

Combining transport Monte Carlo with other simulation methods

Nikita Medvedev

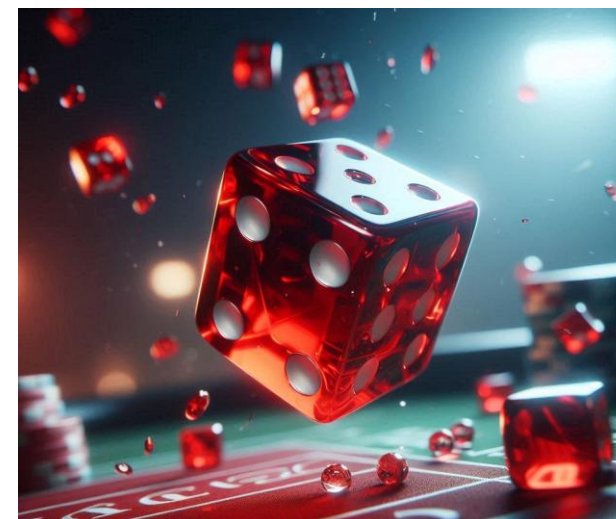
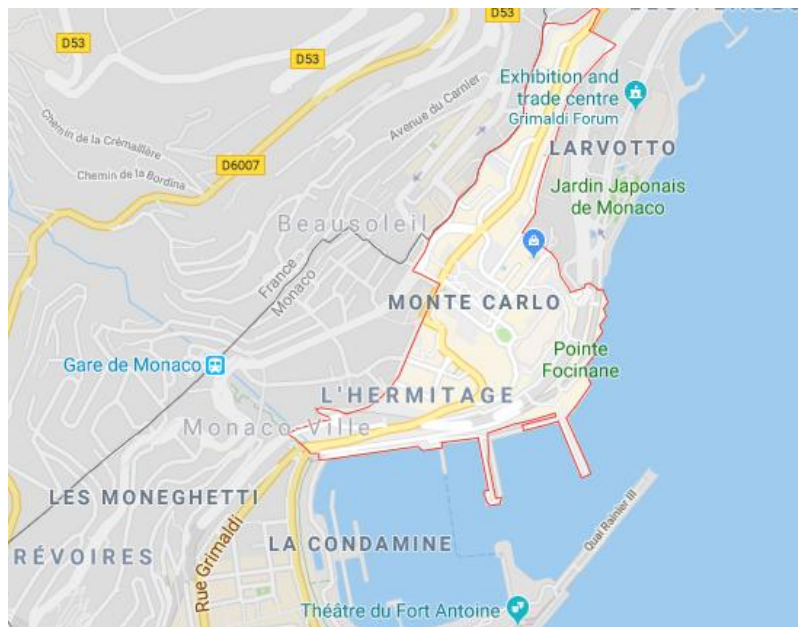


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Outline

1. Introduction: what is MC and applications
2. Example: MC integration
3. Random number generators
4. MC in radiation transport
5. Cross sections
6. Photons, electrons, holes, ions
7. Combining MC with other methods: hybrid, multiscale...
8. Models without feedbacks
9. Models with feedbacks

Introduction: Monte Carlo



N. Metropolis, S. Ulam
"The Monte Carlo Method"
J. Amer. Stat. Assoc. 44, 335 (1949)

Key to Monte Carlo method: randomness

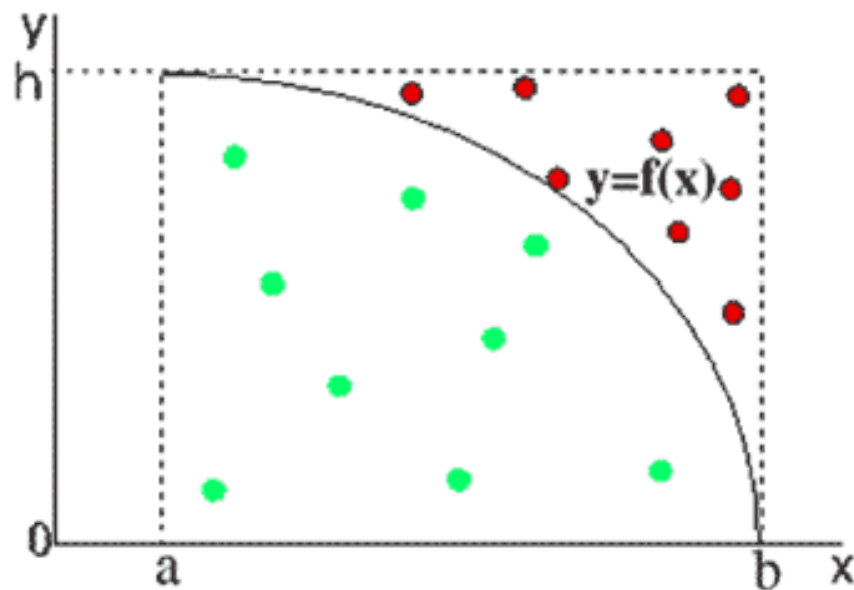
Monte Carlo methods (plural)

There is no “**the**” Monte Carlo (MC) method: huge variety of them

- Metropolis MC (ensemble evolution)
- Biology MC
- Comparison of risk analysis (investment banking)
- Direct simulation MC (gas and fluid flow)
- Dynamic MC (chemistry)
- Kinetic MC (defects in solids)
- Quantum MC
- **MC for particle transport**
- **Event-by-event (analog) MC**
- Condensed collisions MC
- etc.

Example: Monte Carlo integration

Solving integrals with MC method (a.k.a. rejection method)



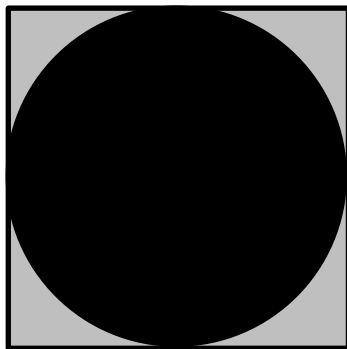
Area under a curve is \sim number of points under it vs. total

$$\int_a^b f(x) dx \approx S_0 \frac{N_{in}}{N}$$

Convergence

Number of iterations defines variance:
$$\sigma_Q = \sqrt{\frac{\text{var}(q)}{N}} = \sqrt{\frac{1}{N} \left[\frac{1}{N} \sum_{i=1}^N q_i^2 - \bar{Q}^2 \right]}.$$

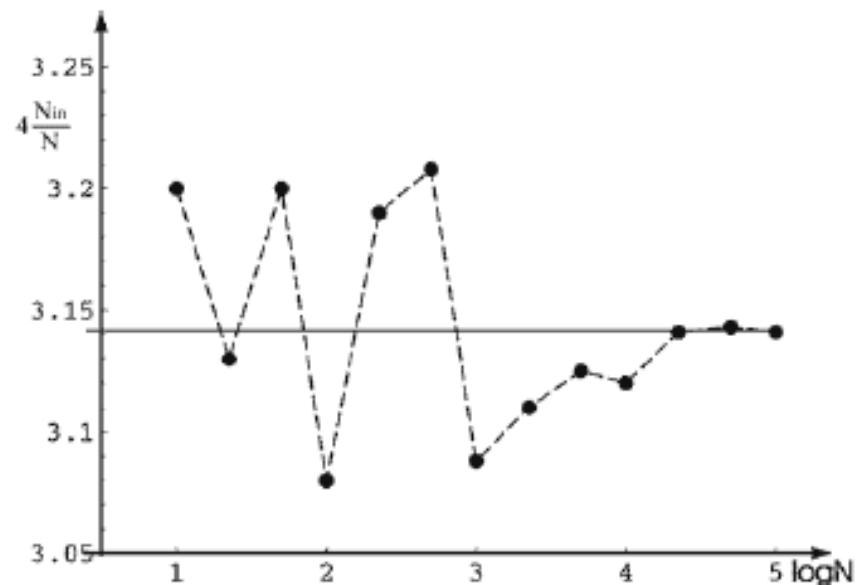
Example: circle in a square



$$x^2 + y^2 < R^2$$

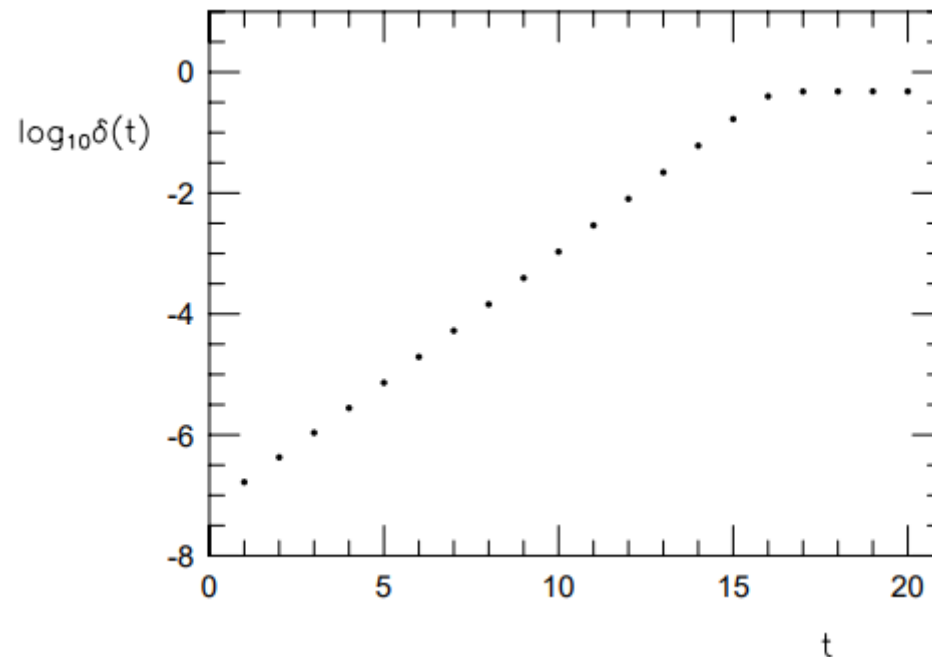
$$S_{in} = S \frac{N_{in}}{N}$$

$$\pi = 4 \frac{N_{in}}{N}$$



Increasing number of iterations decreases variance

Random number generators



Theoretically proven randomness:

ranlux: <http://luscher.web.cern.ch/luscher/ranlux/>

mixmax: <https://mixmax.hepforge.org/>

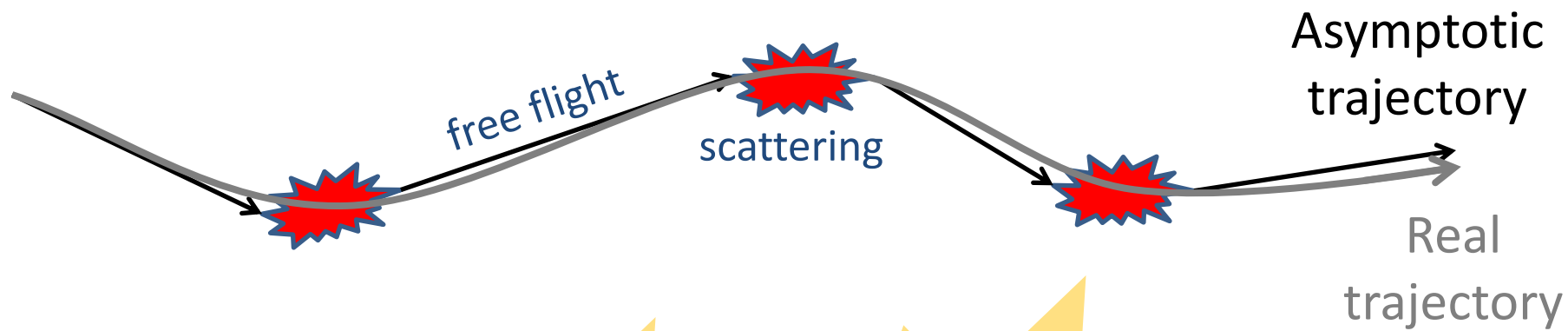
Both are available for FORTRAN and C

Radiation transport

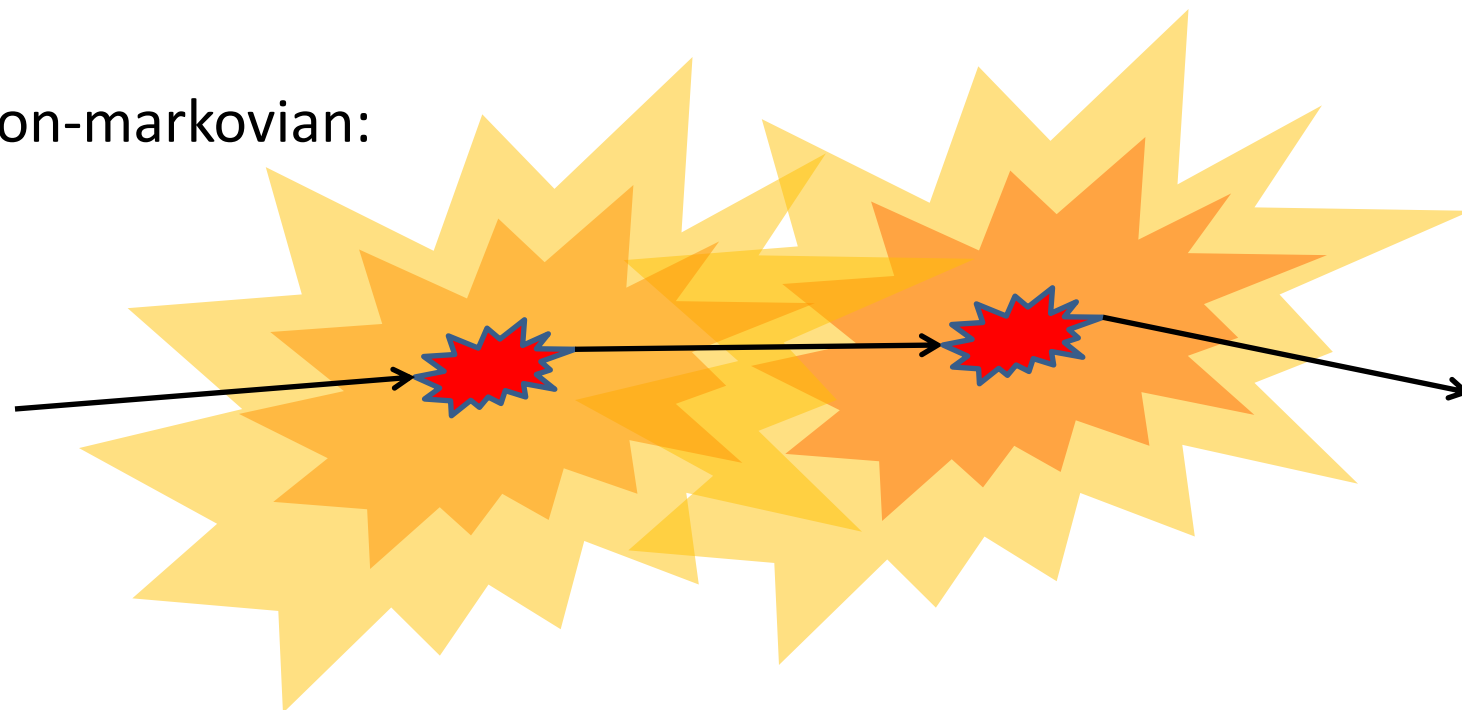
“Tools for investigating electronic excitation: experiment and multi-scale modelling”
Instituto de Fusión Nuclear "Guillermo Velarde", Universidad Politécnica de Madrid
<https://doi.org/10.20868/UPM.book.69109>

Radiation transport

Markov chain: duration of scattering event \ll free flight



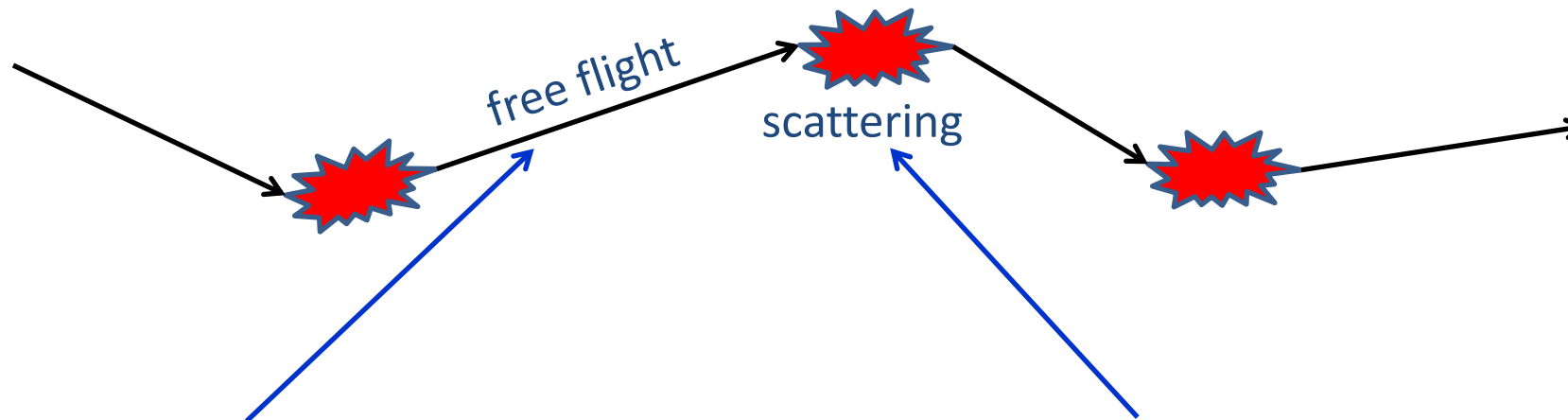
Non-markovian:



Keep in mind quantum speed limit theorem!

Monte Carlo of radiation transport

Quantities defined by probabilities:



Free flight distance

$$l = -\lambda \ln(\gamma)$$

$$\text{Mean free path: } \lambda = \frac{1}{n_e \sigma}$$

Scattering event:

- What kind of event?
- How much energy is lost?
- How is momentum changed?

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<https://doi.org/10.20868/UPM.book.69109>

Monte Carlo: key values

Sampled flight distance: $l = -\lambda \ln(\gamma) \quad \lambda = \frac{1}{n_e \sigma}$.

Mean free path: $\lambda^{-1} = n_{at} \int_{E_{\min}}^{E_{\max}} \int_{q_-}^{q_+} \frac{d^2 \sigma}{d(\hbar\omega) d(\hbar q)} d(\hbar\omega) d(\hbar q)$

Stopping power: $S_e = -\frac{dE}{dx} = n_{at} \int_{E_{\min}}^{E_{\max}} \int_{q_-}^{q_+} \hbar\omega \frac{d^2 \sigma}{d(\hbar\omega) d(\hbar q)} d(\hbar\omega) d(\hbar q)$

Kind of collisions: $P_i = \frac{\sigma_i}{\sum \sigma_i}$

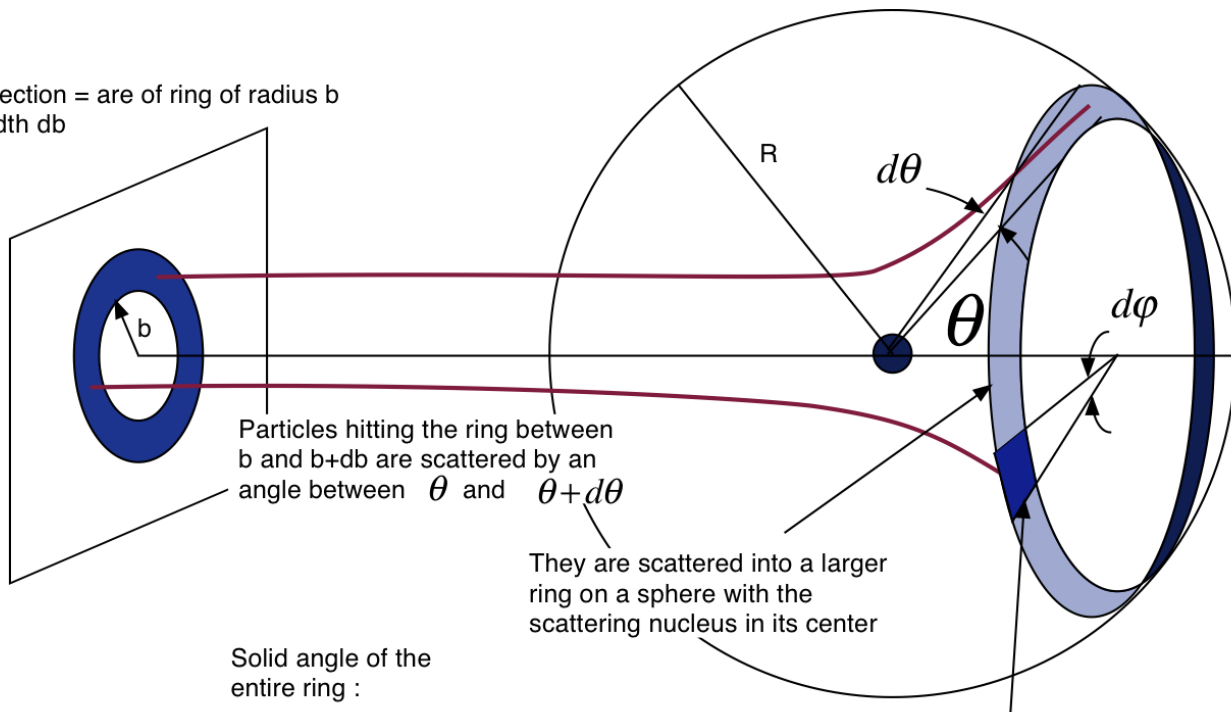
Energy loss in a scattering event: $\gamma \sigma = \int_{E_{\min}}^{\delta E} \int_{q_-}^{q_+} \frac{d^2 \sigma}{d(\hbar\omega) d(\hbar q)} d(\hbar\omega) d(\hbar q)$

In summary: cross sections define everything!

