

## Books:

- *The Young Universe: Primordial Cosmology*,  
edited by R. Taillet (John Wiley & Sons, 2022) ISBN : 1789450322  
→ Chapter 2: Cosmological Microwave Background, by JL
- *Neutrino cosmology*,  
JL, G. Mangano, G. Miele, S. Pastor (Cambridge University Press 2013)  
→ Chapter 5: Cosmological Microwave Background, by JL

## Notes from Master course on advanced Cosmology:

- *The Ingredients of the Universe*, -> [link on Indico page of this school](#)  
JL, course at RWTH Aachen University
  1. Recalls on homogeneous cosmology
  2. Thermal history of the Universe
  3. Linearised gravity
  4. Inflation
  5. CMB anisotropies
  6. Large Scale Structure

# Summary of yesterday: perturbed degrees of freedom and equations of motion

Degrees of freedom	Equation of motion
Gravitational potential $\psi(\eta, \vec{x})$	Einstein 00: $\phi, \psi \leftrightarrow \delta$
Scale factor distortion $\phi(\eta, \vec{x})$	Einstein ij: $(\phi - \psi) \longrightarrow \sigma$
CDM density/velocity $\delta_c(\eta, \vec{x}), \theta_c(\eta, \vec{x})$	continuity $\delta'_c + \text{Euler } \theta'_c$
Baryon density/velocity $\delta_b(\eta, \vec{x}), \theta_b(\eta, \vec{x})$	Continuity $\delta'_b + \text{Euler } \theta'_b$ (incl. Thomson)
Photons $f_\gamma(\eta, \vec{x}, p, \hat{n})$	Boltzmann $\frac{d}{d\eta} f_\gamma = \text{Thomson}$
[ Neutrinos $f_\nu(\eta, \vec{x}, p, \hat{n})$ ]	[ Boltzmann $\frac{d}{d\eta} f_\nu = 0$ ]

# Photon phase-space distribution

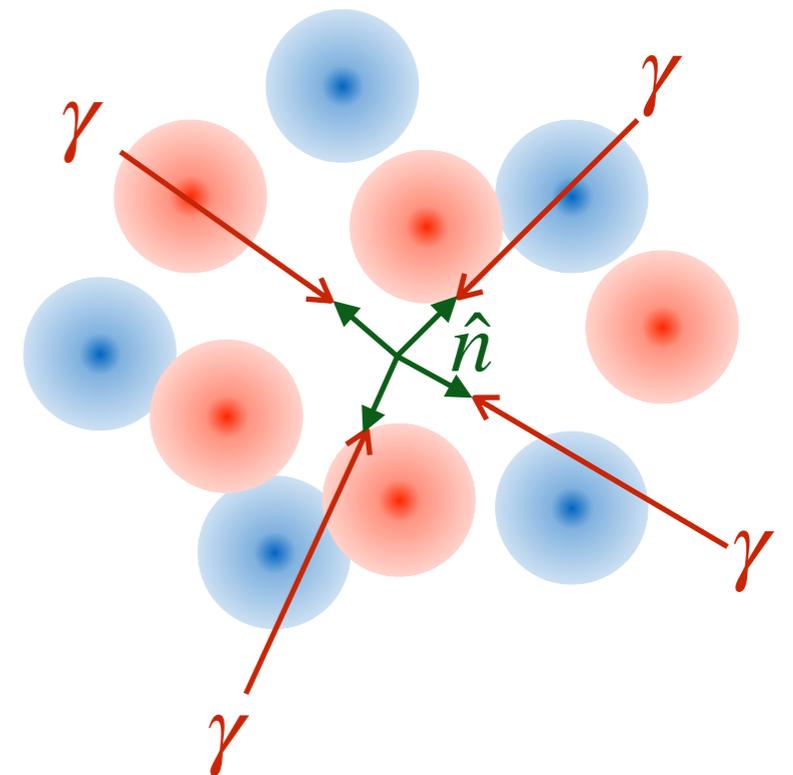
Blackbody shape:  $f_\gamma(\eta, \vec{x}, p, \hat{n}) = \frac{1}{e^{\frac{p}{T(\eta, \vec{x}, \hat{n})}} - 1}$

