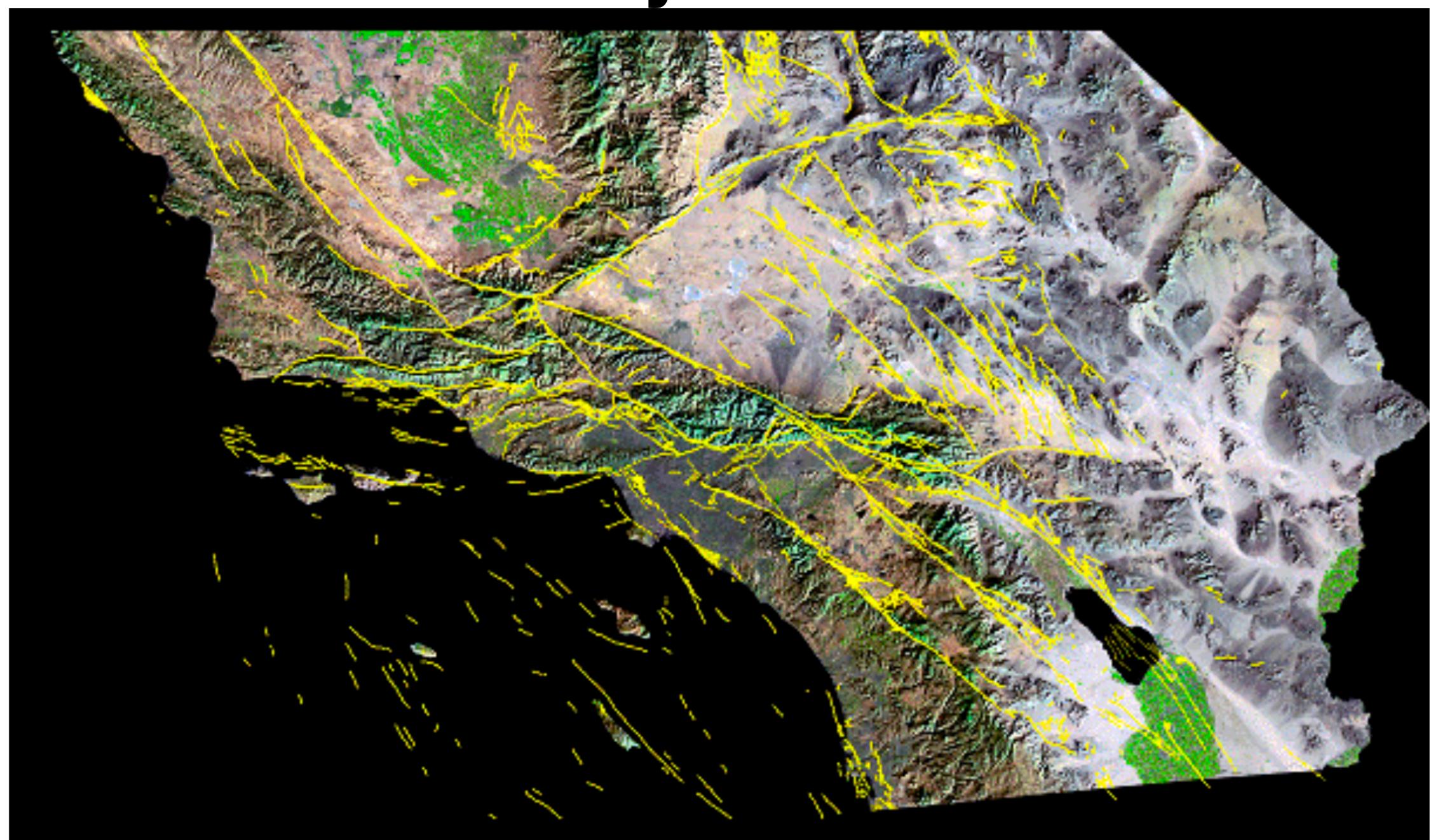
Earthquake Physics and Fault Geometry

Simple Models and Complex Behavior, Chapter 1.

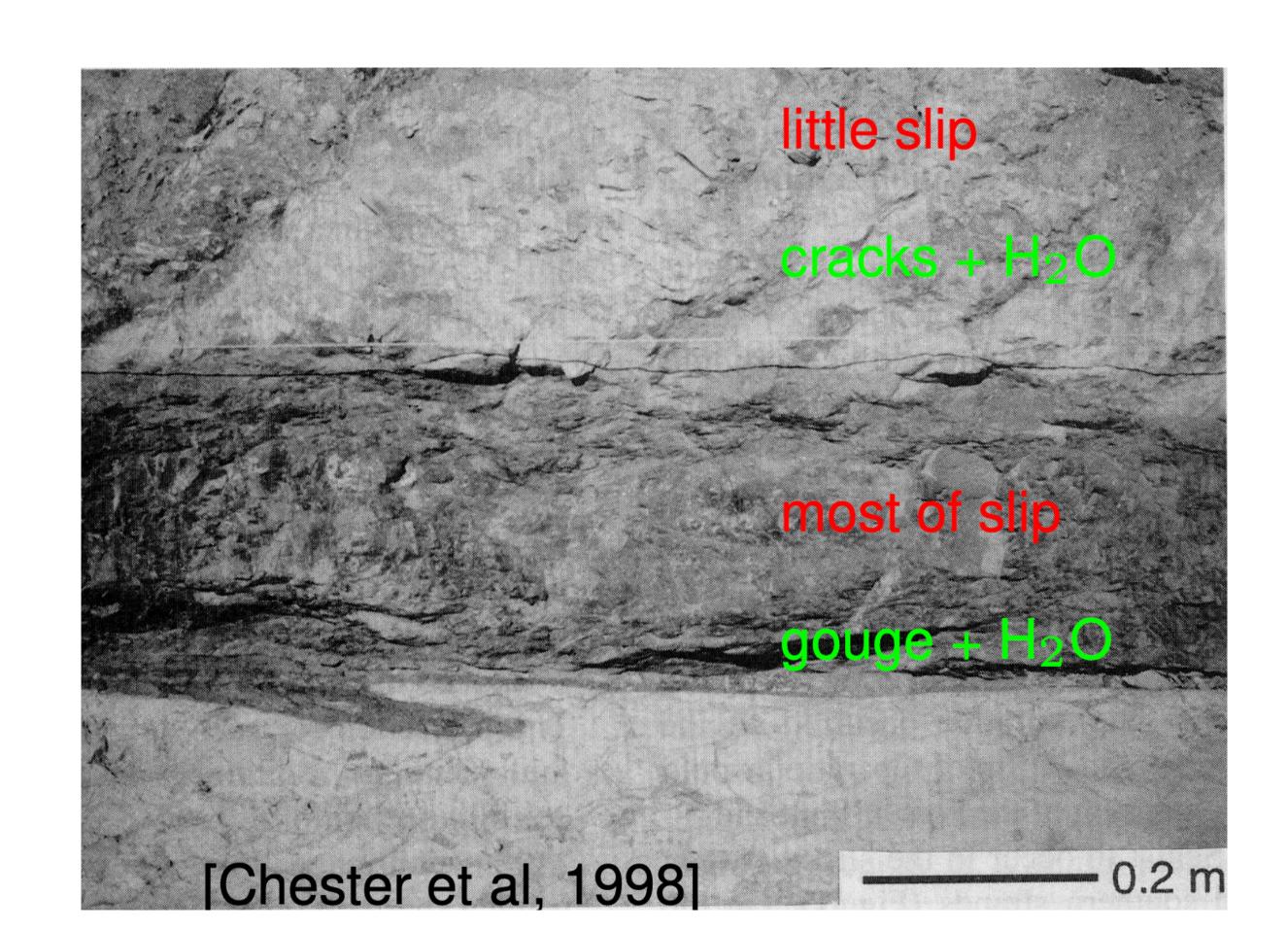
Bruce E Shaw, Columbia University | ICTP Trieste, October 2023

