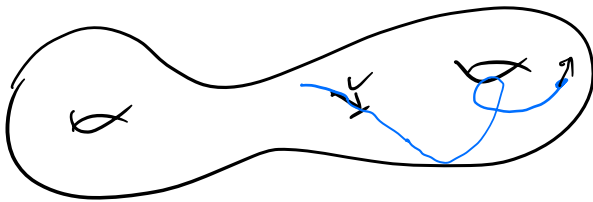


Geodesic flow

Given M manifold, $F = (g_t): SM \rightarrow SM$
 $v \rightarrow t_t v$
 is given by moving at unit speed
 along geodesic determined by v .

A geodesic is a curve which (locally)
 minimizes distance



$v \in SM \iff$ geodesic $c: \mathbb{R} \rightarrow M$ with

$$c_v(0) = v$$

can thus identify SM with
 space of geodesics

$$GM = \{ c: \mathbb{R} \rightarrow M : c \text{ is local isometry} \}$$

(a- also think of $(g_t): GM \rightarrow GM$
 $g_t c$ given by $(g_t c)(s) = c(s+t)$)

d is Riem. distance on M

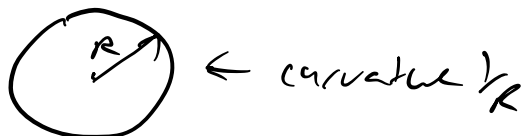
Metric on SM :

$$d(v, w) = \max_{\gamma: S^1 \rightarrow M} \{ d(\gamma_s v, \gamma_s w) \}$$

Curvature Consider a surface

Gauss curvature defined using curvature of curves

curvature



\uparrow
 2^{nd} derivative

On surface at x : take smallest and largest

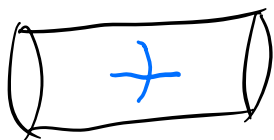
signed curvatures through x , K_1 and K_2

Define $K(x) = K_1 K_2$

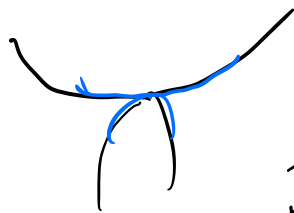
embedded in \mathbb{R}^3
 intersect with plane
 containing normal
 direction. get a curve.
 compute curvature



Sphere $K(x) > 0$



Cylinder
 $K(x) \geq 0$



Saddle/
 mountain pass $K(x) < 0$

Surfaces: completely classified by genus g
 "number of holes"

Gauss-Bonnet : $\int K(x) dV(x) = -4\pi(g-1)$

\therefore Only genus $g \geq 2$ can have $K(x) < 0$
 $\forall x$

Constant -ve curvature Model is \mathbb{H}^2

$$\{x + iy, y > 0, x \in \mathbb{R}\}$$

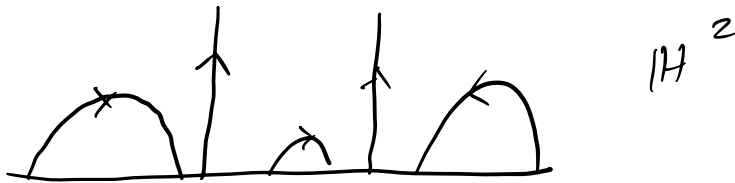
with metric $(ds)^2 = \frac{1}{y^2} (dx^2 + dy^2)$

Given $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{R})$, $z \rightarrow \frac{az + b}{cz + d}$ is an isometry

This is all orientation-preserving isometries
 Since $\begin{pmatrix} a & b \\ c & d \end{pmatrix}, \begin{pmatrix} -a & -b \\ -c & -d \end{pmatrix}$ define same isometry
 isometry group is $PSL(2, \mathbb{R}) = SL(2, \mathbb{R}) / \pm Id$

Knowing the isometries allows us to find the geodesics

- circles centred on real axis
- vertical lines



Projecting to disc model:



Boundary at ∞ - $\partial^\infty \mathbb{H}^2$ is equivalence classes of geodesic rays which stay a bounded distance from each other. Topologically, the boundary is S^1 .

