

Atomic collisions, stopping of ions and defect production

P. Hosemann,
Nuclear Engineering Department
University of California, Berkeley

IAEA school

- 1. Atomic collisions and energy loss mechanisms of high-energy ions*
- 2. Defect production and collision cascades*

References:

- 1. Chapter 17 of D. R. Olander, Fundamental Aspects of Nuclear Reactor Fuel Elements, NTIS (1976)*

What is what?

What is radiation: *Radiation: Radiation is the emission and transfer of energy through space!*

What is radioactive material: *Any material that spontaneously emits any type of ionizing radiation.*

What is radioactive contamination: *Radioactive material in an unwanted Location.*

What is ionization: *The process of removing electrons from atoms to make charged ions.*

What is dose: *The effect of radiation on mater.*

What is an Isotope: *Any of the different forms of an element each having a different mass. Isotopes of an element have nuclei with the same number of protons (atomic number) but different number of neutrons.*

What is dose?

Counts per minute (cpm): is a measure of radioactivity. It is the number of atoms in a given quantity of radioactive material that are detected to have decayed in one minute.

Becquerel (Bq): is a measure for radioactivity. It is the disintegrations per second

Gray (Gy): One gray is the absorption of one joule of radiation energy by one kilogram of matter. **J/kg**

Sivert (Sv): is the SI derived unit of dose equivalent. It attempts to reflect the biological effects of radiation as opposed to the physical aspects, which are characterized by the absorbed dose, measured in grays

Röntgen (roentgen) equivalent in man or rem [rem] is a unit of radiation dose. It is the product of the absorbed dose in roentgens (R) and the biological efficiency of the radiation. 100 Rem equal 1 Sivert; the Sivert is the recommended SI derived unit,

Displacement per atom (dpa): is a unit of radiation dose on materials.

Radiation damage in solids

- ***Radiation damage*** refers to the **localized disruption of the crystal lattice** of a solid (defect production) by high-energy radiation passing through it.
- ***Radiation effects*** are the **consequence of radiation damage on the mechanical and/or physical properties of the solid**
- Solids consist of stationary atoms with:
 - a massive nucleus
 - electrons bound to the nucleus with energies ranging from keV (K-shell, heavy atoms) to eV (ionization potential of outer shell)
- Energetic particles of three types:
 - Neutral elementary particles***
 - **neutrons interact with atomic nuclei of solid atoms**
 - **gamma rays interact with electrons** within the solid (pair production, compton scattering and photo-electric effect)
 - Charged elementary particles***
 - **protons, alpha particles**
 - **electrons (accelerator produced, beta particles & Compton)**
 - High-energy atoms or ions***
 - **fission fragments**
 - **accelerator-produced heavy ions**
 - **recoil atoms (ions) of solid produced from primary collisions**

Interactions between energetic particles and solids

- Energetic charged particles interact independently with the nuclei and the electrons in a solid
- Interactions are characterized by *scattering cross sections*

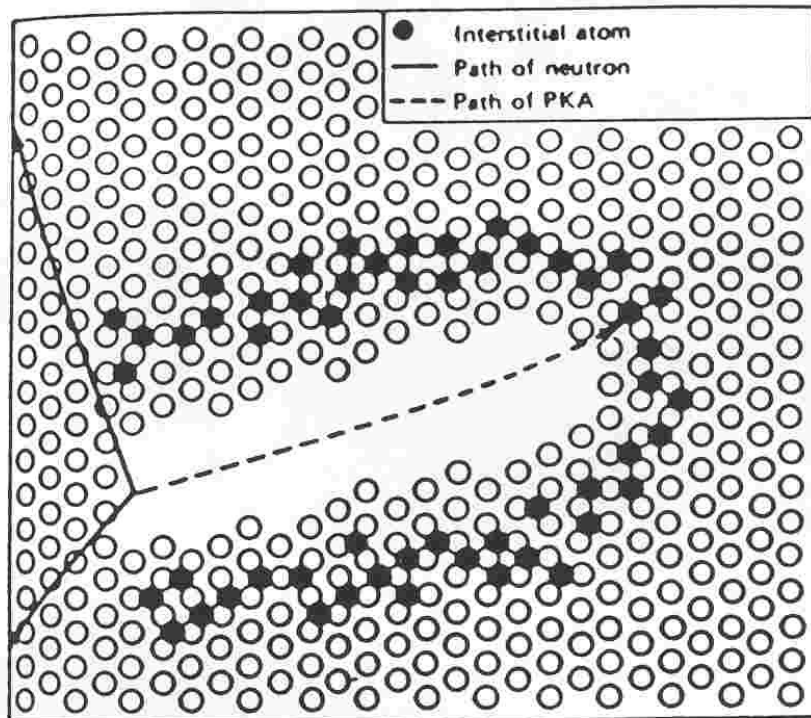
Interactions between *charged particles and electrons in a solid medium*

- Characterized by the stopping power, or energy loss per unit path length
- Electronic stopping is dominant at high energies
- Electrons of the medium are excited (energized) and removed from the nucleus leading to ionization of the medium along ion-track
- Most of the energetic electrons are thermalized and re-captured by nuclei
 - Ionization energy is degraded into heat
- In insulators, some electrons are captured in defect sites, producing changes in electrical and magnetic properties

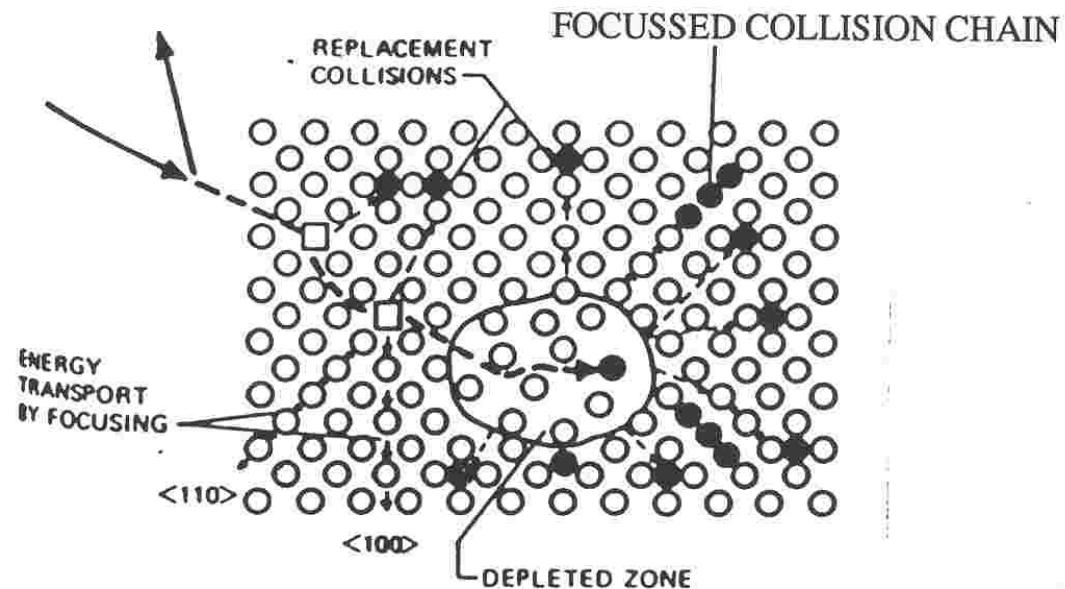
Interactions between *charged particles and nuclei in a solid medium*

- Nuclear interactions predominate at low energies
- Nuclear interactions (scattering) involve:
 - bare nuclei at high particle energy
 - nuclei and bound electrons at low energies
- Nuclear interactions produce permanent atomic displacements within the solid
- Atomic displacements are the underlying cause of physical and mechanical property changes - *radiation damage* and *radiation effects*

Displacement spike: *Early Qualitative Ideas*



Original model of the displacement spike,
J.A. Brinkman, *Amer. J. Phys* 24 (1956) 24.



Later, but still qualitative, version of the
displacement spike (also called depleted zone).

There are ASTM standards!



Designation: E521 – 96 (Reapproved 2009)^{ε1}

Standard Practice for Neutron Radiation Damage Simulation by Charged-Particle Irradiation¹

This standard is issued under the fixed designation E521; the number immediately following the designation indicates the year of original adoption or, in the case of revision, the year of last revision. A number in parentheses indicates the year of last reapproval. A superscript epsilon (ϵ) indicates an editorial change since the last revision or reapproval.

^{ε1} NOTE—Editorial corrections were made in Section 14 in November 2012.



Designation: E798 – 96 (Reapproved 2009)

Standard Practice for Conducting Irradiations at Accelerator-Based Neutron Sources¹

This standard is issued under the fixed designation E798; the number immediately following the designation indicates the year of original adoption or, in the case of revision, the year of last revision. A number in parentheses indicates the year of last reapproval. A superscript epsilon (ϵ) indicates an editorial change since the last revision or reapproval.

of interactions/reactions

The probability of a neutron fission is described by the nuclear cross section σ .

$$r_x = \Phi \sigma_x \rho_A = \Phi \Sigma_x$$

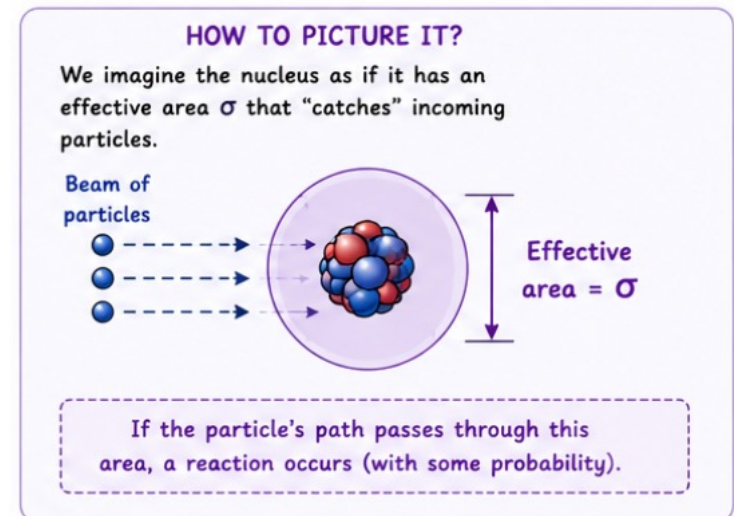
Flux

*Microscopic
cross section*

density of
atoms

*Macroscopic
cross section*

The nuclear cross section depends on the isotope and the energy of the incoming neutron



In nuclear physics, we use "barns" to measure how likely a target is to get hit by a particle!

I'll fire the same shotgun blast at all these barns and see which ones get hit the most!

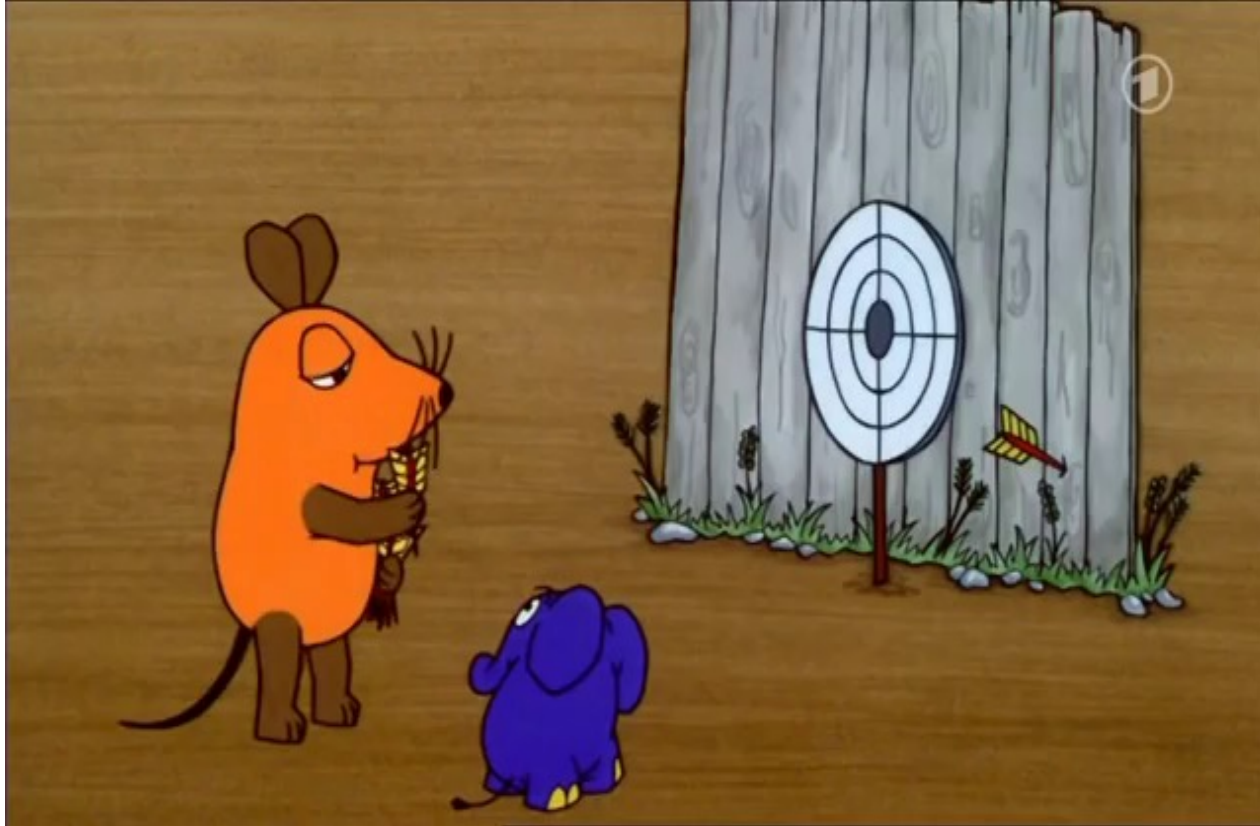
THE MEANING OF THE CROSS SECTION: **A BARN!**

1 barn = 10^{-24} cm^2
(about the size of a barn...
if barns were tiny!)

THE BIGGER THE BARN,
THE GREATER THE CHANCE
OF GETTING HIT!



Cross section collision density



Different types of cross sections

Cross Section Type	Meaning
Total cross section	Probability of ANY interaction
Elastic scattering	Projectile bounces off nucleus
Inelastic scattering	Nucleus excited during collision
Absorption	Particle absorbed by nucleus
Capture	Particle absorbed without fission
Fission	Nucleus splits
Transfer / reaction	Nucleons transferred
Spallation	High-energy fragmentation
Fusion	Two nuclei combine
Charged particle reactions	Proton/alpha-induced reactions
Activation cross section	Produces radioactive isotope

Total Cross Section

$$\sigma_T$$

***Probability for ANY interaction to occur.
Includes: scattering, absorption, fission ,
capture, reactions***

Energy transfer in collisions

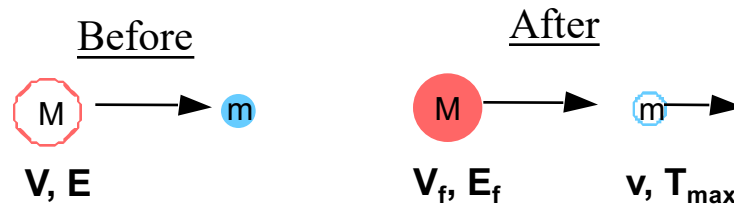
$$R = \frac{\text{collisions}}{\text{unit time} - \text{unit volume}} = N\sigma\phi \quad (2.1)$$

N = atom density of solid (atoms/cm³)

σ = scattering cross-section (cm²)

ϕ = flux of energetic particles (#/(cm²-s))

Head-on collision:



M = mass of projectile particle

m = mass of struck atom

V, E = velocity, energy of projectile particle

$v, T_{\max}$ = velocity, energy of struck particle after collision

Conservation of energy in the head-on collision

$$\frac{1}{2}MV^2 = \frac{1}{2}MV_f^2 + \frac{1}{2}mv^2 \quad \text{or} \quad E = E_f + T_{\max}$$

Conservation of linear momentum in the head-on collision

$$MV = MV_f + mv \quad \text{or} \quad \sqrt{ME} = \sqrt{ME_f} + \sqrt{mT_{\max}}$$

Energy transfer in collisions

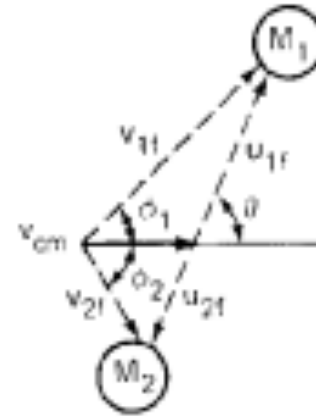
Solving simultaneously by eliminating E_f yields the maximum energy transferred in the (head-on) collision:

$$T_{\max} = \frac{4mM}{(m+M)^2} E \Rightarrow \Lambda E \quad (2.2)$$

More generally, energy transferred to an atom in an elastic collision is given by:

$$T = \frac{1}{2} T_{\max} (1 - \cos\theta) = \frac{1}{2} \Lambda E (1 - \cos\theta) \quad (2.2a)$$

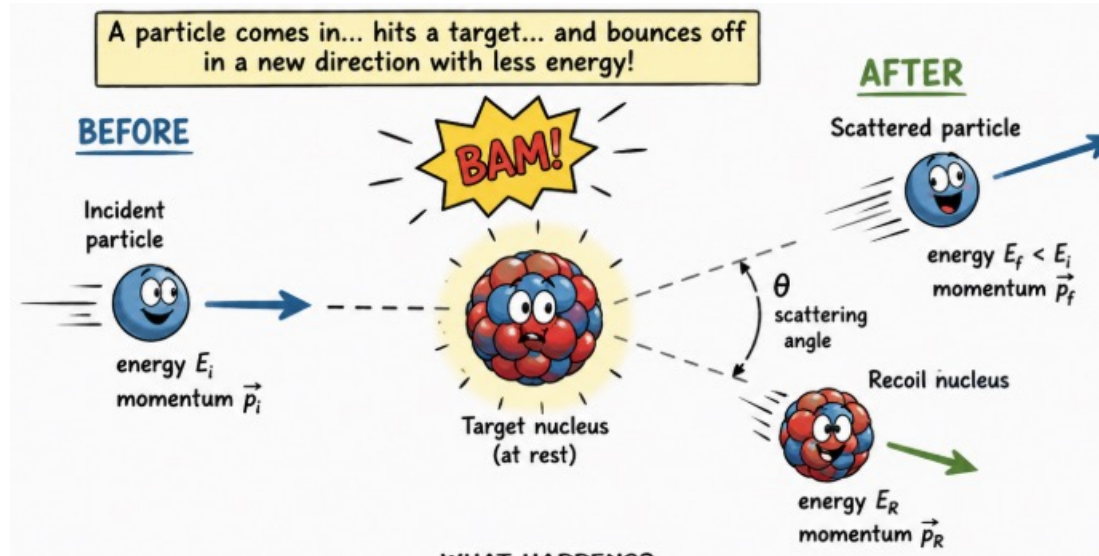
$\theta =$ center of mass scattering angle



- In *hard-sphere* (billiard-ball) collisions, all values of T between 0 and $T_{\max}$ are equally probable.
- The **average energy transferred to the struck atom is $T_{\max}/2$** .
- Energy imparted to struck atoms causes displacements from the lattice sites in the crystal of the solid (*Frenkel pairs*) $\rightarrow$ **Radiation Damage**

Energy transfer in collisions

Example: fission neutron collision with an iron atom

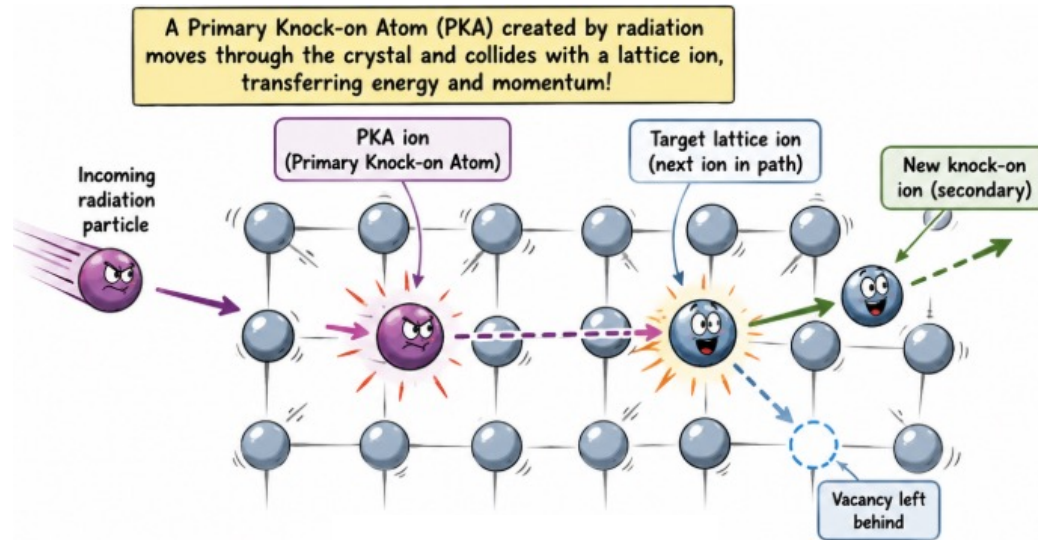


$M = 1$, $E = 1$ MeV (average energy of fission neutron); $m = 56$ (iron);
Substituting into Eq. (2.2):

$$T_{\max} = 0.07 \text{ MeV} = 70 \text{ keV} = 70,000 \text{ eV}$$

- The average energy of the first (iron) atom struck is $0.035 \text{ MeV} = 35 \text{ keV} = 35,000 \text{ eV}$; this struck atom is called the **Primary Knockon Atom (PKA)**

Energy transfer in collisions



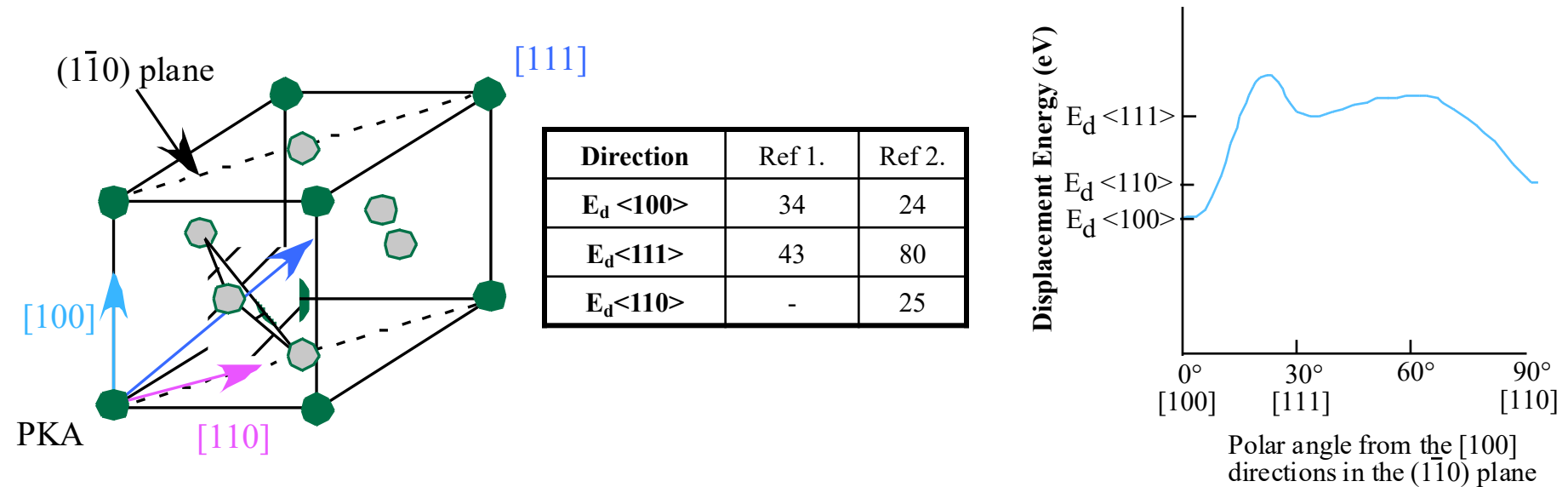
- Since about 25 eV ($E_d = \text{displacement energy}$) is needed to displace a lattice atom from its site in the crystal, the PKA is displaced (often at high velocity), leaving behind a vacant lattice site
- The PKA strikes another lattice atom; since $M = m$ (for atom - atom collisions), the average energy transferred to the second atom (Secondary Knockon Atom, SKA) is one-half of the PKA energy (17.5 keV)
- The average energy of the recoils of the n th generation is one half of that of the previous generation. The process continues as a chain reaction until the recoil energy is less than $E_d \sim 25$ eV.

