



n-acoustic Solitary and Shock Dusty Plasmas



The Abdus Salam
International Centre
for Theoretical Physics



Num Prasad Acharya

Central Department of Physics, Tribhuvan University

Kirtipur, Kathmandu, Nepal

- Presence of micron sized dust grains is ubiquitous in naturally occurring and man-made plasma systems¹
- When the dust particles are exposed to the plasma environment, they are getting charged by collecting plasma particles hitting on their surfaces
- For typical lab plasma, $Z_d = 10^2 - 10^4$
- Plasma system maintains quasineutrality condition, i.e., $n_i = n_e + Z_d n_d$
- Plasma exhibits Debye shielding of the particle

$$\lambda_D = \frac{\lambda_{De} \lambda_{Di}}{\sqrt{\lambda_{De}^2 + \lambda_{Di}^2}}$$

¹(Shukla and Mamun, 2002)

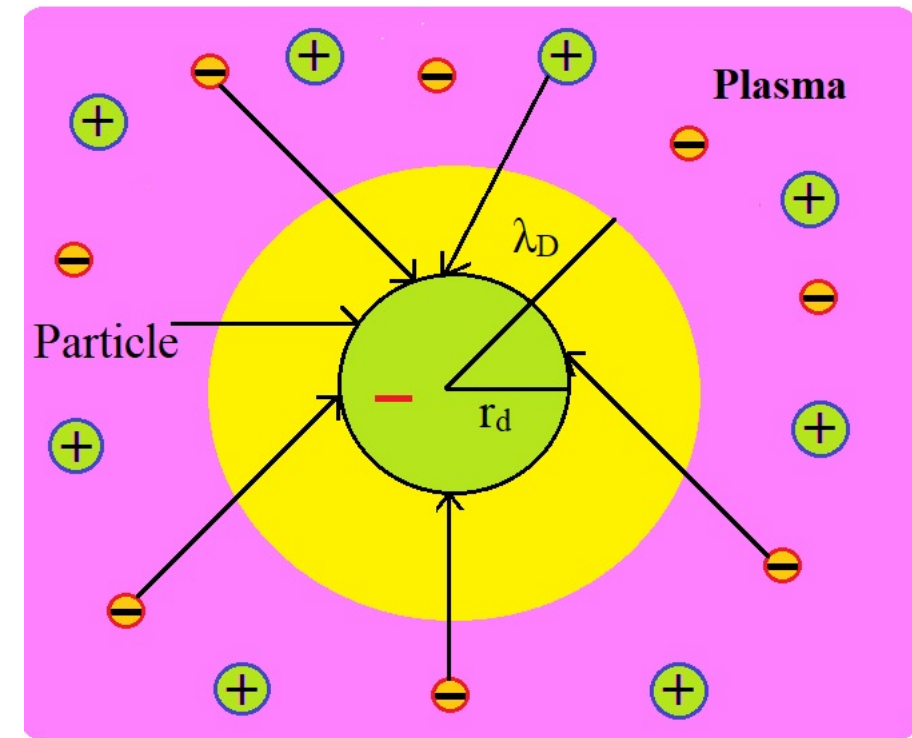


Figure: Schematic diagram of dusty plasma and dust charging process.

- ❖ Dust charge fluctuations affect linear and non-linear wave characteristics²
- ❖ In addition, events happening in cosmic environments such as planetary rings, dust wave instabilities in astrophysical plasmas and Earth's magnetosphere

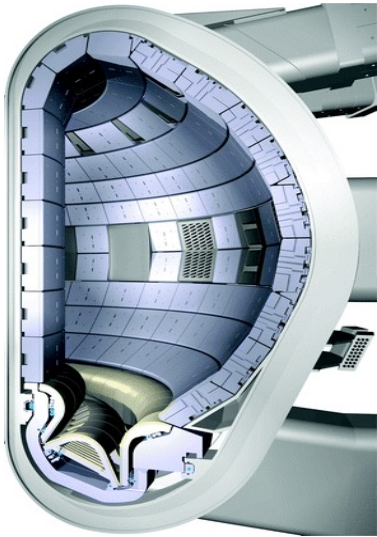


Figure: Divertor part of fusion device³

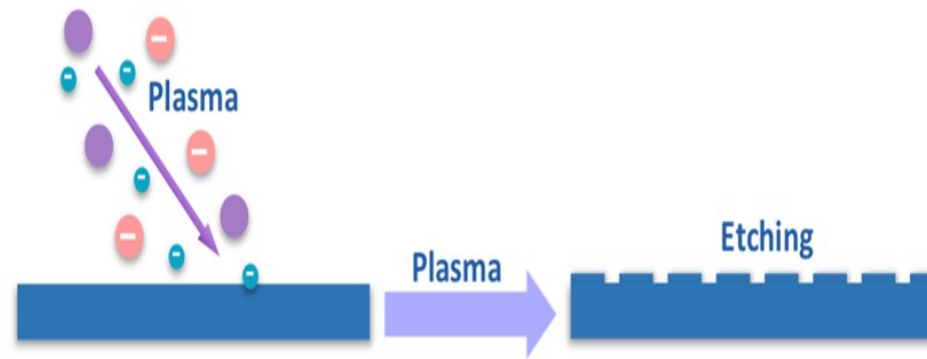


Figure: Plasma etching for semiconductor processing⁴

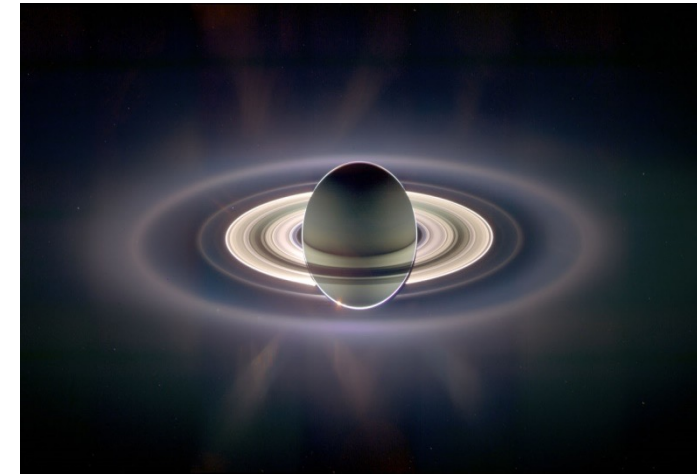


Figure: Saturn rings⁵

² (Das et al., 2017)

