



Investigating the Alternating Structures of TI-NI and TI-S to Realize Different Symmetries of the Superconducting Order Parameter in Weyl Semimetals

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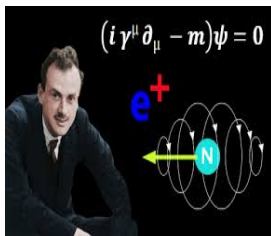
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Introduction

- Dirac equations \rightarrow Dirac fermions
- Weyl equation \rightarrow Weyl fermions
- Weyl semimetals \rightarrow Weyl superconductors
- Method: analytical



Dirac equation

$$\textcircled{1} (i\gamma^\mu \partial_\mu - m)\psi = 0.$$

$$\textcircled{2} \frac{\partial \psi}{\partial t} + \alpha^r \frac{\partial \psi}{\partial x^r} + \beta m \psi = 0.$$

- γ^μ and α^r are 4 component matrices
- satisfies the Clifford algebra.
- $\{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu}$
- predicts $\pm E \rightarrow \mathbf{s}$

Weyl equation

$$\textcircled{1} i\sigma^\mu \partial_\mu \psi_R = 0.$$

$$\textcircled{2} i\bar{\sigma}^\mu \partial_\mu \psi_L = 0.$$

$$\textcircled{3} \frac{\partial \psi}{\partial t} + \sigma^r \frac{\partial \psi}{\partial x^r} = 0.$$

- $\sigma^\mu = 1, \sigma_x, \sigma_y, \sigma_z$
- $\bar{\sigma}^\mu = 1, -\sigma_x, -\sigma_y, -\sigma_z$
- satisfies the Clifford algebra
- $\{\sigma^\mu, \sigma^\nu\} = 2\eta^{\mu\nu}$

Chirality and Helicity

Handedness

Chirality

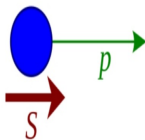
- mathematical concept
- breaks M R symmetry.
- eigenvalue of $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$
- $[\gamma^5, H_W] = 0$



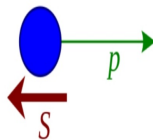
Helicity

- projection of spin in momentum direction
- $h = \mathbf{s} \cdot \hat{\mathbf{p}}$
- $[h, H_D] = 0$
- eigenvalue = ± 1

Right-handed:

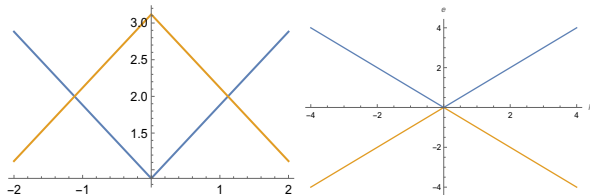


Left-handed:



Weyl Points and Weyl Semimetals

- Energy expansion: $E = \hbar v_F |\vec{k} - \vec{k}_0|$.
- Effective Hamiltonian: $H = \varepsilon_0 \sigma_0 \pm \hbar v_F (\vec{k} - \vec{k}_0)$.
- Weyl Hamiltonian ($\varepsilon_0 = 0$): $H = \pm \hbar v_F (\vec{k} - \vec{k}_0)$.



Important Properties of Weyl Semimetals

Berry curvature and Berry phase

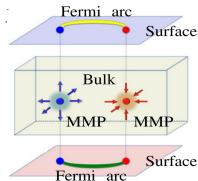
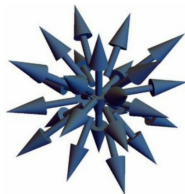
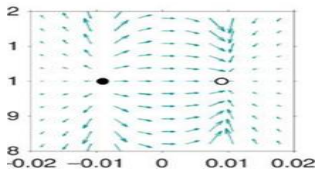
- Berry connection: $\vec{A}(\vec{k}) = -i\sum_{n,occ}\langle u_{n,\vec{k}}|\vec{\nabla}_{\vec{k}}|u_{n,\vec{k}}\rangle$.

- Berry curvature: $B(\vec{k}) = \vec{\nabla}_{\vec{k}} \times \vec{A}(\vec{k})$.

- Berry phase: $\Omega(\vec{k}) = \int d\vec{k}\langle u_{n,\vec{k}}|\vec{\nabla}_{\vec{k}}|u_{n,\vec{k}}\rangle = \pm 2\pi$

$$H = \pm \hbar v_F(\vec{k} - \vec{k}_0)$$

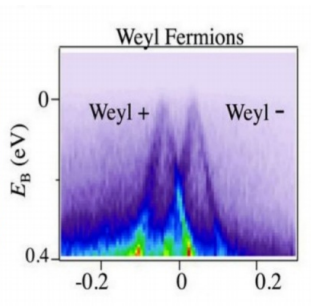
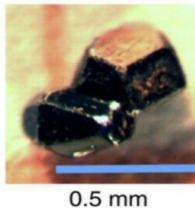
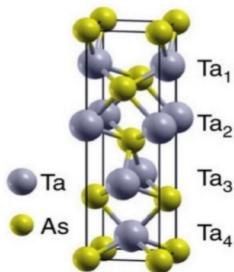
- Fermi arcs and chiral anomaly



Realization of Weyl Semimetals

Experimental achievements

- $TaAs$: Tantalum arsenide
- $A_2Ir_2O_7$: Pyro-chloore iridate
- Na_3Bi : Sodium bismuthide



