Large Scale Structure of the Universe

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Dark Energy Survey & LSST

Legacy Survey of Space and Time

The Abdus Salam International Centre for Theoretical Physics

Giambiagi Winter School on Cosmology
Buenos Aires, July 2023

Five lectures on Cosmology and Large Scale Structure

Lecture I: The average Universe

Lecture II: Distances and thermal history

Lecture III: The perturbed Universe

Lecture IV: Theoretical challenges and surveys

Lecture V: Observational cosmology with LSS

Plan for Lecture II:

II.1 – Distances, horizons, etc

II.2 – Thermal history of the Universe

II.1- Distances in the universe

II.1.1 – Physical and comoving distances

II.1.2 – Luminosity distance, Hubble-Lemaître law, distance ladder and accelerated expansion

II.1.3 – Angular diameter distance

II.1.4 – Hubble radius

II.1.1 – Physical and comoving distances

Universe is spatially homogeneous and isotropic on average.
On average it is described by the FLRW metric (for a spatially flat universe):

$$ds^2 = dt^2 - a(t)^2 \left[dx^2 + dy^2 + dz^2\right]$$

$$r(t) = a(t)\, x$$

physical distance   comoving distance

For galaxies following the “Hubble flow”(expanding as the average universe):

$$\dot{r}(t) = V = \dot{a}(t)x = \frac{\dot{a}(t)}{a(t)}\, a(t)x = H\, r(t)$$

[In reality galaxies will have some “peculiar velocities” due to inhomogeneities – more later]

The question is: how to measure distances and velocities in practice?

One measures redshifts instead of velocities. That’s done with either measuring spectral lines or with photometric techniques – more later.

How to measure distances in practice?

Comoving distance between us (z=0) and an object at redshift z (c=1):

$$ds^2 = 0 \Rightarrow dt^2 = a(t)^2 d\chi^2 \Rightarrow$$

$$\boxed{\chi(z) = \int_0^z \frac{dz'}{H(z')}}$$

Comoving distances can't be directly measured, but are a useful quantity in cosmology.
They depend on cosmology (through H(z)).

Another useful quantity is the horizon at a given redshift z:
co-moving size of the causal region since the big bang ($z=\infty$) until a given redshift z.

$$\boxed{\chi_h = \int_z^{\infty} dz' \frac{c}{H(z')}}$$

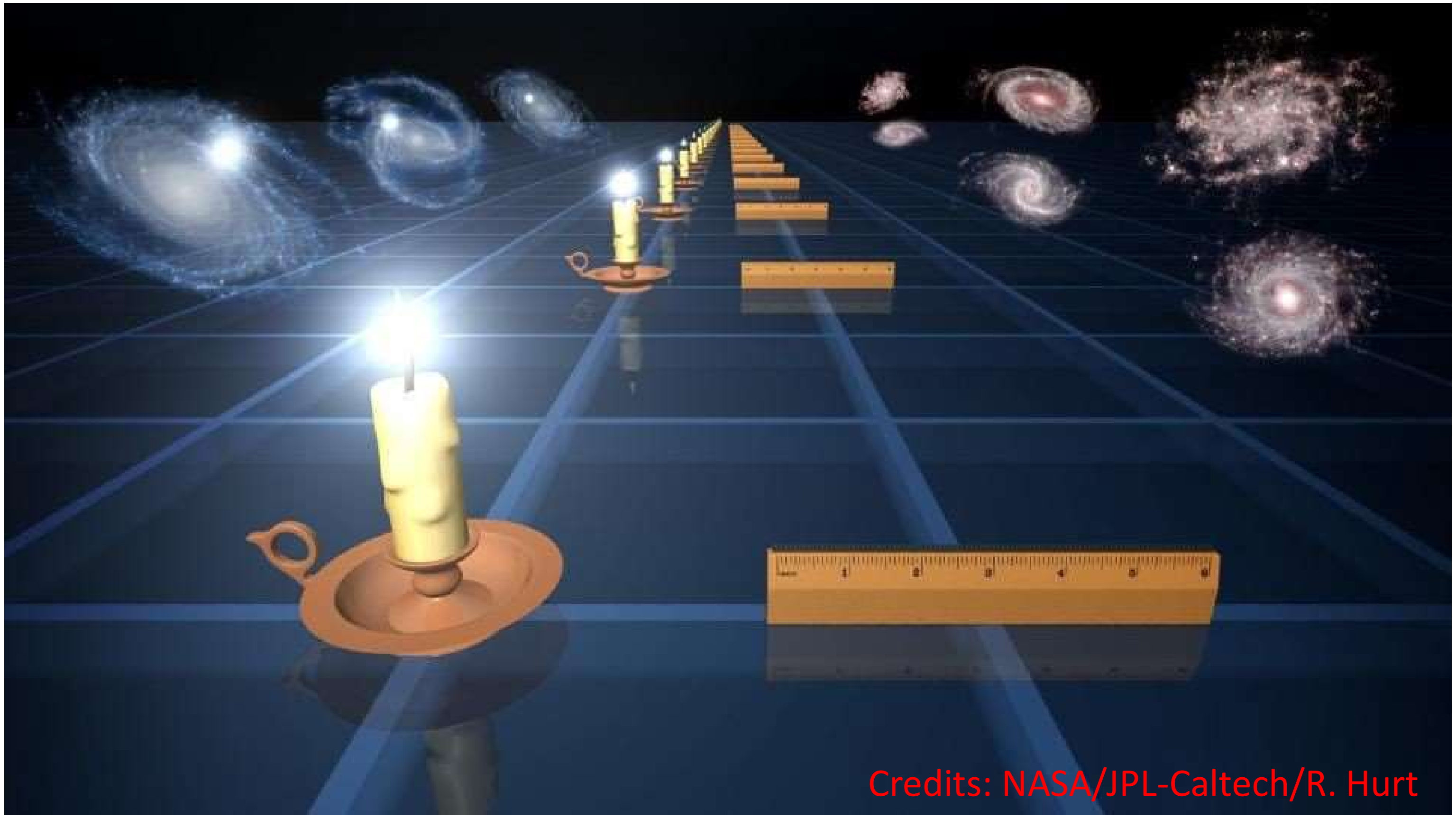
Sound horizon (will be important later): $c \rightarrow c_s$

Sound speed

In the universe it is easy to measure the redshift of objects, using eg spectroscopy.

Measuring distances is much more difficult.

There are 2 ways to measure large distances in the Universe:
from known luminosities (standard candles – eg SNIa) – luminosity distance
from known scales (standard rulers – eg BAO) - angular diameter distance



Credits: NASA/JPL-Caltech/R. Hurt

II.1.2 – Luminosity distance

Luminosity distance (d_L):

$$F = \frac{L}{4\pi d_L^2}$$

Flux of photons received at the detector (a=1)

Known luminosity of the source (energy/time)

In FLRW there are 2 extra source of dilution of the flux:

- redshift of photons (1/(1+z))
- rate of arrival decrease by (1/(1+z)) – time dilation

Therefore: $$\boxed{d_L = (1+z)\chi(z)}$$

We can Taylor expand H(z) to first order:

$$H(z) = H_0 + \frac{dH}{dz} z$$

to obtain the Hubble-Lemaître law:

$$d_L = \frac{z}{H_0} + \mathcal{O}(z^2)$$

Belgian astronomer Georges Lemaître proposed the idea of an expanding universe 2 years before Edwin Hubble's work on receding galaxies. BETTMANN/GETTY IMAGES

Move over, Hubble: Discovery of expanding cosmos assigned to little-known Belgian astronomer-priest

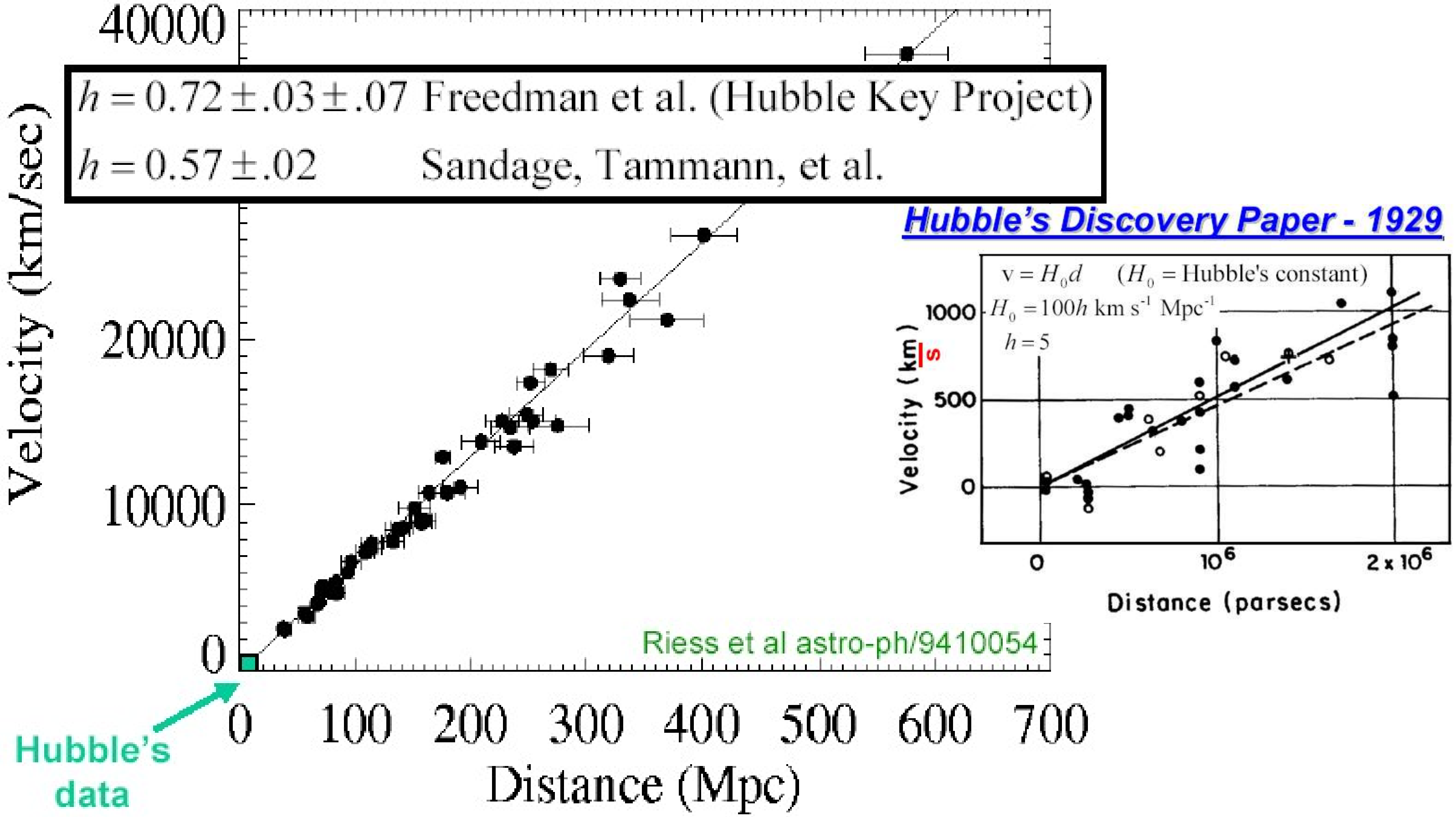
By Daniel Clery | Oct. 29, 2018 , 3:45 PM

Hubble-Lemaître law

Hubble's Law, a cornerstone of cosmology that describes the expanding universe, should now be called the Hubble-Lemaître Law, following a vote by the members of the International Astronomical Union (IAU), the same organization that revoked Pluto's status as a planet. The change is designed to redress the historical neglect of Georges Lemaître, a Belgian astronomer and priest who in 1927 discovered the expanding universe—which also suggests a big bang. Lemaître published his ideas 2 years before U.S. astronomer Edwin Hubble concluded that galaxies farther from the Milky Way recede faster.