Space of geodesics on \mathbb{H}^2 can be identified with

$$\partial^\infty \mathbb{H}^2 \times \partial^\infty \mathbb{H}^2 \times \mathbb{R}$$

↑
where you came from

$$c \in \mathcal{GH}^2$$

↑
where you're going

↑
where $t=0$ is

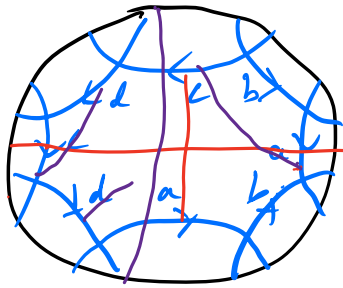


Notation:

$c(x, \beta)$ is unique geodesic with $c(0) = x$
 $c(\infty) = \beta$

$c(m, x)$ is unique geodesic with $c(\infty) = m, c(0) = x$

To get compact constant -ve curvature manifold
we quotient by discrete group of isometries



← gives genus 2

4g-gon gives
genus g

Conversely, every compact constant negative
curvature surface M is \mathbb{H}^2 / Γ

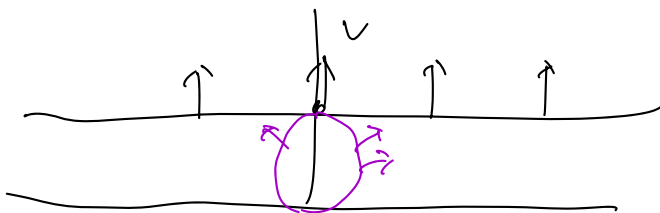
Γ Fuchsian group (cocompact)

stable set on $S\mathbb{H}^2$: For $v \in S\mathbb{H}^2$, consider

$W^{ss}(v) = \{ w \in S\mathbb{H}^2 \text{ with } d(t_t v, t_t w) \rightarrow 0 \text{ as } t \rightarrow \infty \}$
(i.e. the stable manifold)

For $c \in G\mathbb{H}^2$, consider $W^{ss}(c) = \{ c' \in G\mathbb{H}^2 \text{ with } d((t_t)_c, (t_t)_{c'}) \rightarrow 0 \}$
 $W^{ss}(v)$ and $W^{ss}(c)$ are identified

Compute explicitly for vertical vector at $(0, 1)$:



$W^{ss}(v)$ is ^{normal} vector field
over horizontal line

$W^{uu}(v)$ " " over
circle

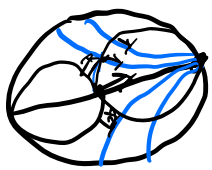
Apply isometries to get all other

$W^{ss}(v)$ and $W^{uu}(v)$

vector fields over

Get horizontal lines and circles tangent
to x -axis

In disc model, we get vector field over sphere
 tangent to ∂H^2



(We can define $W^s(v)$ and $W^u(v)$ this way)

Define $E_v^s = T_v W^s(v)$

$E_v^u = T_v W^u(v)$

$E_v^p = \text{flow direction}$

$E^s \oplus E^u \oplus E^p$ will be Anosov splitting

Hence, Anosov flow.

(Need to know exponential contraction / expansion... soon...)

Haar measure on $PSL(2, \mathbb{R})$ descends to natural volume measure

- Liouville measure on SM

(f_t) is ergodic (Artin, Hopf)

Variable negative curvature

(M, g) M surface genus ≥ 2

is topologically \mathbb{H}^2 / Γ

\tilde{M} is universal cover: \mathbb{H}^2 equipped with

\tilde{g} is still g metric \tilde{g} which projects to g

Hamiltonian flow, so natural volume on SM preserved by $F = \{f_t\}$

- called Liouville measure

- locally it is $\text{Vol} \times \text{leb}_g$

If (M, g) has -ve curvature, it is Anosov

\rightarrow 2 approaches

- geometric argument like before \nearrow use Busemann functions to describe w^s

or - Study Jacobi equation -

Get Anosov property by verifying cone condition -

Anosov: ('60s)

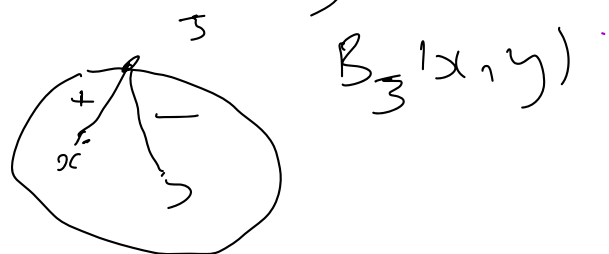
Volume-preserving Anosov diffeos/flows are ergodic (w.r.t. Liouville measure)

- Hopf argument

- Absolute continuity of foliations is the breakthrough

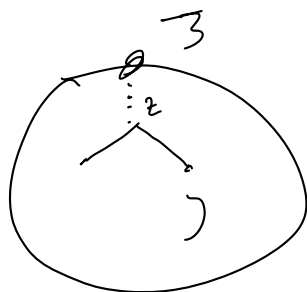
Buseman Functions

- measures relative distance to a point at infinity
i.e. how "out of phase" you are with a reference point traveling to infinity