Up to very good approximation: **preserved** even when leaving thermal equilibrium, but becomes direction-dependent due to gravitational interactions:

redshifting along geodesics:

$$\frac{d \ln(a p)}{d\eta} = \phi' - \hat{n} \cdot \vec{\nabla} \psi$$

dilation  
gravitational Doppler



Then:  $T(\eta, \vec{x}, \hat{n}) = \bar{T}(\eta) (1 + \Theta(\eta, \vec{x}, \hat{n}))$

# Linearised Boltzmann equation

Temperature fluctuation:  $T(\eta, \vec{x}, \hat{n}) = \bar{T}(\eta) (1 + \Theta(\eta, \vec{x}, \hat{n}))$

Monopole and dipole of  $\Theta$  account for local density & bulk velocity:

$$\Theta(\eta, \vec{x}, \hat{n}) = \frac{1}{4} \delta_\gamma(\eta, \vec{x}) + \hat{n} \cdot \vec{v}_\gamma(\eta, \vec{x}) + \text{higher multipoles}$$

Linearised Boltzmann:

$$\Theta' + \hat{n} \cdot \vec{\nabla} \Theta - \phi' + \hat{n} \cdot \vec{\nabla} \psi = -\Gamma_\gamma \left( \hat{n} \cdot (\vec{v}_\gamma - \vec{v}_b) + \text{higher multipoles} \right)$$

dilation

gravitational Doppler

Thomson scattering

Thomson scattering wants to align velocity of photons vs. electron/baryons, and to wash out higher multipoles!

# Boltzmann hierarchy

Start from linearised Boltzmann and perform:

1. Fourier transformation

2. Legendre expansion  $\Theta(\eta, \vec{k}, \hat{n}) = \sum_l (-i)^l (2l + 1) \Theta_l(\eta, \vec{k}) P_l(\hat{k} \cdot \hat{n})$

$$\delta'_\gamma + \frac{4}{3}\theta_\gamma - 4\phi' = 0$$

$$\theta'_\gamma + k^2 \left( -\frac{1}{4}\delta_\gamma + \sigma_\gamma \right) - k^2\psi = \tau'(\theta_\gamma - \theta_b)$$

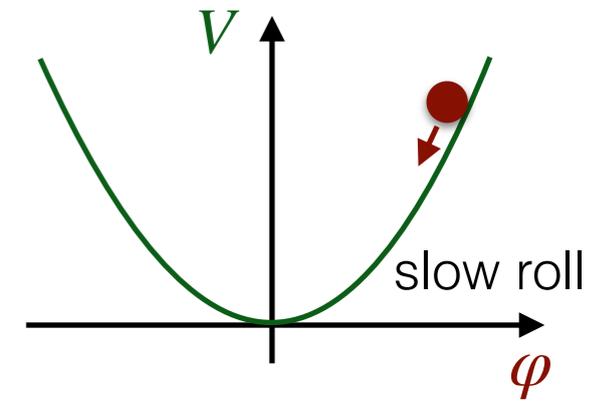
$$\Theta'_l - \frac{kl}{2l+1}\Theta_{l-1} + \frac{k(l+1)}{2l+1}\Theta_{l+1} = \tau'\Theta_l \quad \forall l \geq 2$$

relates to  $\Theta_2$

⇒ Solved together with previous equations by Einstein-Boltzmann solvers

# Stochastic theory of cosmological perturbations

# Initial conditions



Canonical single-field inflation guarantees:

A. **stochastic** perturbations with **independent** Fourier modes

B. **gaussian** statistics for each Fourier mode / each d.o.f.

⇒ described by variance(wavenumber) = **power spectrum**

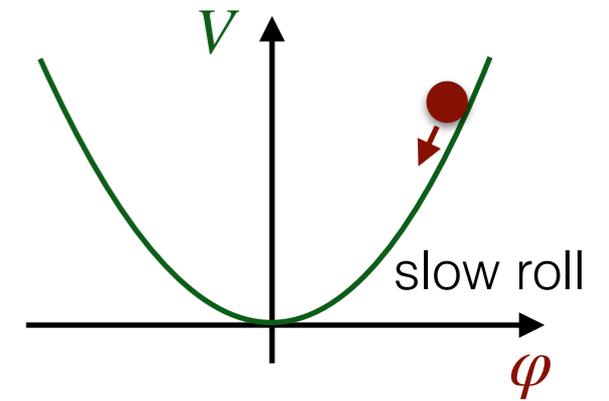
C. for each Fourier mode, all d.o.f. related to each other (fully **correlated**) on super-Hubble scales: “**adiabatic initial conditions**”

e.g. during RD: 
$$-2\psi = -2\phi \underset{\substack{\uparrow \\ \text{Einstein eq.}}}{=} \delta_\gamma = \delta_\nu = \frac{4}{3}\delta_b = \frac{4}{3}\delta_c \underset{\substack{\uparrow \\ \text{Einstein eq.}}}{=} \text{constant}$$