The final tally of the 4060 cast votes, announced today by IAU, was 78% in favor of the name change, 20% against, and 2% abstaining. But the vote was not without controversy, both in its execution and the historical facts it was based on. Helge Kragh, a historian of science at the Niels Bohr Institute in Copenhagen, calls the background notes presented to IAU members "bad history." Others argue it is not IAU's job to rename physical laws. "It's bad practice to retroactively change history," says Matthias Steinmetz of the Leibniz Institute for Astrophysics

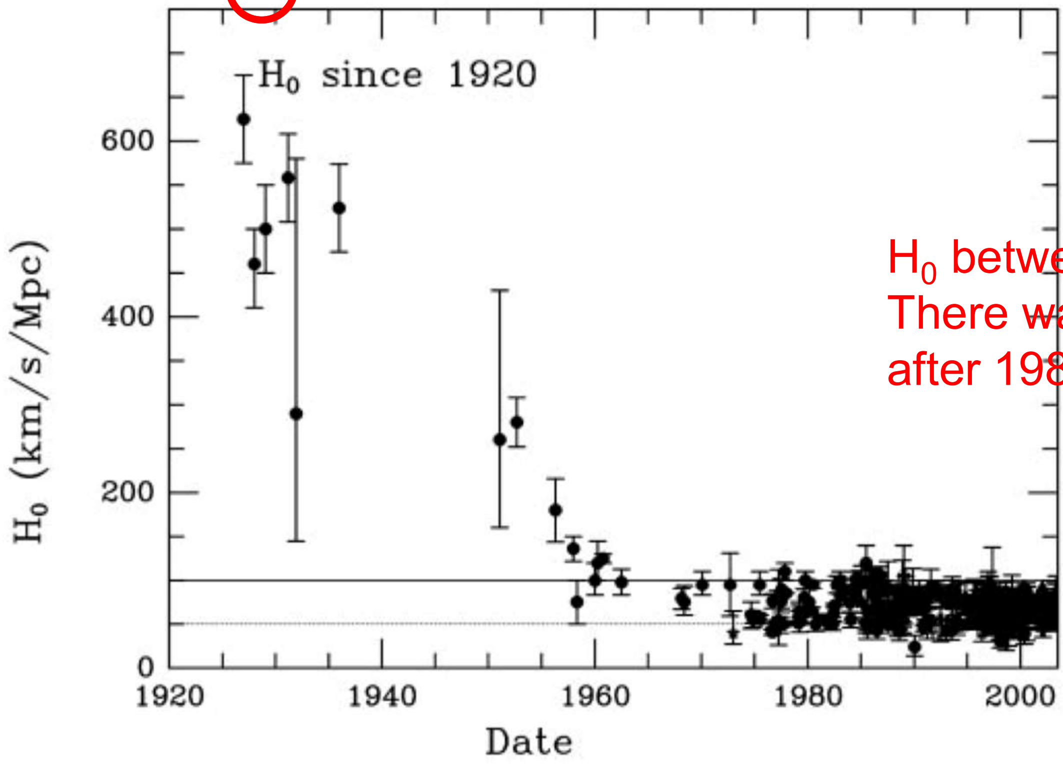
Hubble-Lemaître law



Hubble's diagram and cosmic expansion

Robert P. Kirshner*
Harvard–Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138

Contributed by Robert P. Kirshner, October 21, 2003



H_0 between 50 and 100 km/s/Mpc.
There was a revolution of precision after 1980 and before Planck.

Fig. 2. Published values of the Hubble constant vs. time. Revisions in Hubble's original distance scale account for significant changes in the Hubble constant from 1920 to the present as compiled by John Huchra of the Harvard–Smithsonian Center for Astrophysics. At each epoch, the estimated error in the Hubble constant is small compared with the subsequent changes in its value. This result is a symptom of underestimated systematic errors.

Distance ladder from standard candles: (Cepheids to SNIa) and H_0

Riess et al 1604.01424

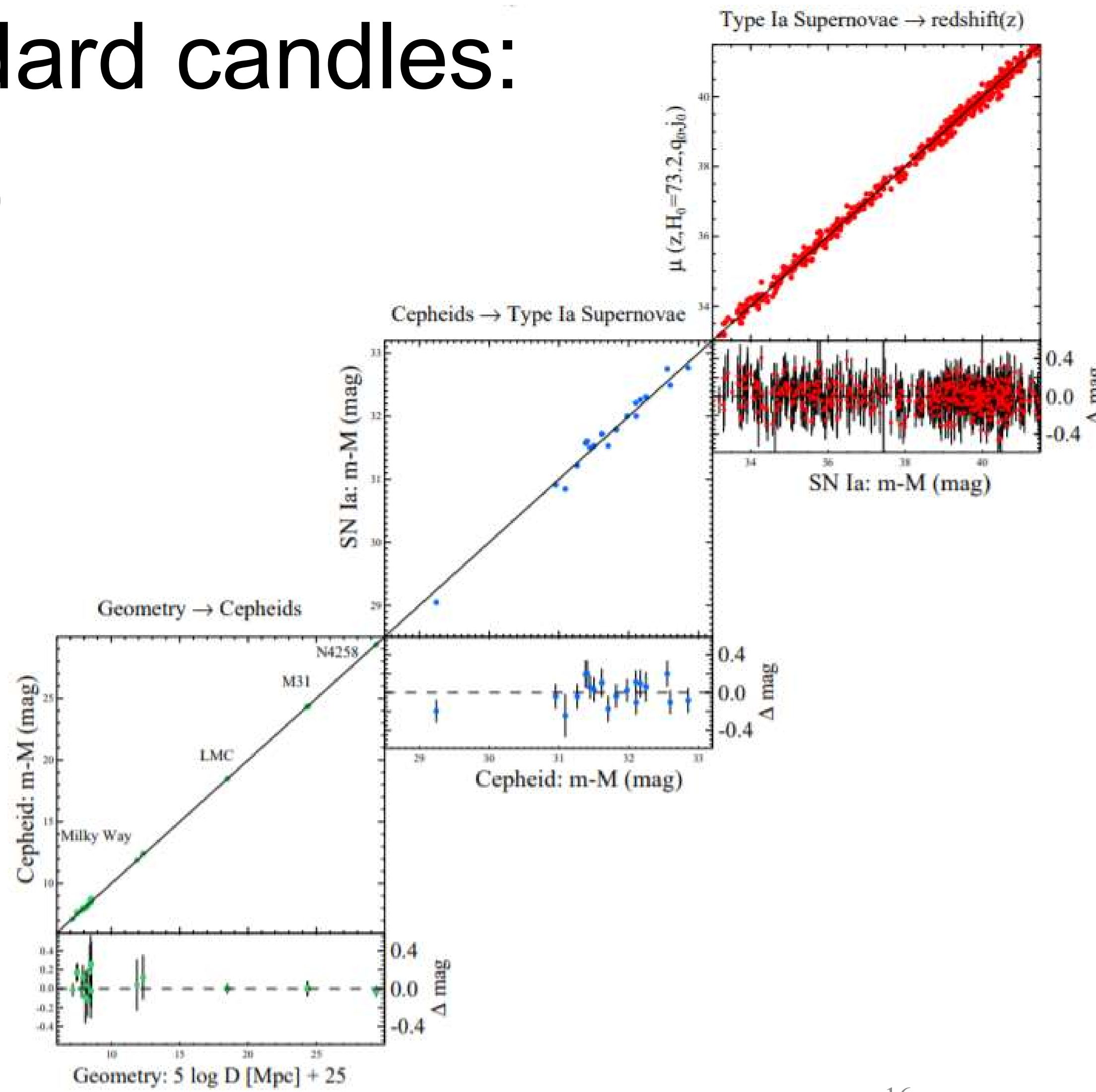
H_0 = 73.24+/-1.74 km/sec/Mpc

Riess et al 1903.07603

H_0 = 74.03 ± 1.42 km/sec/Mpc

Riess et al 2112.04510

H_0 = 73.04 ± 1.04 km/sec/Mpc



Luminosity distance at higher redshifts:

$$\boxed{d_L = (1+z)\chi(z)}$$

Expanding to second order in redshift:

$$H_0 d_L(z) = z + \frac{1}{2}(1-q_0)z^2 + \cdots$$

$$q_0 = -\frac{\ddot{a}}{a\,H_0^2}$$

01 FEBRUARY 1970 • page 34

Cosmology: A search for two numbers

H_0 q_0

Precision measurements of the rate of expansion and the deceleration of the universe may soon provide a major test of cosmological models

