

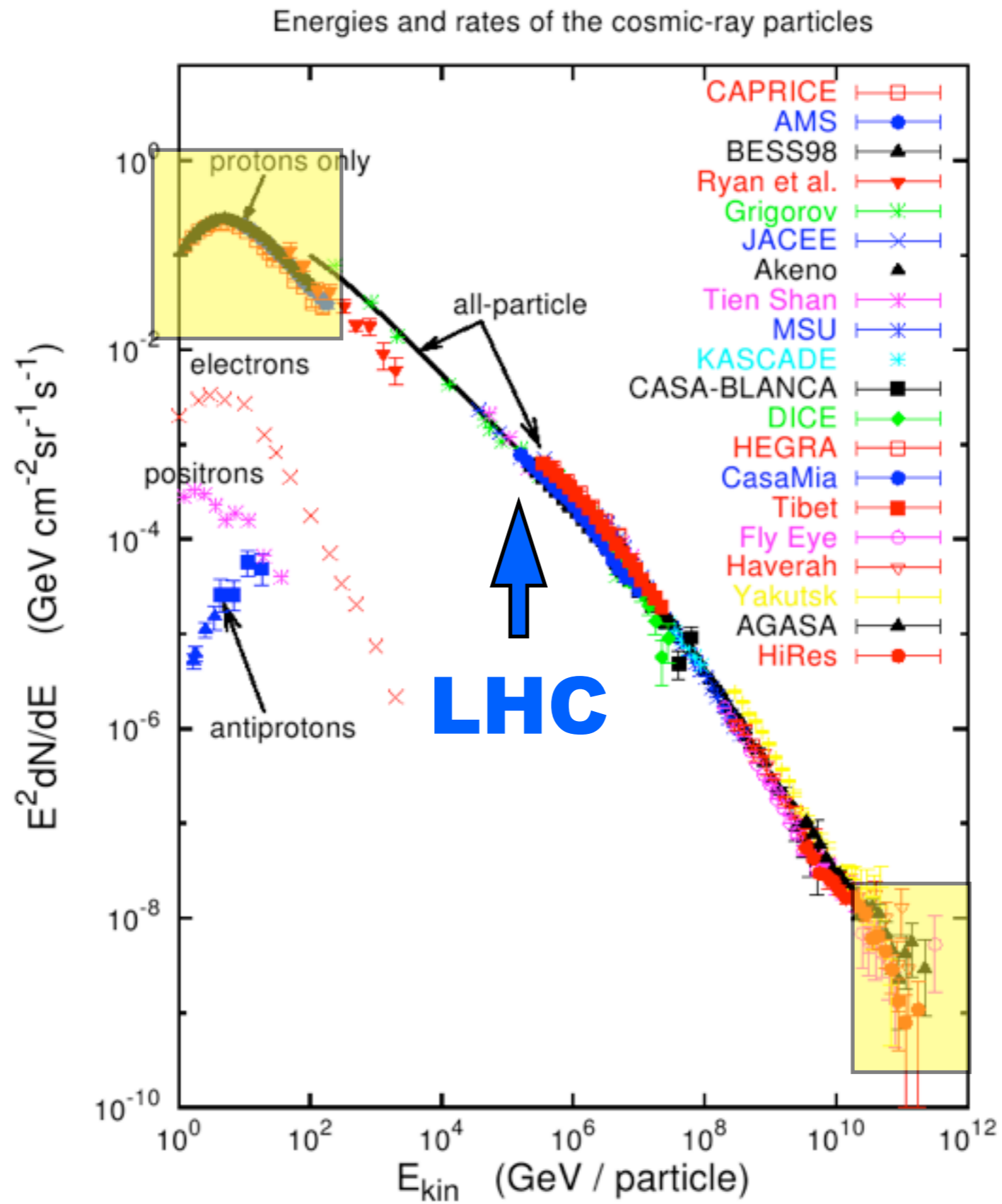
Cosmic-rays propagation in the ISM

Carmelo Evoli (Universität Hamburg)



Trieste | ICTP Workshop | 11th of October 2013

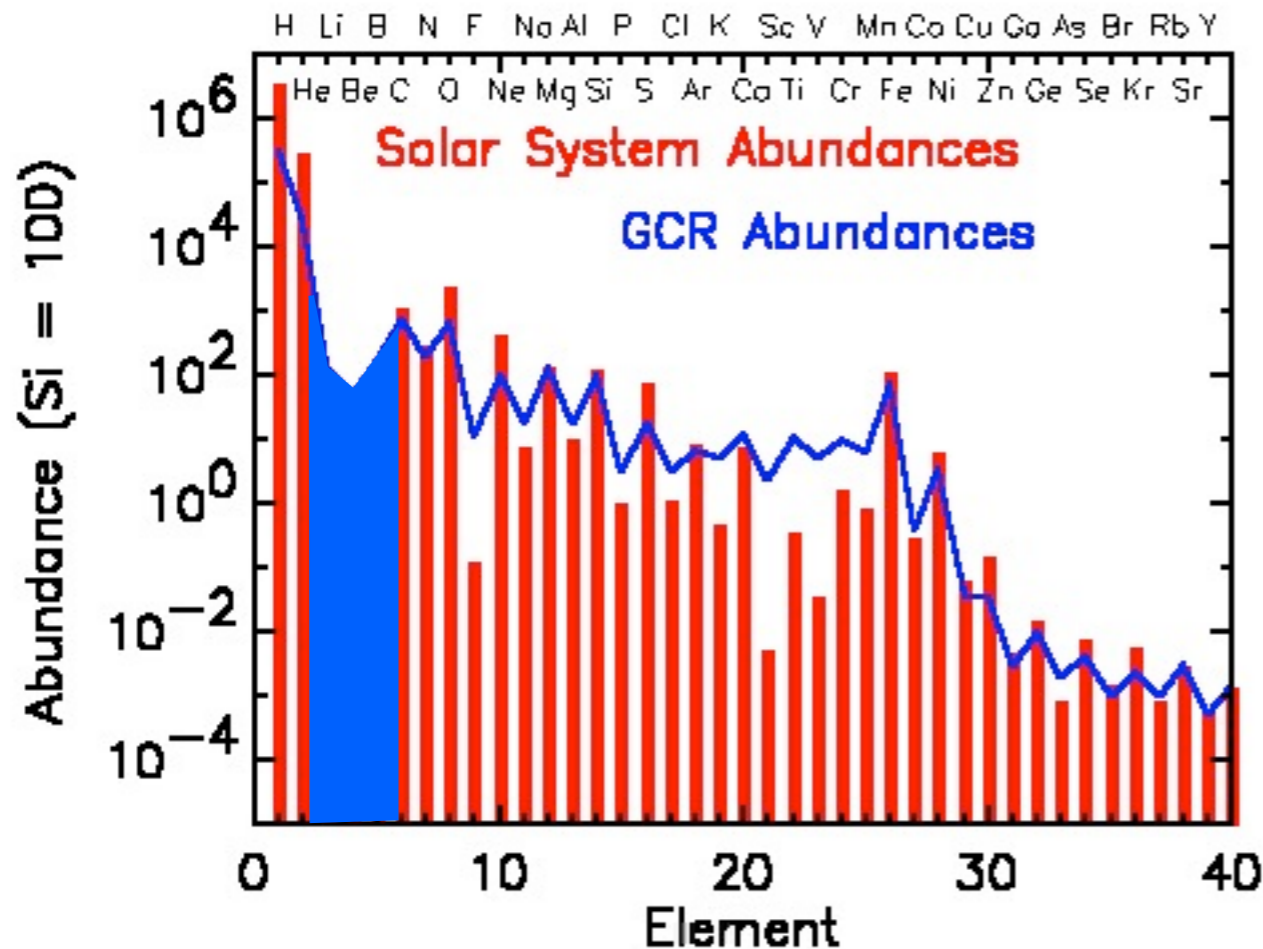
1/cm²/s



1/km²/century

$$L_{\text{SN}} \sim R_{\text{SN}} E_{\text{kin}} \sim 3 \times 10^{41} \text{ erg/s}$$

Secondary / Primary



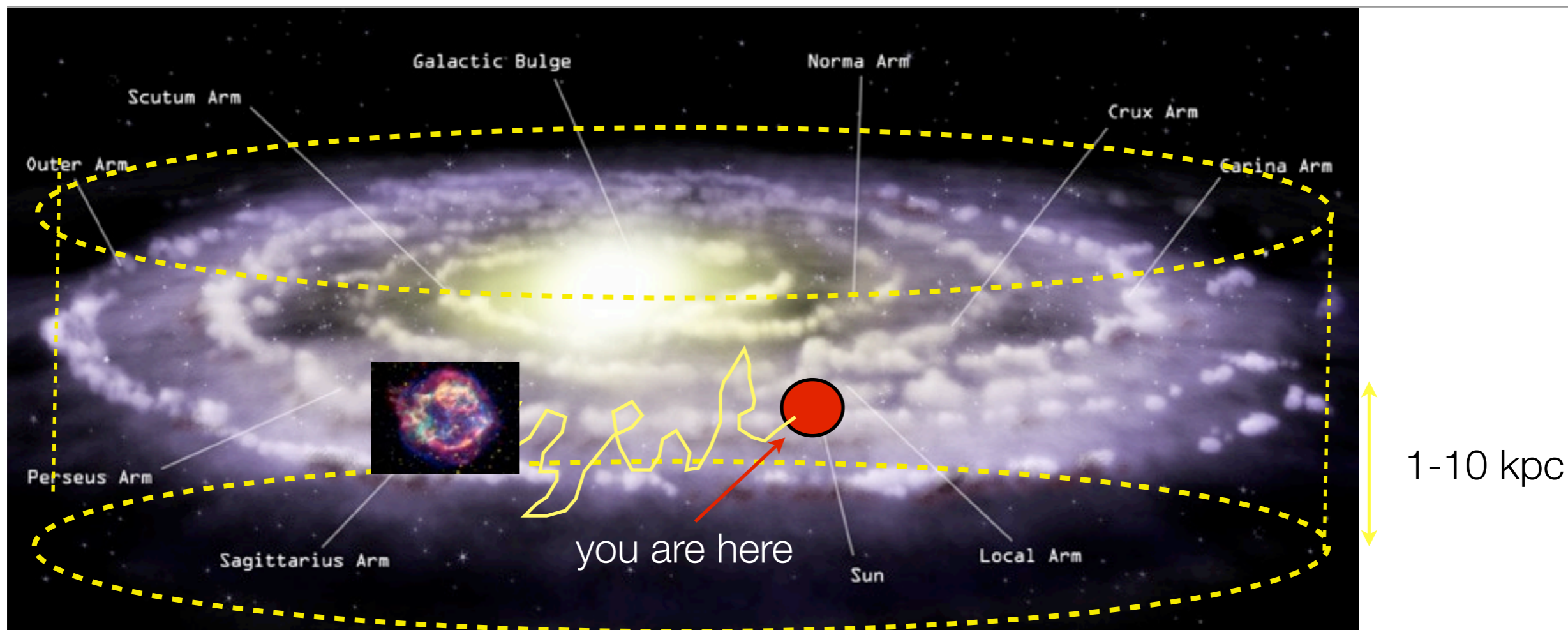
In order to reproduce the measured abundances of stable nuclei, CRs should have traversed: $\sim 10 \text{ g/cm}^2$ material:

$$L = \frac{\text{grammage}}{n_{\text{ISM}} m_p} \sim 10^4 \text{ kpc}$$

\gg Galaxy size!

- **Primary** species are present in sources (CNO, Fe). Produced by stellar nucleosynthesis. Acceleration in SN shocks ($\geq 10^4$ yr).
- **Secondary** species are absent of sources (LiBeB, SubFe). Produced during propagation of primaries.

Galactic Propagation



CRs propagate into the **turbulent** Galactic magnetic field!

The *Larmor* radius of a CR is:

$$r_L(E) = \frac{E}{ZeB} \sim 1 \text{ pc} \left(\frac{E}{10^{15} \text{ eV}} \right) \left(\frac{B}{1 \mu\text{G}} \right)^{-1}$$

for a magnetic field coherence length $\sim 100 \text{ pc}$ \Rightarrow propagation is diffusive up to $\sim 10^{16}$ - 10^{17} eV

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$

$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Source term:

- ▶ assumed to trace the SNR in the Galaxy
- ▶ assumed the same power-law everywhere

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$
$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Spallation cross-section:

- ▶ appearance of nucleus i due to spallation of nucleus j

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$

$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Spallation cross-section:

- ▶ appearance of nucleus i due to spallation of nucleus j
- ▶ total inelastic cross-section: disappearance of nucleus i

Ginzburg & Syrovatsky, 1964

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$

$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Diffusion tensor:

► $D(E) = D_0 (\rho / \rho_0)^\delta \exp(z / z_t)$

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$
$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Energy losses:

- ▶ ionization, Coulomb, synchrotron
- ▶ adiabatic convection

CR Diffusion in the MW

The diffusion equation:

$$\frac{\partial N^i}{\partial t} - \nabla \cdot (D \nabla - v_c) N^i + \frac{\partial}{\partial p} \left(\dot{p} - \frac{p}{3} \nabla \cdot v_c \right) N^i - \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial N^i}{\partial p p^2} =$$
$$Q^i(p, r, z) + \sum_{j>i} c \beta n_{gas}(r, z) \sigma_{ij} N^j - c \beta n_{gas} \sigma_{in}(E_k) N^i$$

Reacceleration:

$$\triangleright D_{pp} \propto \frac{p^2 v_A^2}{D}$$

The ISM turbulence

S. Chandrasekhar, ApJ (1949)

- turbulence type:

HD vs MHD

- assuming constant energy flow (η) between eddies (of length k^{-1}):

$$\epsilon = f(\eta, k) \quad \text{vs} \quad \epsilon = f(\eta, k \cdot v_k/v_A)$$

- a simple dimensional analysis gives for the energy spectrum:

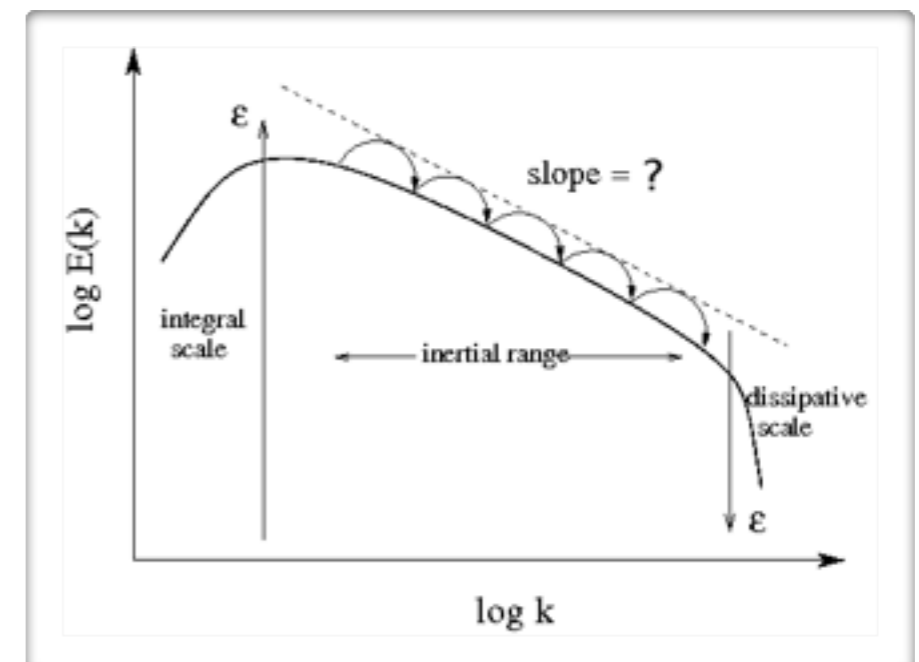
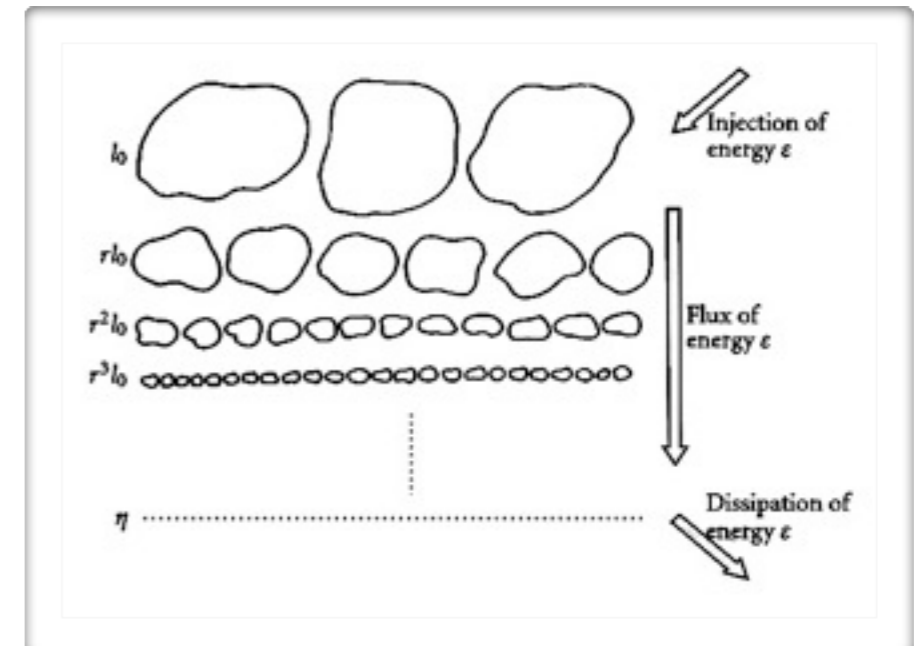
$$\epsilon \propto \eta^{2/3} k^{-5/3} \quad \text{vs} \quad \epsilon \propto \eta^{1/2} v_A^{1/2} k^{-3/2}$$

- finally, diffusion coefficient in QLT is given by:
(i.e. imposing the resonance condition $r_L = 1/k$)

$$D = \rho^{2-5/3} = \rho^{1/3} \quad \text{vs} \quad D = \rho^{2-3/2} = \rho^{1/2}$$

“Kolmogorov”

“Kraichnan”

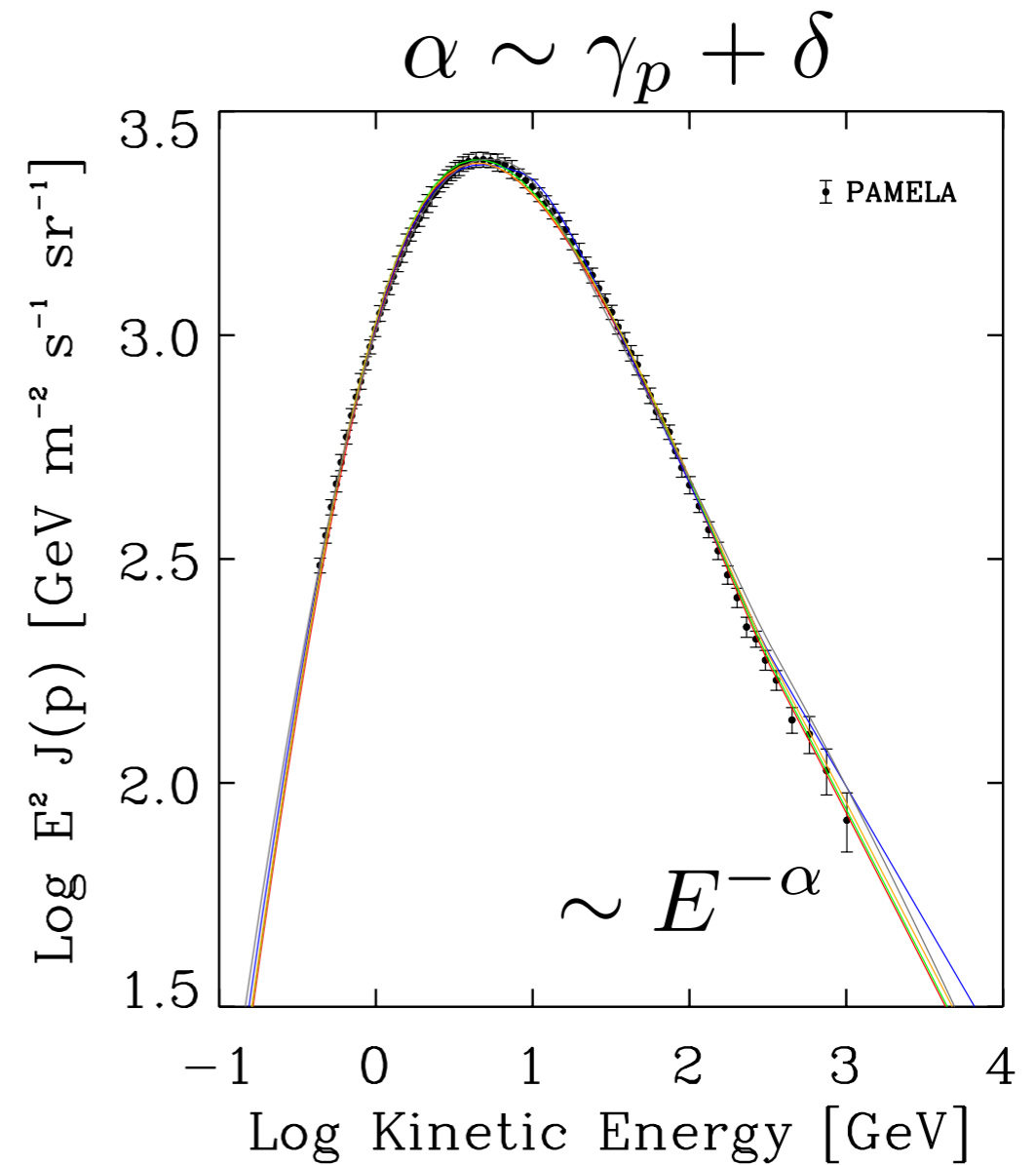
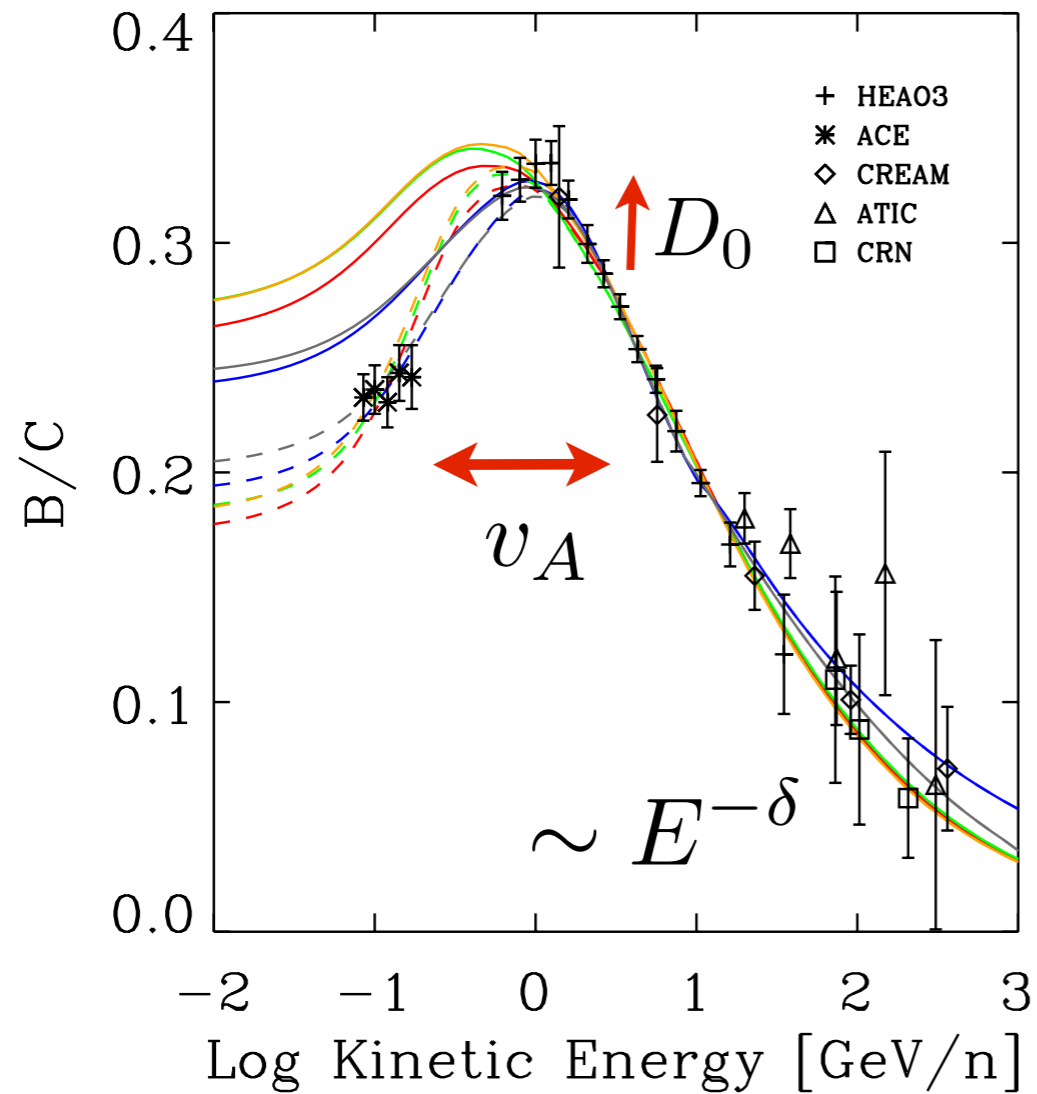




“Local” observables

CE, I.Cholis, D.Grasso, L.Maccione & P.Ullio, PRD, 2012, 1108.0664

Kolmogorov / Kraichnan



Is it possible being not-local?

- we can measure the anisotropy:

$$\delta \propto \nabla n_{\text{cr}}$$

- we can observe diffuse emissions:

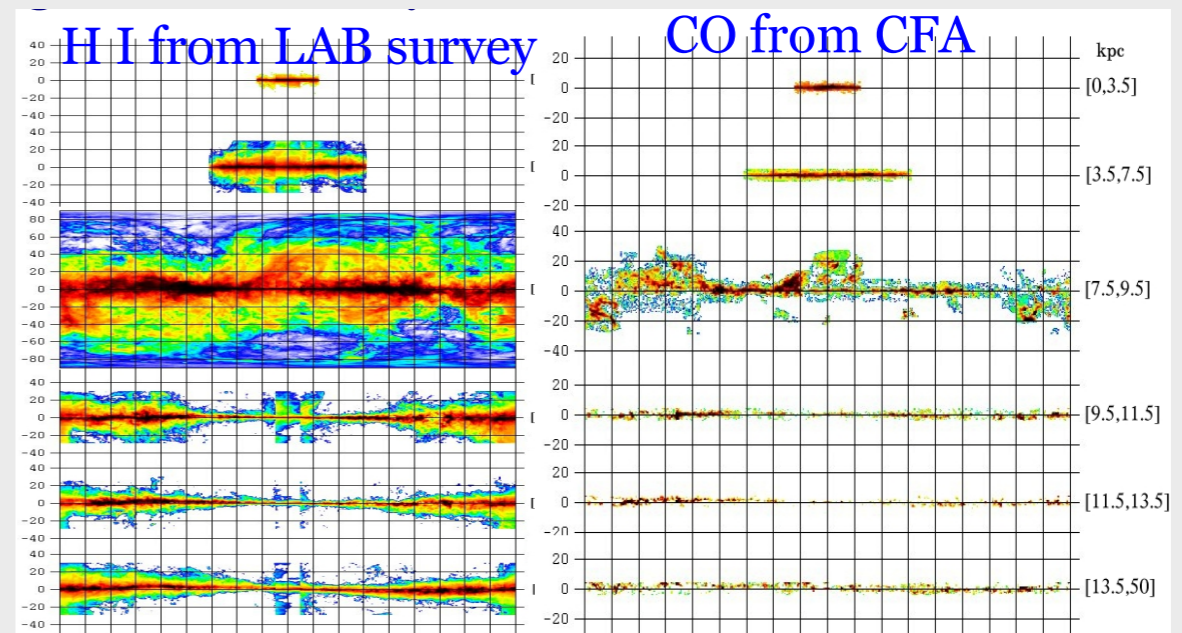
$$\phi_{\gamma} \propto \int_{l_{\text{os}}} n_{\text{cr}} \cdot n_{\text{gas}} dr$$

Atomic (HI):

Most massive component with a large filling factor, $z_{1/2} \sim 200$ pc.