³(<https://f4e.europa.eu/mediacorner/newsview.aspx>) ⁴(<http://www.grinp.com/plasma/chemistry.html>)

⁵(<https://dornsife.usc.edu/news/stories/3063/a-brief-astronomical-history-of-saturns-amazing-rings>)

Motivation

Continuity equation: $\frac{\partial n_i}{\partial t} + \nabla \cdot (n_i \vec{v}_i) = S_i - S_l$

Momentum equation: $m_i n_i \left(\frac{\partial \vec{v}_i}{\partial t} + (\vec{v}_i \cdot \nabla) \vec{v}_i \right) = -en_i \nabla \phi + en_i (\vec{v}_i \times \vec{B}) - m_i n_i \nu_i \vec{v}_i - \nabla \cdot \vec{P}$

Poisson's equation: $\varepsilon_0 \nabla^2 \phi = en_e - en_i - q_d n_{d0}$

Quasineutrality condition: $en_{e0} = en_{i0} + q_{d0} n_{d0}$

Density distribution for electrons: $n_e = n_{e0} \left\{ 1 + (q-1) \frac{e\phi}{k_B T_e} \right\}^{(3q-1)/(2q-2)}$

Dust charge equation: $\frac{dq_d}{dt} = I_{Ti} + I_e$

Ion current: $I_{Ti} = \pi r_d^2 en_i \left(\frac{8k_B T_i}{\pi m_i} \right)^{\frac{1}{2}} \left[1 - \frac{eq_d}{r_d k_B T_i} + 0.4 \left(\frac{eq_d}{r_d k_B T_i} \right)^2 \frac{\lambda_D}{\lambda_{in}} \right]$

Electron current: $I_e = -\pi r_d^2 en_{e0} \left(\frac{8k_B T_e}{\pi m_e} \right)^{\frac{1}{2}} B_q \left[1 + (q-1) \frac{eq_d}{r_d k_B T_e} + (q-1) \frac{e\phi}{k_B T_e} \right]^{\frac{2q-1}{q-1}}$