Allan R. Sandage

Mount Wilson and Palomar Observatories

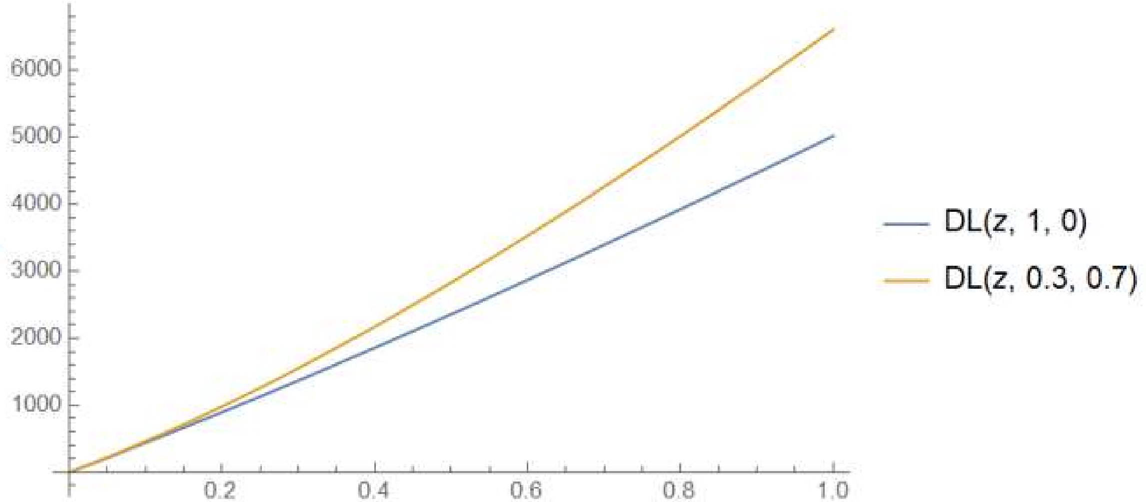
$$\boxed{d_L = (1+z)\chi(z)}$$

We plot $d_L(z)$ (exact expression) for $0<z<2$ for a flat Universe with

a. $\Omega_m = 1$ and $\Omega_\Lambda = 0$
b. $\Omega_m = 0.3$ and $\Omega_\Lambda = 0.7$

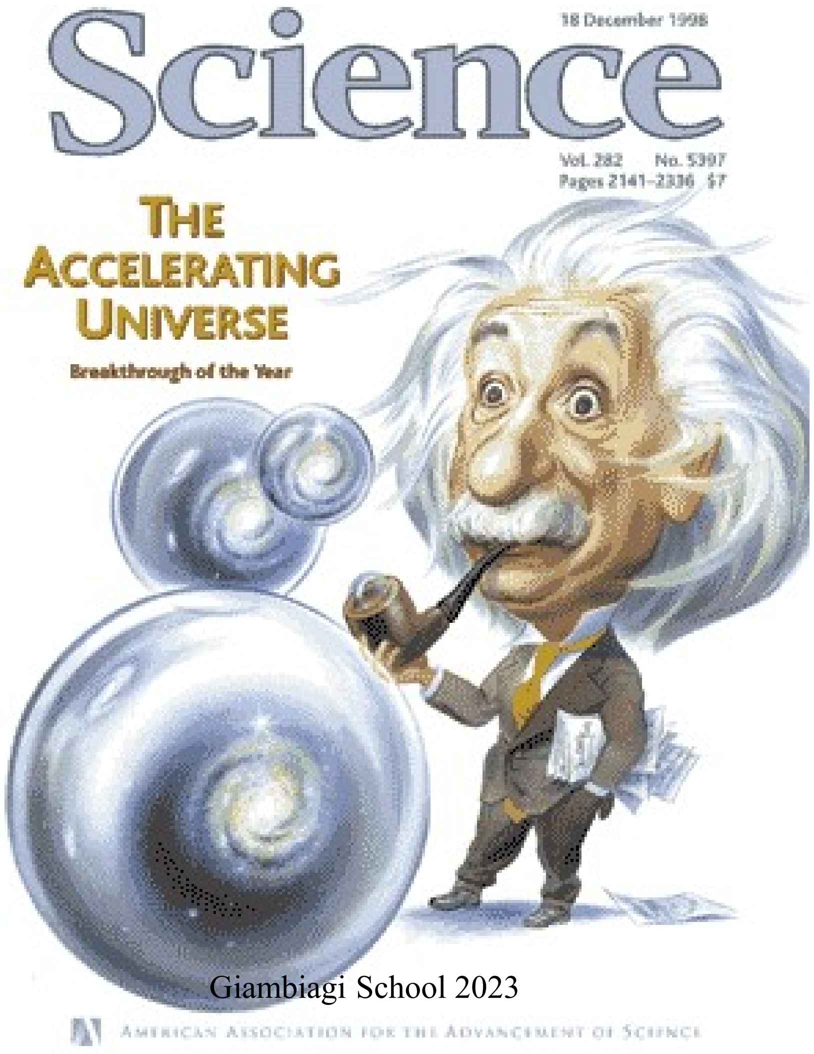
```
In[127]:= H0 = 70. (* km/s/Mpc *);
          c = 300000. (*km/s*);

In[130]:= H[z_, Om_, OΛ_] := H0 √(Om (1 + z)^3 + OΛ)

In[131]:= DL[z_, Om_, OΛ_] := c (1 + z) NIntegrate[1 / H[zp, Om, OΛ], {zp, 0, z}]
                                          integra numéricamente

In[137]:= Plot[{DL[z, 1, 0], DL[z, 0.3, 0.7]}, {z, 0, 1}, PlotLegends → "Expressions"]
          gráfico                                      legenda do gráfico
```

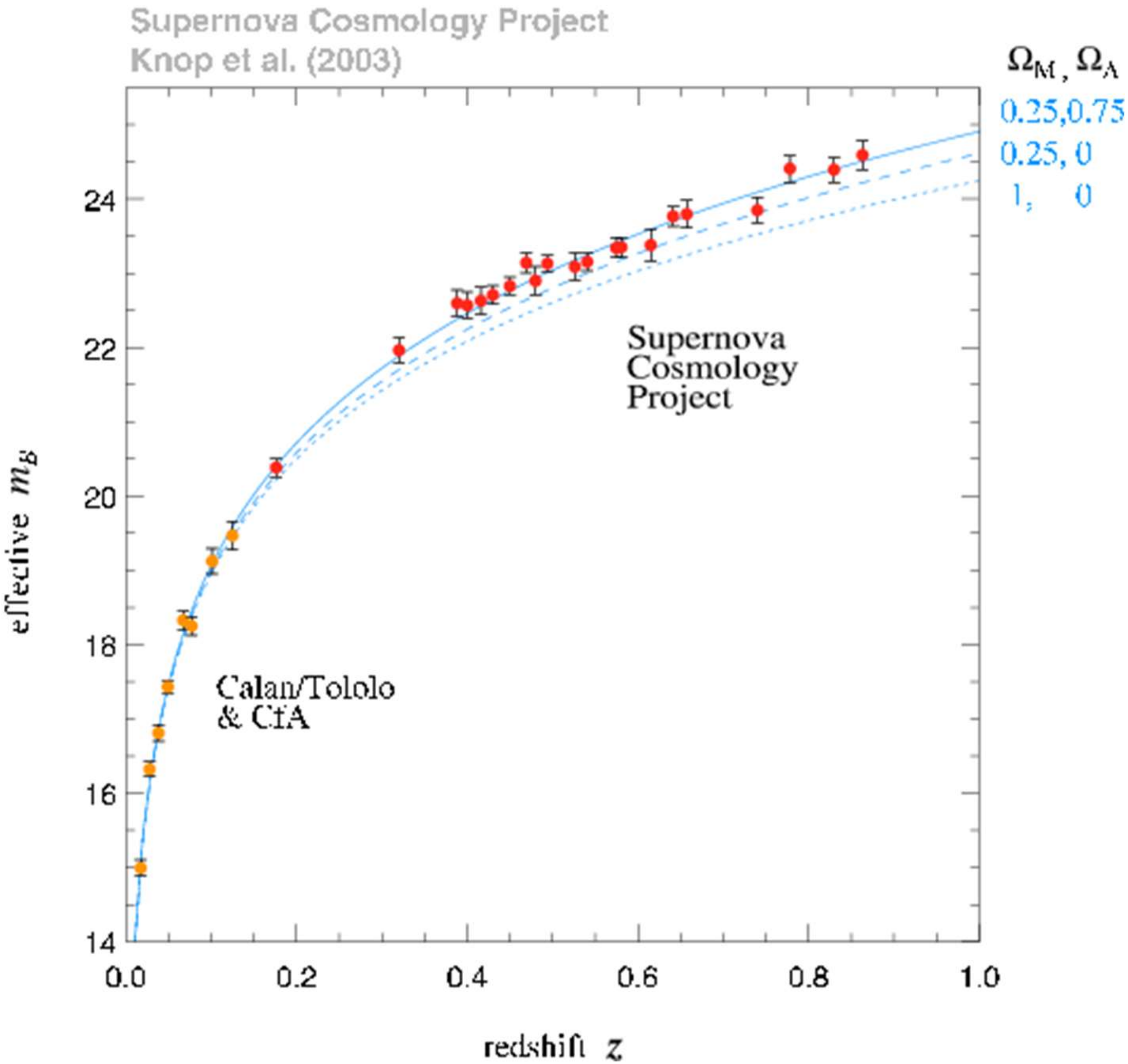
Out[137]=



D_L is larger for a Universe with Λ -> objects **with same z look fainter**. This is how the accelerated expansion of the Universe was discovered in 1998 using SNIa

The big surprise in 1998:





Supernova Cosmology Project
Knop et al. (2003)
Ω_M, Ω_Λ
0.25,0.75
0.25, 0
1, 0
Supernova
Cosmology
Project
Calan/Tololo
& CfA
effective m_B
redshift z
24
22
20
18
16
14
0.0
0.2
0.4
0.6
0.8
1.0

The Nobel Prize in Physics 2011
Saul Perlmutter, Brian P. Schmidt, Adam G. Riess

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The Nobel Prize in Physics 2011