Displacement Energy

Because of the crystal structure of atoms in a solid, the displacement energy required to produce a Frenkel pair depends on PKA direction

- E_d represents an average over all possible directions

Consider a face-centered cubic (fcc) crystal structure (Cu, Ni, stainless steel):



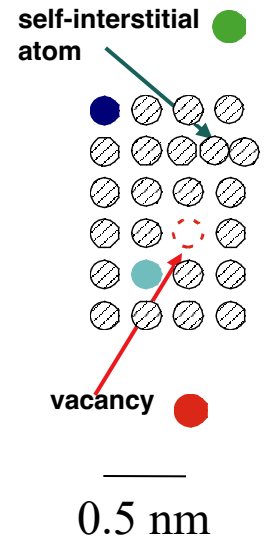
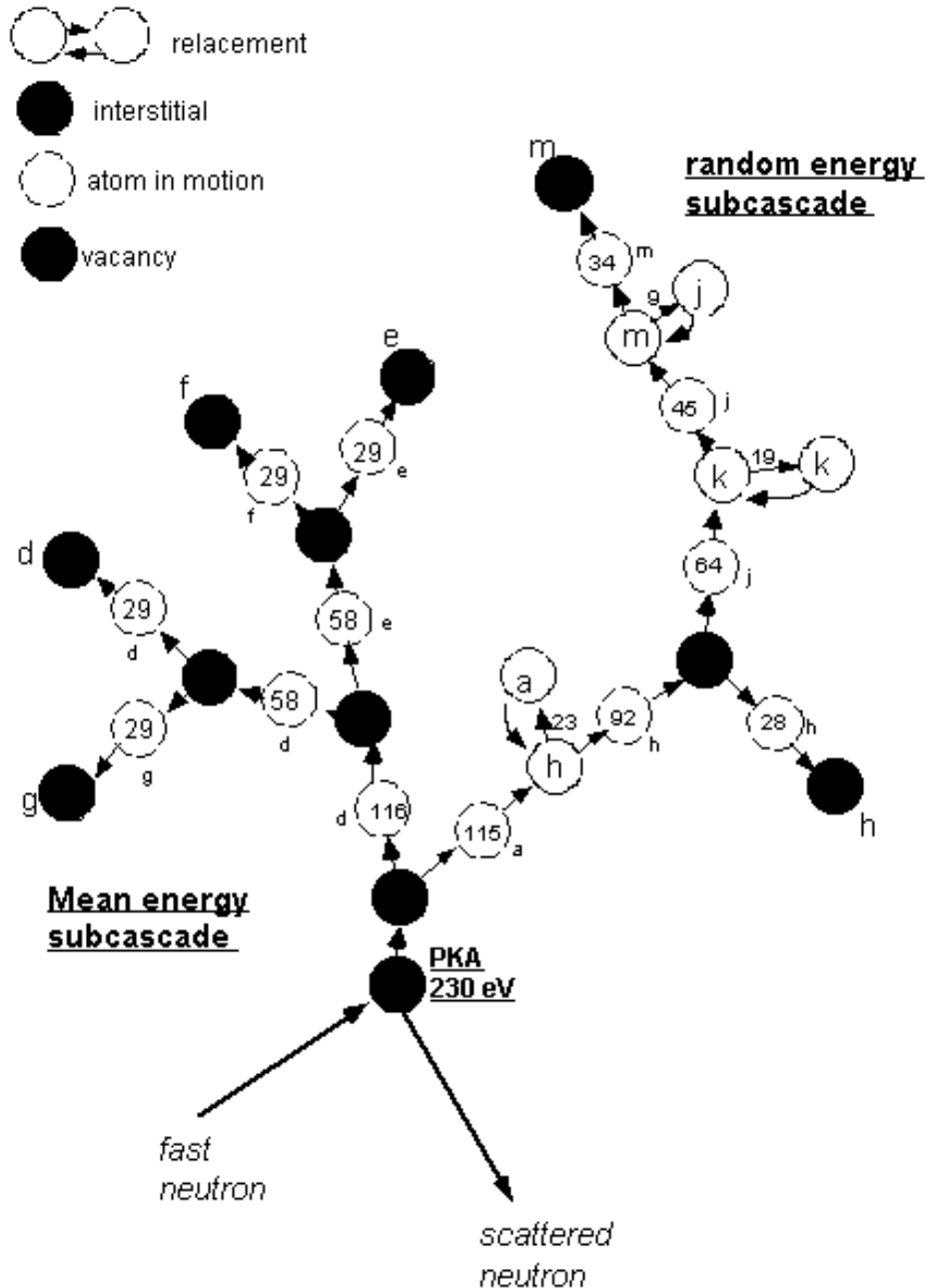
• E_d is much greater than the energy to create a V-SIA pair by equilibrium processes, since the collision is highly non-equilibrium.

• Average value of $E_d \sim 25 - 50$ eV applies to most metals

1. H.B. Huntington, *Phys Rev* **93** (1954) 1414.

2. J.B. Gibson, A.N. Goland, M. Milgram, G.H. Vineyard, *Phys Rev* **120** (1960) 1229.

Example of cascade by low energy PKA's



Lattice displacements due to a single PKA

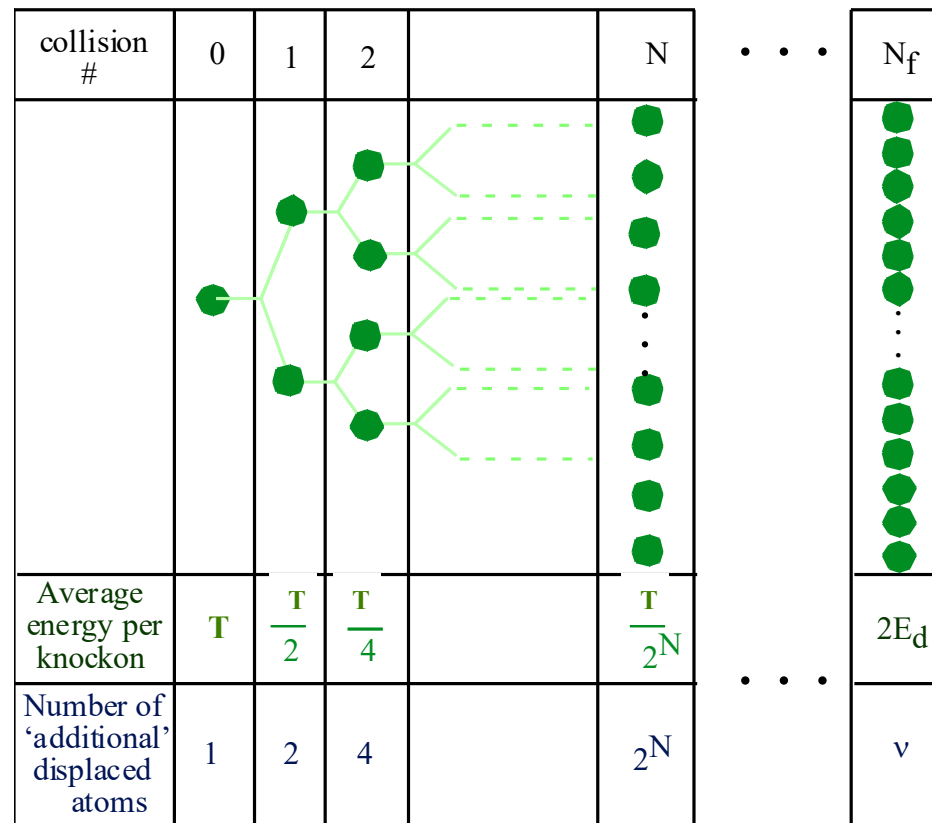
- Thus, the single neutron initiates a collision cascade that contains n atomic displacements (*Frenkel pairs* or vacancy-(self)-interstitial pairs)

Because the PKA energy is so large compared to the displacement energy, each neutron generates many tens and even hundreds of V-I pairs

- $T \leq E_d$ - no displacements
- $E_d \leq T \leq 2E_d$ - only one PKA displaced ($T' < E_d$)
- For $T \geq 2E_d$

- Cascade ceases when average knockon energy is $2E_d = T/2^{N_f}$
- Final number of displaced atoms is $\nu = 2^{N_f}$

$$v = \frac{T}{2E_d} \quad (2.3)$$

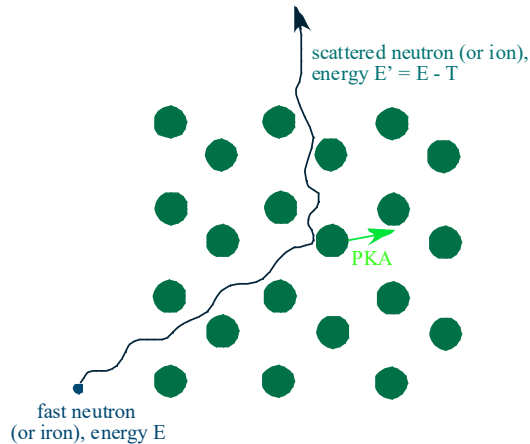


For iron example: $v = 35000/2/25 = 700$ displacements

Number of displaced atoms from nuclear collisions

A fast neutron (or ion) creates Primary Knockon Atoms (PKA) within a solid as a result of (primarily) elastic collisions with energy T .

Primary collision:



If $T > E_c$ - **consider** electronic energy loss (Bethe-Bohr or Lindhard) - **no atomic displacements**

For $T < E_c$ - **consider** atomic displacements (nuclear stopping) which produces vacancies and self-interstitial atoms

E = initial neutron energy

PKA = primary knockon atom = atom initially struck by energetic particle

SKA = secondary knockon atom

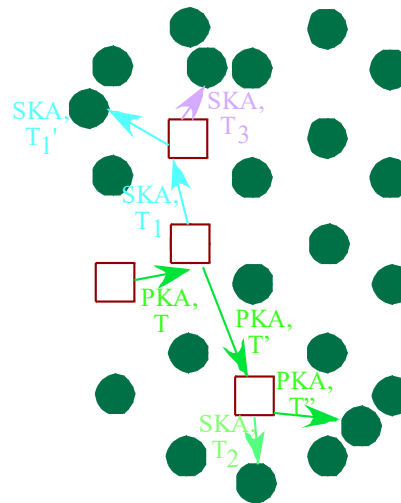
T (or E) = energy transferred to PKA

E' = neutron energy following collision

T_j = energy transferred to struck atom j

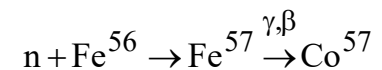
E_d = displacement energy = minimum energy transferred to lattice atom to produce permanent displacements (~ 25 eV)

E_b = energy to form Frenkel (V + SIA) pair (~ 5 eV)



Bombarding particles (material affected):

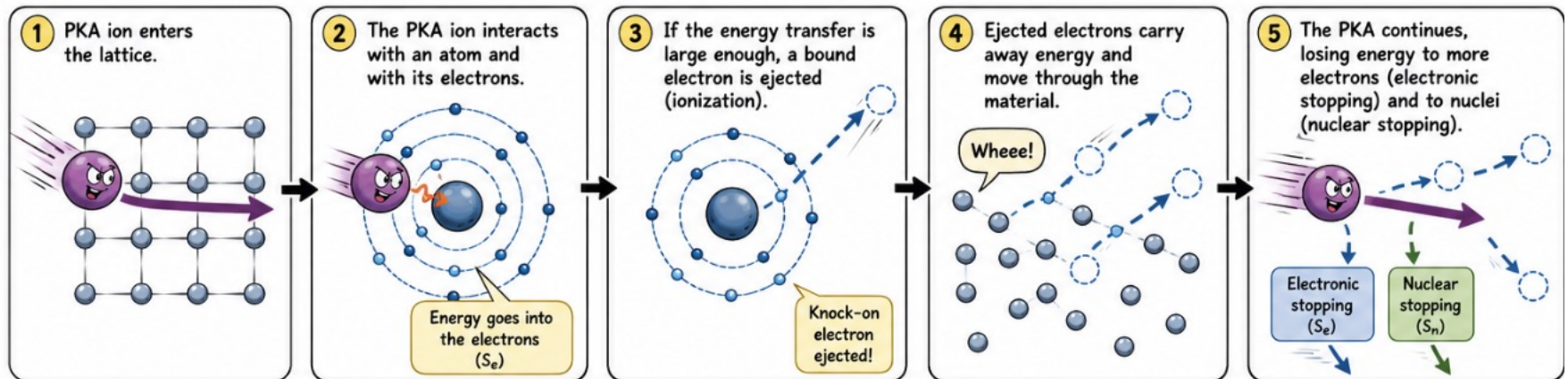
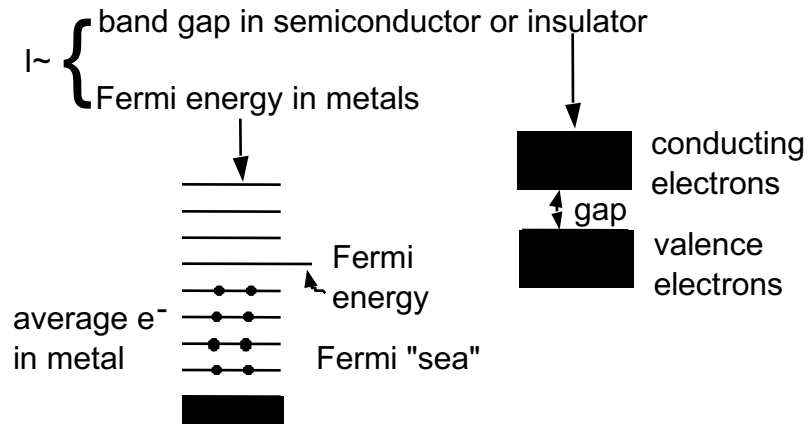
- Fission fragments (fuel UO_2)
- fast neutrons (cladding, structures)
- α -particles and protons (fusion reactor first wall)
- photons (structures - damage by Compton scattering)
- thermal neutrons:



Damage caused by recoil of Co^{57} nucleus due to momentum transfer to γ

Energy loss to electrons (ionization of the solid)

- I = binding energy of electrons in medium (*like ionization potential*)



Energy loss to electrons (ionization of the solid)

- In order for energy transfer **from moving atom** (or ion) **to electron** in medium to occur, a head-on collision of atom (or ion) and electron **must transfer minimum energy I**.
- From Eq. (2.2) with $T_{\max} = I$, $m = m_e$, $E_c =$ minimum ion energy for transferring enough energy to an electron to remove it from the atom:

$$I = \frac{4m_e M}{(m_e + M)^2} E_c = \frac{4m_e}{M} E_c \quad \text{or} \quad E_c = \frac{M}{4m_e} I \quad (2.11)$$

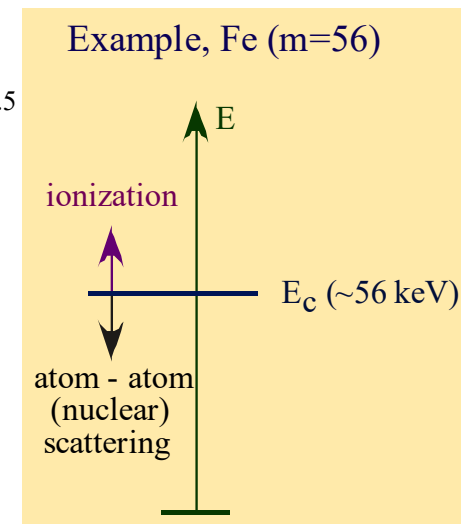
$$\text{for } I = 2 \text{ eV, } \underline{E_c = 10^3 M \text{ (eV)}}$$

- From the initial energy of the ion to E_c , the primary energy loss mechanism is **electronic stopping, also simply called ionization**

$$\left(\frac{dE}{dx} \right)_e = k \sqrt{E} = \text{electronic stopping power of the medium}; k = 8\sigma_e N \left(\frac{m_e}{M_1} \right)^{0.5}$$

Simplified model of the interaction

- for $E > E_c$, all energy loss is by **ionization**, no displacements
- for $E < E_c$, all energy loss by **displacement** collisions with atoms



Calculation dpa in Fe for example

Iron example:

σ = scattering cross section for 1 MeV neutrons $\sim 3 \times 10^{-24} \text{ cm}^2$

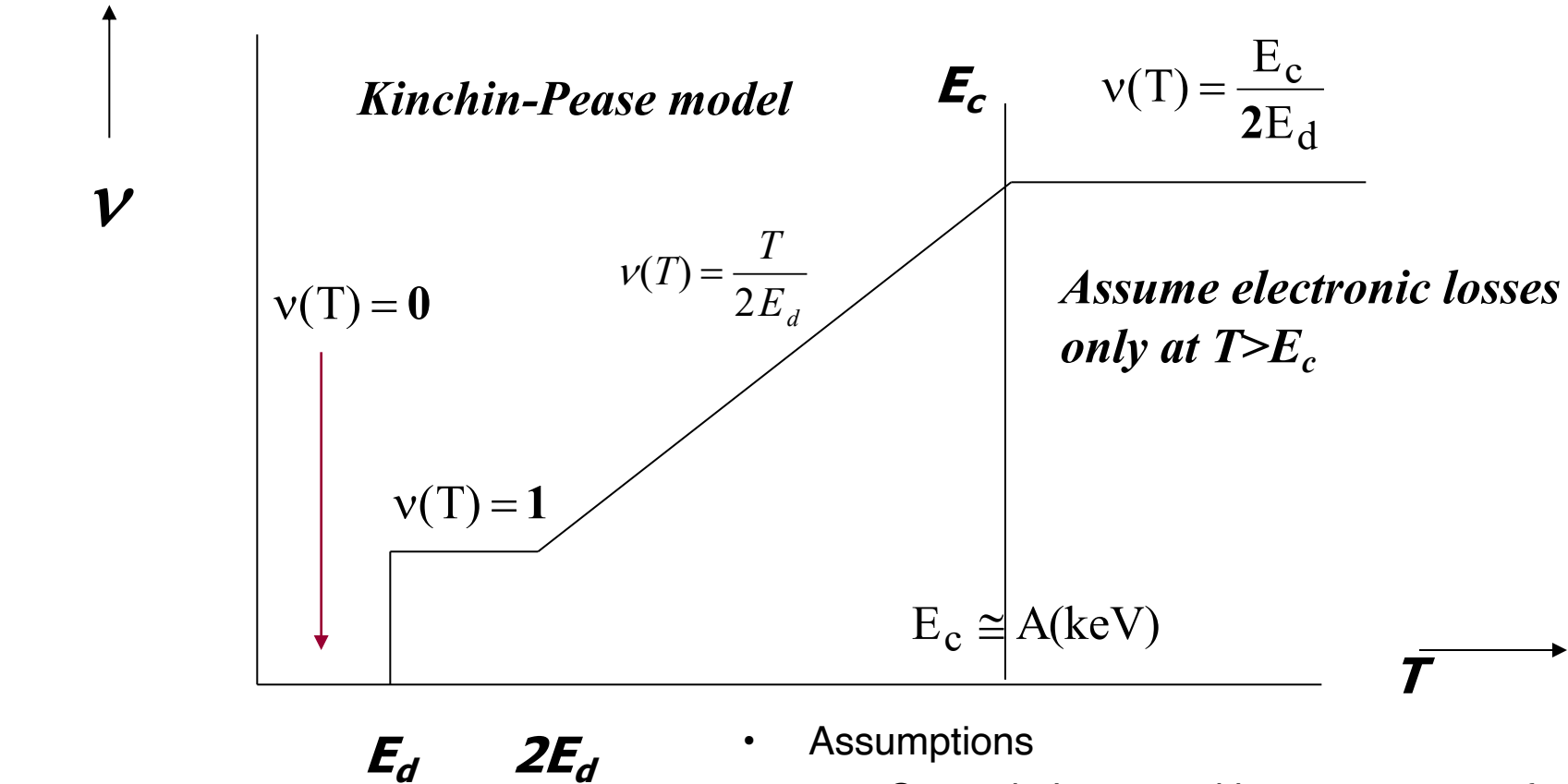
N = atom density of iron atoms in solid $\sim 8 \times 10^{22} \text{ atoms/cm}^3$

ϕ = fast neutron flux $\sim 5 \times 10^{13} \text{ n/(cm}^2\text{s)}$

- neutron - iron atom collision rate = $10^{13} \text{ collisions/cm}^3\text{/s}$
- displacement production rate = $(10^{13})(700) = 7 \times 10^{15} \text{ disp/cm}^3\text{/s}$ (see earlier slide)
- in 3 years (10^8 s): $7 \times 10^{23} \text{ displacements/cm}^3$
- divide by atom density of iron to obtain “dose” in displacements per atom (dpa)

$$\text{dpa} = \frac{7 \times 10^{23}}{8 \times 10^{22}} \approx 9$$

Every atom is ripped off its lattice about 9 times!!
Why doesn't the metal turn to powder???



