Primordial Black Holes generated from SUSY Breaking in Non-oscillatory Inflation Models

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Outline

Supergravity Framework

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Sgoldstinoless Model

$$W = \left(1 - \frac{S}{\sqrt{3}}\right)^3 f(Z)$$

$$K = K_1 \left(Z, \overline{Z}\right) - 3 \log \left[1 - \frac{|S|^2}{3} + \frac{|S|^4}{\Lambda^2}\right],$$

(L. Heurtier, A.M., L. Wacqwez, 2023)

- f(Z) is a holomorphic function: $\overline{f(Z)} = f(-\overline{Z})$, $f(0) \neq 0$, S is the goldstino superfield. (Dall'Agata and Zwirner 2014)
- The nilpotent limit: $S^2 = 0 \Rightarrow$ scalar potential $V_{\text{nil}}(z)$.
- $\bullet\,$ Inflaton effective potential after integrating out $S\,$

$$V_{\rm eff}(z) = V_{\rm nil}(z) - \frac{\Lambda^2 V_{\rm nil}(z)^2}{2 \left[6m_{3/2}^2 + \Lambda^2 V_{\rm nil}(z) \right]}.$$

Sgoldstinoless Model

- The gravitino mass: $m_{3/2}^2 \simeq f (i z)^2$.
- Sgoldstino mass

$$m_s^2 = \frac{36 \, m_{3/2}^2 + 6 \, \Lambda^2 \, V_{\mathsf{nil}}}{\Lambda^2} \,,$$

• Large UV scale, integrate out the sgoldstino with:

$$\langle s
angle = rac{\Lambda^2 V_{
m nil}}{2\sqrt{3} \left[\Lambda^2 V_{
m nil} + 6 \, m_{3/2}^2
ight]}$$

- $\bullet\,$ The nilpotent limit is equivalent to sending the UV scale $\Lambda \to 0$
- The nilpotent limit: $S^2 = 0 \Rightarrow$ scalar potential $V_{\rm nil}(z)$.

Model: Canonical Kähler

$$K_{1} = \frac{1}{2} \left(Z + \overline{Z} \right)^{2}, \qquad f(Z) = f_{0} - \sqrt{V_{0}} \log \left(1 + e^{2\sqrt{2} i b Z} \right).$$
$$V_{\text{nil}} = -f' \left(iz \right)^{2} = \frac{8 b^{2} V_{0}}{\left(e^{2b\varphi} + 1 \right)^{2}}.$$



Features

- Cosmological Seesaw Mechanism: Two different plateaus with a hierarchical difference $\sim V_0$, .
- Inflation tail: $\varphi \ll -1 \Rightarrow H^2 \sim V_0/3$
- DE tail: $V \ll 1$ as $\varphi \gg 1$, and $m_{3/2} \sim f_0$
- Corrections due to integrating out the sgoldstino
 - Bumps, dips or inflection points in the inflation tail.
 - These Corrections (proportional to $\Lambda^2)$ vanish as $\varphi\gg 1$
 - Corrections depend on SUSY breaking f_0 and the UV scale Λ
 - For specific values of these two parameters in order to obtain Ultra Slow Roll (USR) phase, during inflation, leading to the formation of PBHs and GWs.

Cosmological History

- Inflation phase: along the Inflation tail $\varphi \ll -1$, for 50-60 efolds.
- Reheating Mechanism: Gravitational Reheating

$$\rho_{\text{grav}}|_{t=t_{\text{end}}} = \frac{g_{\star,\text{end}}q}{1440\pi^2} \left(\frac{H_{\text{end}}}{M_{\text{Pl}}}\right)^2 \rho_{\text{end}}, \qquad (1)$$

• Kination era: Energy density dominated almost entirely by the kinetic energy of φ . $w_{\rm kin} \approx 1$, $\rho_{\rm kin} \propto a^{-6}$, $a \propto t^{1/6}$ and $H = \frac{1}{3t}$.

