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Comb-assisted Spectroscopy

Paolo DE NATALE

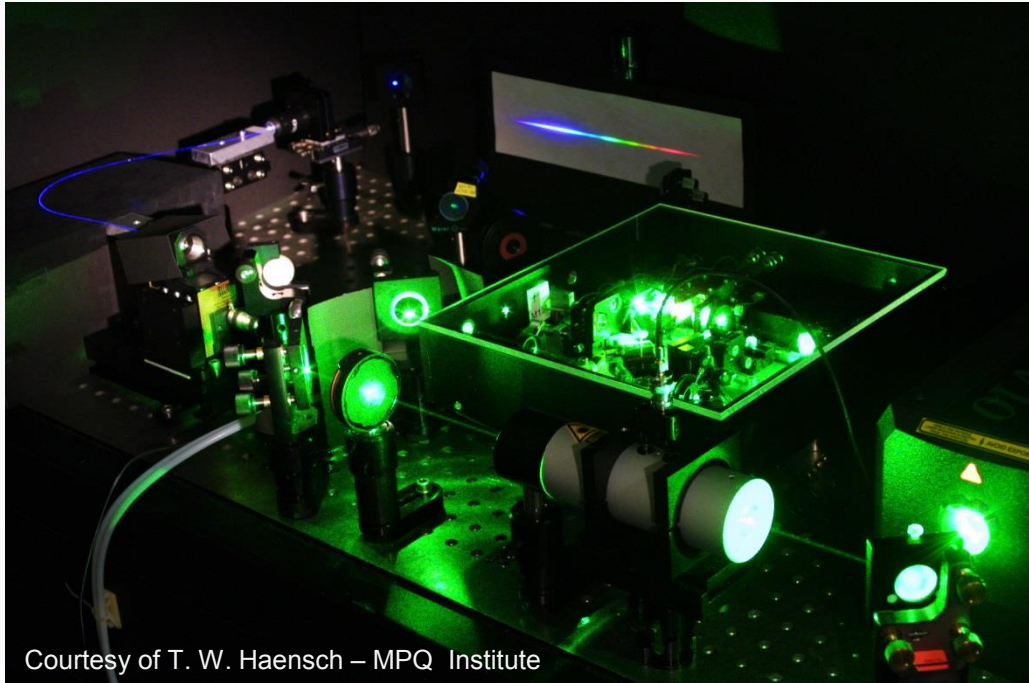


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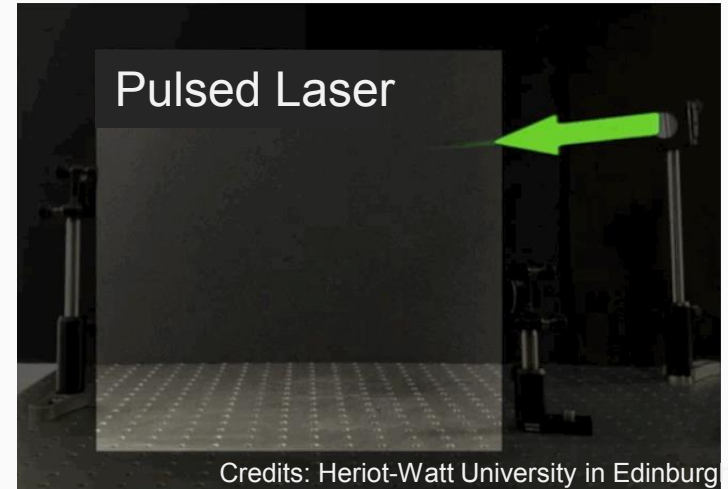
The laser frequency Comb



Theodor W. Haensch
Nobel Prize 2005



Courtesy of T. W. Haensch – MPQ Institute



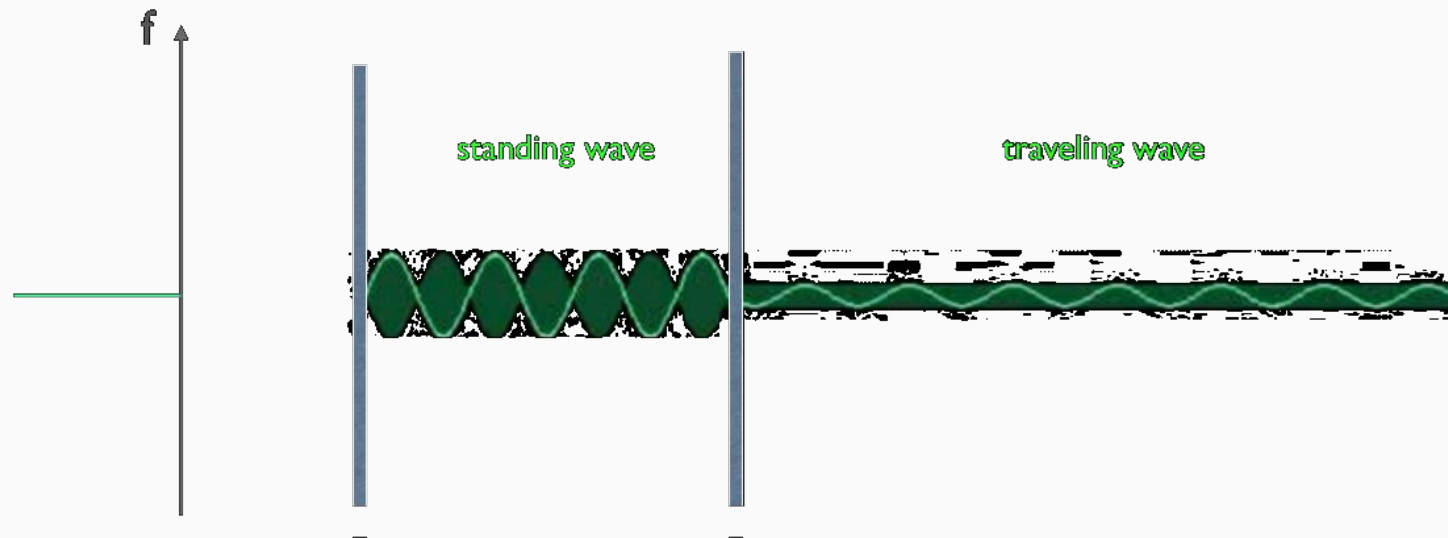
Credits: Heriot-Watt University in Edinburg

- A simple instrument to measure frequencies from 0.1 THz to 1000 THz.
- A phase-coherent link between the Optical and RF regions
- A clockwork to transfer the resonant frequency of Atomic Clocks