(Comes from  $A(\eta, \vec{x}) = \bar{A}(\eta + \delta\eta(\vec{x})) = \bar{A}(\eta) + \bar{A}'(\eta) \delta\eta(\vec{x})$ )

perturbation  $\delta A(\eta, \vec{x})$   
in adiabatic case

# Primordial power spectrum



Canonical single-field inflation guarantees:

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B. **gaussian** statistics for each Fourier mode / each d.o.f.

⇒ described by variance(wavenumber) = **power spectrum**

C. for each Fourier mode, all d.o.f. related to each other (fully **correlated**) on super-Hubble scales: “**adiabatic initial conditions**”

⇒ need power spectrum for single degree

of freedom, e.g. **curvature perturbation**  $\mathcal{R} \equiv \phi - \frac{a'}{a} \frac{v_{\text{tot.}}}{a^2}$  in Newt. Gauge

⇒ **Primordial spectrum**:  $\langle \mathcal{R}(\eta_i, \vec{k}) \mathcal{R}^*(\eta_i, \vec{k}') \rangle = \delta_D(\vec{k}' - \vec{k}) P_{\mathcal{R}}(k)$

D. **Power law, nearly scale-invariant** spectrum:  $P_{\mathcal{R}}(k) = \frac{2\pi^2}{k^3} A_s \left( \frac{k}{k_*} \right)^{n_s - 1}$

# Transfer functions

For each Fourier mode  $\vec{k}$ :

- all perturbations  $\rightarrow$  system of linear coupled differential equations
- adiabatic ICs  $\rightarrow$  single constant of integration  $\mathcal{R}(\eta_{\text{ini}}, \vec{k})$
- $\forall A \in \{ \phi, \psi, \delta_X, \theta_X, \Theta_\ell, \dots \}$

