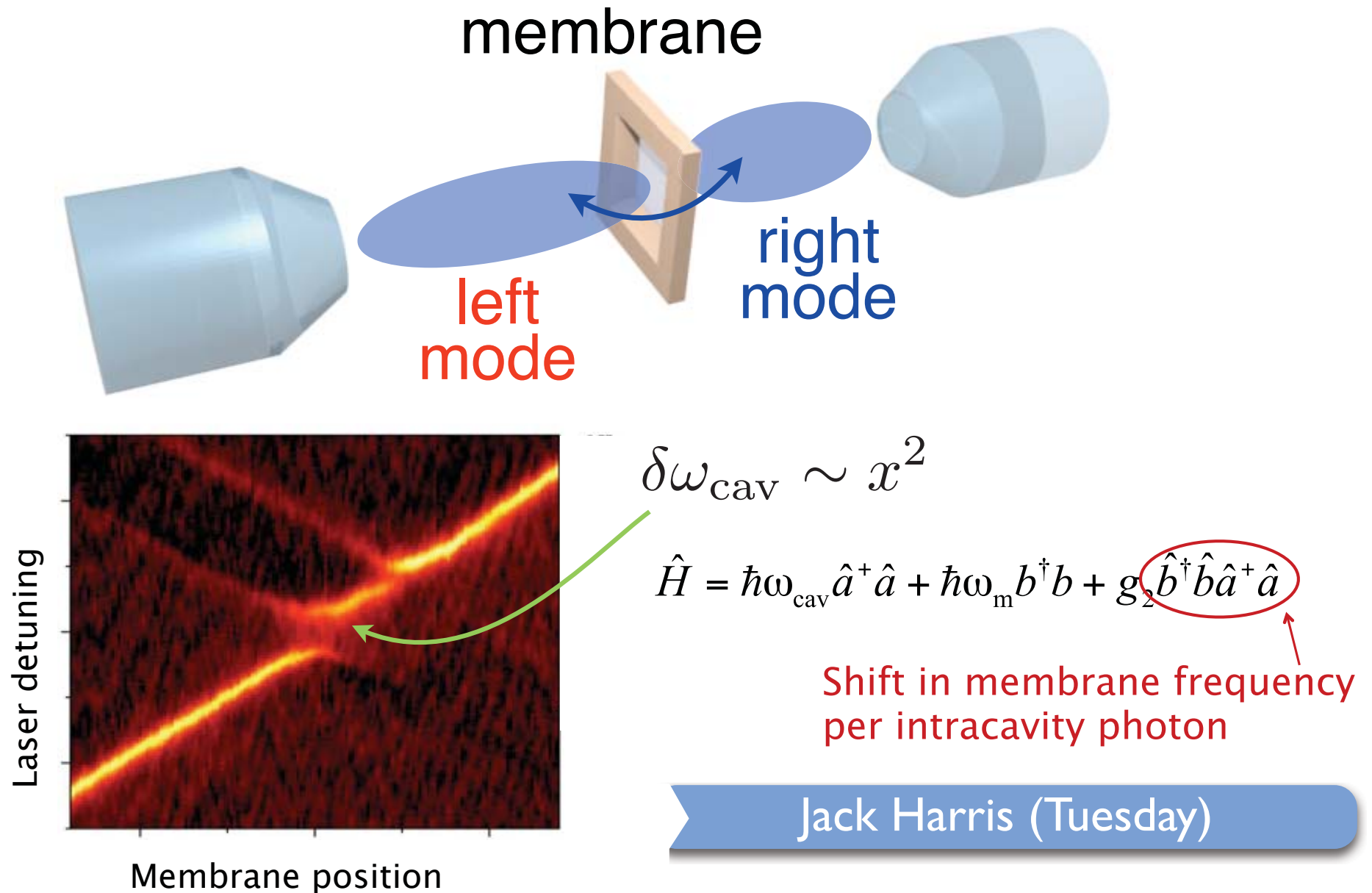
Optomechanical system:
Intensity-dependent
optical resonance
(=Kerr medium)

Amir Safavi-Naeini (Wednesday)

Thomas Purdy (Monday)

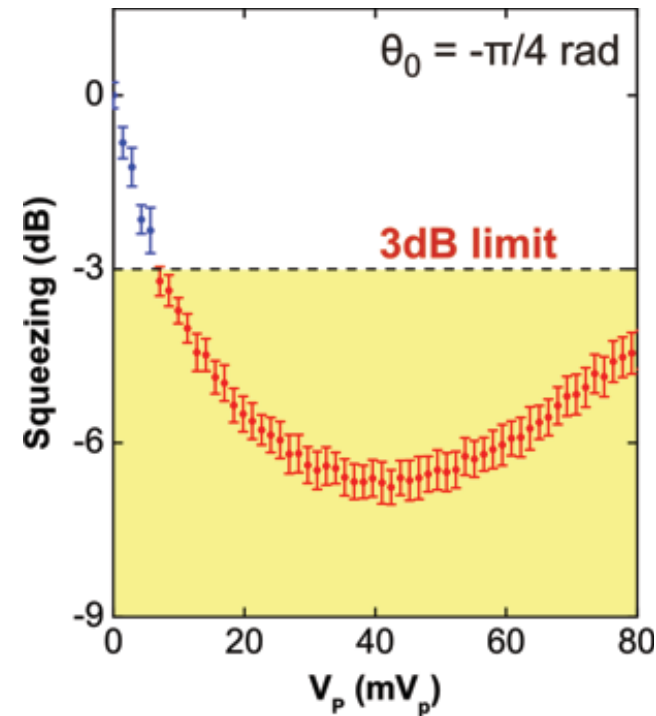
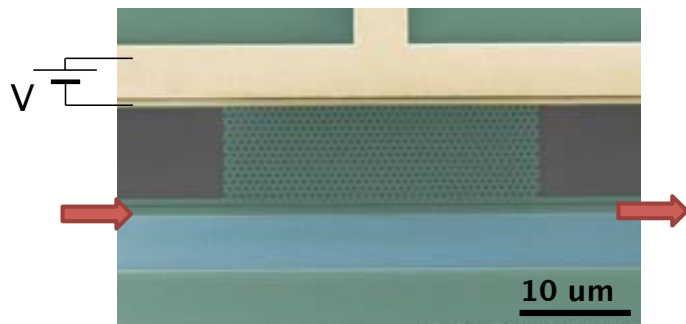
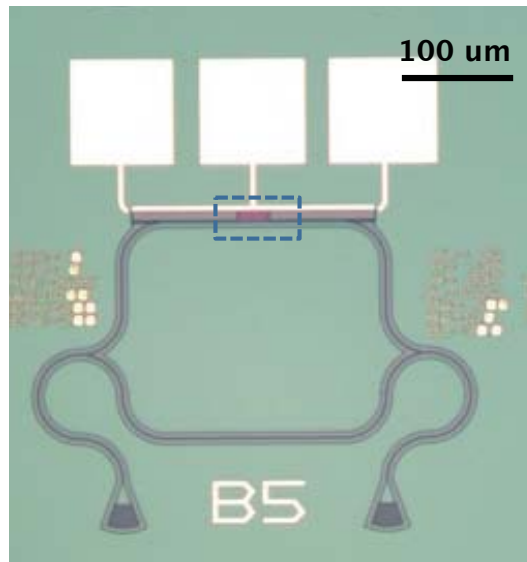
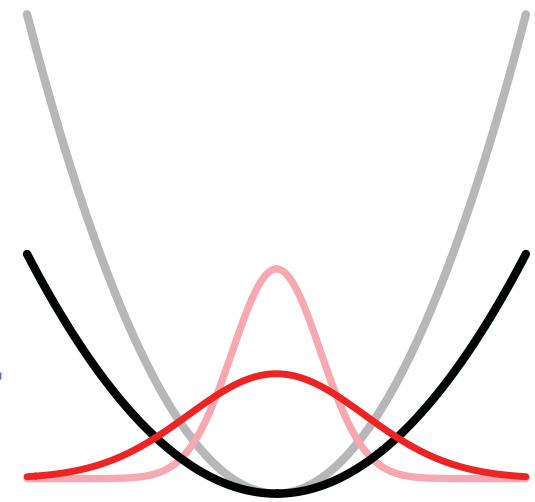
Measuring x^2 instead of x