Cross sections

Cross sections

cross section = area of ring of radius b and width db



Particles hitting the ring between b and $b+db$ are scattered by an angle between θ and $\theta+d\theta$

They are scattered into a larger ring on a sphere with the scattering nucleus in its center

Solid angle of the entire ring :

$$d\Omega = \frac{2\pi R \sin(\theta) R d\theta}{R^2} = 2\pi \sin(\theta) d\theta$$

solid angle of small area:

$$d\Omega = \frac{d\phi R \sin(\theta) R d\theta}{R^2} = \sin(\theta) d\theta d\phi$$

Area, where particles should meet to scatter
(for given energy, into given angle)

Pic from:

<http://hep.physics.wayne.edu/~harr/courses/5210/w15/lecture29.htm>

Probability of scattering \sim cross section

Cross sections

L. Van Hove, Phys. Rev. **95**, 249 (1954)

Plane waves: $|k_{i,f}^e\rangle = V^{-\frac{1}{2}} \exp(i\mathbf{k}_{i,f}^e \cdot \mathbf{r})$ <- Fourier transform

Dynamic structure factor:

$$S_{ab}(q, \omega) = \sum_i P_i \langle i | \int \frac{dt}{2\pi} \exp(i\omega t) \int d\mathbf{r} \int d\mathbf{r}' \exp(-i\mathbf{q}(\mathbf{r} - \mathbf{r}')) \hat{n}_a(\mathbf{r}, t) \hat{n}_b(\mathbf{r}', 0) | i \rangle$$

First Born approximation: $\frac{d^2\sigma_{e-at}}{d(\hbar q)d(\hbar\omega)} = \underbrace{\frac{q}{2\pi\hbar^4} \frac{1}{v^2} \left(\frac{4\pi e^2}{q^2}\right)^2}_{\text{Scattering on individual atom}} \underbrace{(Z)^2 S_{ii}(\omega, q)}_{\text{Collective behavior of the system}}$

Scattering on individual atom

Collective behavior of the system

Cross sections

Fluctuation dissipation theorem:

$$\text{Im} \left[-\frac{1}{\varepsilon(\omega, q)} \right] = \frac{4\pi e^2}{q^2 \hbar} \left(1 - e^{-\frac{\hbar\omega}{T}} \right) S(\omega, q)$$

$$\frac{d^2\sigma}{d(\hbar\omega)d(\hbar q)} = \frac{2(Z_{eff}(v)e)^2}{n_{at}\pi\hbar^2v^2} \frac{1}{\hbar q} \left[1 - \exp\left(-\frac{\hbar\omega}{k_B T}\right) \right]^{-1} \text{Im} \left(\frac{-1}{\varepsilon(\omega, q)} \right)$$

Individual atom

Thermal factor

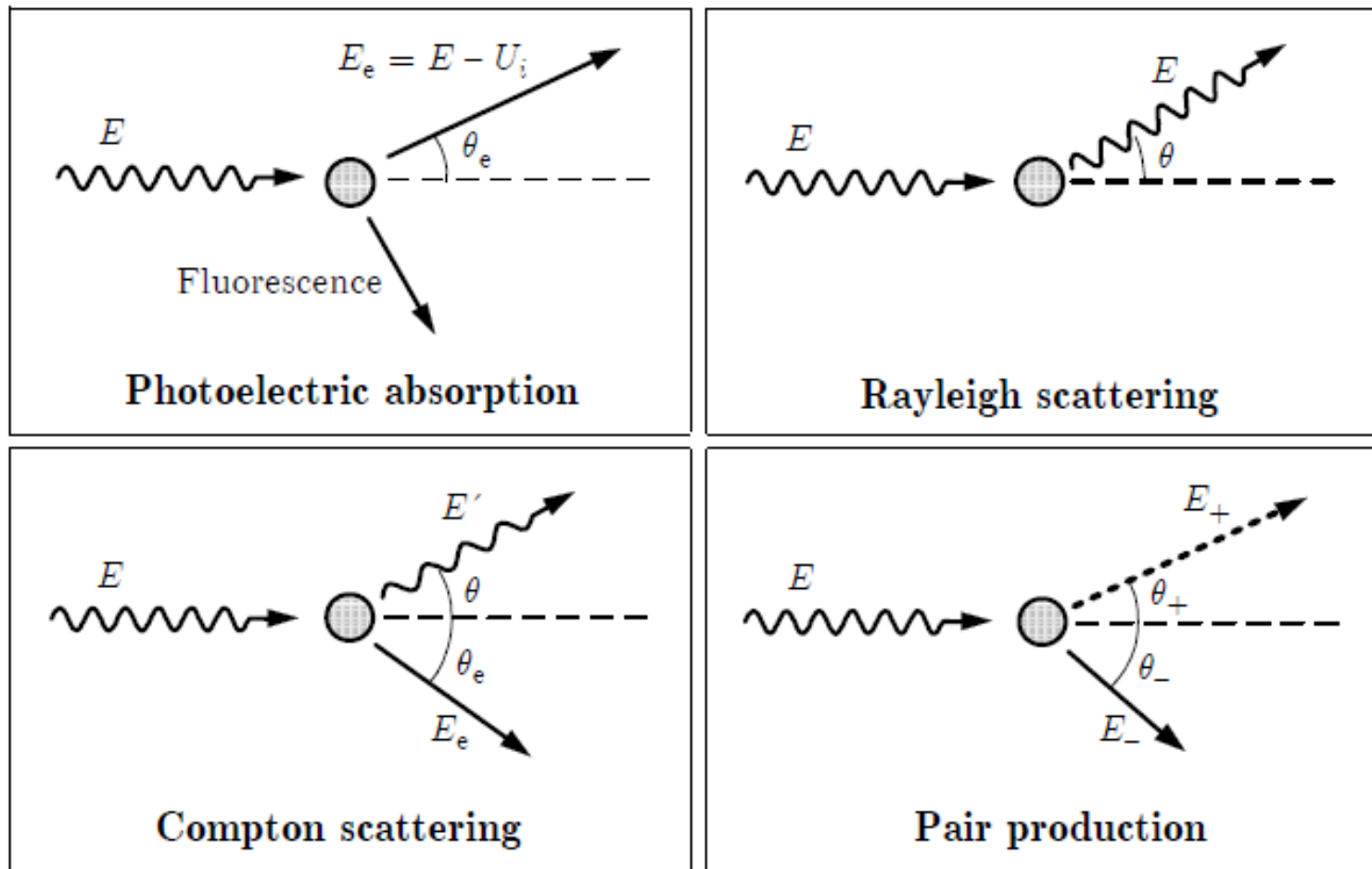
Loss function

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Photons

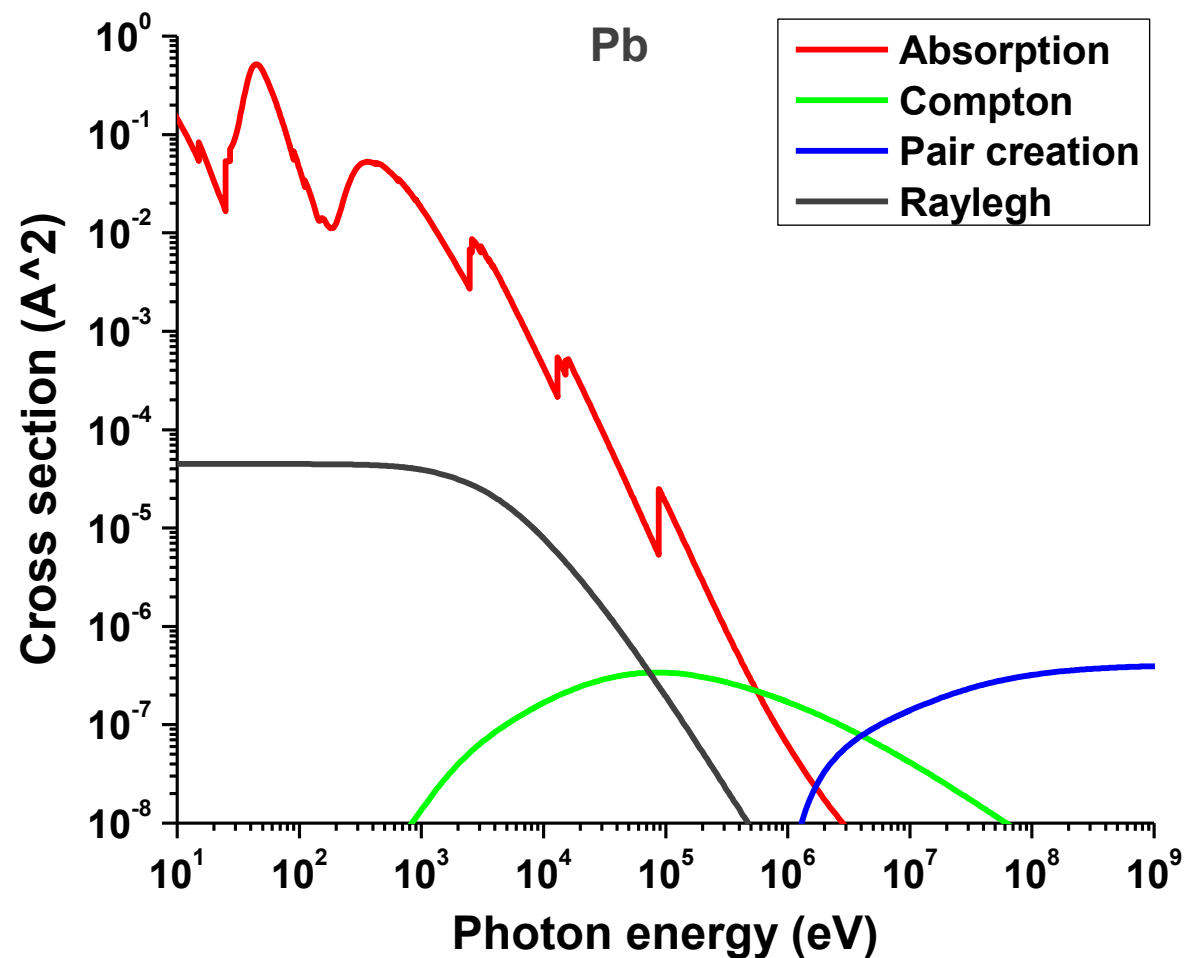
Photons

Scattering processes:



[F. Salvat, J.M. Fernandez-Varea, J. Sempau "PENELOPE 2014 - A Code System for Monte Carlo Simulation of Electron and Photon Transport" \(2014\)](#)

Photon cross sections



EPICS2023 database: <https://www-nds.iaea.org/epics/>

Henke's tables: http://henke.lbl.gov/optical_constants/atten2.html

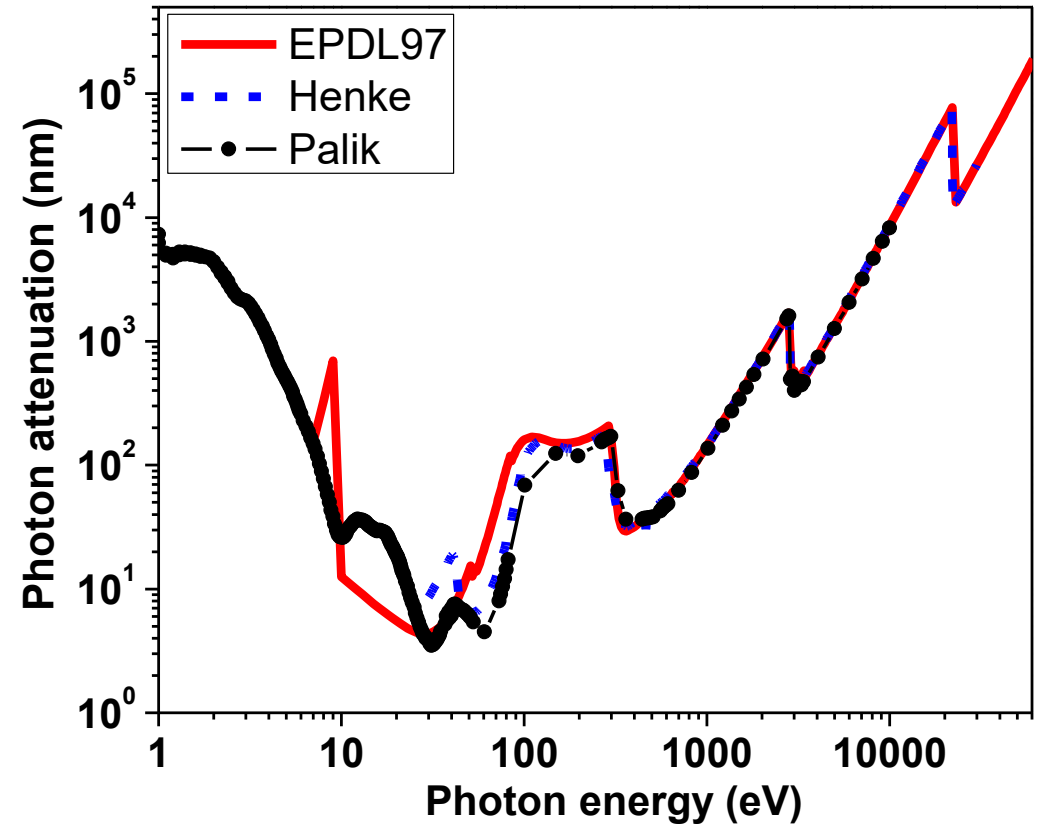
Photoabsorption: collective effects

In solids – band effects at low energies

Experimental optical coefficients:

$$\text{Im} \left[\frac{-1}{\varepsilon(\omega, q=0)} \right] = \frac{2nk}{(n^2 - k^2)^2 + (2nk)^2}$$

$$\text{Im} \left[\frac{-1}{\varepsilon(\omega, q=0)} \right] = \frac{c}{\lambda_{ph} \omega}$$

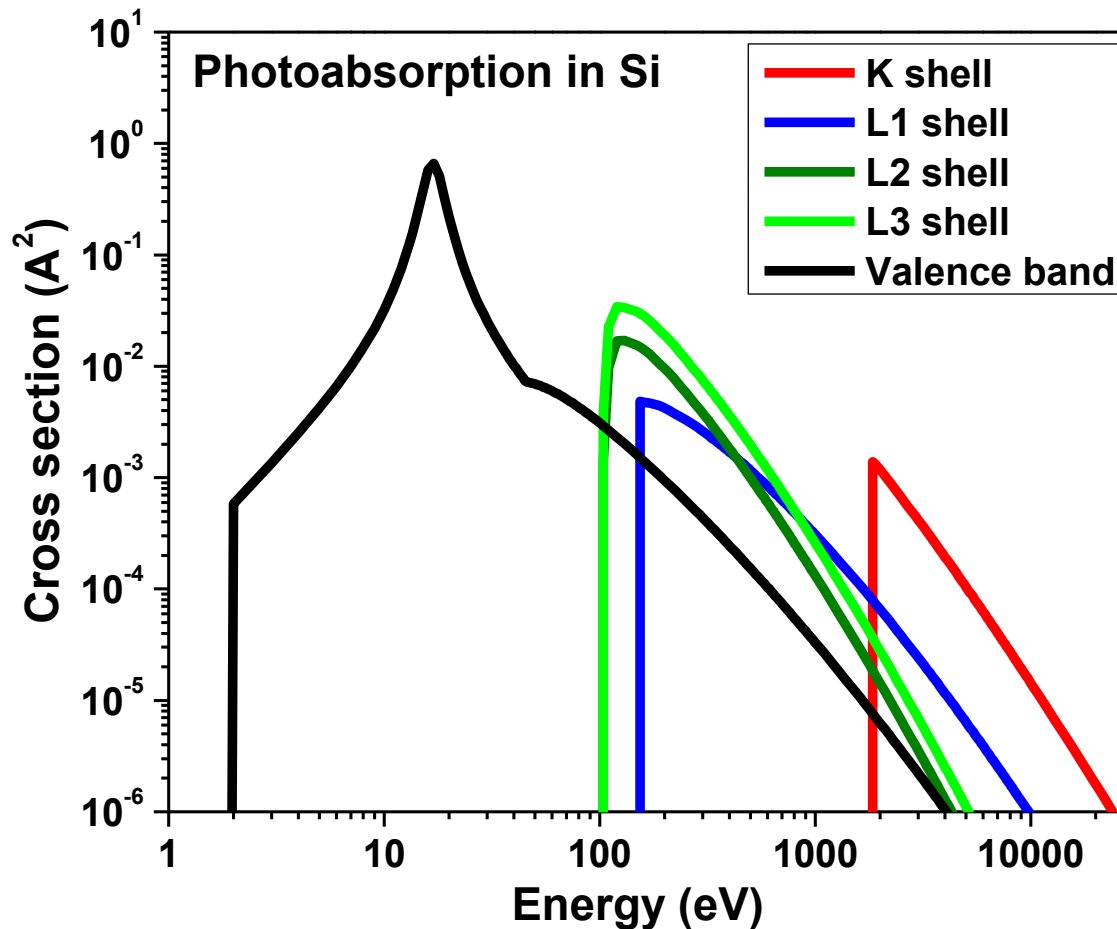


[Milov et al., JOSA B 35, B43 \(2018\)](#)

For photon energies below ~100 eV atomic cross sections differ from solids

Photoabsorption: collective effects

Atomic shell choice: $P_i = \frac{\sigma_i}{\sum \sigma_i}$



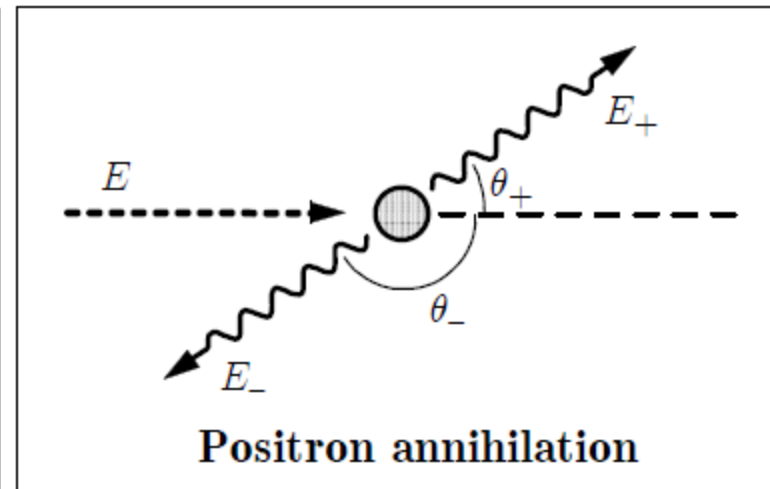
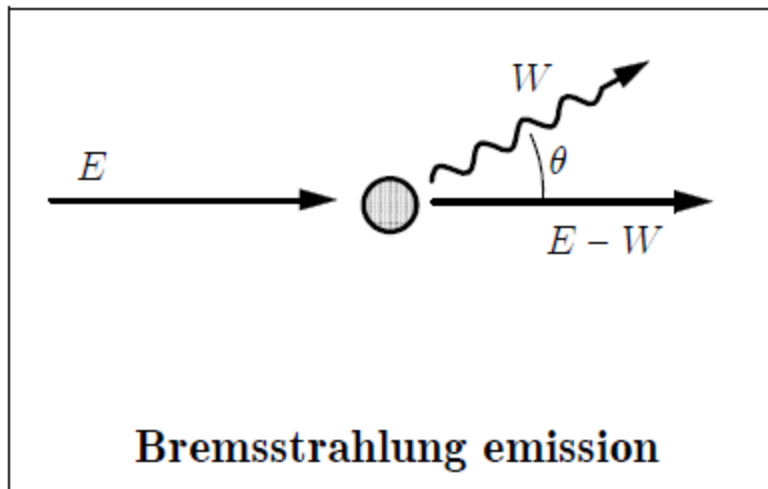
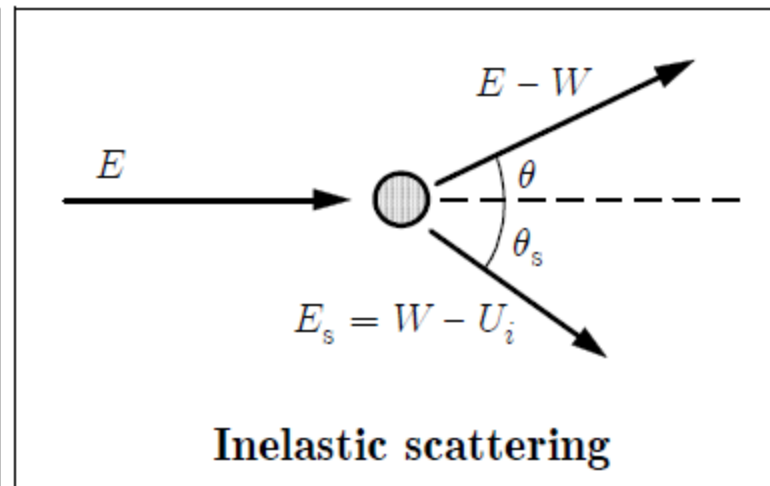
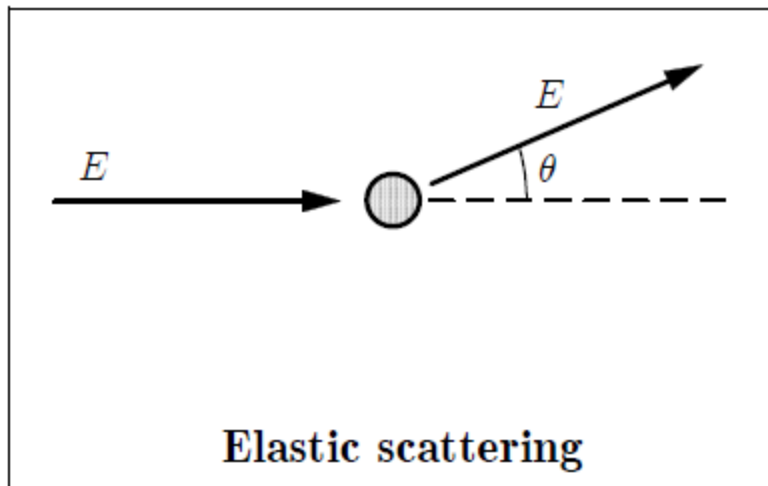
Photons: deepest shell ionization is most probable