$$A(\eta, \vec{k}) = T_A(\eta, k) \mathcal{R}(\eta_i, \vec{k})$$

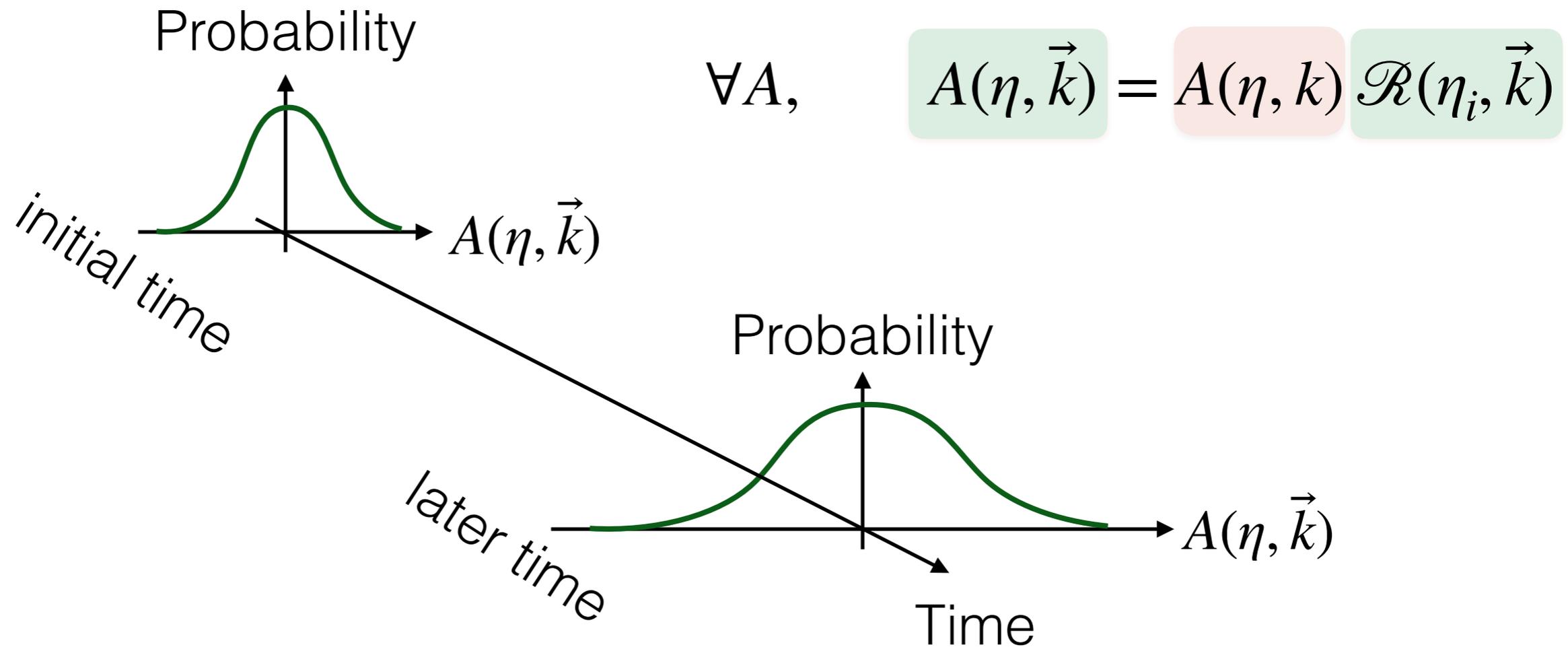
stochastic Fourier mode

stochastic IC

Deterministic solution of e.o.m. normalised to  $\mathcal{R} = 1$   
 = transfer function of  $A$

Isotropic background  $\Rightarrow$  depends only on  $k$   
 $\Rightarrow$  denoted later as  $A(t, k)$

# Linear transport of probability



Linearity of solutions  $\Rightarrow$  probability shape always preserved  
(standard model: Gaussian)

$\Rightarrow$  variance evolves like square of transfer function

# Power spectrum

Adiabatic initial conditions

⇒ for any perturbation at any time:

$$\langle A(\eta, \vec{k}) A^*(\eta, \vec{k}') \rangle = A(\eta, k) A^*(\eta, k') \langle \mathcal{R}(\eta_i, \vec{k}) \mathcal{R}^*(\eta_i, \vec{k}') \rangle$$

$$= |A(\eta, k)|^2 P_{\mathcal{R}}(k) \delta_D(\vec{k} - \vec{k}')$$

transfer function of  $A$

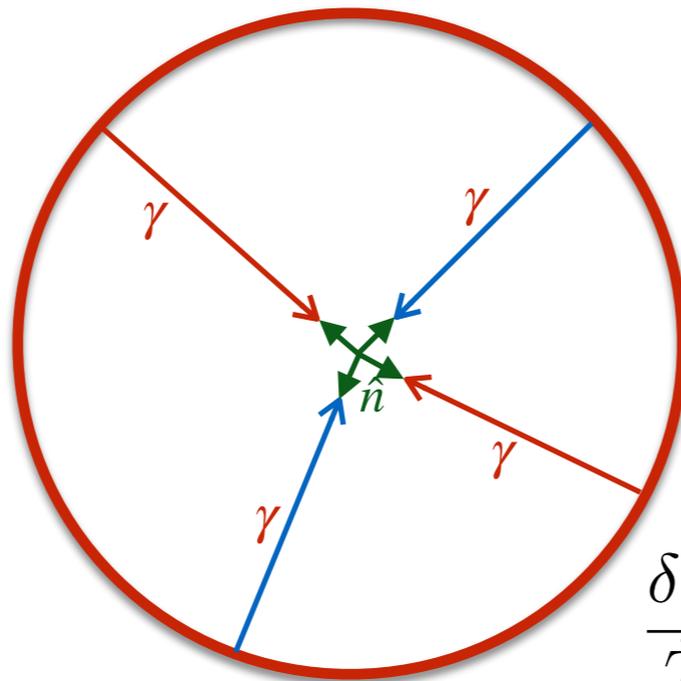
power spectrum  $P_A(\eta, k)$  of  $A$  at  $\eta$

primordial curvature spectrum

# Spectrum of temperature anisotropies

# Temperature multipoles

$g(\eta)$  very peaked at  $\eta_{\text{dec}}$   
 $\Downarrow$   
 last scattering sphere