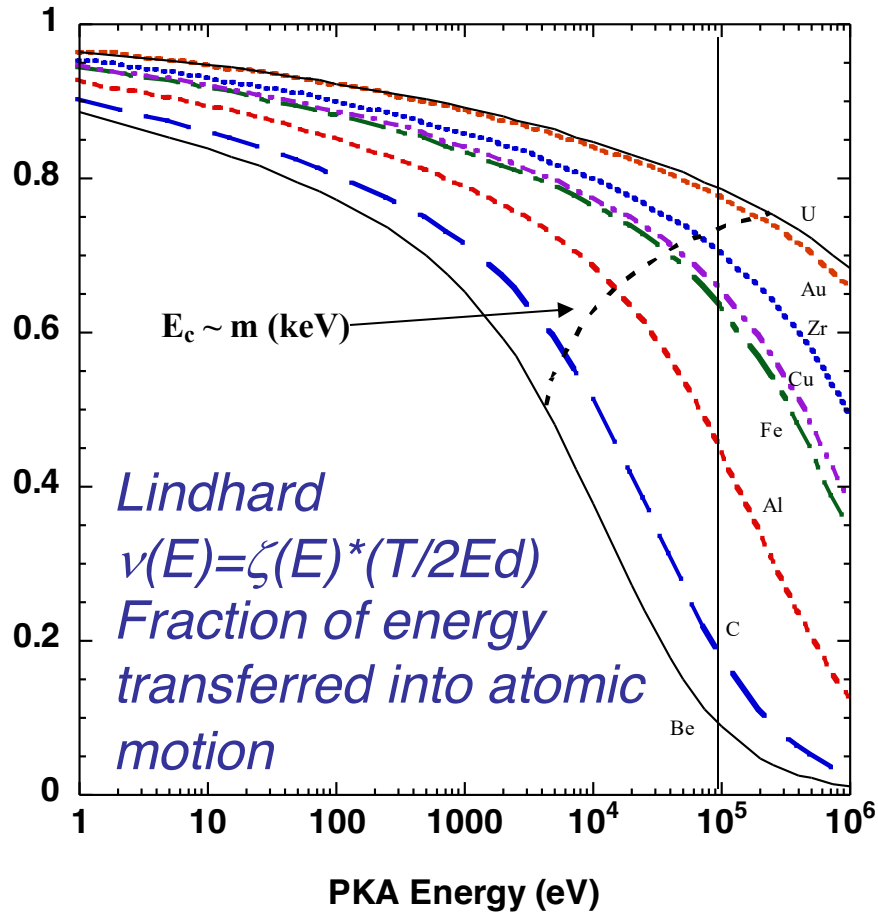
$$\nu(T) = \begin{cases} 0 & \text{for } T < E_d \\ 1 & \text{for } E_d < T < 2E_d \\ \frac{T}{2E_d} & \text{for } 2E_d < T < E_c \\ \frac{E_c}{2E_d} & \text{for } E_c < T \end{cases}$$

• Assumptions

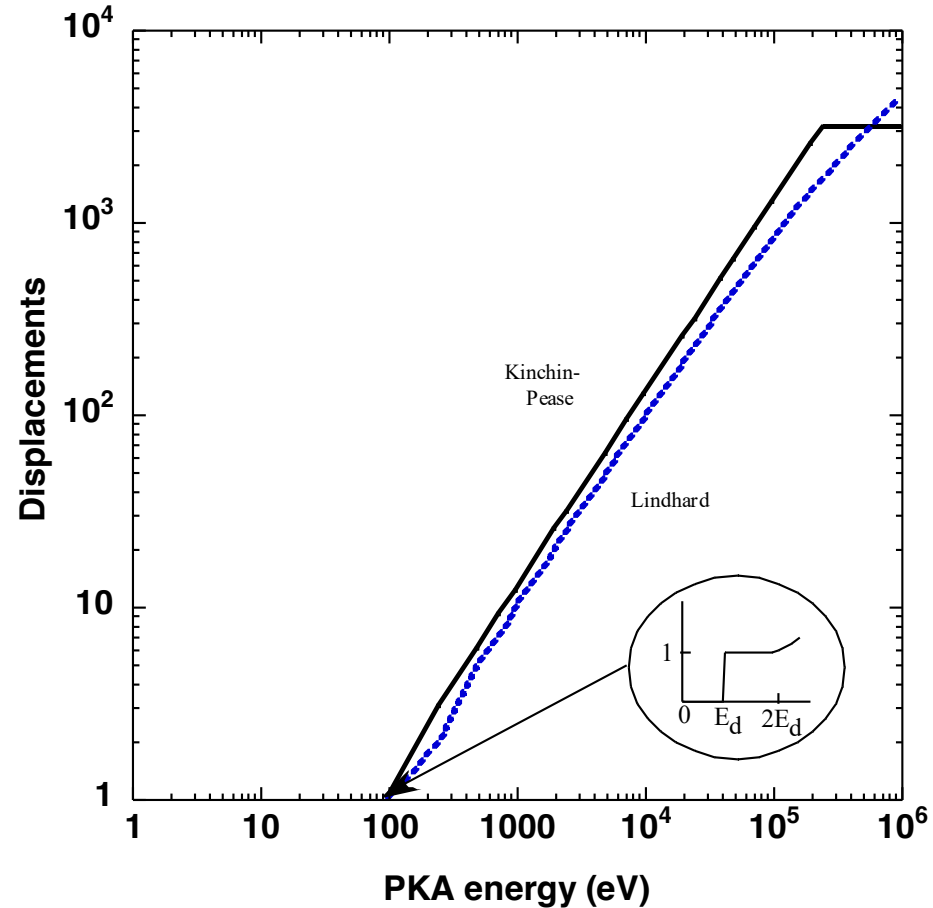
- Cascade is created by a sequence of two-body elastic collisions between atoms
- The energy E_d consumed in displacing an atom is neglected in the energy balance
- Energy loss by electronic stopping is treated by a cutoff energy
- Energy-transfer cross section is given by a hard-sphere model
- The arrangement of atoms in the solid is random; effects due to crystal structure are neglected

Lindhard electronic energy loss

Damage efficiency ζ



Displacements in Fe



ζ = fraction of energy deposited in the material due to collision interaction

Norgett–Robinson–Torrens (NRT)

KP- systematically overpredicted damage

Detailed collision simulations showed that the KP equation predicted about 20–30% too many displaced atoms.

Why?: Because real cascades are more complicated than ideal binary collisions.

For example:

- *replacement collisions reduce the number of vacancies,*
- *energy is transferred to several atoms simultaneously,*
- *some recoil energy becomes lattice vibrations,*
- *not every collision splits energy exactly in half,*
- *some atoms receive just below E_d .*

The Norgett–Robinson–Torrens (NRT) model first calculates the damage energy
 T_d

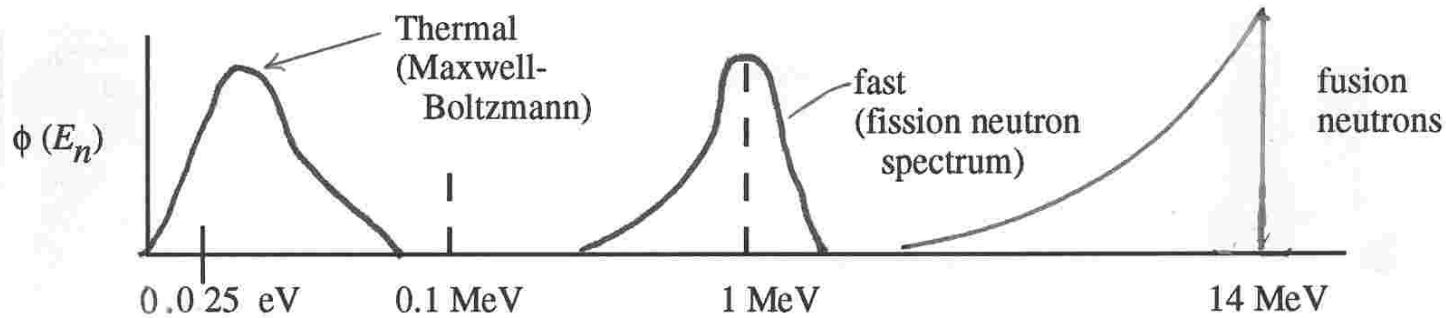
after continuously subtracting electronic stopping using the Lindhard partition.
Then

$$N_d = 0.8 \frac{T_d}{2E_d}$$

Fast neutron scattering & primary collision rate

• E_n = neutron energy

$\phi(E_n)dE_n$ = flux of neutrons with energies between E_n and $E_n + dE_n$ ($n/[cm^2-s-eV]$)



$$\Phi = \text{total fast flux} = \int_{0.1 \text{ MeV}}^{\infty} \phi(E_n) dE_n \quad \frac{\text{neutron}}{\text{cm}^2 - \text{s}} \quad (2.4)$$

Neutron scattering

$n + \text{lattice atom} \rightarrow \text{PKA, energy } T$

$\sigma_n(E_n, T)$ = neutron differential energy transfer scattering cross section (elastic *and* inelastic)

The total neutron scattering cross section:

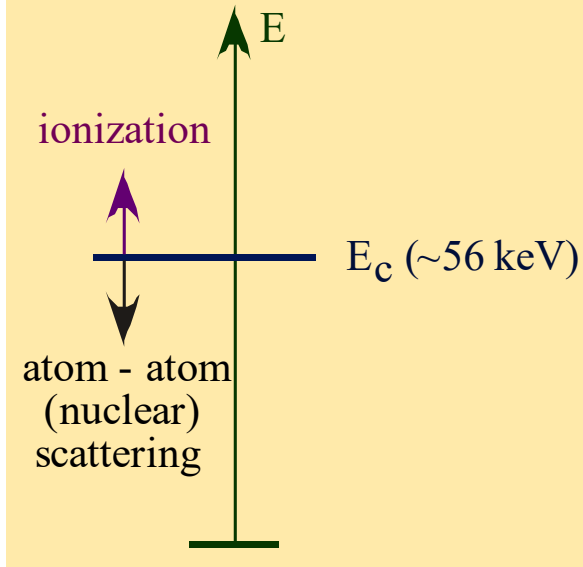
$$\sigma_s(E_n) = \int_0^{\Lambda E_n} \sigma_n(E_n, T) dT$$

For elastic scattering:

$$T_{\max} = \Lambda E_n \quad \bar{T} = \frac{1}{2} \Lambda E_n \quad \Lambda = \frac{4m}{(m+1)^2} \approx \frac{4}{m}$$

Fast neutron interactions

Example, Fe ($m=56$)



Fast reactor neutrons

$$\bar{E}_n = 1 \text{ MeV}$$

$$\bar{T} = \frac{1}{2} T_{\max}$$

$$\approx \frac{1}{2} \frac{4}{m} \bar{E}_n = \frac{2\bar{E}_n}{m} = \frac{2 \times 1000}{56} (\text{keV})$$

$$\bar{T} \approx 36 \text{ keV}$$

Fusion neutrons (14 MeV)

$$\bar{T} = 500 \text{ keV}$$

Displacement Cross Section

Weight the differential scattering cross section with the number of displacements produced by the PKA of energy E :

$$\sigma_d(E_n) = \int_{E_d}^{T_{\max} = \Lambda E_n} \sigma_n(E_n, T) v(T) dT \quad (2.5)$$

For isotropic, elastic scattering:

$$\sigma_n(E_n, T) = \frac{\sigma_s(E_n)}{\Lambda E_n} \quad (2.6)$$

Fast neutron interactions

Approximate displacement function:

$$\sigma_d = \frac{\sigma_s}{\Lambda E_n} \int_{E_d}^{T_{max}} v(T) dT = \frac{\sigma_s}{\Lambda E_n} \left[\int_0^{E_c} \left(\frac{T}{2E_d} \right) dT + \int_{E_c}^{\Lambda E_n} \left(\frac{E_c}{2E_d} \right) dT \right] = \left(\frac{E_c^2}{4\Lambda E_n E_d} + \frac{E_c(\Lambda E_n - E_c)}{2\Lambda E_n E_d} \right) \sigma_s \quad (2.7)$$

**For iron: A (= m) = 56 ($\Lambda = 0.069$); $E_d = 40$ eV; $E_n = 1$ MeV [fission neutrons]
 $E_n = 14$ MeV [fusion neutrons]**

($E_c^2 + \Lambda E_n \dots$) = 415 displaced atoms per neutron collision

($E_c^2 + \Lambda E_n \dots$) = 680 displaced atoms per neutron collision

Using Equ above with $\sigma_s = 3$ barns:

$\rightarrow \sigma_d \sim 1245$ barns [fission] $\sigma_d \sim 2040$ barns [fusion]

Relationship between number of displacements and Dose Rate

Displacement Rate:

$$R_d = N \int_{E_d/\Lambda}^{\infty} \sigma_d(E_n) \phi(E_n) dE_n \quad \frac{\text{displacements}}{\text{cm}^3 - \text{s}} \quad (2.8)$$

atom density (atoms/cm³) = N = 1/ Ω (Ω = atomic volume)

$$K = \frac{R_d}{N} = \frac{\text{displacements}}{\text{atom} - \text{s}} = \int_{E_d/\Lambda}^{\infty} \sigma_d(E_n) \phi(E_n) dE_n = \frac{\text{dpa}}{\text{s}} \quad (2.9)$$

Fast neutron interactions - neutron displacement rate

Approximation:

Evaluate σ_d at $\bar{E}_n$ = average neutron energy (> 0.1 MeV):

$$K = \sigma_d \Phi \quad (2.10)$$

Example:

Fission neutrons $\sigma_d \approx 1250 \text{ barns} = 1.25 \times 10^{-21} \text{ cm}^2$
 $\Phi \approx 5 \times 10^{13} \text{ n/(cm}^2\text{-s)}$

in 5 yrs $\sim 1.6 \times 10^8 \text{ s}$, $\text{dpa} = (1.25 \times 10^{-21})(5 \times 10^{13})(1.6 \times 10^8) \approx \underline{10}$

Each atom in the solid is displaced 10 times in 5 years!!

For D-T fusion $\sigma_d \approx 2040 \text{ barns} \approx 2.0 \times 10^{-21} \text{ cm}^2$
 $\Phi \approx 5 \times 10^{13} \text{ n/(cm}^2\text{-s)}$

in 5 yrs $\sim 1.6 \times 10^8 \text{ s}$, $\text{dpa} = (2.0 \times 10^{-21})(5 \times 10^{13})(1.6 \times 10^8) \approx \underline{16}$

Each atom in the solid is displaced 16 times in 5 years (actually too low)

At high neutron energies

- *inelastic scattering* begins: $n + m \rightarrow m^* + n'$
- *anisotropic scattering* dominates: Eq. (2.6) is no longer applicable
- (n,p) and (n,α) transmutations occur

Computer Simulation of Neutron Damage: Molecular Dynamics of Displacement Cascades

- A completely numerical representation of damage due to fast neutrons can be achieved by **following the atomic trajectories in a crystallite of N lattice atoms**, when one atom is given kinetic energy characteristic of a neutron scattering event
- The force on the i^{th} atom in the crystal is the sum of forces from all neighboring atoms in the crystal, with Newton's 2nd law of motion applied to each atom (mass m):

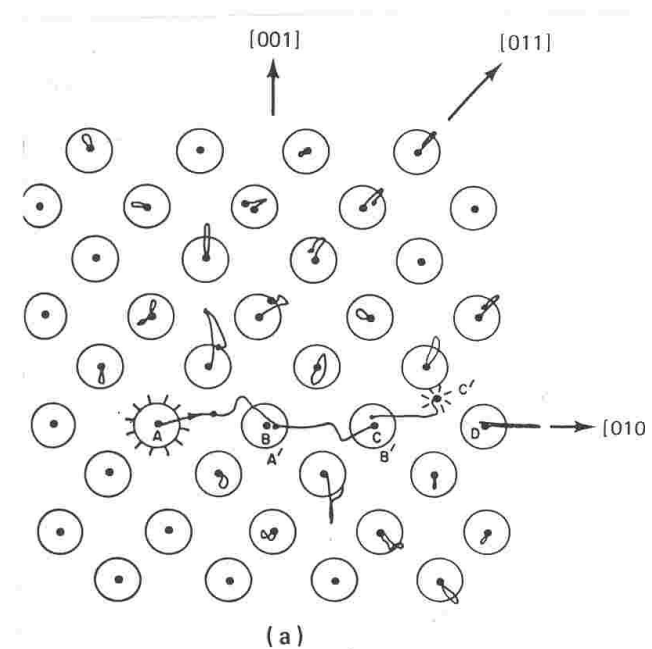
$$m \frac{d^2 \vec{x}_i}{dt^2} = \vec{F}_i = - \frac{\partial V(r_{i1}, r_{i2}, r_{ij} \dots r_{iN})}{\partial \vec{r}_i} \quad (2.12)$$

$\vec{x}_i$ = position of the atom

r_{ij} = separation between atoms i and j

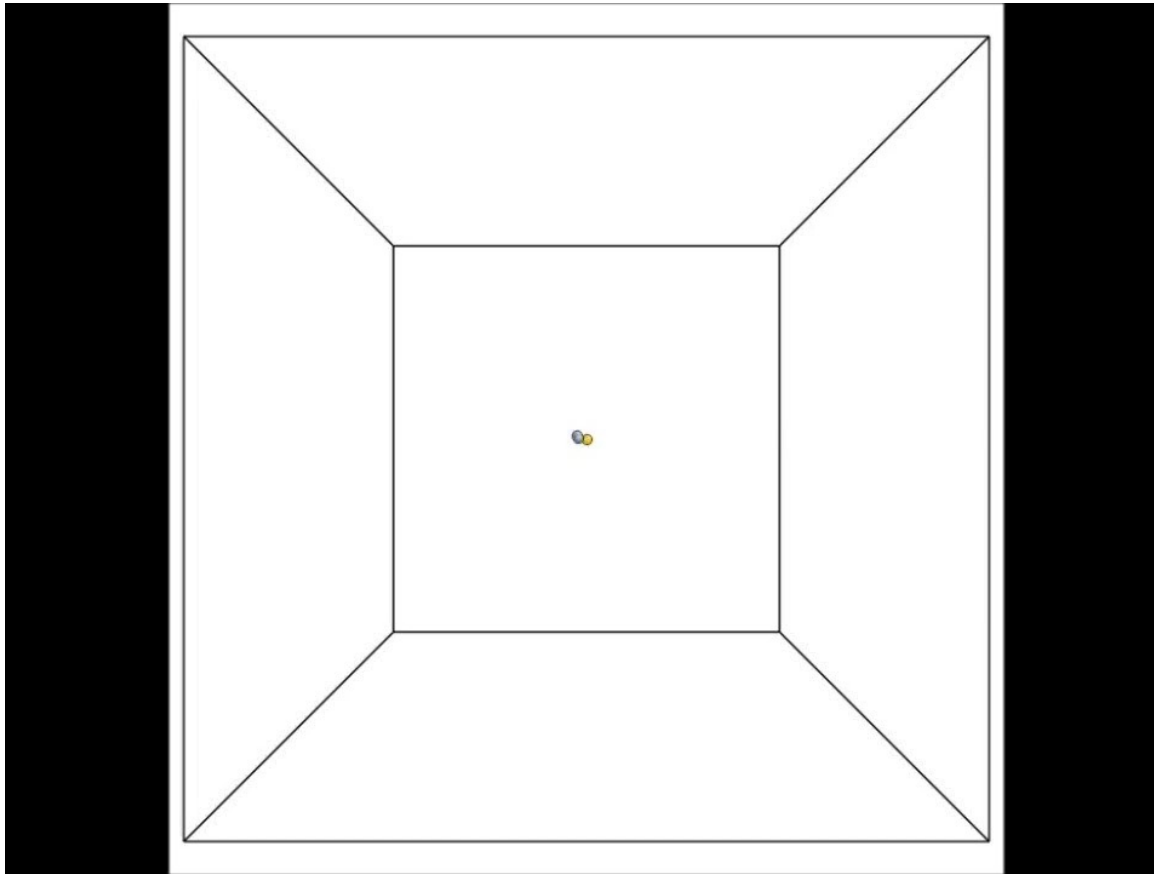
$V(r_{ij})$ = potential energy between i - j pair

- The system of $3N$ (coupled) differential equations is solved for $\vec{x}_i$
- Note that the displacement energy, E_d , need not be specified a priori
- Shown to the right are early simulation results of a 40 eV PKA in the (100) plane of Cu

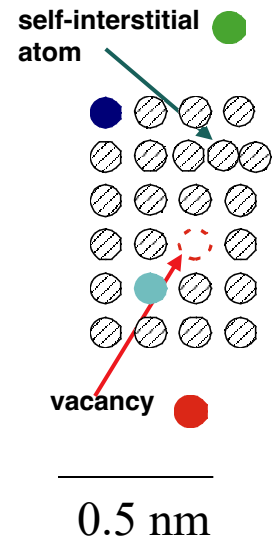


Damage structure: high energy PKA

Many atomic displacements (V-SIA) created at higher PKA energy and atomic displacement structure is clustered and spatially correlated (high energy PKAs typically result from heavy ion or fast neutron irradiation)



20 keV PKA in Fe:
the average energy
transferred by a 550 keV
neutron



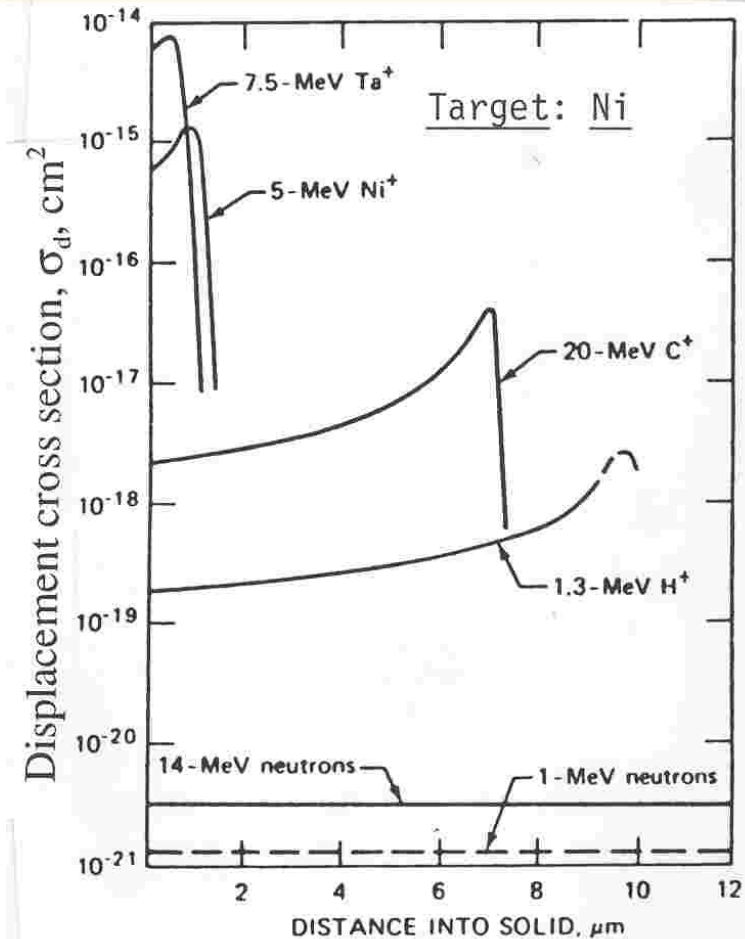
Comparison of ion and fast neutron damage

Since the energy transfer cross section for ions is an atomic cross section ($\sim 10^{-17} \text{ cm}^2$) while that for neutrons is a nuclear cross section ($\sim 10^{-24} \text{ cm}^2$),
 $(\sigma_d)_{\text{ions}} \gg (\sigma_d)_{\text{neutrons}}$

Examples:

- 1) 20 MeV C⁺: $I = 10^{14} \text{ cm}^{-2}\text{sec}^{-1}$
gives $\sim 4 \times 10^{-3} \text{ dpa/s}$ in a thin layer;
dpa varies with penetration depth
- 2) neutrons give $\sim 6 \times 10^{-8} \text{ dpa/s}$;
dpa distribution is uniform

Characteristic	Neutrons	Ions
Flux, $\text{cm}^{-2}\text{s}^{-1}$	5×10^{13} (isotropic)	10^{11} - 10^{12} (beam)
dose rate, dpa/s	$\sim 10^{-7}$	10^{-5} - 10^{-4}
Time for 1 dpa	$\sim 4 \text{ mo.}$	3 - 30 hr
Temp., °C	~ 300	adjustable
Macro. damage distr.	uniform	peak at end of range
Micro. damage distr.	inhomogeneous	inhomogeneous
Irradiated area	large	$\sim 1 \text{ cm}^2$
Irradiation depth	bulk	1 - 10 μm nr. Surf.
Free defect fraction	3 - 10%	$\sim 30\%$
σ_d , cm^2	$\sim 10^{-21}$ Eq(1.52)	10^{-18} - 10^{-15} Eq(1.89)

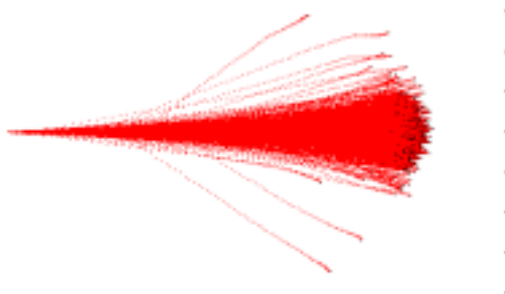


SRIM *The Stopping and Range of Ions in Matter*

SRIM is a group of programs which calculate the stopping and range of ions (up to 2 GeV/amu) into matter using a quantum mechanical treatment of ion-atom collisions

SRIM is a Kinetic Monte-Carlo program (so no time dependence) that uses the potentials you've talked about in class to simulate radiation damage in a material

Calculates most aspects of the energy loss of ions in matter, produces tables of stopping powers, range and straggling distributions for any, ion at any energy in any elemental target and is able to do targets with complex multi-layer configurations



SRIM *INPUT*

Material composition ; Material configuration (SRIM can handle layers of differing composition)
Displacement energy of each element in each layer for every material
SRIM will recommend a displacement energy
Density of each layer of material; SRIM can calculate this for you – not necessarily accurate
Incident ion type and energy
Type of calculation

Types of SRIM calc:

Ion Distribution and Quick Calculation of Damage

Don't care about details of target damage or sputtering; Quick statistical estimates based on the Kinchin-Pease Model

Detailed Calculation with full Damage Cascades

Follows every recoil until its energy drops below the lowest displacement energy of any target atom; All collision damage to the target is analyzed.