[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Electrons / ions / charged particles

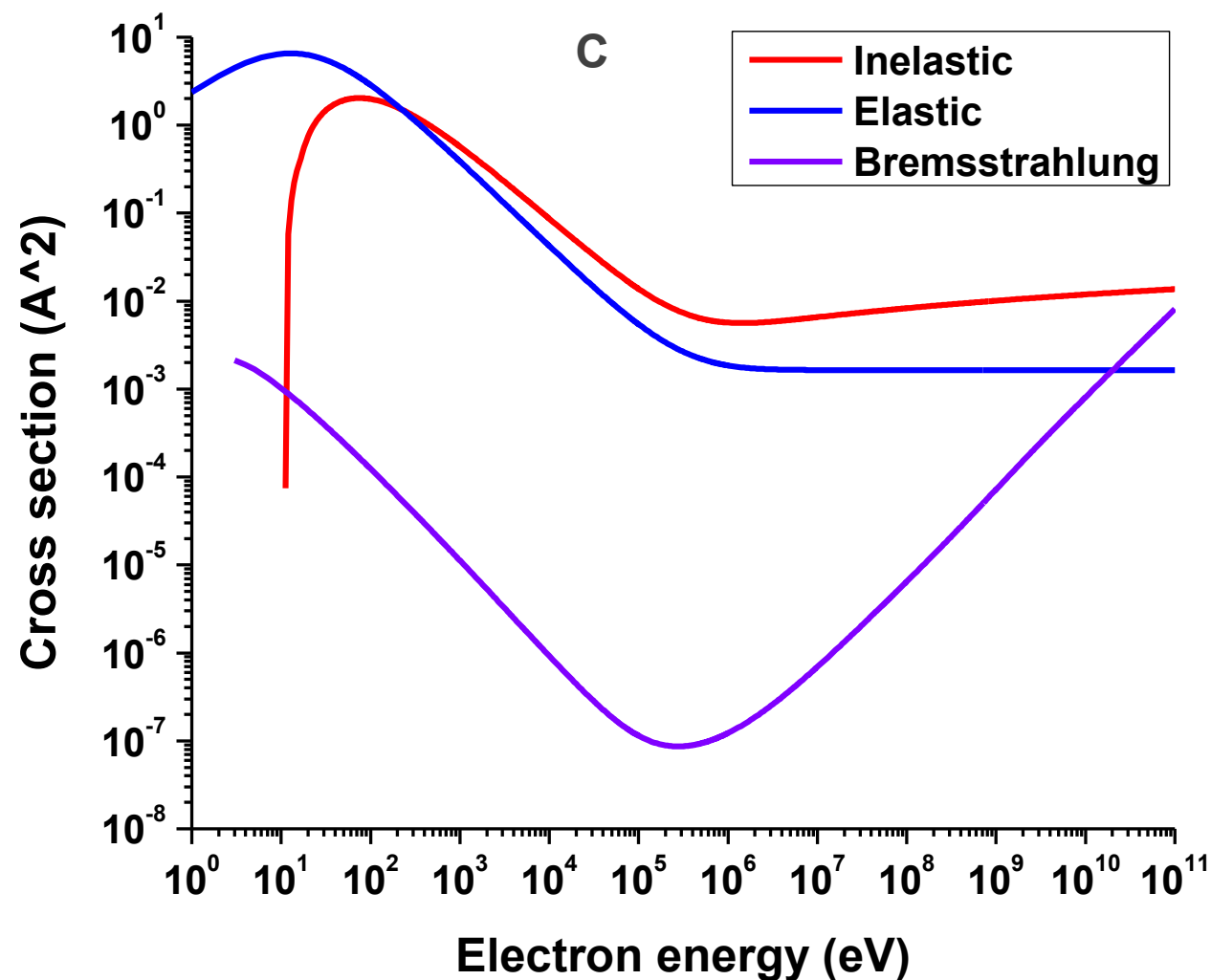
Electrons

Processes



[F. Salvat, J.M. Fernandez-Varea, J. Sempau "PENELOPE 2014 - A Code System for Monte Carlo Simulation of Electron and Photon Transport" \(2014\)](#)

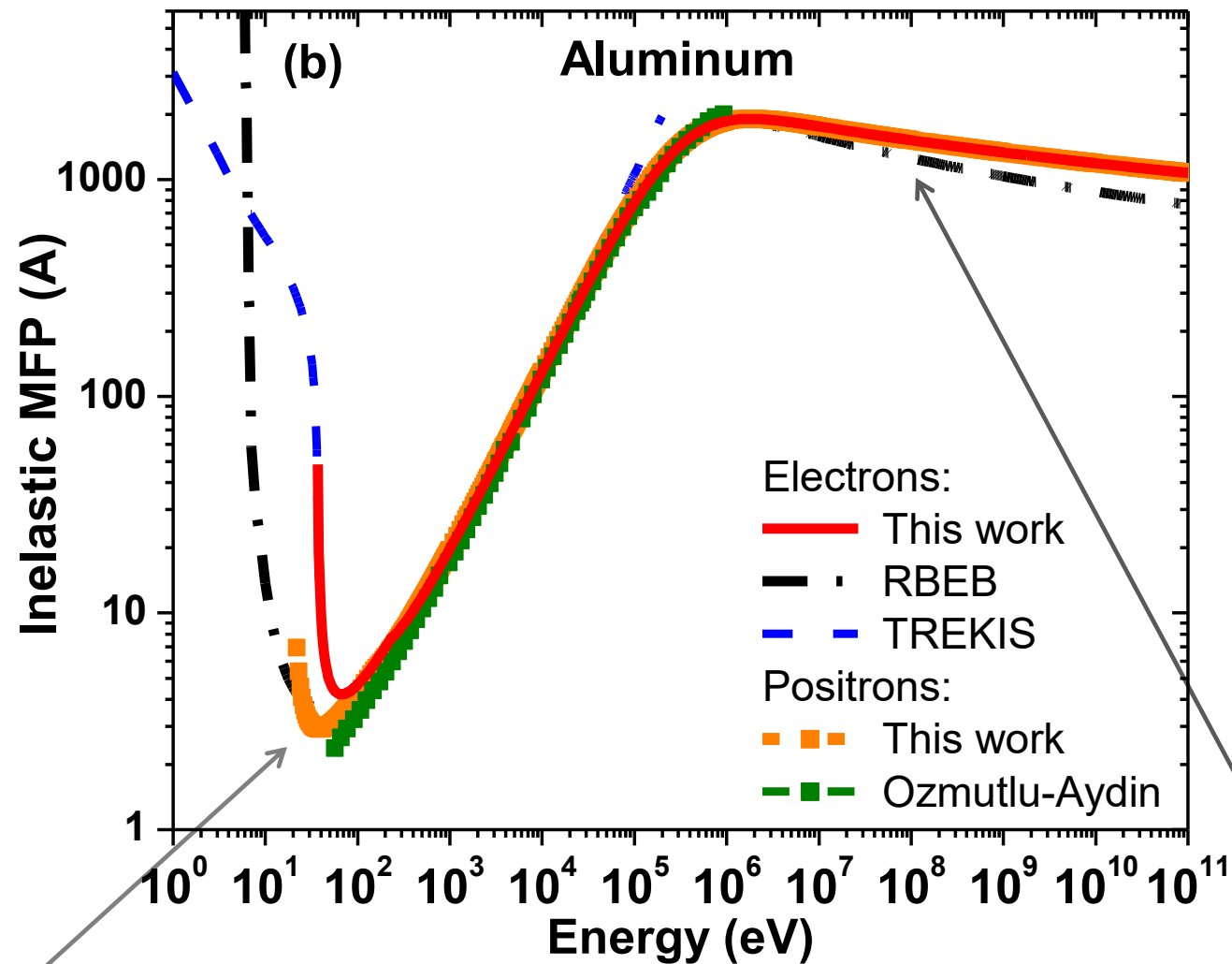
Electrons



Inelastic : impact ionization, electron-electron, electron-plasmon

Elastic : electron-atom, electron-phonon

Electrons: “inelastic” scattering

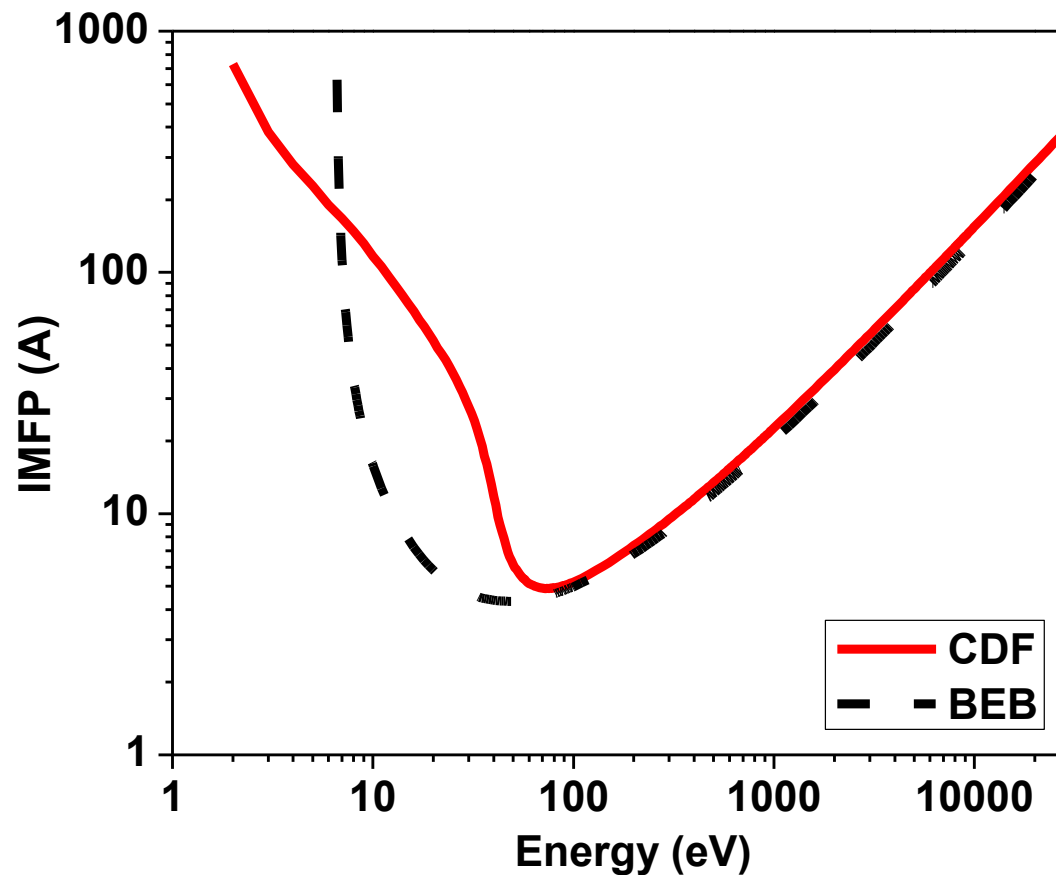


Scattering on plasmons (collective)

Scattering on individual atoms

[Medvedev, *et al.* J. Phys. D 53, 235302 (2020)]

Caution!

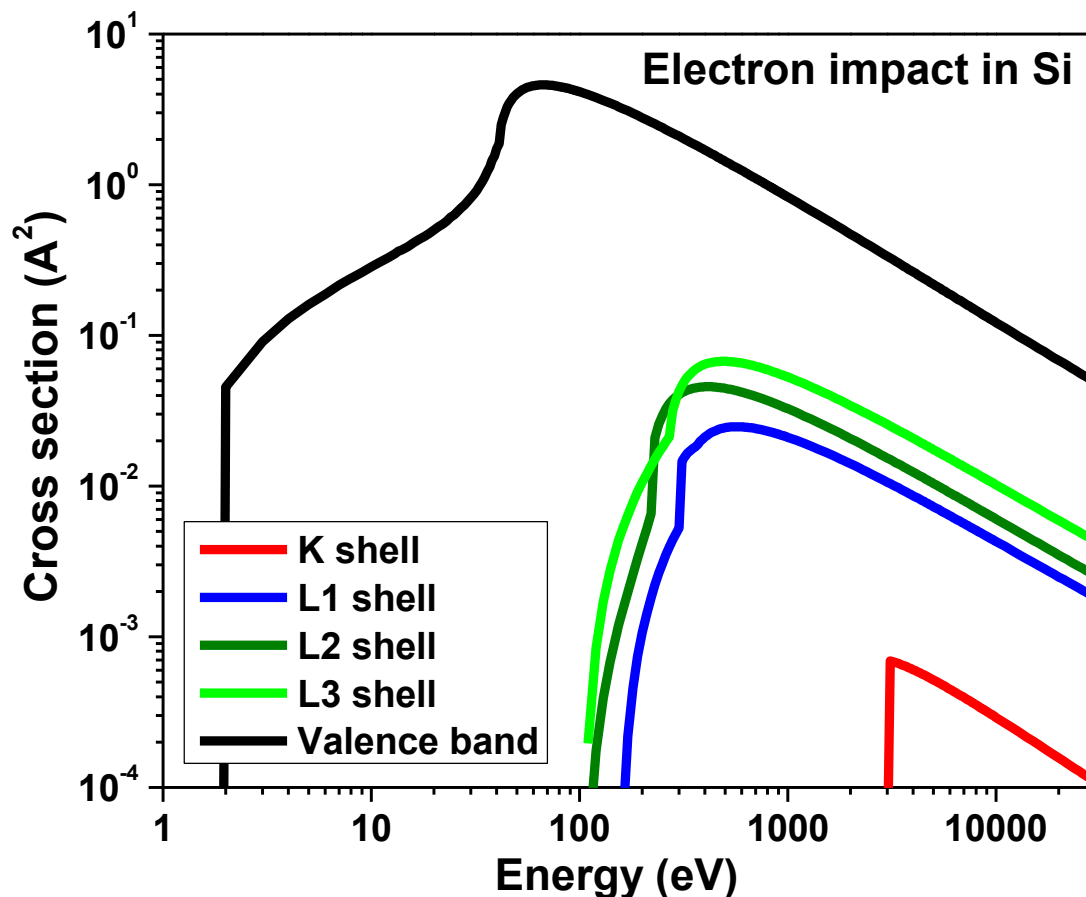


Standard MC codes often use atomic cross sections
That only works for very fast particles!

[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Impact ionization (“inelastic”)

Atomic shell choice: $P_i = \frac{\sigma_i}{\sum \sigma_i}$

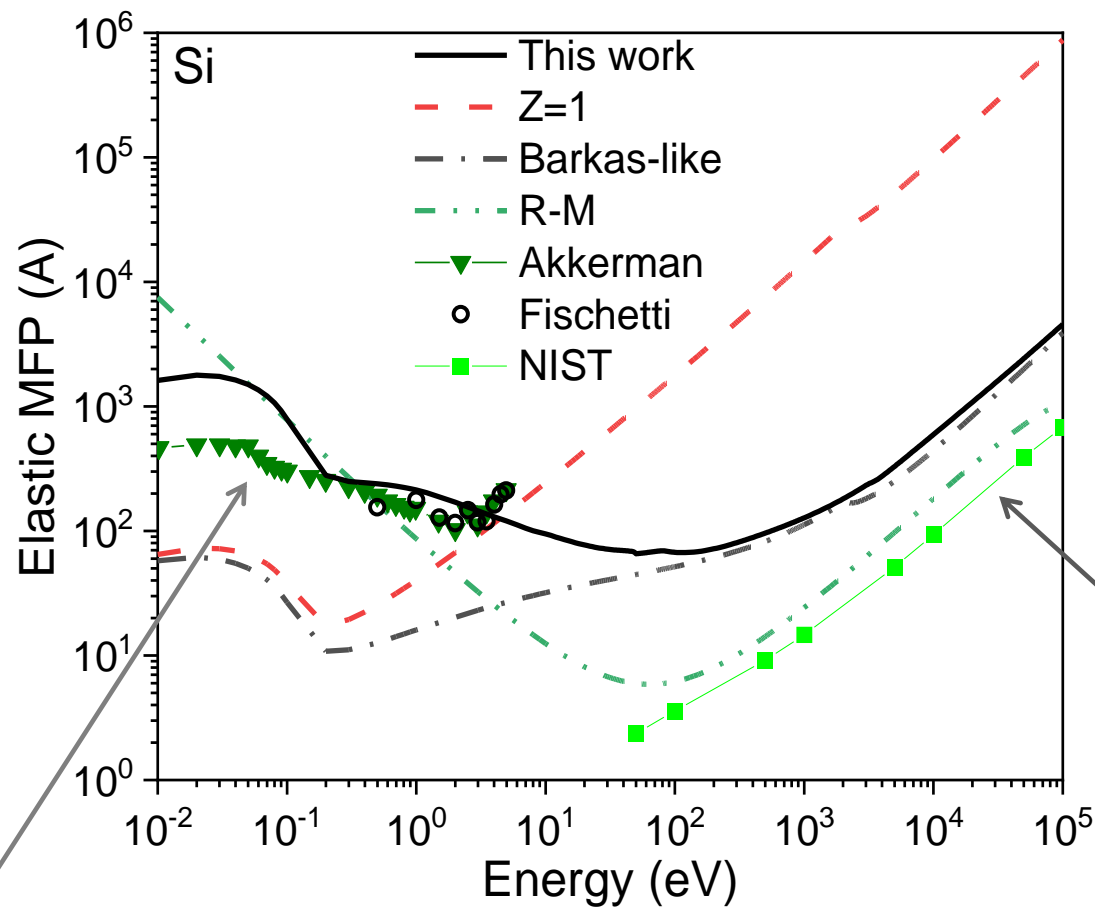


Electrons: highest shell ionization is most probable!