$$\varphi(N) = \varphi_{\rm kin} + \sqrt{6} \, M_{\rm Pl} \, \left(N - N_{\rm kin} \right),$$

• Radiation Domination: SM radiation starts dominating over the scalar field energy density.

$$\varphi_{\rm rad} = \varphi_{\rm kin} + \sqrt{\frac{2}{3}} M_{\rm Pl} \ln\left(\frac{H_{\rm kin}}{H_{\rm rad}}\right), \quad H_{\rm rad}^2 = \frac{2\rho_{\rm rad}}{3M_p^2}.$$
 (2)

• Matter Domination: Matter-radiation equality at $N_{
m eq} = \log\left(rac{\Omega_{
m rad,0}}{\Omega_{
m m,0}}
ight)$,

where the matter energy density begins to dominate the universe.

Cosmological History and Observables



- We have scanned over f_0, Λ, V_0, b
- Large SUSY breaking scale $f_0 > 10^{-8}$ in order to avoid long tracking regime.



Primordial Black Holes

- Breaking of SUSY at a scale comparable to the inflation scale typically produces inflection points along the inflation potential.
- \bullet We are interested in USR regime around inflection point: $V'\approx 0$,



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Primordial Black Holes

The mass of a PBH that forms during the kination dominated era

$$M(k) = 4\pi\gamma \left(\frac{\pi^2 g_*^{\rm eq}}{45}\right)^{\frac{1}{1+3w}} \left(\frac{g_s^{\rm eq}}{g_s(T_{kin})}\right)^{\frac{3w-1}{3(3w+1)}} T_{kin}^{-\frac{3w-1}{3w+1}} \left(\frac{a_{\rm eq}T_{\rm eq}}{k}\right)^{\frac{3(1+w)}{3w+1}}$$

• The fractional abundance of PBHs

$$\begin{split} f_{\mathsf{PBH}}(M) &= \frac{\Omega_{\mathsf{PBH}}(M)}{\Omega_c} \\ &= \frac{\gamma}{T_{\mathsf{eq}}} \left(\frac{g_s(T_{\mathrm{kin}})}{g_s(T_{\mathsf{eq}})}\right)^{1/3} \left(\frac{\Omega_m h^2}{\Omega_c h^2}\right) \left(\frac{90}{\pi^2 g_*(T_{\mathrm{kin}})}\right)^{\frac{w}{1+w}} (4\pi\gamma)^{\frac{2w}{1+w}} T_{\mathrm{kin}}^{\frac{1-3w}{1+w}} \beta(M) M^{-\frac{2w}{1+w}} \end{split}$$

• Modes that re-enter the horizon then form black holes with the usual mass and dark matter fraction spectrum: sending $T_{\rm kin} \rightarrow T_{\rm eq}$ and $w \rightarrow 1/3$.

Primordial Black Holes

$$\begin{split} \text{Mass Fraction: } \beta(k) &= 2 \int_{\delta_c(k)}^{\infty} \frac{1}{\sqrt{2\pi\sigma^2(k)}} e^{-\delta^2/2\sigma^2(k)} \mathsf{d}\delta \,, \\ \delta_c(k) &= \frac{3(1+w)}{5+3w} \sin^2\left(\frac{\pi\sqrt{w}}{1+3w}\right) \,. \end{split}$$



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Scalar Induced Gravitational Waves (SIGW)

Preliminary Results



Conclusions

- Sgoldstinoless models of non-oscillatory inflation provide an intersting framework that can combine Infation, SUSY breaking and DE.
- Integrating out the sgoldstino with a finite mass, has many interesting cosmological implications.
- SUSY breaking effects with finite sgodstino mass can generate an inflection point in the inflation tail, that may feature an ultra slow roll phase before the end of inflation.
- This may be translated into an enhancement in the primordial power spectrum, that may generate primordial Black Holes with reasonable abundnace accounting for a DM fraction and can source gravitational waves.
- Kination phase has interesting implications where the PBH mass fraction is enhanced more than fractions formed during radiation.