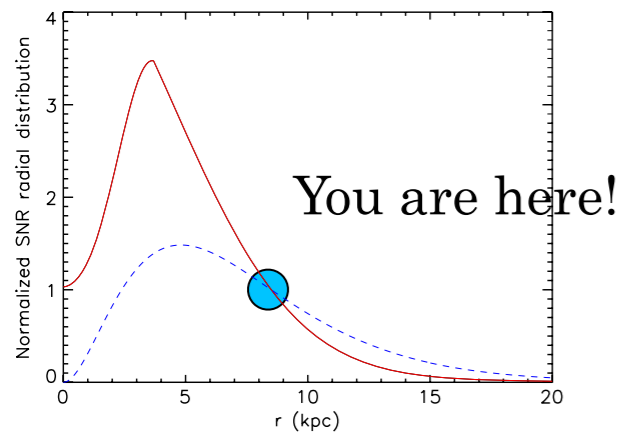
Molecular (H₂):

The most dense component, very clumpy, $z_{1/2} \sim 100$ pc (derived from the **CO!**)



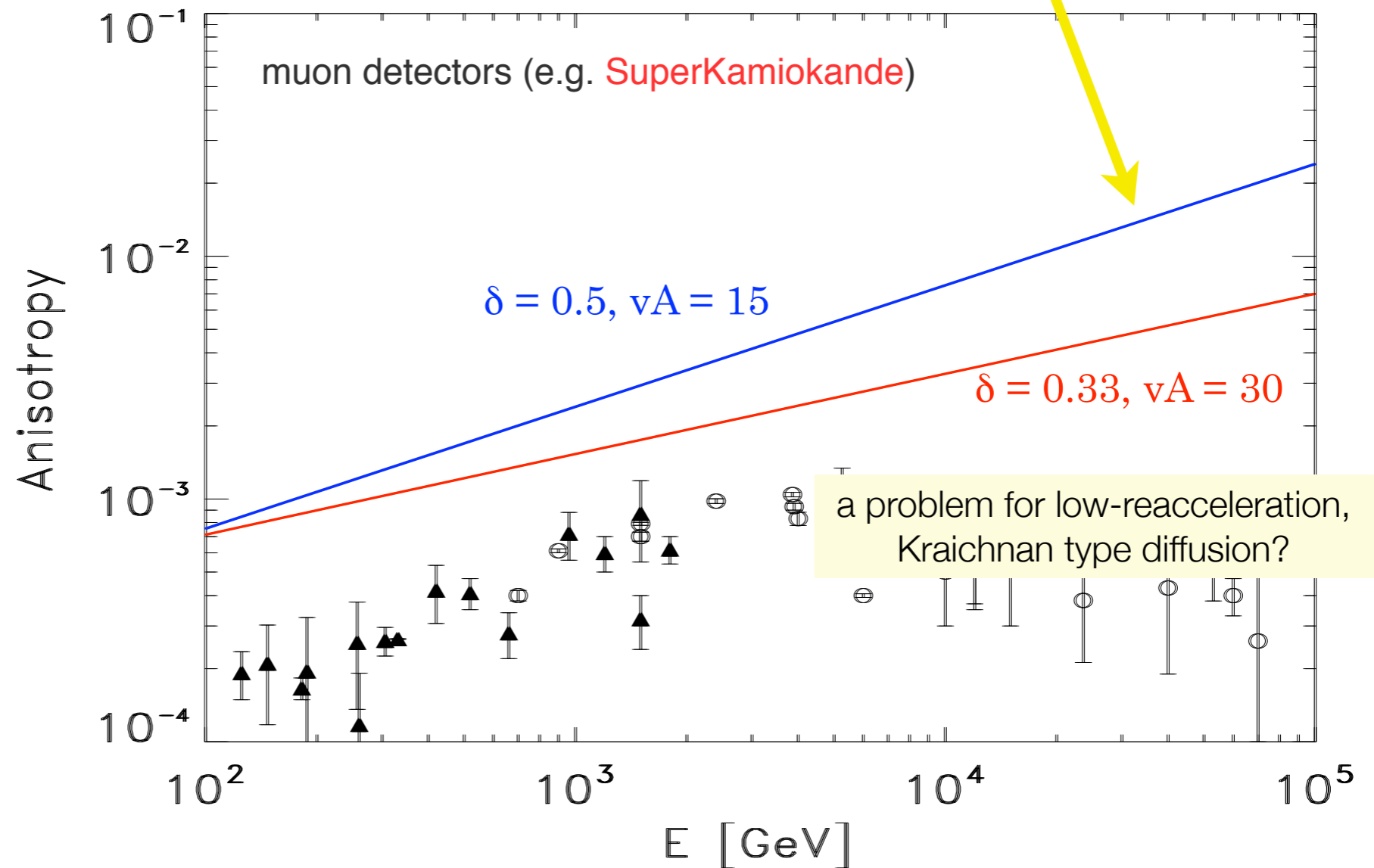
The *anisotropy* problem

Macro Collaboration, PRD, 2003; Super-Kamiokande Collaboration, PRD, 2007



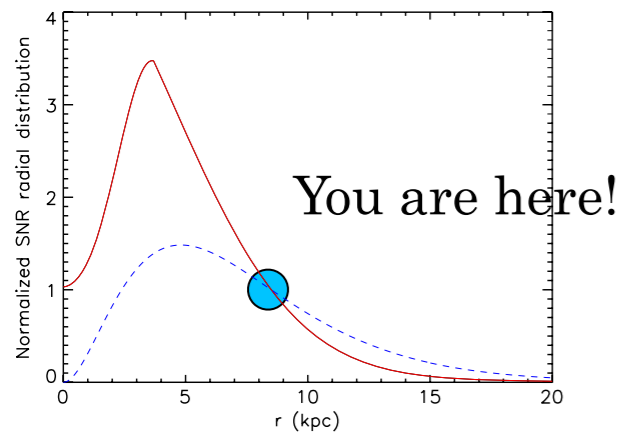
Since CR sources are more abundant in the inner Galaxy, a dipole anisotropy is expected towards the Galactic center:

$$\delta_{\vec{x}} = \frac{3D(E)}{c} \frac{\nabla_{\vec{x}} n_{CR}(E, \vec{r}, t)}{n_{CR}}$$



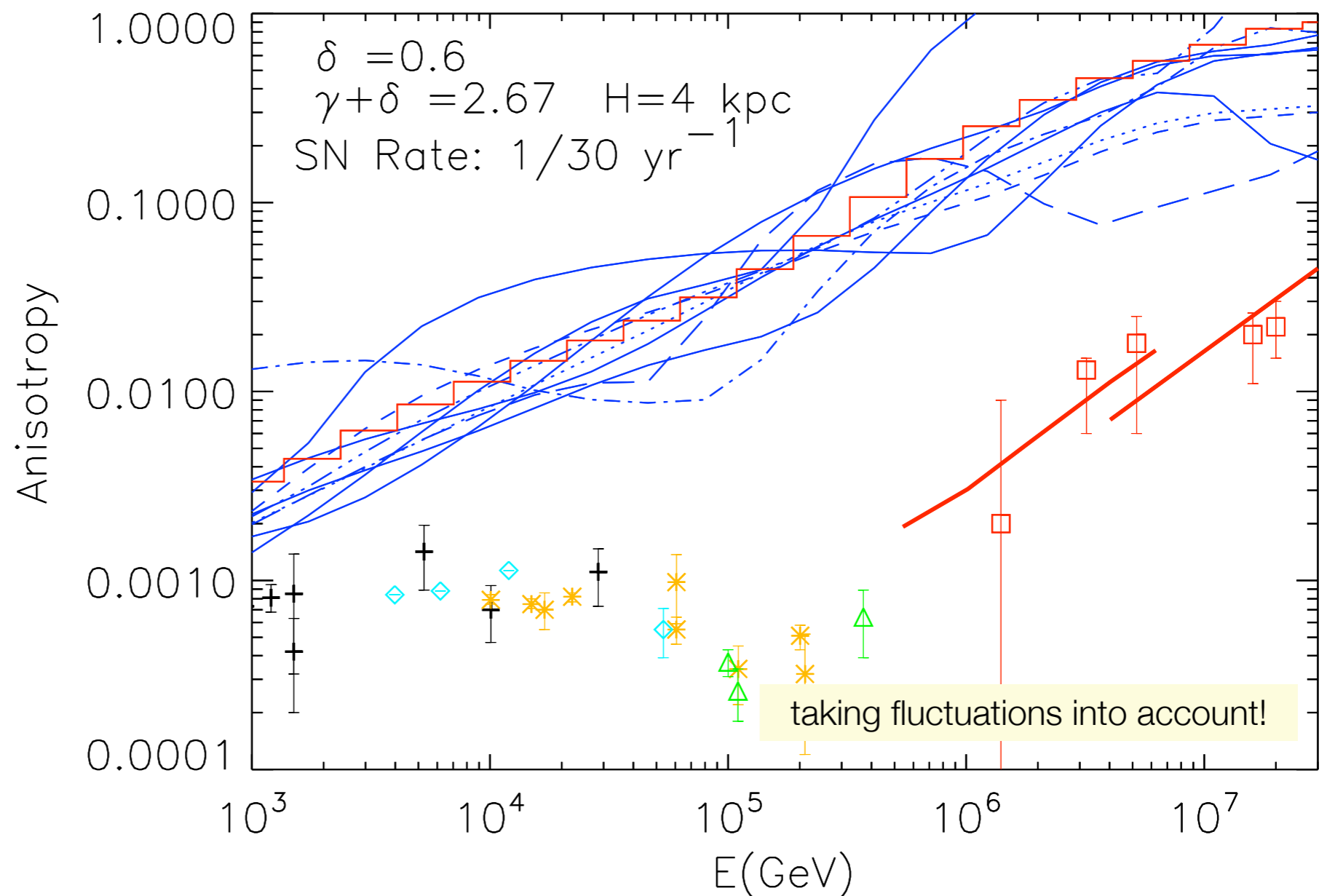
The *anisotropy* problem

P. Blasi and E. Amato, JCAP, 2012



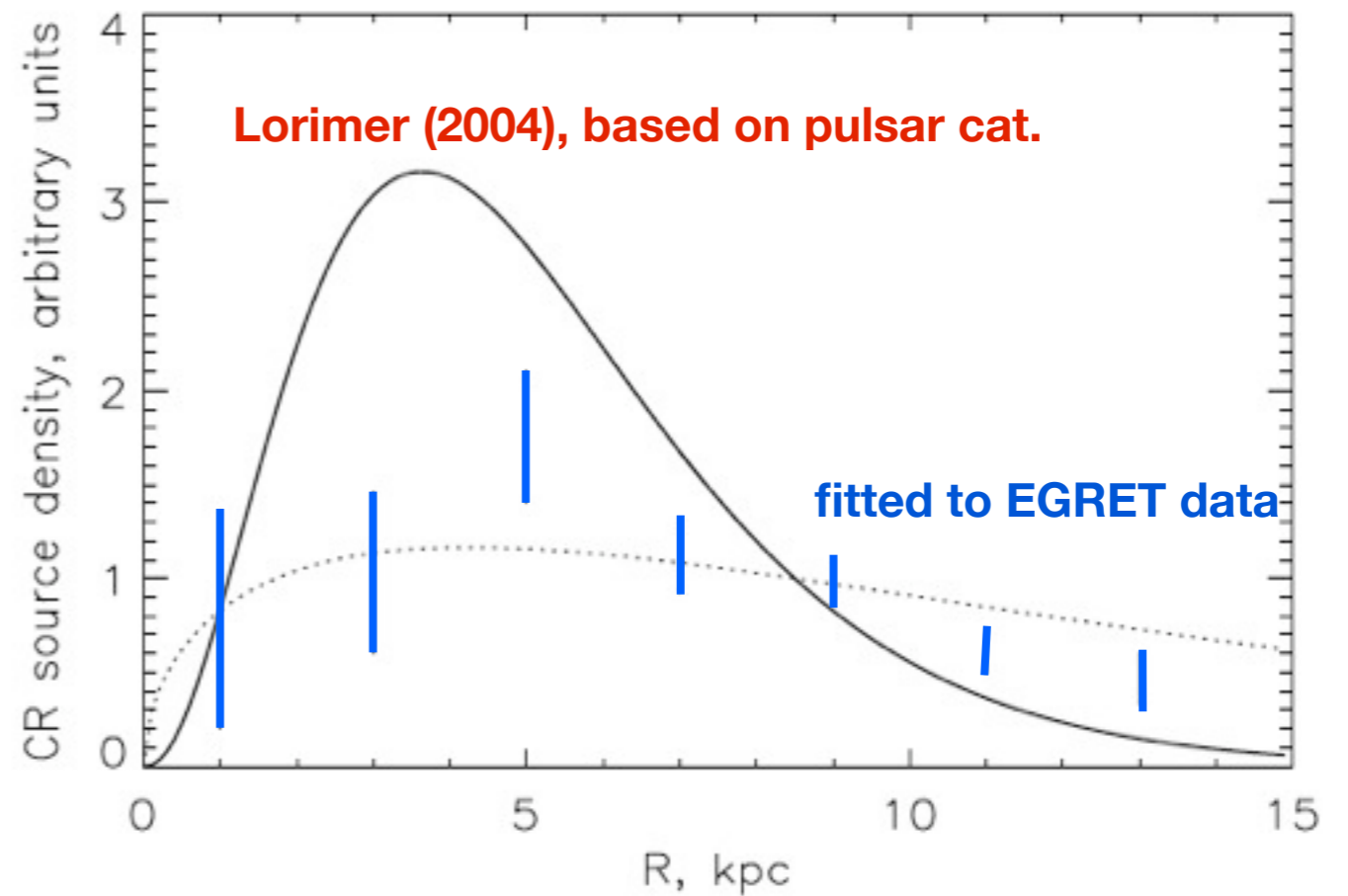
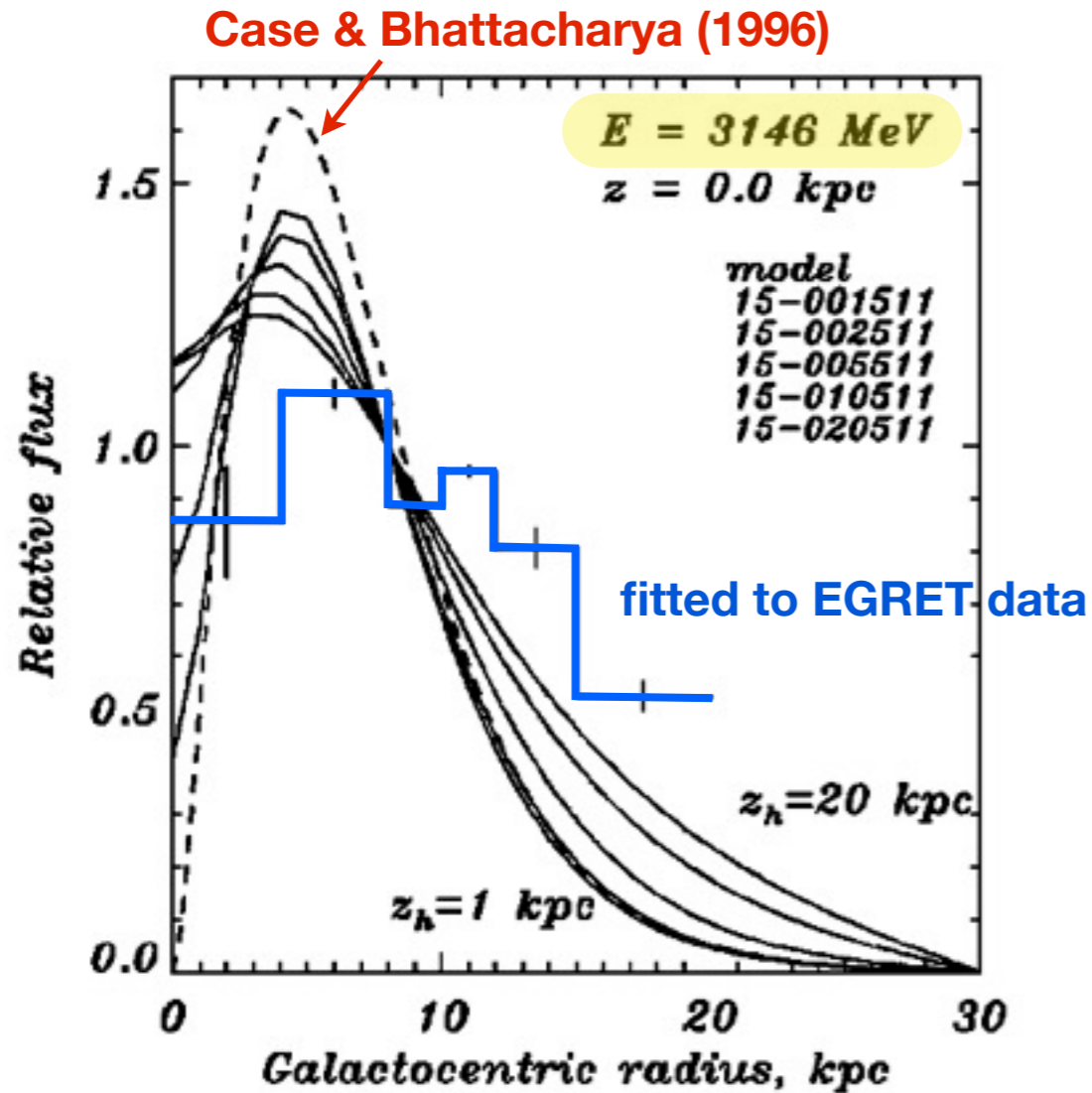
Since CR sources are more abundant in the inner Galaxy, a dipole anisotropy is expected towards the Galactic center:

$$\delta_{\vec{x}} = \frac{3D(E)}{c} \frac{\nabla_{\vec{x}} n_{CR}(E, \vec{r}, t)}{n_{CR}}$$



The *gradient* problem

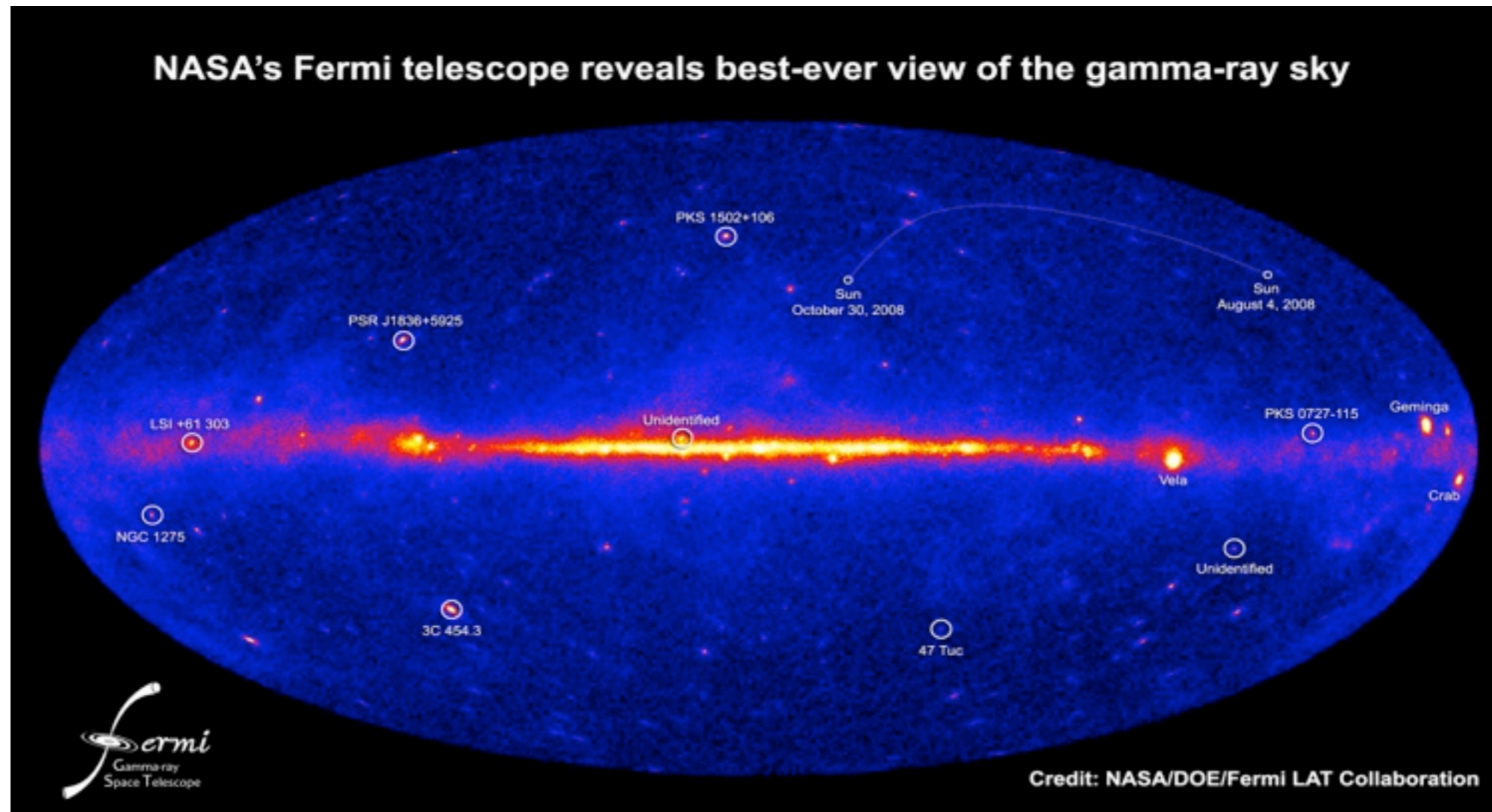
Strong & Mattox, A&A, 1996; Strong et al., ApJ, 2000



- CR distribution inferred from gamma-ray data (method goes back to SAS-2/COS-B era) is **flatter** than that computed assuming the observed **SNR** (source) profile.

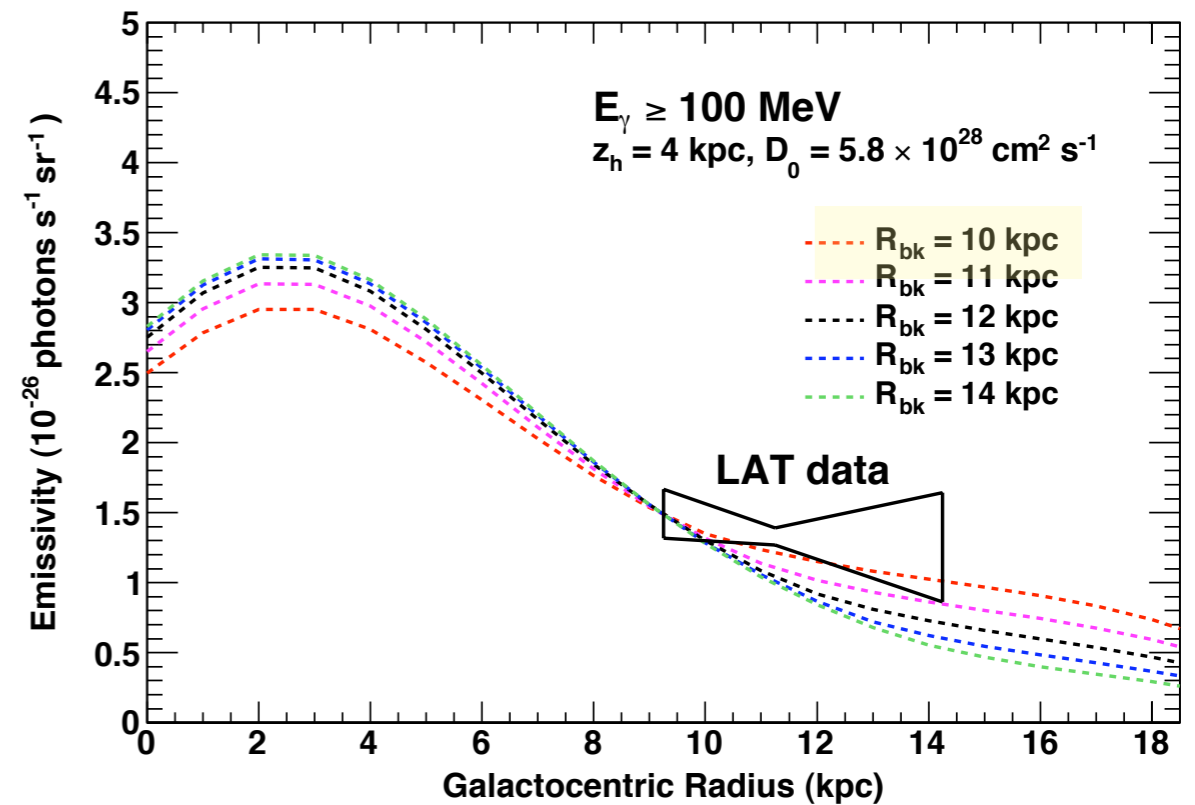
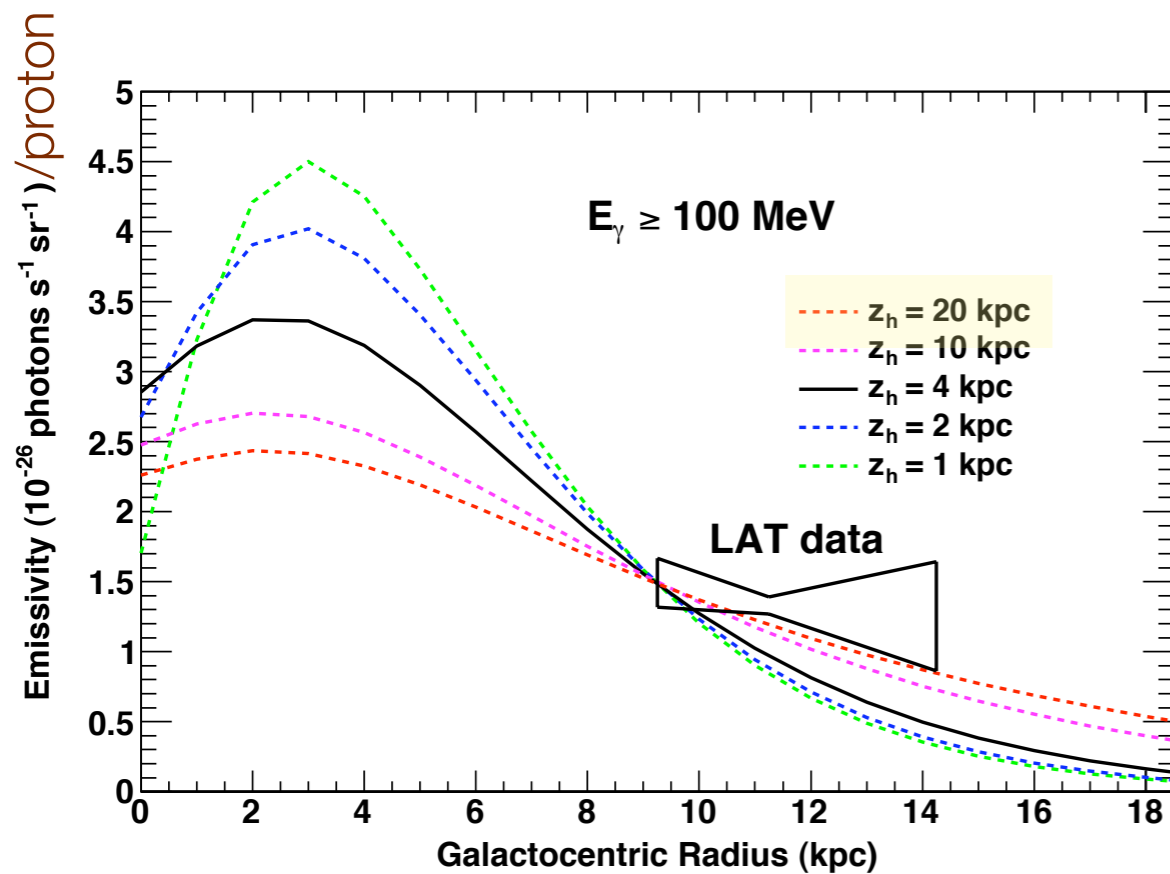
The *gradient* problem in the FERMI era

- The extremely accurate gamma ray maps that Fermi is providing are useful to trace the CR distribution throughout all the Galaxy!



The *gradient* problem in the FERMI era

Fermi Collaboration, ApJ, 2011



FERMI detected **more** γ 's than a prediction based on SNR distribution and standard CR halo:
more CR sources, more "dark gas" or larger CR halo?