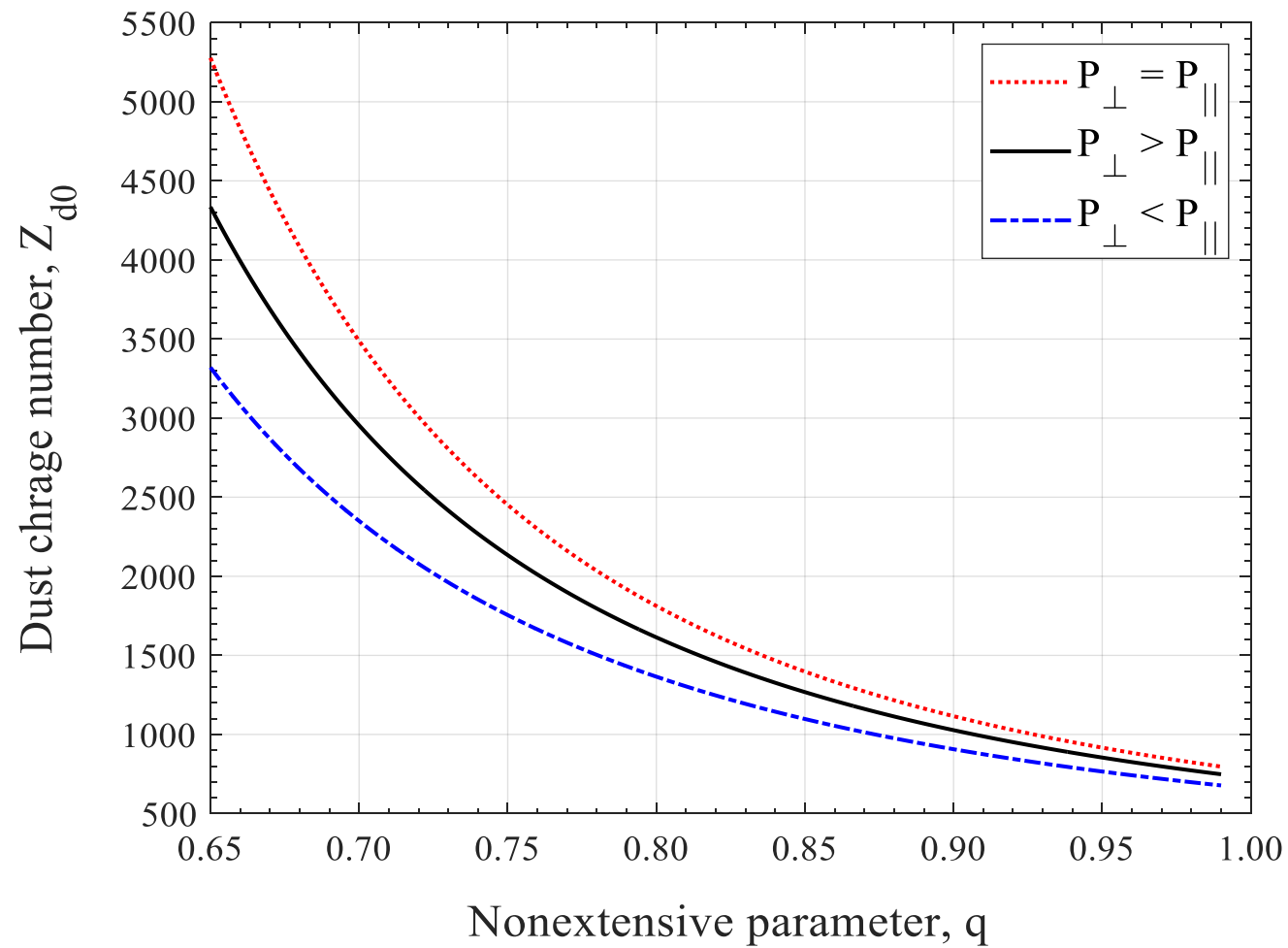


Figure: The variation of the equilibrium dust-charge number is shown against the nonextensive parameter q for different values of the anisotropic pressure: $P_{\perp} = P_{\parallel} = 0.5$, $P_{\perp}(0.3) > P_{\parallel}(0.1)$, and $P_{\perp}(0.1) < P_{\parallel}(0.3)$.

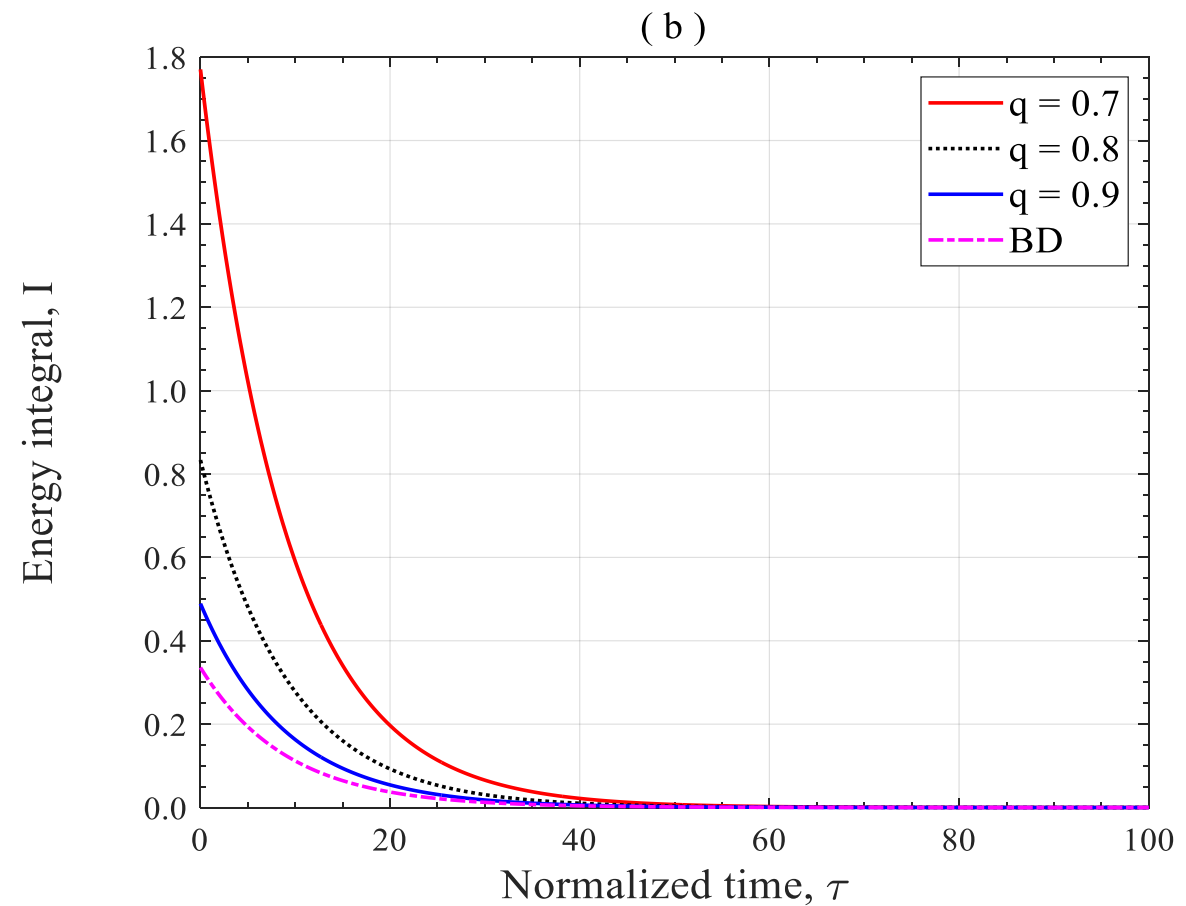
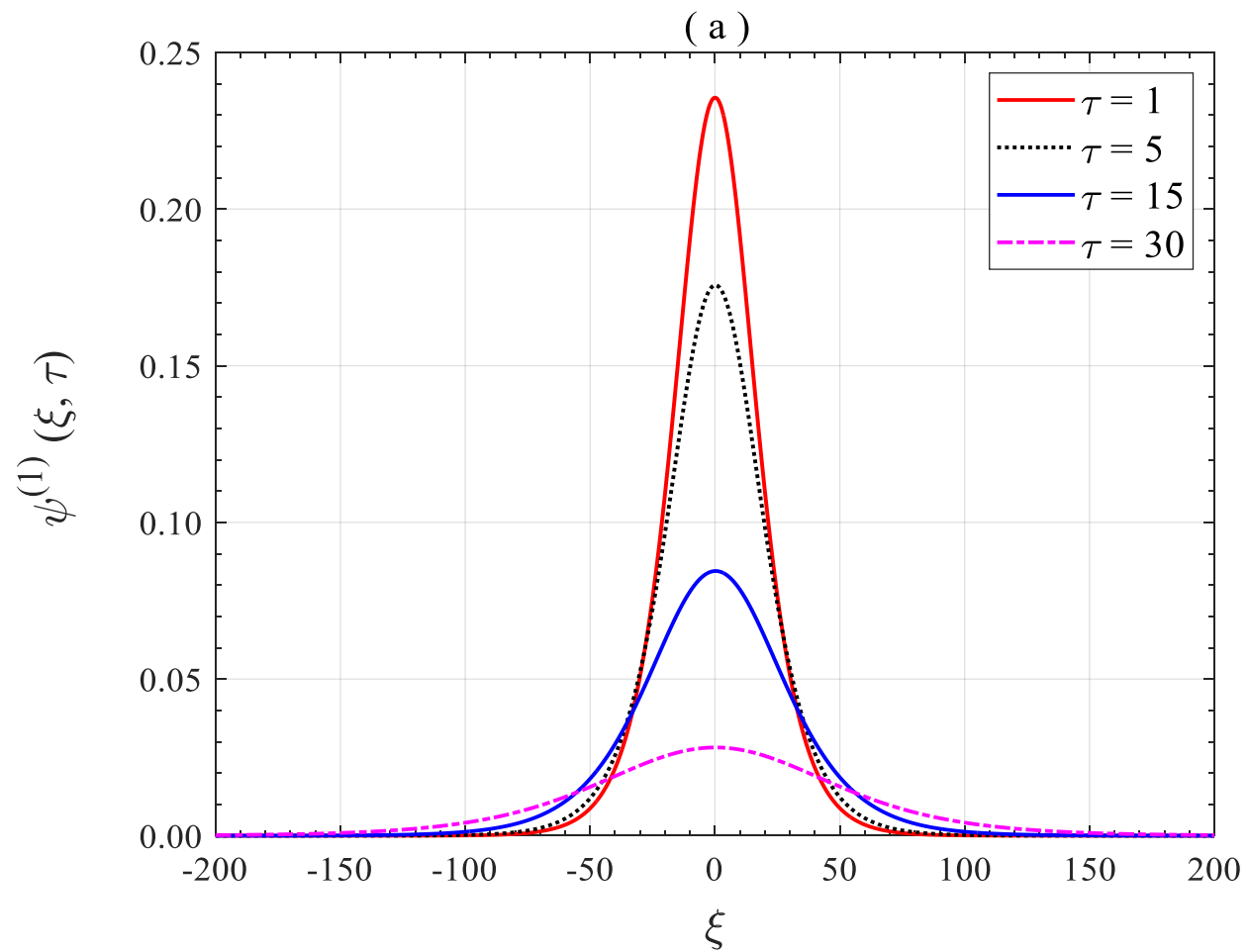