Surface Sputtering / Monolayer Collisions

A full treatment of sputtering is contained in the above calculation for full damage cascades. *TRIM contains the concept of a "free flight path", in which ions make large jumps between collisions and any target damage is spread over the jumped target atoms.*

Read Me

TRIM (Monte Carlo Ranges)

Type of TRIM Calculation

DAMAGE

Detailed Calculation with full Damage Cascades

?

TRIM Demo

?

Restore Last TRIM Data

?

Basic Plots

Ion Distribution with Recoils projected on Y-Plane

?

? ION DATA

PT

H

Hydrogen

1

1.008

5000

?

0

? TARGET DATA

Input Elements to Layer 1

Layers

Add New Layer

?

Add New Element to Layer

Compound Dictionary

?

Layer Name

Width

Density
(g/cm³)

Compound
Corr

Gas

Symbol

Name

Atomic
Number

Weight
(amu)

Atom
Stoich or %

Damage (eV)
Disp

Latt

Surf

X Stainless Steel 100 8 1

X PT Cr Chromium 24 51.99 8 08.0 25 3 4.1

X PT Fe Iron 26 55.84 24 24.0 26 2 4.3

X PT Ni Nickel 28 58.69 18 18.0 25 3 4.4

Special Parameters

Name of Calculation

H (10) into Layer 1

Stopping Power Version

SRIM-2003

?

Output Disk Files

? ☐ Ion Ranges

? ☐ Backscattered Ions

? ☐ Transmitted Ions/Recoils

? ☐ Sputtered Atoms

? ☐ Collision Details

? ☐ Special "XYZ File" Increment (eV)

☐ Resume saved
TRIM calc.

☐ Use TRIM-96
(DOS)

Save Input &
Run TRIM

Clear All

Calculate Quick
Range Table

Main Menu

Problem Solving

Quit

? AutoSave at Ion # 10000

? Total Number of Ions 99999

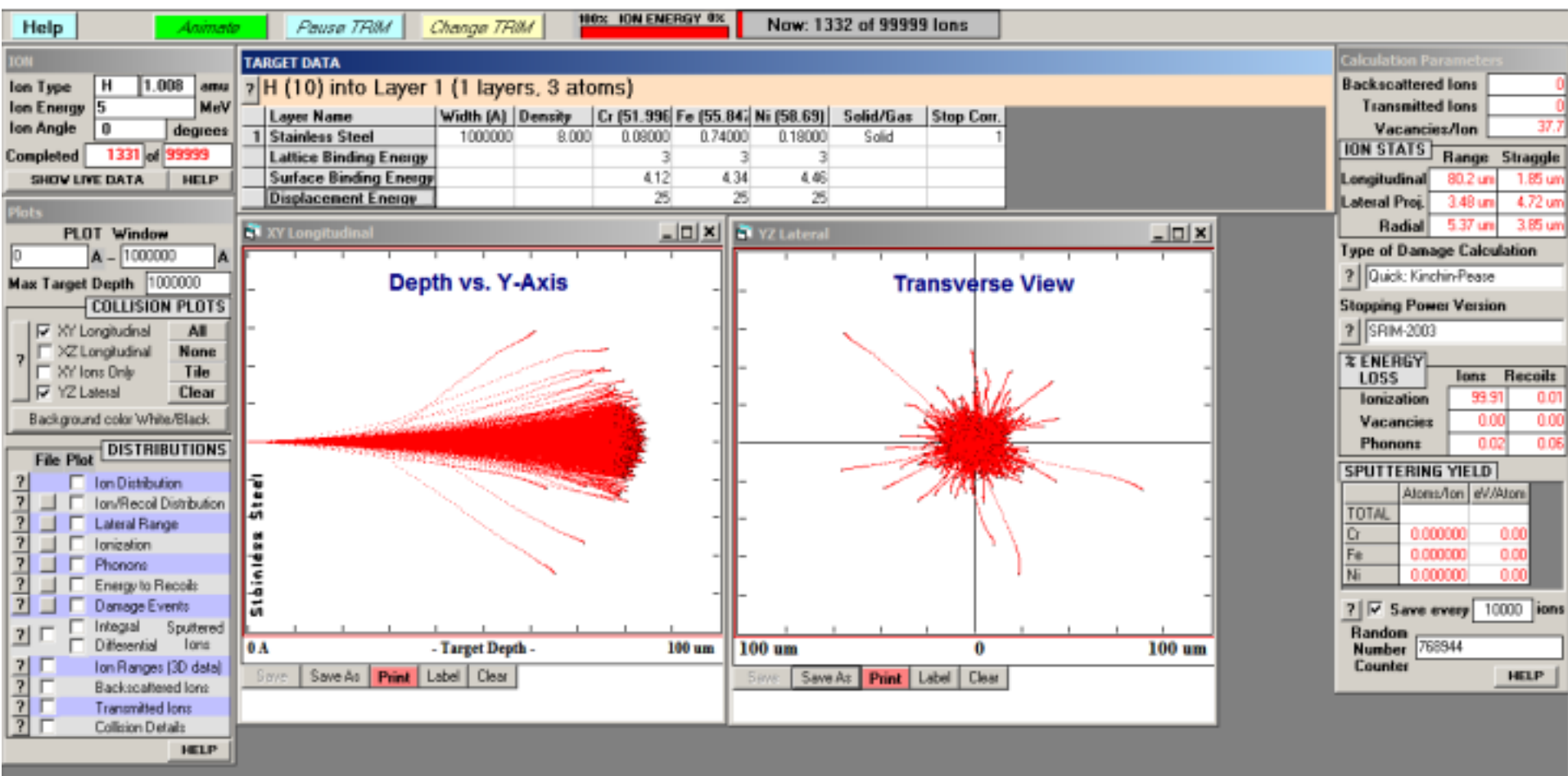
? Random Number Seed

Plotting Window Depths

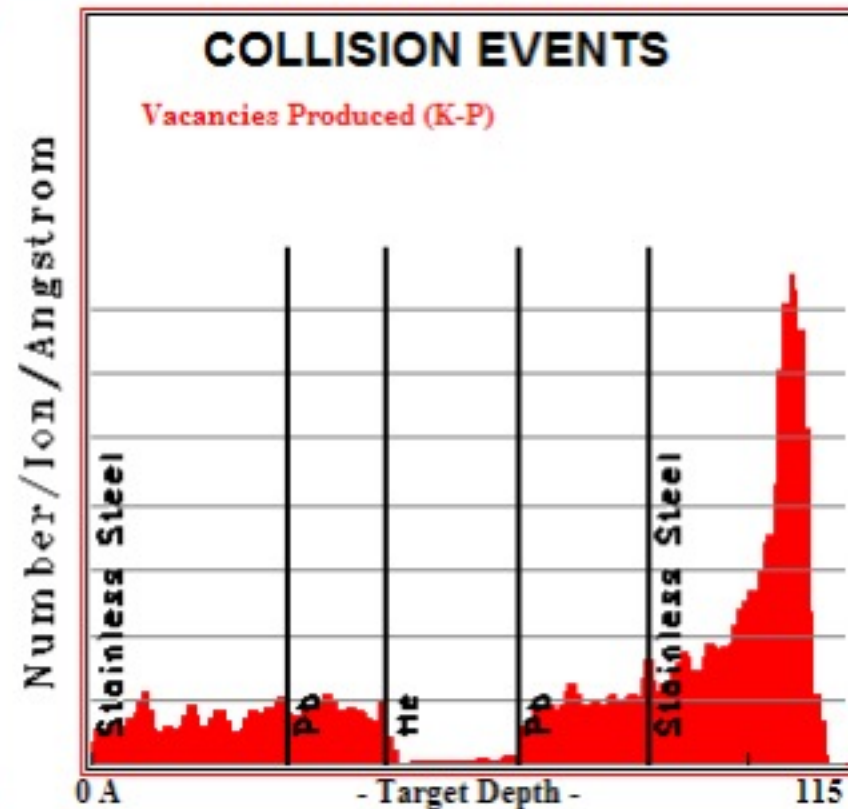
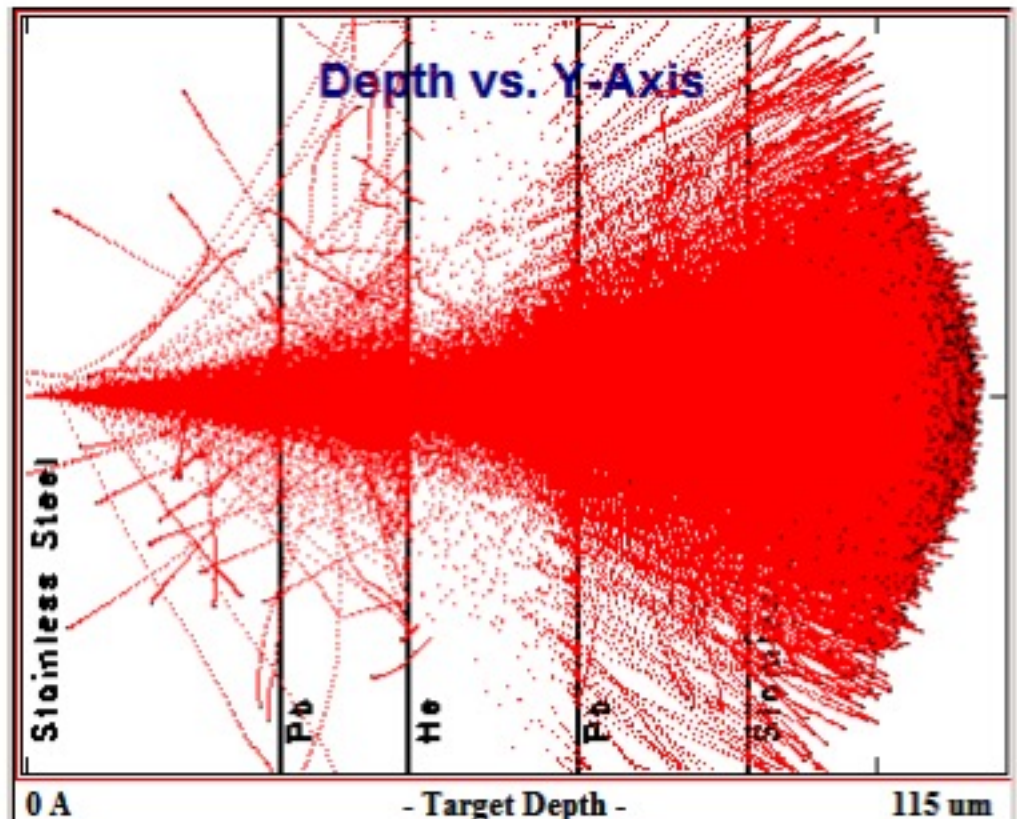
Min 0

Max 1000000

SRIM Output



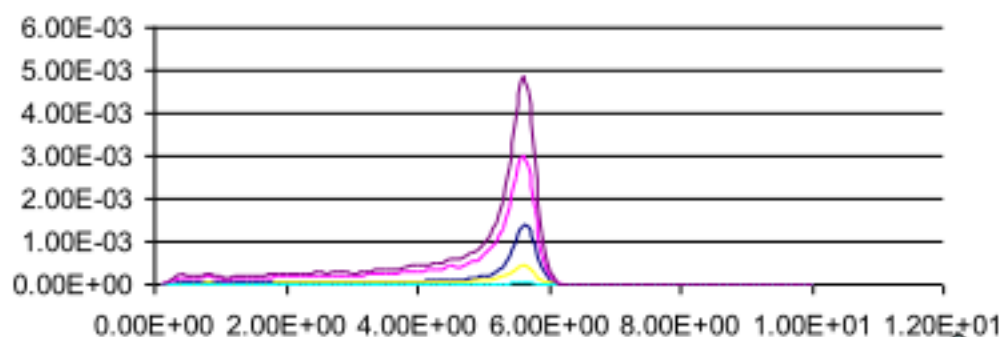
Multilayer structures



SRIM is a very powerful and easy to use tool.

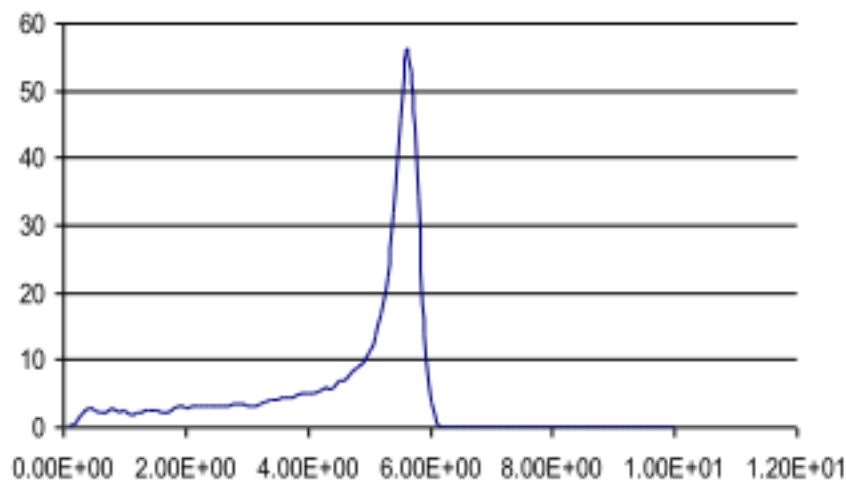
Multi layer systems are possible as well as angled impacts.

Dpa calc. using SRIM



Vacancies produced in the material by incoming ion calculated using SRIM.

dpa in the material caused by the incoming ion with a given beam current, beam density, surface area, etc.



Vacancy output

```
===== H (3500) into Ta =====
SRIM-2003.26

=====
Ion and Target VACANCY production
See SRIM Outputs\TDATA.txt for calc. details
=====
===== TRIM Calc.= H(3.5 MeV) ==> Ta( 50 um) =====
=====
See file : SRIM outputs\TDATA.txt for calculation data
Ion      = H      Energy = 3500 keV
===== TARGET MATERIAL =====
Layer 1 : Ta
Layer width = 5.E+05 A :
Layer # 1- Density = 5.524E22 atoms/cm3 = 16.60 g/cm3
Layer # 1- Ta = 100 Atomic Percent = 100 Mass Percent
=====
Total Ions calculated = 6984.00
Total Target Vacancies = 75 /Ion
Total Target Displacements = 83 /Ion
Total Target Replacement Collisions = 7 /Ion

!!!! NOTE : 2nd column below is number of Primary Knock-Ons !!!!
              ( PKO are number of Target Atoms Recoiling from the Ion. )
=====
Table Units are >>>> Vacancies / Angstrom / Ion <<<<
=====
TARGET
DEPTH      H      Ta
(Ang.)     knock-Ons  Vacancies
-----
500001.E-02 3665.52E-09 1431.84E-08
100000.E-01 1323.02E-08 8287.51E-08
150000.E-01 1131.16E-08 6847.08E-08
200000.E-01 1239.98E-08 7903.78E-08
250000.E-01 1199.89E-08 7468.50E-08
-----
```

Sum of all ions

SRIM *dpa* calc.

$$\text{SRIM result} = \frac{\text{vacancies}}{\text{ion} \cdot \text{Angstrom}}$$

$$\text{dpa} = \text{Vac}/(\text{ion } \text{\AA}) * \text{number of ions} / \text{number density}$$

Number of ions:

$$\left. \begin{array}{l} \text{Beam current } I = A = C/\text{sec} \\ 1e = 1.6 * 10^{-19} \text{ C} \end{array} \right\} I * t / (1.6 * 10^{-19}) = \text{ions}$$

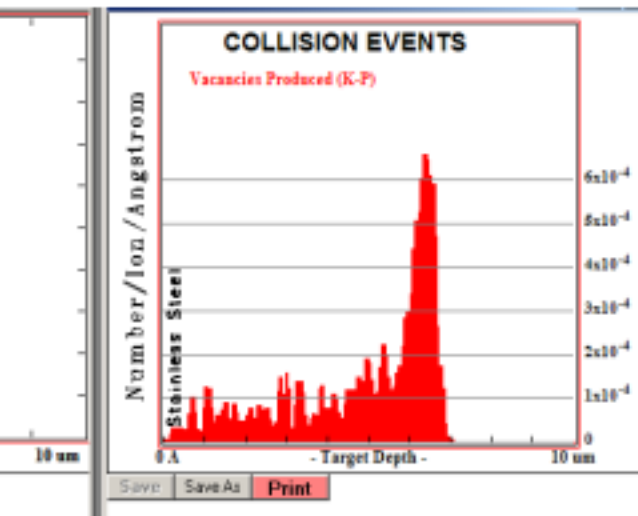
Consider ion charge state!!
Consider irradiation area!!

Number density:

$$\text{Atomic density} * \text{surface irradiated area} = \text{Atoms/cm}^3 * \text{cm}^2 = \text{Atoms/cm}$$

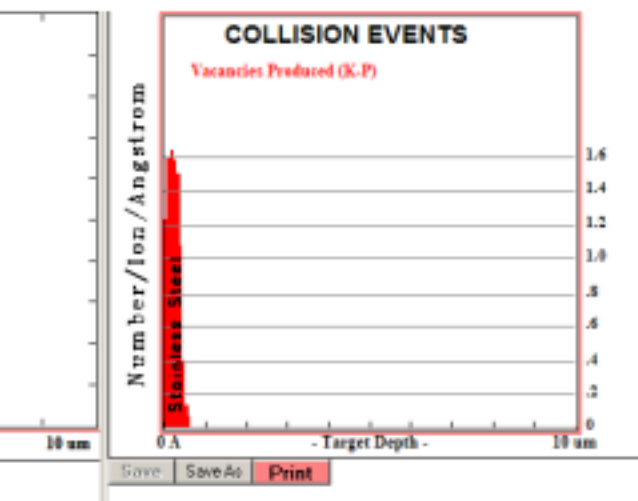
WATCH UNITS !!!!!!!!!!!!!!!!!!!!!!!!!!!!!

Ion comparison with SRIM



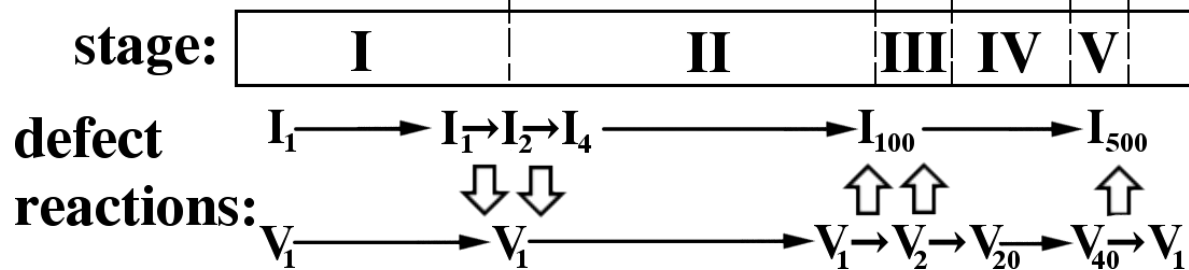
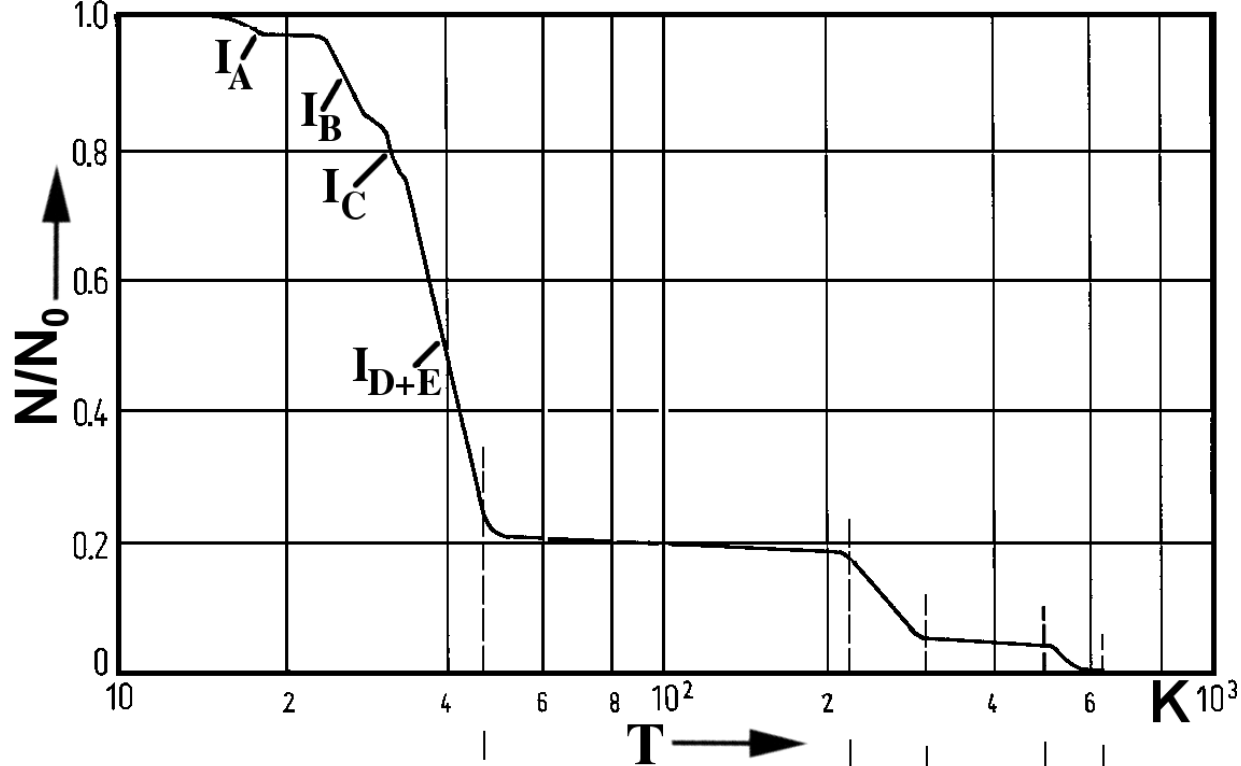
1MeV protons in stainless steel

Heavier ions cause more damage!
Lighter ions are going deeper into the material than heavier ions
Heavier ions cause higher dose rate.



1MeV Fe²⁺ in stainless steel

So we created displacements.
So What??



Five recovery stages in electron-irradiated Cu I and V denote single interstitials and vacancies, respectively; and $I2$ and $V2$, di-interstitials.

Stage I: The lower temperature substages IA , IB , and IC are due to collapse of close Frenkel pairs. ID is due to correlated recombination and IE due to uncorrelated recombination.

Stage II: Interstitial clusters grow, leading to identifiable small interstitial loops.