Fault Systems



Faults

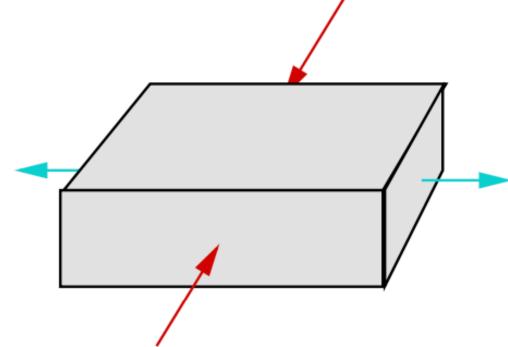




Equations of Motion





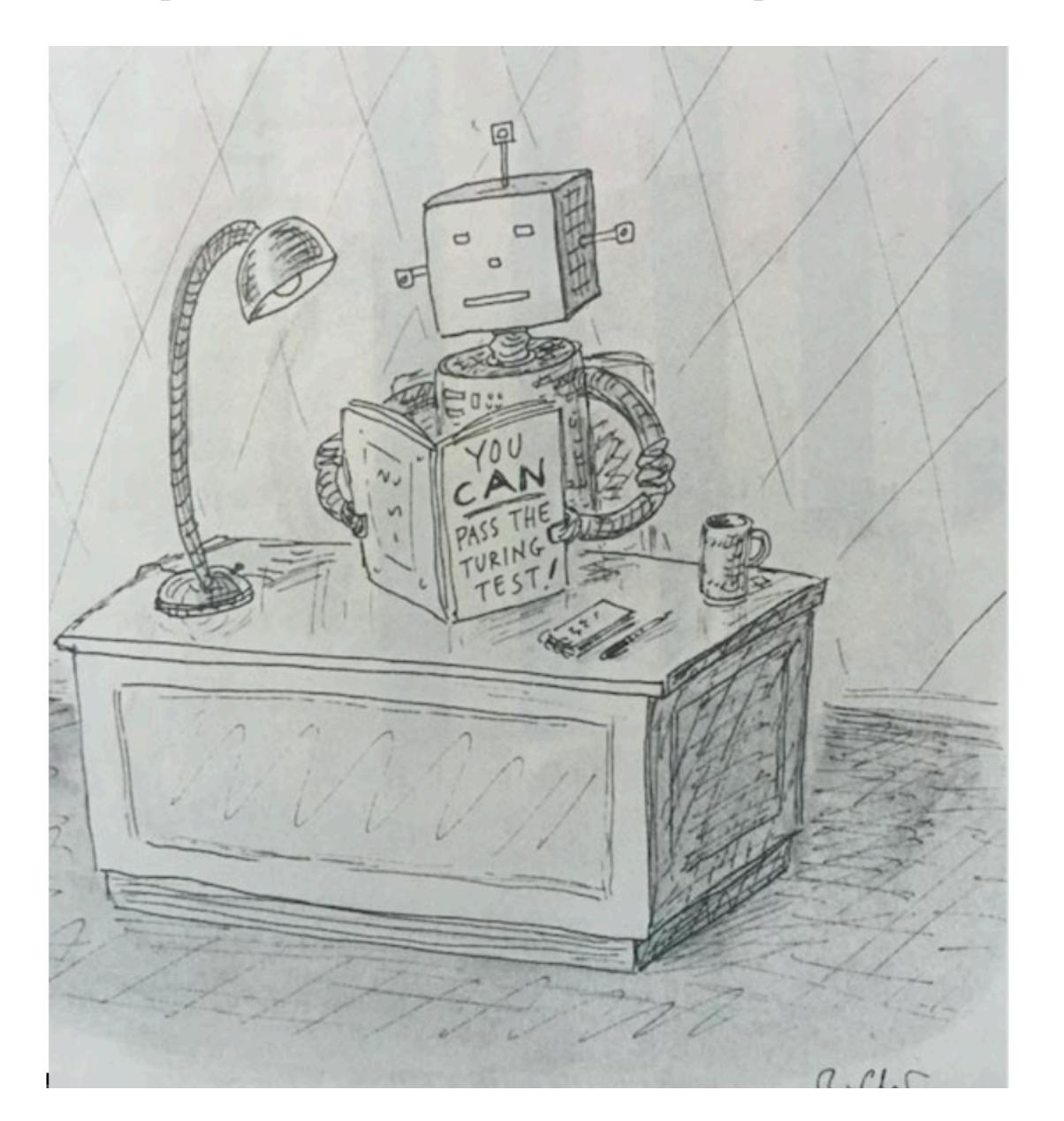


? Boundary Conditions

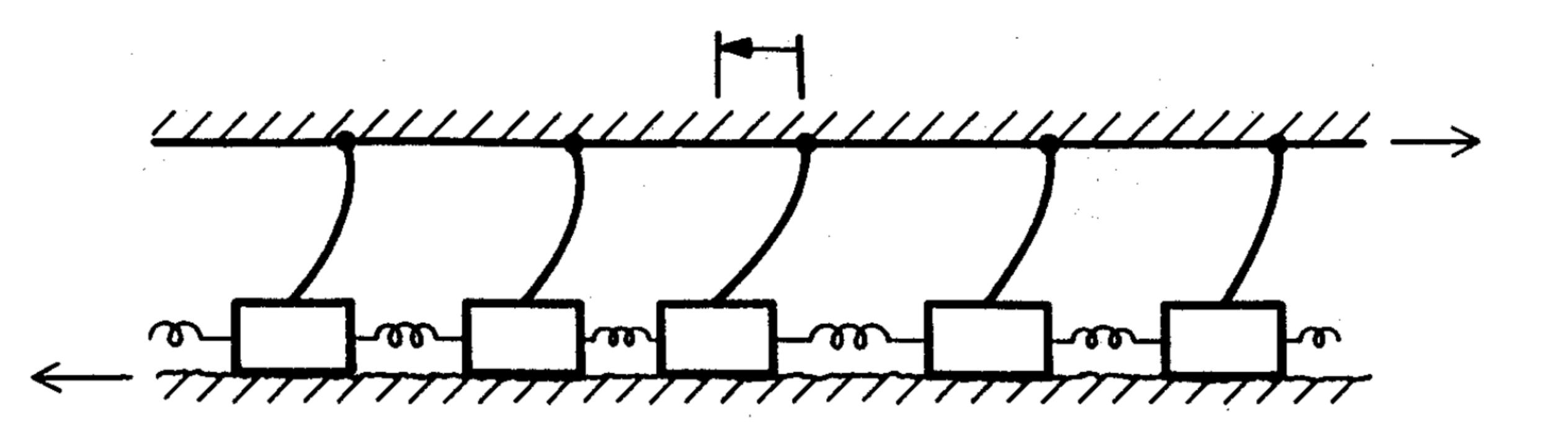
Condition on Boundary?

- Shape of Boundary?
- Boundary Layer??

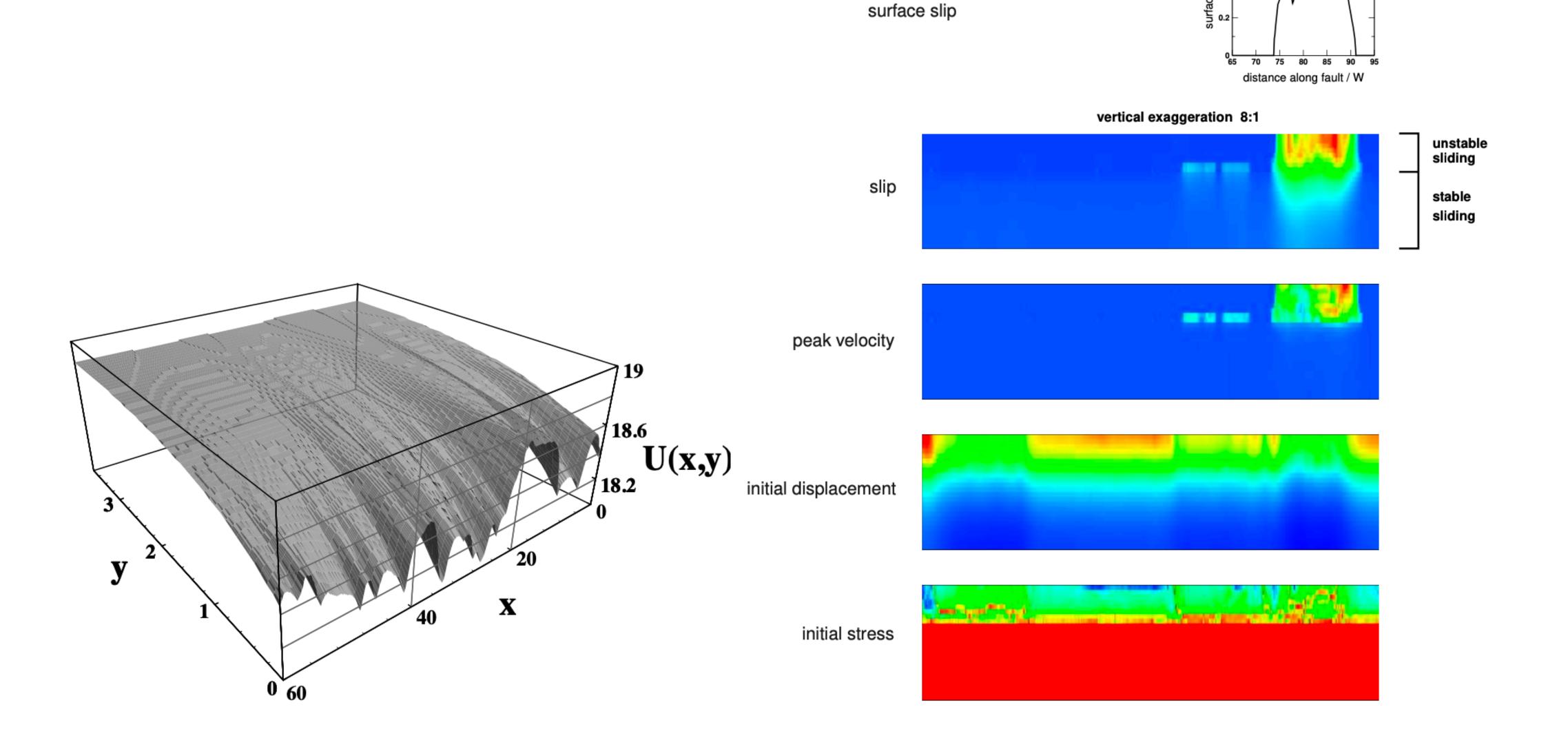
Goal: Develop a model that can pass the Turing Test