Figure: Profiles of the DIA soliton [subplot (a)] and the soliton energy [subplot (b)] are shown at different times (τ) and for different values of q as in the legends.

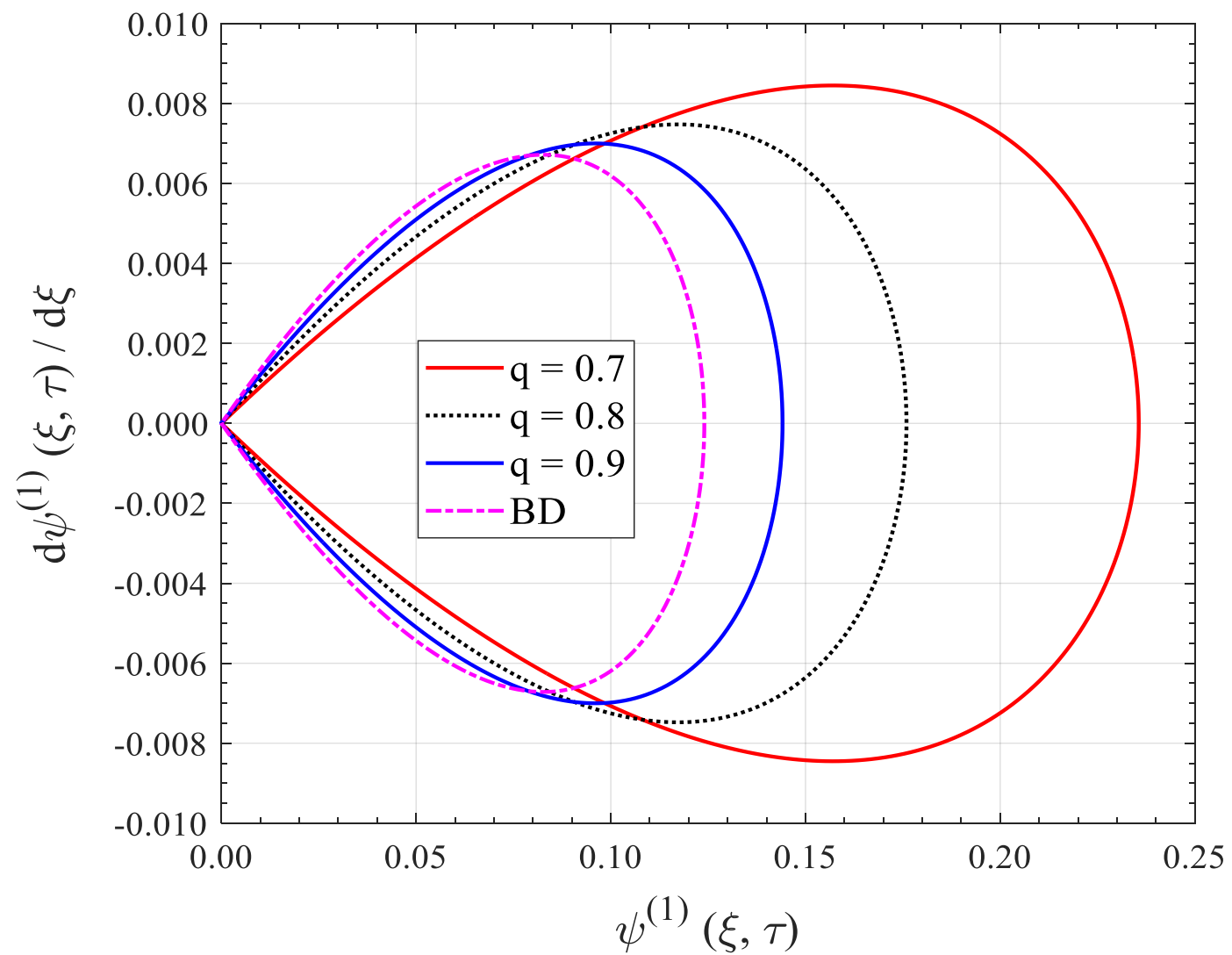


Figure: Phase portraits are shown for different values of the nonextensive parameter q .