TI-NI Hetero Structure

$$H = \sum_{\mathbf{k}_\perp, ij} [v_F \tau^z (\hat{\mathbf{z}} \times \boldsymbol{\sigma}) \cdot \mathbf{k}_\perp \delta_{i,j} + m \sigma^z \delta_{i,j} + \Delta_S \tau^x \delta_{i,j} + \frac{1}{2} \Delta_D \tau^+ \delta_{i,j+1} + \frac{1}{2} \Delta_D \tau^- \delta_{i,j-1}] c_{\mathbf{k}_\perp i}^\dagger c_{\mathbf{k}_\perp j}$$

- $\Delta_D < \Delta_S \rightarrow$ NI
- $\Delta_D > \Delta_S \rightarrow$ TI
- $\Delta_S = \pm \Delta_D \rightarrow$ Dirac node

$m=0$:

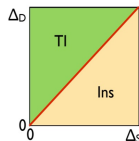
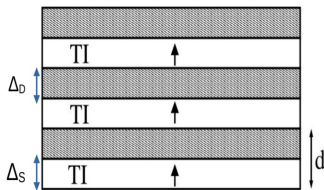
$$E_\pm^2 = v_F^2 (k_x^2 + k_y^2) + \Delta^2(k_z).$$

$$\Delta^2(k_z) = \Delta_S^2 + \Delta_D^2 + 2\Delta_S \Delta_D \cos(k_z d).$$

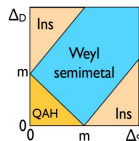
$m \neq 0$:

$$E_\pm^2 = v_F^2 (k_x^2 + k_y^2) + [m \pm \Delta(k_z)]^2.$$

$$(\Delta_S - \Delta_D)^2 < m^2 < (\Delta_S + \Delta_D)^2.$$



(a) $m=0$



(b) $m \neq 0$

TI-S Hetero-Structure

Scalar order parameter

$$H = \sum_{\vec{k}_{\perp,ij}} c_{\vec{k}_{\perp,i}}^{\dagger} H_{ij} c_{\vec{k}_{\perp,j}} + H_{SC}.$$

$$H_{ij} = v_F \tau^z (\hat{z} \times \vec{\sigma}) \cdot \vec{k}_{\perp} \delta_{i,j} + m \sigma^z \delta_{i,j} + t_S \tau^x \delta_{i,j} \\ + \frac{1}{2} t_D \tau^+ \delta_{i,j+1} + \frac{1}{2} \tau^- \delta_{i,j-1}.$$

$$H = \sum_{\vec{k}l=\pm} c_{\vec{k},l}^{\dagger} H_l c_{\vec{k}l} + \sum_{l=\pm} H_{SC,l}.$$

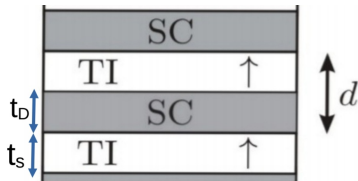
$$H_{\pm} = v_F (\hat{z} \times \vec{\sigma}) \cdot \vec{k} + M_{\pm}(k_z) \sigma^z.$$

$$M_{\pm}(k_z) = m \pm \sqrt{t_S^2 + t_D^2 + 2t_S t_D \cos(k_z d)}.$$

$$(t_S - t_D)^2 < m^2 < (t_S + t_D)^2.$$

- $\sigma^{\pm} = \tau^z \sigma^{\pm}$

- $\tau^{\pm} = \sigma^z \tau^{\pm}$



$$\vec{k} = (0, 0, \frac{\pi}{2})$$

$$\vec{k} = (0, 0, \frac{\pi}{d} \pm k_0)$$

$$\vec{k} = (0, 0, \frac{\pi}{d} \pm k_{\pm}^{\Delta})^T$$

Different Phases of Superconductivity

$$H_- = \frac{1}{2} \sum_{\vec{k}} \psi_{\vec{k}}^\dagger [(v_F(\hat{z} \times \vec{\sigma}) \cdot \vec{k} + M_-(k_z)\sigma^z) \mathbf{I}_{\vec{k}} + \frac{1}{2} \sigma^z (\Delta \kappa^+ + \Delta^* \kappa^-)] \psi_{\vec{k}}$$

$$H_- = \frac{1}{2} \sum_{\vec{k}, n=\pm} \phi_{\vec{k}, n}^\dagger H_-^{n\Delta} \phi_{\vec{k}, n}$$

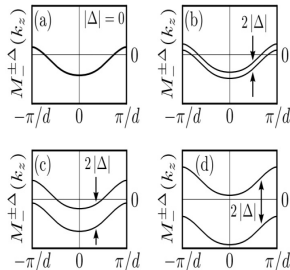
$$H_-^{\pm\Delta} = v_F(\hat{z} \times \vec{\sigma}) \cdot \vec{k} + M_-^{\pm\Delta}(k_z)\sigma^z.$$

$$M_-^{\pm\Delta}(k_z) = (m \pm |\Delta|) - \sqrt{t_S^2 + t_D^2 + 2t_S t_D \cos(k_z d)}.$$

$$m > |\Delta| : (t_S - t_D) < m \pm |\Delta| < (t_S + t_D).$$

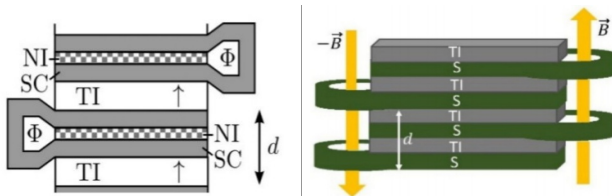
$$m < |\Delta| : (t_S - t_D) < |\Delta| \pm m < (t_S + t_D).$$

- $m > 0$
- $t_S > t_D$



Different Symmetries of Superconductivity

- applying π phase difference \rightarrow pseudo-scalar
- applying ϕ and $-\phi \rightarrow$ two components order parameter
(pseudo scalar and scalar)



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Thank you for your attention.