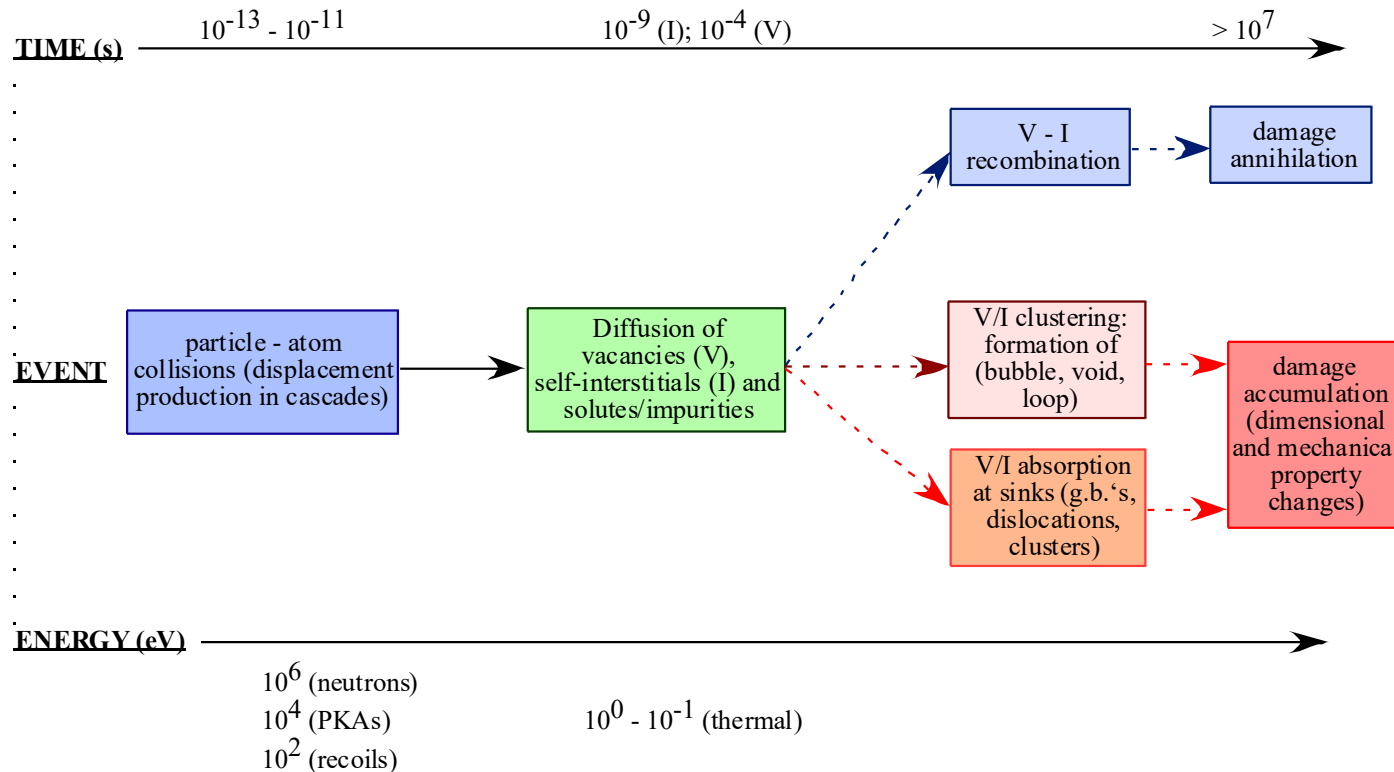
Stage III: Vacancies migrate and adulate at interstitial clusters.

Stage IV: The vacancy clusters surviving stage III grow in size. The vacancy clusters dissociate thermally.

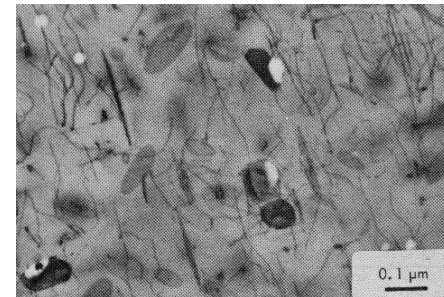
Fate of the displaced atoms

- Recombination of vacancies and interstitials: $V + I \rightarrow \text{null}$ (perfect lattice)
- **Vacancies** coalesce into:
 - small clusters (depleted zones or *displacement spikes*) - in fcc metals, these are usually stacking fault tetrahedra (SFT); in bcc metals, these are ultra-small, sub-nanometer cavities (nanovoids)
 - large clusters (voids or cavities)
- **Interstitials** coalesce into **dislocation loops**
- It is these point defect clusters (along with radiation induced/enhanced diffusion, precipitation and segregation) that are responsible for mechanical (hardening and embrittlement) and/or physical property changes

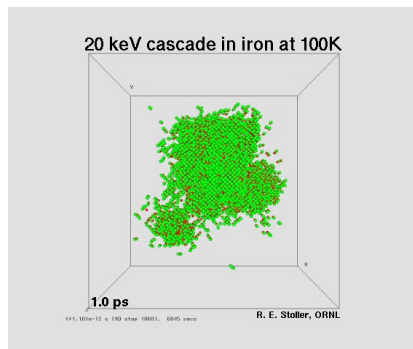
Radiation effects are governed by the ultimate fate of point defects (& transmutants)



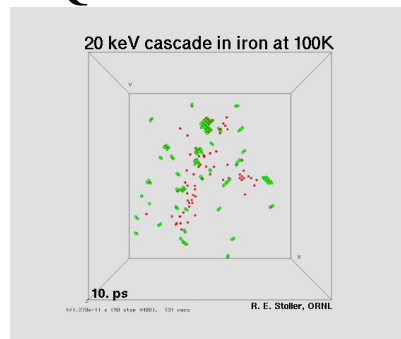
Voids, I-loops
& precipitates
in stainless steel



Nascent cascade



“Quenched” cascade

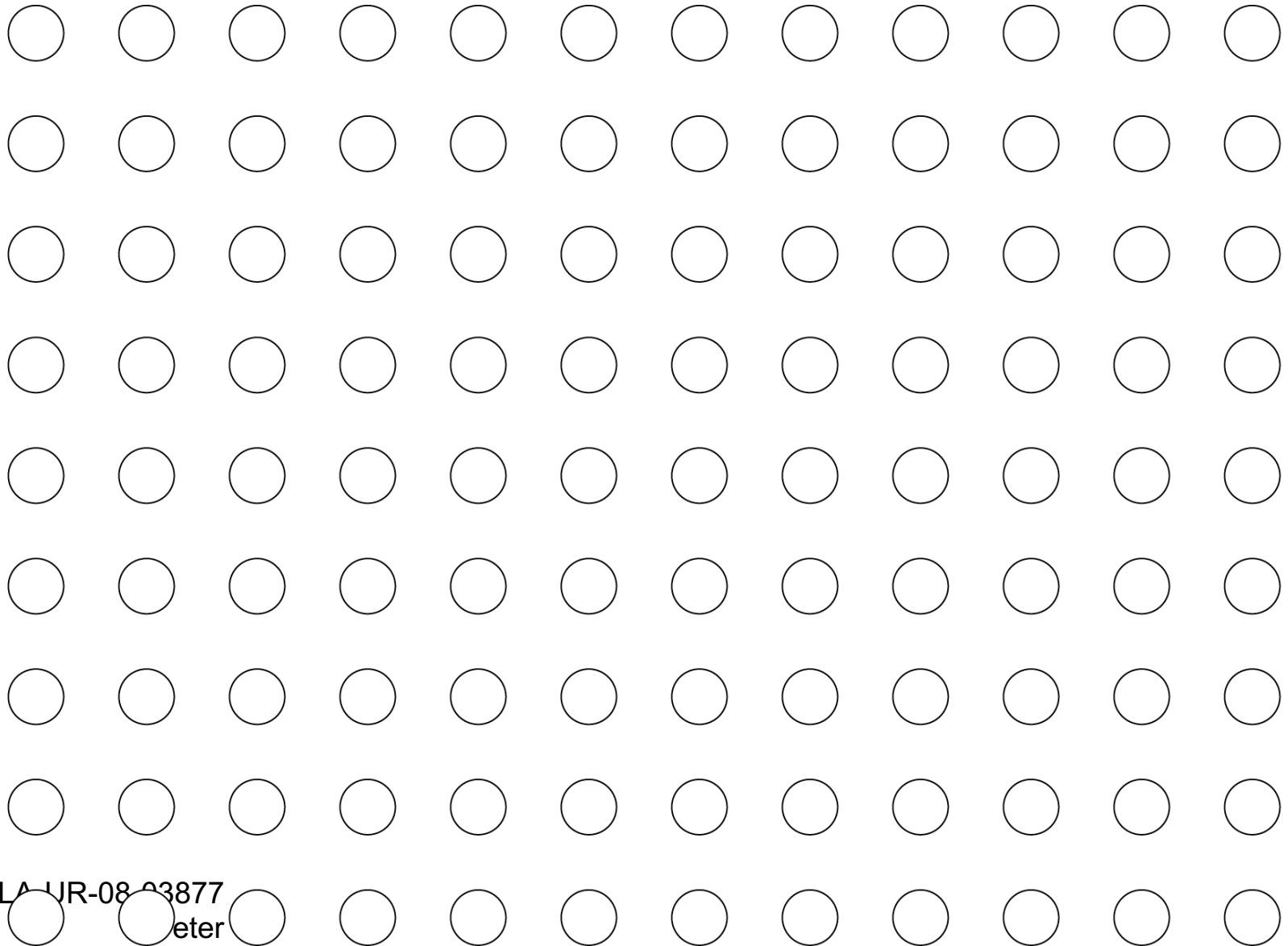


Number of defects:

Nascent cascade: $N_i = N_v = v(T)$

“Quenched” cascade: $N_i' = N_v' < v(T)$

Formation of dislocation loops in a perfect lattice



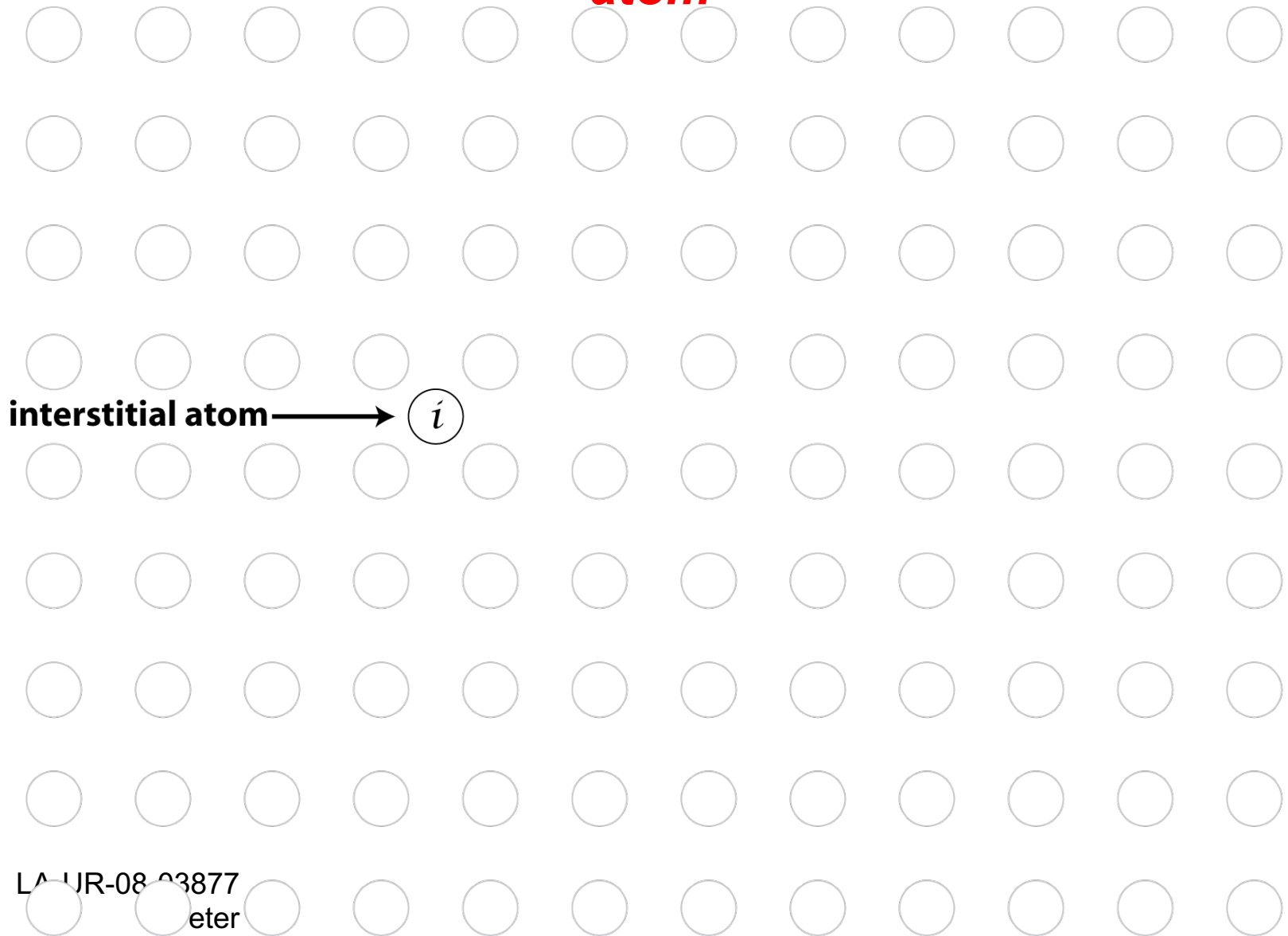
LA-UR-08-03877

eter

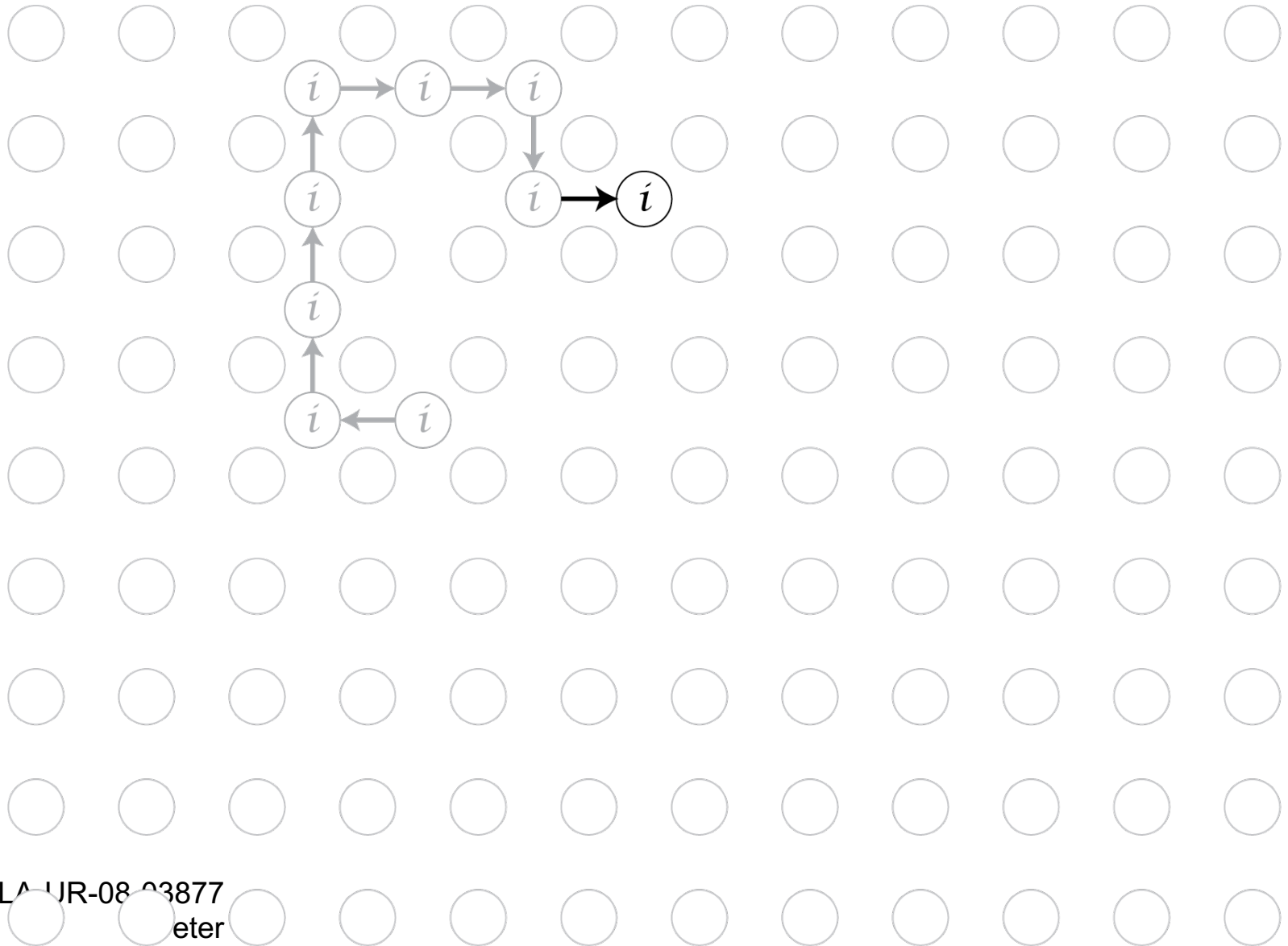
Hosemann

Perfect Lattice + Interstitial Point Defect

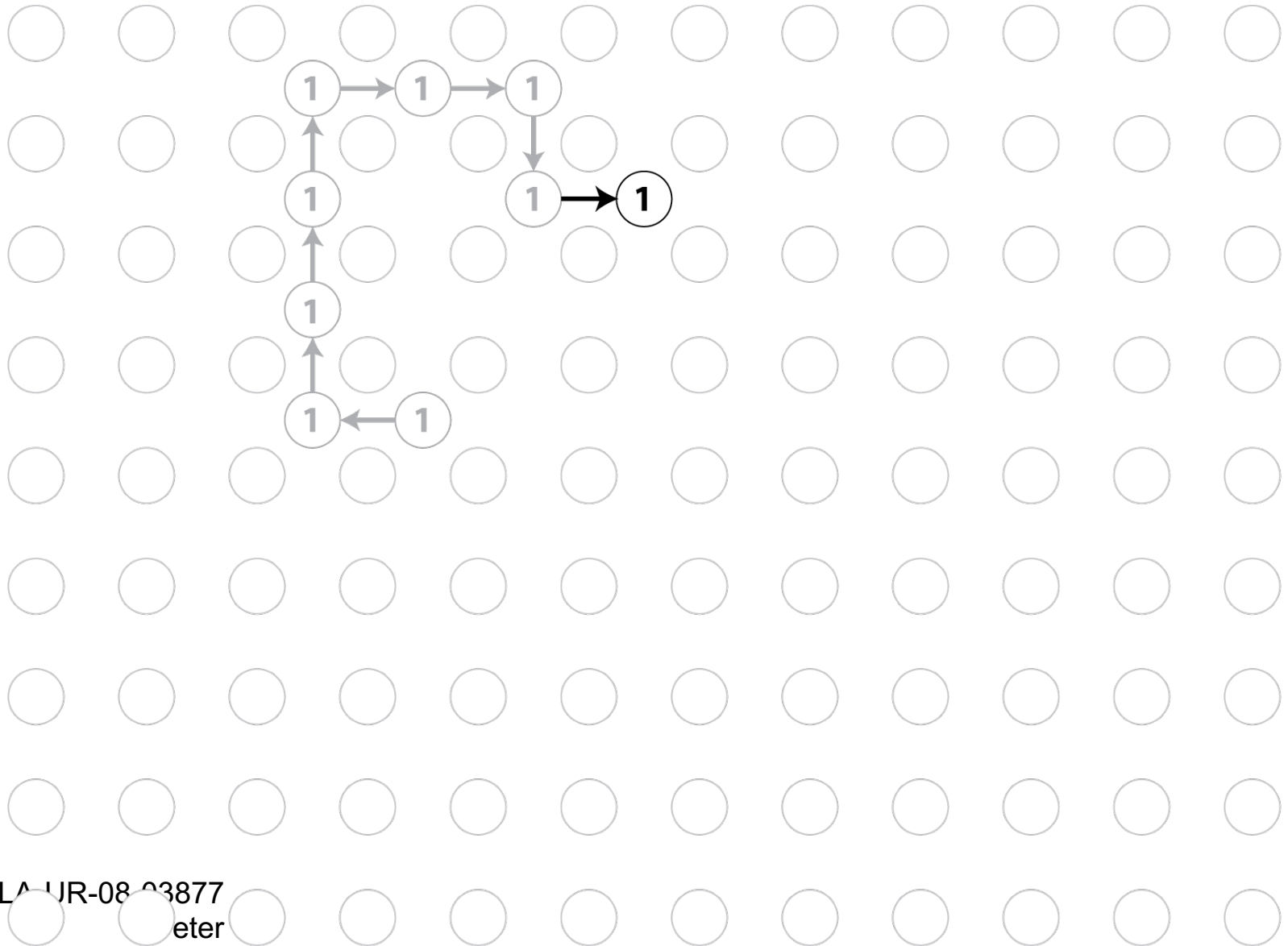
In a monoatomic solid, this would be called a self-interstitial atom



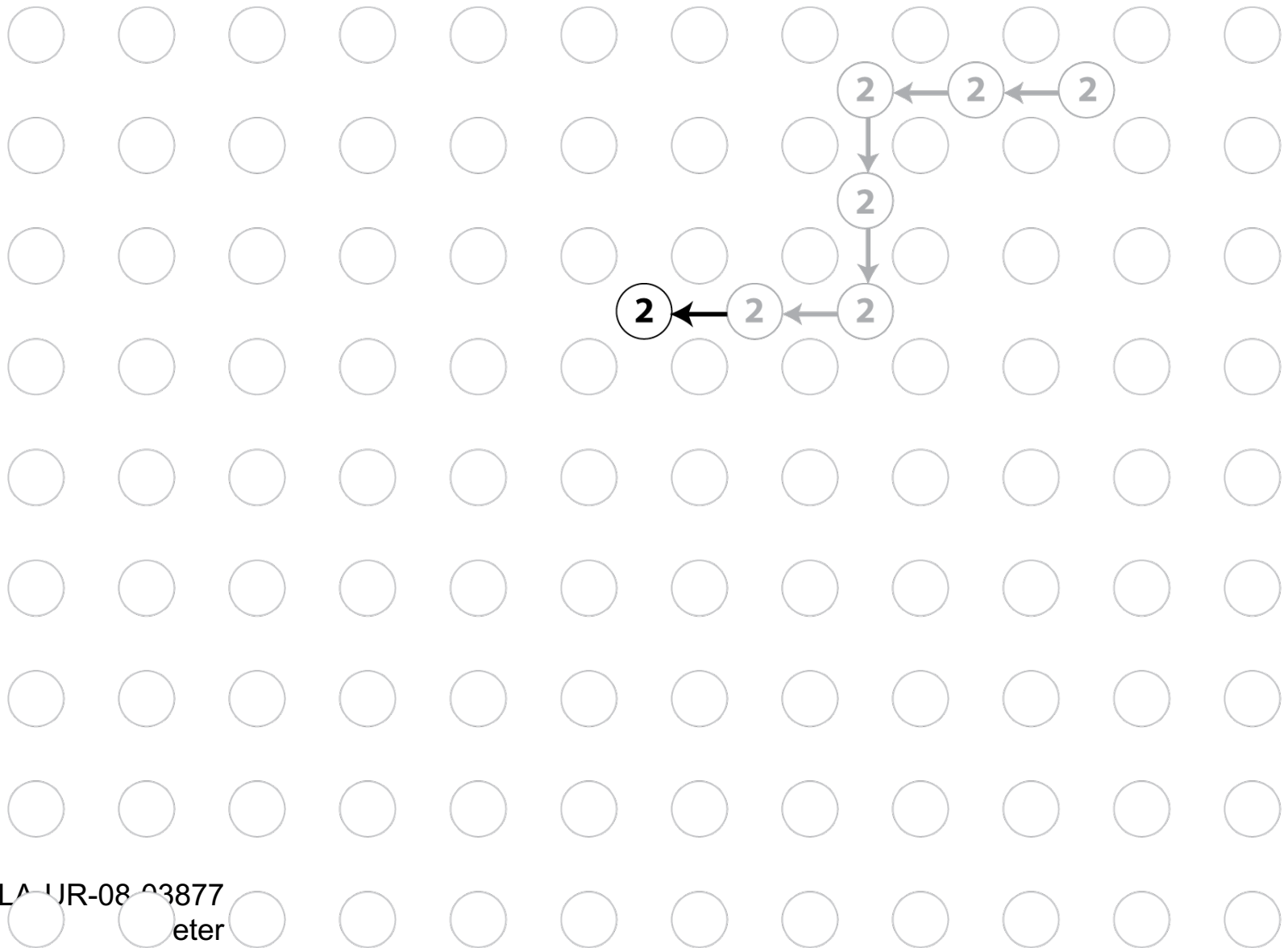
Perfect Lattice + Freely-Migrating Self-Interstitial Atom (SIA)