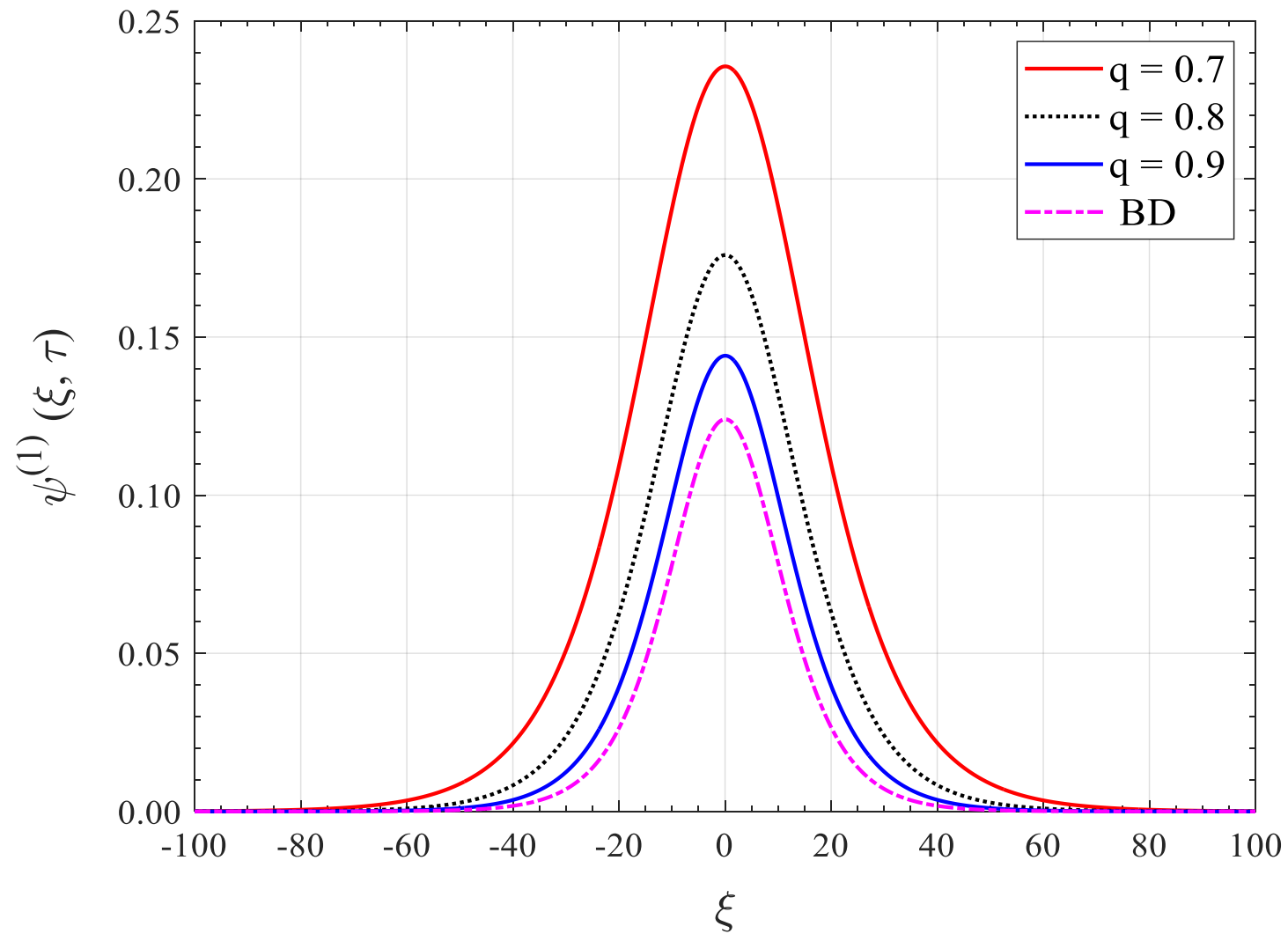


Fig. Profiles of the damped DIA soliton are shown for different values of the nonextensive parameter q .