Photo: U. Montan

Saul Perlmutter

Prize share: 1/2

Photo: U. Montan

Brian P. Schmidt

Prize share: 1/4

Photo: U. Montan

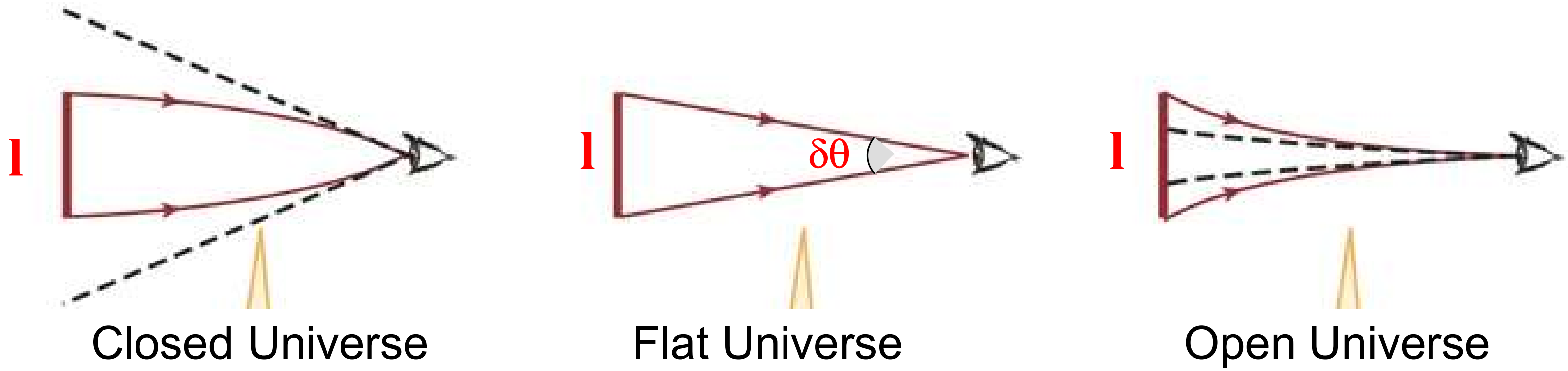
Adam G. Riess

Prize share: 1/4

The Nobel Prize in Physics 2011 was divided, one half awarded to Saul Perlmutter, the other half jointly to Brian P. Schmidt and Adam G. Riess *"for the discovery of the accelerating expansion of the Universe through observations of distant supernovae".*

II.1.3 – Angular diameter distance

Angular diameter distance (d_A) is related to the angle subentended by a physical scale (l)



For a flat universe: $d_A = \frac{l}{\delta\theta},\ l = a\chi\delta\theta \Rightarrow$

$$\boxed{d_A = \frac{1}{1+z}\chi(z)}$$

Is there a favored physical scale in the universe?

Yes: the “acoustic horizon scale at decoupling” (r_s) – more later

This physical scale sets the angular scale for the fluctuations in the cosmic microwave background (CMB) and in the distribution of galaxies that are formed much later (baryon acoustic oscillation – BAO)!

arXiv.org > astro-ph > arXiv:2107.04646

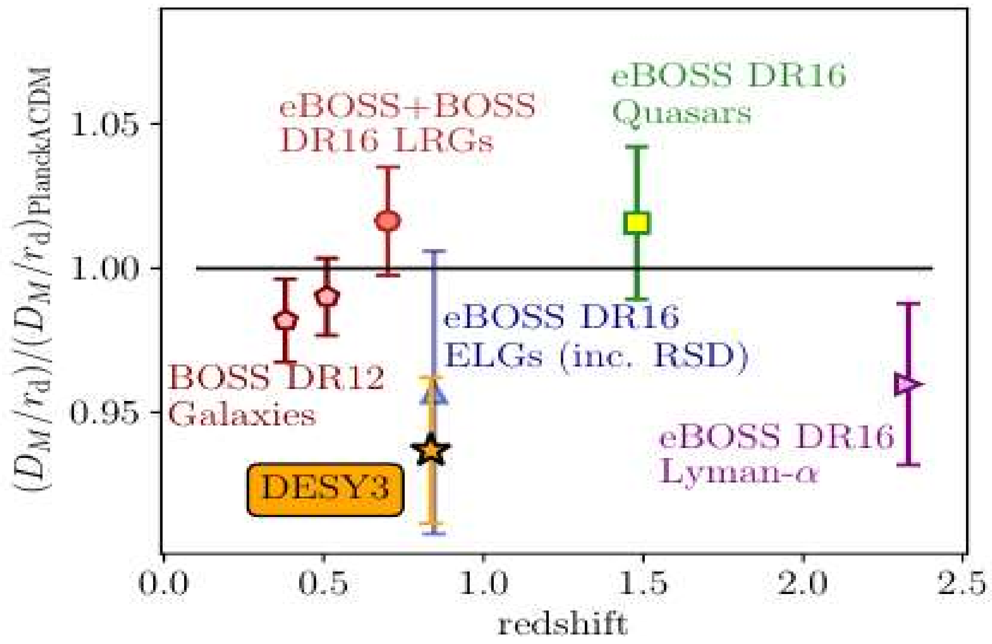
Search...

Help | Ad

Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 9 Jul 2021]

Dark Energy Survey Year 3 Results: A 2.7% measurement of Baryon Acoustic Oscillation distance scale at redshift 0.835



DES-Y3 BAO sample:
7 million galaxies with
$0.6 < z_{phot} < 1.1$ in an area of ~4100 deg^2

2.3σ below Planck

II.1.4 – Hubble radius ("Hubble horizon")

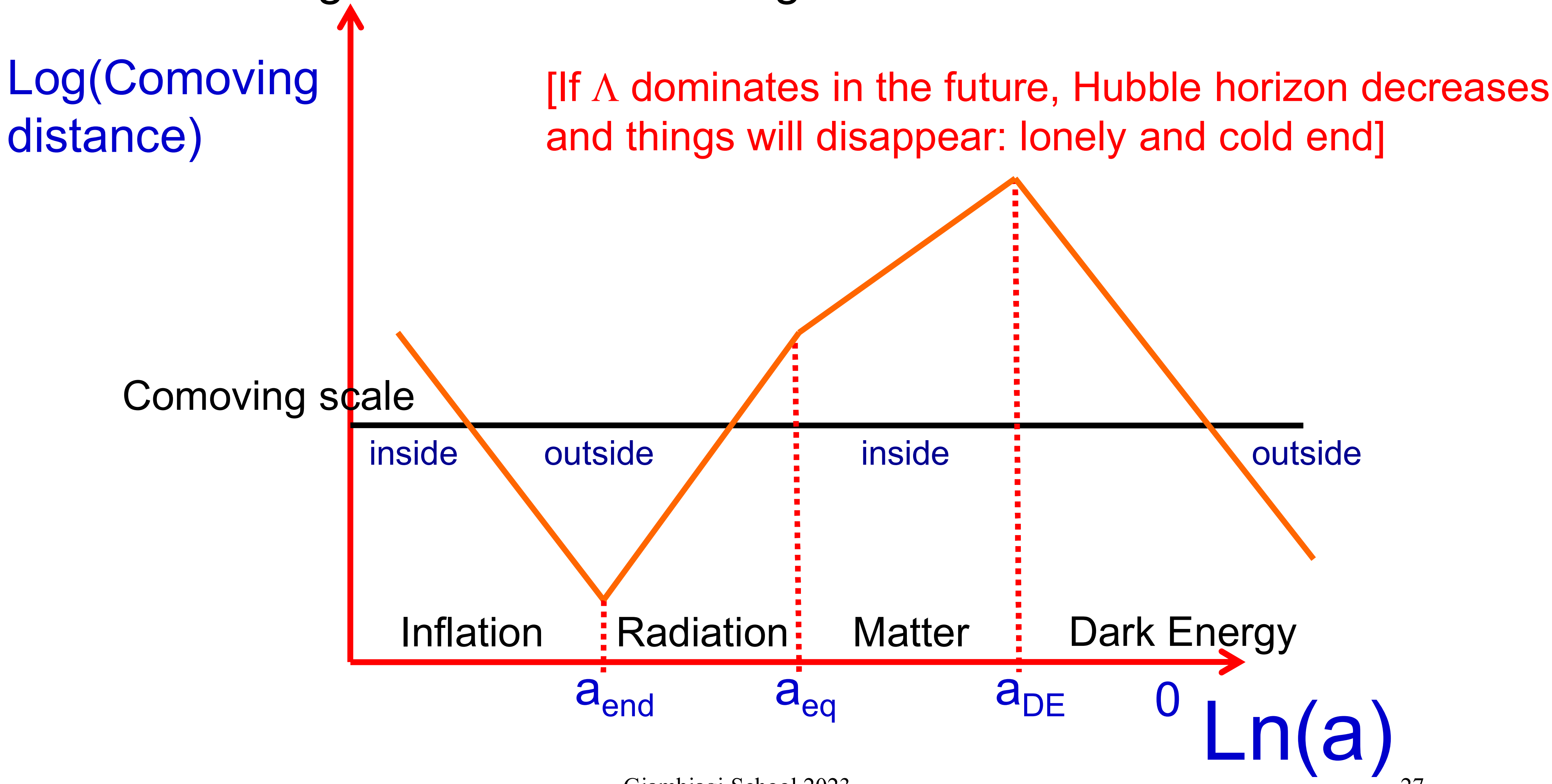
Another useful scale: characteristic distance light can travel in a Hubble time (c=1):

$$R_H = \frac{1}{H(t)}$$

Comoving Hubble radius: $$r_H = \frac{1}{aH} = \frac{1}{\dot{a}}$$

- Radiation dominated $\quad r_H \propto a$
- Matter dominated $\quad r_H \propto a^{1/2}$
- Λ dominated $\quad r_H \propto 1/a$

Comoving Hubble radius during the evolution of the Universe



II.2- Thermal history of the Universe

II.2.1 – Brief review of thermodynamics

II.2.2 – Effective number of relativistic degrees of freedom

II.2.3 – Relation between scale fator and temperature

II.2.4 – Temperature-time relationship

II.2.5 – The origin of the cosmic microwave background

II.2.6 – The standard ruler in the sky

II.2.1 – Brief review of thermodynamics

Number density and energy density of a dilute, weakly interacting gas with g internal degrees of freedom:

$$n = \frac{g}{(2\pi)^3} \int d^3p\, f(\vec{p})$$

$$\rho = \frac{g}{(2\pi)^3} \int d^3p\, E(\vec{p}) f(\vec{p})$$

E is the energy of a state, f(p) is its phase-space distribution and g is number of internal degrees of freedom
(eg g=2 for photons, g=16 for gluons, g=12 for quarks, etc).

$$E = \sqrt{|\vec{p}|^2 + m^2}$$

Phase-space distribution for one species in kinetic equilibrium (+ for FD, - for BE), k_B=1, μ chemical potential:

$$f(\vec{p}) = \frac{1}{e^{(E-\mu)/T} \pm 1}$$

Relativistic limit (T>>m) and T>>μ

$$\rho = \left(\frac{\pi^2}{30}\right) g T^4 * \begin{cases} 1 \text{ (Bose - Einstein)} \\ \frac{7}{8} \text{ (Fermi - Dirac)} \end{cases}$$

$$n = \frac{\zeta(3)}{\pi^2} g T^3 * \begin{cases} 1 \text{ (Bose - Einstein)} \\ \frac{3}{4} \text{ (Fermi - Dirac)} \end{cases} \qquad \zeta(3) = 1.202\cdots$$

Exercise: compute the number of CMB photons (T=2.73 K) in 1 cm^3

Non-relativistic limit (T<<m) and μ=0 [same for B-E and F-D]

$$n = g\left(\frac{mT}{2\pi}\right)^{3/2} e^{-m/T}$$

$$\rho = mn$$

Exponential Boltzmann suppression

For more than one species: density of relativistic particles in the Universe is set by the effective number of relativistic degrees of freedom g_*:

$$\rho_r = \frac{\pi^2}{30} g_* T^4$$

$$g_* = \sum_{\text{bosons}} g_i \left(\frac{T_i}{T}\right)^4 + \sum_{\text{fermions}} \frac{7}{8} g_i \left(\frac{T_i}{T}\right)^4$$

We assumed that in principle the particles can have a different temperature than T (the photon's temperature). If they are in chemical equilibrium then same T.

