Introduction to the EXAFS data analysis

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Material almost integrally taken from <u>Carlo Meneghini</u>: EXAFS tutorial at Synchrotron Radiation school of Duino 2011

Characteristics of a XAS spectrum



E(eV)

XAFS study: from experiment to results



XAFS study: from experiment to results



Considerations:

- 1) Proposal submission + proposal evaluation + beamtime scheduling = 6 to 12 months
- 2) Difficult to have new beamtime in case of proposal failure

- Check the proposal submission deadlines

- discuss your experiment with local contacts



Choose properly the experimental set-up & sample preparation

• For massive concentrated samples: TRANSMISSION



inhomogeneities, holes, not parallel surfaces, etc...

• For thin concentrated or thin diluted samples: FLUORESCENCE



Self absorption, detector linearity, Bragg reflections

Choose properly the data collection strategy

- Acquisition time per point
- Single scan or repeated scans
- ΔE or Δk step





Measure reference samples



Check data quality constantly during the experiment

• Evaluate signal/noise ratio



Check data quality constantly during the experiment

• Check for:





XAFS analysis: from experiment to results



Choose the best spectra and useful data regions



do not use data beyond 13000 eV!

De-glitch



Align





Average



Preliminary data treatment is boring, it may be long... While you are waiting for your data collection to finish...

Do it on already collected data!! You will save your time at home!!

XAFS analysis: from experiment to results





pre-edge line + post-edge line

Normalized data









 μ_{o} calculation

1) Define E_0



- E_0 will allow to set the starting point of $\chi(k)$.
- It is generally taken at the maximum of the 1st derivative of the absorption

2) Calculate μ_0

 μ_{o} is the bare atom atomic background.

It is calculated empirically as a smooth curve across the data.

Different XAFS data analysis softwares apply different (<u>equivalent</u>) approaches

3) Subtract μ_0 from μ

Fourier Transform



FT shows more intuitively the main structural features in the real space: the FT modulus represent a pseudo-radial distribution function (RDF)

|FT| peaks represent interatomic correlation

Peak position are not the true correlation distances due to the phase shift effect

Fourier Transform - window size effect













Minor effects are given by type of windows (Hanning, Kaiser-Bessel, Sine) and apodization



DO NOT DO THE OPPOSITE ERROR



R (Å)

Fourier Filtering

Fourier filtering allows isolating contributions of selected regions of the FT



XAFS analysis: from experiment to results



Structural refinement



Structural refinement

Choose a model

How to find a model structure

How to visualize the structure

How to calculate distances and geometries

https://icsd.fizkarlsruhe.de

Database for inorganic structures



http://database.iem.ac.ru/mincryst/

ICSD database





ATOMS on the Web

http://millenia.cars.aps.anl.gov/cgi-bin/atoms/atoms.cgi

		Run ATOMS	Clear	Reset					
	Gold								
Titles									
Operational Parameters									
Space	Fm-3m		6	Edge:	-				
Group: Output -	1								
Type:	feff6.inp	• Shift:							
Lattice Constants and Angles									
<u>A:</u>	4.08	<u>B:</u>	4.08	<u>C:</u>	4.08				
<u>Alpha:</u>	90	Beta:	90	Gamma:	90				
Run ATOMS Clear Reset									
Table of Crystallographic Sites									

Table of Crystallographic Sites							
Cent. Element	X	<u>¥</u>	<u>Z</u>	Tag			
• 1 Au 💌	0	0	0	Au			



XAFS data analysis softwares

http://www.xafs.org/

www.ixasportal.net

http://cars9.uchicago.edu/ifeffit/ Click DOWNLOADS Click <u>ifeffit-1.2.11.exe</u>

http://bruceravel.github.io/demeter/



Step for reducing measured data to $\mu(E)$ and then to $\chi(k)$:

- 1. convert measured intensities to $\mu(E)$
- 2. subtract a smooth pre-edge function, to get rid of any instrumental background, and absorption from other edges.
- 3. normalize $\mu(E)$ to go from 0 to 1, so that it represents 1 absorption event
- 4. remove a smooth post-edge background function to approximate $\mu_0(E)$ to isolate the XAFS χ .
- 5. identify the threshold energy E_0 , and convert from E to k space: $k = \frac{\sqrt{2m(E E_0)}}{\frac{1}{2}}$
- 6. weight the XAFS $\chi(k)$ and Fourier transform from k to R space.
- 7. isolate the $\chi(k)$ for an individual "shell" by Fourier filtering.



Converting raw data to $\mu(E)$





Transmission

The absorption is measured directly by measuring what is transmitted through the sample

$$I = I_0 e^{-\mu(E)t}$$
$$\mu(E)t = \alpha = \ln I_0 / I_1$$

Fluorescence

The re-filling the deep core hole is detected. Typically the fluorescent X-ray is measured

$$\alpha \propto I_F/I_0$$

Tecniche di caratterizzazione con luce di sincrotrone



Pre-edge subtraction and normalization



Pre-edge subtraction

We subtract away the background that fits the *pre edge* region. This gets rid of the absorption due to other edges (say, the Fe L_{III} edge).