Simplest earthquake model



2D & 3D earthquake model



Equations

$$\frac{\partial^2 U}{\partial t^2} = \nabla^2 U$$

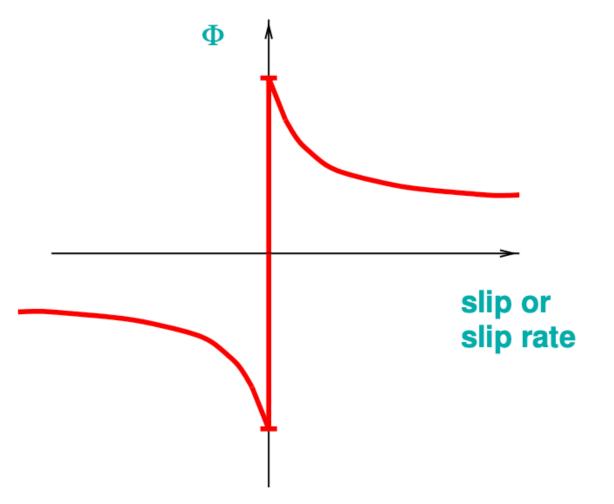
Wave equation Bulk

$$\left| rac{\partial U}{\partial t}
ight|$$
 loading

 $= \nu \qquad \nu \ll 1 \text{ slow uniform loading}$

$$\left. \frac{\partial U}{\partial n} \right|_{\text{fault}} = \Phi$$

Fault Boundary Condition



If frictional weakening:

⇒ chaotic slip sequences

Frictional weakening → complexity



weakening → complexity

100.0 X

200.0

strengthening → periodic

100.0 X

1.0 0.5 0.0 -1.0 0.0 -1.0 0.0

200.0

Frictional weakening from frictional heating

$$\Phi = N\mu$$

$$\mu = \begin{cases} \left[-\mu_0 , \mu_0 \right] & \frac{\partial S}{\partial t} = 0; \\ -\mu_0 (1 - \sigma) , \mu_0 (1 - \sigma) & \frac{\partial S}{\partial t} < 0, \frac{\partial S}{\partial t} > 0 \end{cases}$$

$$N = N_0 - \alpha Q$$

$$\frac{\partial Q}{\partial t} = -\gamma Q + \Phi \frac{\partial S}{\partial t}$$

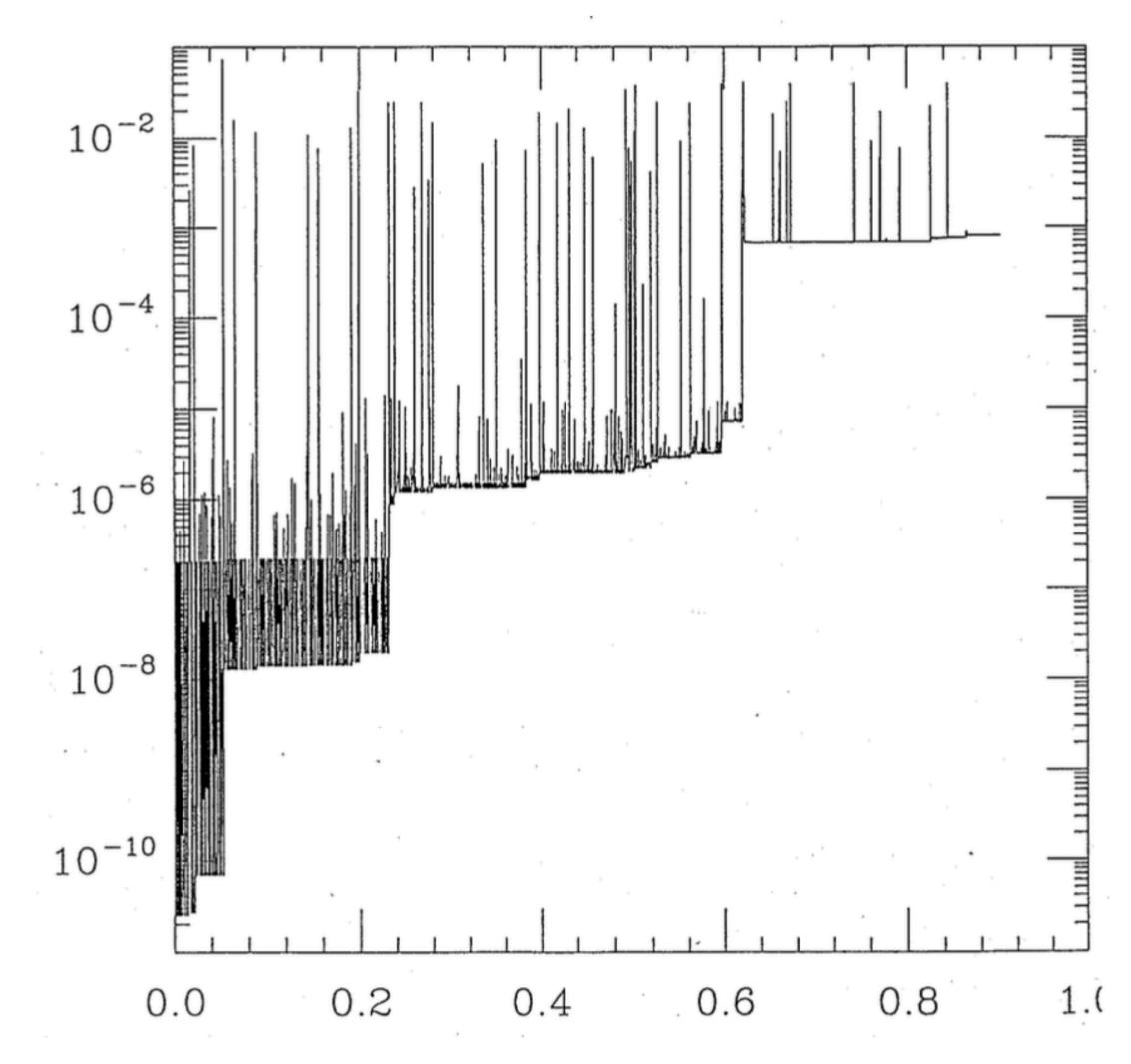
 $\gamma \ll 1 \longrightarrow \text{slip weakening}$:

$$\Phi = N_0 \mu e^{-\alpha \mu (S - S_0)}$$

 $\gamma\gg 1$ — velocity weakening:

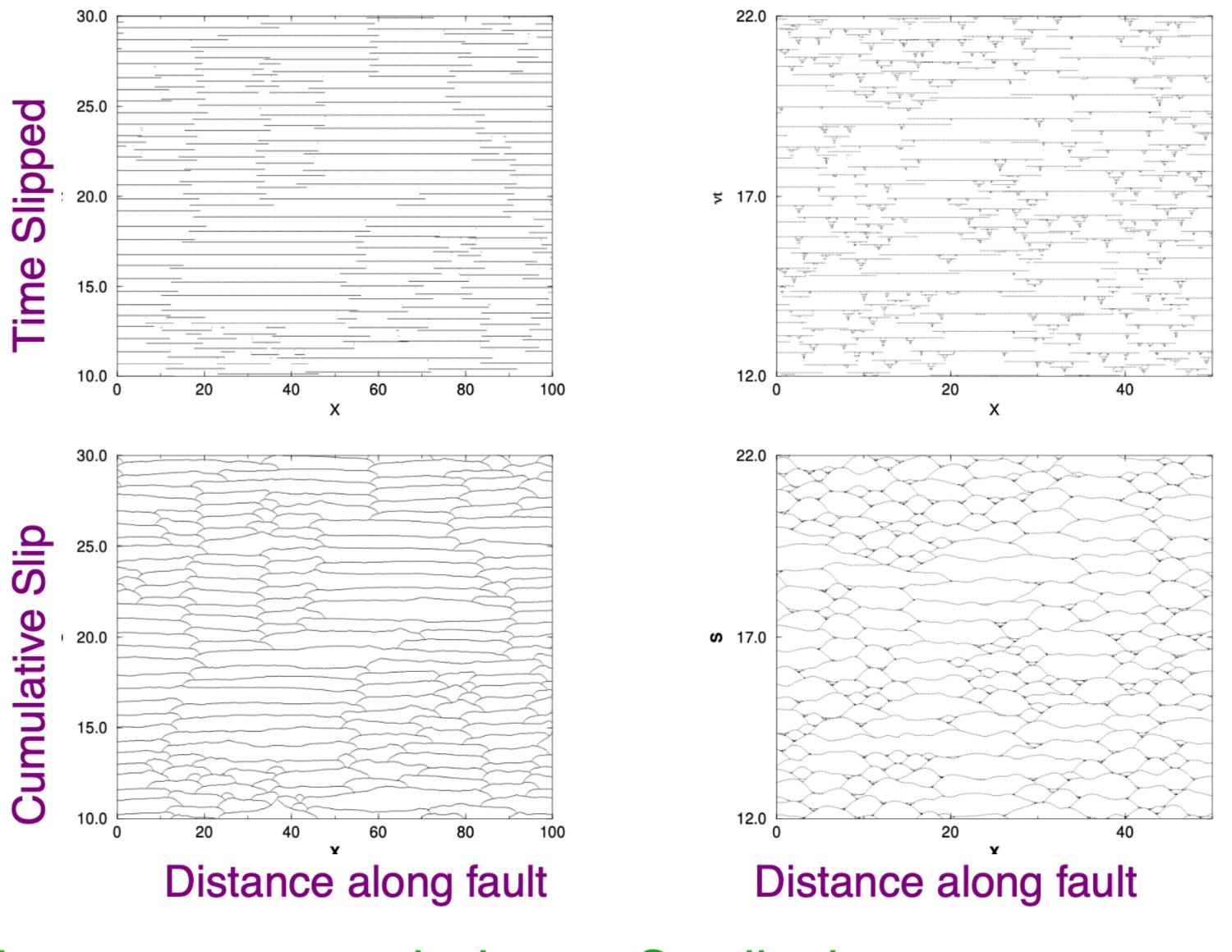
$$\Phi = \frac{N_0 \mu}{1 + \frac{\alpha}{\gamma} \mu \frac{\partial S}{\partial t}}$$

Chaos



- Exponential divergence
- Huge Lyopunov exponents
- Divergence happens during large events
- Rules out predicting past next large event, but not next large event

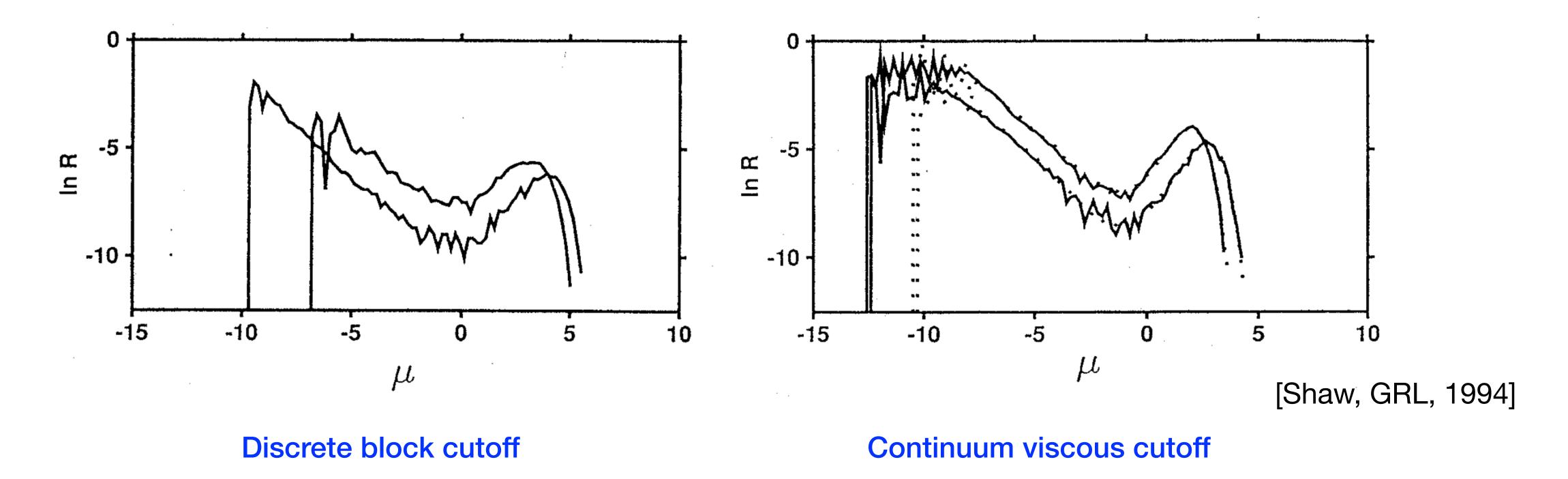
Self-Organized Slip Complexity



Large event complexity
Generic

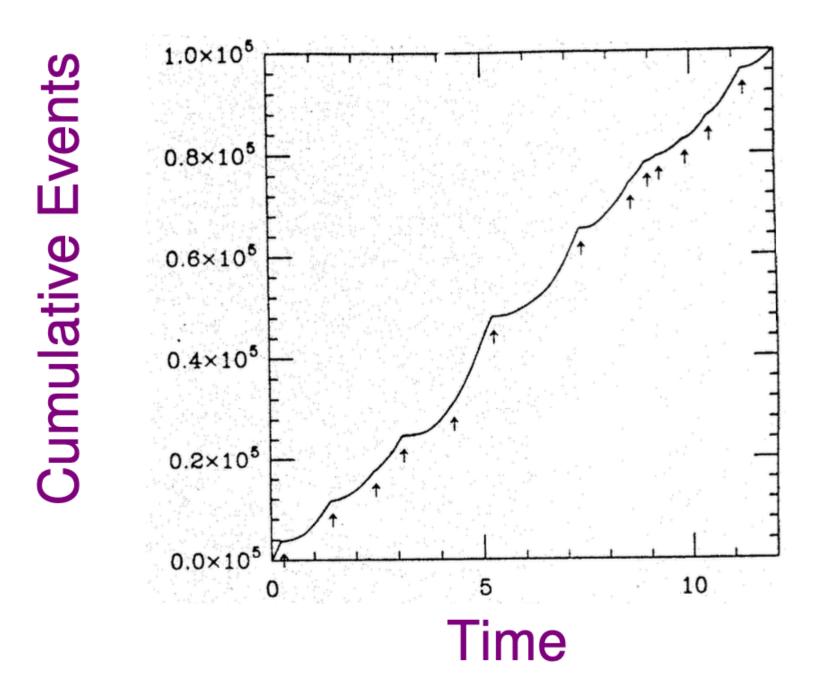
Small + large events near Critical value

Complexity and the continuum



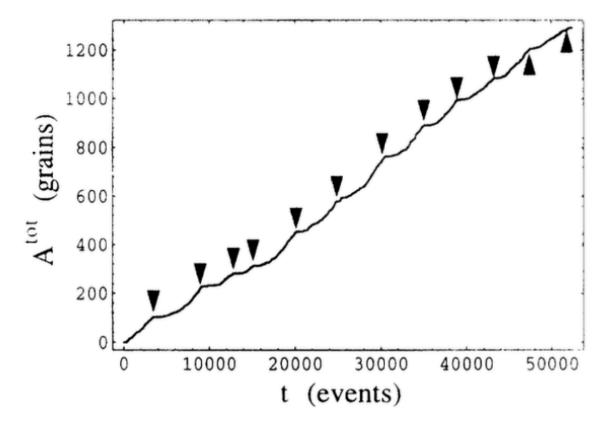
- Distribution of magnitudes depends on small cutoff scale.
- Scale of small cutoff matters at small and large scales.
- But whether discrete or continuum process at cutoff get similar behavior.