[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Electrons: “elastic” scattering

$$\frac{d^2\sigma_{e-at}}{d(\hbar q)d(\hbar\omega)} = \frac{2e^2}{n_i\pi\hbar^2v^2} \frac{1}{\hbar q} \left(1 - e^{-\frac{\hbar\omega}{k_B T}}\right)^{-1} \left[Z - Z_I f_I(\tilde{q}) - N_{VB} \left(1 - \frac{1}{|\varepsilon_{VB}(\omega, \tilde{q})|}\right) \right]^2 \text{Im} \left[\frac{-1}{\tilde{\varepsilon}_{at}(\omega, q)} \right]$$



Scattering on phonons (collective)

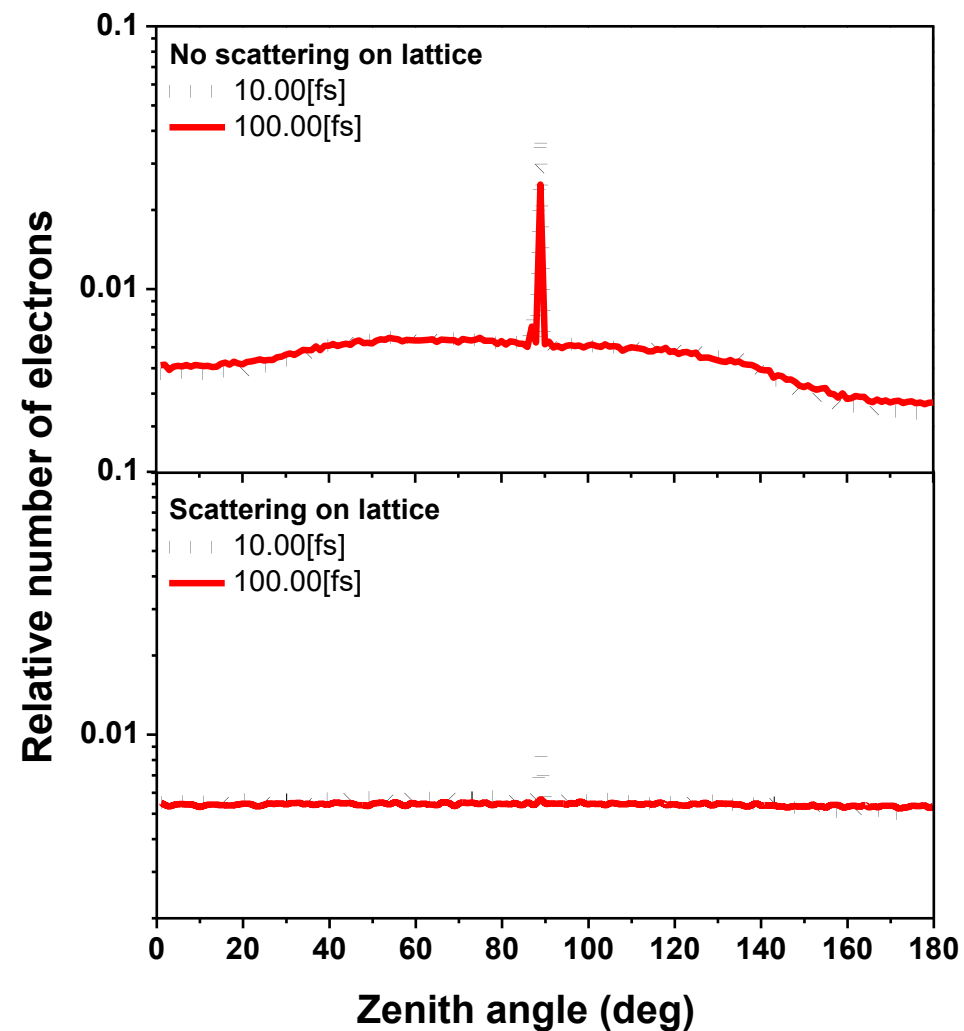
Scattering on individual atoms

[Medvedev, *et al.* J. Appl. Phys. 137, 015903 (2025)]

Electrons: “elastic” scattering

$$E_{max} = 4E_e \frac{M_{ion}m_e}{(M_{ion} + m_e)^2} \sim 4E_e \frac{m_e}{M_{ion}}$$

Elastic scattering transfers small amount
of energy,
but alters momentum a lot!



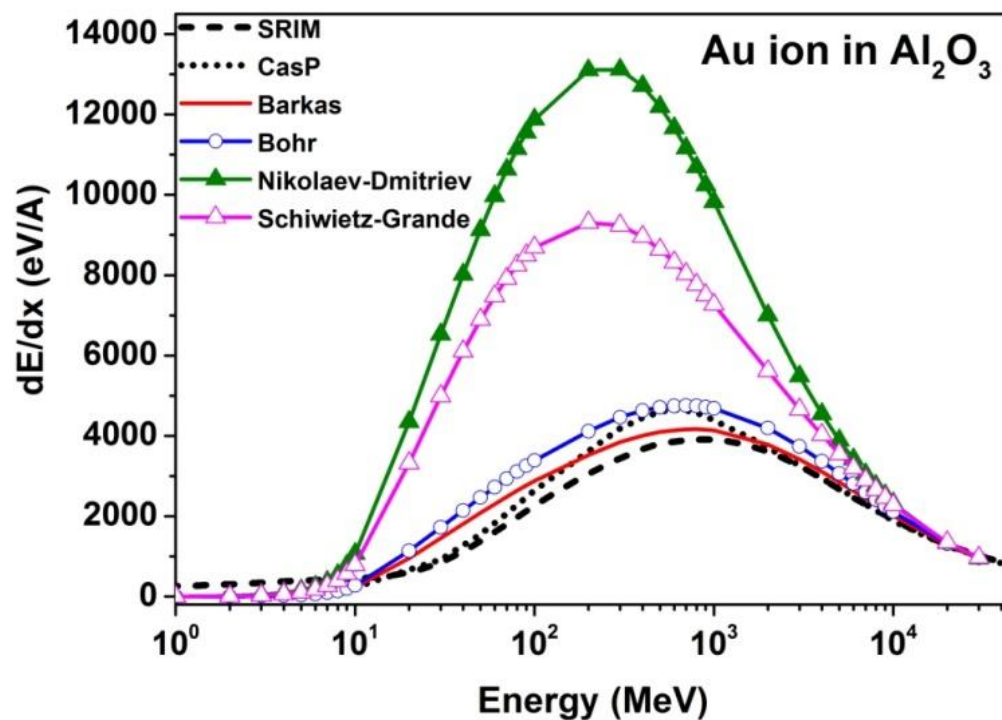
[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Ions

*Ions are like electrons,
but heavy and highly charged*

$$\frac{d^2\sigma}{d(\hbar\omega)d(\hbar q)} = \frac{2(Z_{eff}(v)e)^2}{n_{at}\pi\hbar^2v^2} \frac{1}{\hbar q} \left[1 - \exp\left(-\frac{\hbar\omega}{k_B T}\right)\right]^{-1} \text{Im}\left(\frac{-1}{\varepsilon(\omega, q)}\right)$$

Barkas formula: $Z_{eff}(v) = Z_{ion} \left[1 - \exp\left(-\frac{v}{v_0} Z_{ion}^{-2/3}\right)\right]$



For ions, the Born approximation requires effective charge!

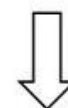
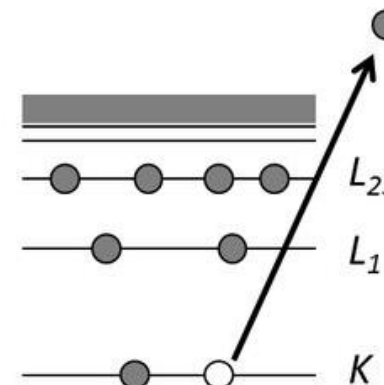
[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Holes

Core holes

Core holes (deep-shell):

Atomic excitation
(ionization)

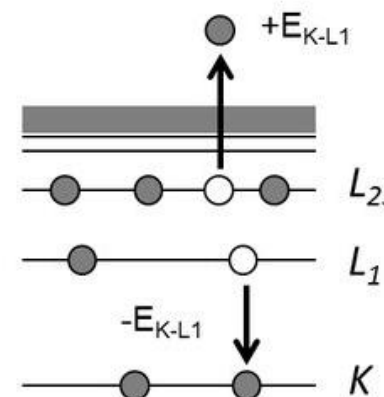


Auger and radiative decay times

EPICS2023 database:

<https://www-nds.iaea.org/epics/>

Auger transition
(KL_1L_{23})



Pic from: Yu-Pu Lin “Functionalization of two-dimensional nanomaterials based on graphene”

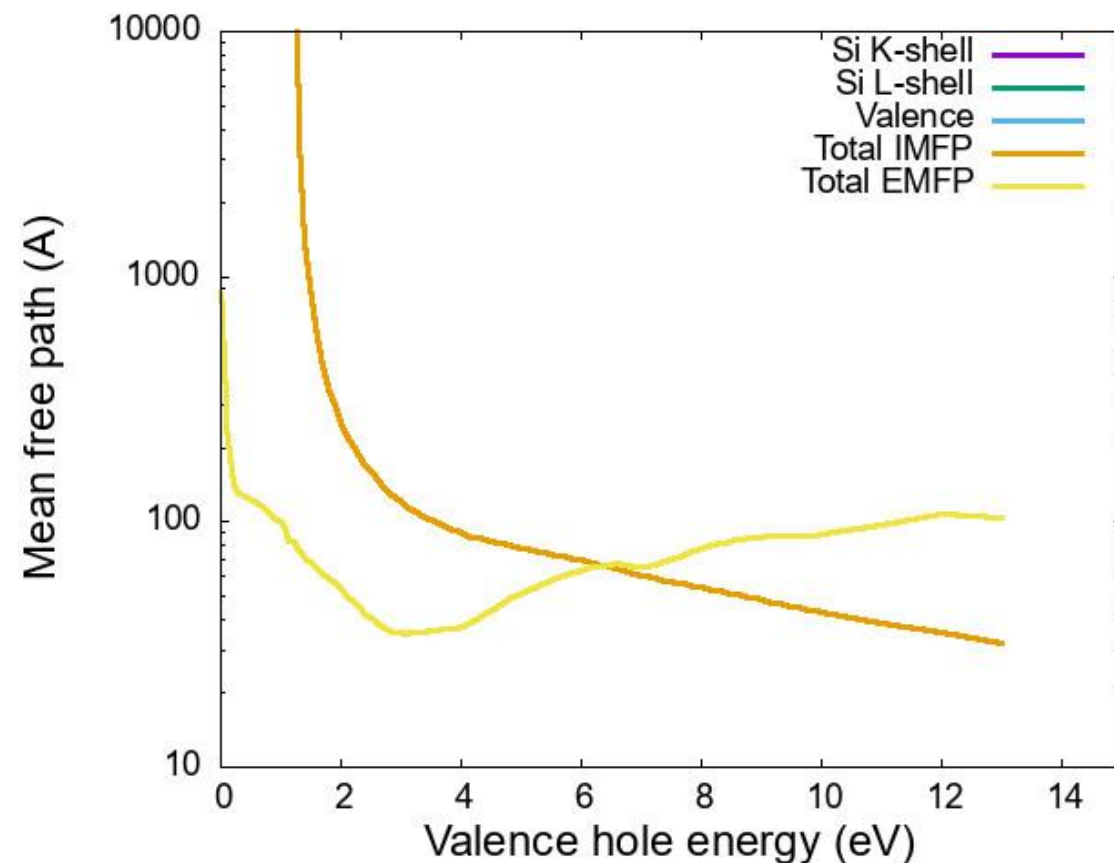
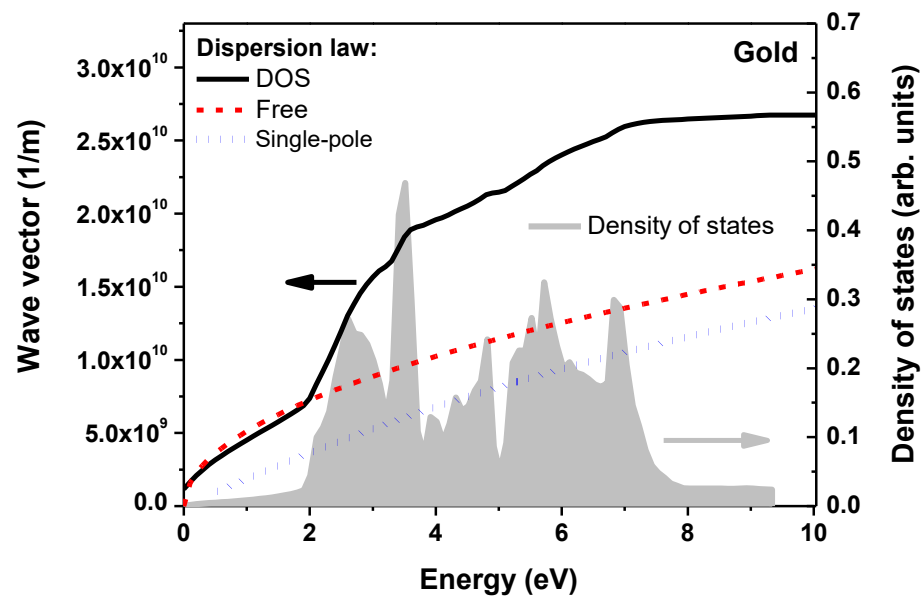
Valence holes

Valence holes in solids are mobile: analogous to electrons (almost)

Standard values are only valid at the top of valence band

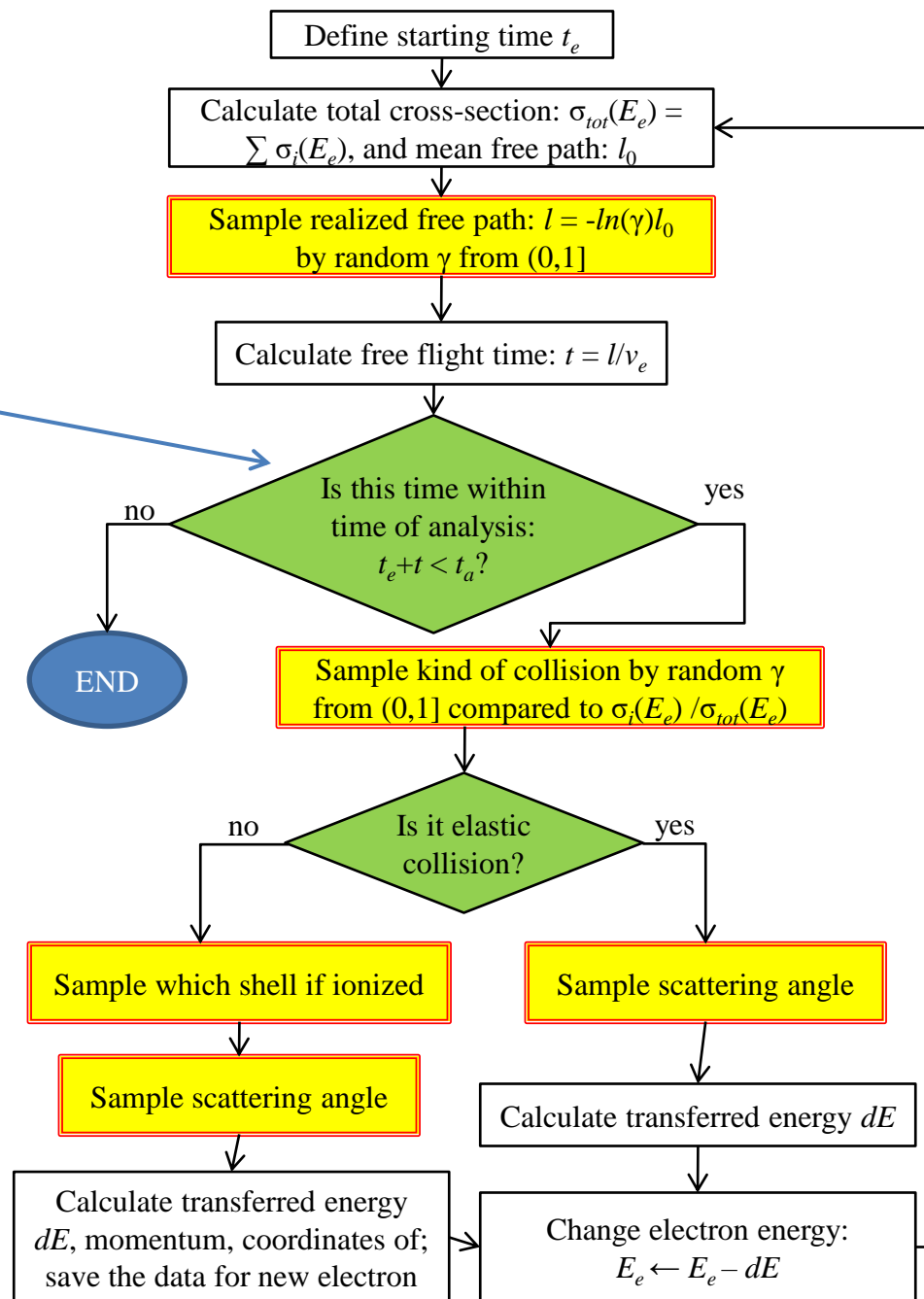
Suggested extension – mass from DOS

$$D(E) = \frac{s}{2\pi^2} q^2(E) \frac{dq}{dE} \quad q(E) = \sqrt[3]{\frac{6\pi^2}{s} \int_0^E d\varepsilon \cdot D(\varepsilon)}$$



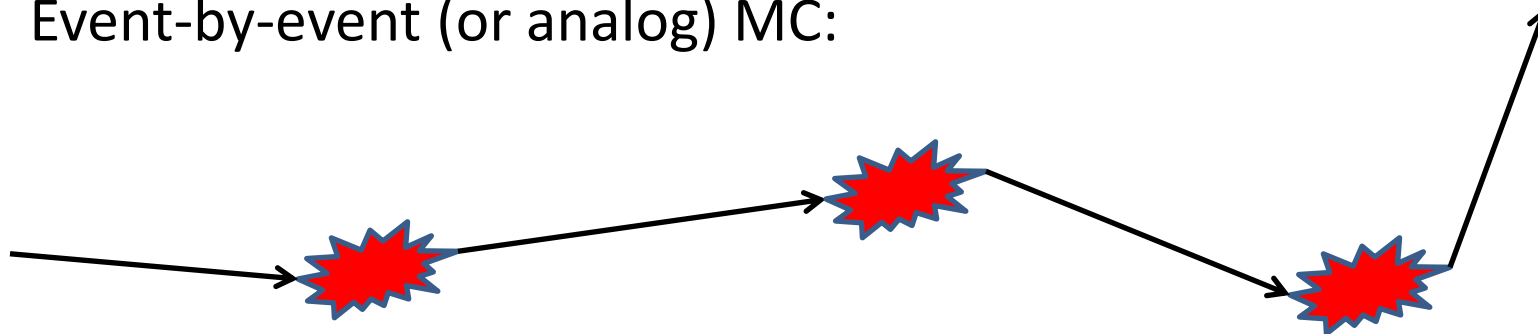
MC codes and algorithms

Energy < cut-off

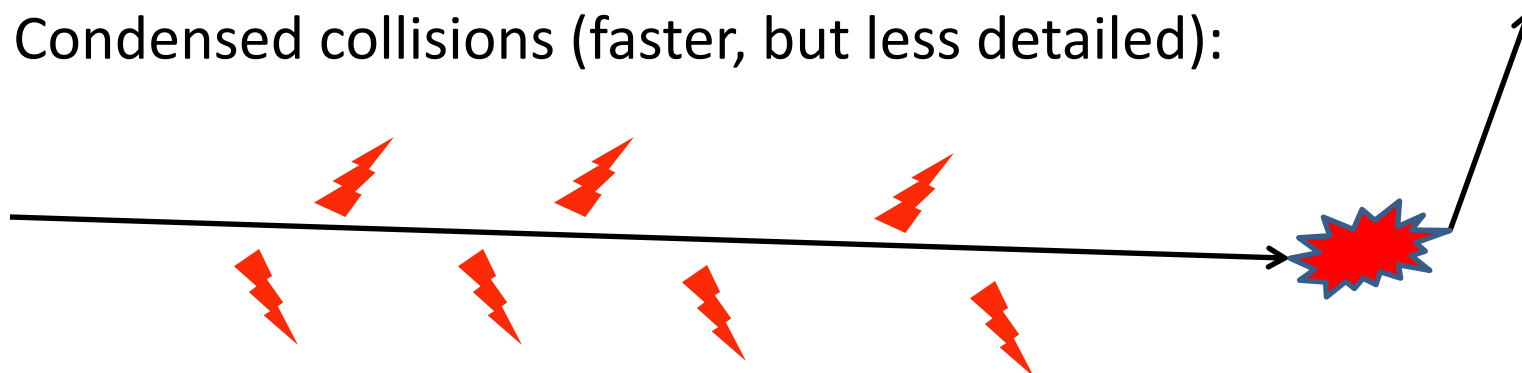


MC algorithms

Event-by-event (or analog) MC:



Condensed collisions (faster, but less detailed):



