







Scalar Field Damping in Thermal Plasma

Mubarak Abdallah

In collaboration with Marco Drewes

ICTP-EAIFR

Catholic University of Louvain (UCLouvain)



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Motivation

- This is part of an ongoing study on Inflation.
- Hot big bang initial conditions of the radiation dominated universe were created during reheating.
- Reheating temperature is related to the dissipation rate

$$T_R \equiv \left(\frac{90}{8\pi^3 q_*}\right)^{1/4} \sqrt{\Gamma M_P}.$$
 1305.0267

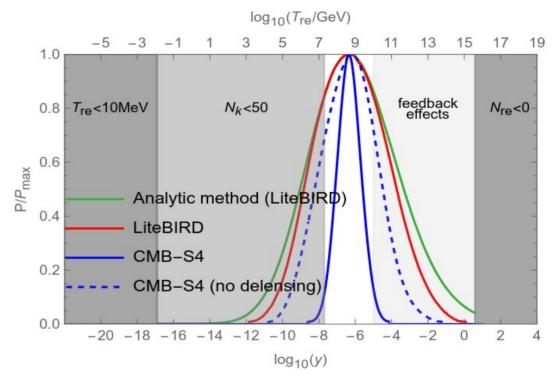
 Correction to M. Drewes, J. U. Kang, Nucl. Phys. B 875 (2013) 315–350, arXiv.1305.0267 [hep-ph]

Outline

- Importance of Reheating.
- Introduce the tools from thermal field theory.
- Apply it to a specific model & calculate the dissipation rate. Show how our results differ from 1305.0267.
- Show possible thermal effects which arises.
- Summery and future steps.

Measuring the reheating temperature

• Reheating temperature T_{re} could be measured (constrained) by future experiments



Marco Drewes,Lei Ming,Isabel Oldengott, arXiv:2303.13503 [hep-ph]

Self-energy and decay rate

• For a general linear Inflation couplings $\phi \mathcal{O}[\mathcal{X}_i]$, We define the self-energies at leading order in the tiny Inflaton couplings as

$$\Pi^{>}(x_1, x_2) = \langle \mathcal{O}(x_1)\mathcal{O}(x_2)\rangle, \ \Pi^{<}(x_1, x_2) = \langle \mathcal{O}(x_2)\mathcal{O}(x_1)\rangle.$$

The spectral selfe-nergy
$$\Pi^{-}(x_1, x_2) = \Pi^{>}(x_1, x_2) - \Pi^{<}(x_1, x_2)$$
,

Determines the total relaxation rate $\Gamma_{m{q}} \sim \Pi_{m{q}}^-(\Omega_{m{q}})$

Spectral Function

• Define the spectral function $\Delta^-(x_1, x_2)$

$$\Delta^{-}(x_1, x_2) = i \left(\langle \phi(x_1) \phi(x_2) \rangle - \langle \phi(x_2) \phi(x_1) \rangle \right).$$

The Dyson-Schwinger equation

$$(\Box_1 + m^2)\Delta_{\mathcal{C}}(x_1, x_2) + \int_{\mathcal{C}} d^4x' \Pi_{\mathcal{C}}(x_1, x') \Delta_{\mathcal{C}}(x', x_2) = -i\delta_{\mathcal{C}}(x_1 - x_2) .$$

The Spectral Density

The spectral density is given as

$$\rho_{\mathbf{q}}(\omega) = \frac{-2\mathrm{Im}\Pi_{\mathbf{q}}^{R}(\omega) + 2\omega\epsilon}{(\omega^{2} - m^{2} - \mathbf{q}^{2} - \mathrm{Re}\Pi_{\mathbf{q}}^{R}(\omega))^{2} + (\mathrm{Im}\Pi_{\mathbf{q}}^{R}(\omega) + \omega\epsilon)^{2}}.$$

The quasi-particle dispersion relation, for weakly coupled theories

$$\omega^2 - \mathbf{q}^2 - m^2 - \operatorname{Re}\Pi_{\mathbf{q}}^R(\omega) = 0,$$

In the Zero width approximation

$$\rho_{i\mathbf{p}}^{0}(p_0) = 2\pi \mathcal{Z}_{i\mathbf{p}} \operatorname{sign}(p_0) \delta(p_0^2 - \Omega_{i\mathbf{p}}^2) + \rho_{i\mathbf{p}}^{\operatorname{cont}}(p_0).$$

The Model

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - \frac{1}{2} m^2 \phi^2 + \sum_{i=1}^{2} \left(\frac{1}{2} \partial_{\mu} \chi_i \partial^{\mu} \chi_i - \frac{1}{2} m_i^2 \chi_i^2 - \frac{\lambda_i}{4!} \chi_i^4 \right) - g \phi \chi_1 \chi_2 + \mathcal{L}_{\text{bath}}.$$

1305.0267

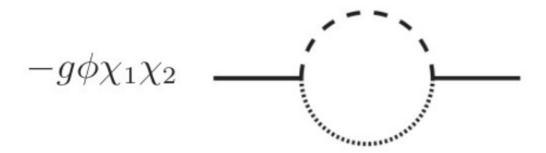
• The χ_i fields receive thermal masses

$$M_i^2 = m_i^2 + \lambda_i \int \frac{d^3 \mathbf{p}}{(2\pi)^3} \frac{f_B(\omega_i)}{2\omega_i} \approx m_i^2 + \frac{\lambda_i}{24} T^2.$$