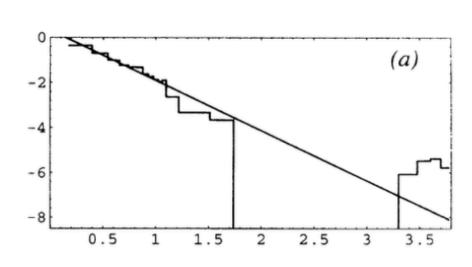
Buildup of small events before the big one



[Shaw et al, 1992]

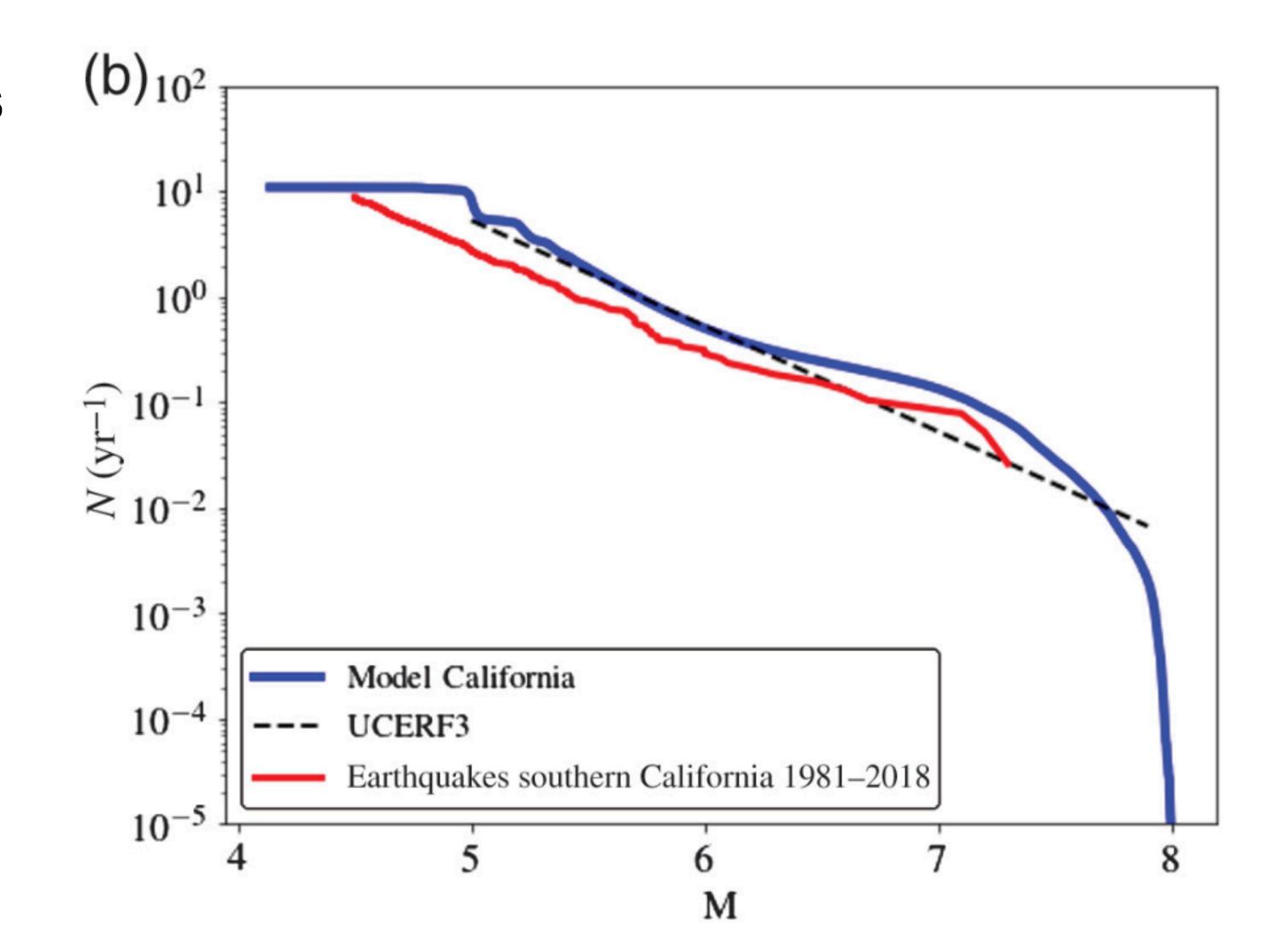
Avalanches in Sand:





[Rosendahl et al, 1994]

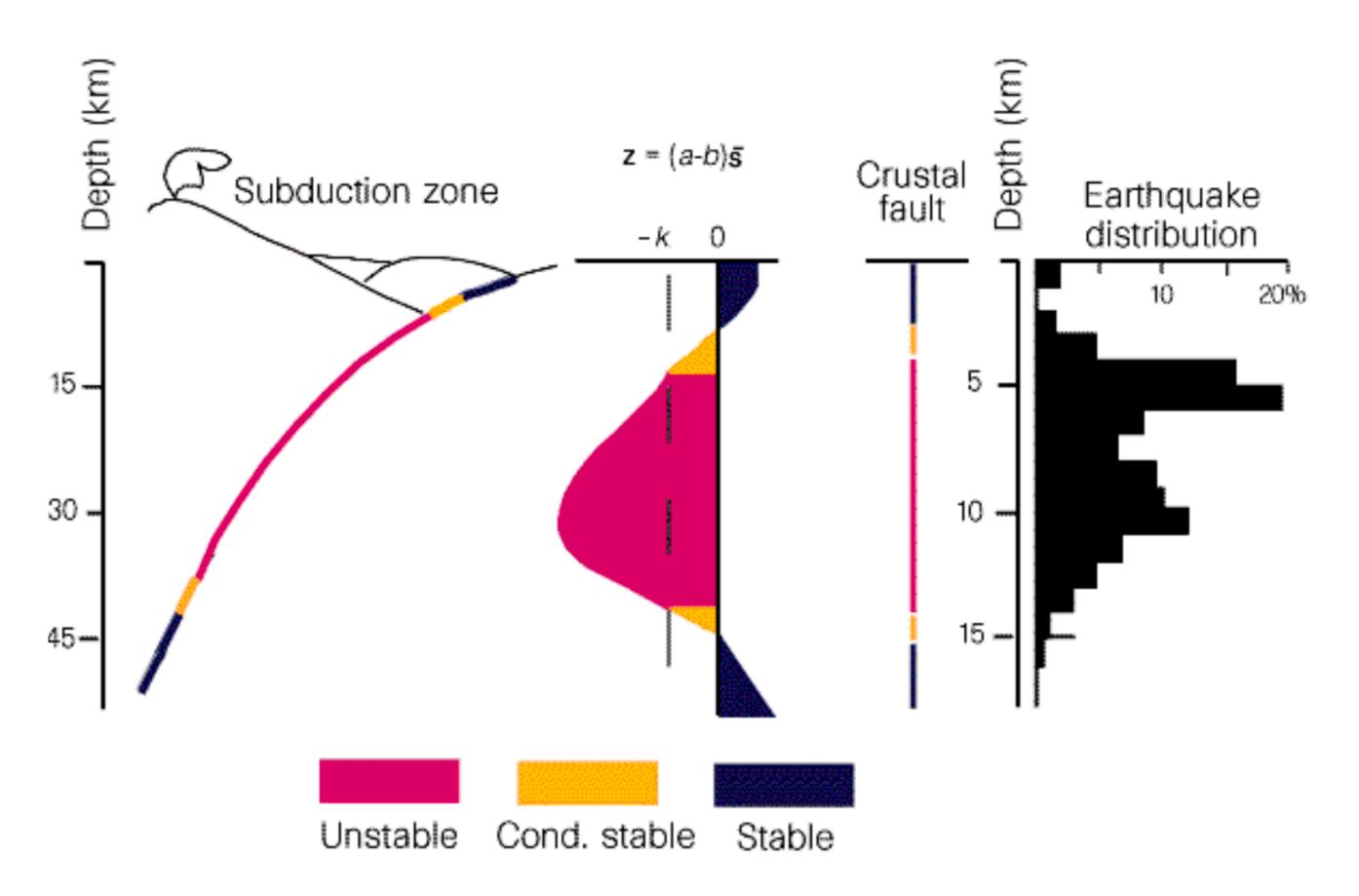
Distribution of sizes



- Excess of large events above extrapolated small event rate
- See this in 1D, 2D, 3D

How Deep Below the Seismogenic Zone do Large Earthquakes Rupture?