Comb working principle

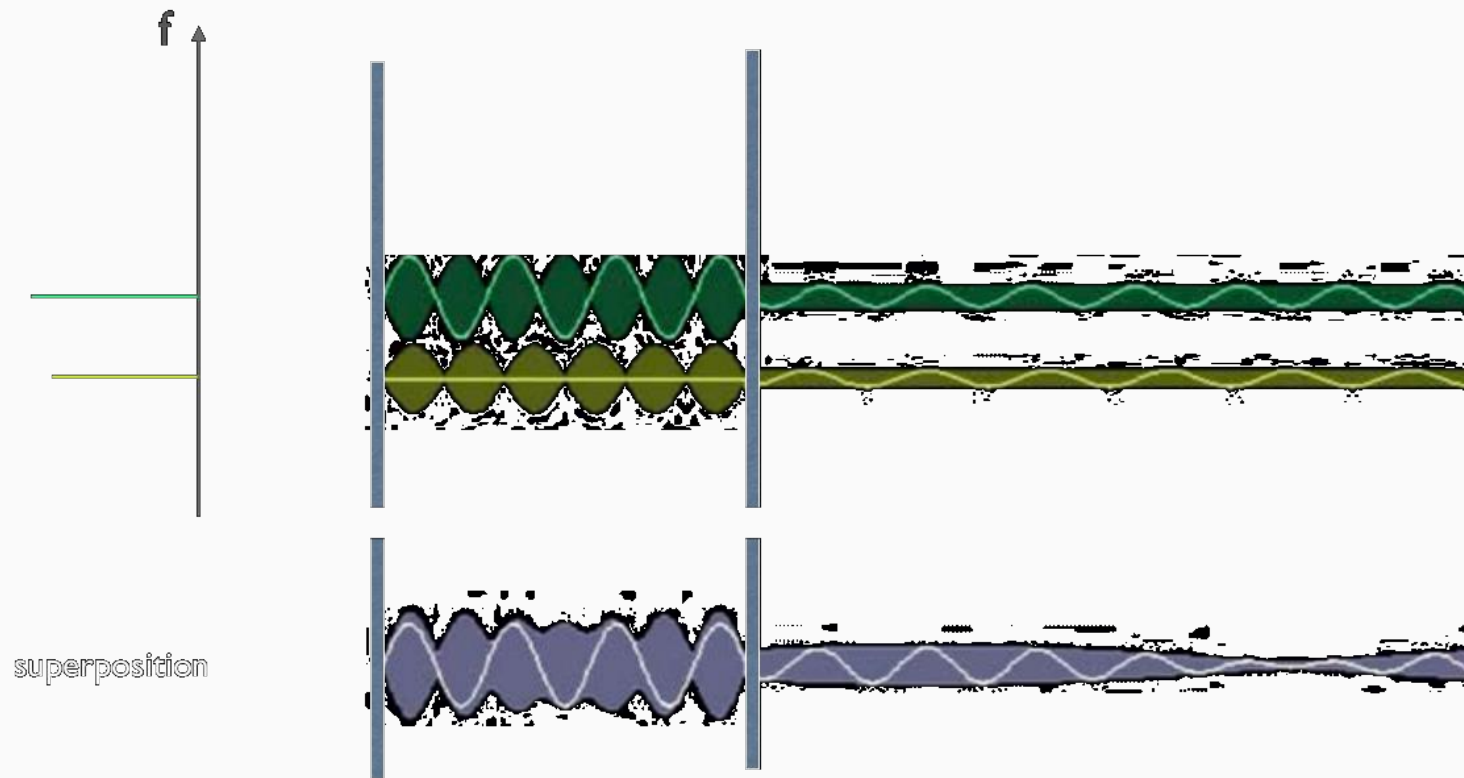
single mode





Comb working principle

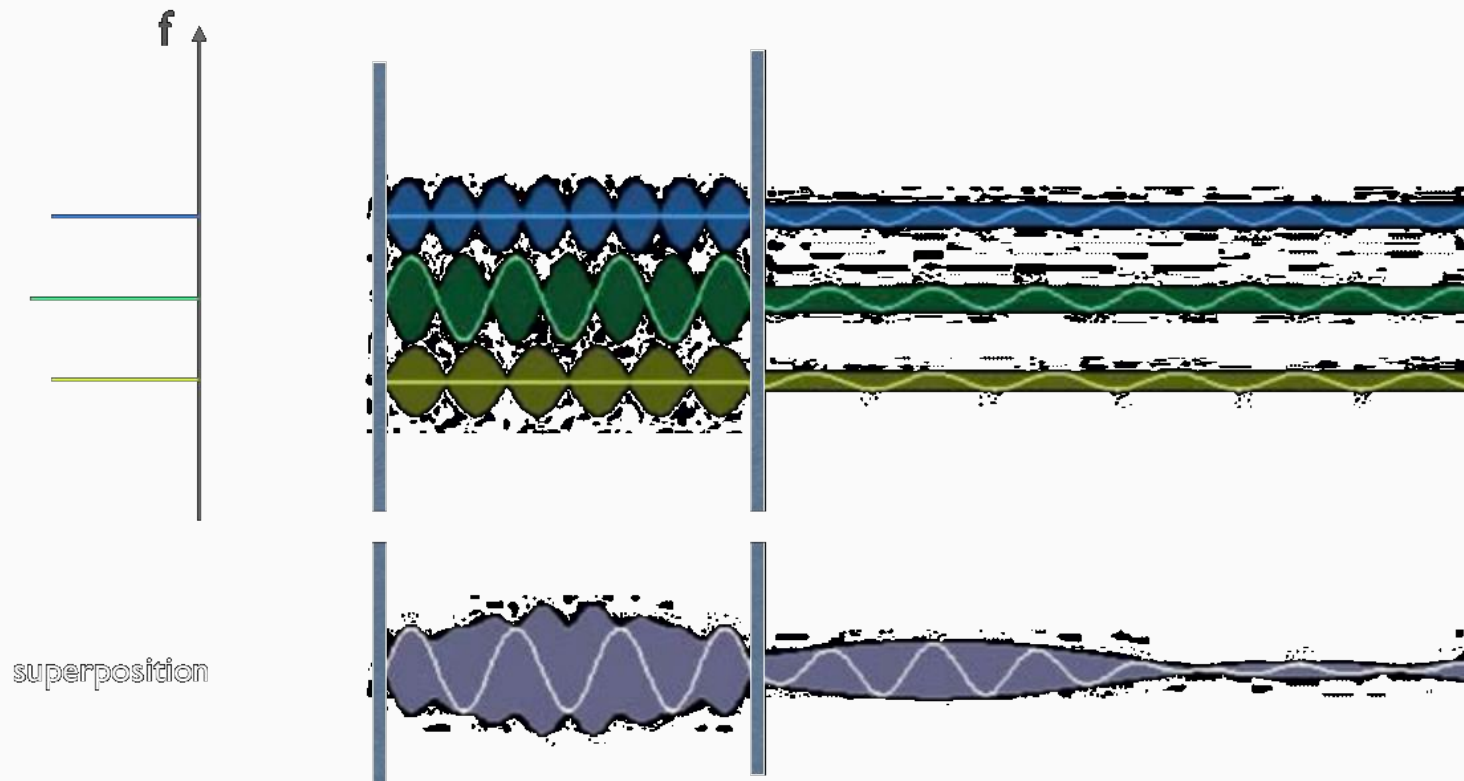
two modes





Comb working principle

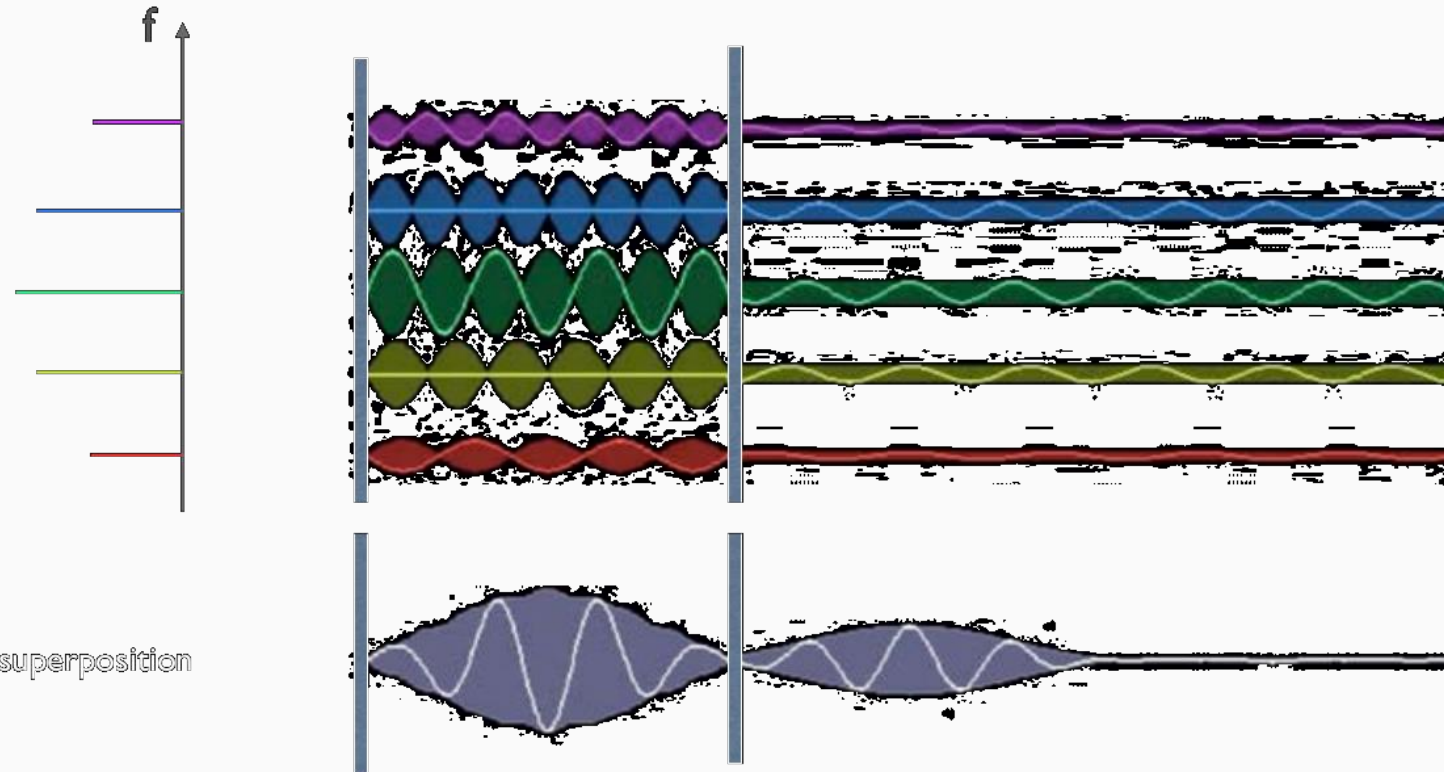
three modes





Comb working principle

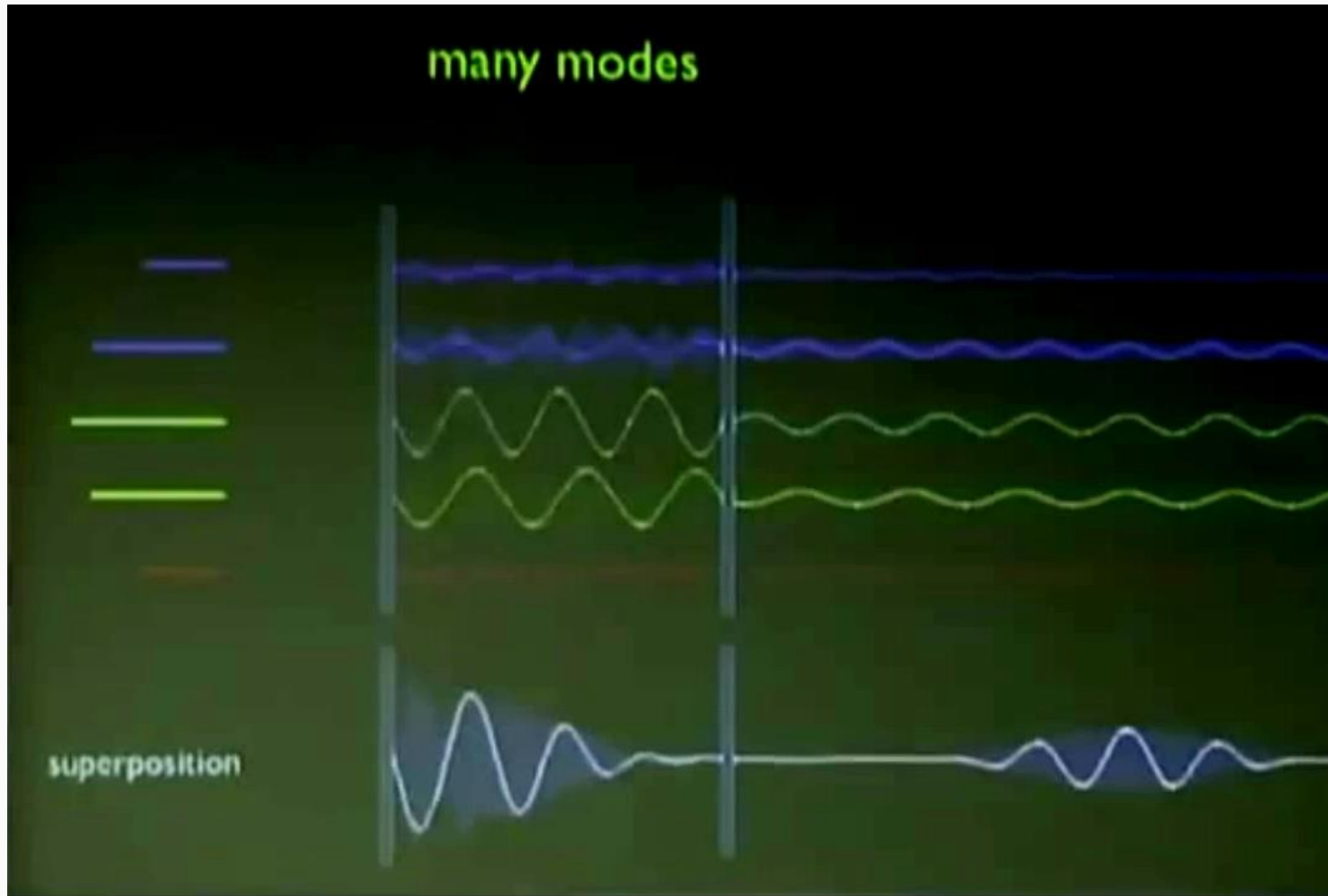
many modes





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Comb working principle

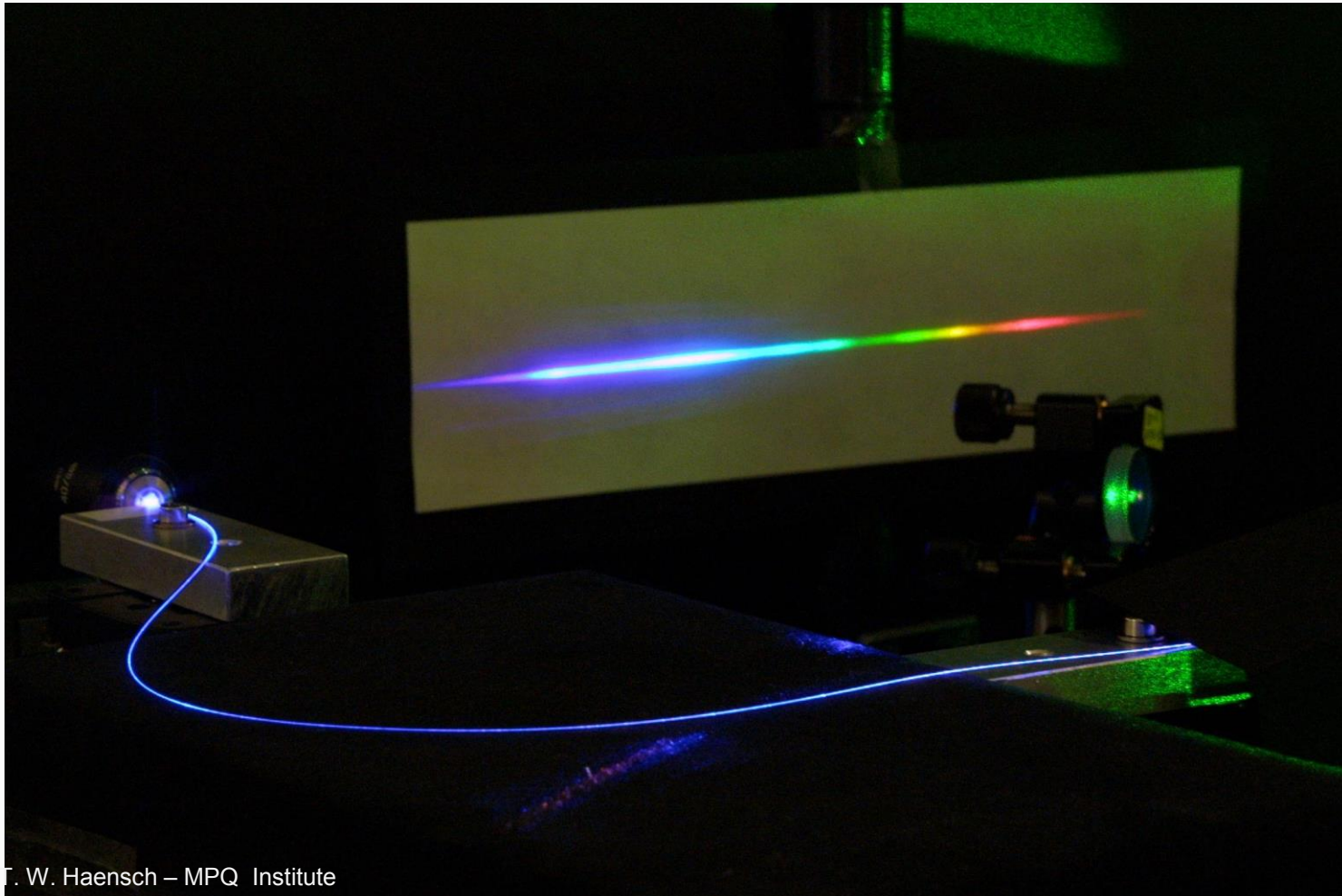


Courtesy of T. W. Haensch – MPQ Institute



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Comb working principle



Courtesy of T. W. Haensch – MPQ Institute

Principles of Optical Frequency Combs