(Ultimate goal: phonon number detection)



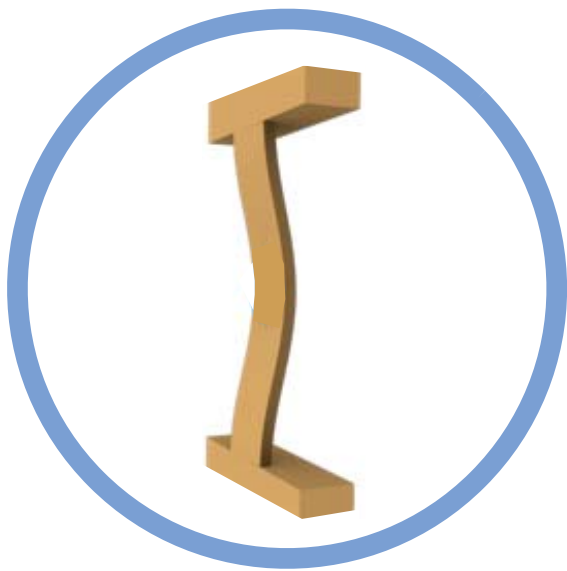
Sensitive measurement can be used for feedback!

here: use feedback to optimize squeezing of a thermal mechanical state
general trick: time-dependent modulation of spring constant produces squeezing

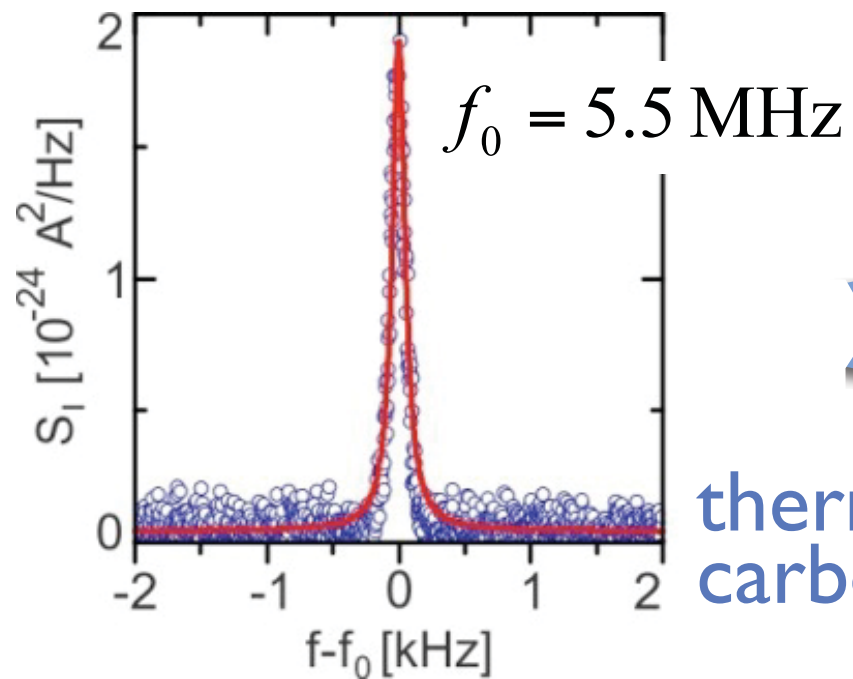
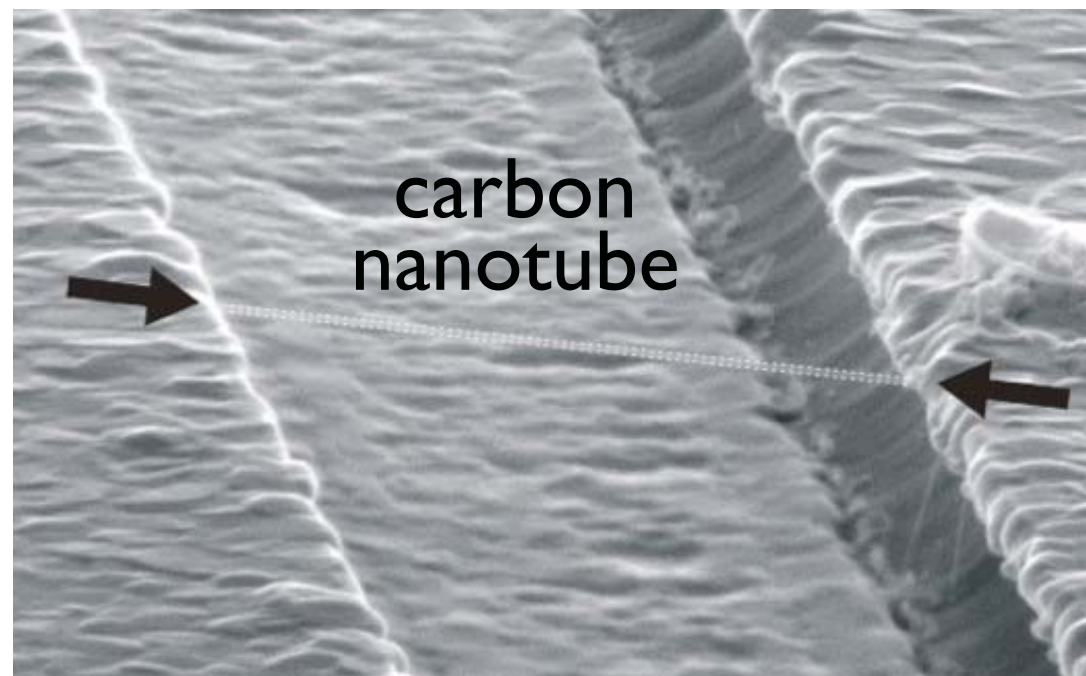


