

Unramified Gromov–Witten and Gopakumar–Vafa invariants

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Plan

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- Wall-crossing in Gromov–Witten theory

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- Bigger picture

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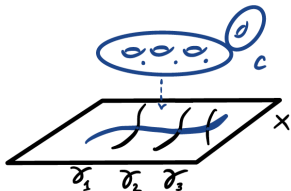
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Algebraic geometers want to interpret $\langle \gamma_1 \cdots \gamma_n \rangle_{g, \beta}^{GW}$ as **counts** of curves of genus g in a class β passing through γ_i .



Problems

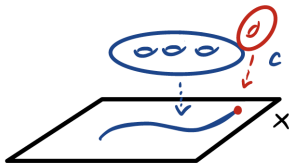
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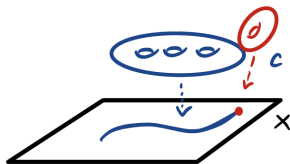
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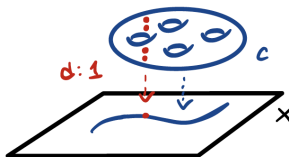
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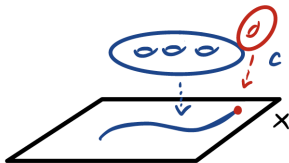
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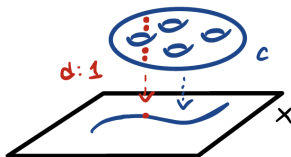
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This leads to **overcounts** in GW invariants, i.e. lower-genus counts contribute to higher-genus ones. Moreover, they are **rational** numbers.

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In this case, there are no **multiple covers**. When $\beta \neq d\beta'$ and $c_1(X) \cdot \beta = 0$, these formulas coincide.

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All together this solves all (geometric) problems of GW theory. However, a **direct** construction of GV invariants is very much desired.

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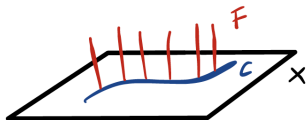
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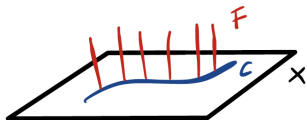
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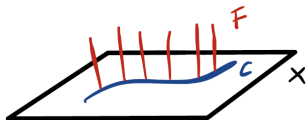


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However, it is not clear how to prove that Maulik–Toda invariants are equal to GV invariants in general. Can be verified in some instances by explicitly computing the invariants. **Very difficult** to compute.

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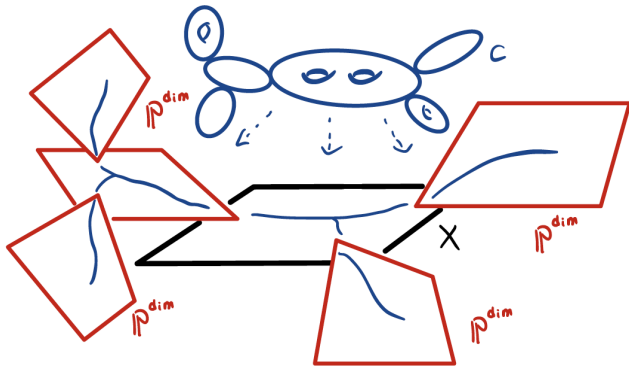
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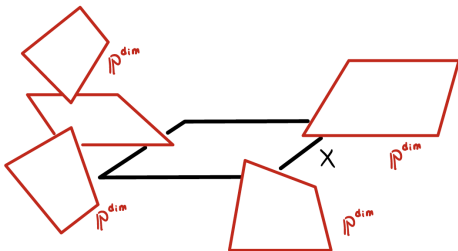
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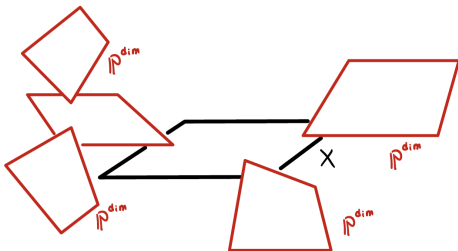
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Like in Gromov–Witten theory, we can define spaces of unramified maps to FM degenerations of X ,

$$u\overline{M}_{g,n}(X, \beta) = \{f: (C, p_1, \dots, p_n) \rightarrow W \mid f \text{ is unramified}\} / \sim .$$

Spaces $u\overline{M}_{g,n}(X, \beta)$ are very similar to $\overline{M}_{g,n}(X, \beta)$. If $\dim(X) = 1$, these are spaces of covers of a curve, known as **Hurwitz** spaces.