A new approach

D. De Marco, P. Blasi & S. Todor, JCAP, 2007

- there are some regions in the **inner** Galaxy where a much higher density of CR sources – and hence turbulence – is present.
- According to both quasi-linear theory and numerical simulations:

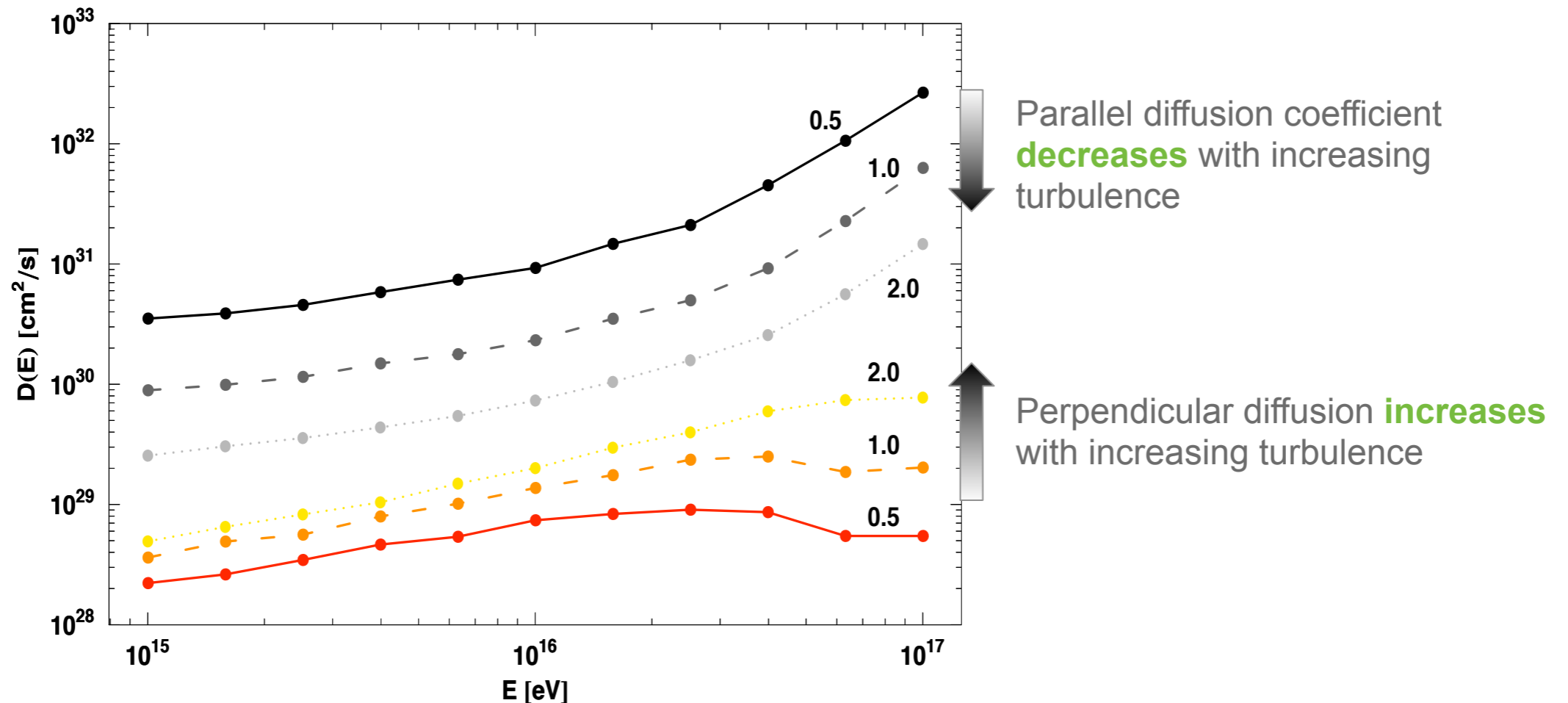
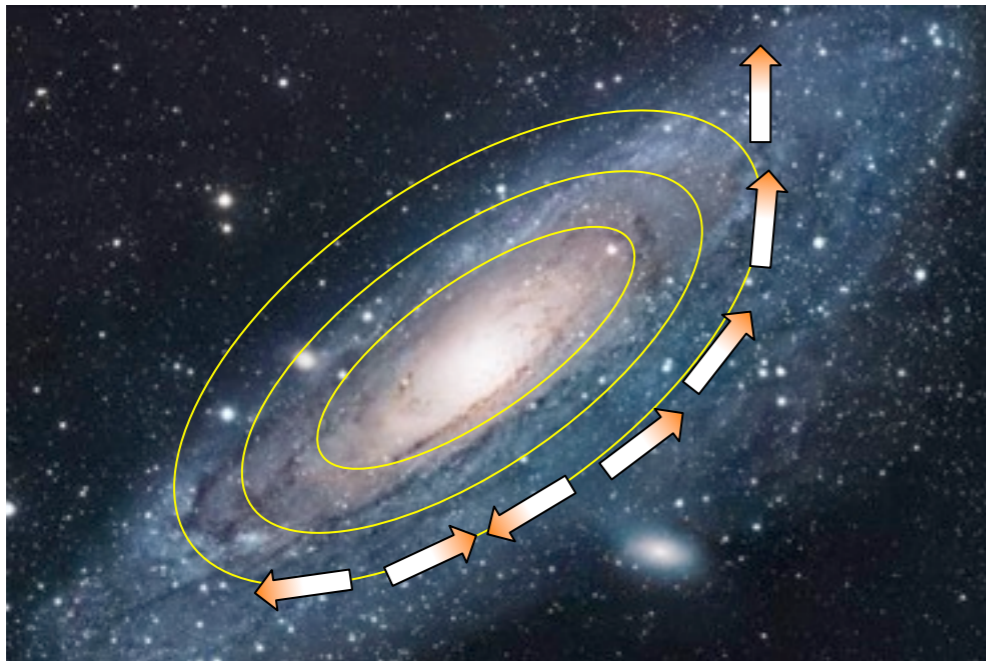


Figure 3. Parallel and perpendicular diffusion coefficients as a function of energy for three levels of turbulence. The upper three lines are the parallel diffusion coefficients, while the bottom three represent the perpendicular one. The level of turbulence, $\delta B/B_0$ is given by the numbers attached to the lines.

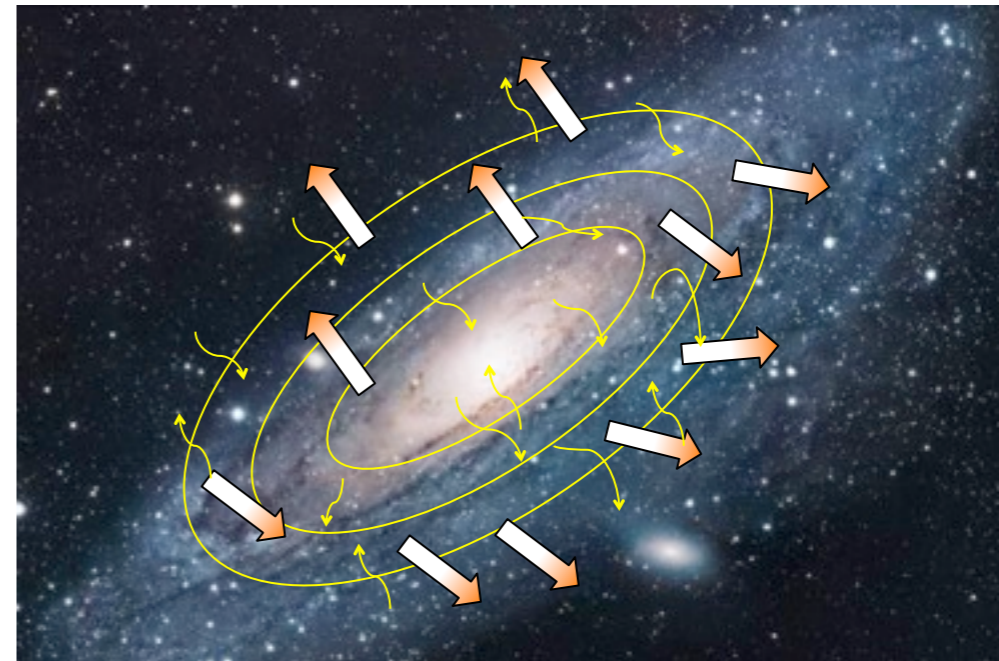
A new approach

- How do the diffusion coefficient depends on turbulence?

If the turbulent field is very low:



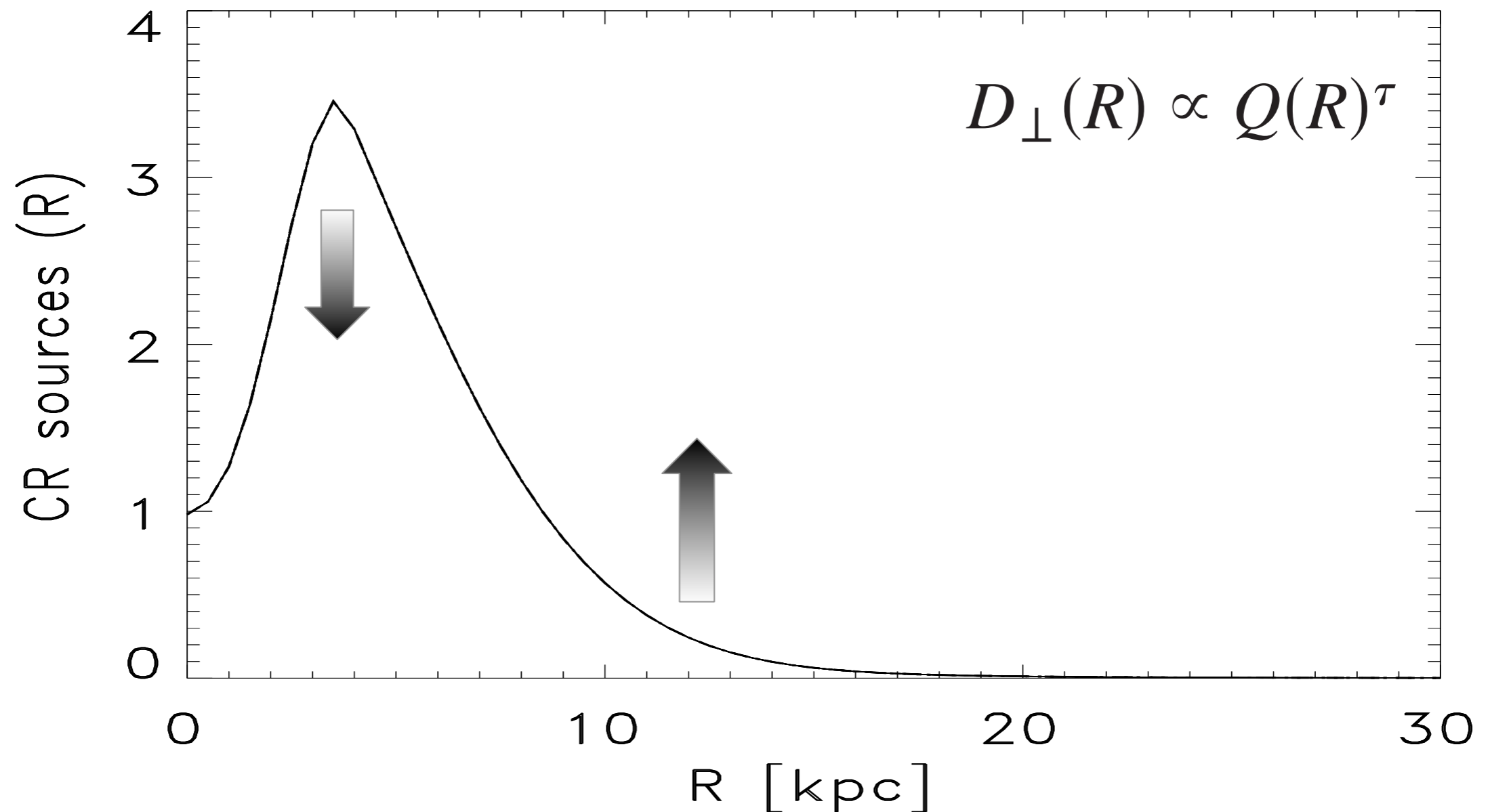
If the turbulent component is comparable to the regular field:



- In the inner galaxy, where turbulence is high, the parallel and perp. diffusion are similar values and the perpendicular escape is the dominant one:

$$\frac{T_{\parallel}}{T_{\perp}} \simeq \left(\frac{R_{\text{arm}}}{H} \right)^2 \quad \frac{D_{\perp}}{D_{\parallel}} \simeq 4 \times 10^2 \left(\frac{H}{4 \text{ kpc}} \right)^{-2} \frac{D_{\perp}}{D_{\parallel}}$$

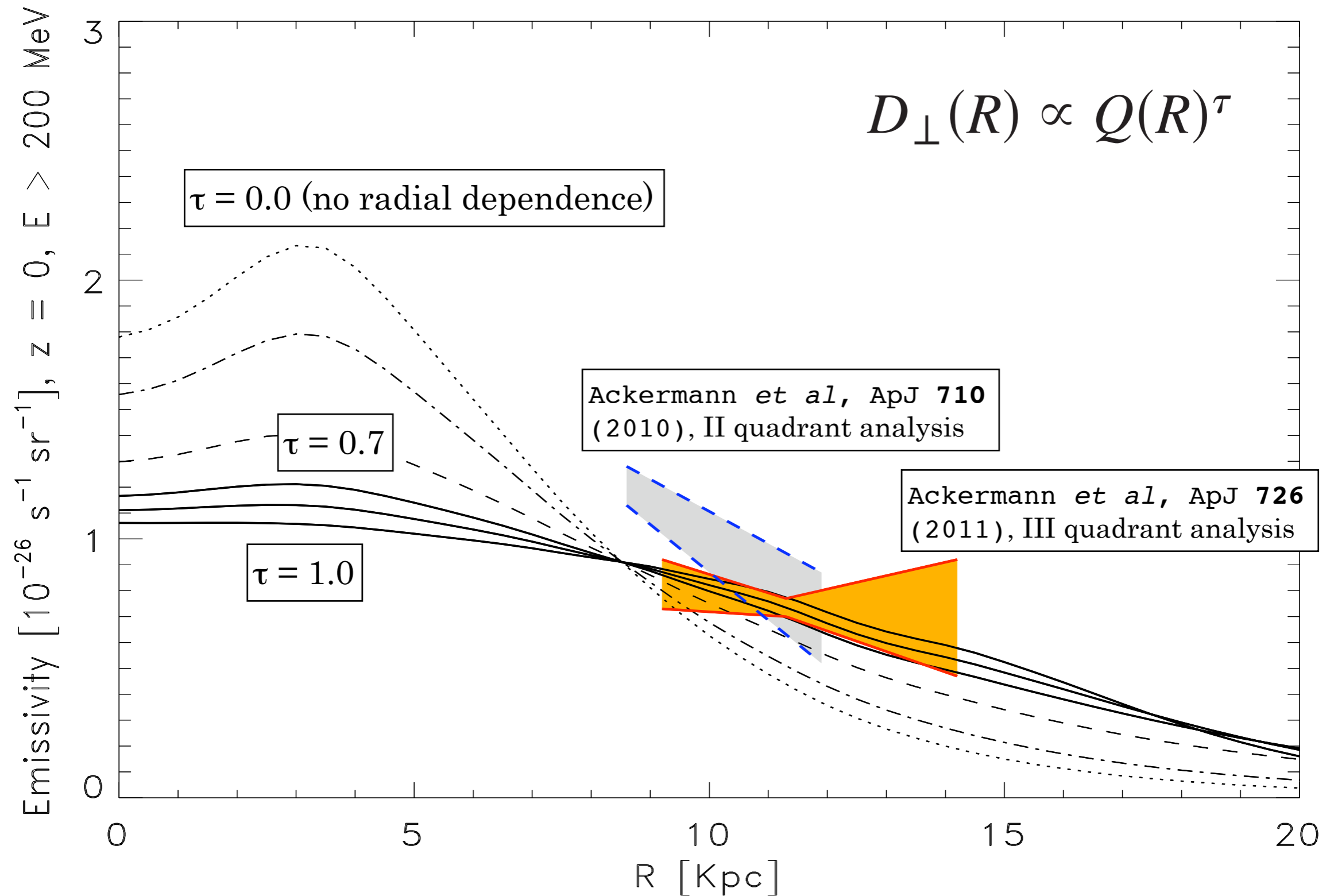
How to solve the gradient problem



- In the regions where CR sources are more abundant turbulence is higher then perpendicular escape is faster, more CR are removed.