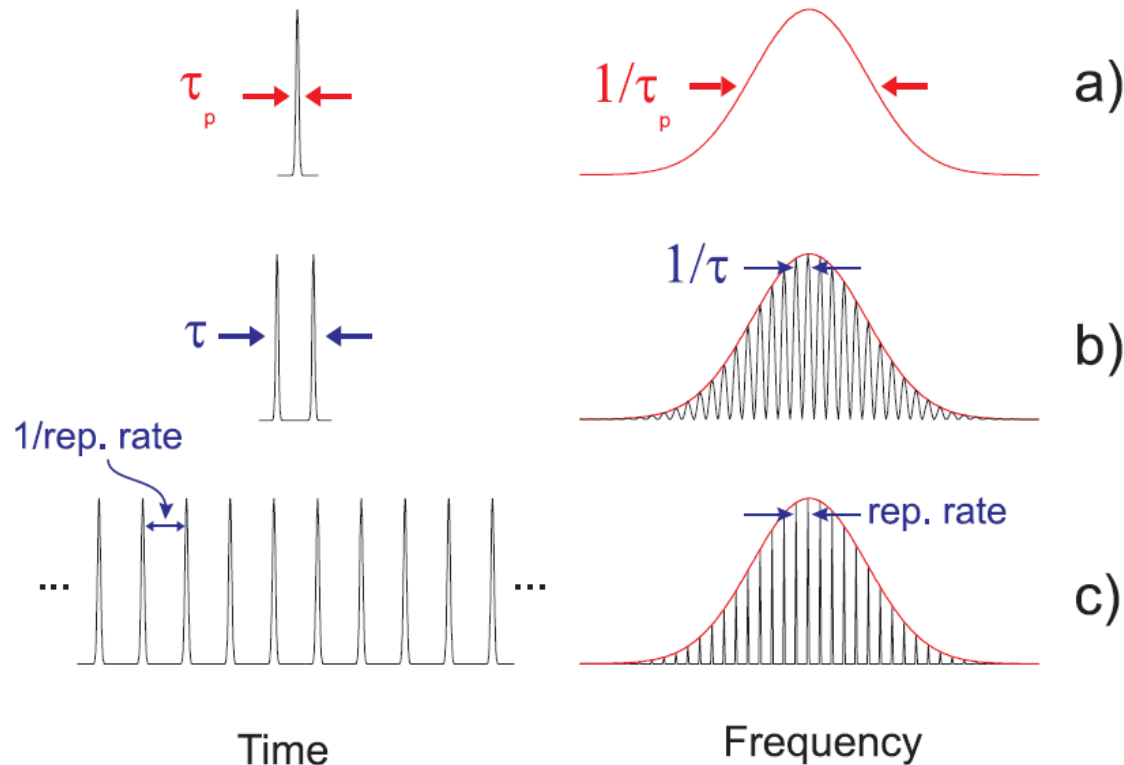


FIGURE 6.1

A single laser pulse has a frequency bandwidth which scales with the inverse of its duration τ_p (a). If one uses a pair of phase-locked pulses, delayed by a time τ , the resulting spectrum maintains the broad envelope of width $1/\tau_p$ but with a sinusoidal modulation of spectral period $1/\tau$ (b). This width sets the new instrumental resolution. If one uses an infinite sequence of pulses, locked in phase and equally delayed in time of a time interval τ , the spectrum breaks up in a *comb* of very narrow lines (the *teeth*) equally spaced by a frequency interval $1/\tau$ (c).