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- $\deg(f^*H) = 2g - 2$. The **genus** parameter becomes a **degree** parameter!

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Theorem (N.'24)

The Conjecture is true. Moreover, the formula above is a **wall-crossing formula**.

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Given a map $f: (C, p_1, \dots, p_n) \rightarrow W$, to each point $x \in W$, and to each bubble $\mathbb{P}^{\dim} \subset W$, we can associate **weights**,

$$w(x) \in \mathbb{Z}, \quad w(\mathbb{P}^{\dim}) \in \mathbb{Z}.$$

The weight $w(x)$ measures how far the map is from being **unramified** over x , while $w(\mathbb{P}^{\dim})$ measures how **big** the curve is over \mathbb{P}^{\dim} .

For example, a contracted component of genus g over x contributes $2g - 2$ to $w(x)$, so does a curve of genus g over \mathbb{P}^{\dim} for $w(\mathbb{P}^{\dim})$.

Given $\epsilon \in \mathbb{R}_{>0}$. The map f is **ϵ -unramified**, if

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Spaces $\overline{M}_{g,n}^\epsilon(X, \beta)$ are the same for each ϵ in a chamber. As we cross a wall between chambers, the space $\overline{M}_{g,n}^\epsilon(X, \beta)$ changes abruptly.

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Each value of ϵ provides a stability condition, such that $\mathbb{R}_{>0}$ is partitioned in **chambers** and **walls**.

Spaces $\overline{M}_{g,n}^\epsilon(X, \beta)$ are the same for each ϵ in a chamber. As we cross a wall between chambers, the space $\overline{M}_{g,n}^\epsilon(X, \beta)$ changes abruptly.

If $\epsilon \ll 1$, then

$$\overline{M}_{g,n}^\epsilon(X, \beta) = \overline{M}_{g,n}(X, \beta).$$

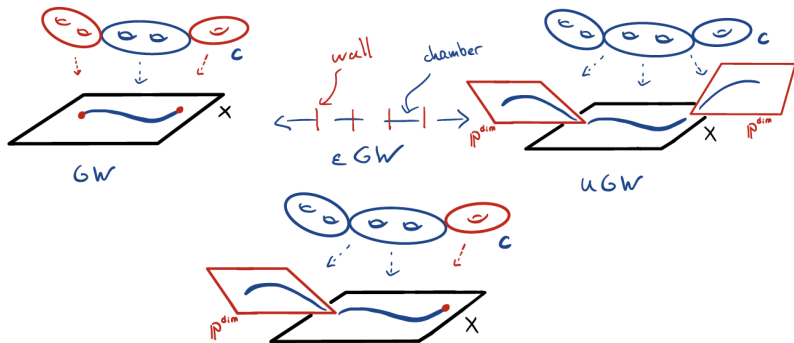
If $\epsilon > 1$, then

$$\overline{M}_{g,n}^\epsilon(X, \beta) = u\overline{M}_{g,n}(X, \beta).$$

ϵ -unramified Gromov–Witten theory

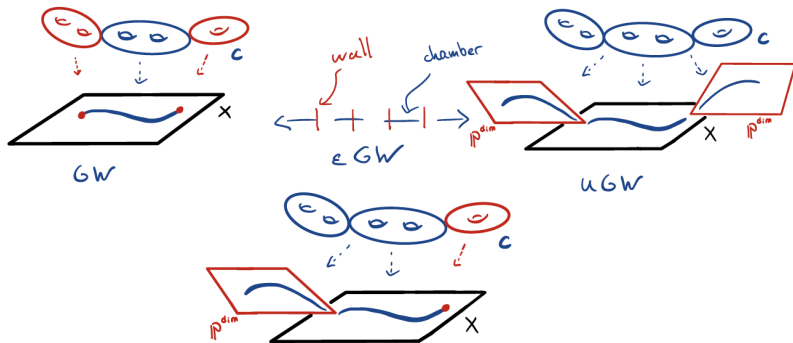
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Overall, we get the following picture:



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It allows us to compare theories across a single **wall** instead of comparing theories for extremal values of ϵ .

Variation of $\epsilon \implies$ **Wall-crossing formulas.**

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$$\text{Hodge integrals} := \int_{\overline{M}_{g,m}} \prod_i \psi_i^{k_i} \cdot \prod_j \lambda_j^{k_j}.$$

$$\psi_i = c_1(L_i), \quad L_i|_C = T_{p_i}^* C,$$

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The contribution from Hodge integrals is **universal**, i.e. depends very little on X .