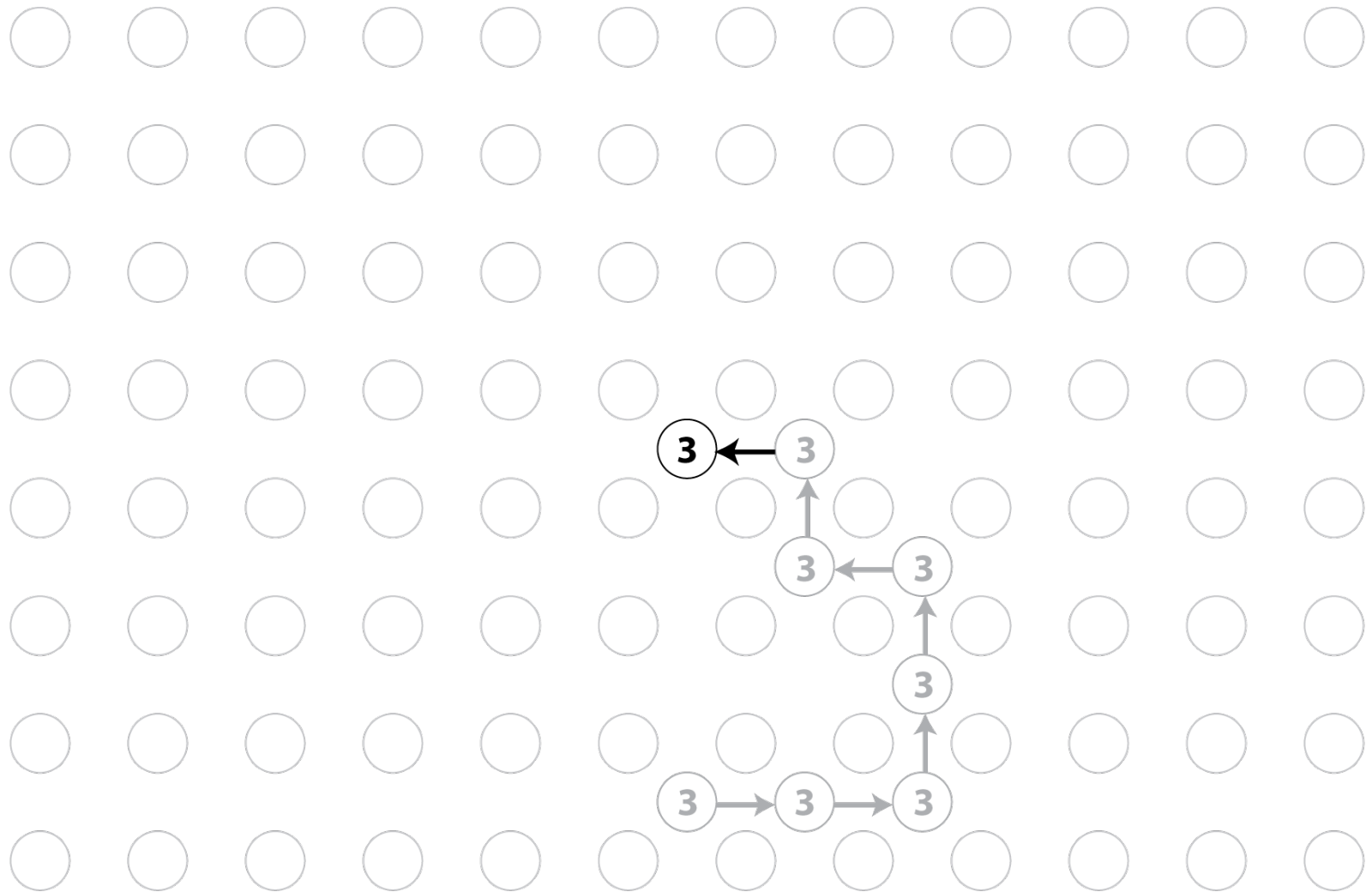
Perfect Lattice + Freely-Migrating SIA #1



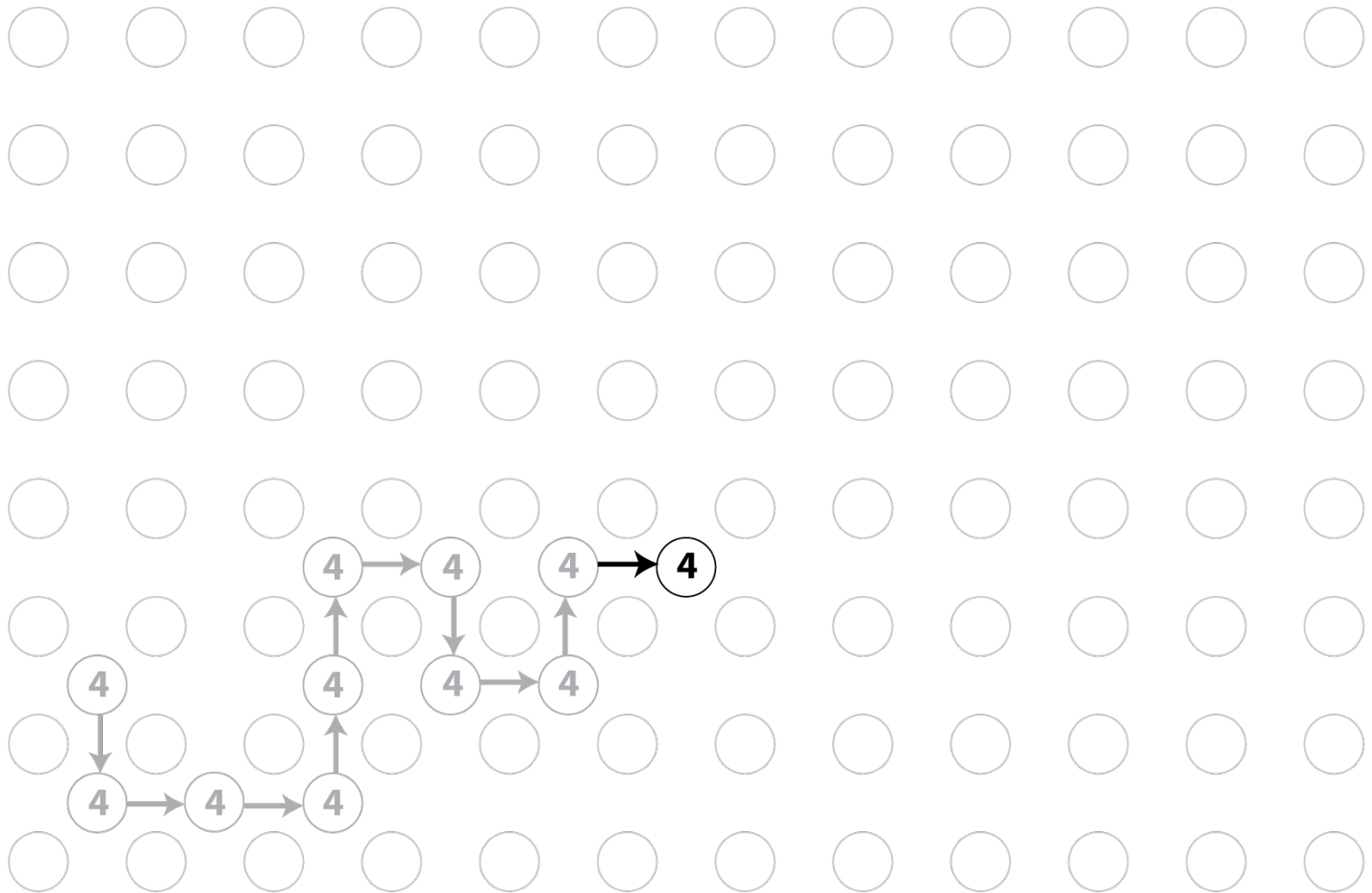
Perfect Lattice + Freely-Migrating SIA #2



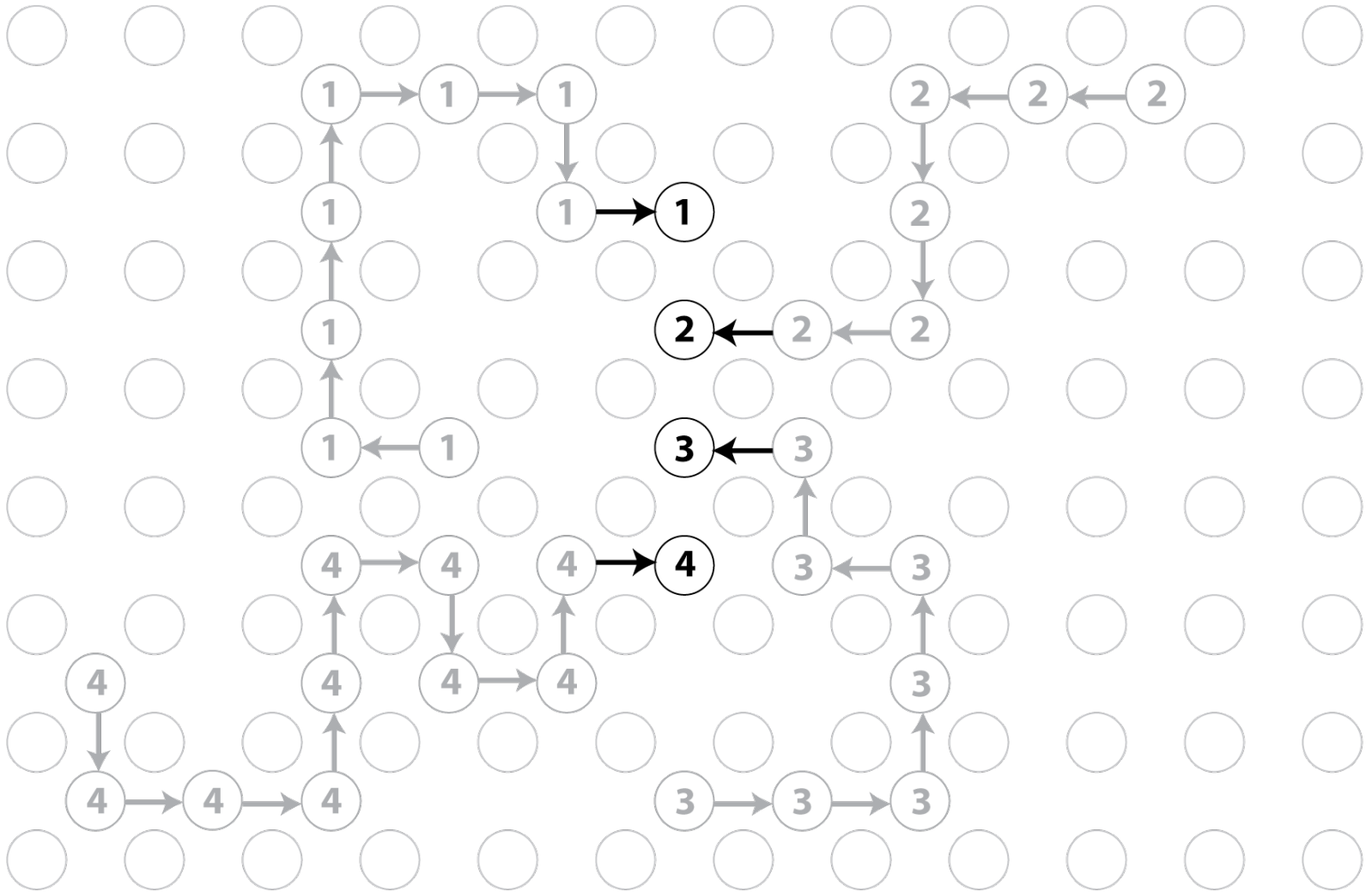
Perfect Lattice + Freely-Migrating SIA #3



Perfect Lattice + Freely-Migrating SIA #4

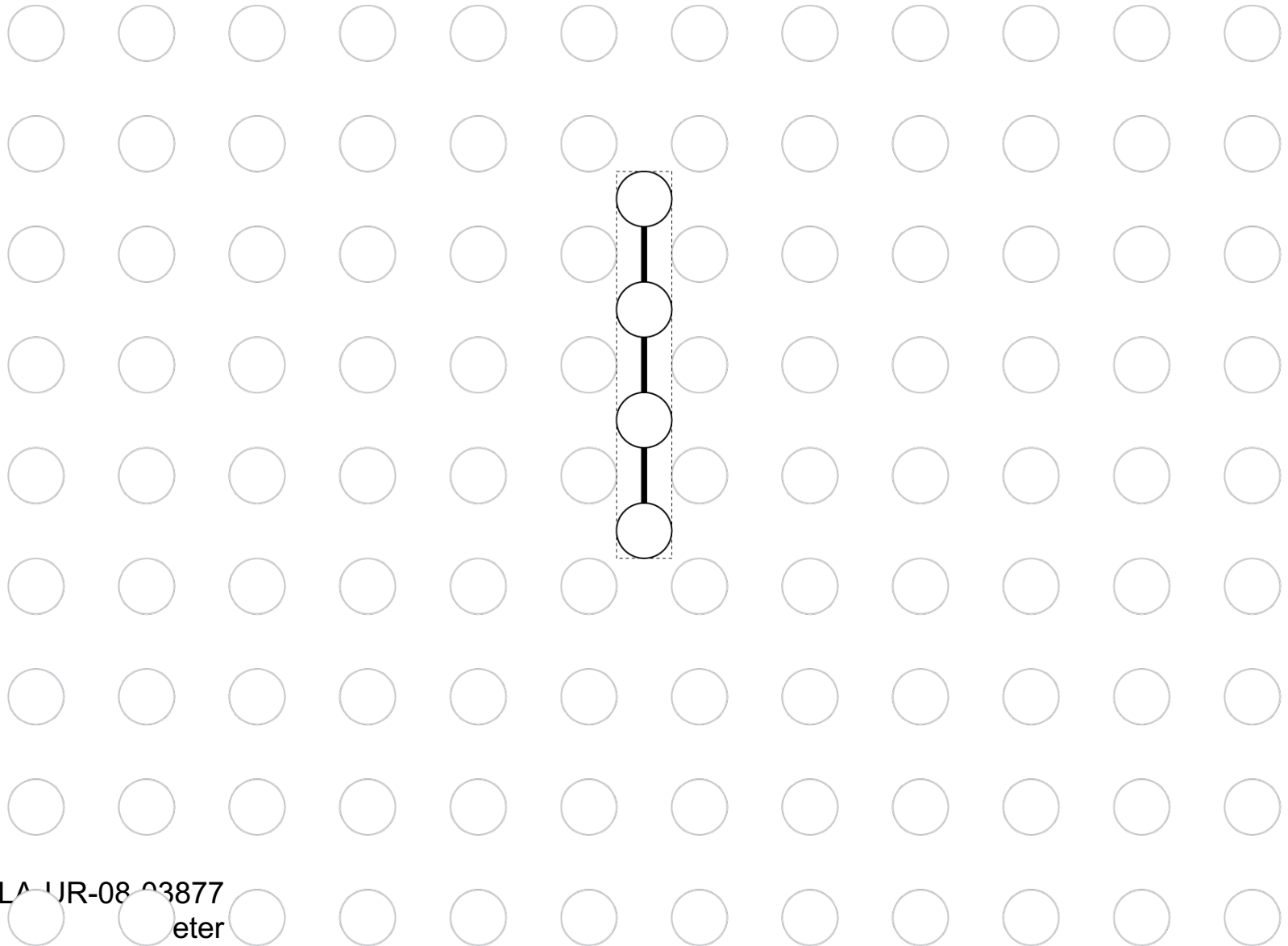


Perfect Lattice + Several Freely-Migrating SIAs

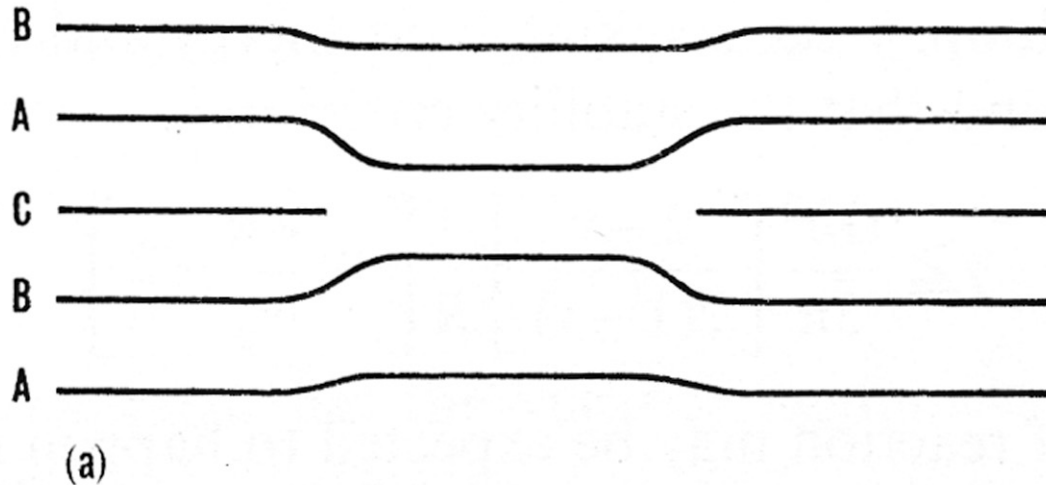


Perfect Lattice + Interstitial Loop

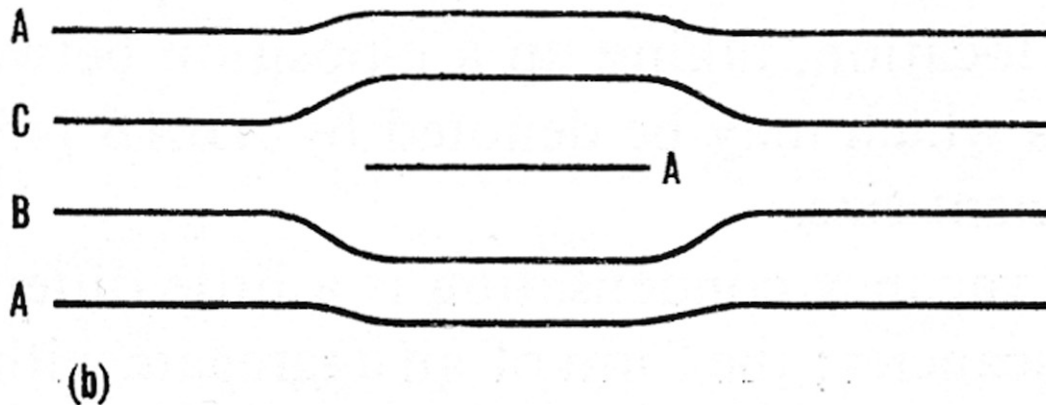
Loop condensed from freely-migrating SIAs



Edge-on View of Loop Dislocation loops



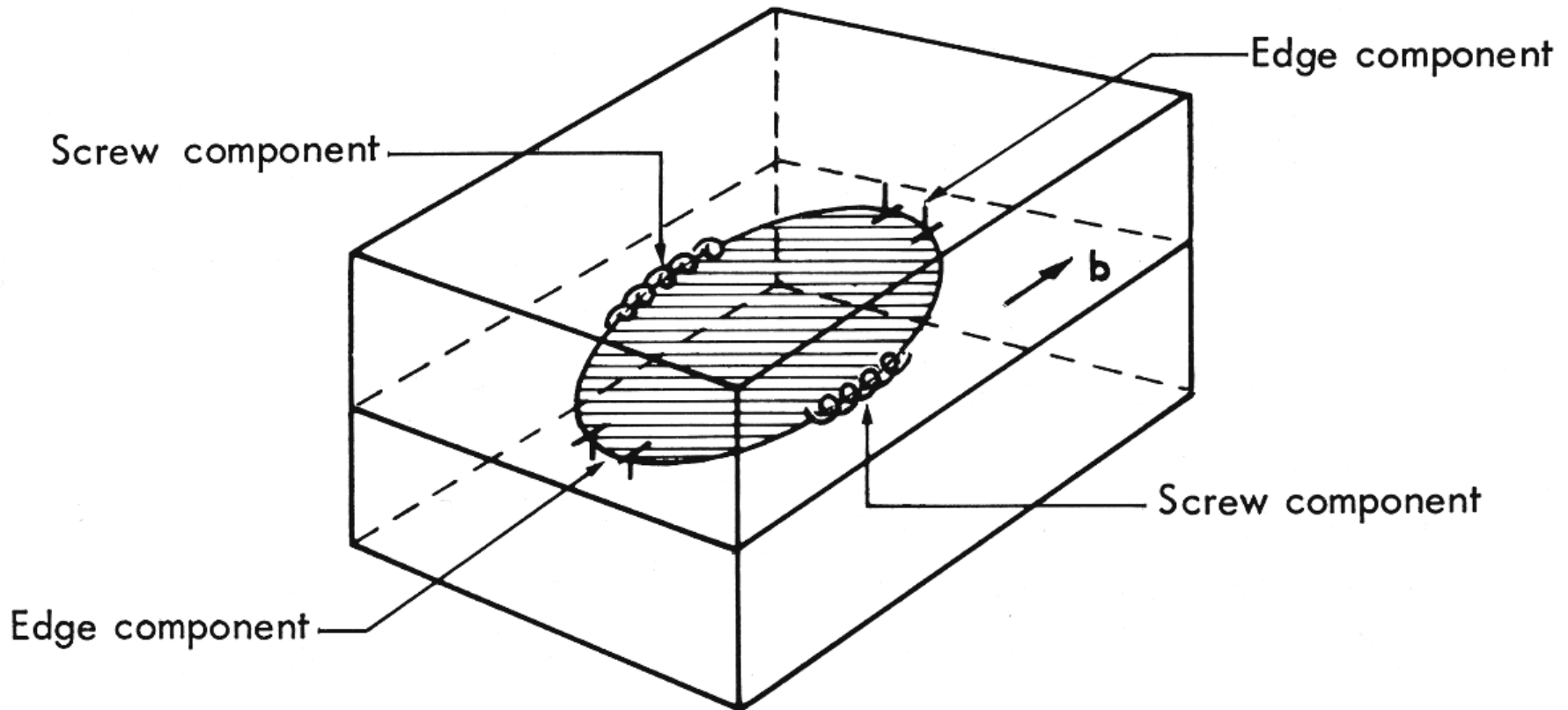
**(a) Condensation
of vacancies to
form a vacancy
loop
(a missing plane of
atoms)**



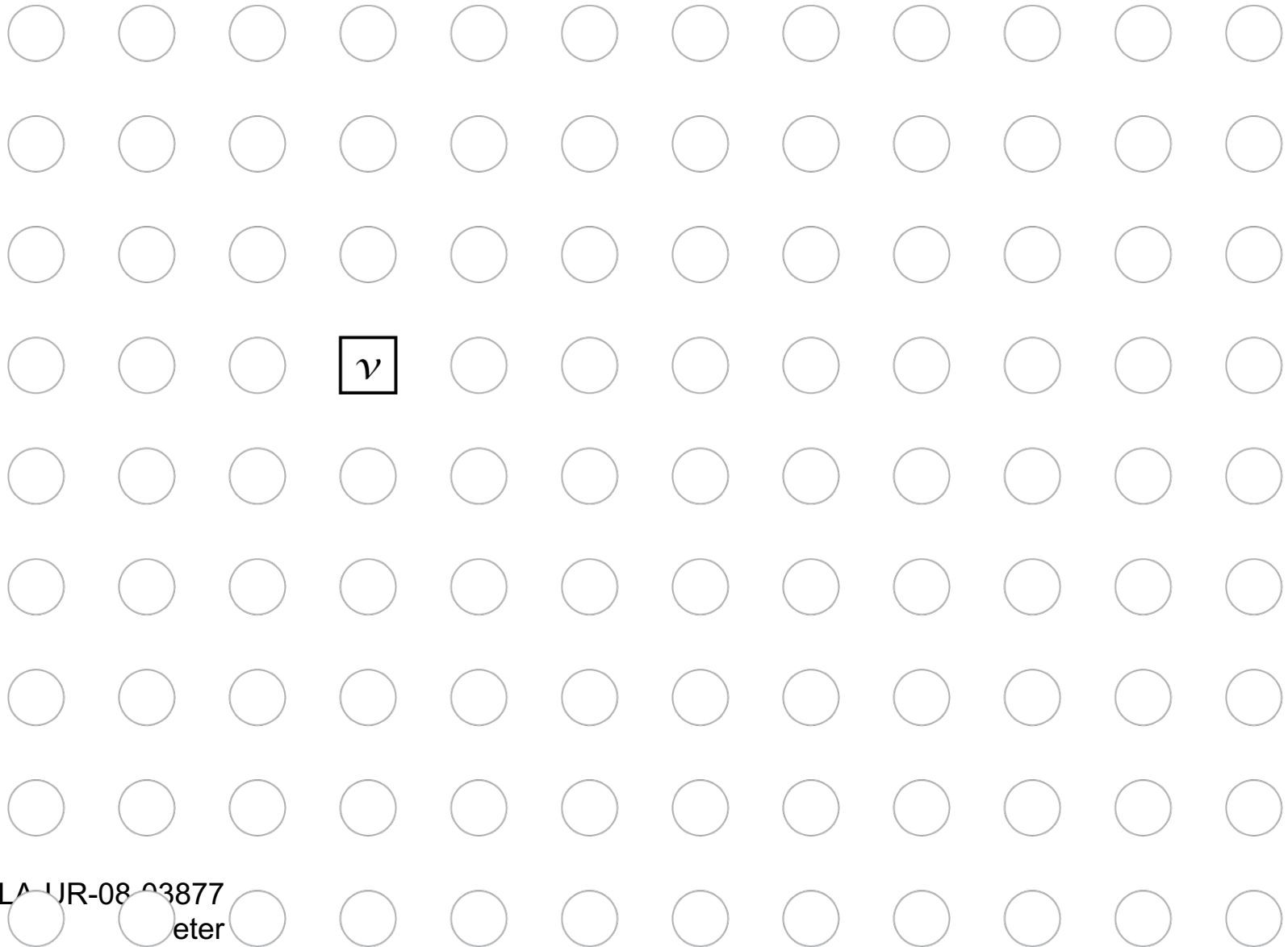
**(b) Condensation
of interstitials to
form an interstitial
loop (an extra
plane of atoms)**