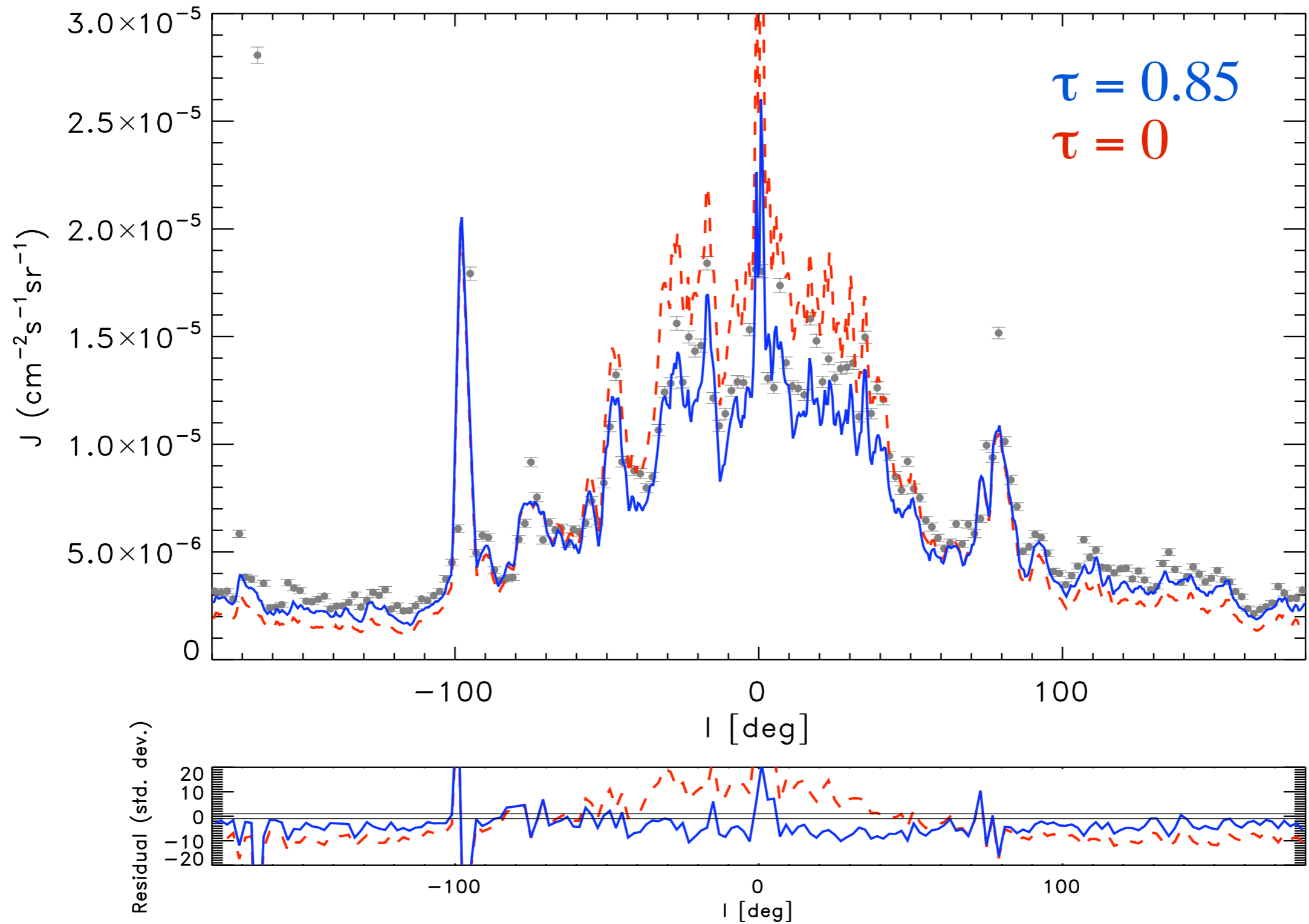
Results

CE, D. Gaggero, D. Grasso & L. Maccione, PRL, 2012



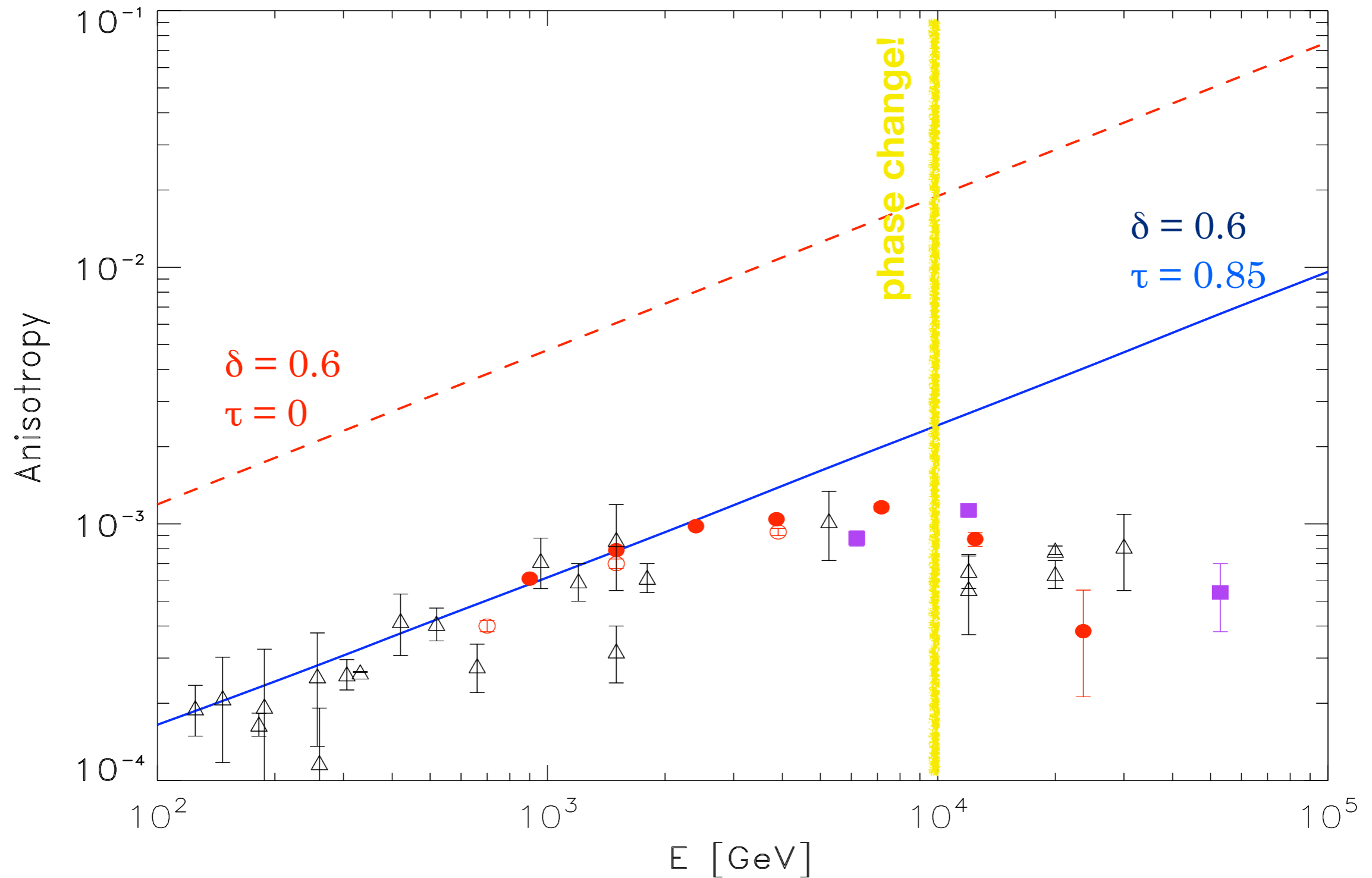
Results

CE, D. Gaggero, D. Grasso & L. Maccione, PRL, 2012



Anisotropy prediction

CE, D. Gaggero, D. Grasso & L. Maccione, PRL, 2012



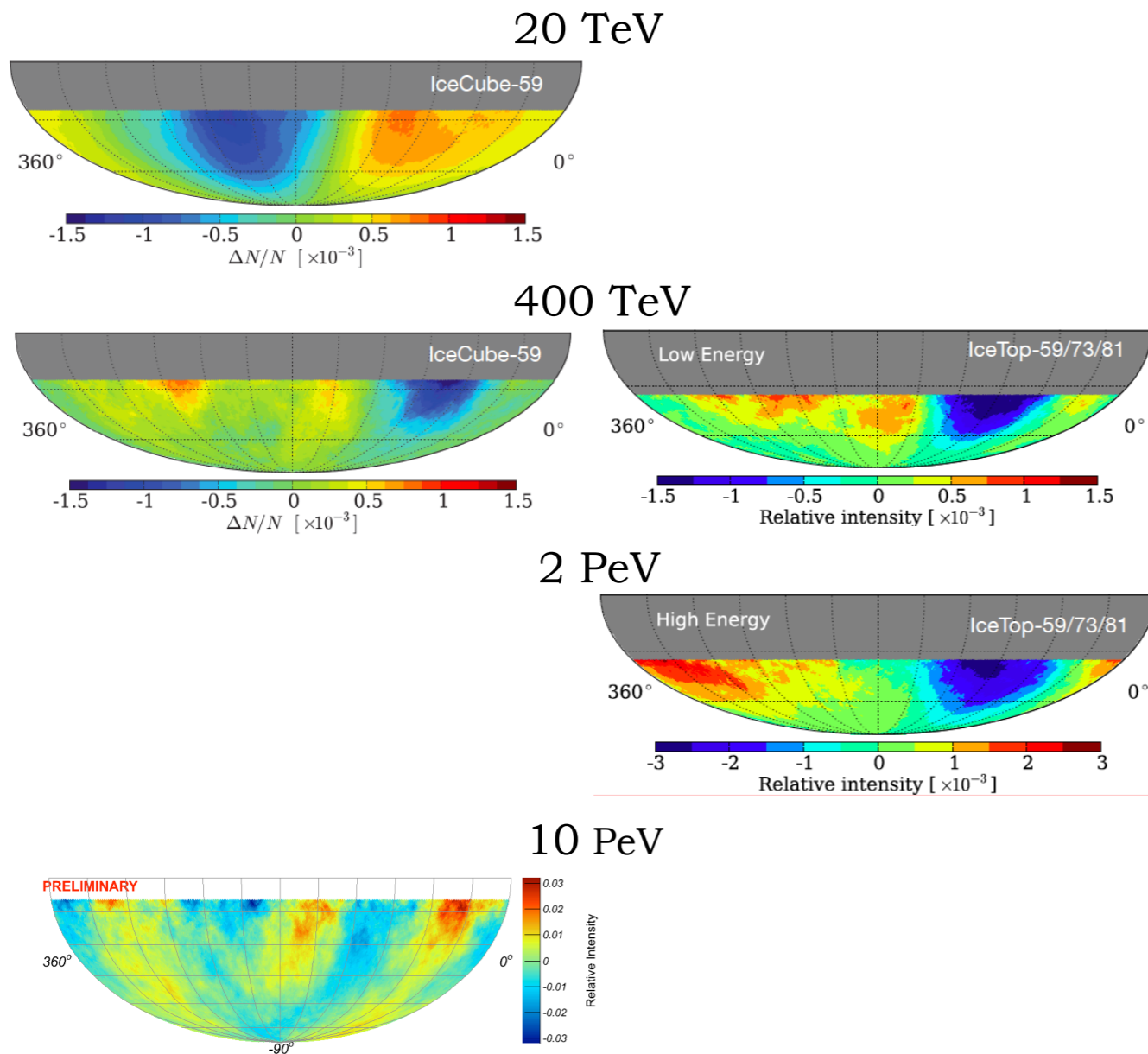
Large Scale Anisotropy with IceCube / IceTop

IceCube
muon bundles > 1 TeV

IceTop
CR showers > 100 TeV

IceCube, ApJ 746, 33 (2012)

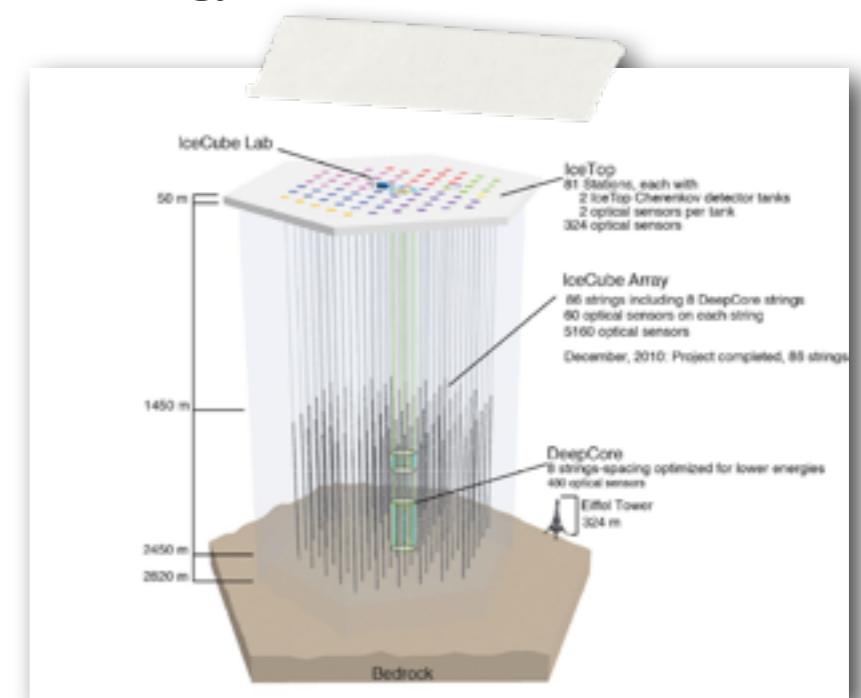
IceCube, ApJ 765, 55 (2013)



topology changes
between 20 - 400 TeV

anisotropy is not dipole

amplitude increases
with energy



Is it possible to move further from QLT?