Principles of Optical Frequency Combs

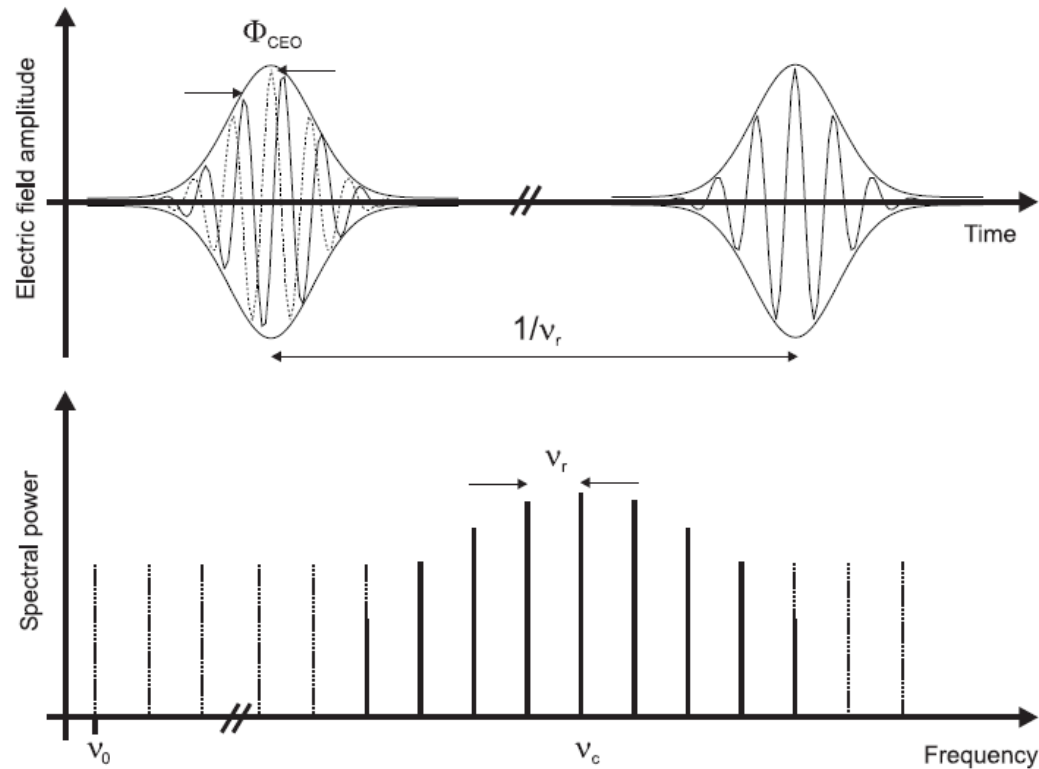
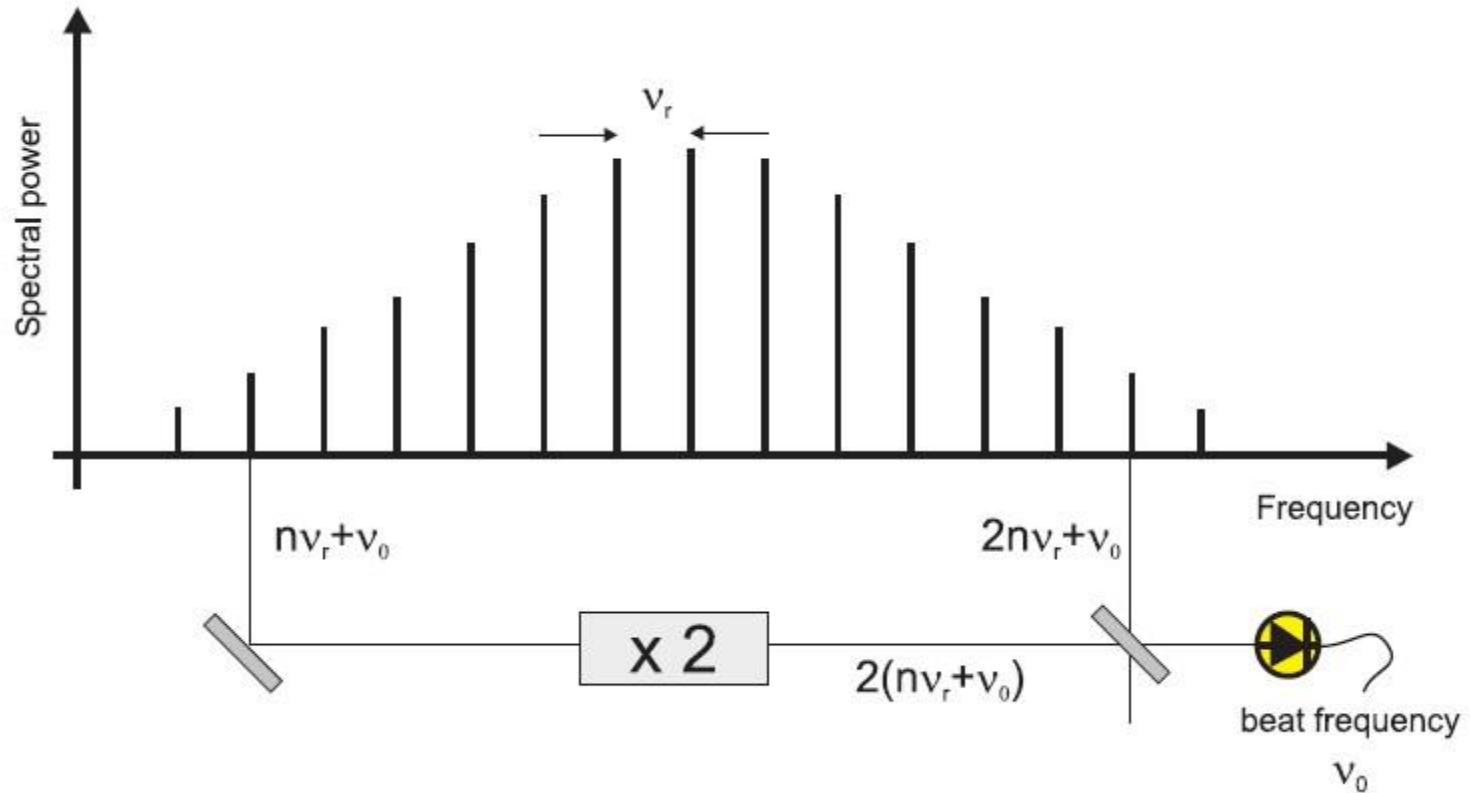


FIGURE 6.6

Scheme of a series of ultrashort laser pulses emitted from a mode-locked laser. Successive pulses are exact replicas of each other apart from a phase factor ϕ_{CEO} . This translates in the frequency domain to a comb-like structure of narrow spectral modes separated by the laser repetition frequency $\nu_r = 1/\tau_r$ and with an offset from zero frequency given by ν_0 .



Principles of Optical Frequency Combs



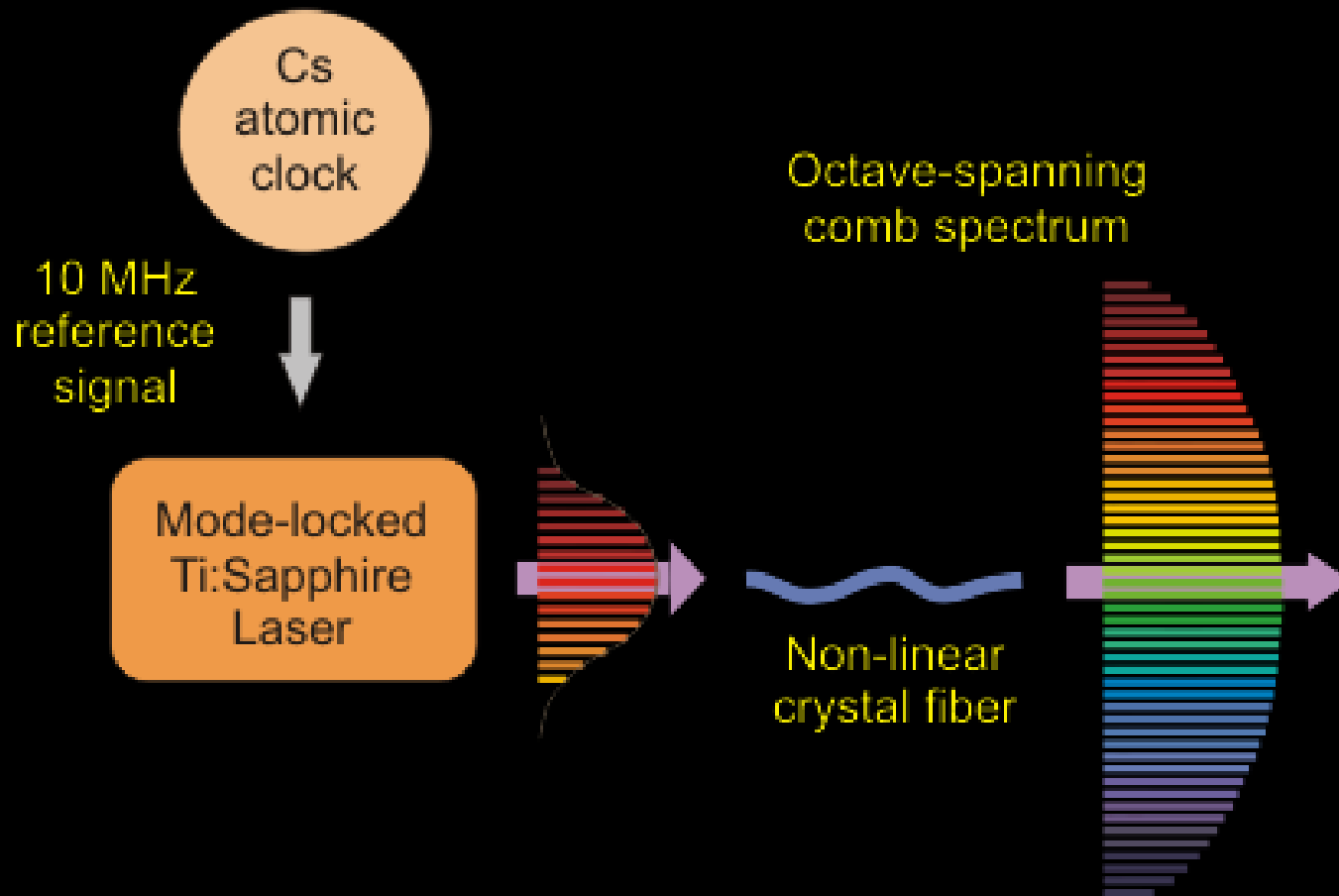


Born for Precise Frequency Measurements

for a review see, e.g.:

Optical comb generators for laser frequency measurement

Maddaloni et al. Meas. Sci. Technol. 20 (2009) 052001

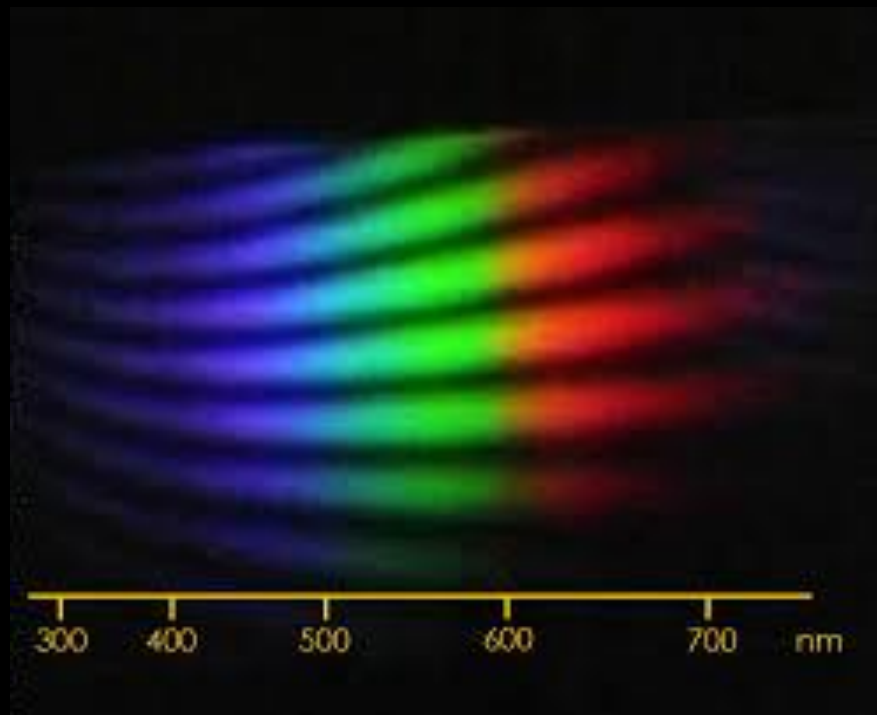




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The first evidence of comb behaviour

Experiment performed at LENS,
Firenze, Italy in 1998



July 15, 2000 / Vol. 25, No. 14 / OPTICS LETTERS 1049

Phase-locked white-light continuum pulses: toward a universal optical frequency-comb synthesizer

Marco Bellini

*Istituto Nazionale di Ottica Applicata, Largo E. Fermi 6, 50125 Florence, Italy, and
European Laboratory for Non Linear Spectroscopy and Istituto Nazionale di Fisica della Materia, Largo E. Fermi 2,
50125 Florence, Italy*

Theodor W. Hänsch

*European Laboratory for Non Linear Spectroscopy, Istituto Nazionale di Fisica della Materia,
and Department of Physics, University of Florence, Largo E. Fermi 2, 50125 Florence, Italy, and
Max-Planck-Institut für Quantenoptik, P.O. Box 1513, D-85740 Garching, Germany*

Received January 31, 2000

We demonstrate that two white-light continuum pulses that are independently generated by phase-locked ultrashort laser pulses are locked in phase and show surprisingly clear and stable Young interference fringes. The experiment shows that the two generated continua emit essentially in phase and that random phase jitter can remain negligible. This result is not only of interest for studies of nonlinear field-matter interactions but also suggests that such white-light continuum pulses can be used to realize a broad frequency comb for absolute frequency measurements from the IR to the UV. © 2000 Optical Society of America

OCIS codes: 190.0190, 190.2620, 190.5940, 300.6320, 320.7110.

Direct Link between Microwave and Optical Frequencies with a 300 THz Femtosecond Laser Comb

Scott A. Diddams,* David J. Jones, Jun Ye, Steven T. Cundiff, and John L. Hall†

JILA, University of Colorado, and National Institute of Standards and Technology, Boulder, Colorado 80309

Jinendra K. Ranka and Robert S. Windeler

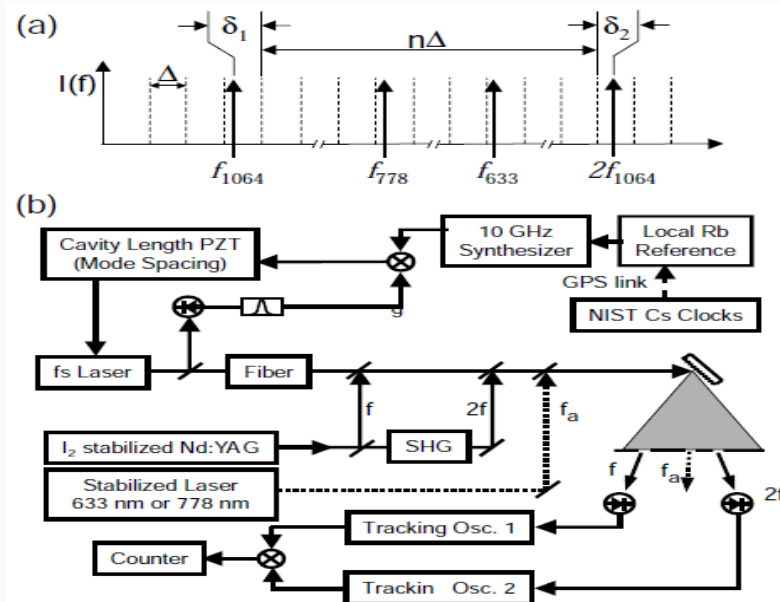
Bell Laboratories, Lucent Technologies, 700 Mountain Avenue, Murray Hill, New Jersey 070974

Ronald Holzwarth, Thomas Udem, and T.W. Hänsch

Max-Planck-Institut für Quantenoptik, 85748 Garching, Germany

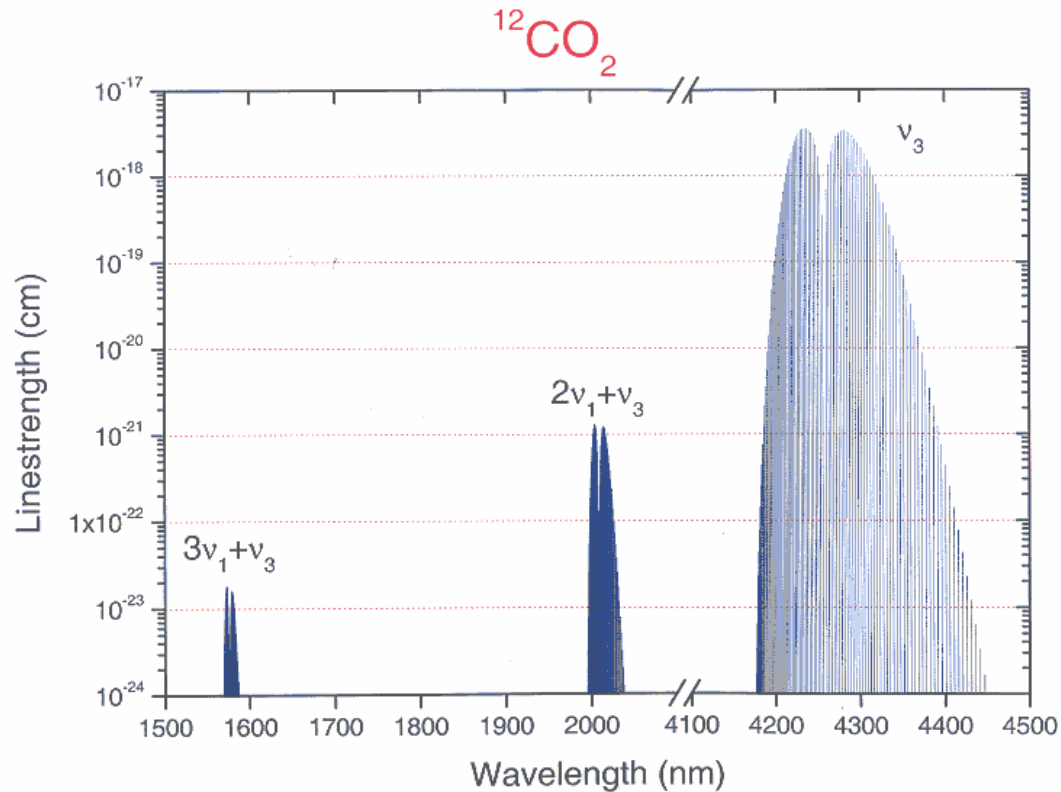
(Received 17 February 2000)

We demonstrate a great simplification in the long-standing problem of measuring optical frequencies in terms of the cesium primary standard. An air-silica microstructure optical fiber broadens the frequency comb of a femtosecond laser to span the optical octave from 1064 to 532 nm, enabling us to measure the 282 THz frequency of an iodine-stabilized Nd:YAG laser directly in terms of the microwave frequency that controls the comb spacing. Additional measurements of established optical frequencies at 633 and 778 nm using the same femtosecond comb confirm the accepted uncertainties for these standards.





Intensity Scaling of Molecular Absorption



.....and natural combs

1256 OPTICS LETTERS / Vol. 27, No. 14 / July 15, 2002

**Low-power Lamb-dip spectroscopy of
very weak CO₂
transitions near 4.25 μm**



IEEE JOURNAL OF QUANTUM ELECTRONICS,
VOL. 29, NO. 10, OCTOBER 2693 (1993)

Wide-Span Optical Frequency Comb Generator for Accurate Optical Frequency Difference Measurement

Motonobu Kourogi, Ken'ichi Nakagawa and Motoichi Ohtsu,

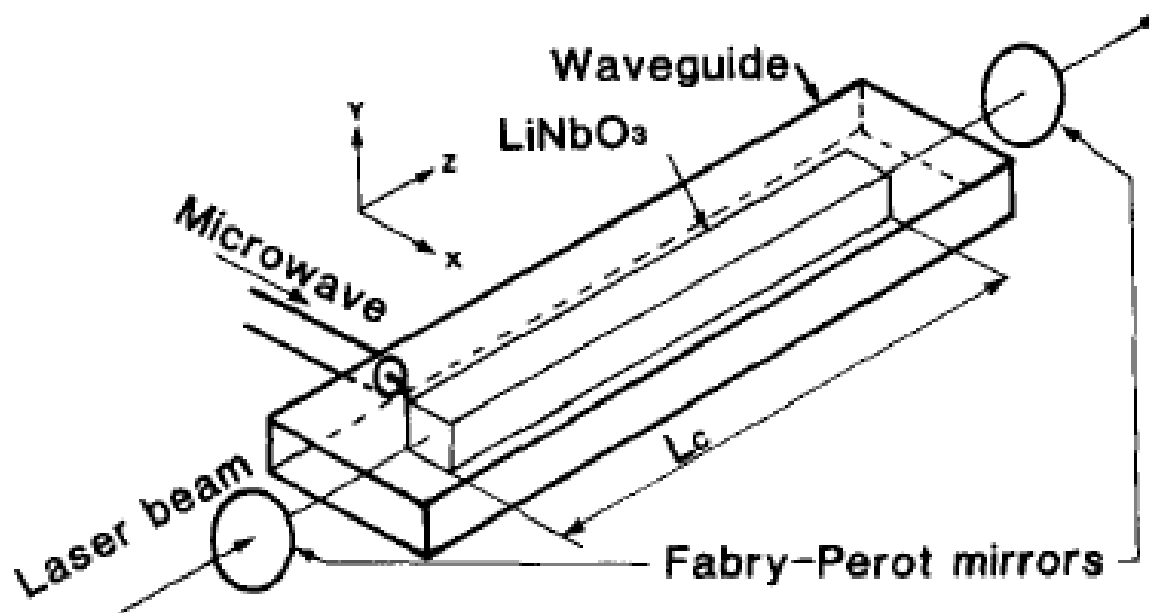


Fig. 2. The construction of the present optical frequency comb generator.