Dislocation Loops II

Top-View of Loop



Void formation: Perfect Lattice + Vacancy Point Defect

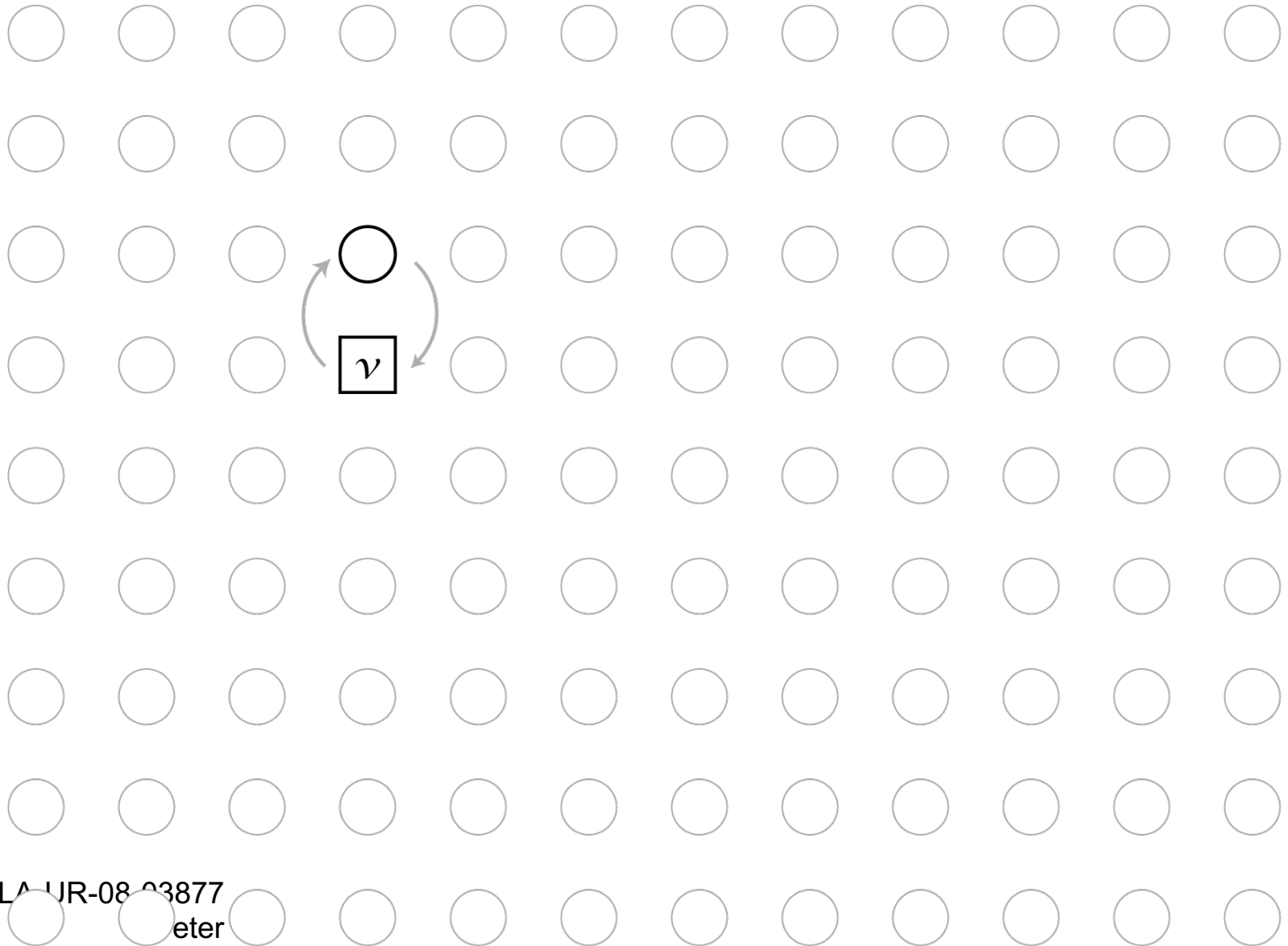


LA-UR-08-03877

eter

Hosemann

Vacancy Before Migration

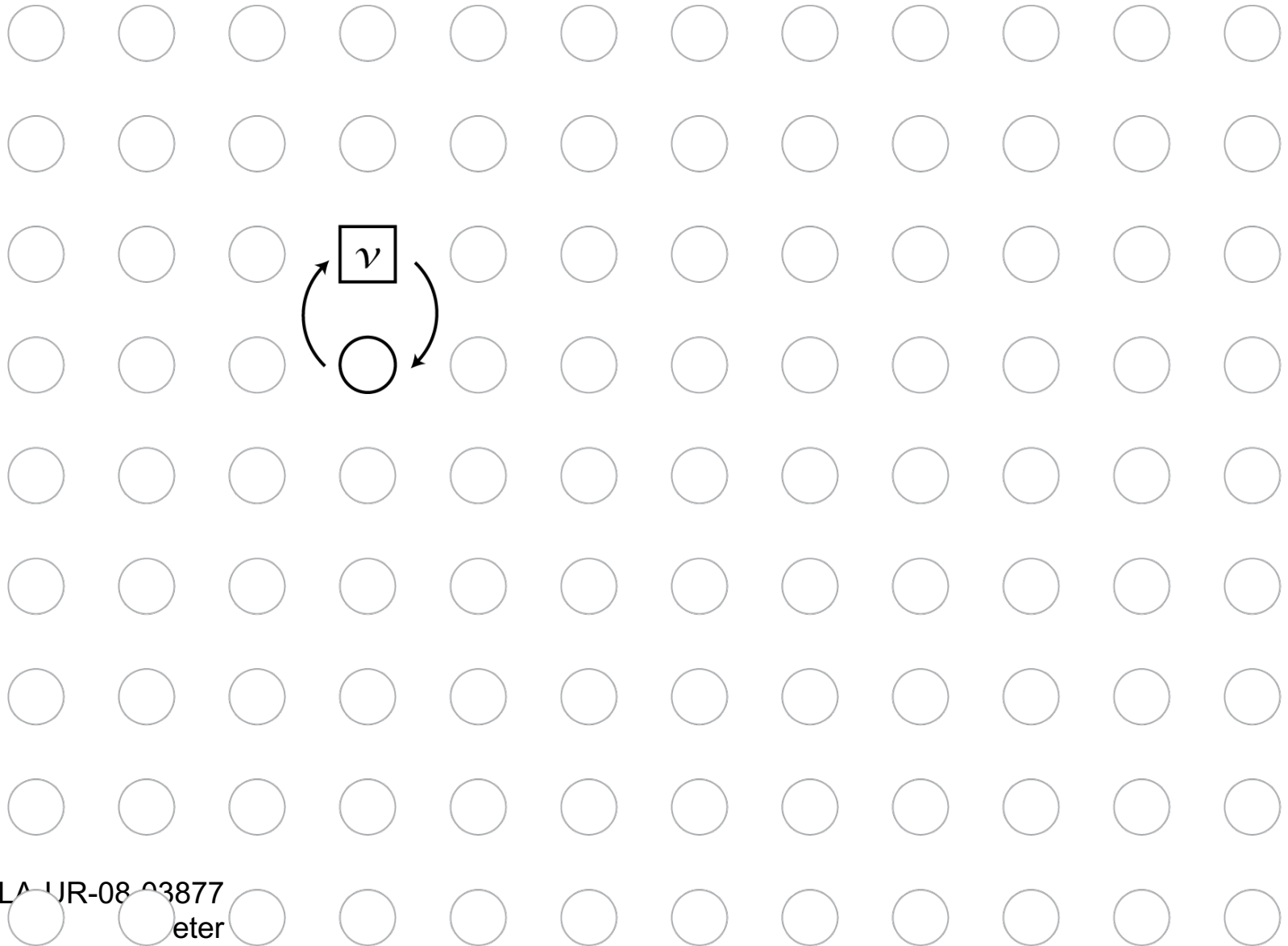


LA-UR-08-03877

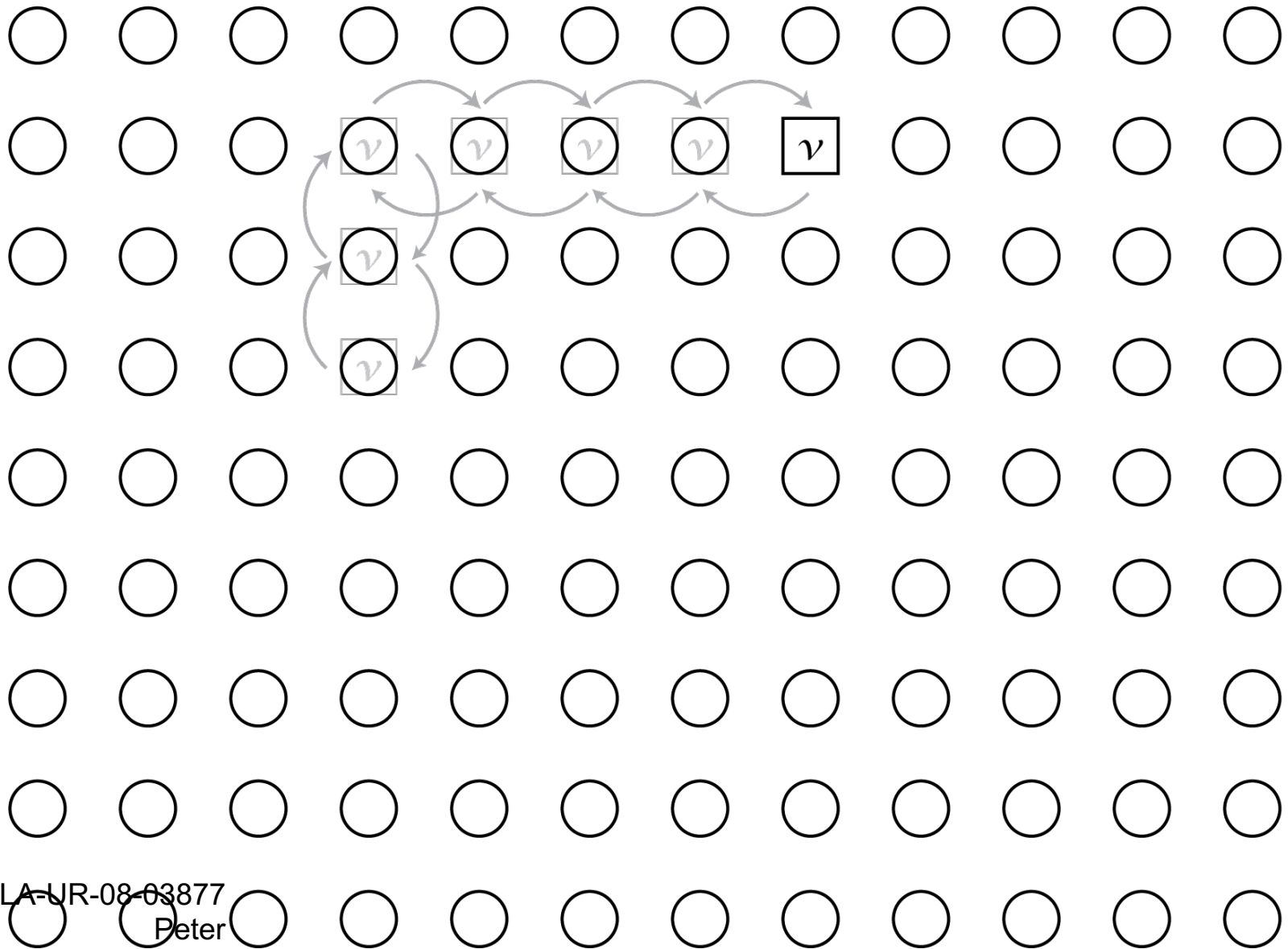
eter

Hosemann

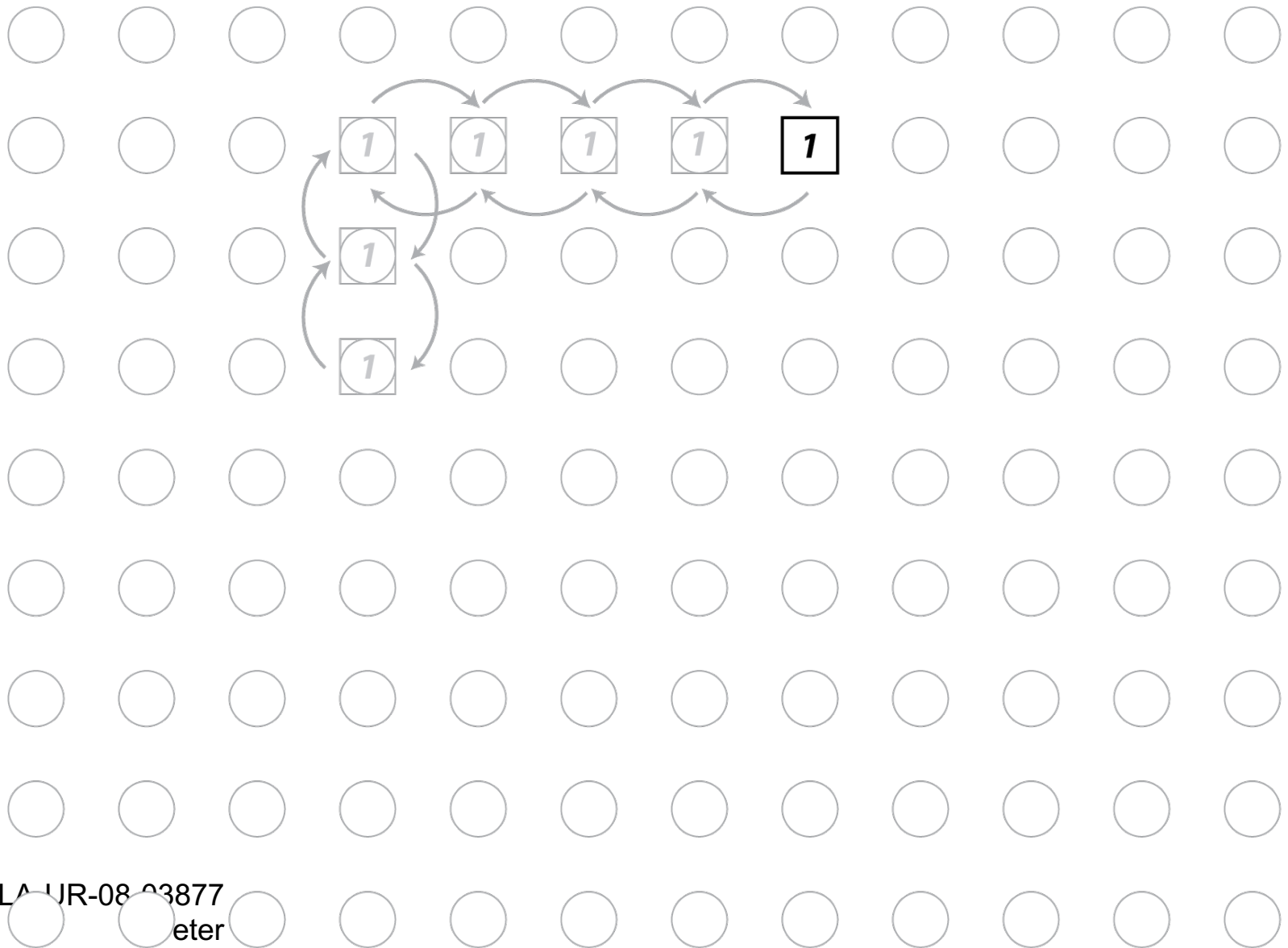
Vacancy After Migration



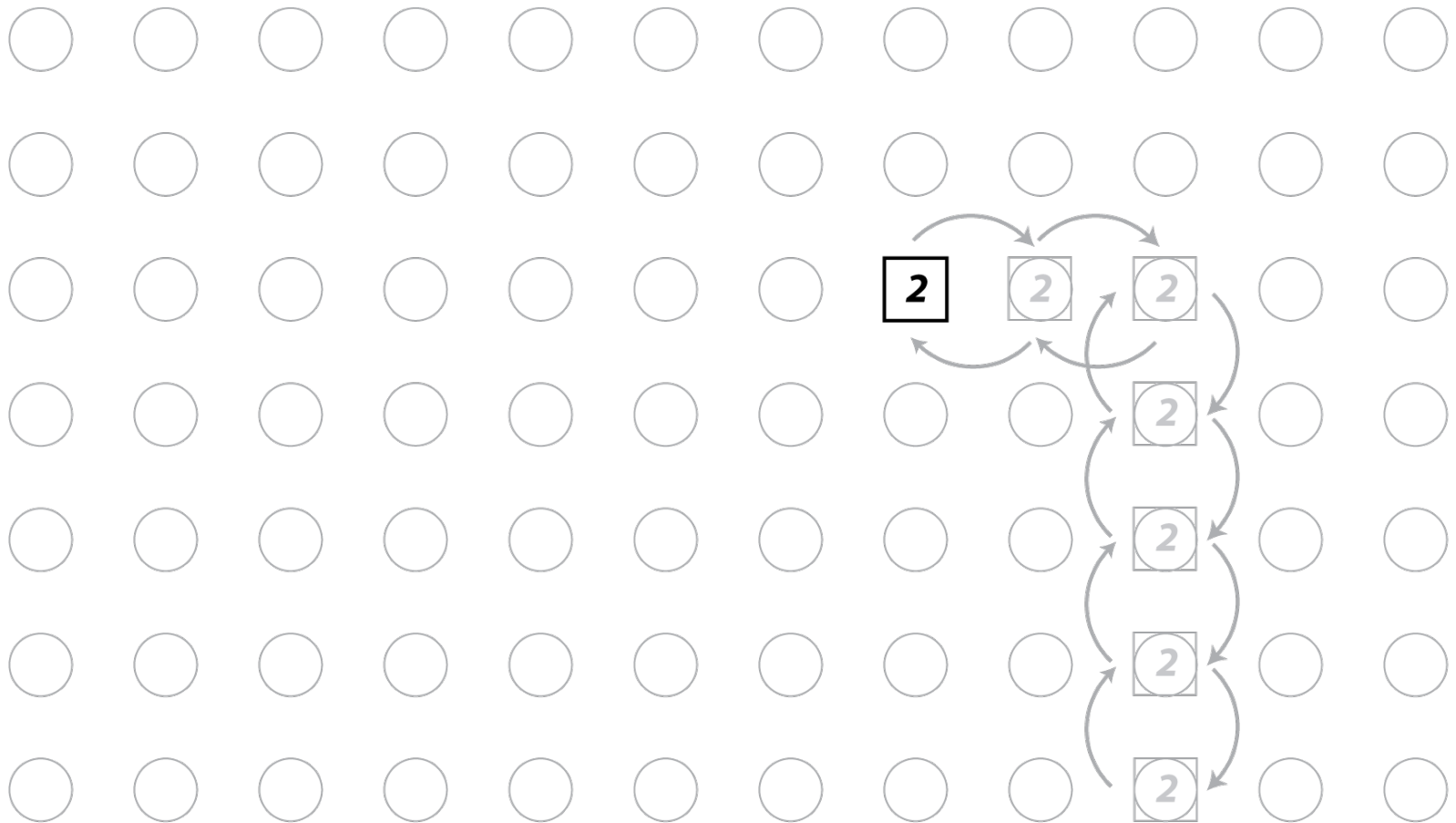
Perfect Lattice + Freely-Migrating Vacancy