H. Yan & A. Lazarian, ApJ, 2008

- particle's pitch angle follows the variation of the *turbulent* magnetic field due to conservation of the adiabatic invariant:

$$\frac{\Delta v_{\parallel}}{v_{\perp}} = \frac{\langle (B - B_0)^2 \rangle^{1/4}}{B_0^{1/2}}$$

- resonance function has a Gaussian broadening:

$$R_n^{\text{NLT}}(k_{\parallel} v_{\parallel} - \omega \pm n\Omega) = \frac{\sqrt{\pi}}{k_{\parallel} \Delta v_{\parallel}} \exp \left[-\frac{(k_{\parallel} v_{\parallel} - \omega \pm n\Omega)^2}{k_{\parallel}^2 \Delta v_{\parallel}^2} \right]$$

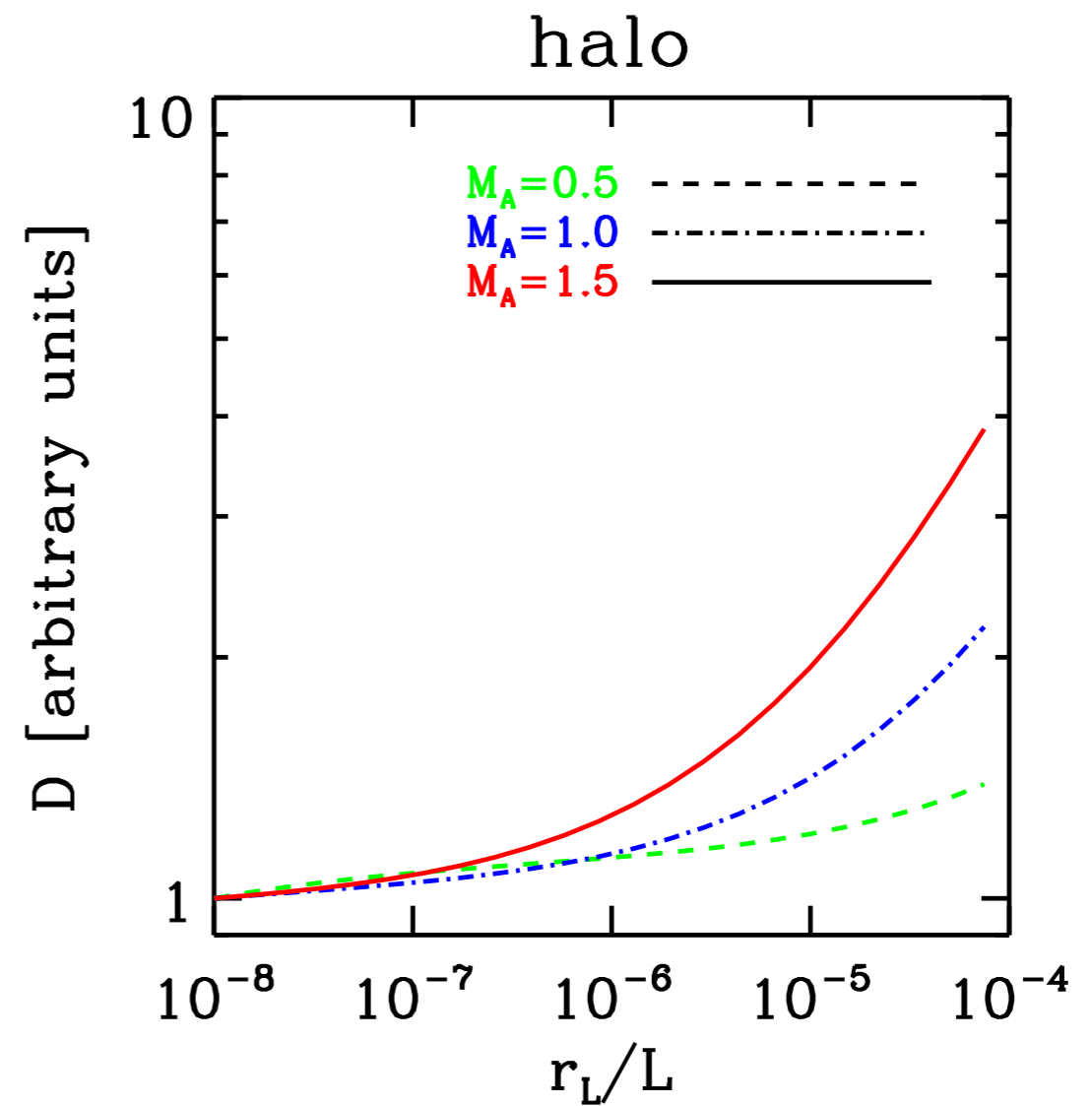
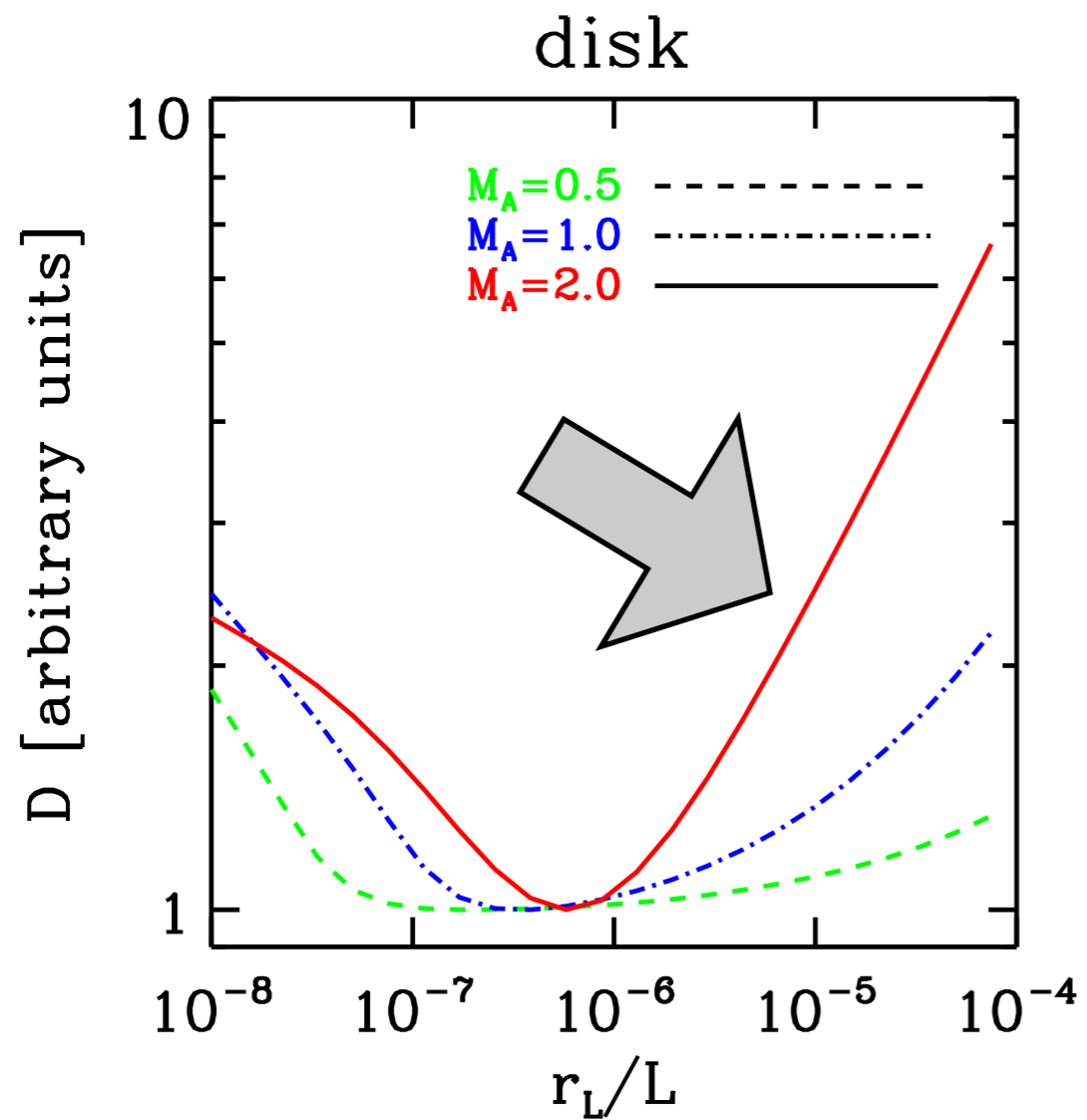
- damping mechanisms make diffusion environment-dependent:

$$D_{\mu\mu} = \frac{\Omega^2(1 - \mu^2)}{B_0^2} \int d^3 k R_n^{\text{NLT}}(\mathbf{k}) \left[\frac{k_{\parallel}^2}{k^2} J_n'^2(w) I^F(\mathbf{k}) \right]$$

Results (preliminary!)

halo = collisionless damping
disk += viscous damping

CE & H. Yan, *in preparation*



Results (preliminary!)

halo = collisionless damping
disk += viscous damping

CE & H. Yan, *in preparation*

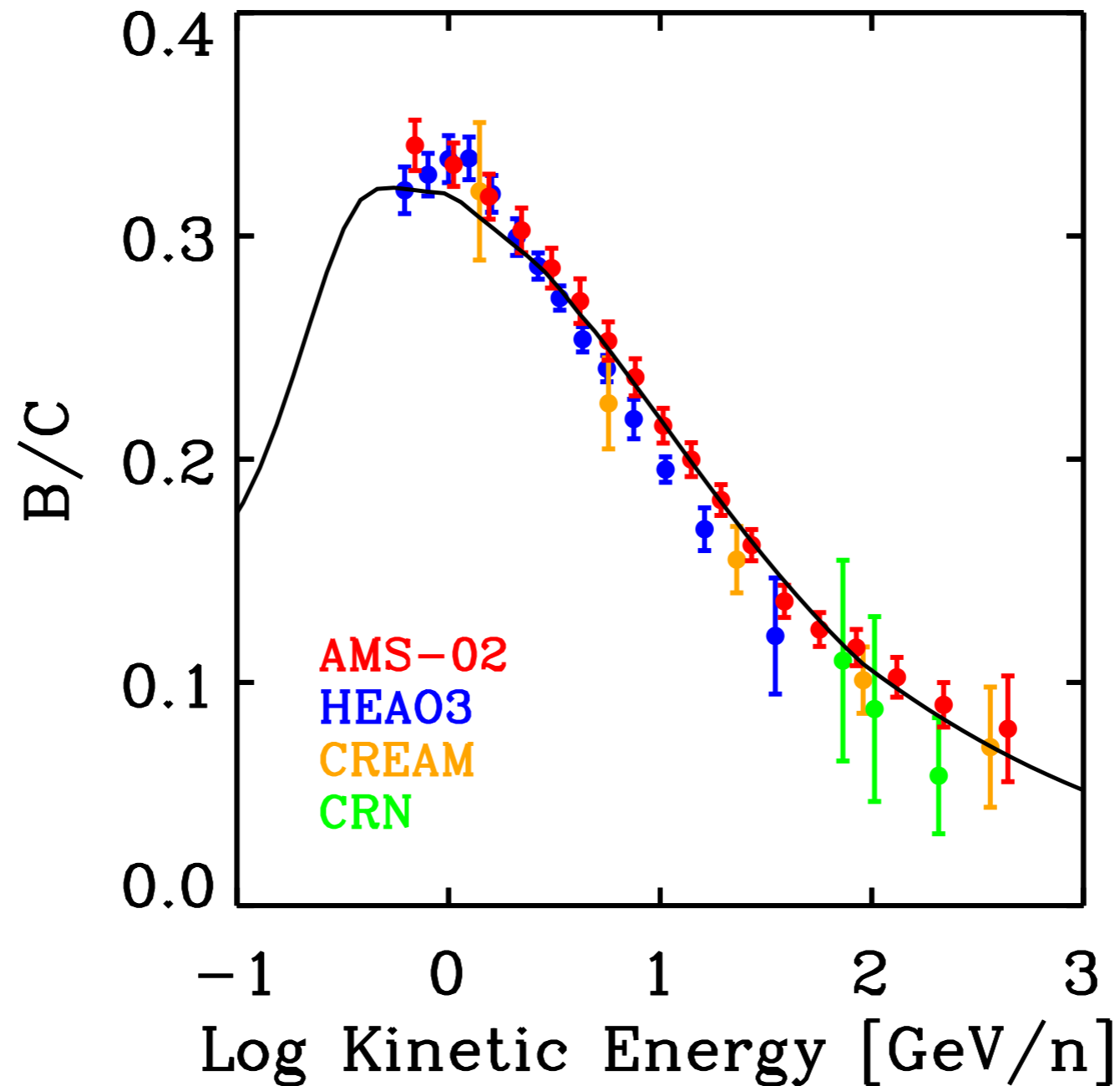


Figure 4. Comparison of our model with $M_A = 2$ for the disk and $M_A = 1$ in the halo and modulated with a 100 GV potential against B/C data. See Table 1 for a reference list of the experimental data.

Conclusions I

- A **position-dependent** perpendicular diffusion coefficient that traces regions of the Galaxy where turbulence is higher offers a natural explanation of not-local CR observations.
- Upcoming (**PLANCK**, LOFAR...) **synchrotron data** can further support this scenario.
- Future diffusion codes should take into account the complexity of the Galaxy, and allow full **3D** simulations, **anisotropic diffusion**, and realistic distributions of **sources**, gas, magnetic fields, especially in the local environments.

The logo features a stylized dragon head in profile, facing right, with a circular pattern of dots on its snout. The word "DRAGON" is written in a large, black, serif font to the right of the dragon head.