$$\frac{\delta T}{\bar{T}}(\hat{n}) = \Theta(\eta_0, \vec{\sigma}, -\hat{n}) = \sum_{lm} a_{lm} Y_{lm}(\hat{n})$$

inversion + Fourier + Legendre  $\Rightarrow$   $a_{lm} = (-i)^l \int \frac{d^3 \vec{k}}{2\pi^2} Y_{lm}(\hat{k}) \Theta_l(\eta_0, \vec{k})$

stochastic, Gaussian  $\longleftarrow$  stochastic, Gaussian

correlation/variance  $\Rightarrow \langle a_{lm} a_{l'm'}^* \rangle = \delta_{ll'}^K \delta_{mm'}^K \left[ \frac{2}{\pi} \int dk k^2 \Theta_l^2(\eta_0, k) P_{\mathcal{R}}(k) \right]$

photon transfer function primordial spectrum

# Temperature power spectrum

Defined as:

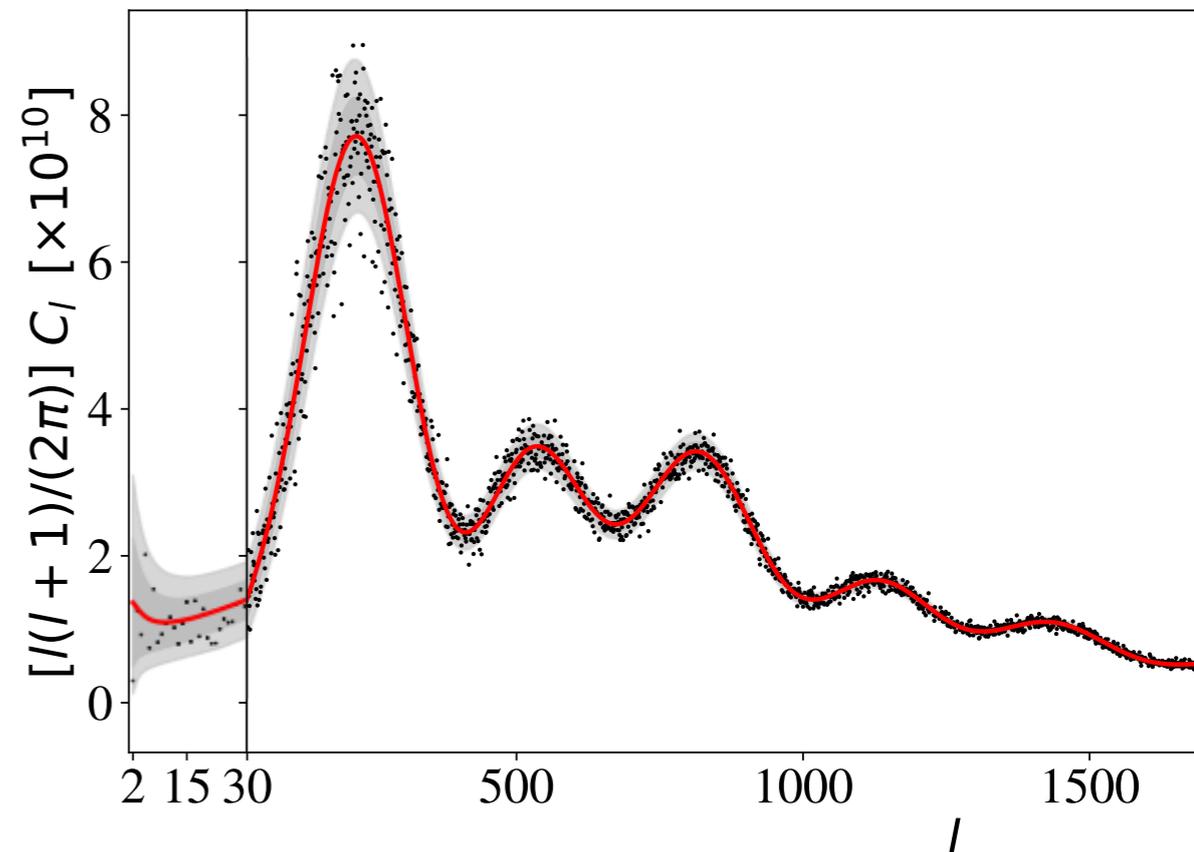
$$C_l = \langle a_{lm} a_{lm}^* \rangle = \frac{2}{\pi} \int dk k^2 \Theta_l^2(\eta_0, k) P_{\mathcal{R}}(k)$$

photon transfer function     primordial spectrum

theory ↔ observations

Estimator:  $\hat{C}_l(a_{lm}) \equiv \frac{1}{2l+1} \sum_{-l \leq m \leq l} |a_{lm}|^2$