Quantum cascade laser frequency combs

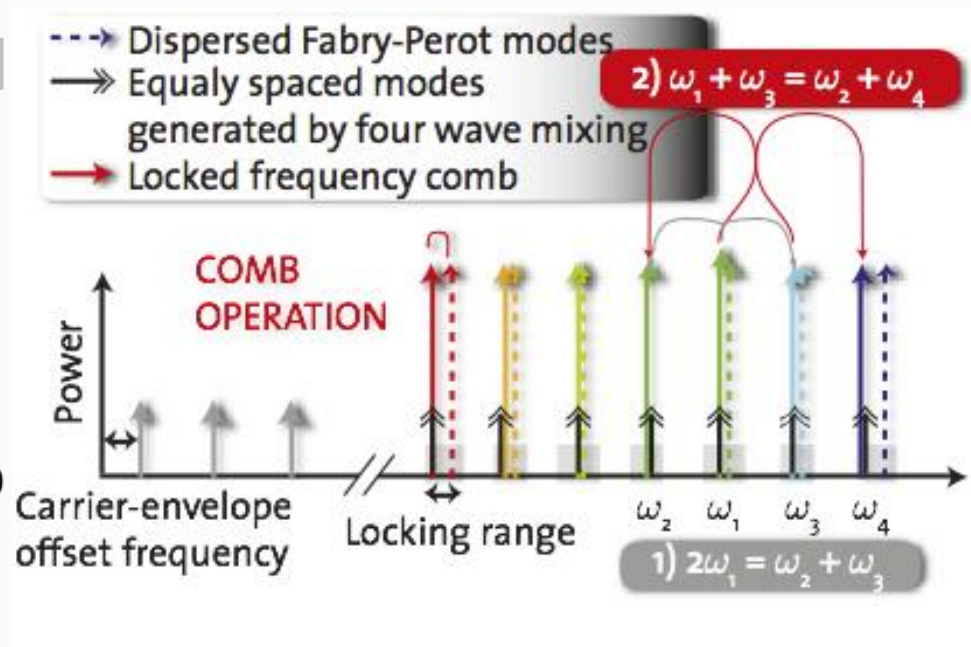
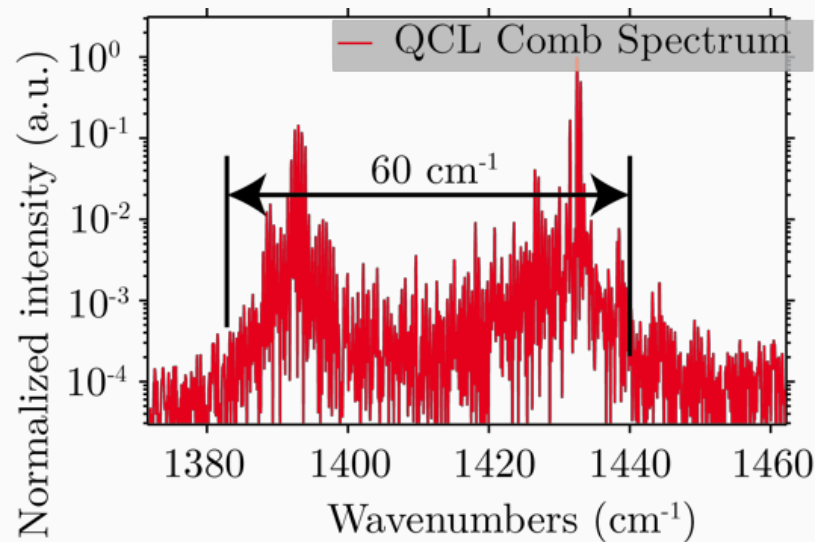


LETTER

doi:10.1038/nature11620

Mid-infrared frequency comb based on a quantum cascade laser

Andreas Hugi¹, Gustavo Villares¹, Stéphane Blaser², H. C. Liu³ & Jérôme Faist¹



(3) E^3

NL



Quantum cascade laser frequency combs

ARTICLES

PUBLISHED ONLINE: 11 MAY 2014 | DOI: 10.1038/NPHOTON.2014.85

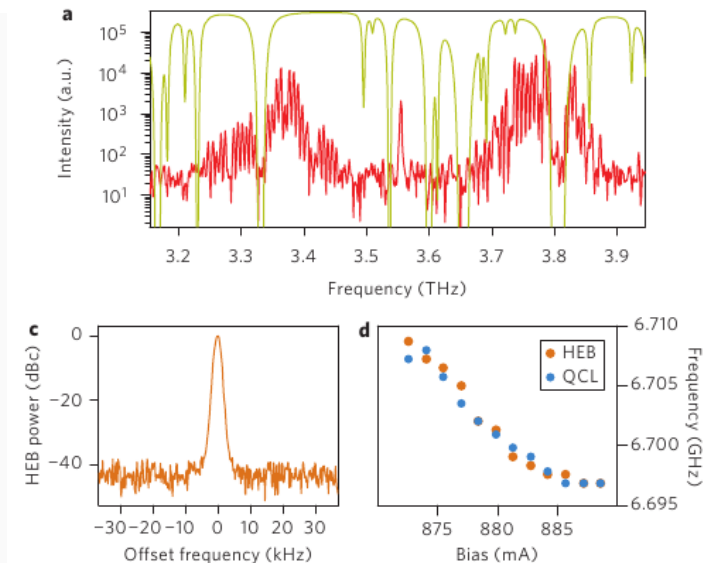
nature
photonics

Terahertz laser frequency combs

David Burghoff^{1*}, Tsung-Yu Kao¹, Ningren Han¹, Chun Wang Ivan Chan¹, Xiaowei Cai¹, Yang Yang¹, Darren J. Hayton², Jian-Rong Gao^{2,3}, John L. Reno⁴ and Qing Hu¹

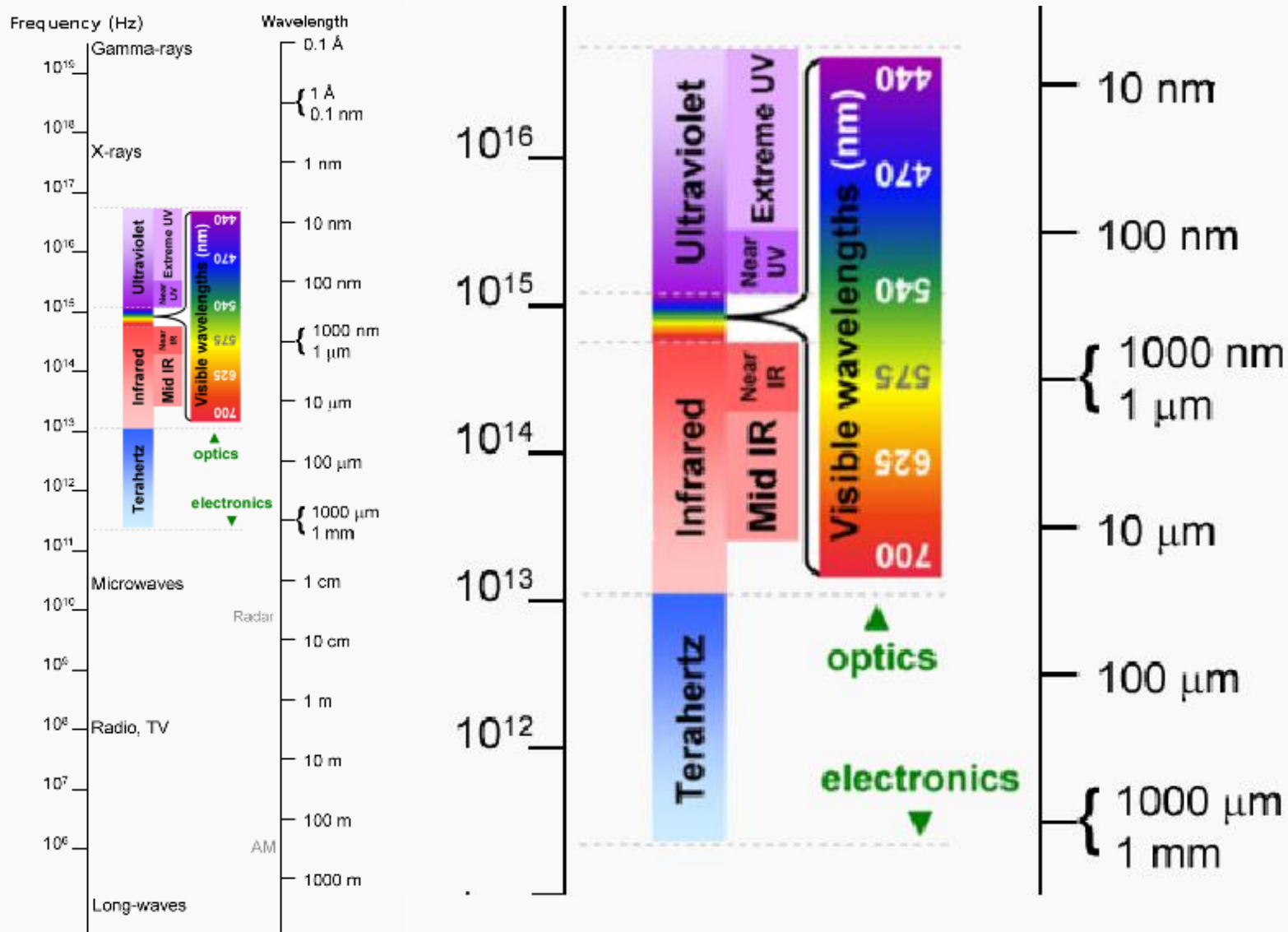
Figure Continuous-wave spectrum and beat notes. **a**, Spectrum of a THz QCL comb biased to 0.9 A at a temperature of 50 K. The device is 20 μm wide, 5 mm long and emits 4.7 mW at 45 K. Atmospheric absorption is shown in yellow.

c, HEB-detected beat-note offset relative to 6.80 GHz, stabilized. The linewidth of 1.53 kHz is limited by the instrument resolution. **d**, Bias dependence of the beat-note frequency measured by the QCL and by the HEB. A repetition rate tuning of 12 MHz is possible, giving a total frequency shift of 6.80 GHz at 3.8 THz, approximately the mode spacing.