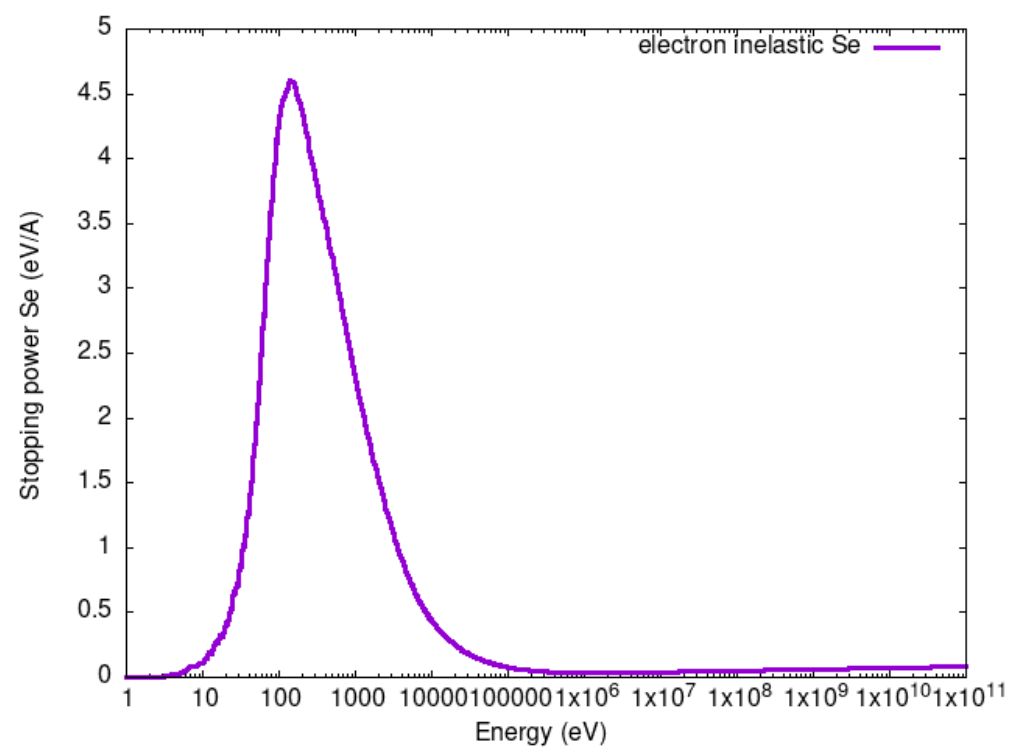
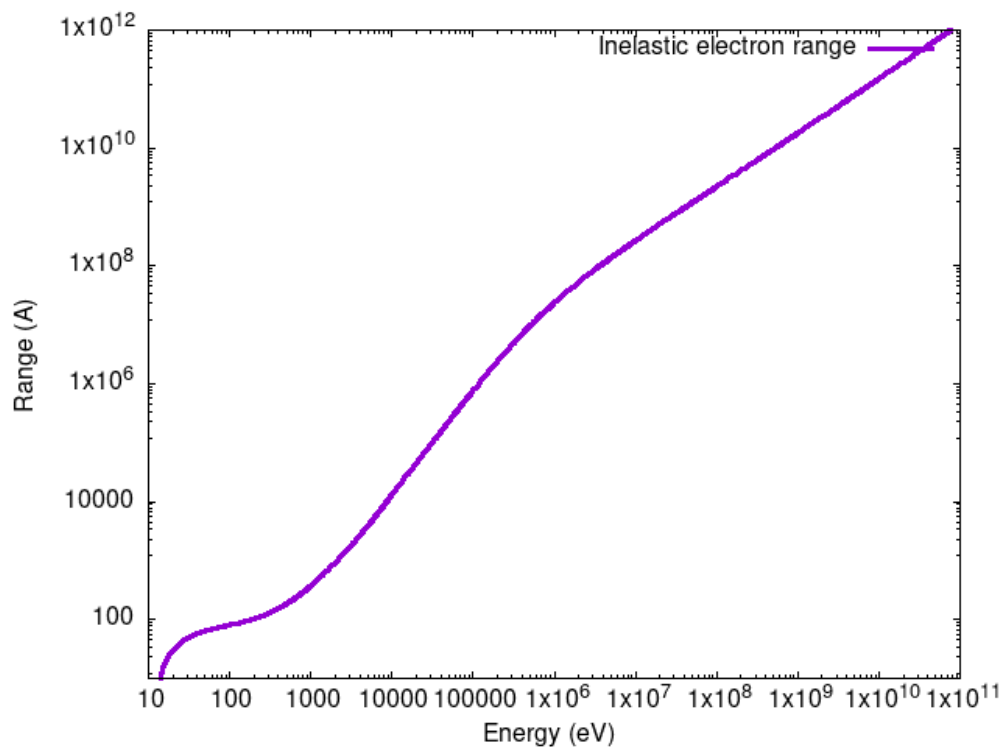
For combined simulations, only event-by-event works

Typical MC outputs: range

Range: total distance until coming to a “stop” (down to 10 eV):

$$R_0(E) = \int_{10\text{eV}}^E \frac{dE'}{S(E')}$$

$$S_e = -\frac{dE}{dx} = n_{at} \int_{E_{\min}}^{E_{\max}} \int_{q_-}^{q_+} \hbar\omega \frac{d^2\sigma}{d(\hbar\omega)d(\hbar q)} d(\hbar\omega)d(\hbar q)$$



Typical MC outputs: cumulative dose

Energy cut off is often used (beware, may be unphysical!)

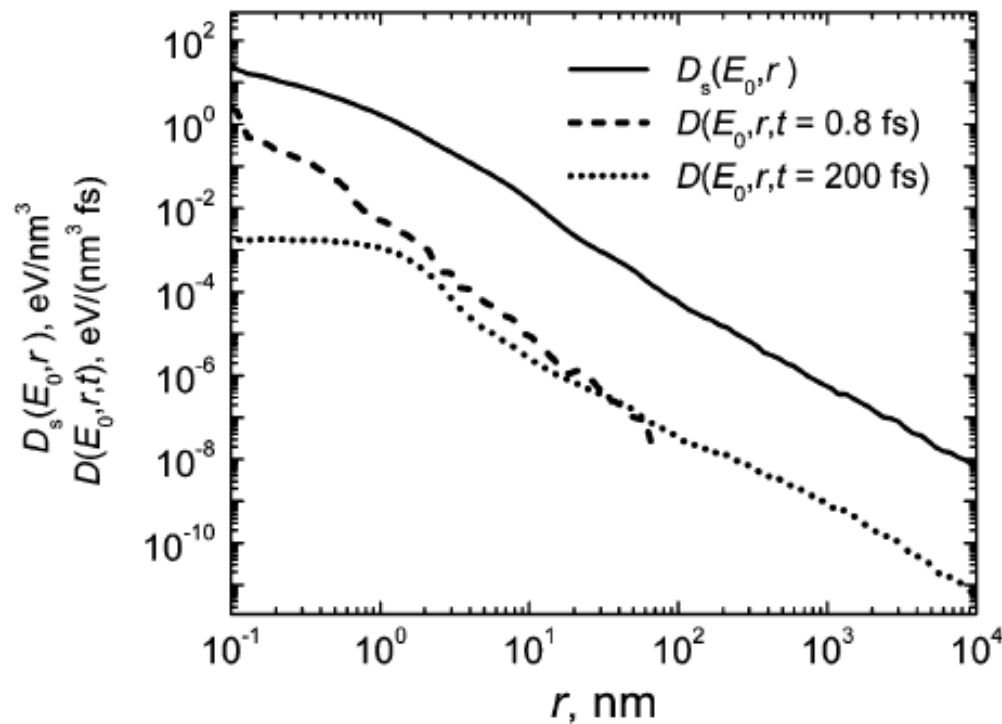


Fig. 2. Differential spatial deposited energy distributions at two discrete times for 100 MeV/amu Ne ions. The upper line is the energy deposition distribution integrated over time.

Cumulative dose does not correspond to real dose at any time instant

MC codes

Standard MC codes for radiation transport:

GEANT4: <https://geant4.web.cern.ch/>
FLUKA: <http://www.fluka.org/fluka.php>
PENELOPE: <http://pypenelope.sourceforge.net>
PHITS: <https://phits.jaea.go.jp/>
MCNP: <https://mcnp.lanl.gov/>
TART: <http://redcullen1.net/homepage.new/tart2016.htm>

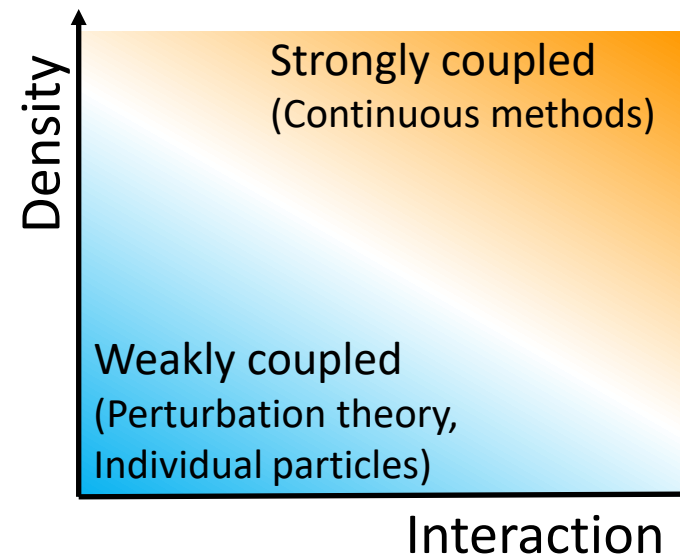
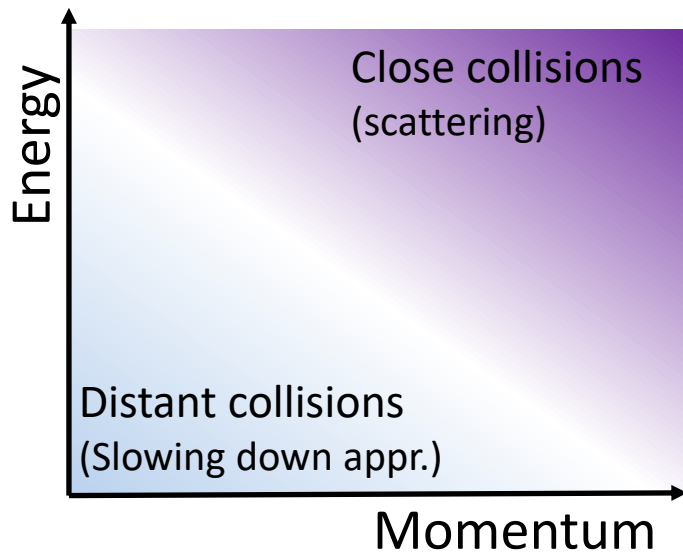
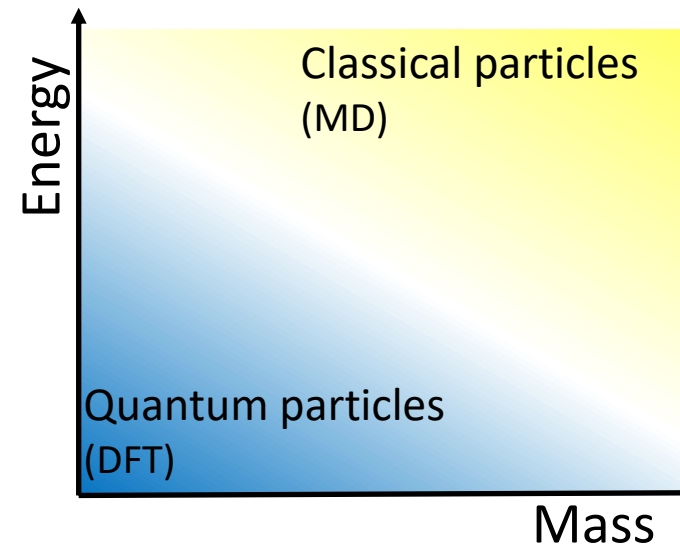
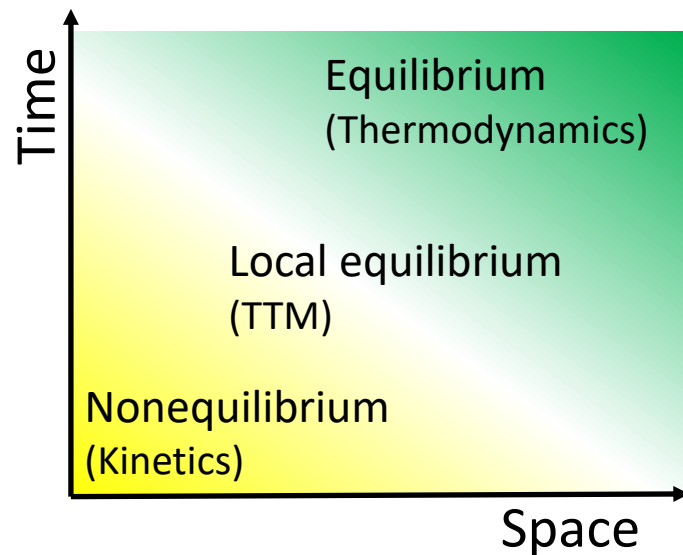
Our MC codes:

TREKIS-3: <https://github.com/N-Medvedev/TREKIS-3>
TREKIS-4: <https://github.com/N-Medvedev/TREKIS-4>



Combining MC codes with others

Combined / hybrid / multiscale approaches

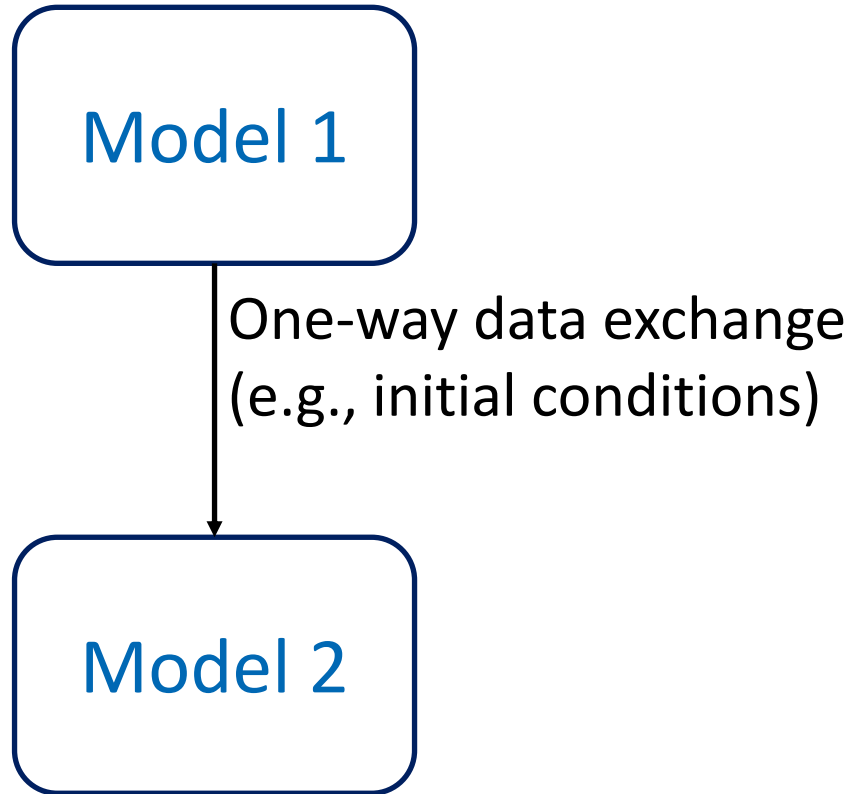


Apostolova et al., "Tools for investigating electronic excitation: experiment and multi-scale modelling"

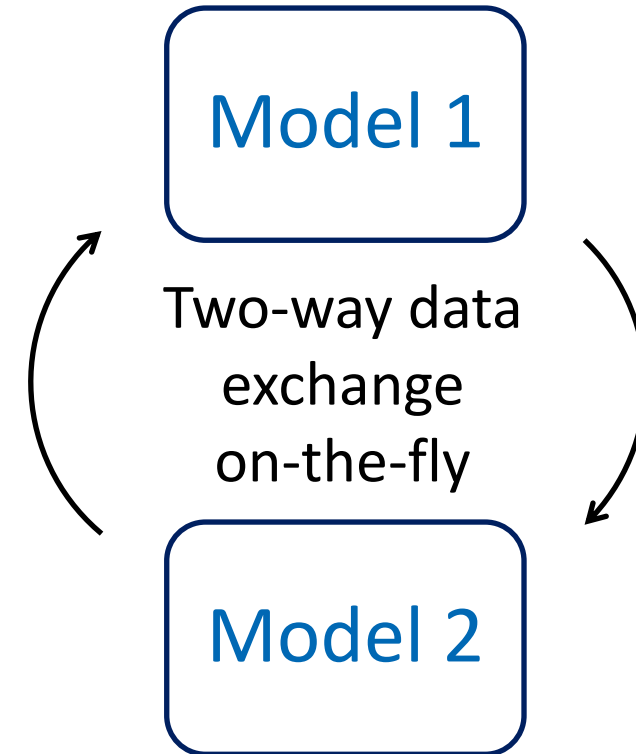
<https://doi.org/10.20868/UPM.book.69109>

Combining MC codes

Model without feedback



Model with feedback

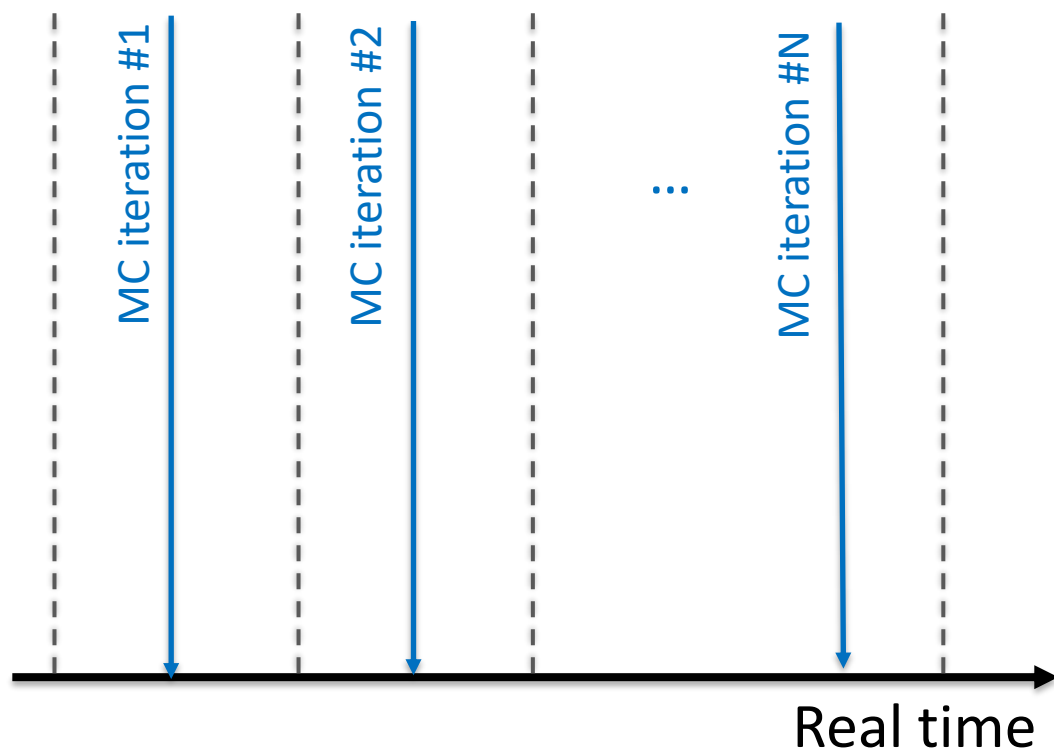


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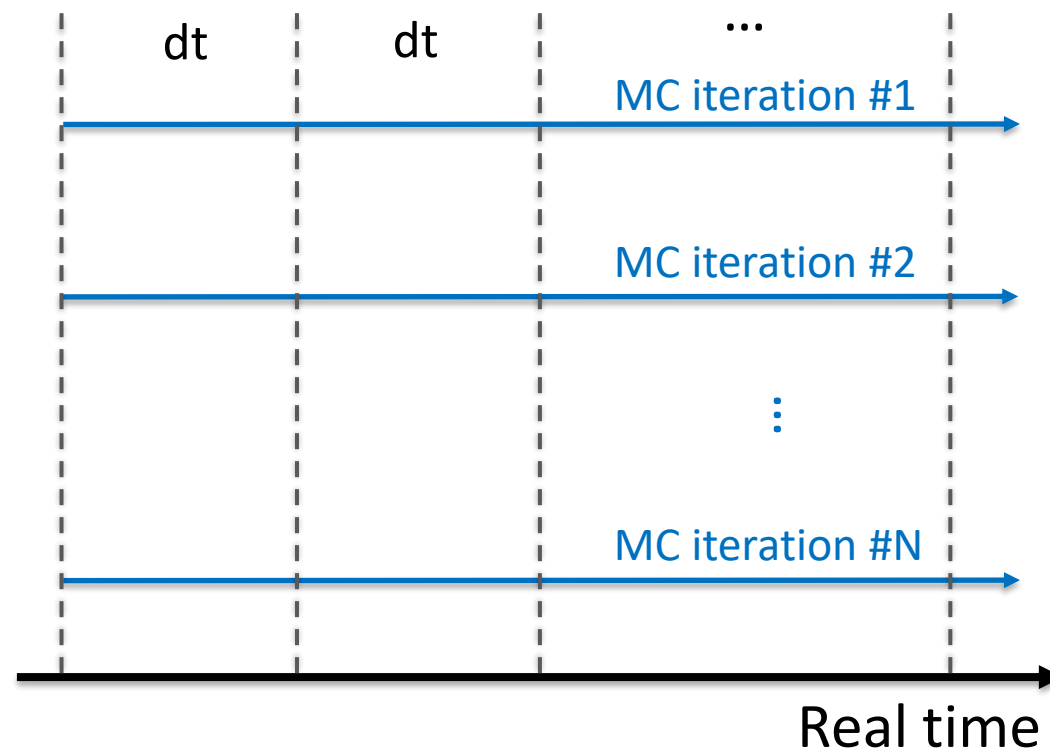
Combining MC codes

