

Cosmic Inflation

[ICTP summer school 2026]

classical / background : rapid expansion of space time $\rightarrow$ homogeneity on large scales

(quantum) fluctuations : quantum fluctuation become stretched $\rightarrow$ classical $\rightarrow$ seeds for inhomogeneities (CMB, Galaxies)

↓ lectures by Marko & Matthew

Plan

- 1) Motivation: Big Bang Puzzles
- 2) Classical background dynamics: slow roll inflation
- 3) Quantum fluctuations I: density perturbation
- 4) Quantum fluctuations II: gravitational waves
- 5) Tests of cosmic inflation

Literature

- Daniel Baumann, 0907.5424, TASI lectures, "Inflation" (+ "Advanced Cosmology")
- Scott Dodelson "Modern Cosmology"
- [• Julien Lesgourgues, 1302.4640, "Cosmological perturbations"
- Michele Maggiore, gr-qc/9909001, "Gravitational Waves, Vol I"

1) FRW Cosmology

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G T_{\mu\nu}$$

metric ↔ matter

Einstein's equations

↪ homogeneity & isotropy ↪

$$ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - kr^2} + r^2 d\Omega^2 \right]$$

$$T^\mu_\nu = \text{diag}(\rho, -p, -p, -p)$$

perfect fluid

↪ ↪

Friedmann equations:

$$[H_p = 1]$$

$$(\mp 1) \left(\frac{\dot{a}}{a} \right)^2 = H^2 = \frac{1}{3} \rho - \frac{k}{a^2}$$

$$\dot{} = \frac{d}{dt}$$

$$(\mp 2) \left(\frac{\ddot{a}}{a} \right) = \dot{H} + H^2 = -\frac{1}{6} (\rho + 3p)$$

$$\hookrightarrow \dot{\rho} = -3H(\rho + p)$$

$\omega \equiv p/\rho$

continuity equation

equation of state parameter

$$\Gamma(\mp 1) \rightarrow \dot{\rho} = \frac{d}{dt} \left(3H^2 + 3k/a^2 \right)$$

$$= 6H\dot{H} - 2 \frac{3k}{a^2} H$$

$$\dot{H} = -\frac{1}{6}(\rho + 3p) - \underbrace{H^2}_{\frac{1}{3}\rho - \frac{k}{a^2}}$$

$$= H \left[-\rho - 3p - 2\rho + \frac{6k}{a^2} - \frac{6k}{a^2} \right]$$

$$= -3H(\rho + p)$$

□]

$$\left[\begin{aligned} \rightarrow \frac{d\rho}{da} &= \frac{d\rho}{dt} \frac{1}{\dot{a}} = - \frac{3H(\rho+p)}{\dot{a}} = - \frac{3\rho}{a} (1+\omega) \\ \rightarrow \frac{d \ln \rho}{d \ln a} &= -3(1+\omega) \end{aligned} \right] \text{ is covered by MoM}$$

$$\rightarrow \rho \propto a^{-3(1+\omega)}$$

$\omega=0$ ($p=0$)	$\rightarrow \rho = a^{-3}$	matter
$\omega = 1/3$	$\rightarrow \rho = a^{-4}$	radiation
$\omega = -1$	$\rightarrow \rho = \text{const.}$	cc / inflation

"hot big bang": $k=0, \omega=0, 1/3$

$$\hookrightarrow a(t) \propto t^{\frac{2}{3(1+\omega)}}$$

$\rightarrow$ singularity "big bang" at $a(t)=0$

2) Horizon problem

Why is the CMB so isotropic?

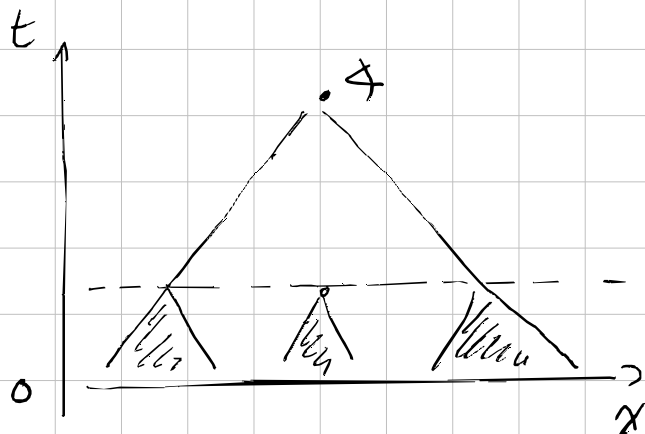
past light-cone, region of causal contact

$$\chi = \int_0^t \frac{dt'}{a(t')} \underset{H = \frac{\dot{a}}{a}}{=} \int_0^a \frac{da}{Ha^2} = \int_0^a da \underbrace{a}_{1/\dot{a}} \underbrace{H}_{\text{comoving Hubble horizon}}$$

$$(F2): \ddot{a} = -\frac{1}{6} a g(1+3w) < 0 \quad \text{for } w > -1/3$$

$$\rightarrow t \uparrow \quad \dot{a} \downarrow \quad \frac{1}{\dot{a}} = \frac{1}{aH} \uparrow \quad \text{for } w > -1/3$$