Electromagnetic Spectrum

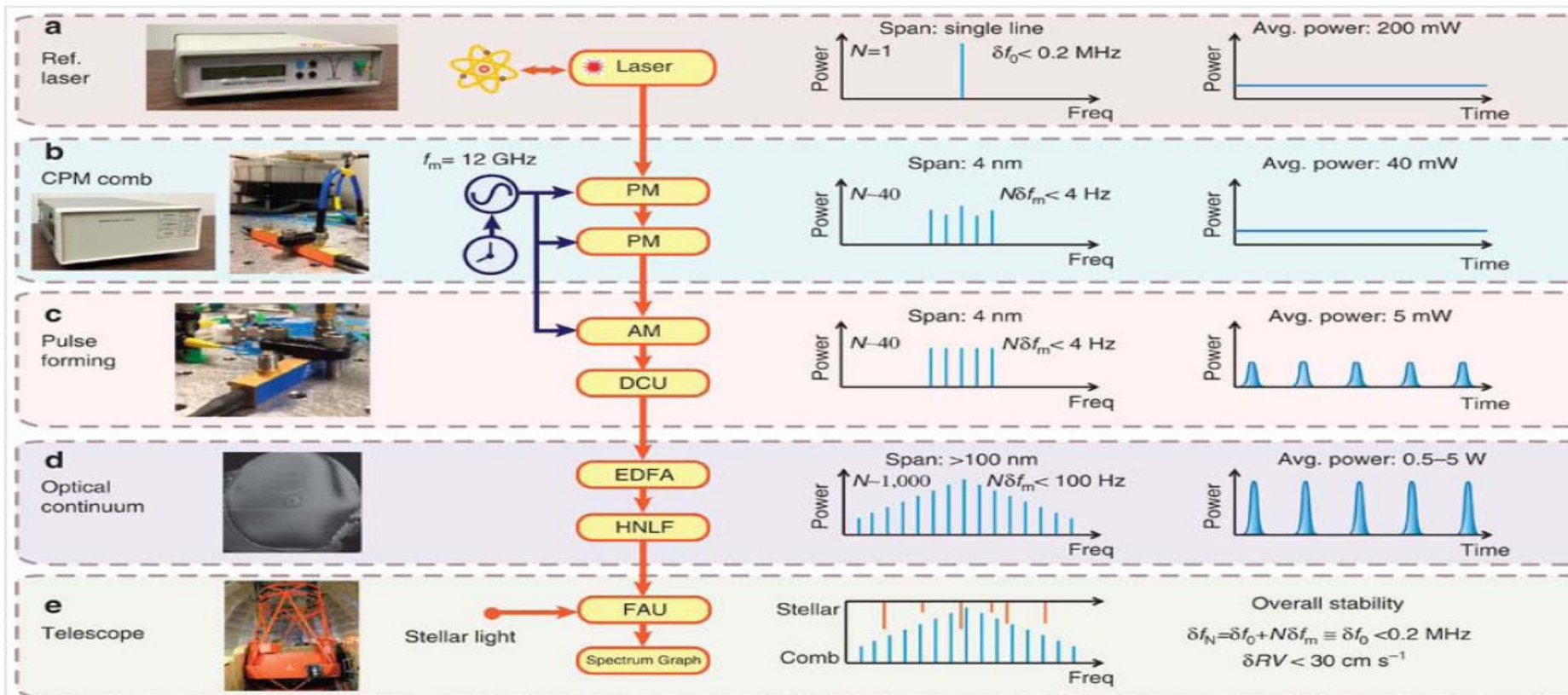


From

Demonstration of a near-IR line-referenced electro-optical laser frequency comb for precision radial velocity measurements in astronomy

X. Yi, K. Vahala, J. Li, S. Diddams, G. Ycas, P. Plavchan, S. Leifer, J. Sandhu, G. Vasisht, P. Chen, P. Gao, J. Gagne, E. Furlan, M. Bottom, E. C. Martin, M. P. Fitzgerald, G. Doppmann & C. Beichman

Nature Communications 7, Article number: 10436 doi:10.1038/ncomms10436

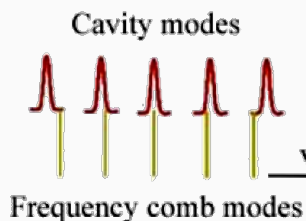
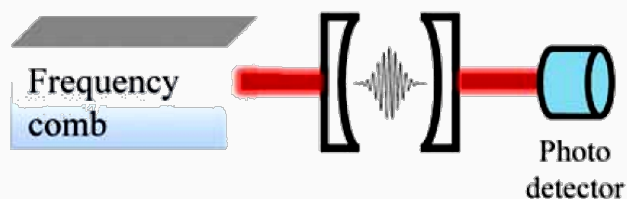
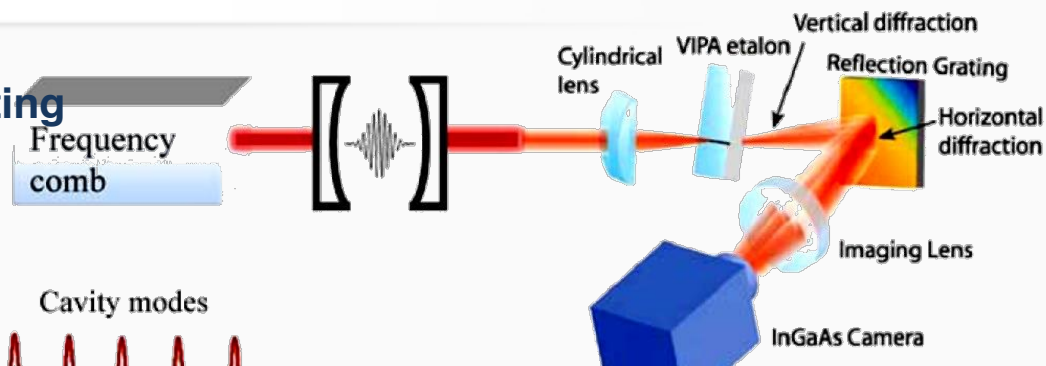


Vertically, the first column contains images of key instruments. (a–e) The images are reference laser, Rb clock (left) and phase modulator (right), amplitude modulator, highly nonlinear fibre and telescope. A simplified schematic set-up is in the second column. Third and fourth columns present the comb state in the frequency and temporal domains. The frequency of N -th comb tooth is expressed as $f_N = f_0 + N \times f_m$, where f_0 and f_m are the reference laser frequency and modulation frequency, respectively. N is the number of comb lines relative to the reference laser (taken as comb line $N=0$), RV is radial velocity and δf_N , δf_0 and δf_m are the variance of f_N , f_0 and f_m . (a) The reference laser is locked to a molecular transition, acquiring stability of 0.2 MHz, corresponding to 30 cm s^{-1} RV . (b) Cascaded phase modulation (CPM) comb: the phase of the reference laser is modulated by two phase modulators (PM), creating several tens of sidebands with spacing equal to the modulation frequency. The RF frequency generator is referenced to a Rb clock, providing stability at the sub-Hz level ($\delta f_m < 0.03 \text{ Hz}$ at 100 s). (c) Pulse forming is then performed by an amplitude modulator (AM) and dispersion compensation unit (DCU), which could be a long single mode fibre (SMF) or chirped fibre Bragg grating (FBG). (d) After amplification by an erbium-doped fibre amplifier (EDFA), the pulse undergoes optical continuum broadening in a highly nonlinear fibre (HNLF), extending its bandwidth $>100 \text{ nm}$. (e) Finally the comb light is combined with stellar light using a fibre acquisition unit (FAU) and is sent into the telescope spectrograph. The overall comb stability is primarily determined by the pump laser.