Standard MC method
(combining without feedback)



Output: distribution at final time

MC methods as part of hybrids with feedback



Output: distribution at each timestep

Apostolova et al., "Tools for investigating electronic excitation: experiment and multi-scale modelling"

<https://doi.org/10.20868/UPM.book.69109>

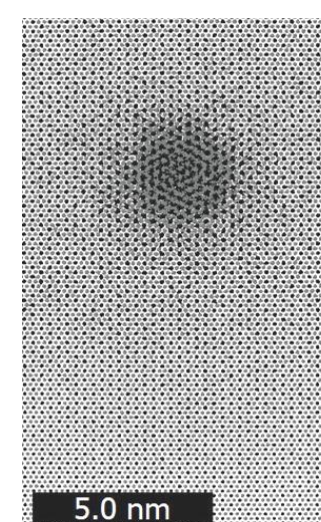
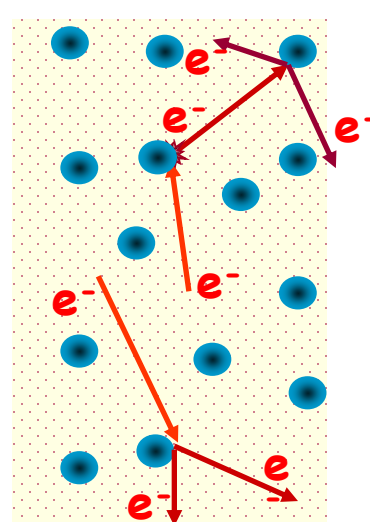
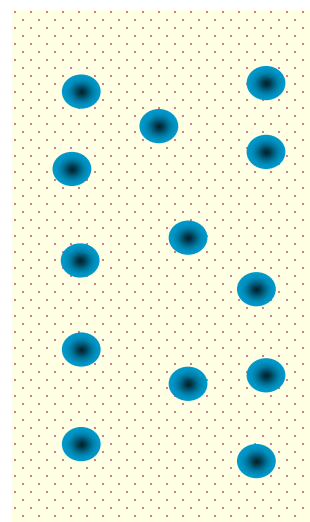
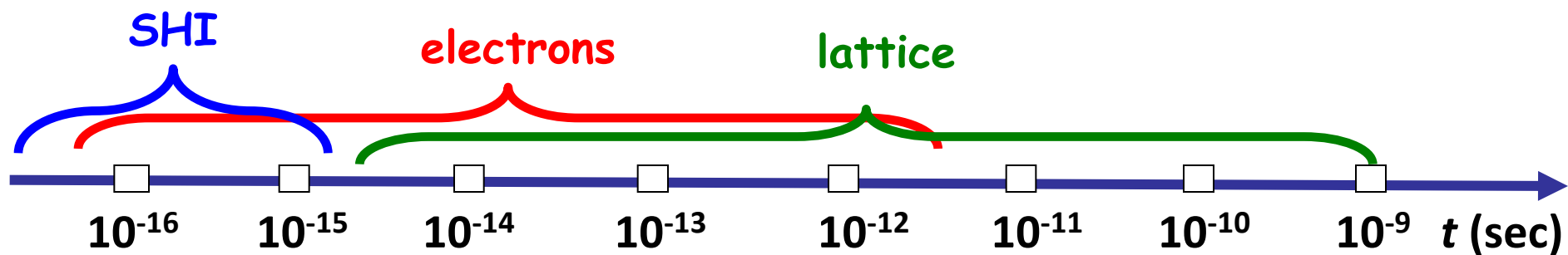
Example: Models without feedback

Example: ion beam



Time-Resolved Electron Kinetics in SHI-Irradiated Solids

Event-by-event Monte Carlo code for simulation of



[R. Rymzhanov]



Example: ion beam



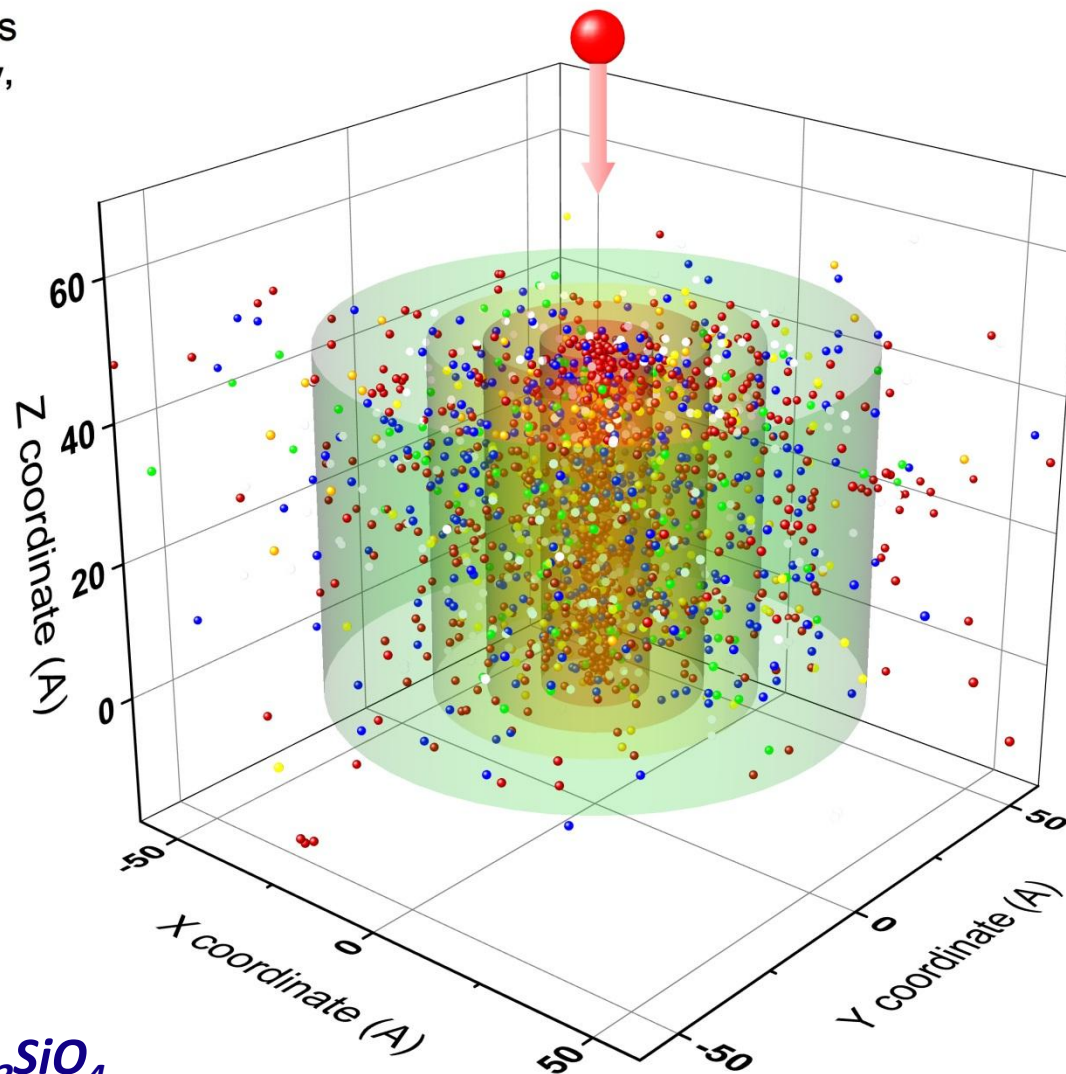
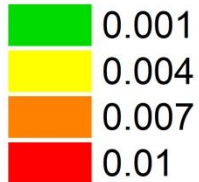
● Electrons

Electron energy,
eV



○ Lattice

Lattice energy
eV/atom

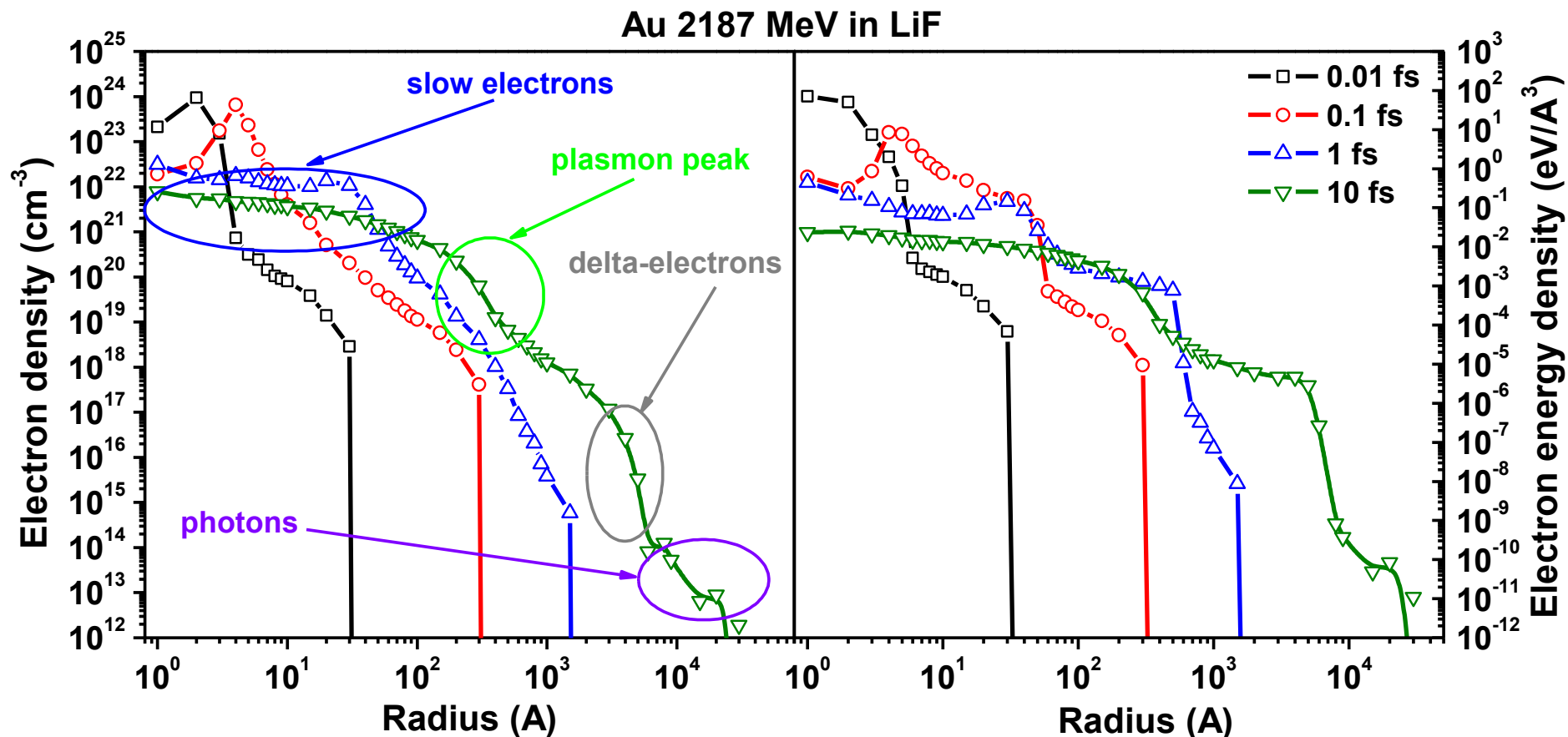


2 fs after 130 MeV Xe in Mg_2SiO_4

Pic: courtesy of R. Rymzhanov

<https://github.com/N-Medvedev/TREKIS-3>

Example: ion beam



Front waves: ballistic transport

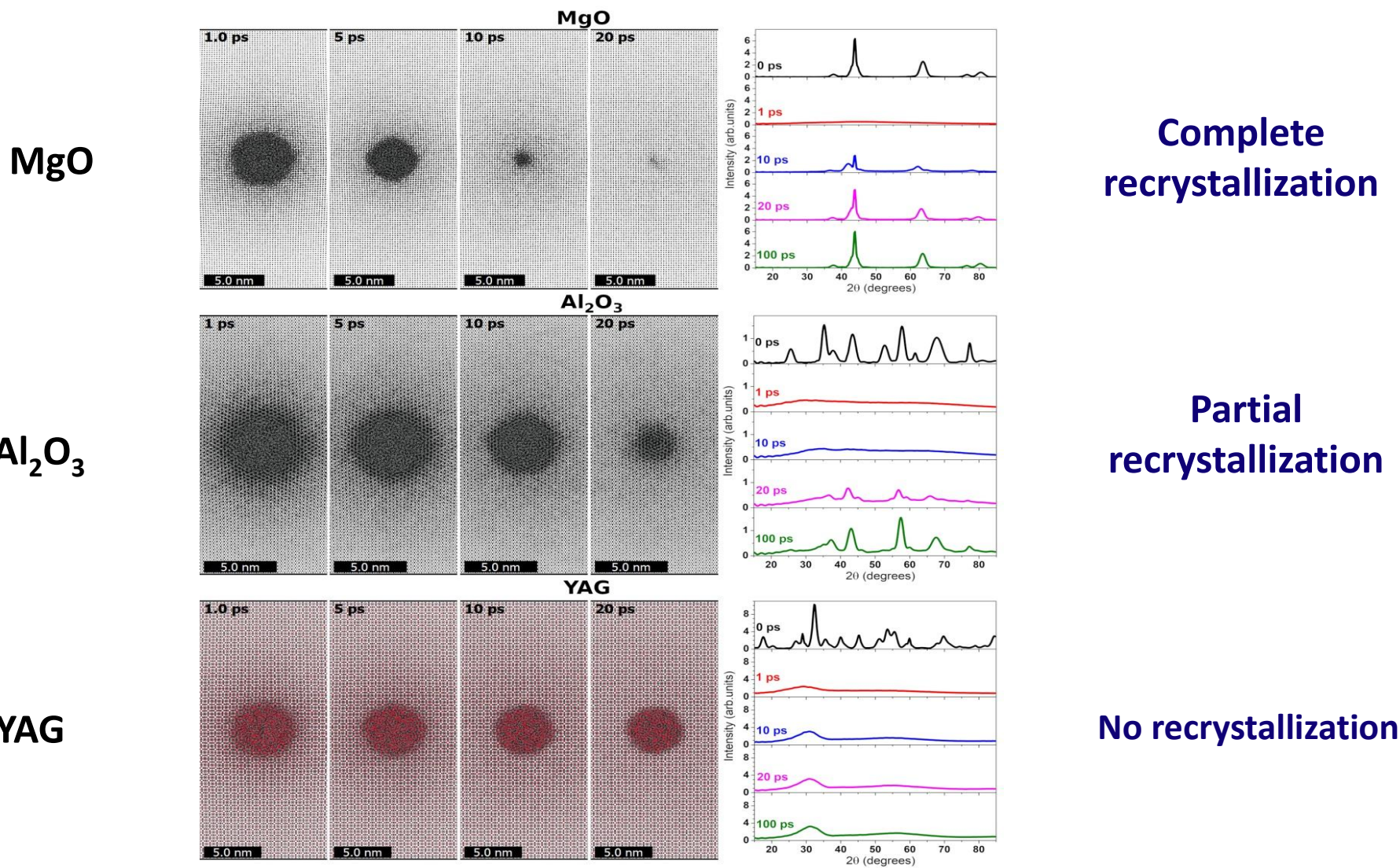
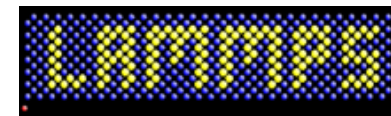
Slow electrons near track center: diffusive behavior



Example: ion beam

TREKIS³

output ⇒ input for

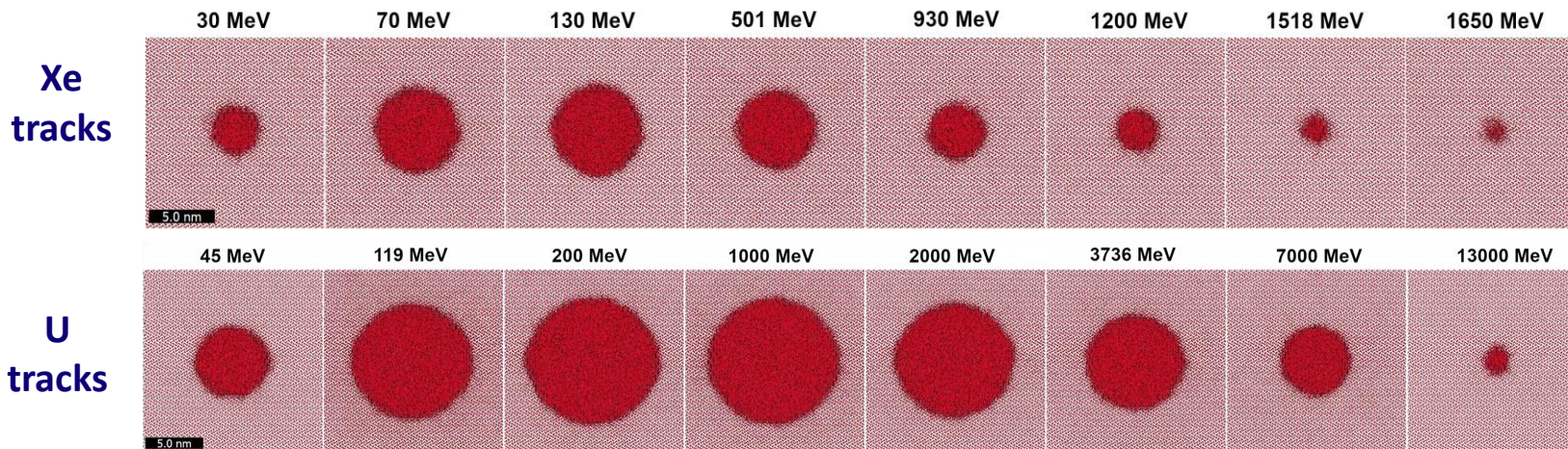
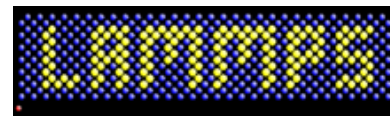


[Rymzhanov *et al.*, Scientific Reports 9, 3837 (2019)]