Normalization

We estimate the *edge step*, $\mu_0(E_0)$ by extrapolating a simple fit to the above $\mu(E)$ to the edge.



Determination of E₀



Derivative and E₀

We can select E_0 roughly as the energy with the maximum derivative. This is somewhat arbitrary, so we will keep in mind that we may need to refine this value later on.





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Post-edge background

- We do not have a measurement of $\mu_0(E)$ (the absorption coefficient without neighboring atoms).
- We approximate μ₀(E) by an adjustable, smooth function:
 a *spline*.
- A flexible enough spline should not match the $\mu(E)$ and remove all the EXAFS. We want a spline that will match the *low frequency* components of $\mu_0(E)$.



X(k), k-weighting



χ(k)

The raw EXAFS $\chi(k)$ usually decays quickly with k, and difficult to assess or interpret by itself. It is customary to weight the higher k portion of the spectra by multiplying by k² or k³.

k-weighted $\chi(k)$: $k^2\chi(k)$

χ(k) is composed of sine waves,
so we'll Fourier Transform from k
to R-space.
To avoid "ringing", we'll multiply
by a window function.



Fourier Transform: χ(R)



χ(R)

The Fourier Transform of $k^2(k)$ has 2 main peaks, for the first 2 coordination shells: Fe-O and Fe- Fe. The Fe-O distance in FeO is 2.14Å, but the first peak is at 1.66Å. This shift in the first peak is due to the *phase-shift*, $\delta(k)$: sin[2kR + $\delta(k)$].

A shift of -0.5 Å is typical.

$\chi(R)$ is complex:

The FT makes c(R) complex. Usually only the amplitude is shown, but there are really oscillations in c(R). Both real and imaginary components are used in modeling.



Fourier filtering



 χ (R) often has well separated peaks for different "shells".

This shell can be isolated by a Filtered Back-Fourier Transform, using the window shown for the first shell of FeO.

This results in the filtered χ(k) for the selected shell. Many analysis programs use such filtering to remove shells at higher R.

Beyond the first shell, isolating a shell in this way can be difficult.

Tecniche di caratterizzazione con luce di sincrotrone



The information content of EXAFS

• The number of parameters we can reliably measure from our data is limited: $N \approx \frac{2\Delta k \Delta R}{N}$

where Δk and ΔR are the k- and R-ranges of the usable data.

- For the typical ranges like k = [3.0, 12.0] Å⁻¹ and R = [1.0, 3.0] Å, there are ~ 11 parameters that can be determined from EXAFS.
- The "Goodness of Fit" statistics, and confidence in the measured parameters need to reflect this limited amount of data.
- It is often important to constrain parameters R, N, σ^2 for different paths or even different data sets (different edge elements, temperatures, etc)
- Chemical Plausibility can also be incorporated, either to weed out obviously bad results or to use other knowledge of local coordination, such as the Bond Valence Model (relating valence, distance, and coordination number).
- Use as much other information about the system as possible!



Modeling the first shell of FeO - 1

FeO has a rock-salt structure.

To model the FeO EXAFS, we calculate the scattering amplitude f(k) and phase-shift $\delta(k)$, based on a guess of the structure, with Fe-O distance R = 2.14 Å (a regular octahedral coordination). We will use these functions to *refine* the values R, N, σ^2 , and E_0 so our model EXAFS function matches our data.





Fit results N = 5.8 \pm 1.8 R = 2.10 \pm 0.02 Å E₀ = -3.1 \pm 2.5 eV σ^2 = 0.015 \pm 0.005 Å ².

Modeling the first shell of FeO - 2



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1st shell fit in k space

The 1st shell fit to FeO in k space. There is clearly another component in the XAFS

1st shell fit in R space

 $|\chi(R)|$ and Re[$\chi(R)$] for FeO (blue), and a 1st shell fit (red).



Modeling the second shell of FeO - 1

To add the second shell Fe to the model, we use calculation for f(k) and $\delta(k)$ based on a guess of the Fe-Fe distance, and refine the values R,N, σ^2 . Such a fit gives a result like this:



 $|\chi(R)|$ data for FeO (blue), and fit of 1st and 2nd shells (red).

The results are fairly consistent with the known values for crystalline FeO:

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6 O at 2.13Å, 12 Fe at 3.02Å .
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Fit results (uncertainties in parentheses):

Shell	Ν	${f R}$ (A)	σ^2 (Ų)	ΔE_0 (eV)
Fe-O	6.0(1.0)	2.10(.02)	0.015(.003)	-2.1(0.8)
Fe-Fe	11.7(1.3)	3.05(.02)	0.014(.002)	-2.1(0.8)



Modeling the second shell of FeO - 2

Other views of the data and two-shell fit:





The Fe-Fe EXAFS extends to higher-k than the Fe-O EXAFS. Even in this simple system, there is some *overlap* of shells in R-space. The agreement in $\text{Re}[\chi(R)]$ look especially good – this is how the fits are done. The modeling can get more complicated than this



Tecniche di caratterizzazione con luce di sincrotrone