Menno Poot (Friday)

Mechanical resonators from carbon



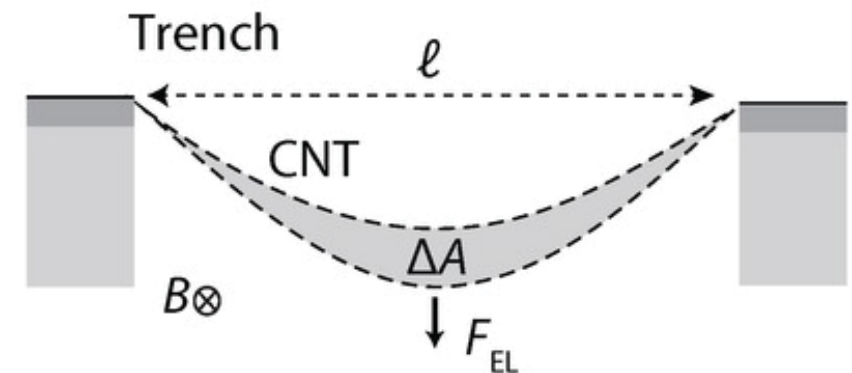
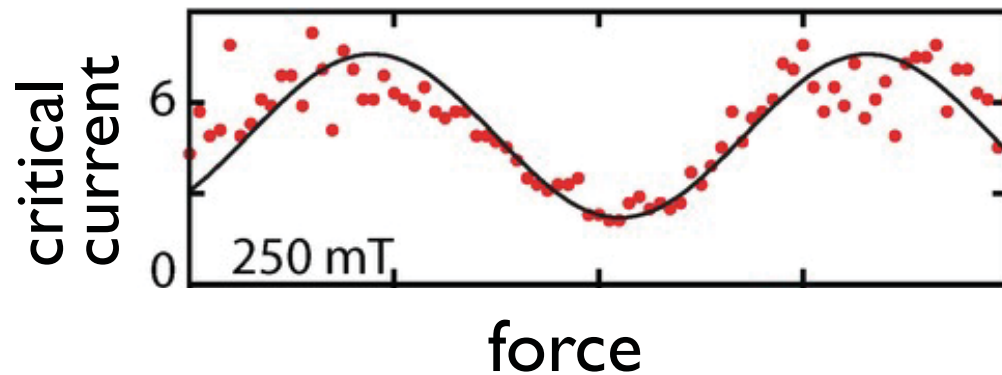
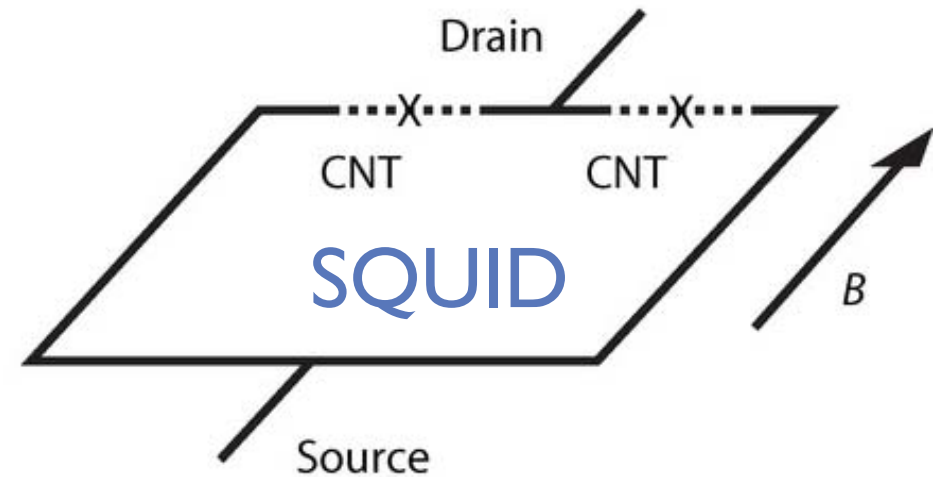
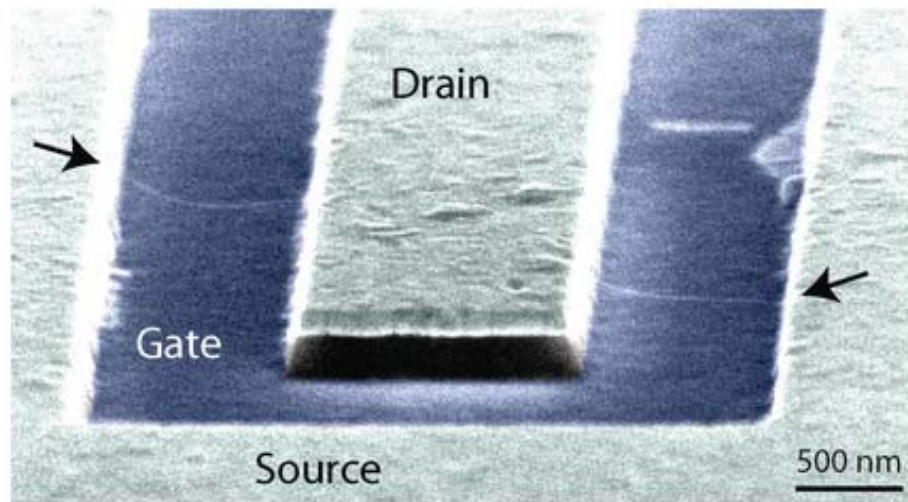
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Adrian Bachtold (Wednesday)