DRAGON

- ▶ solve the **diffusion equation** on a 3D (r,z,E) grid (now also **4D!**)
- ▶ realistic distributions for **sources** and **ISM**
- ▶ different models for fragmentation **cross sections**
- ▶ position dependent, **anisotropic** diffusion
- ▶ independent injection spectra for each nuclear species
- ▶ speed and memory high-performances (full C++)
- ▶ **public**: <http://dragon.hepforge.org>

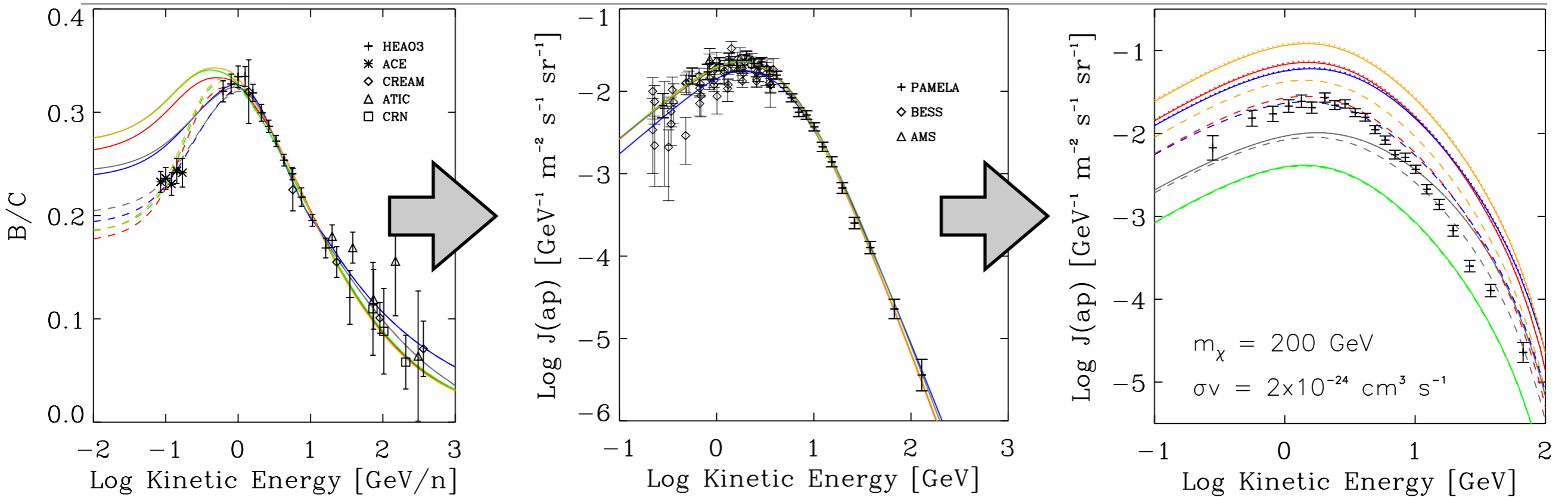
see Daniele's talk!

Why should I care about it?



Playing with anti-protons from DM

CE, I.Cholis, D.Grasso, L.Maccione & P.Ullio, PRD, 2012, 1108.0664



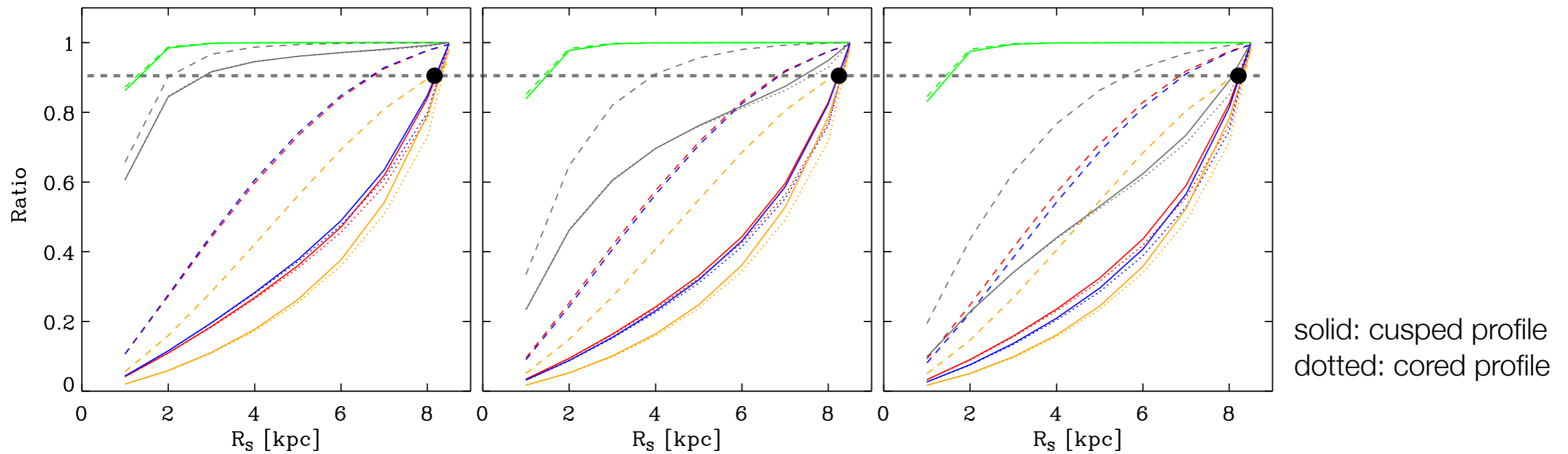
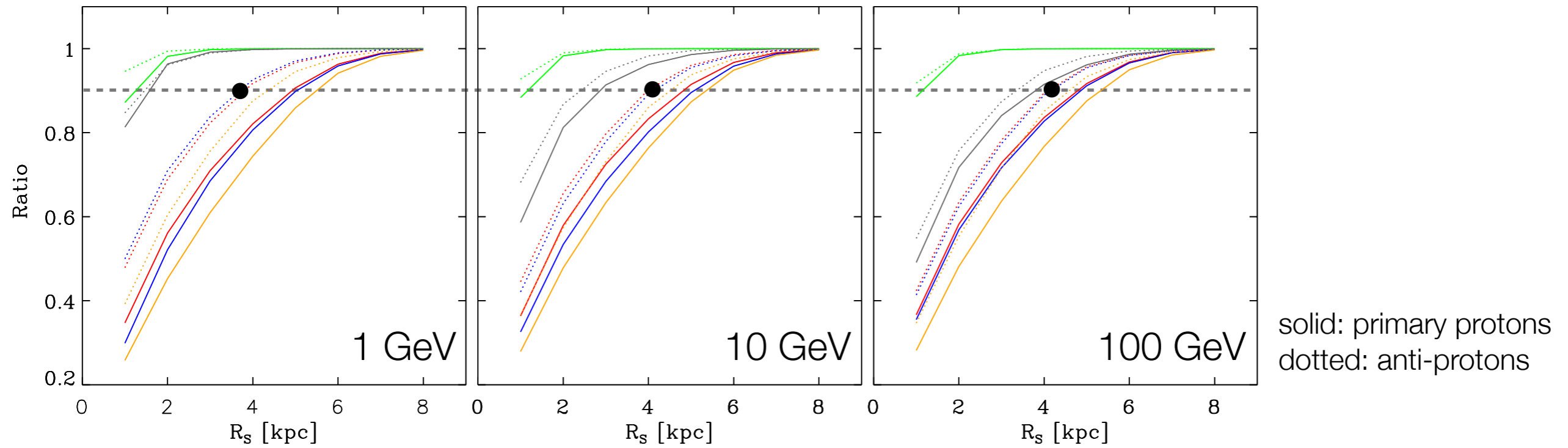
30%

~1 oom

Model	z_t (kpc)	δ	$D_0(10^{28} \text{ cm}^2/\text{s})$	η	v_A (km/s)	γ	dv_c/dz (km/s/kpc)	$\chi_{B/C}^2$	χ_p^2	Φ (GV)	χ_p^2	Color in Figs.
KRA	4	0.50	2.64	-0.39	14.2	2.35	0	0.6	0.47	0.67	0.59	Red
KOL	4	0.33	4.46	1.	36.	1.78/2.45	0	0.4	0.3	0.36	1.84	Blue
THN	0.5	0.50	0.31	-0.27	11.6	2.35	0	0.7	0.46	0.70	0.73	Green
THK	10	0.50	4.75	-0.15	14.1	2.35	0	0.7	0.55	0.69	0.62	Orange
CON	4	0.6	0.97	1.	38.1	1.62/2.35	50	0.4	0.53	0.21	1.32	Gray

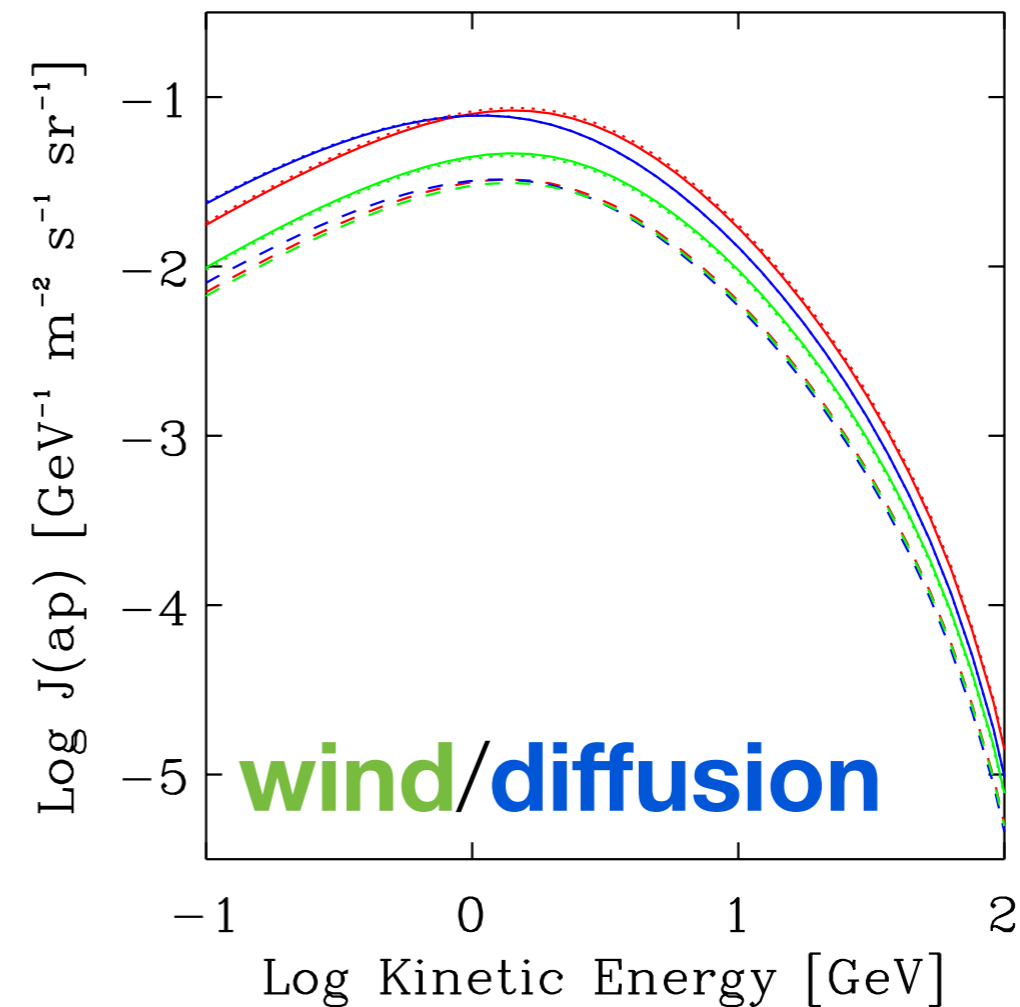
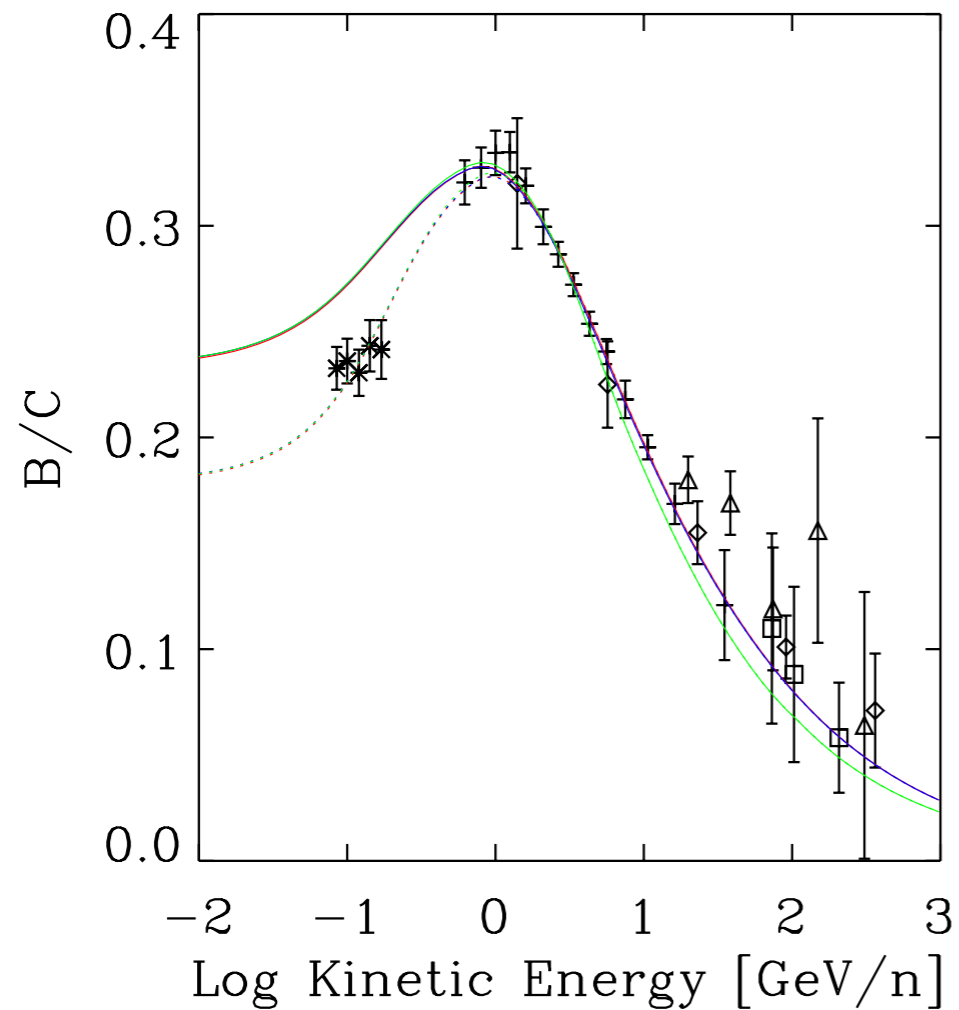
It's all about locality!

CE, I.Cholis, D.Grasso, L.Maccione & P.Ullio, PRD, 2012, 1108.0664



Unavoidable uncertainties?

CE, I.Cholis, D.Grasso, L.Maccione & P.Ullio, PRD, 2012, 1108.0664



multiwave/messenger is the solution!

Conclusions II

- If your CR source is the “galactic-center” do not trust too much “global models” that are “locally” tuned...
- ...or at least “put a warning” about uncertainties due to not-locality.
- A lot of work is still ahead for understanding ISM propagation **before** to be robust in DM searches with charged particles!