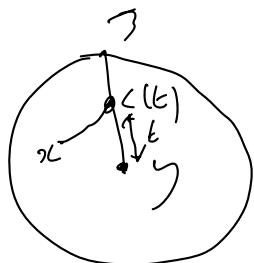
Define

$$B_3(x, y) = \lim_{z \rightarrow \infty} (d(x, z) - d(y, z))$$



Any $z \rightarrow \infty$

More often we think of y fixed and let $z \rightarrow \infty$ along $c(y, \infty)$



so

$$B_3(x, y) = \lim_{t \rightarrow \infty} (d(x, c(t)) - t)$$

where $c(t) = c(y, \infty)(t)$

(Let's think of y as a fixed origin " ∞ ".)

We often write $B_{\underline{0}, \underline{3}}(x) = B_{\underline{3}}(x, \underline{0})$.

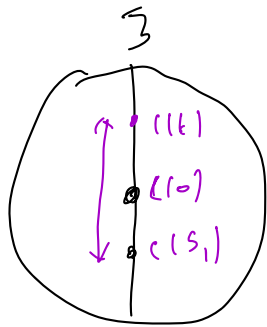
$$B_{\underline{0}, \underline{3}}(\cdot) : \tilde{X} \rightarrow \mathbb{R}$$

$$x \mapsto \lim_{t \rightarrow \infty} (d(x, c(t)) - t)$$

where $c = c(\underline{0}, \underline{3})$

Properties:

1. For $c = c(\underline{0}, \underline{3})$, $B_{\underline{0}, \underline{3}}(c(s)) = -s$



$$s_1 < 0 : d(c(s_1), c(t)) - t = t + |s_1| - t = -s_1$$

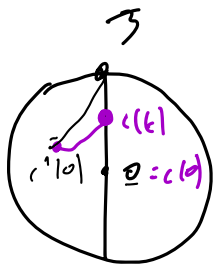
$s_2 > 0$, for t large



$$d(\underline{0}, c(t)) - t = (t - s_2) - t = -s_2$$

" $B_{\underline{0}, \underline{3}}(x) \begin{cases} > 0 \\ < 0 \end{cases}$ when x is "behind" $\underline{0}$
 "ahead" $\underline{0}$ "

2. Let $c = c(0, 3)$ and let $c' \in C^m$ with $d(c(t), c'(t)) \rightarrow 0$
 as $t \rightarrow \infty$. Then $c'(0) = 3$, by definition.
 and $B_{0,3}(c'(0)) = 0$



check: $B_{0,3}(c'(0)) = \lim (d(c'(0), c(t)) - t)$

$\leq \lim (d(c'(0), c'(t)) + d(c'(t), c(t)) - t)$

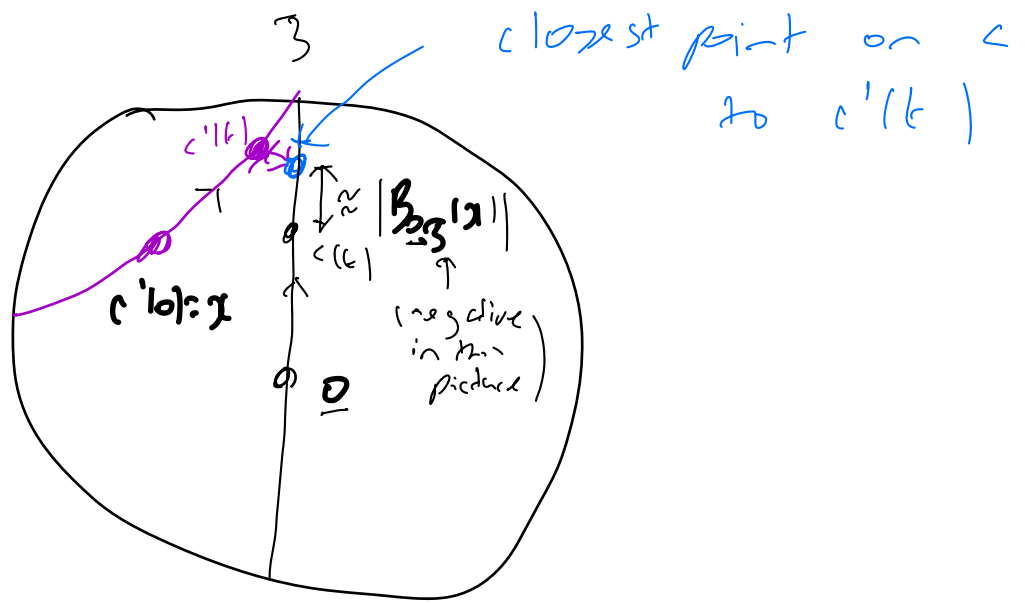
$= \lim d(c'(t), c(t)) = 0$

Thus $B_{0,3}(c'(0)) = 0$

3. For $x \in X$, $B_{0,3}(x)$ can be thought of as taking the geodesic $c' = c'(x, 3)$

and flowing towards 3 . Eventually the closest distance between c and c' becomes negligible. $B_{0,3}(x)$ is what is left over. It's like a signed distance in flow direction