thermal motion of a
carbon nanotube

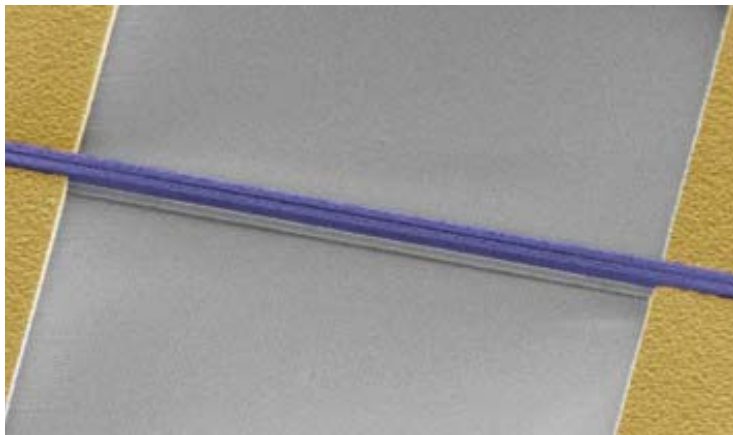
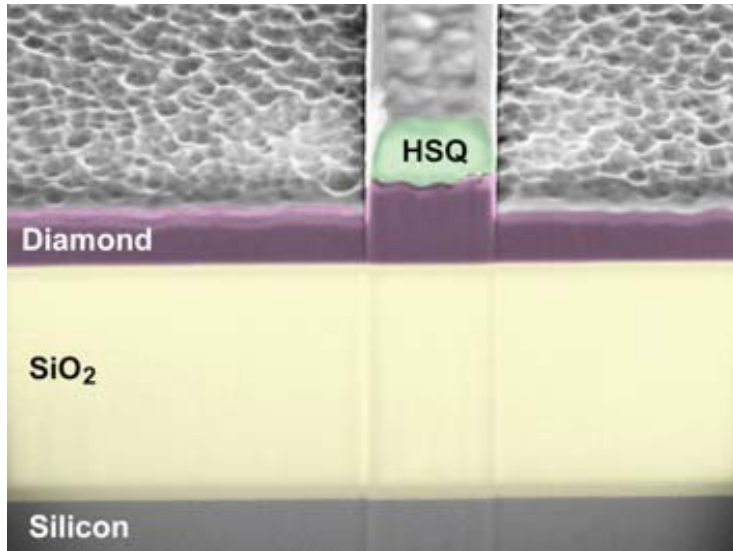
Carbon nanotubes: very low mass, strong quantum zero-point fluctuations – couple to other quantum devices!



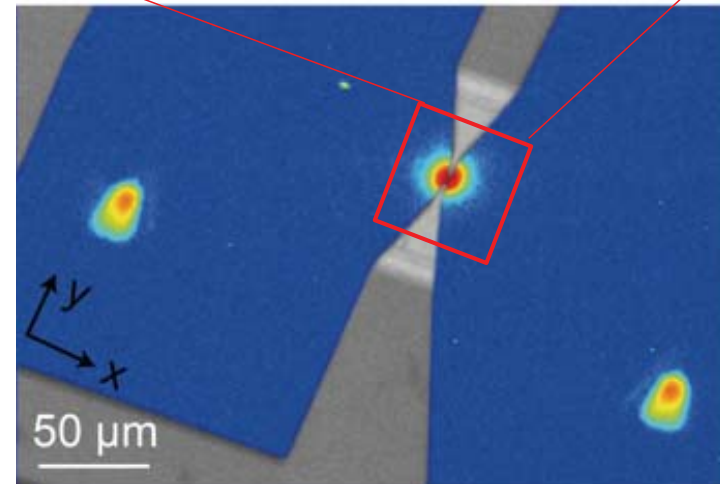
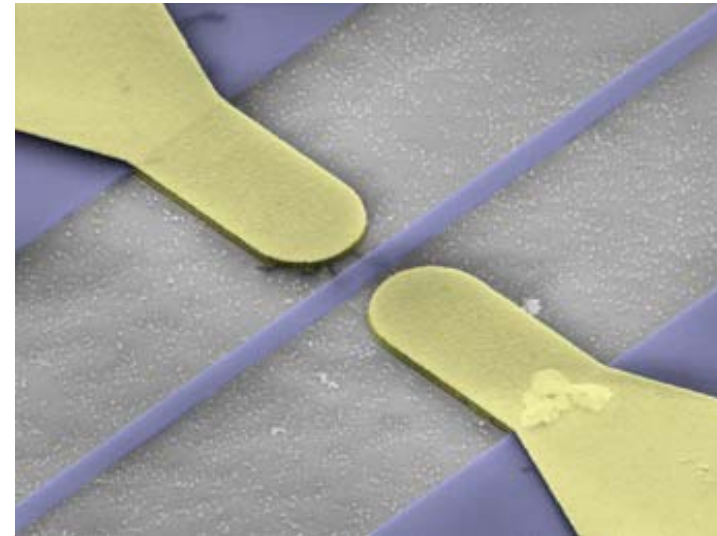
carbon nanotube motion
coupled to a
superconducting circuit
(SQUID)

Carbon nanotubes or diamond in photonic circuits

Diamond nanophotonic circuits



Waveguide integrated carbon nanotubes



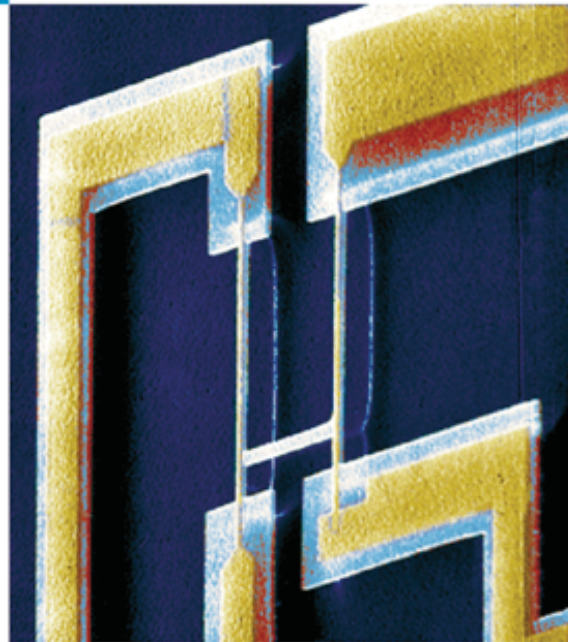
The Quantum Regime

(still mostly theory, but first experiments exist)



JULY
2005

PHYSICS TODAY



The quantum
mechanic's toolbox

Putting Mechanics into Quantum Mechanics

Nanoelectromechanical structures are starting to approach the ultimate quantum mechanical limits for detecting and exciting motion at the nanoscale. Nonclassical states of a mechanical resonator are also on the horizon.