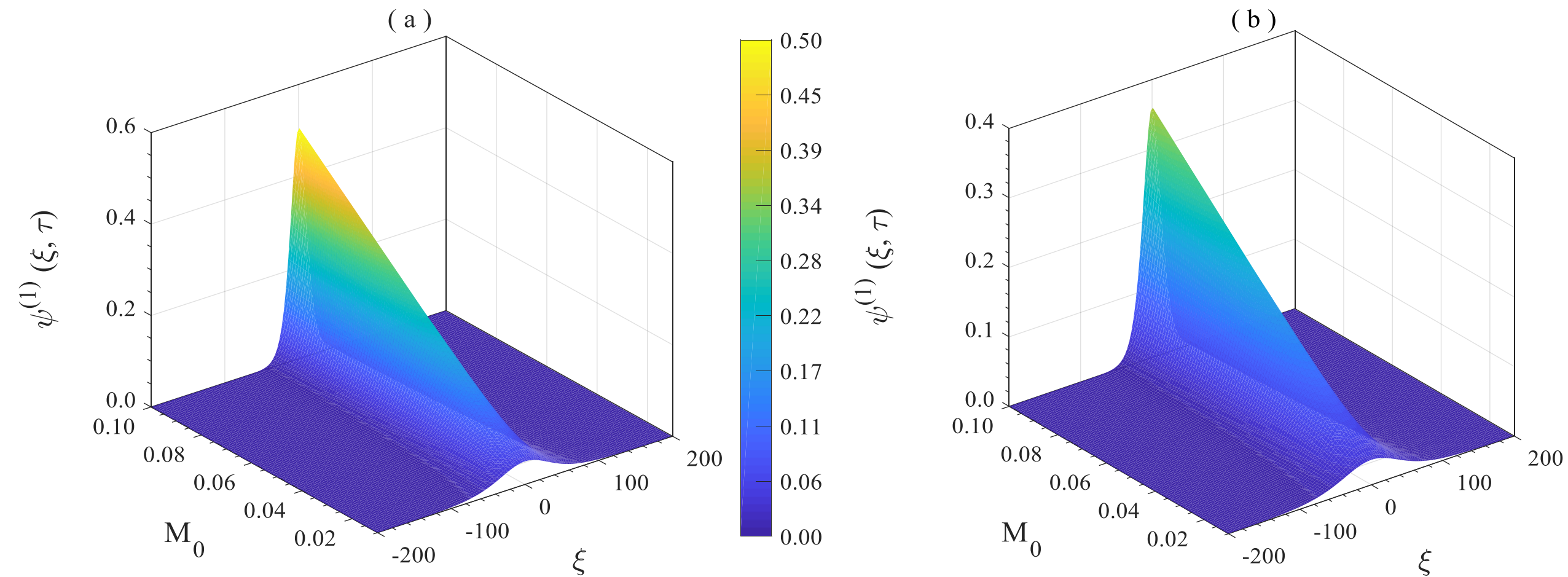


Fig. Profiles of the damped DIA solitons are shown in absence [subplot (a)] and presence [subplot (b)] of collision enhancement ion current.

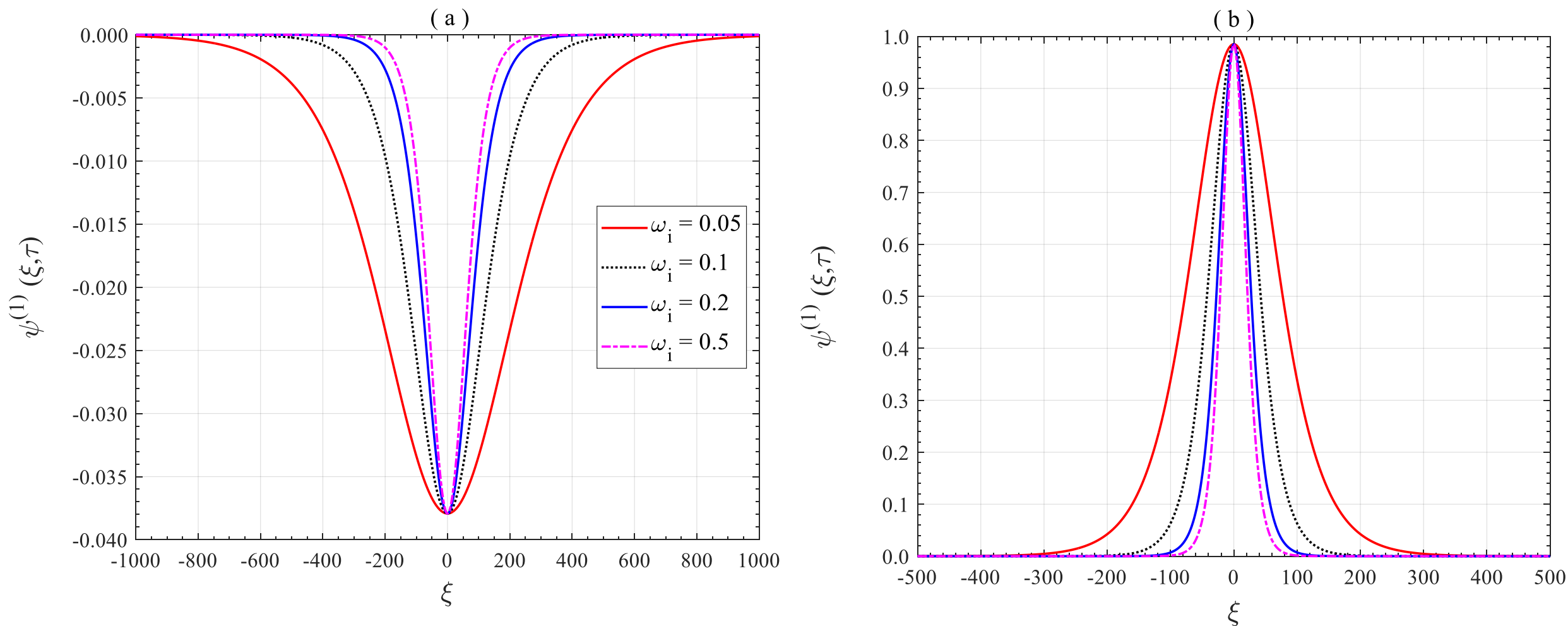


Figure: Profiles of the rarefactive [subplot (a) for $q=0.8$] and compressive [subplot (b) for $q=0.98$] DIA solitons are shown for different values of the ion gyrofrequency ω_i as in the legend.