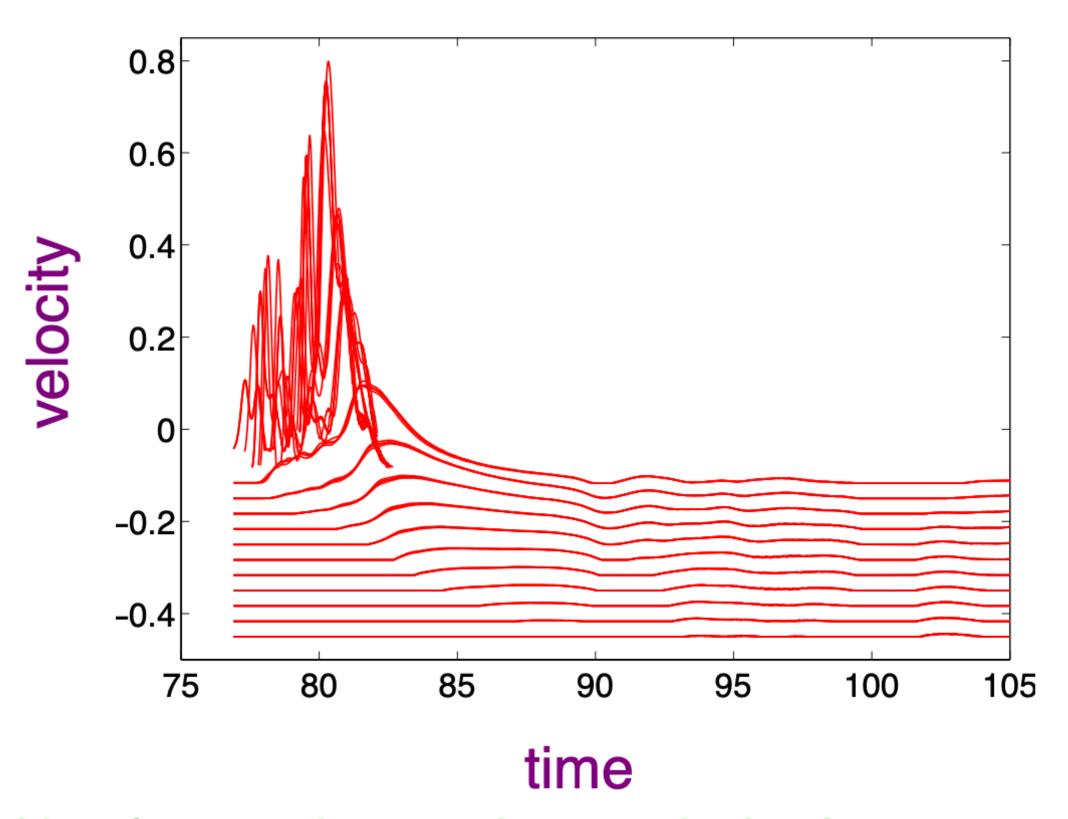
Faults and Stability of Sliding



3D Model

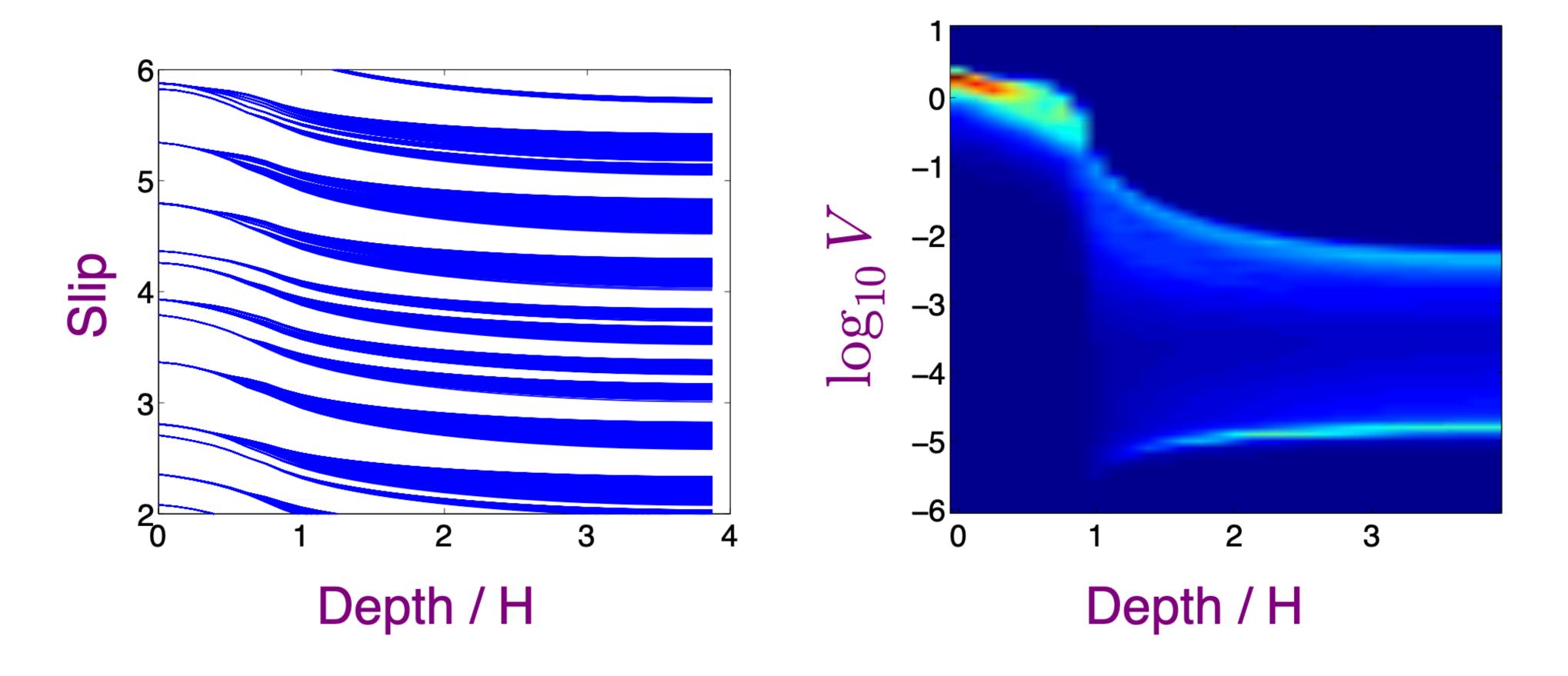
surface slip surface slip 0 65 70 90 95 80 75 85 distance along fault / W vertical exaggeration 8:1 unstable sliding slip stable sliding peak velocity initial displacement

Velocities During Slip Events



- High and low frequencies at seismogenic depths
- Only low frequencies below seismogenic depths
- Coseismic slip at depth