Keith C. Schwab and Michael L. Roukes

Everything moves! In a world dominated by electronic devices and instruments it is easy to forget that all measurements involve motion, whether it be the motion of electrons through a transistor, Cooper pairs or quasiparticles through a superconducting quantum interference device (SQUID), photons through an optical interferometer—or the simple displacement of a mechanical element.

achieved to read out those devices, now bring us to the realm of quantum mechanical systems.

The quantum realm

What conditions are required to observe the quantum properties of a mechanical structure, and what can we learn when we encounter them? Such questions have received

Schwab and Roukes, Physics Today 2005

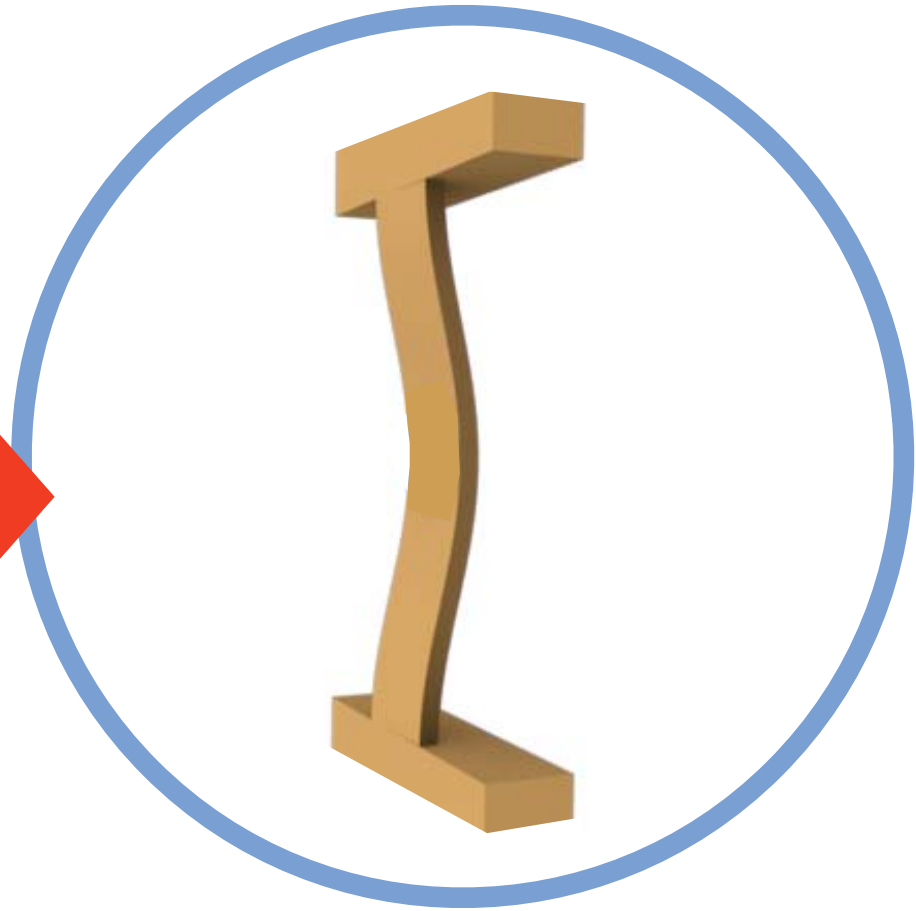
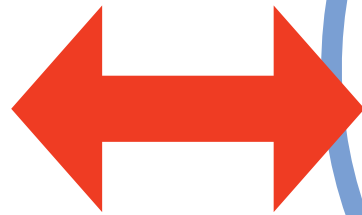
- **nano-electro-mechanical systems**

Superconducting qubit coupled to nanoresonator: Cleland & Martinis 2010

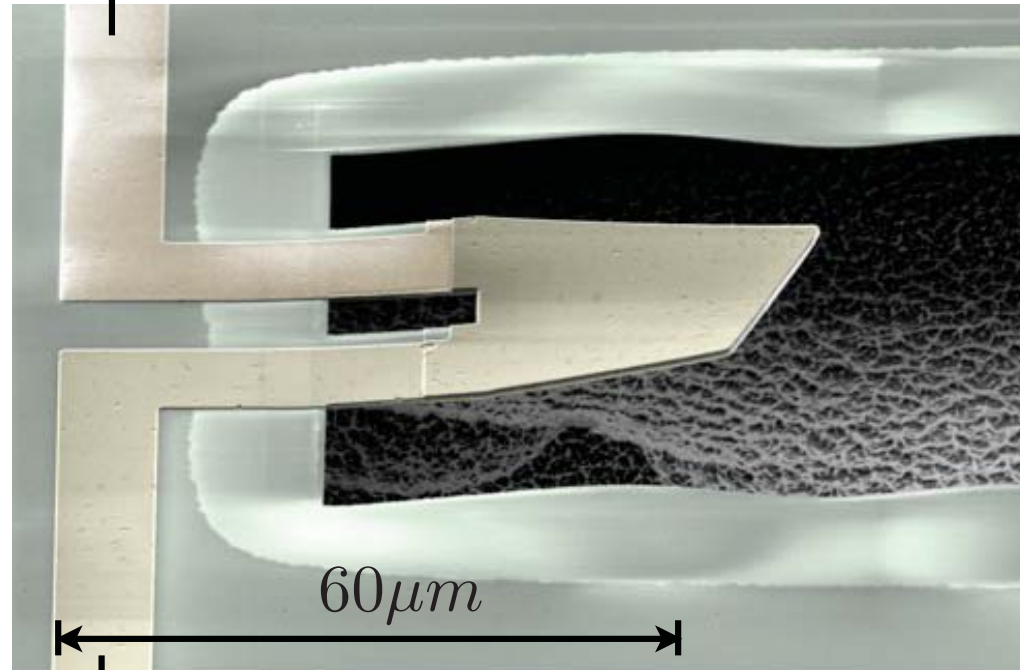
- **optomechanical systems**

Laser-cooled to ground state: Teufel et al in microwave circuit 2011,
Painter group in photonic crystal 2011