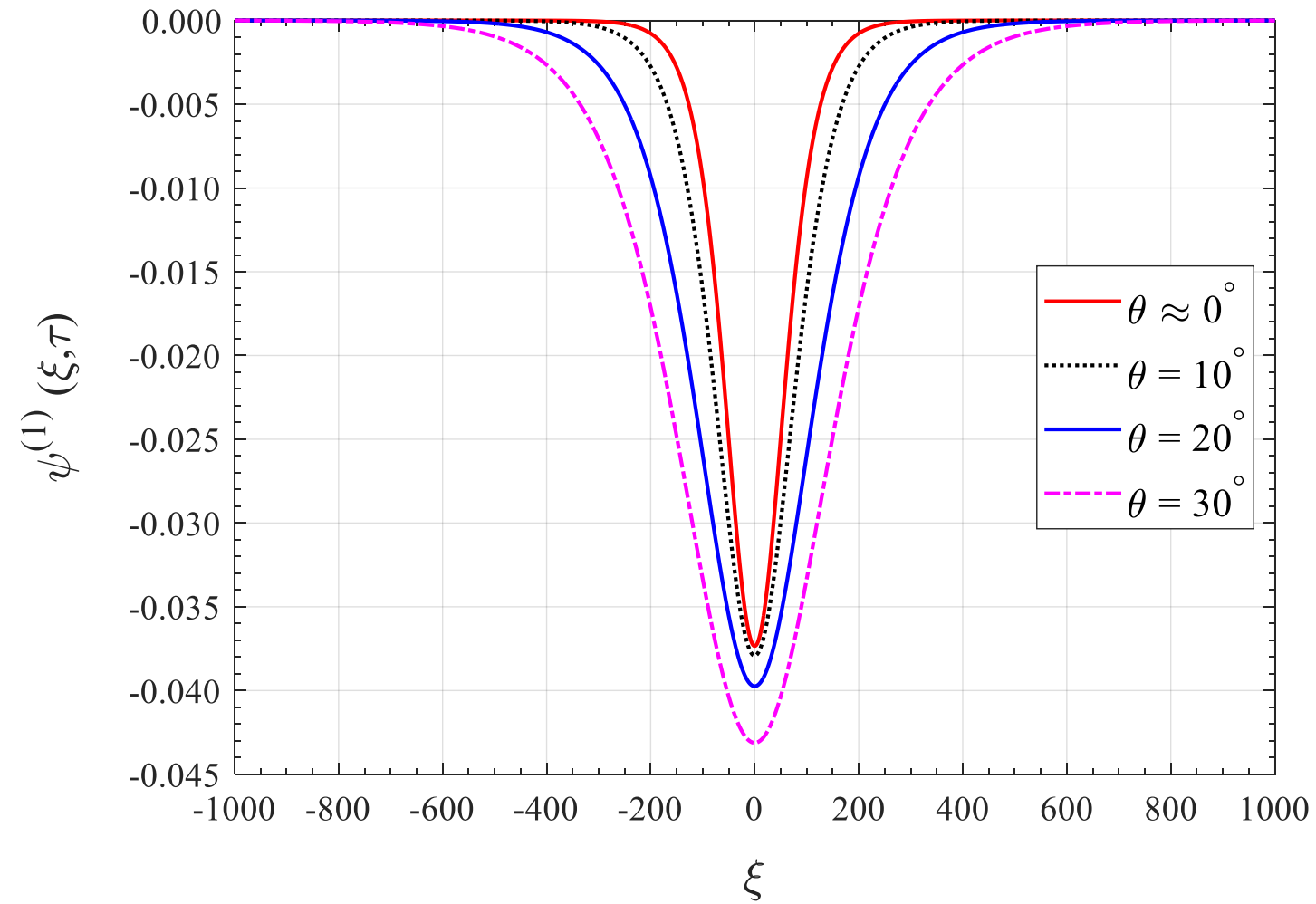


Figure: Effects of the obliqueness of wave propagation on the profiles of rarefactive DIA solitons are shown for different values of θ .