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Theorem (N.'24)

In all dimensions, we have

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- invariants on the right from the bar are **relative** invariants, i.e. invariants associated to markings on the **target**;

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In dimension three, we need to compute **triple** Hodge integrals,

$$\text{Hodge integrals} = \int_{\overline{M}_{g,1}} \frac{c(\mathbb{E}^\vee z \otimes TX)}{(z - H - \psi_1)} \in H^*(\mathbb{P}(TX))[z^\pm],$$

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Plugging Hodge integrals into the wall-crossing formula, we obtain

$$\sum_g \langle \gamma_1 \cdots \gamma_n \rangle_{g,\beta}^{GW} \lambda^{2g-2} = \sum_g \langle \gamma_1 \cdots \gamma_n \rangle_{g,\beta}^{uGW} \lambda^{2g-2} \left(\frac{\sin(\lambda/2)}{\lambda/2} \right)^{2g-2+c_1(X) \cdot \beta}.$$

the exponent “ $2g - 2 + c_1(X) \cdot \beta$ ” arises due to **dilaton** and **divisor** equations.

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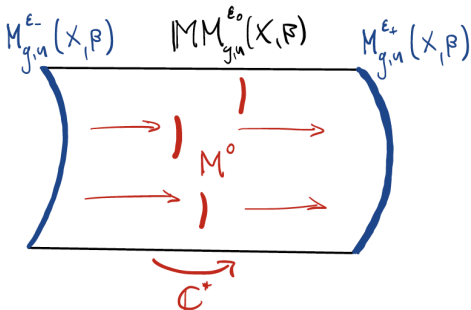
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We construct a **master** space $\mathbb{M}\overline{M}_{g,n}^{\epsilon_0}(X, \beta)$ with a \mathbb{C}^* -action, such that the \mathbb{C}^* -fixed locus contains both $\overline{M}_{g,n}^{\epsilon_+}(X, \beta)$ and $\overline{M}_{g,n}^{\epsilon_-}(X, \beta)$, and **something else** M^0 .

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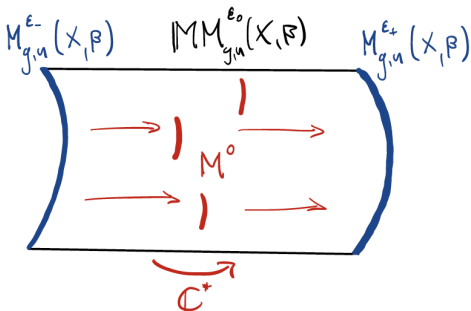
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We then apply the equivariant localisation formula and take the equivariant residue, M^0 are responsible for the **wall-crossing** terms.

Wall-crossing formulas in dimension one

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In dimension one,

Unramified Gromov–Witten theory=Hurwitz theory.

Hurwitz theory counts **covers** of curves with prescribed **ramifications**.
For a collection of ramification profiles $\{\eta^1, \dots, \eta^n\}$, $H_{X,d}(\eta^1, \dots, \eta^n)$ is the associated degree- d Hurwitz number of X .

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The wall-crossing recovers it in a slightly different but equivalent form, expressing $\overline{(k + 1)}$ in terms of Hodge integrals, N.–Schimpf'24.

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A similar phenomenon holds for Gromov–Witten theories of $S \times C$ and of an orbifold symmetric product $Sym_d(S)$, N'22.

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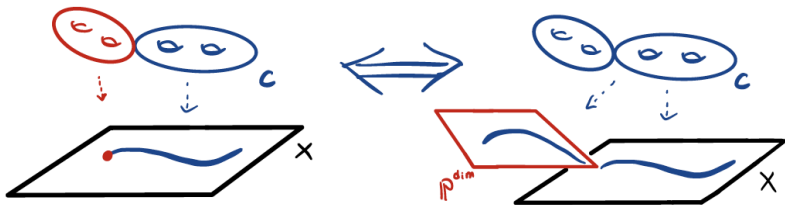
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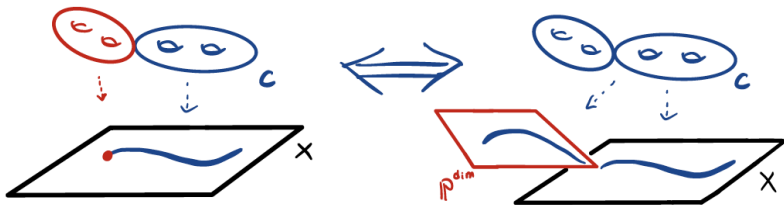


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Instead of a map from a curve, there could be a **sheaf**.

Thank you