**two-
level
system
(qubit)**



**Josephson
phase
qubit**

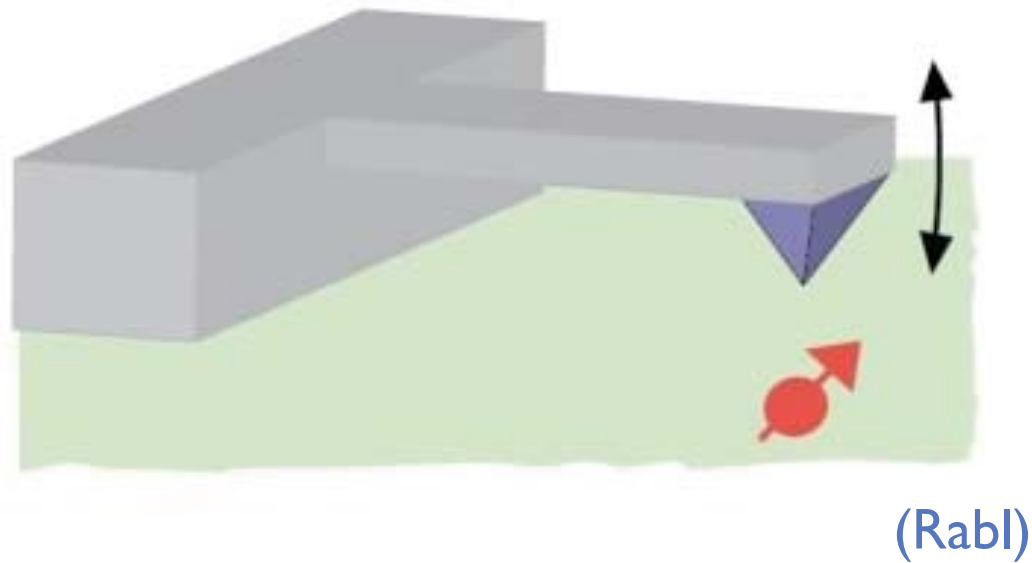


piezoelectric nanomechanical resonator

(GHz @ 20 mK: ground state!)

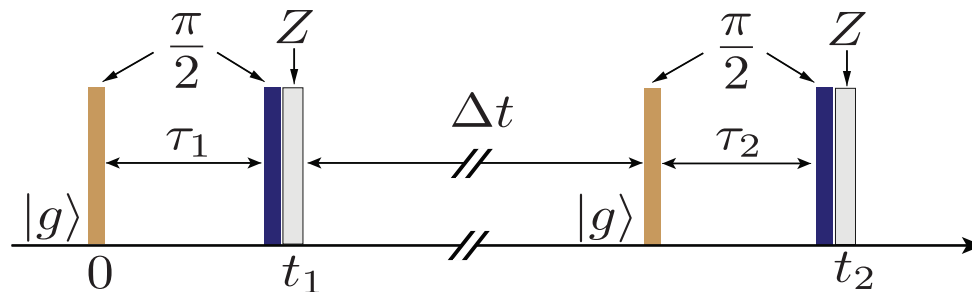
swap excitation between qubit and
mechanical resonator in a few ns!

Nanomechanical resonator coupled to spin



Two-level system as a probe of a mechanical resonator

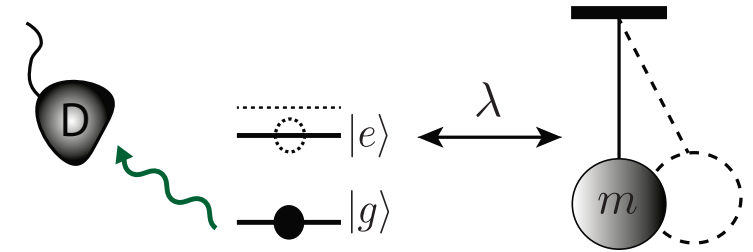
probe quantum superpositions of a macroscopic resonator via multiple Ramsey measurements:



Correlations between subsequent measurement outcomes violate the Leggett-Garg inequality

$$C(t_1, t_2) + C(t_2, t_3) - C(t_1, t_3) \leq 1$$

and can be used for other fundamental tests of quantum mechanics !



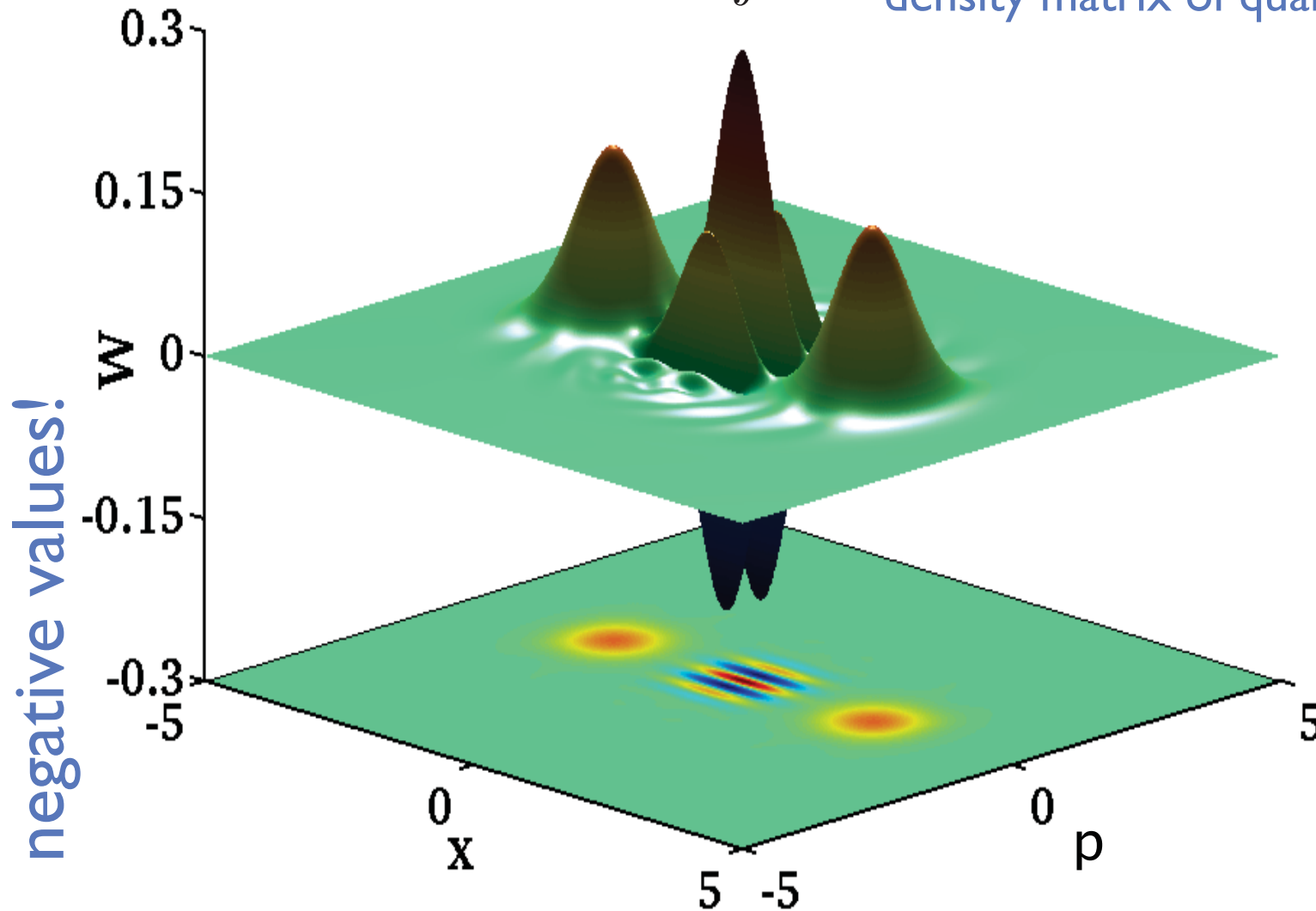
$$\Rightarrow C(t_1, t_2) = \langle Z(t_2)Z(t_1) \rangle$$