$\rightarrow$ fraction of universe in causal contact increases with time



CMB decoupling

to "singularity"

CMB as observed today $\approx 10^5$ regions which were never in causal contact

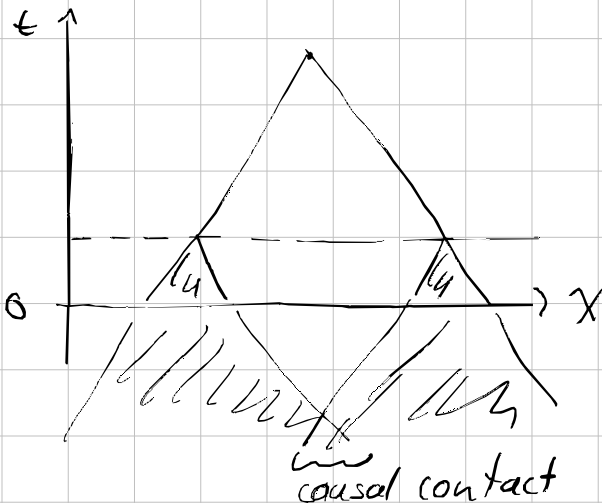
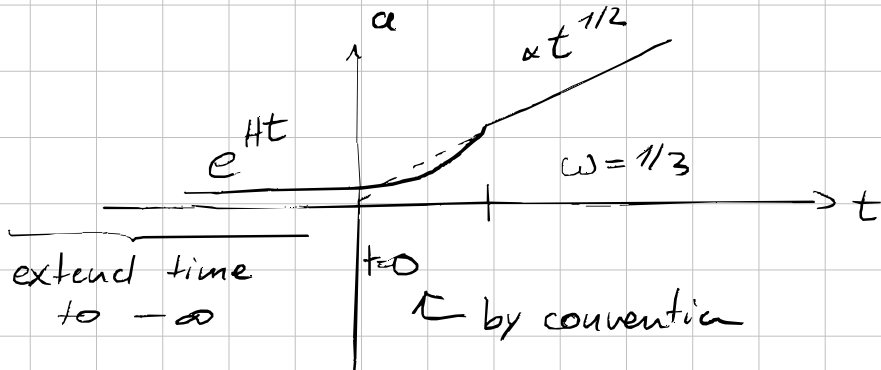
$\rightarrow$ why do they all have the same T ?

resolution

need $\omega < -1/3$: $1/aH$ decreases w. time,
later observer sees "less" than early observer

e.g. $\omega = -1 \rightarrow \rho \propto a^{-3(1+\omega)} = \text{const} \quad / \quad \rho_c = 0$

$$\rightarrow \frac{\dot{a}}{a} = H = \text{const} \rightarrow a \propto e^{Ht}$$



3) Flatness problem

$$(F1) \rightarrow 1 - \underbrace{\frac{8(a)}{3H^2(a)}}_{\equiv \Omega(a)} = -k \underbrace{\frac{1}{(aH)^2}}_{\substack{\text{grows w. time} \\ \text{for } w > -1/3}}$$

$$k = \{-1, 0, 1\}$$

Today, $\Omega(a_0) \simeq 1$ (observed)

$\hookrightarrow$ unstable point for $k \neq 0$

$\hookrightarrow \Omega(a_{\text{BBN}}) - 1 \leq \mathcal{O}(10^{-16})$

$\rightarrow$ why is the universe so flat?
 $k=0$ or fine-tuning?

Also solved by decreasing comoving horizon in the past, i.e. $w < -1/3$.

- $\downarrow$
- negative pressure?
 - how do we achieve this?

Bonus: also explains characteristic structure of CMB fluctuations.

4) slow-roll inflation

• Consider real scalar field ϕ

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{2} R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right)$$

$$= S_{EH} + S_\phi$$

$$\rightarrow T_{\mu\nu}^{(\phi)} = -\frac{2}{\sqrt{-g}} \frac{\delta S_\phi}{\delta g^{\mu\nu}} = \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left(\frac{1}{2} \partial^\sigma \phi \partial_\sigma \phi - V(\phi) \right)$$

$(-, +, +, +)$ $\underbrace{\quad}_{-\partial_0 \phi \partial_0 \phi + \partial_i \phi \partial_i \phi}$

Assume $\phi(\vec{x}, t) = \phi(t)$ homogeneous, $\partial_i \phi = 0$

$$\hookrightarrow T_{\mu\nu}^{(\phi)} = \text{diag}(S, -p, -p, -p) \quad \text{with}$$

$$S_\phi = \frac{1}{2} \dot{\phi}^2 + V(\phi), \quad p_\phi = \frac{1}{2} \dot{\phi}^2 - V(\phi)$$

$$\rightarrow \omega_\phi = \frac{p_\phi}{S_\phi} = \frac{\frac{1}{2} \dot{\phi}^2 - V(\phi)}{\frac{1}{2} \dot{\phi}^2 + V(\phi)} \xrightarrow{\dot{\phi}^2 \ll V} -1 < -1/3$$

"slow roll" □