$g_*(T)$ changes when mass thresholds are crossed as T decreases and particles become non-relativistic. At high T (>200 GeV) $g_*^{(SM)} \sim 100$.

http://www.damtp.cam.ac.uk/user/db275/Cosmology/Lectures.pdf

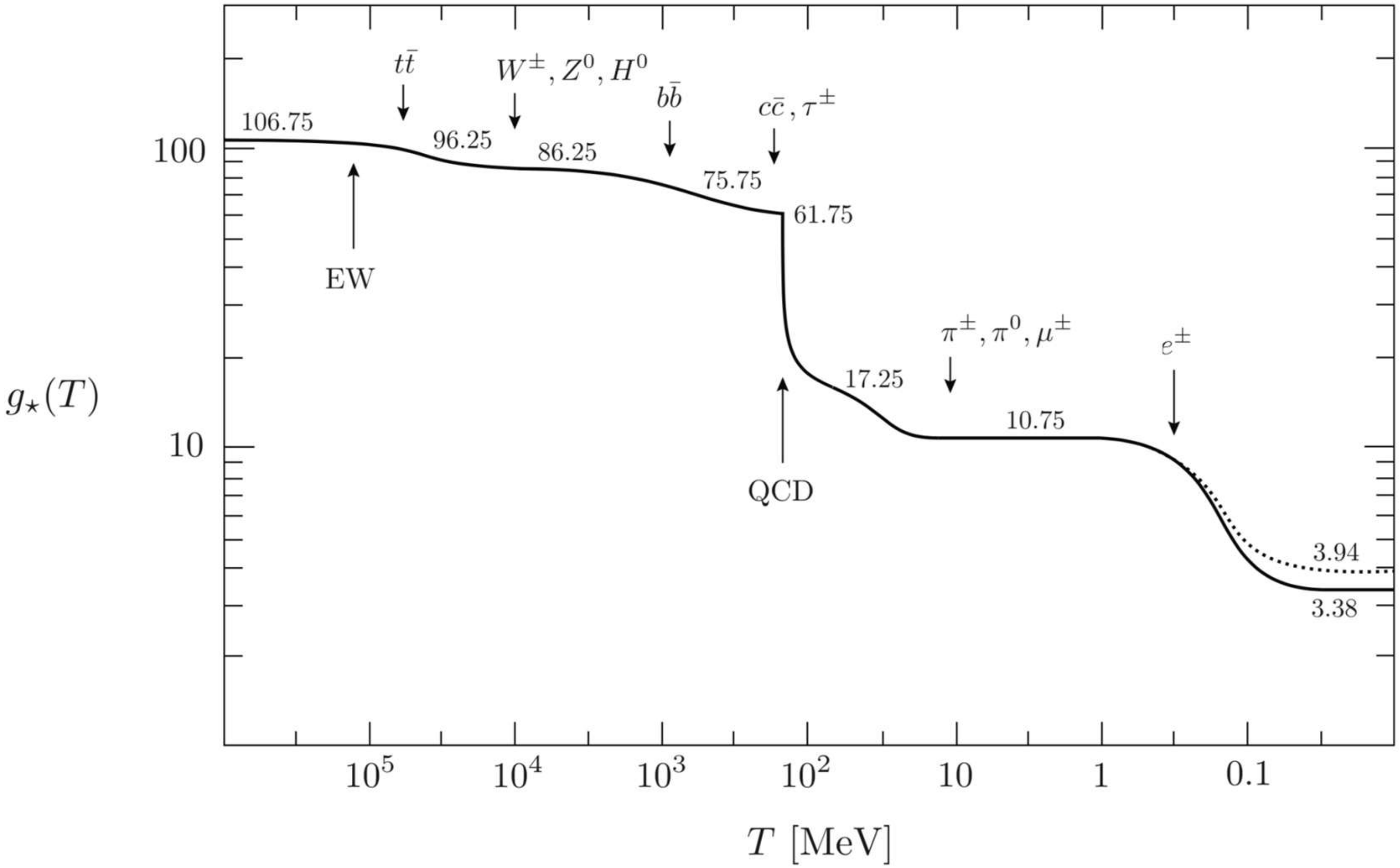


Figure 3.4: *Evolution of relativistic degrees of freedom $g_\star(T)$ assuming the Standard Model particle content. The dotted line stands for the number of effective degrees of freedom in entropy $g_{\star S}(T)$.*

II.2.2 – Effective number of relativistic degrees of freedom: N_{eff}

$$\rho_r = \rho_\gamma \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\rm eff} \right]$$

2306.05460

$$N_{\rm eff} = N_{\rm eff}^{\rm SM} + \Delta N_{\rm eff} \qquad N_{\rm eff}^{\rm SM} = 3.043$$

ΔN_{eff} parametrizes new relativistic degrees of freedom: Dark Radiation

Possible solution to the Hubble tension? See, eg, 2306.15067

II.2.3 – Relation between scale fator and temperature

Naivelly one can find the relation between the scale fator and the temperature as:

$$\rho_r \propto T^4; \; \rho_r \propto a^{-4} \Rightarrow a \propto T^{-1}$$

More precisely, conservation of entropy implies:

$$T \propto g_*^{-1/3} a^{-1}$$

When a mass threshold of a particle species is crossed, those particles annihilate into photons and increase the temperature.

II.2.4 – Temperature-time relationship

From Friedmann's 1st equation for a radiation-dominated era:

$$H = \sqrt{\frac{\rho_r}{3\tilde{M}_{\mathrm{Pl}}^2}} \sim \frac{T^2}{M_{\mathrm{Pl}}}$$

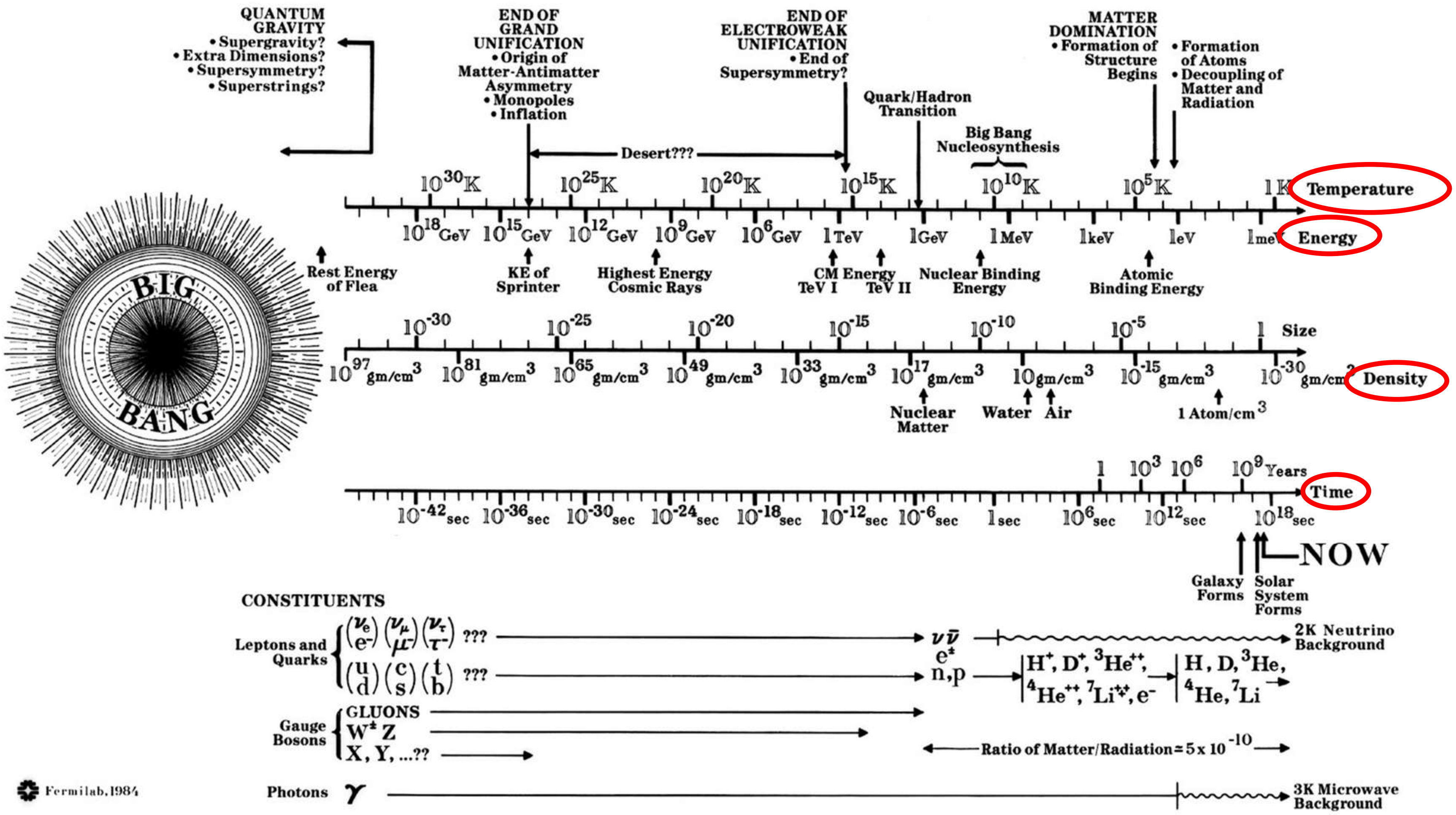
and $\quad H = \dfrac{\dot{a}}{a} \propto t^{-1}$