Leading order contribution

Leading order contribution to the relaxation rate



$$\Gamma_{\mathbf{q}} = -ig^2 \int \frac{d^4p}{(2\pi)^4} \left(1 + f_B(p_0) + f_B(\omega - p_0)\right) \rho_{1\mathbf{p}}(p_0) \rho_{2\mathbf{q}-\mathbf{p}}(\omega - p_0).$$

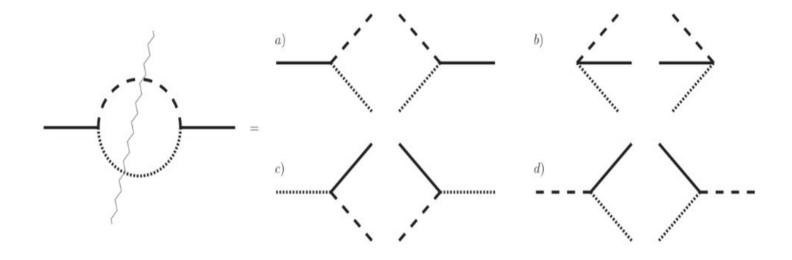
Diagrammatic Interpretation of the decay rate

In the Zero width approximation for quasi-particles

$$\Gamma_{\mathbf{q}} = -ig^2 \int \frac{d^3\mathbf{p}}{(2\pi)^2} \frac{1}{4\Omega_2\Omega_1}$$

$$\times \left[\left((f_1+1)(f_2+1) - f_1 f_2 \right) \left(\delta(\omega - \Omega_1 - \Omega_2) - \delta(\omega + \Omega_1 + \Omega_2) \right) + \left((f_1+1)f_2 - (f_2+1)f_1 \right) \left(\delta(\omega - \Omega_1 + \Omega_2) - \delta(\omega + \Omega_1 - \Omega_2) \right) \right]$$

Diagrammatic Interpretation



Result

$$\Pi_{0}^{-}(\omega) = 2i \left(\mathcal{D}_{0}^{[g]}(\omega) + \mathcal{S}_{0}^{[g]}(\omega) \right)$$

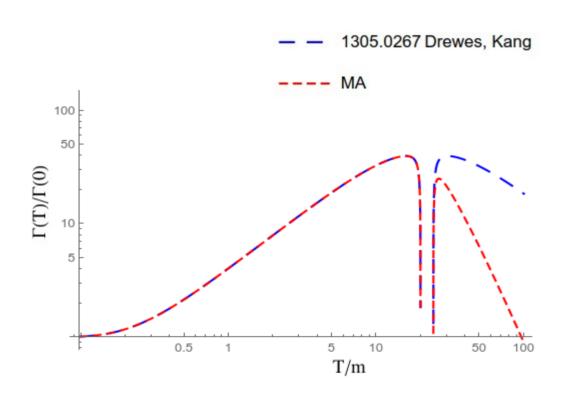
$$\mathcal{D}_{0}^{[g]}(\omega) = \frac{A}{\bar{\omega}_{1} + \bar{\omega}_{2}} \Theta(\omega - M_{1} - M_{2})$$

$$\mathcal{S}_{0}^{[g]}(\omega) = \frac{A}{\bar{\omega}_{1} - \bar{\omega}_{2}} \left(\Theta(M_{1} - \omega - M_{2}) - \Theta(M_{2} - \omega - M_{1}) \right)$$

$$A = -\frac{g^{2}}{8\pi} \sqrt{\bar{\omega}_{1}^{2} - M_{1}^{2}} \left(1 + f(\bar{\omega}_{1}) + f(\bar{\omega}_{2}) \right)$$

$$\bar{\omega}_{1} = \frac{\omega^{2} + M_{1}^{2} - M_{2}^{2}}{2\omega} \qquad \bar{\omega}_{2} = \frac{\omega^{2} - M_{1}^{2} + M_{2}^{2}}{2\omega}$$

Behavior of the decay rate



- The rate Γ as a function of T, normalized to its zero temperature value.
- Allowed regimes as a function of temperature

$$M \geqslant M_1 + M_2 \qquad \phi \leftrightarrow \chi_1 \chi_2$$

 $|M_1 - M_2| \stackrel{\circ}{\geqslant} M. \quad \chi_i \phi \leftrightarrow \chi_j$

For this choice of parameters $\lambda 1 = 0.01$, $\lambda 2 = 1$, m1 = m2 = 0.001m.

Summary

- On-shell contribution to the relaxation rate gives rise to processes such as decays $\phi \leftrightarrow \chi_1 \chi_2$, and inverse decays $\chi_i \phi \leftrightarrow \chi_j$
- These processes are relevant to the reheating process.
- The damping rate can be used to estimate the dissipation rate during inflation itself (interesting for warm inflation).
- These kinematic effects can be relevant for Baryogenesis, dark matter production.
- Next step is to include higher order and study landau damping effects.

Murakoze Cyane