Depth Dependent Slip Rate Distributions

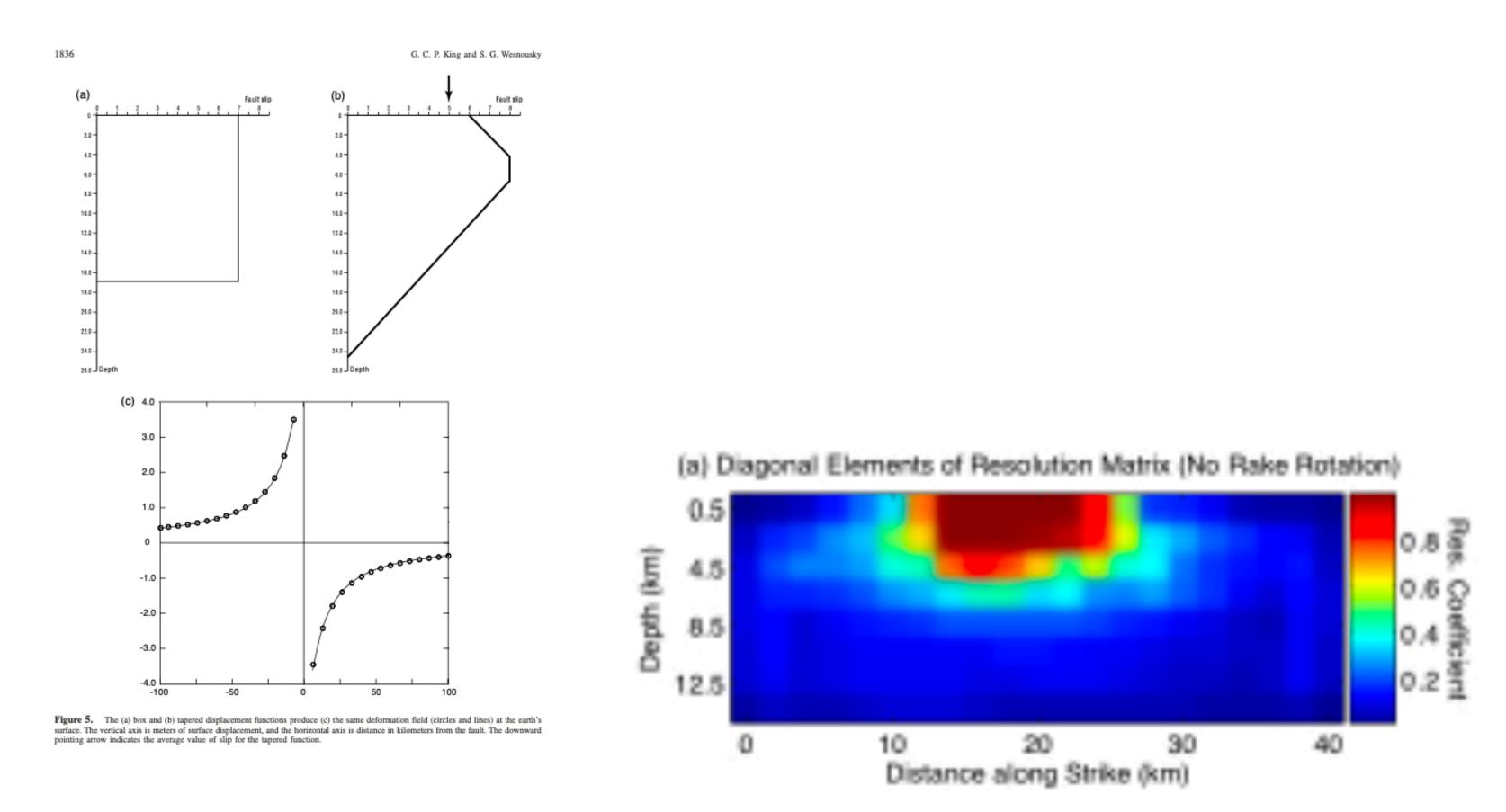


Cumulative Slip

Distribution of Slip Rates

Increasing creep fraction with depth; but coseismic still nonzero

GPS Inversion Ambiguity



[King and Wesnousky, 2007]

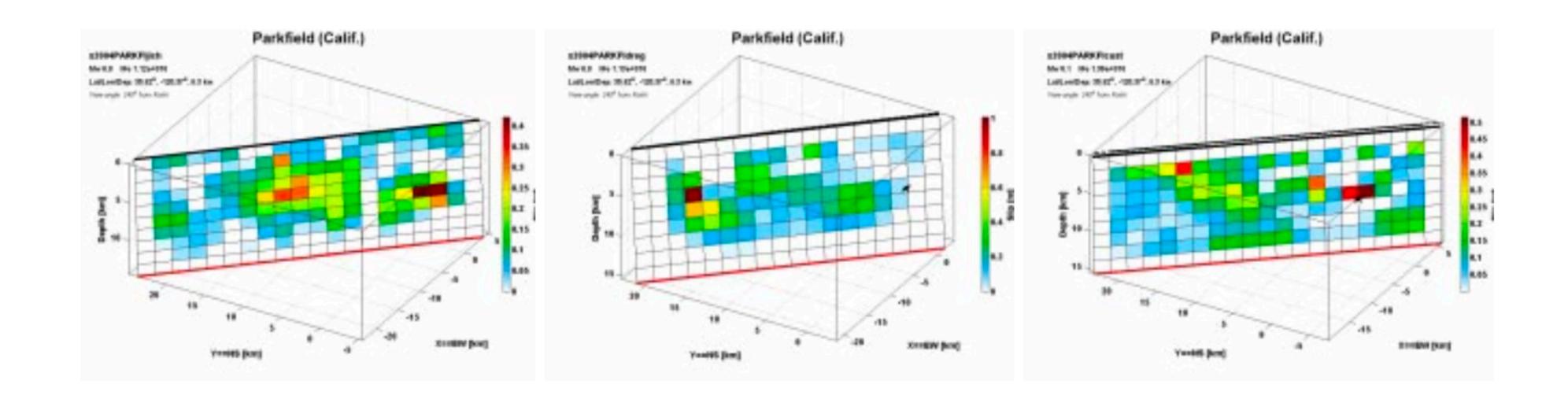
[Page et al, 2008]

Ambiguous depth resolution

Poor resolution for Parkfield

Seismic Inversion Ambiguity

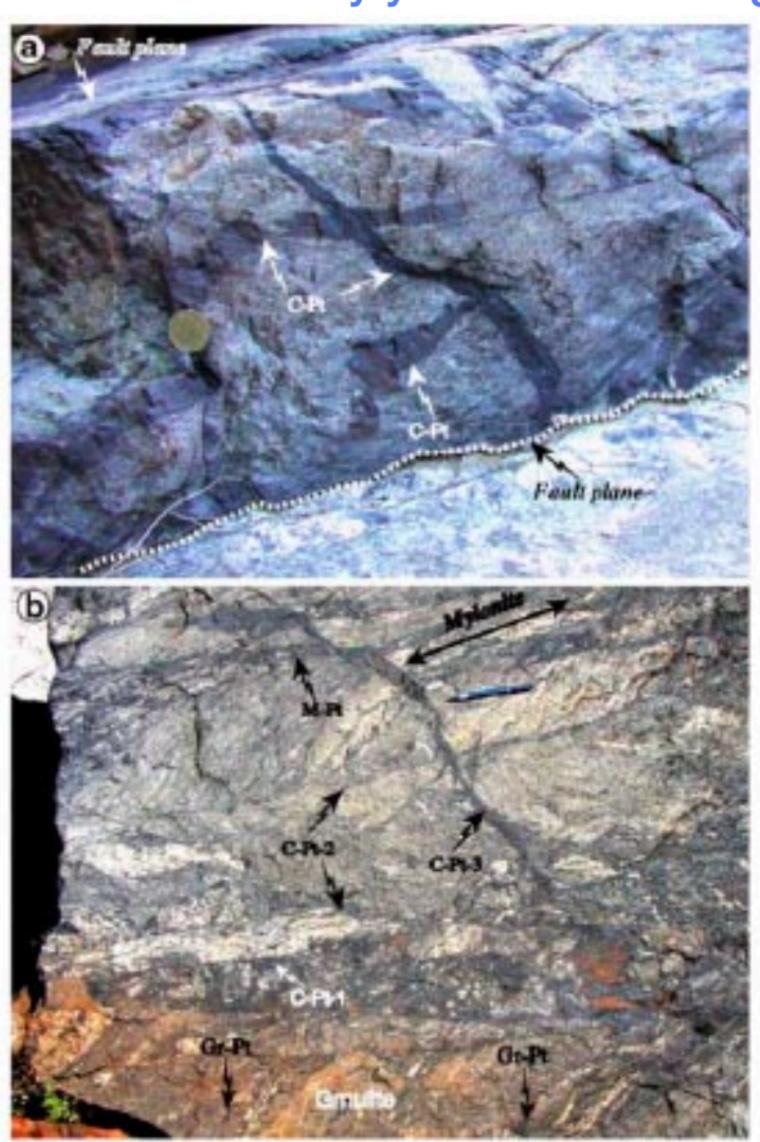
Inversions for slip in 2004 M6.0 Parkfield earthquake (images from [Mai, 2007] database)



• Big differences even in extremely well networked event

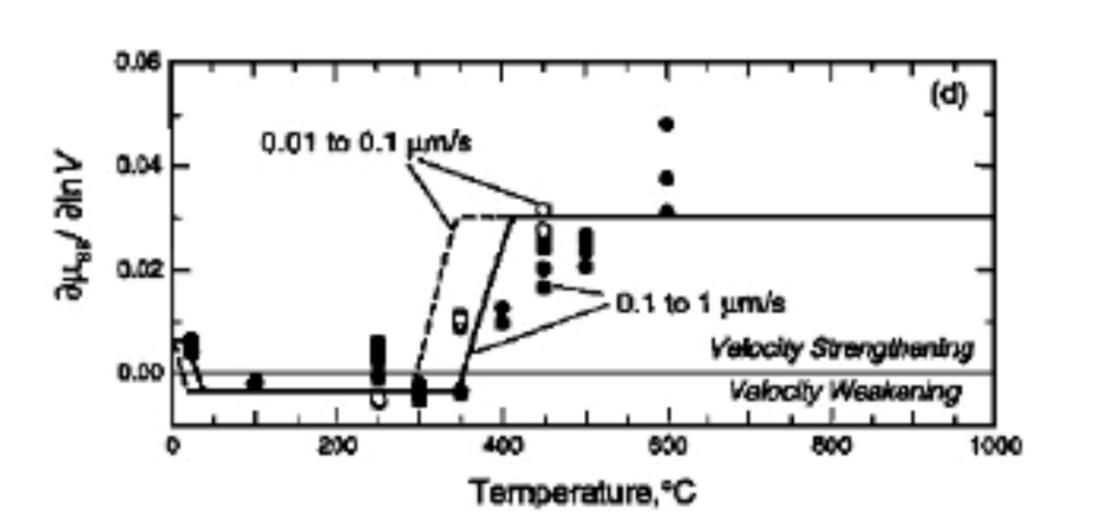
Exhumed Rocks: Fast Deep Slip

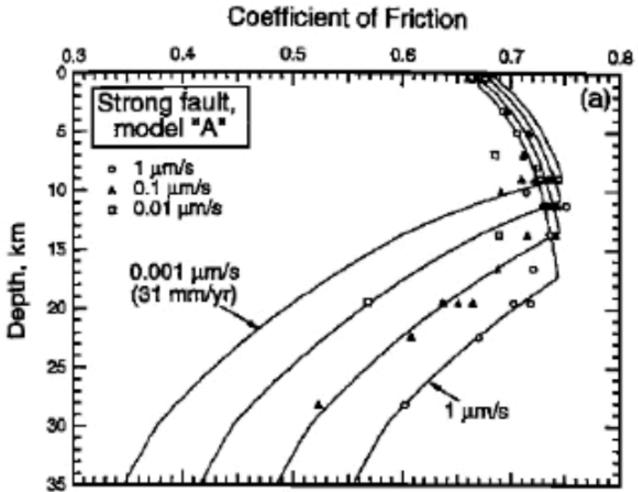
from Pseudotachylytes Crosscutting Mylonites