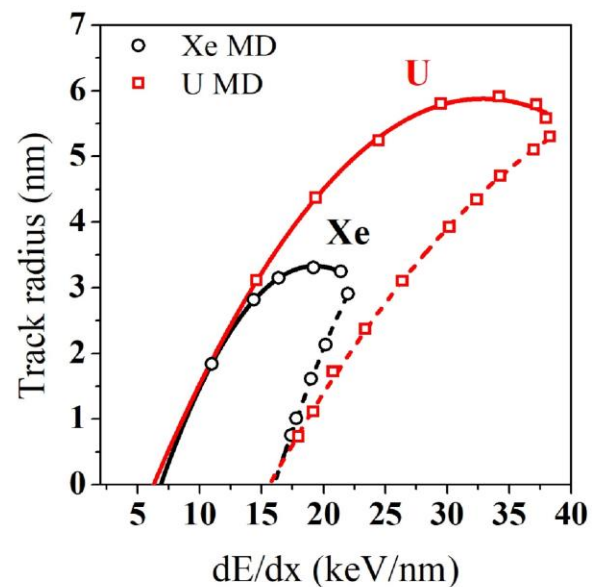
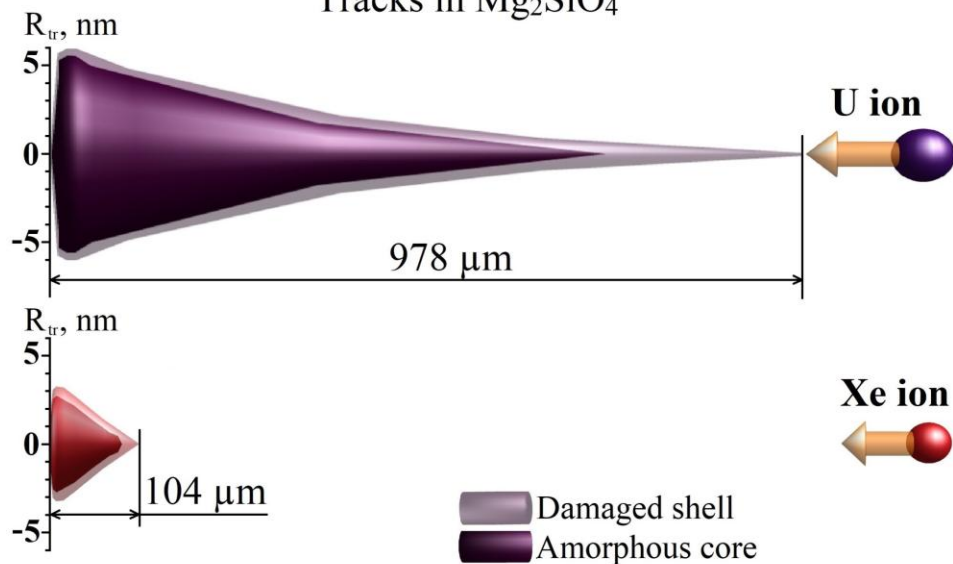
Example: ion beam



output \Rightarrow input for



Tracks in Mg_2SiO_4

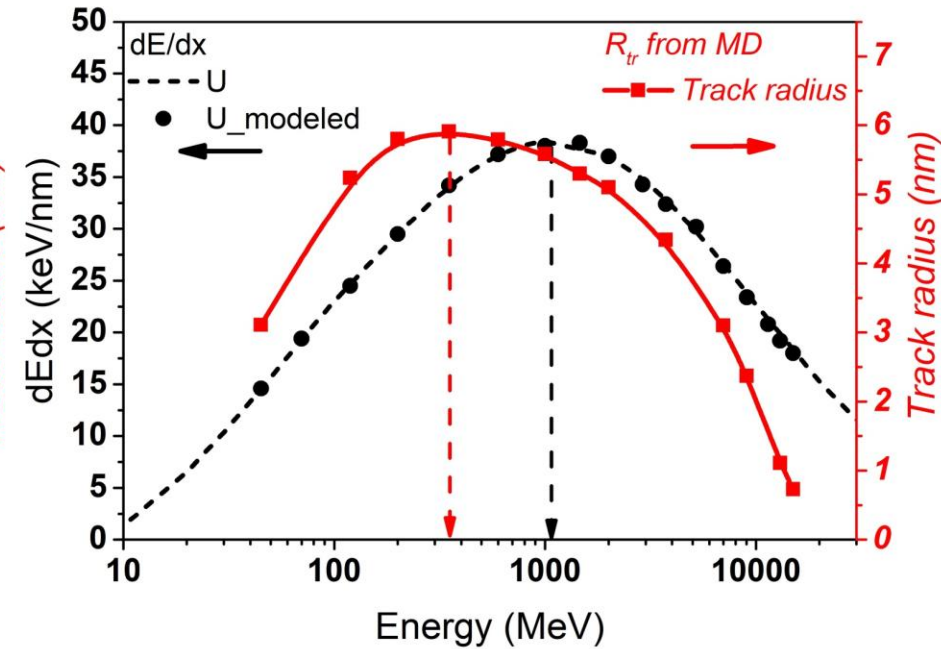
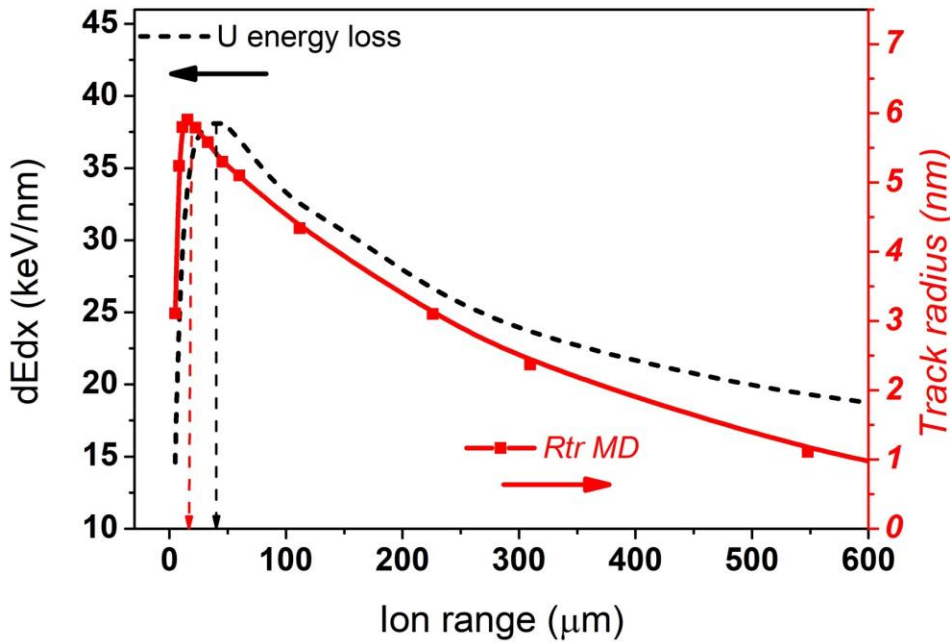
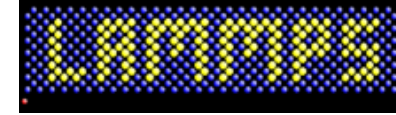


[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

Example: ion beam



output \Rightarrow input for



Maximal damage does not coincide with the maximal dose
Knowing the dose from MC is not sufficient to predict damage!

[Medvedev *et al.*, J. Appl. Phys. 133, 100701 (2023)]

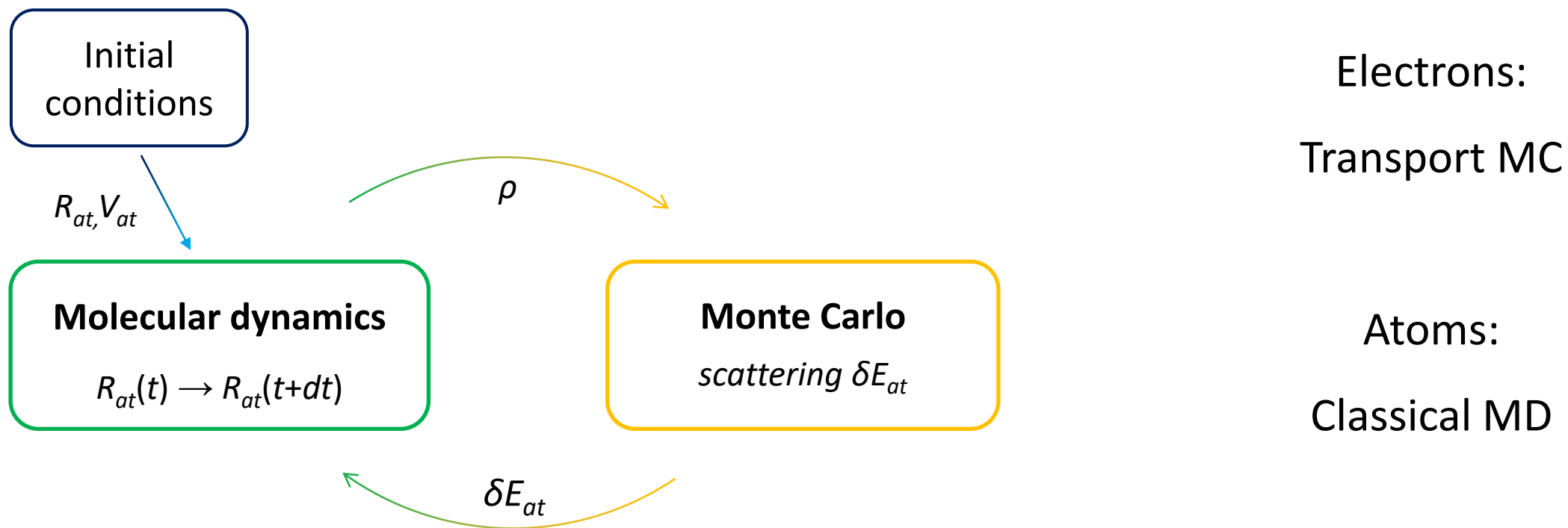
Example: Models with feedback

Example: ion beam



Time-Resolved Kinetics in Irradiated Solids

Generalization of TTM-MD to nonequilibrium electrons

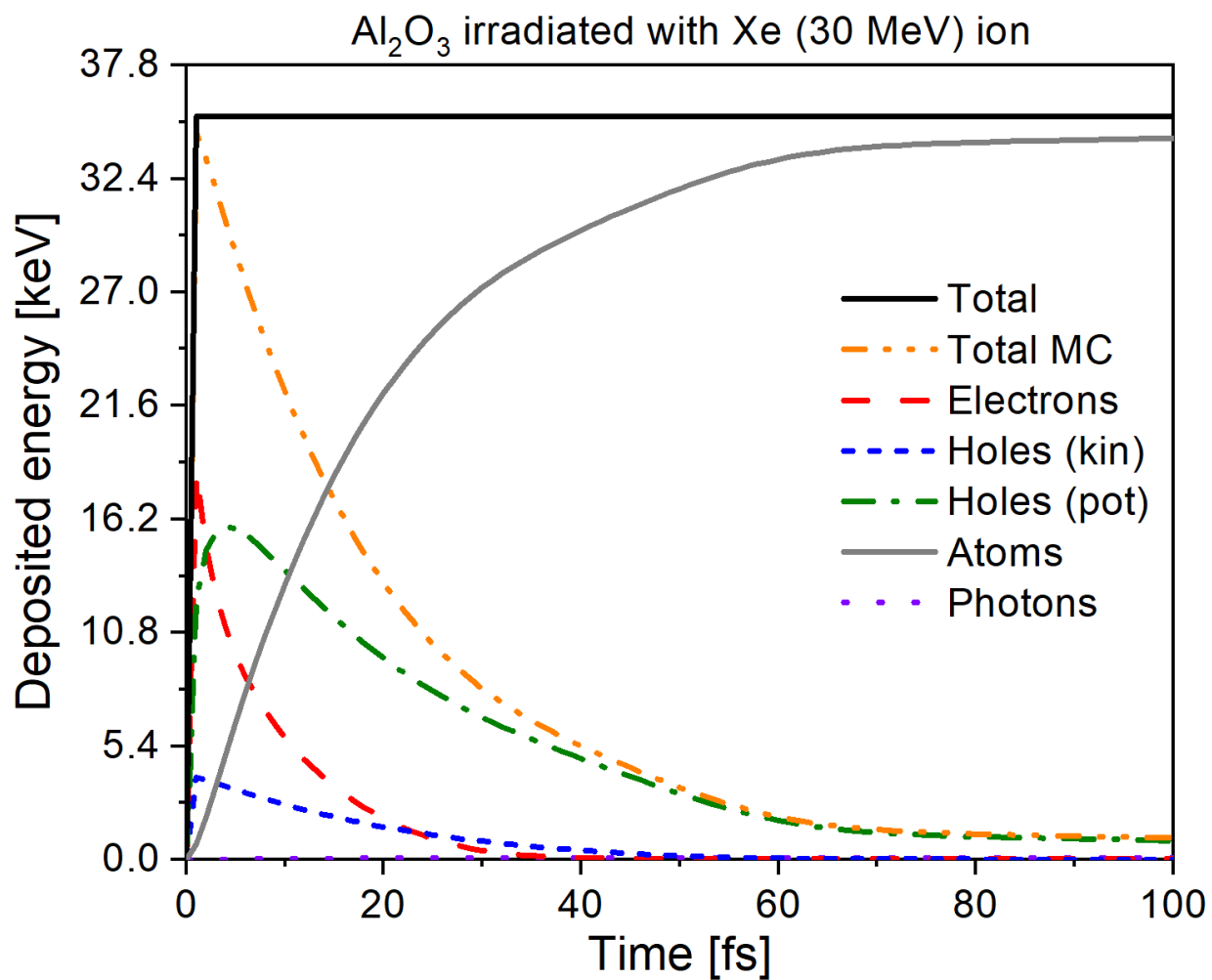


On-the-fly exchange of information between MC and MD

<https://github.com/N-Medvedev/TREKIS-4>



Example: ion beam

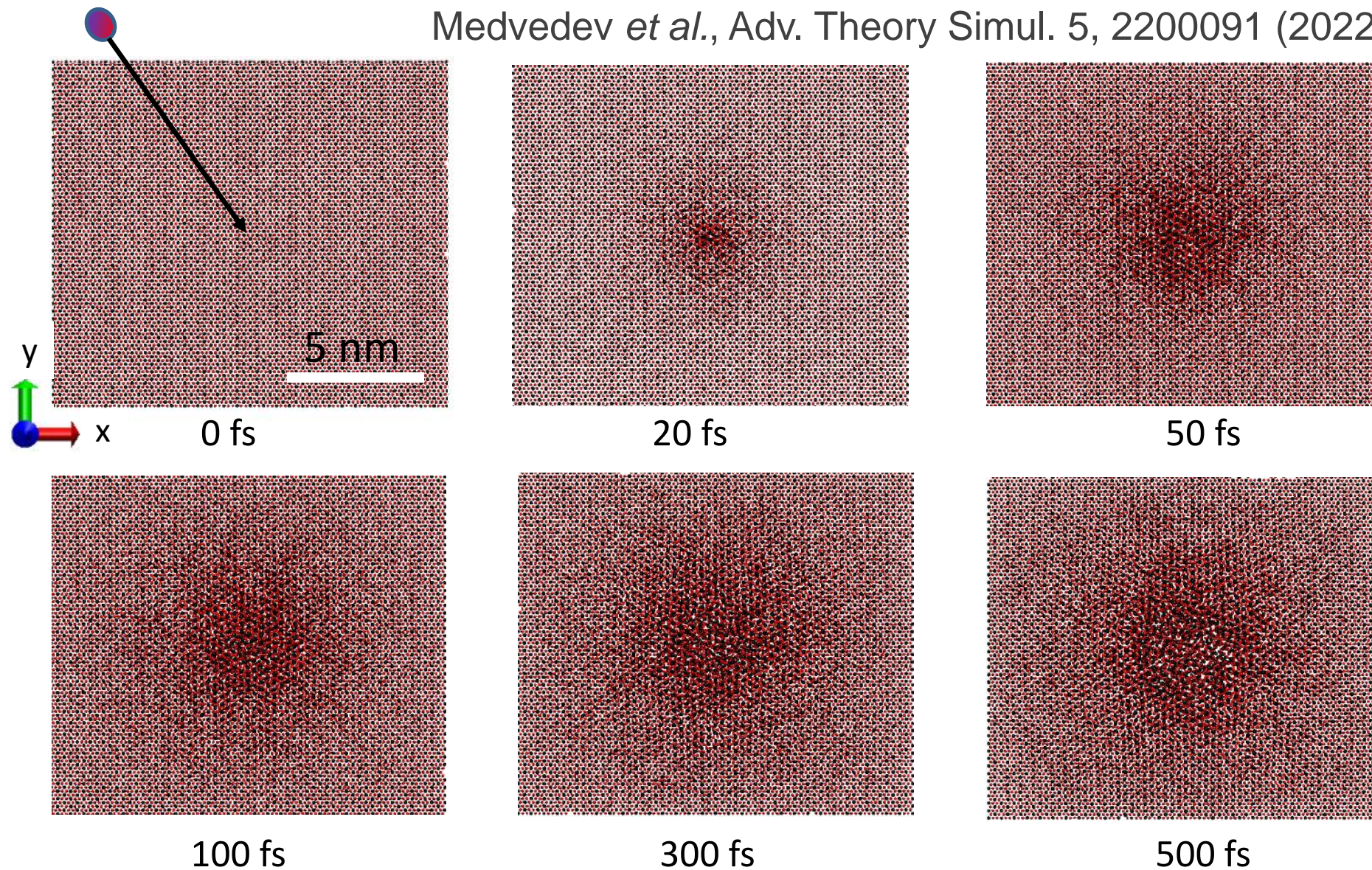
TREKIS⁴

On-the-fly exchange of information between MC and MD

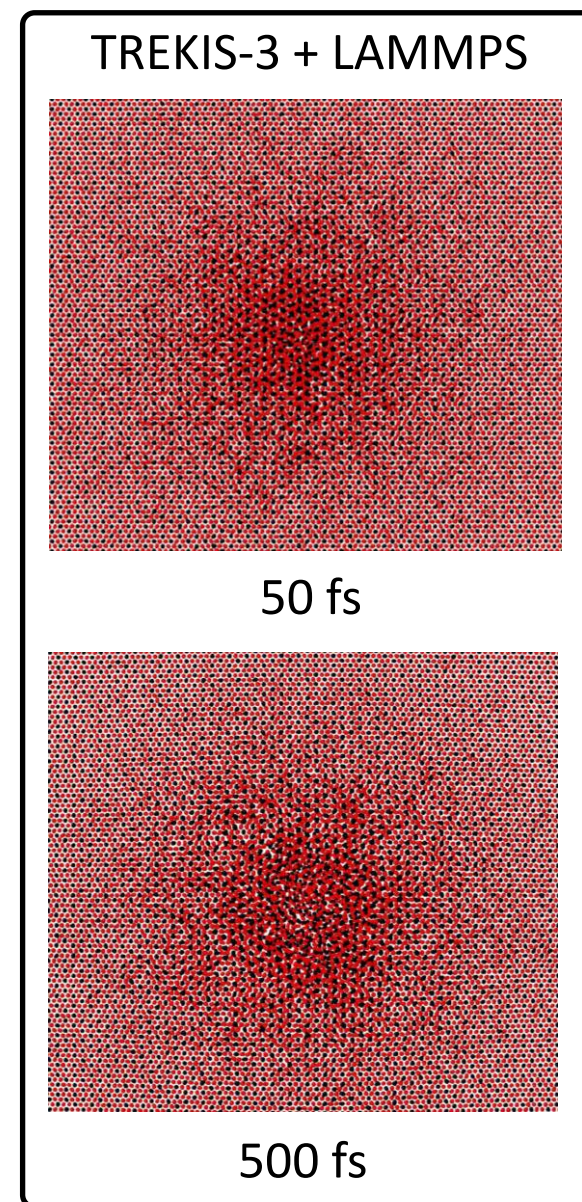
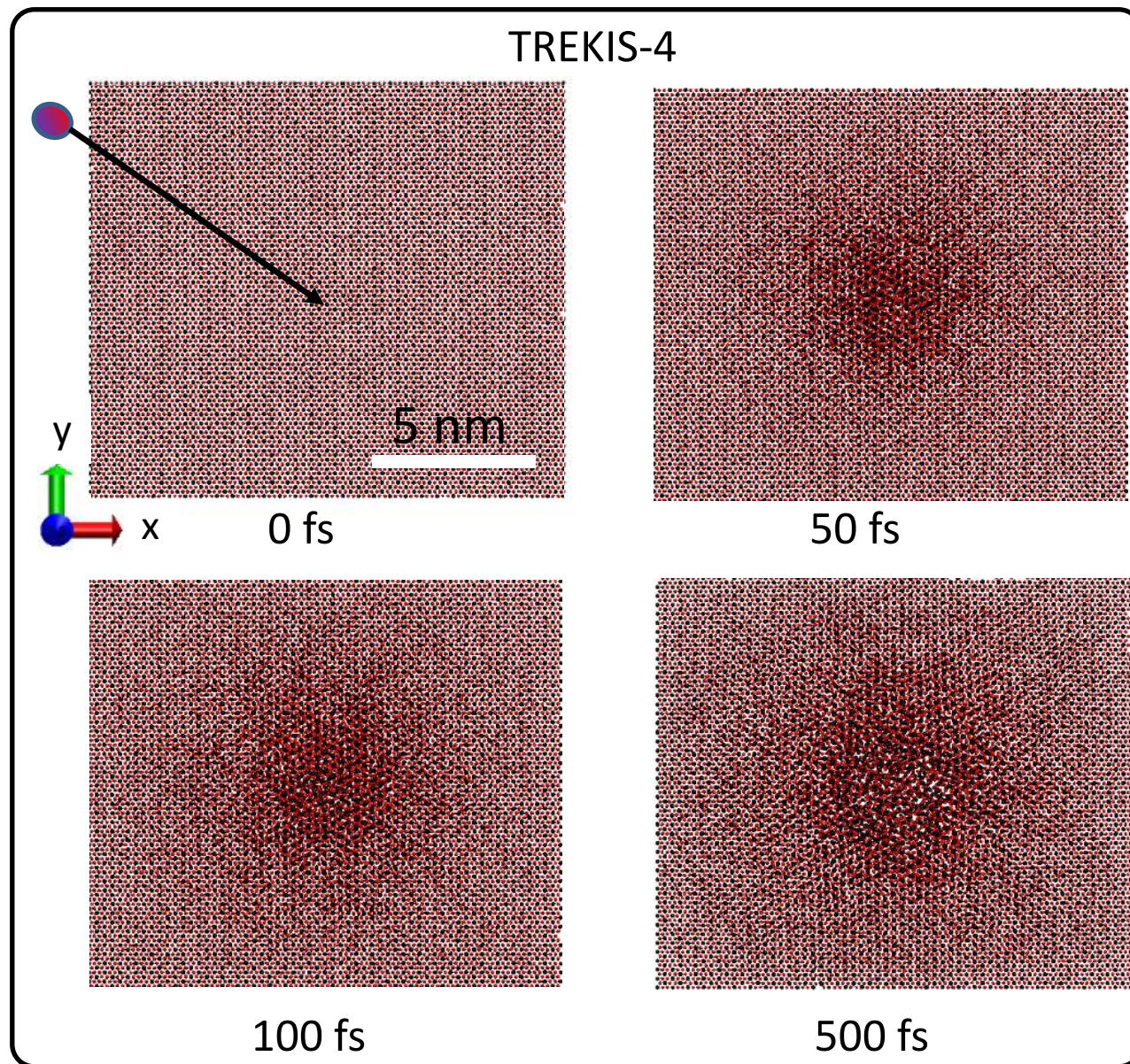
<https://github.com/N-Medvedev/TREKIS-4>