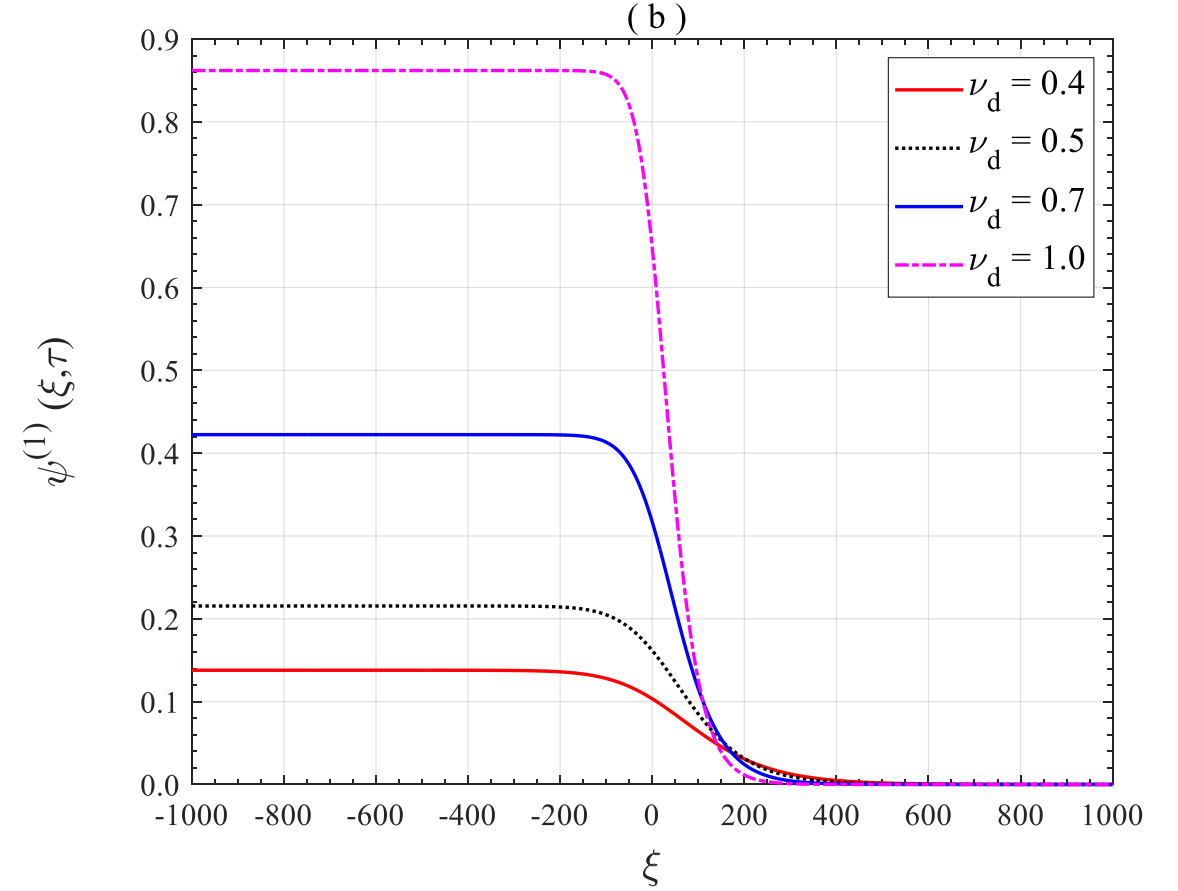
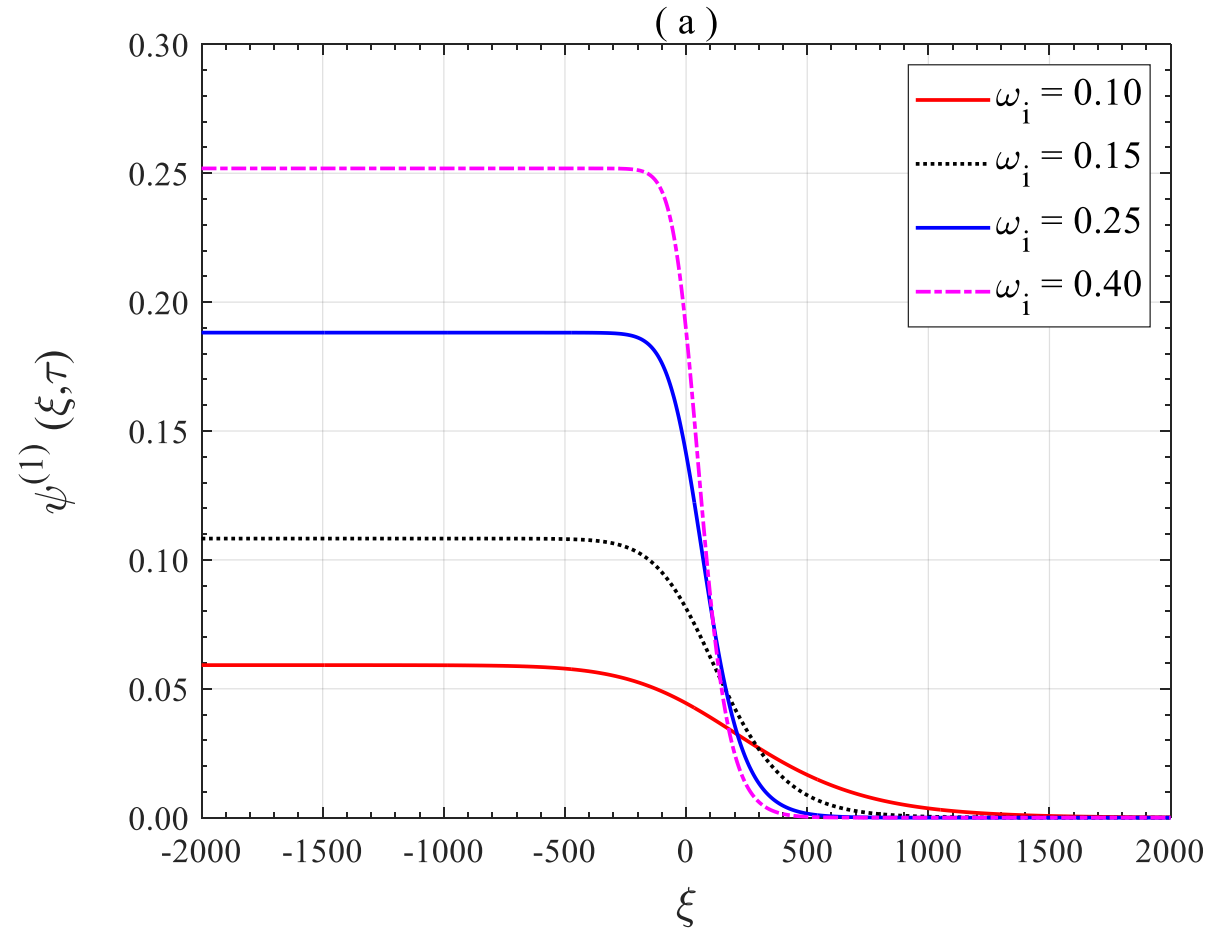


Figure: Monotonic shock profiles for different values of the (a) ion gyrofrequency, and (b) dust charge fluctuation rate with $q=0.8$.

Conclusion

- At equilibrium, the dust-charge number can achieve a maximum value in collisionless isotropic nonthermal dusty plasmas with super extensive electrons. The effects of the ion pressure anisotropy and the ion-neutral collision are to reduce the dust-charge number significantly.
- The DIA solitary waves get damped due to the effects of the adiabatic dust-charge variation and ion-dust and ion-neutral collisions.
- The increment of q-nonextensive parameter leads to the decrement of soliton amplitude and width,
- A deviation from the limit of adiabatic dust-charge variation, i.e., either (i) $v_{ch} < \omega_{pd}$ or (ii) $v_{ch} > \omega_{pd}$, relevant for space plasmas, can lead to the evolution of either (i) DIA damped solitary waves or (ii) DIA shocks.
- The qualitative features of the damped DIA solitons by the effects of the magnetic field (ω_i) and the propagation angle θ remain similar to those for laboratory plasmas.
- A transition of DIA shocks from rarefactive to compressive types can occur as we approach from nonthermal states (with superextensive electrons) to thermal ones (with Boltzmann distributed electrons).
- It is noted that the effects of the magnetic field (ω_i) and the dust-charge fluctuation rate (v_d) on the profiles of DIA shocks are similar.

References

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