Nonclassical mechanical states

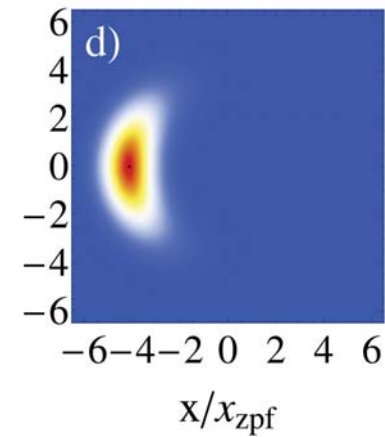
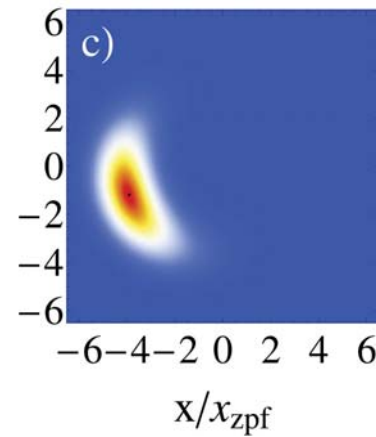
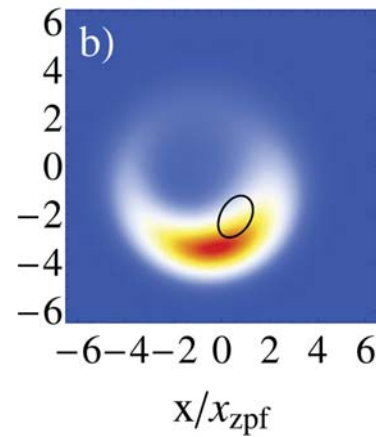
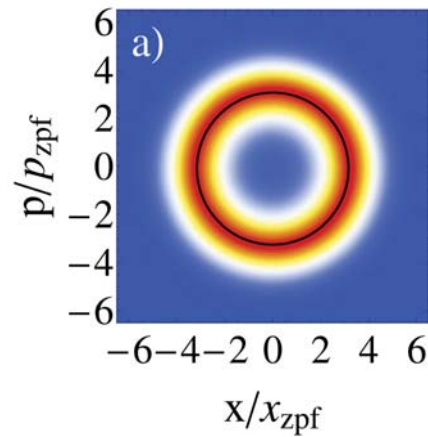
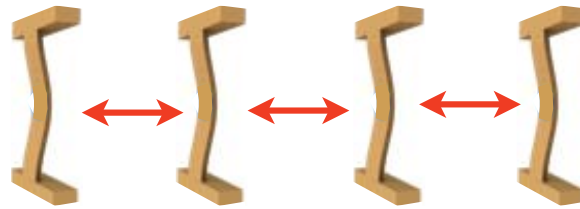
Wigner phase space density:

$$W(x, p) \propto \int dy e^{ipy/\hbar} \rho(x - y/2, x + y/2)$$

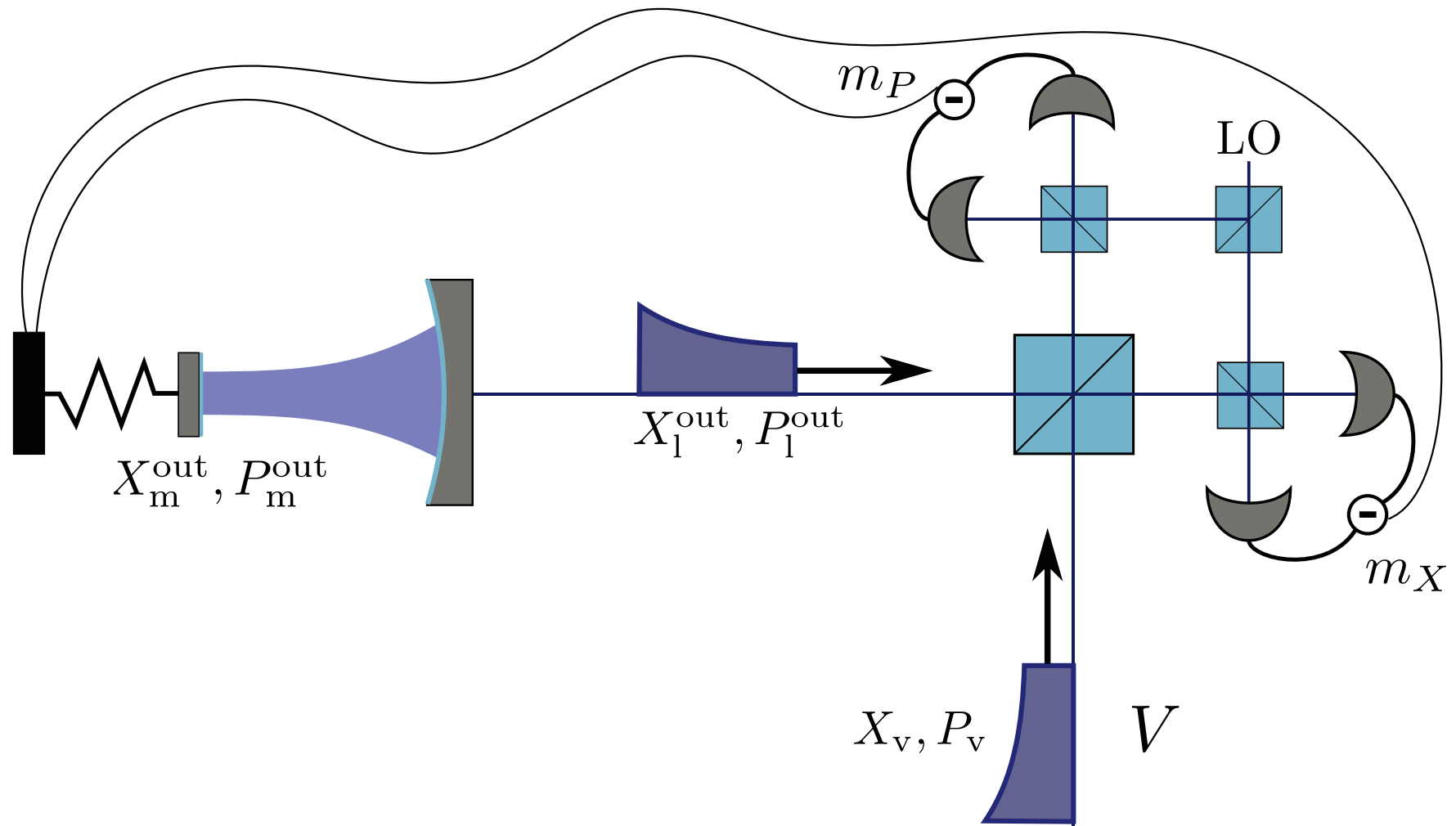
density matrix of quantum state



Synchronization between multiple resonators in the quantum regime

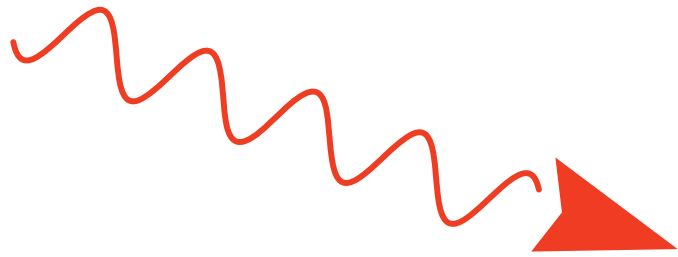


Optomechanical control & entanglement with light pulses

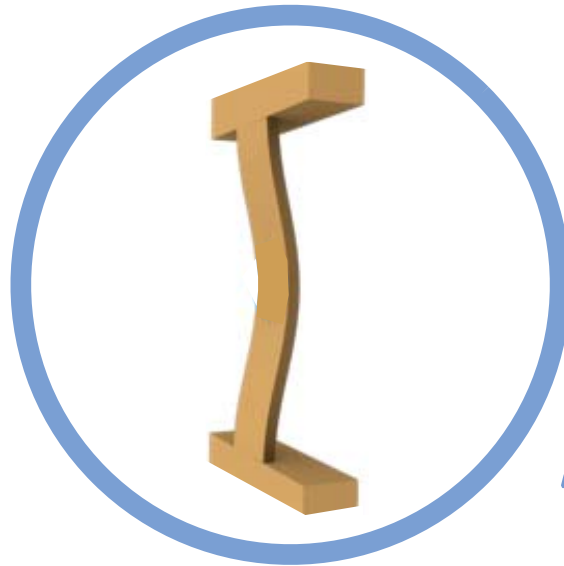


A quantum interface: Taking quantum information from microwave to optical

microwave field



GHz

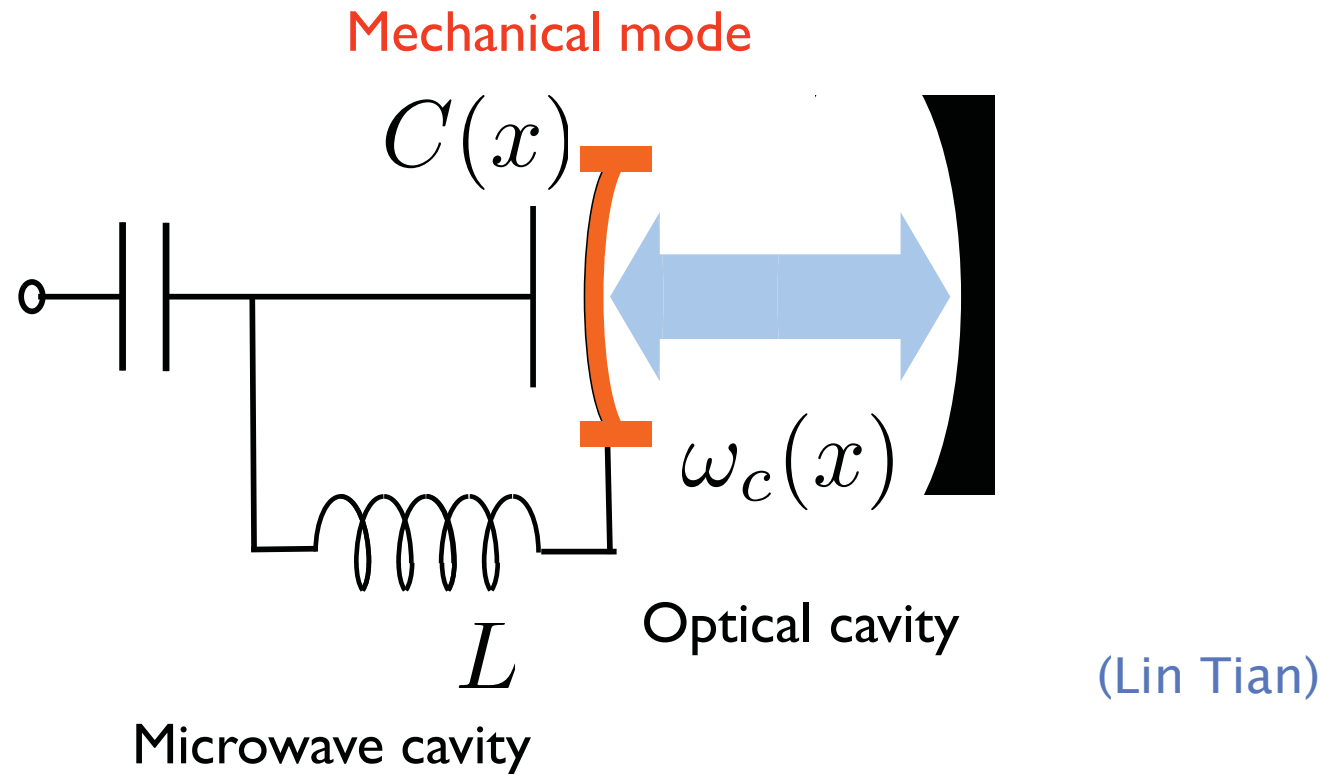


optical field



100 THz

Mechanical mode connects resonators with different frequencies



Connect different parts of a hybrid quantum network
Achieve quantum operations through the mechanical mode

Experiment:

Andrew Cleland (Monday)

Theory:

Lin Tian (Wednesday)

Summary