Example: ion beam

Medvedev *et al.*, Adv. Theory Simul. 5, 2200091 (2022)

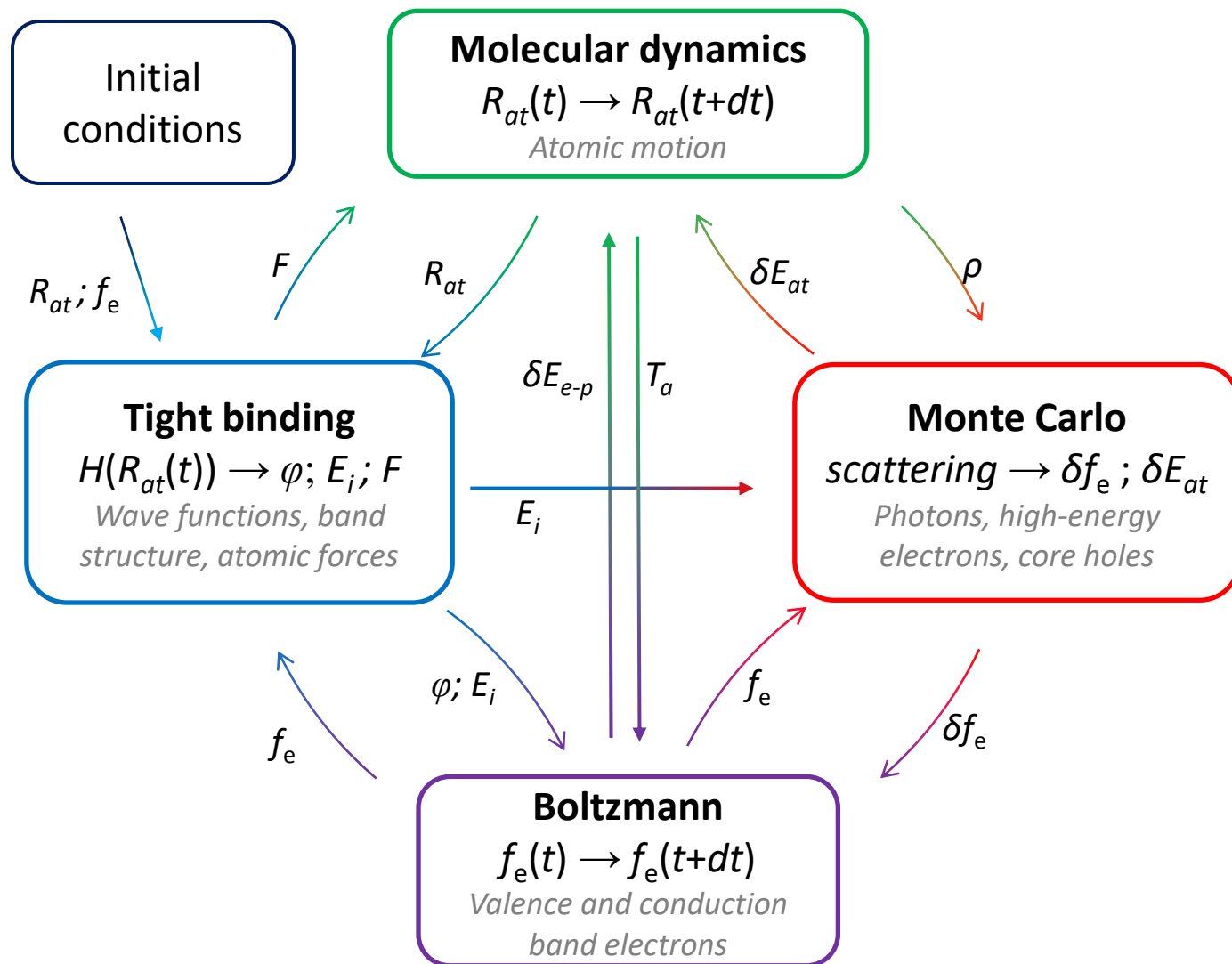


Ion beams: with vs. without feedback



Example: X-ray laser pulse

X-ray-induced Thermal And Nonthermal Transitions



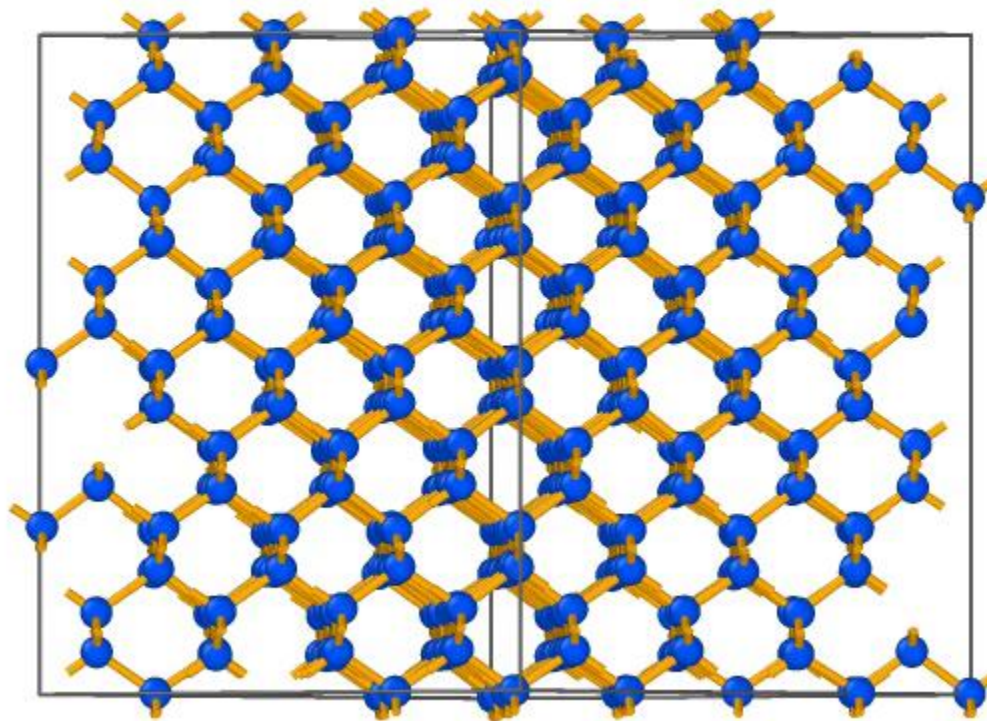
Includes :

- photon absorption
- excitation of electrons
- electron kinetics
- Auger-cascades of core holes
- equilibration of the electronic ensemble
- electron-ion (phonon) coupling
- evolution of band structure
- changes in the interatomic potential
- nonthermal melting
- atomic dynamics
- possible cooling (via thermostats)

Example: X-ray laser pulse

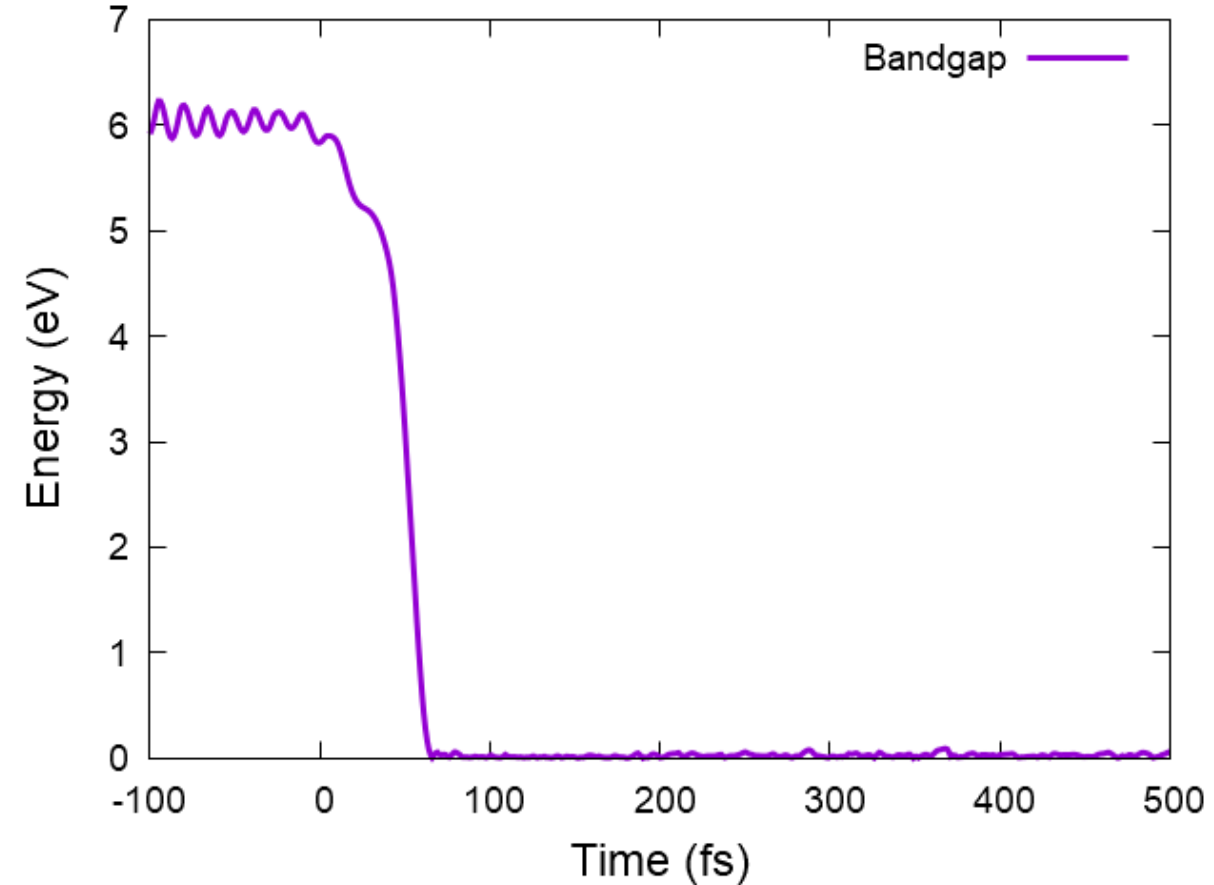
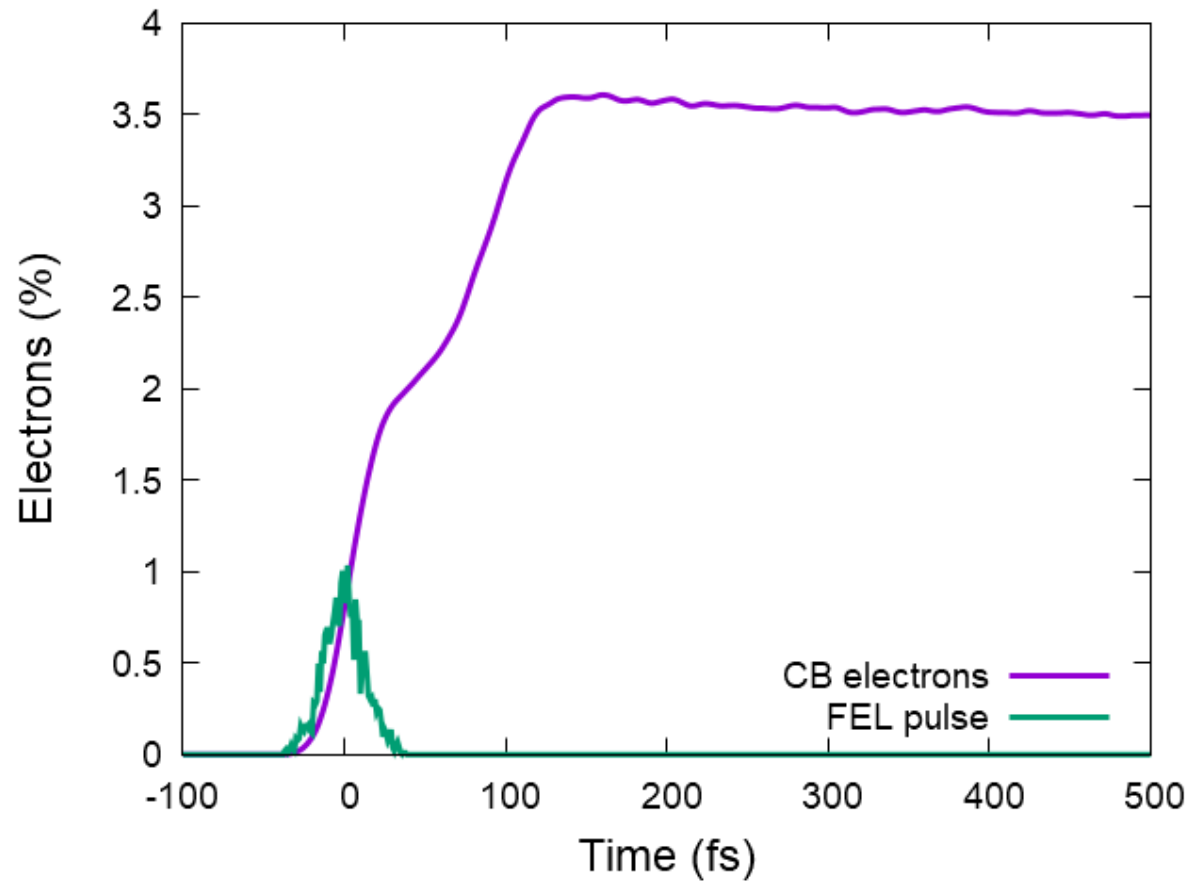


Diamond graphitization: irradiated with
10 fs pulse, 92 eV photons, 2 e V/atom dose



<https://github.com/N-Medvedev/XTANT-3>

Example: X-ray laser pulse

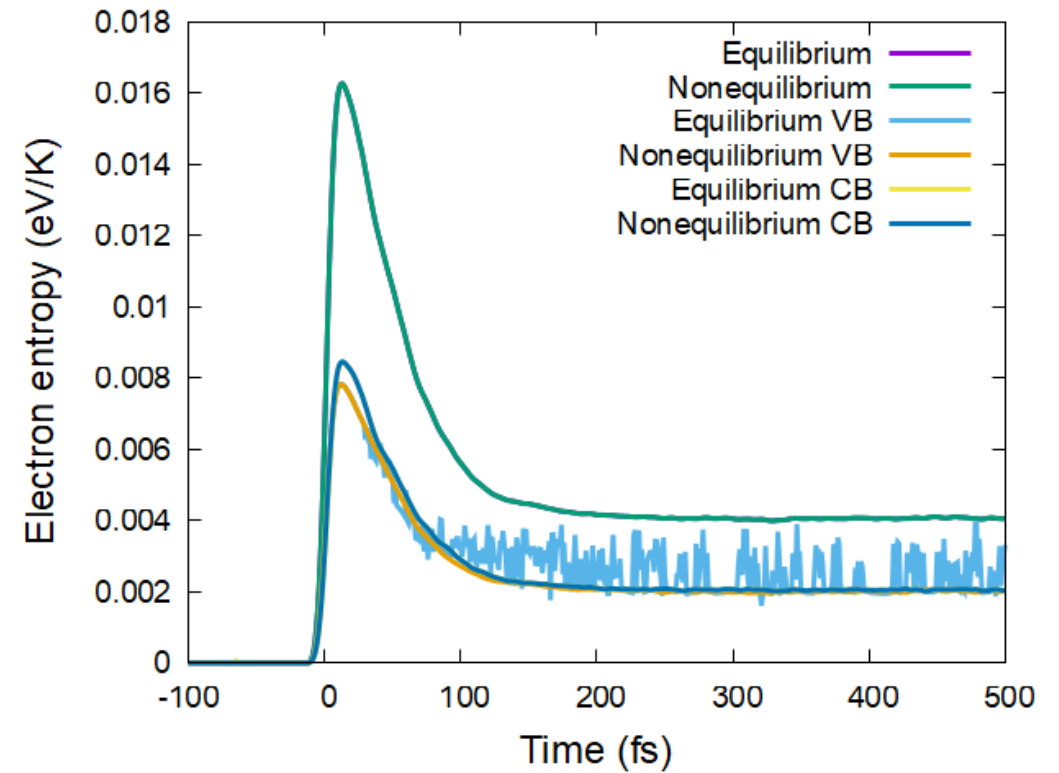
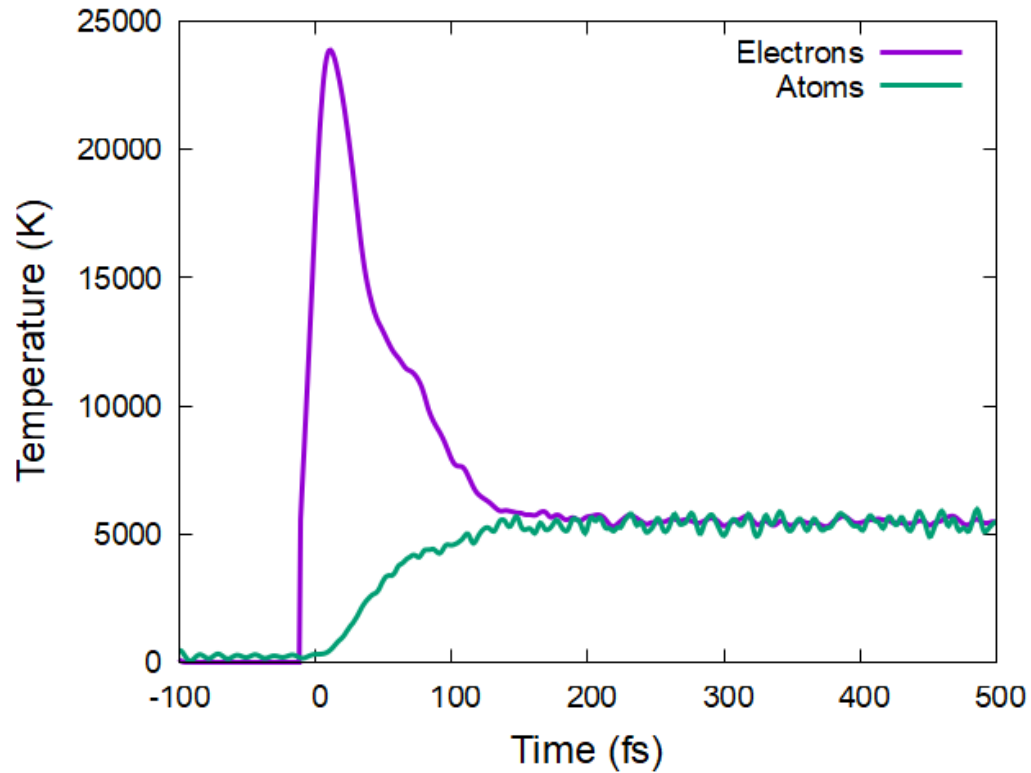


Electron density $\sim 1.5\%$ (dose ~ 0.7 eV/atom) \Rightarrow phase transition

Bandgap collapse induces ultrafast phase transition

N. Medvedev *et al.*, *4open* **1**, 3 (2018)

Example: X-ray laser pulse



Simultaneous kinetics of electrons and atoms, not possible to model without feedbacks!

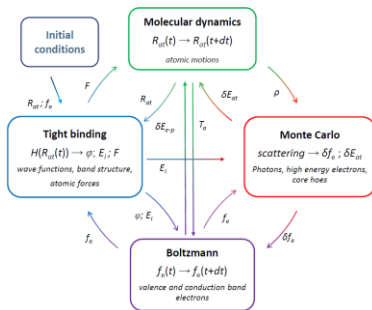
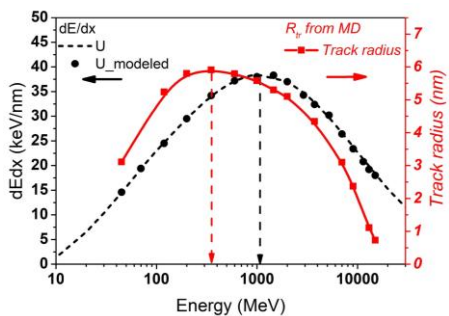
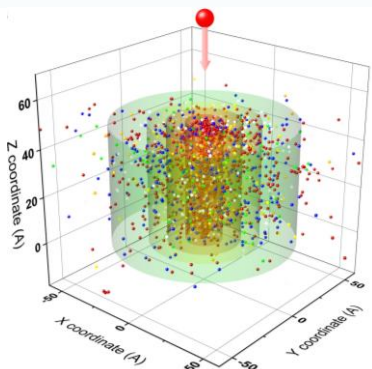
To sum up...

Transport Monte Carlo method describes evolution of the dose distribution

But! It is not sufficient to describe material damage

MC needs to be combined with other methods, describing material response (classical MD, *ab-initio* MD, etc.)

Hybrid models with feedback are needed; models without feedbacks occasionally work too...



Our codes



Time-Resolved Electron Kinetics in SHI-Irradiated Solids
<https://github.com/N-Medvedev/TREKIS-3>



Time-REsolved Kinetics in Irradiated Solids
<https://github.com/N-Medvedev/TREKIS-4>



X-ray-induced Thermal And Nonthermal Transitions
<https://github.com/N-Medvedev/XTANT-3>



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