Cosmic variance:  $\langle (\hat{C}_l - C_l)^2 \rangle = \frac{2}{2l+1} C_l^2$



# Physics of temperature anisotropies

# “Line-of-sight” integral in Fourier space

Boltzmann hierarchy  $\Rightarrow$  formal solution Zaldarriaga & Harari [astro-ph/9504085](https://arxiv.org/abs/astro-ph/9504085):

$$\Theta_l(\eta_0, \vec{k}) = \int_{\eta_{\text{ini}}}^{\eta_0} d\eta \left\{ g(\Theta_0 + \psi) j_l(k(\eta_0 - \eta)) \right. \\ \left. + g k^{-1} \theta_b j'_l(k(\eta_0 - \eta)) \right. \\ \left. + e^{-\tau} (\phi' + \psi') j_l(k(\eta_0 - \eta)) \right\}$$

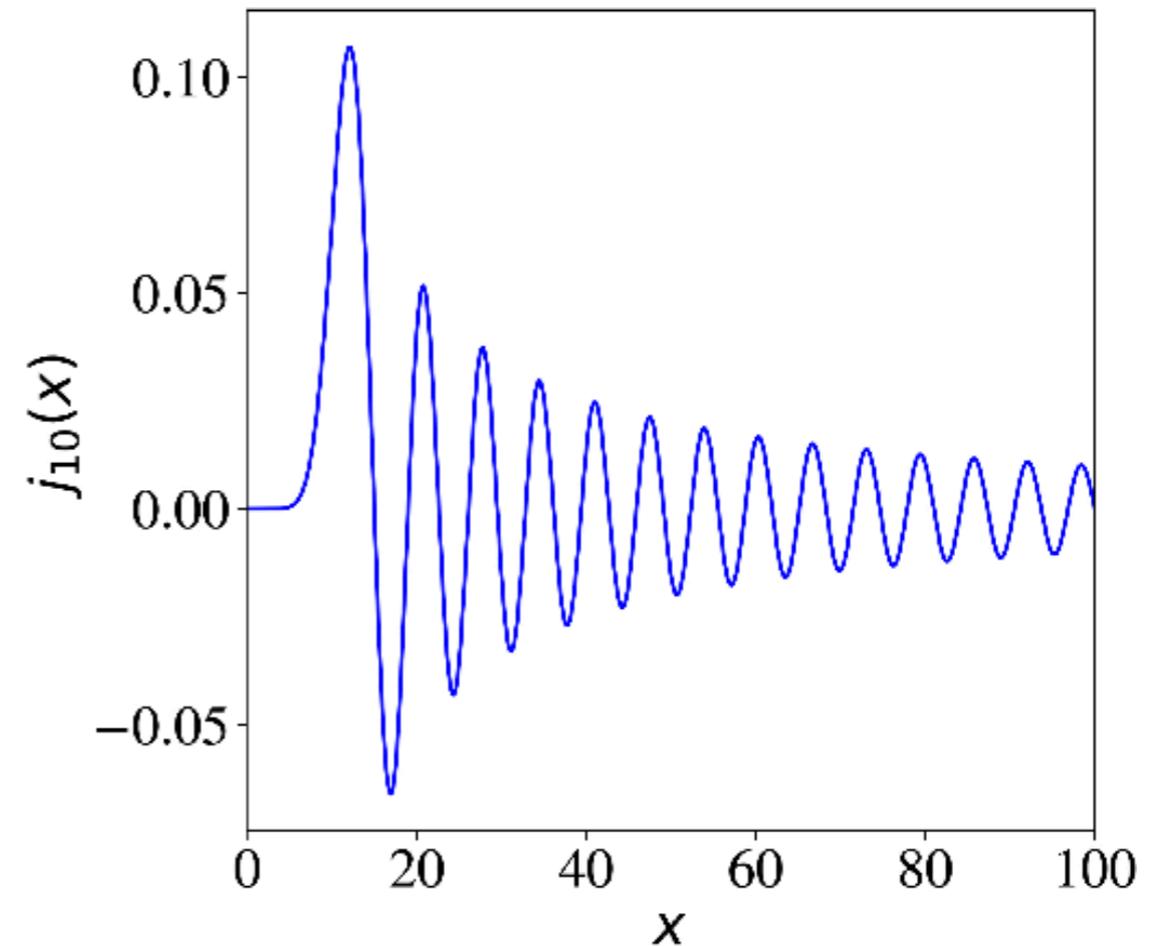
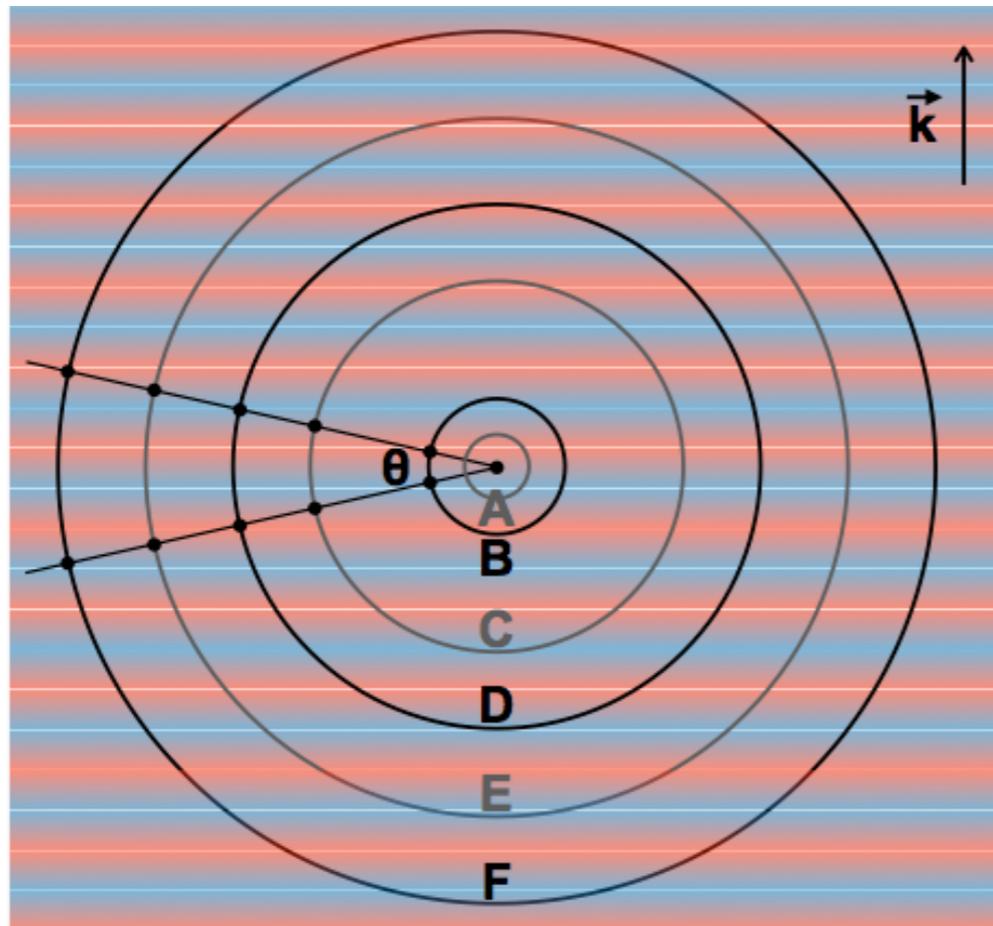
valid both for  
single mode  $\vec{k}$  or  
transfer function with  $k$

structure:  $\int d\eta f(\eta) A(\eta, \vec{k}) j_\ell(k(\eta_0 - \eta))$

“Physical effects relevant at times described by  $f(\eta)$   
imprint CMB photon anisotropies described in Fourier space by  $A(\eta, \vec{k})$ ,  
that project to multipole space according to  $j_\ell(k(\eta_0 - \eta))$ ”

# Angular projection of Fourier modes

Role of  $j_\ell(k(\eta_0 - \eta))$  ?



$$\text{Main contribution: } \theta = \frac{\pi}{l} = \frac{\lambda/2}{d_a} = \frac{a(\eta) \pi/k}{a(\eta) (\eta_0 - \eta)} \Leftrightarrow l = k(\eta_0 - \eta)$$

Other contributions: harmonics