Perfect Lattice + Freely-Migrating Vacancy **#1**

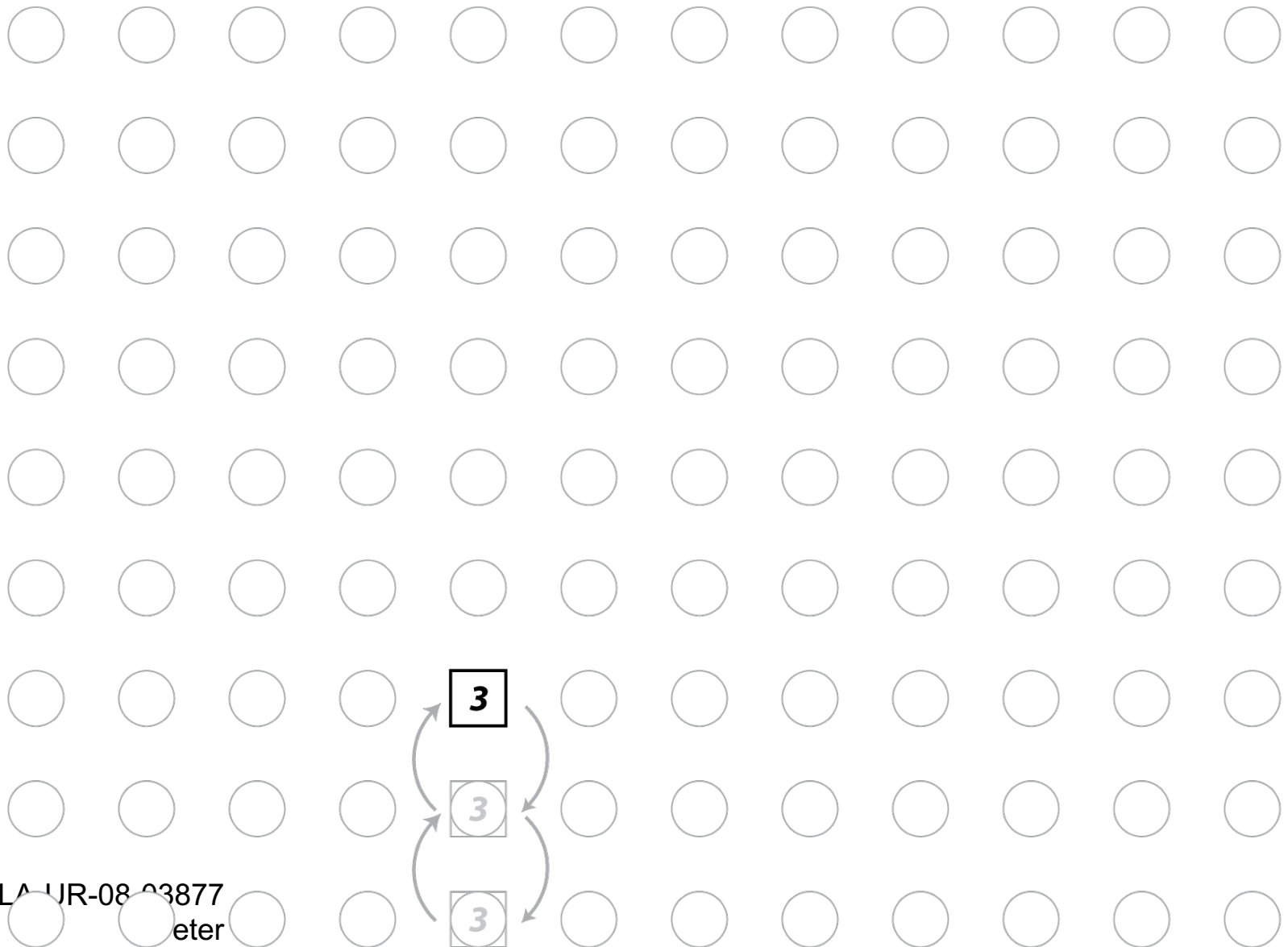


Perfect Lattice + Freely-Migrating Vacancy #2



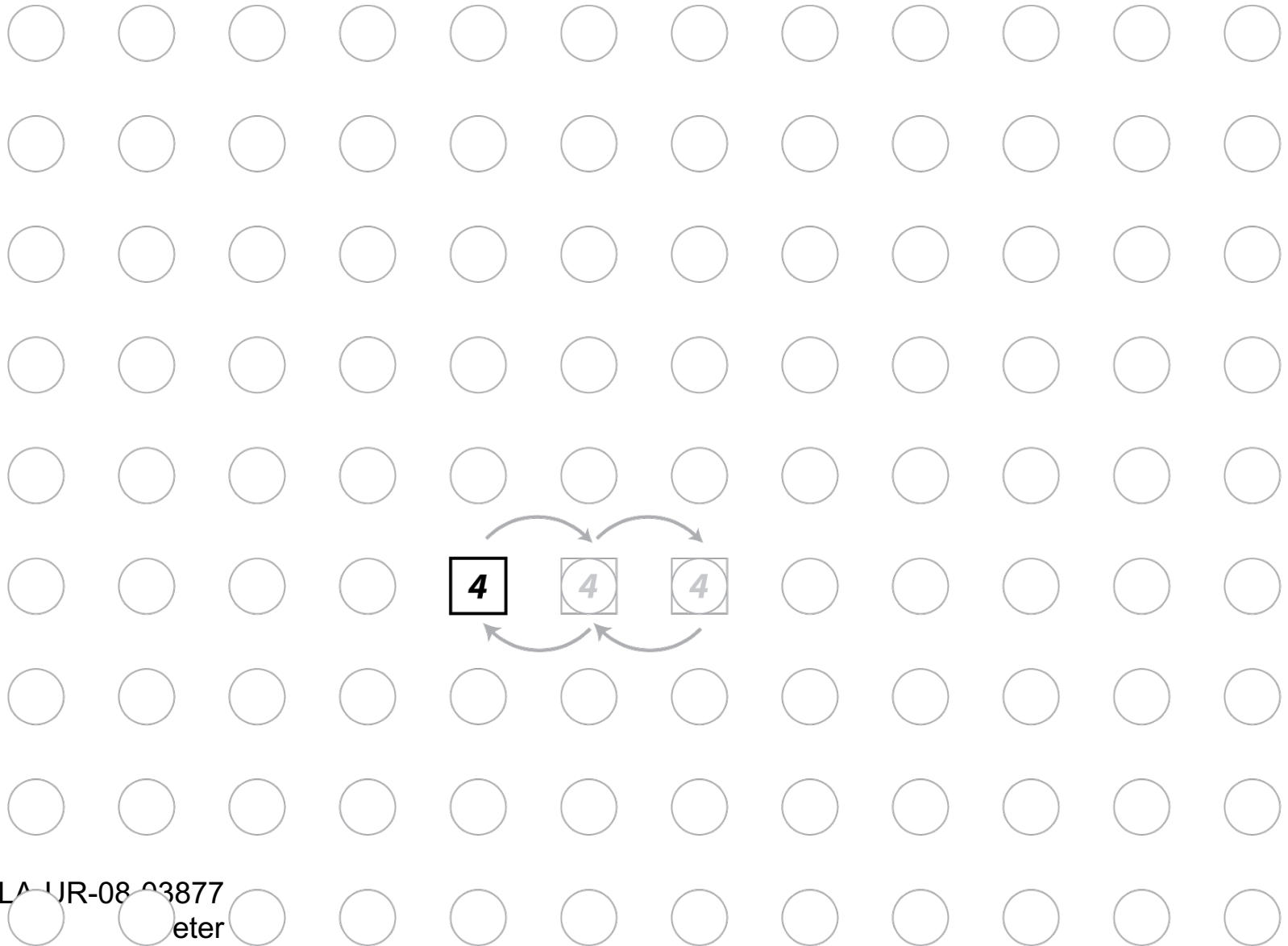
Perfect Lattice + Freely-Migrating Vacancy

#3

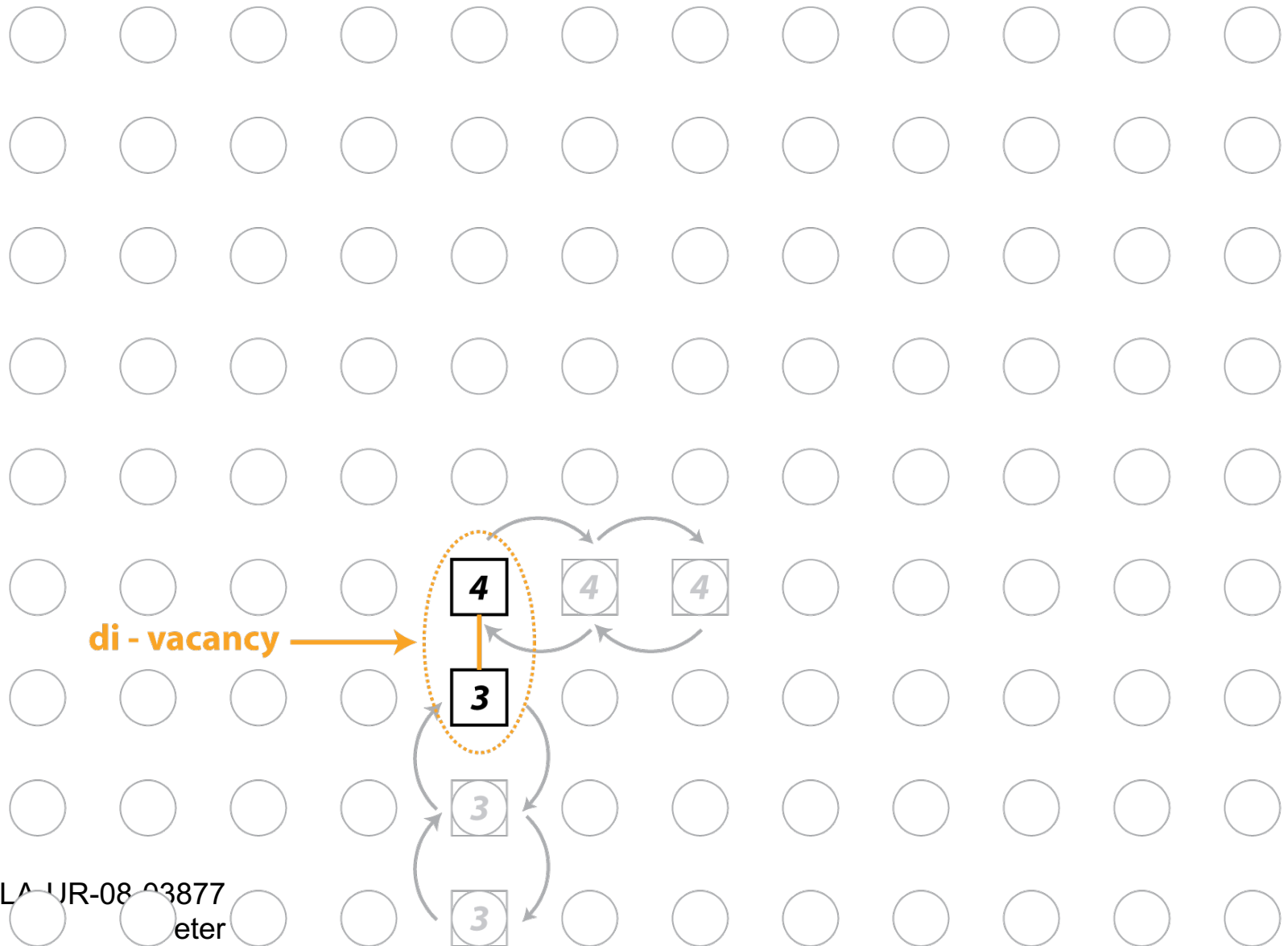


Perfect Lattice + Freely-Migrating Vacancy

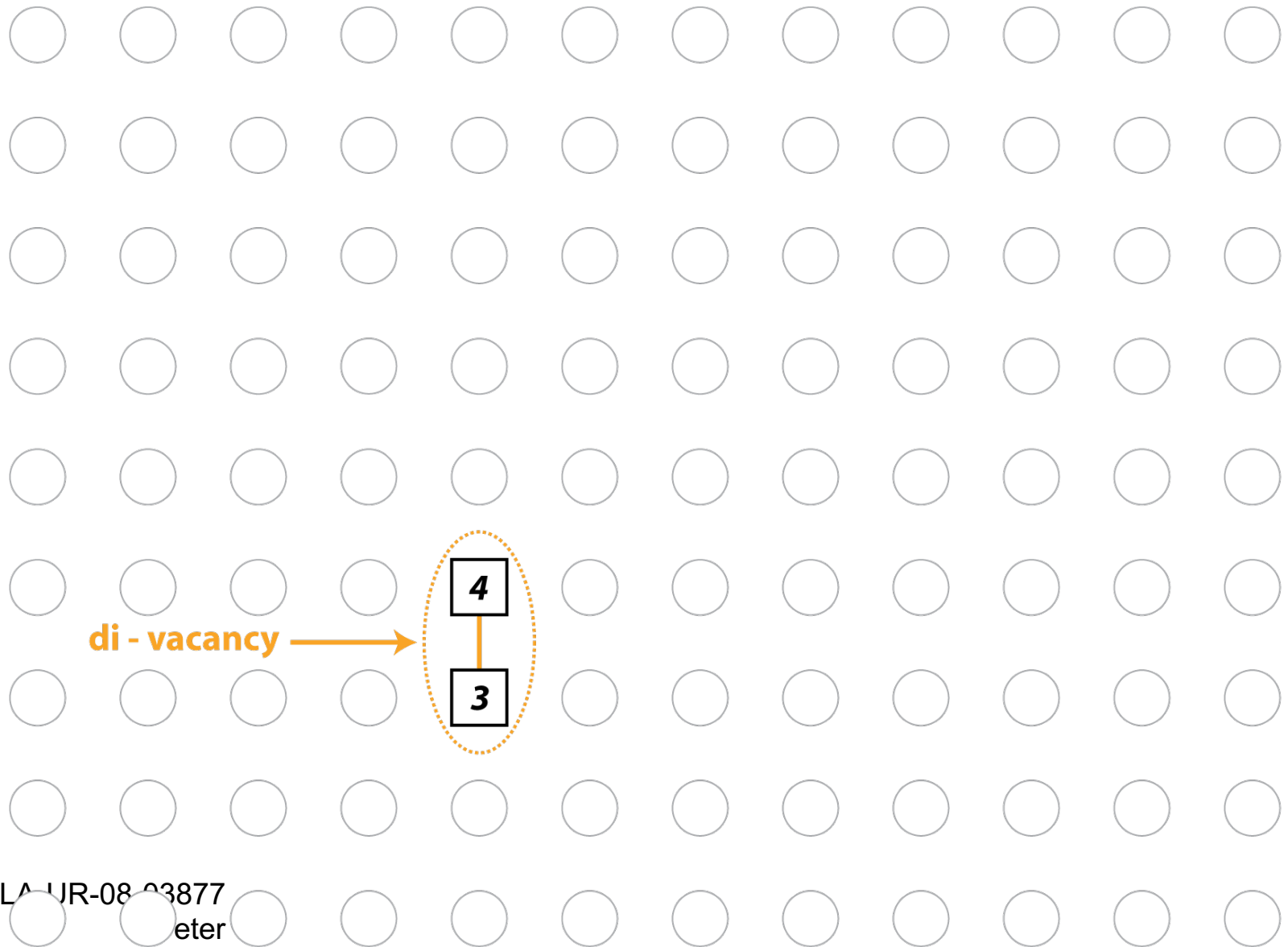
#4



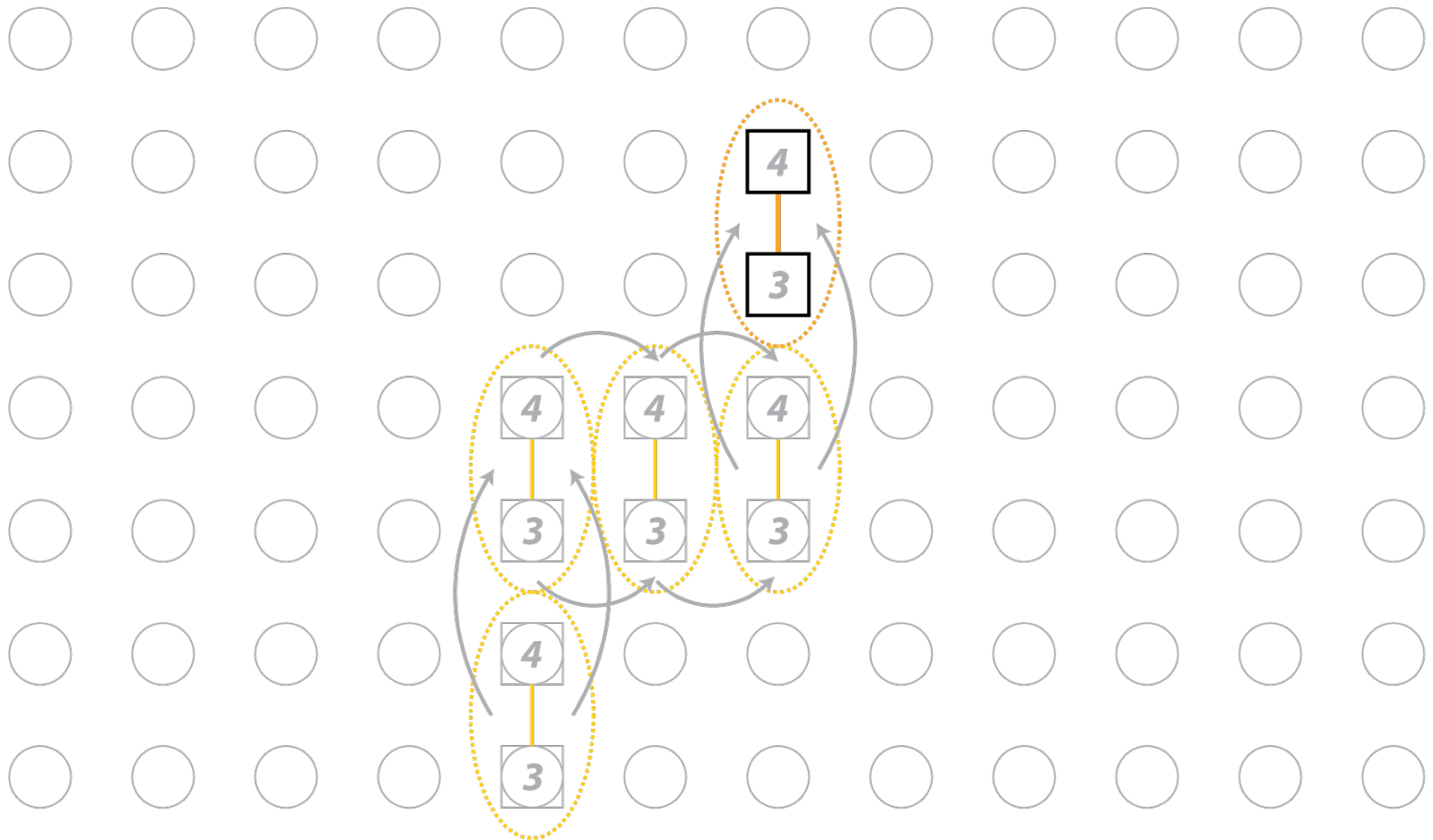
Di - Vacancy Formation



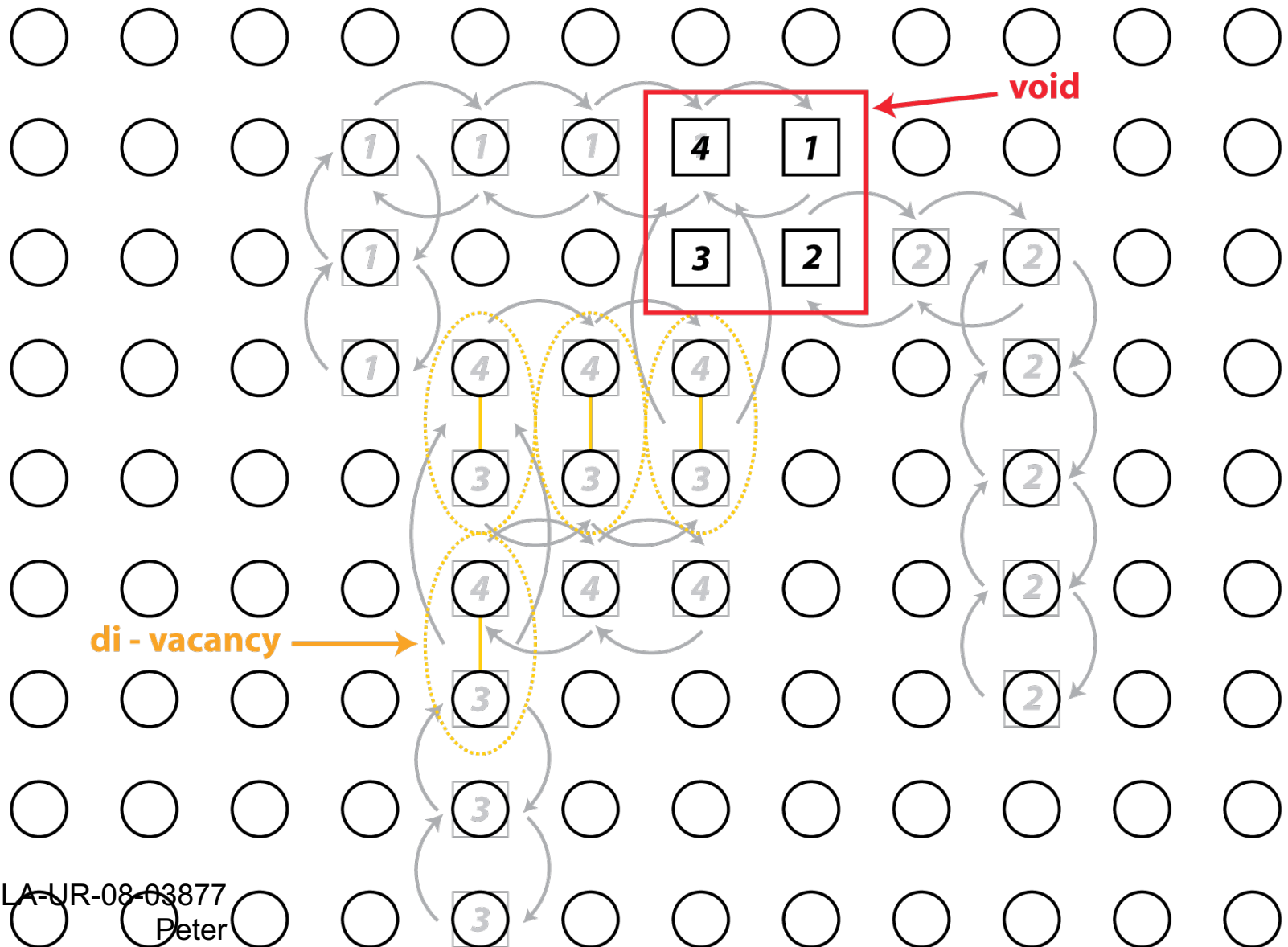
Di - Vacancy Formation



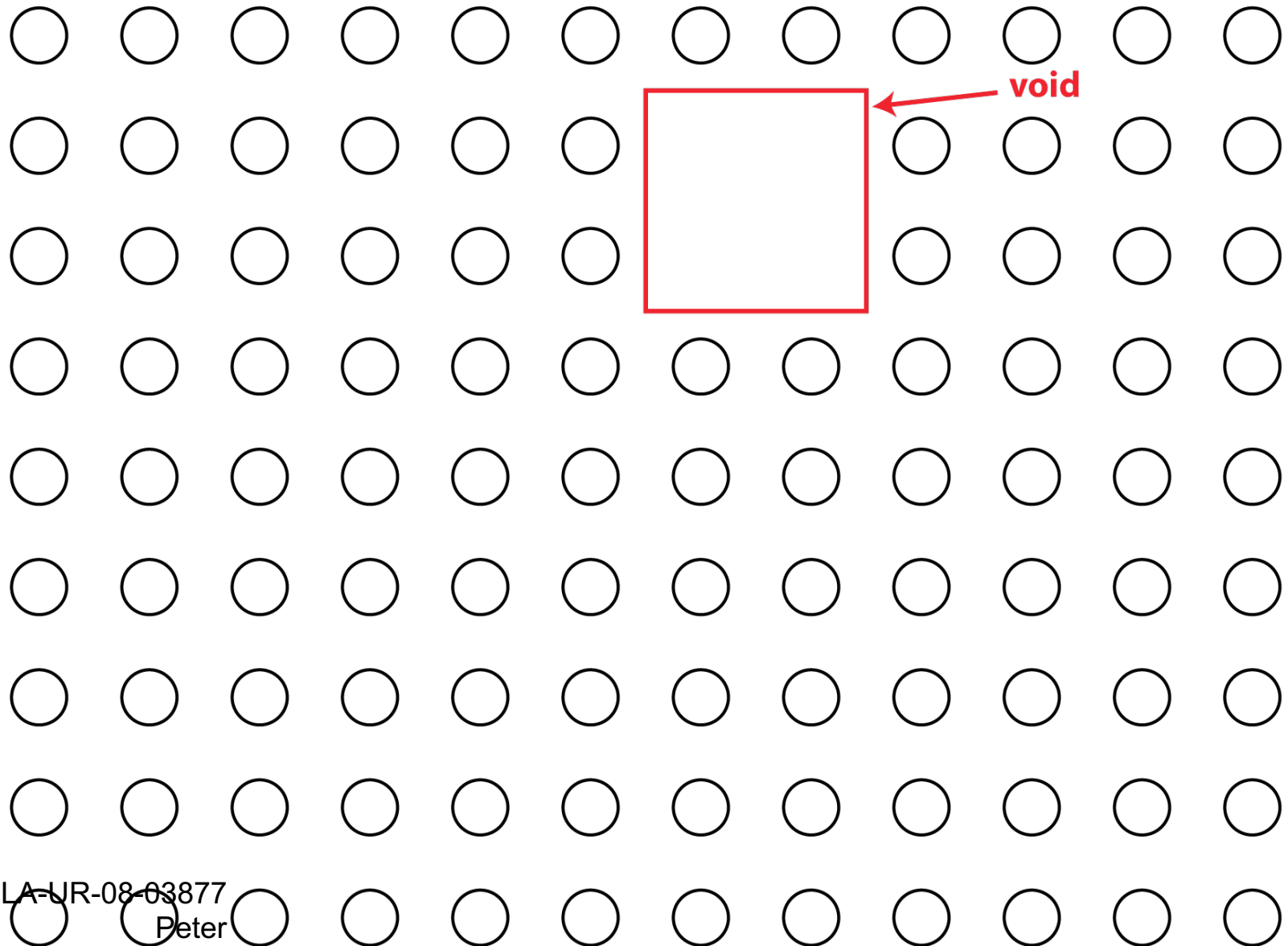
Di - Vacancy Migration



Freely-Migrating Vacancy Point Defects - Void Nucleation



Freely-Migrating Vacancy Point Defects - Void Formation



LA-UR-08-03877

Peter

Hosemann

Void swelling effects