i.e. $B_{0,3}(x)$ is (asymptotically) how much time lag $c'(x, 3)$ has compared with $c(0, 3)$



- $B_{p,z}(\cdot)$ is Lipschitz and convex (see wiki - Busemann functions)

Busemann functions let us define $w^s(c)$

For a point p and $z \in \partial X$

define $H_{p,z} = \{x : B_{p,z}(x) = 0\}$.

$H_{p,z}$ is the horosphere for z

at p . Alternate construct - take ball (radius) k around $c(t)$. Take limit



"horosphere" means "sphere touching infinity". This construction shows this



The (strong) stable set of geodesics for

$$C = C(p, 3) \quad \text{is}$$

$$W^{ss}(C) := \left\{ c' : c'(\infty) = 3, c'(0) \in H_{p,3} \right\}$$

In neg. curv. this agrees with dynamical

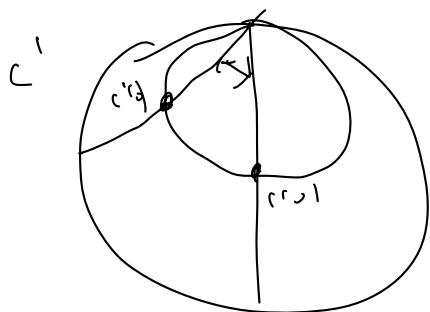
$$\text{defn. } \left\{ c' : d(c(t), c'(t)) \rightarrow 0 \right\}$$

This is because:

"Really requires neg. curv. not true in non-pos. curv."

Can have flat strip with $B(\cdot) = 0$

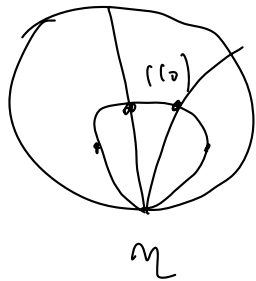
Geometry fact: In neg. curv., two asymptotic rays converge exponentially in nearest distance



Thus for $c' \in W^{ss}(C)$, dist. between geodesics $\rightarrow 0$ and "time lag" is 0

The unstable set of geodesics for $c(m, 0)$ is

$$W^{uu}(c) = \left\{ c' : c'(t \rightarrow \infty) = \eta, c(0) \in H_{p, \eta} \right\}$$



For $v \in S^1 \mathbb{H}^2$, we thus constructed $W^{uu}(v)$ and $W^{ss}(v)$ on which distance is expanded (resp. contracted) exponentially

This construction upstairs projects down to SM

$$\text{Given } T^*M = E^s \oplus E^u \oplus E^0$$

Anosov flow

Qn: What did we gain by Busemann function definition of $W^{ss}(v)$?

Ans: regularity

and explicit characterization we can understand by further analysis of Busemann functions

— The definition

$$W^{ss}(x) = \{y : d(f_t x, f_t y) \rightarrow 0\}$$

can always be made in any setting!! But it may be useless. Why is $W^{ss}(x) \neq \emptyset$? Why does it have any structure?

- The Busemann functions have nice regularity and hence so does $H_{p,3}$ and hence \Rightarrow does $W^{ss}(U)$ defined using $H_{p,3}$
- Lipschitz regularity at $B_{3,0}(x)$ is easy in x
- Hölder regularity as 3 varies is also elementary (if a little harder)

[Corresponds to Hölder regularity of $x \rightarrow E_x^s$ in hyperbolic dynamics]

- For a non-positive curvature manifold, the Busemann functions and hence $H_{p,3}$ are regular (geometry gives that each $W^{ss}(U)$)

is at least a C^2 manifold)

↑ Smoothness was needed
to define

$$E_V^S = T W^S(V), \text{ etc.}$$

• Note Another advantage
of the geometry approach
is W^S , W^{uc} and
defined for some
interesting settings beyond
the manifold case

e.g. $CAT(-1)$ spaces

↑ flow is no longer
Anosov in the conventional
smooth sense but there is enough

Hölder regularity to
study it dynamically
from a metric point
of view.