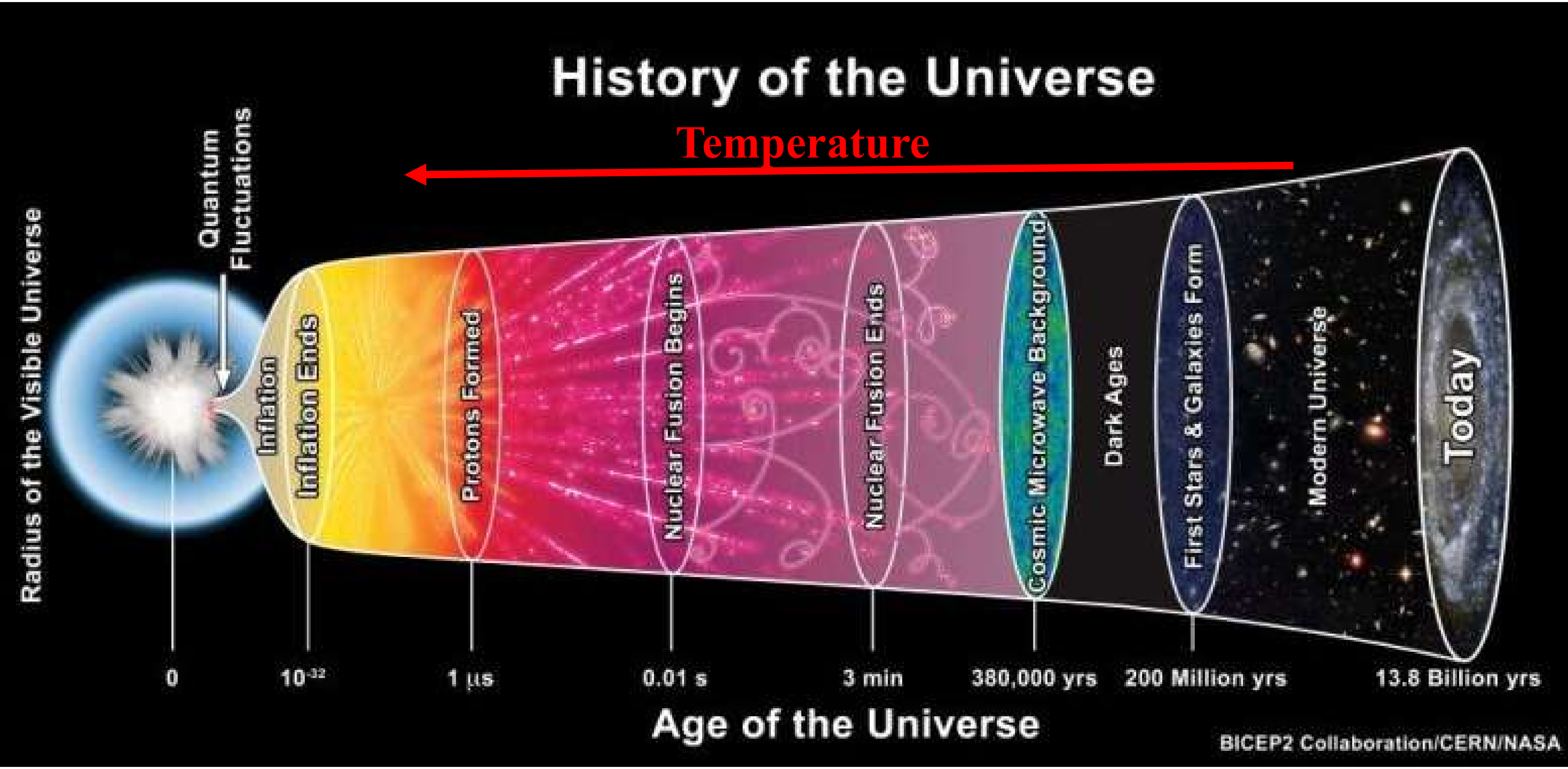
one finds: $\quad T \propto t^{-1/2}$

Putting numbers: $\quad T(\mathrm{MeV}) \simeq 1.5 g_*^{-1/4} t(\mathrm{s})^{-1/2}$

Thermal history of the Universe

Kolb & Turner




History of the Universe
Temperature
Radius of the Visible Universe
Quantum Fluctuations
Inflation
Inflation Ends
Protons Formed
Nuclear Fusion Begins
Nuclear Fusion Ends
Cosmic Microwave Background
Dark Ages
First Stars & Galaxies Form
Modern Universe
Today
0
10^-32
1 μs
0.01 s
3 min
380,000 yrs
200 Million yrs
13.8 Billion yrs
Age of the Universe
BICEP2 Collaboration/CERN/NASA

II.2.5 – The origin of the cosmic microwave background
(Much more in Julien's lectures!!)

When the universe is very hot there are no atoms! It is fully ionized. It is opaque to EM radiation. Light and matter are tightly coupled.

Decoupling of light occurs when the Universe cools down - protons and eletrons can combine to form hydrogen atoms: recombination epoch.

The cosmic microwave background (CMB) is generated after photons decouple: last scattering surface.

Naivelly one may think that this happens at a temperature $T_{binding}$=13.6 eV

However, one must take into account that there are many more photons than protons in the Universe!

One should study the reaction:

$$H + \gamma \longleftrightarrow e^- + p^+$$

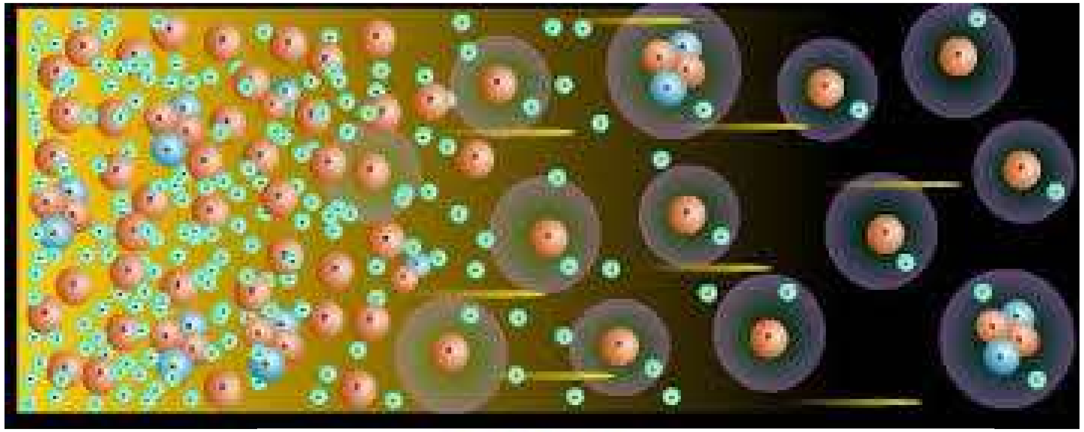
and find at what temperature hydrogen stops being destroyed by photons. Correct treatment is to analyze a Boltzmann equation.

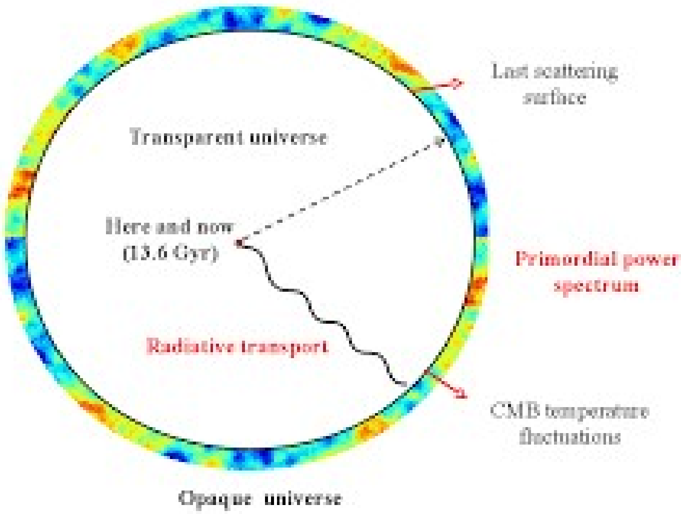
Result: T_{rec}~0.26 eV ~ 3000 K which implies:

z_{rec}~ T_{rec}/T_0~ 3000/2.7~ 1100

Time ~ 380,000 years

After decoupling the Universe is neutral and becomes transparente to photons: “last scattering suface” at z=1100.