[Lin, 2008]

Sliding Rate Effect on Lab Friction

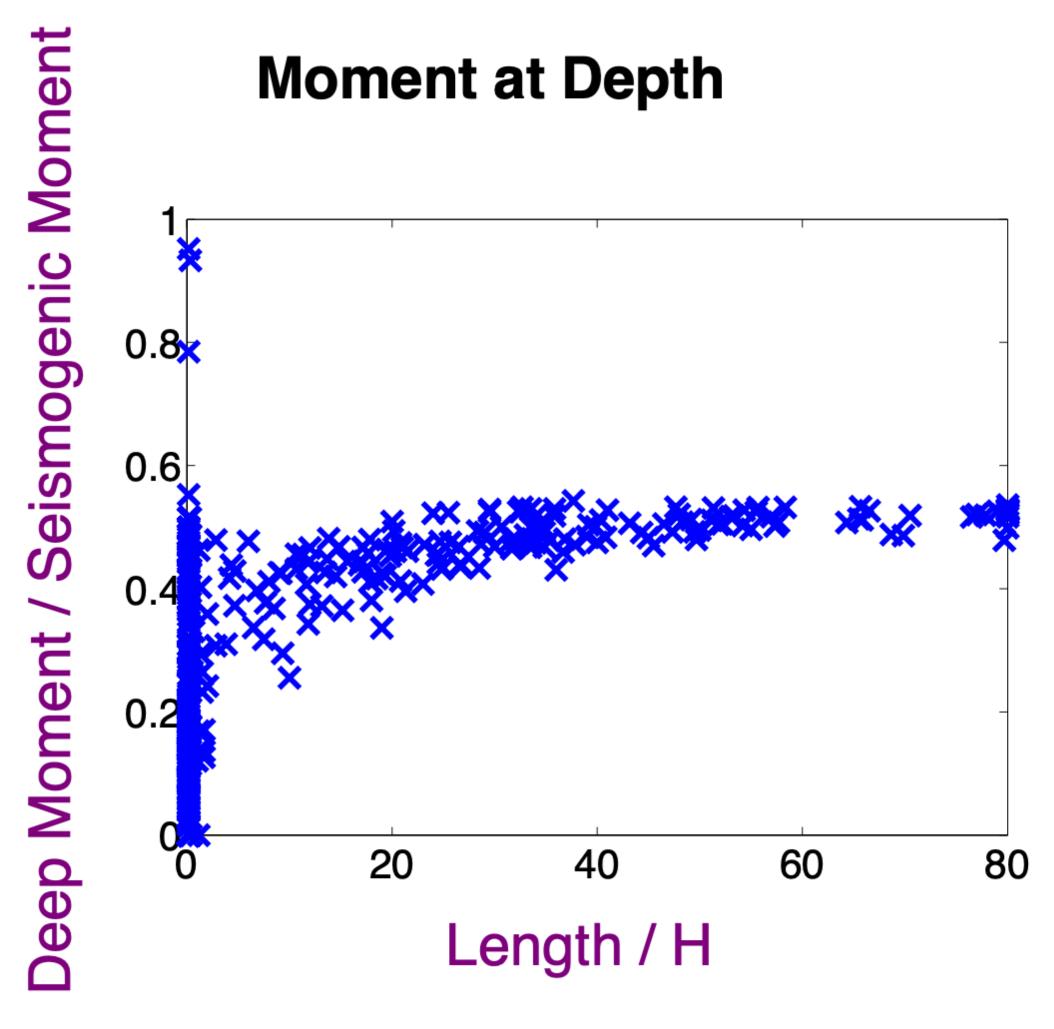




[Blanpied et al, 1995]

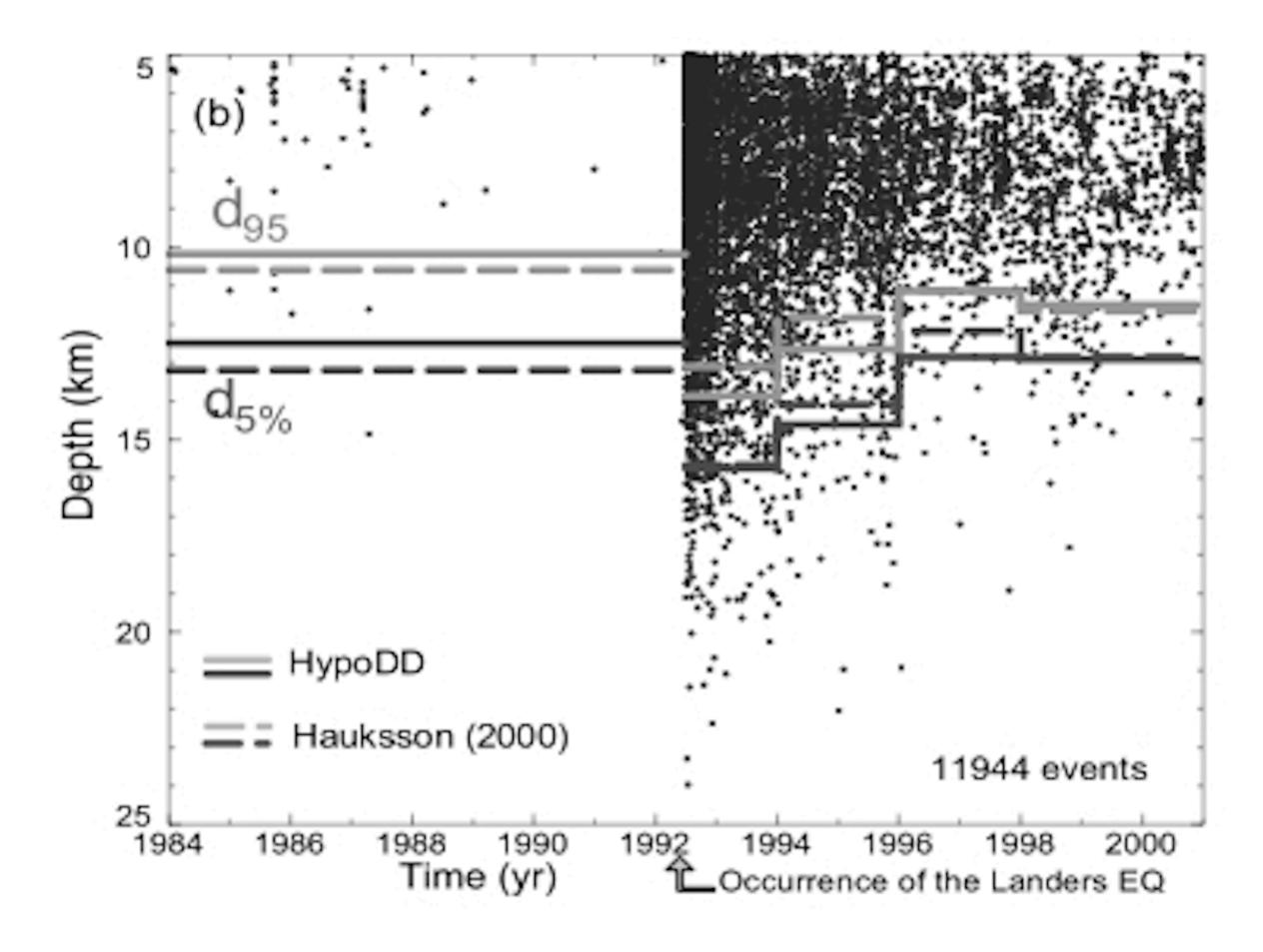
- Unstable sliding deepens at faster sliding rates
- Potentially big effect

3D Model shows significant fraction of slip below seismogenic layer



ullet Half of seismogenic moment, or 1/3 total moment at depth

Deepening of aftershocks following large events



[Rolandone, et al., 2004]

Has implications for scaling relations we will see later