II.2.6 – The standard ruler in the sky

The comoving sound horizon at decoupling: r_s

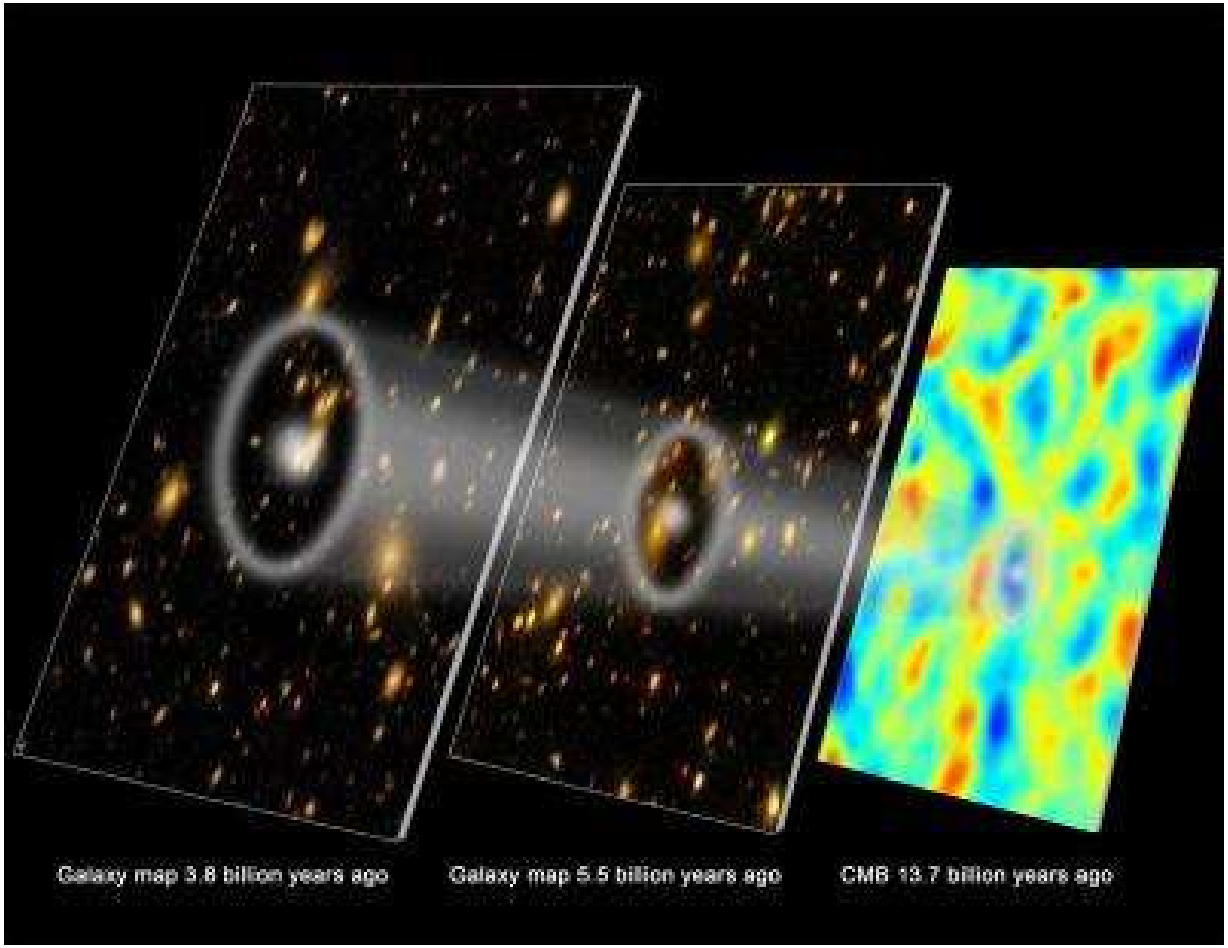
$$r_s = \int_{z_{dec}}^{\infty} \frac{c_s}{H(z)}$$

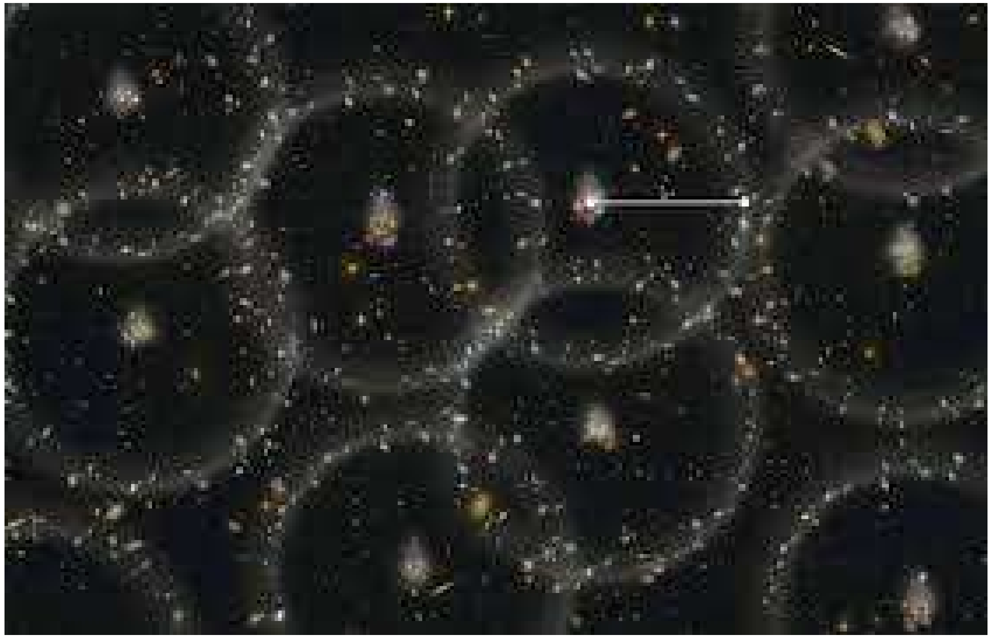
This is the standard ruler measured in CMB and BAO

c_s: speed of perturbations in the coupled baryon-photon fluid – for relativistic fluids

$$c_s^2 = \frac{\delta P}{\delta \rho} \approx \frac{1}{3}$$

Exercise: show that $r_s \sim 150$ Mpc





Artist impression of baryonic acoustic oscillations imprinting a mean galaxy separation of 150 Mpc on the cosmic galaxy distribution.
Credit: Eric Huff, the SDSS-III team, and the South Pole Telescope team. Graphic by Zosia Rostomian

End of second lecture

Extra slides (time permitting)

II.3- The Hubble tension (crisis)

The comoving sound horizon at decoupling ($z \sim 1100$) sets a physical scale in the Universe, both in the fluctuations of the CMB and the BAO.

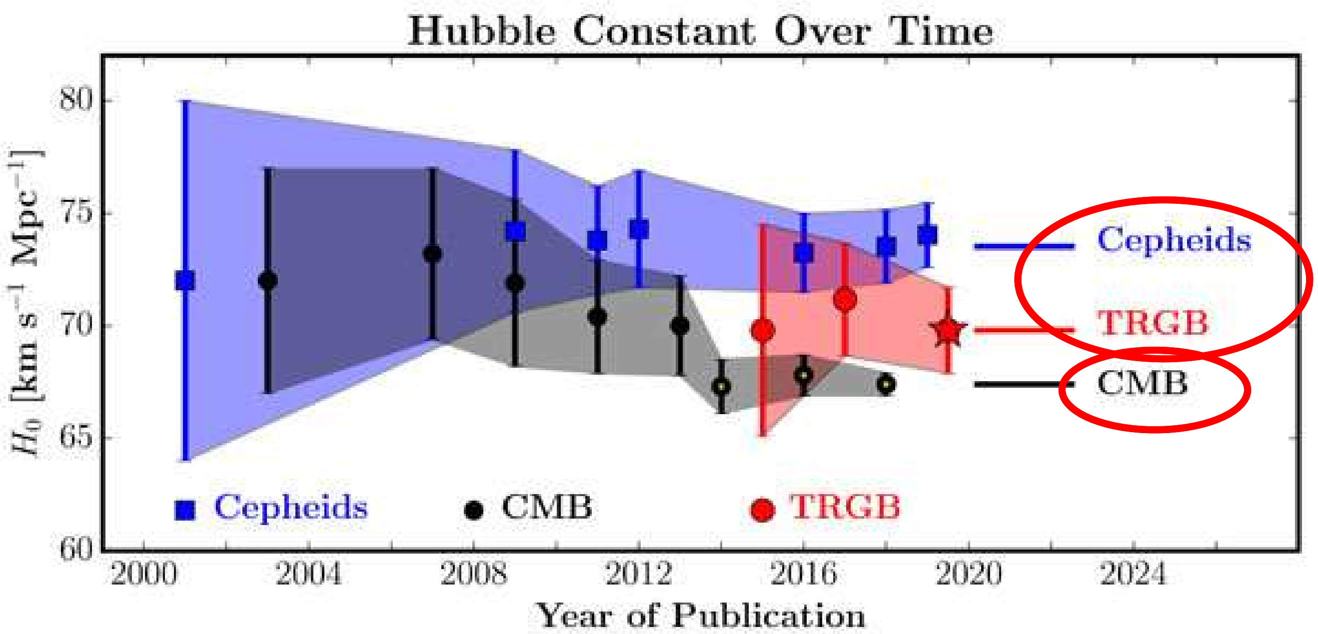
A measurement of the angular distance related to this scale can also allow for an indirect (model dependent) estimate of the Hubble constant.

This indirect measurement performed by the Planck satellite is very precise.

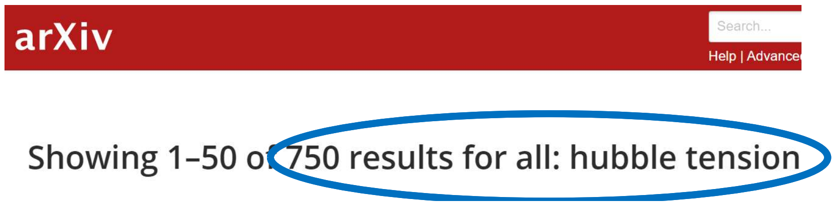
Greater precision brings greater possibilities for disagreement.
In fact, there is a 4-6 σ tension with the local measurements.

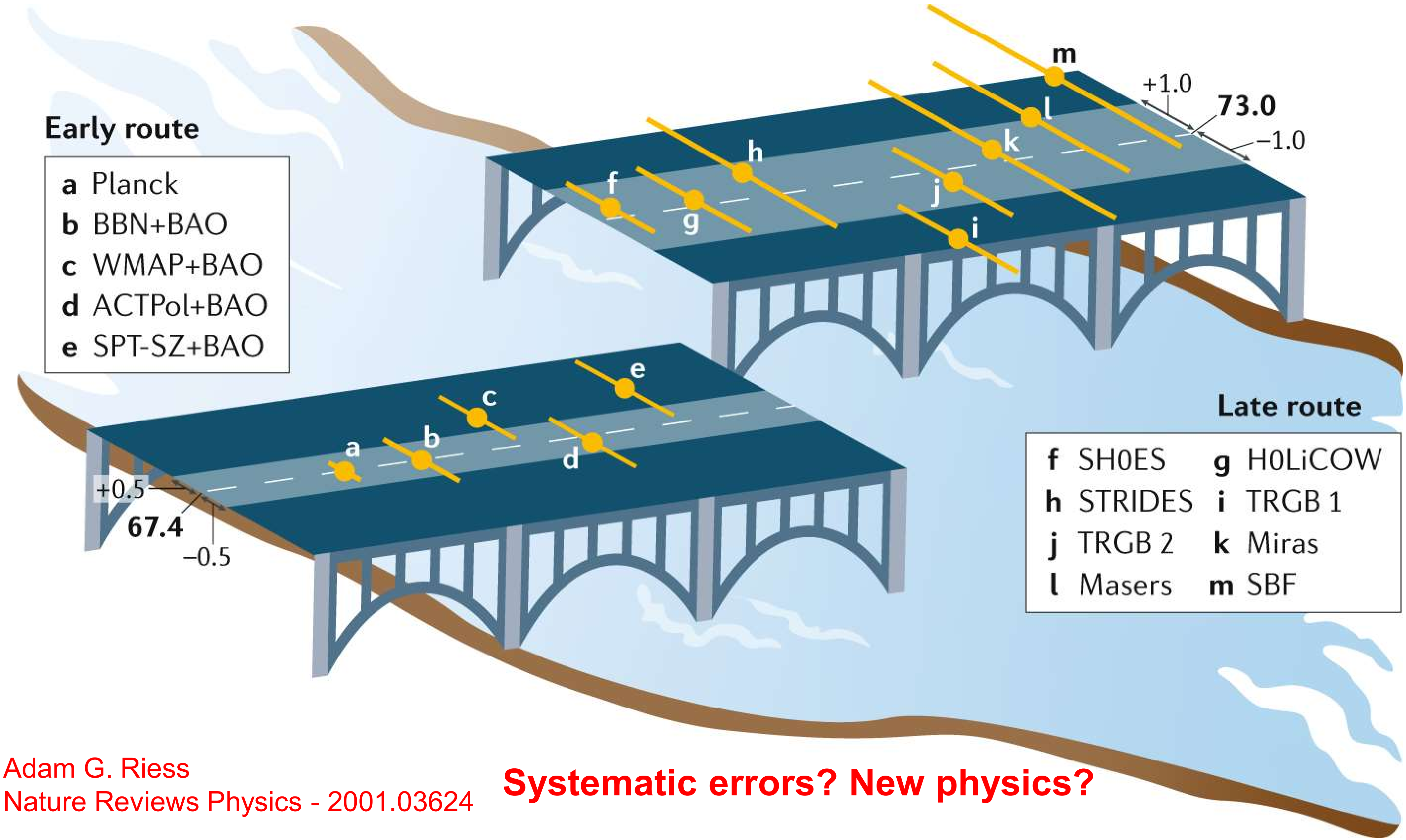
There was a recent revolution in the measurement of H_0 with great precision (~1%)! The **Hubble tension** ~4-6σ!

First crack in the standard ΛCDM model?



Freedman et al.,1907.05922





Adam G. Riess
Nature Reviews Physics - 2001.03624

Systematic errors? New physics?

II.3.1 – The CMB measurement

CMB provides an indirect (model dependent) estimate of the Hubble constant. It's a result of a complicated fit to the CMB angular power spectrum with several parameters.

However, we can have a rough idea by looking at the physical quantity measured by Planck satellite: the angular acoustic horizon scale of the CMB fluctuations, θ_* (~1^0)

The angular acoustic scale is measured with high precision (0.03%) by Planck:

$$100\ \theta_* = 1.0411 \pm 0.0003$$

$$\theta_* = \frac{r_s}{d_A(z_{dec})}$$

$$r_s(z_{dec}) = \int_{z_{dec}}^{\infty} \frac{c_s dz}{H(z)}$$

Early times

(expansion rate around decoupling era)

$$d_A(z_{dec}) = \int_{0}^{z_{dec}} \frac{c\, dz}{H(z)}$$

Late times

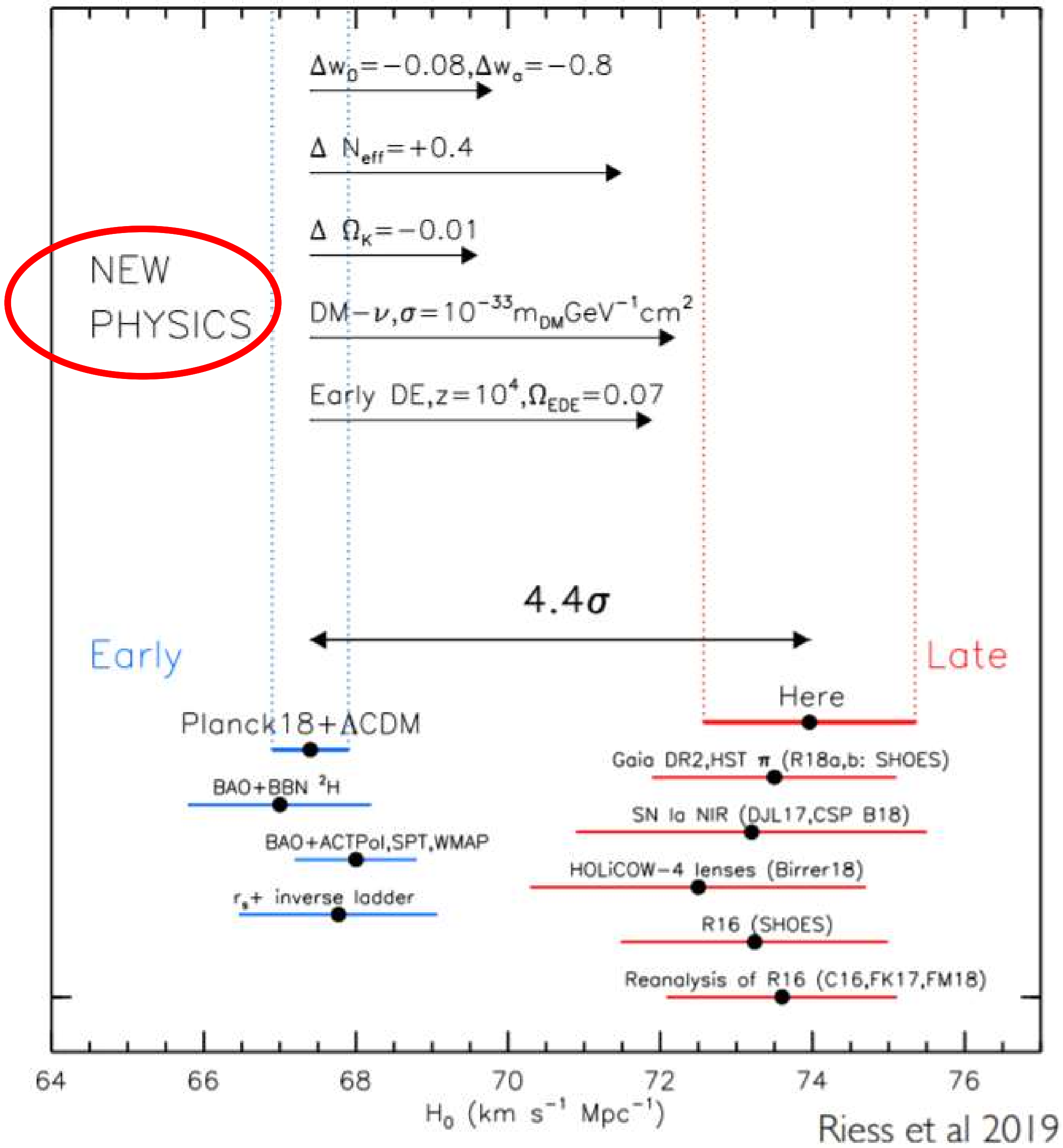
(expansion rate around present era)

Obs: comoving angular diameter distance

Attempts to reconcile CMB with local measurements: New Physics!

If there is an extra contribution to the energy density (=faster expansion rate) with respect to ΛCDM around the recombination era then in order to keep θ_* fixed requires **a larger value of H_0**

New relativistic degrees of freedom, early dark energy, decaying dark matter,...



$\Delta w_0=-0.08, \Delta w_a=-0.8$
$\Delta\ N_{eff}=+0.4$
$\Delta\ \Omega_K=-0.01$
$DM-\nu, \sigma=10^{-33} m_{DM} GeV^{-1} cm^2$
Early DE, $z=10^4, \Omega_{EDE}=0.07$
NEW PHYSICS
4.4σ
Early
Late
Here
Planck18+ΛCDM
BAO+BBN ^{2}H
BAO+ACTPol,SPT,WMAP
r_s+ inverse ladder
Gaia DR2,HST π (R18a,b: SHOES)
SN Ia NIR (DJL17,CSP B18)
HOLiCOW-4 lenses (Birrer18)
R16 (SHOES)
Reanalysis of R16 (C16,FK17,FM18)
64
66
68
70
72
74
76
H_0 (km s^{-1} Mpc^{-1})
Riess et al 2019

Models abound:

In the Realm of the Hubble tension − a Review of Solutions 2103.01183

E. Di Valentino, O. Mena, S. Pan, L. Visinelli, W. Yang, A. Melchiorri, D. F. Mota, A. G. Riess, J. Silk

tension $\lesssim 1\sigma$ *"Excellent models"*	tension $\lesssim 2\sigma$ *"Good models"*	tension $\lesssim 3\sigma$ *"Promising models"*
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]
Vacuum Metamorphosis [402]	Bulk viscous models [340, 341]	Decaying Ultralight Scalar [538]
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	BD-ΛCDM [852]
Sterile Neutrinos [433]	Phantom Braneworld DE [378]	Metastable Dark Energy [314]
Decaying Dark Matter [481]	PEDE [391, 392]	Self-Interacting Neutrinos [700]
Neutrino-Majoron Interactions [509]	Elaborated Vacuum Metamorphosis [401]	Dark Neutrino Interactions [716]
IDE [637, 639, 657, 661]	IDE [659, 670]	IDE [634–636, 653, 656, 663, 669]
DM - Photon Coupling [685]	Interacting Dark Radiation [517]	Scalar-tensor gravity [855, 856]
$f(\mathcal{T})$ gravity theory [812]	Decaying Dark Matter [471, 474]	Galileon gravity [877, 881]
BD-ΛCDM [851]	DM - Photon Coupling [686]	Nonlocal gravity [886]
Über-Gravity [59]	Self-interacting sterile neutrinos [711]	Modified recombination [986]
Galileon Gravity [875]	$f(\mathcal{T})$ gravity theory [817]	Effective Electron Rest Mass [989]
Unimodular Gravity [890]	Über-Gravity [871]	Super ΛCDM [1007]
Time Varying Electron Mass [990]	VCDM [893]	Axi-Higgs [991]
ΛΛCDM [995]	Primordial magnetic fields [992]	Self-Interacting Dark Matter [479]
Ginzburg-Landau theory [996]	Early modified gravity [859]	Primordial Black Holes [545]
Lorentzian Quintessential Inflation [979]	Bianchi type I spacetime [999]	
Holographic Dark Energy [351]	$f(\mathcal{T})$ [818]	

Table B2. Models solving the H_0 tension with R20 within 1σ, 2σ and 3σ considering *Planck* in combination with additional cosmological probes. Details of the combined datasets are discussed in the main text.

Jury is still out on the possible solutions to the Hubble tension or crisis... lots of works

New physics vs Systematic errors