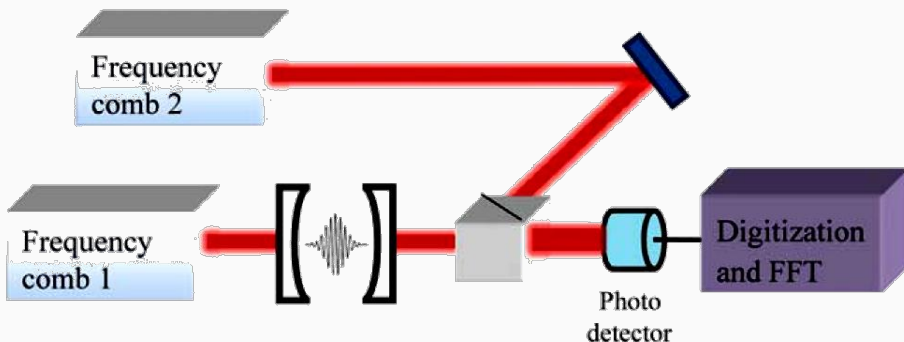
Direct Comb Spectroscopy

Comb Dispersion by VIPA and Grating



Comb Vernier Spectroscopy

Fourier-Transform Analysis



Dual Comb Spectroscopy



Advanced Optical Fiber-based spectrometers

- Broad Absorptions in Liquids
- Strain (DL/L)
- Temperature variations

Combs for spectroscopy in liquids

201116-2 Avino *et al.*

Appl. Phys. Lett. 102, 201116 (2013)

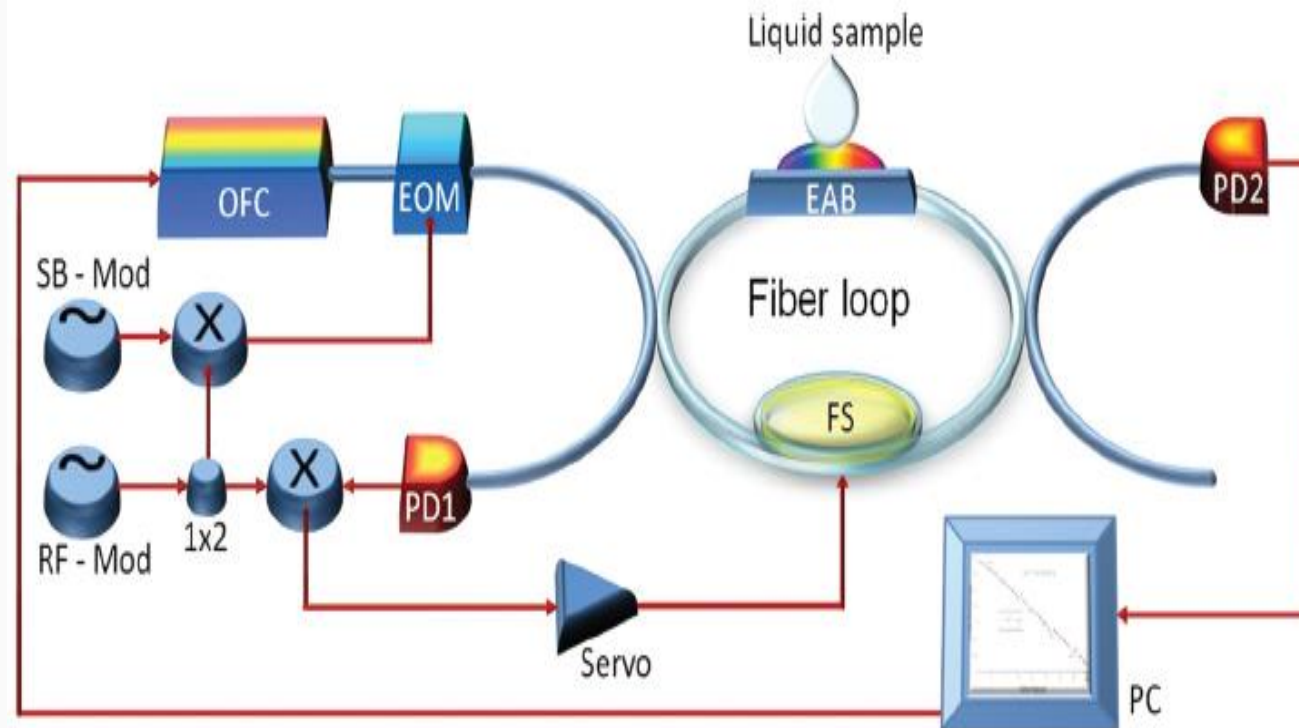
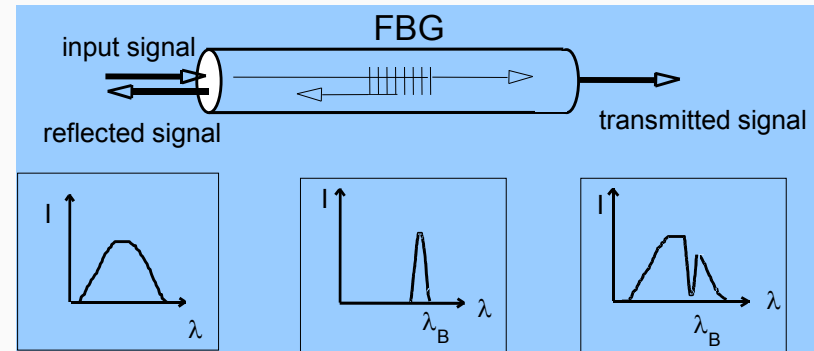


FIG. 1. Experimental layout. OFC: optical frequency comb; EOM: electro-optic phase modulator; SB – Mod: phase modulation to switch the beam off; RF – Mod: phase modulation for PDH locking. PD: photodiode; PC: personal computer; Servo: locking electronics. EAB: evanescent-wave access block; FS: fiber stretcher.

Operating Principle of FBGs

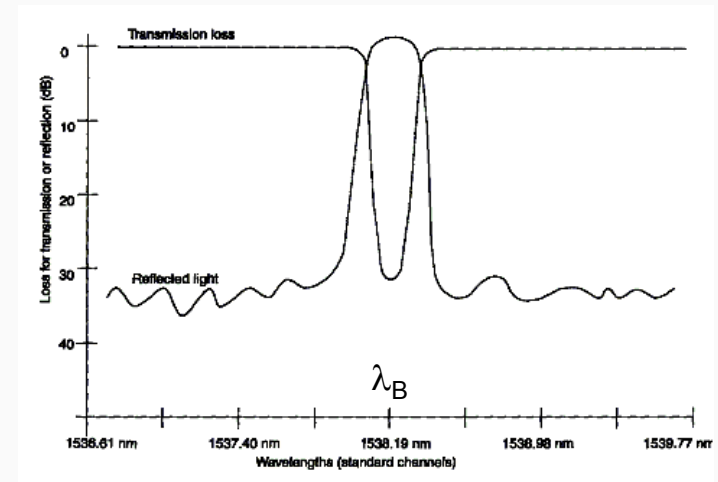
Periodic modulation
of the refractive index
in the core of an optical fiber



Bragg
max reflectivity

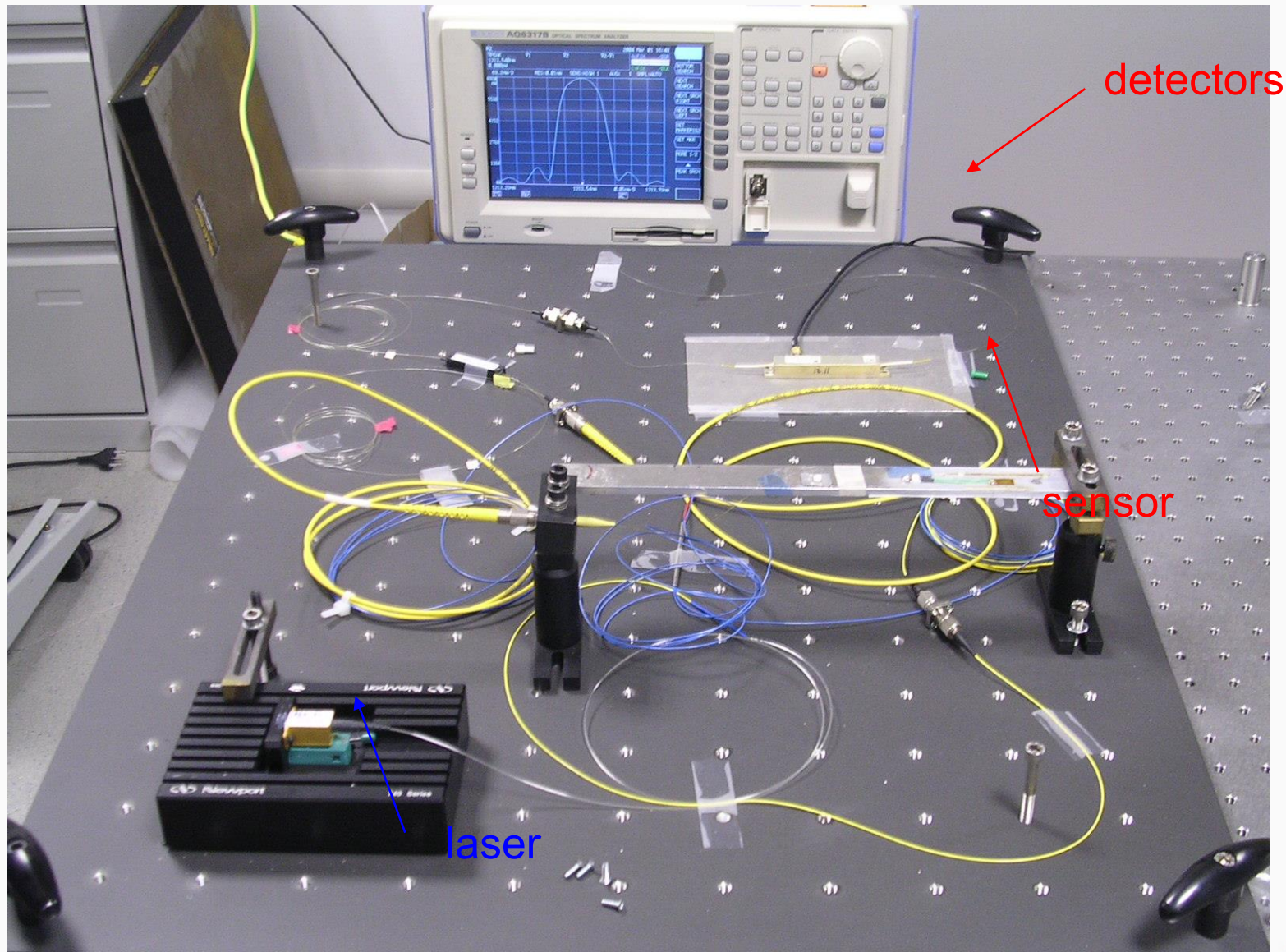
$$\lambda_B = 2n_{eff}\Lambda$$

**Any change in the grating pitch
and/or refractive index results in
($\Delta L/L=10^{-6} \Rightarrow 1 \mu\epsilon$) a shift of λ_B
 $\sim 10^{-3} \text{ nm}/\mu\epsilon$**





Set-up for Deformation Sensor





FBG-based sensors

Capabilities:

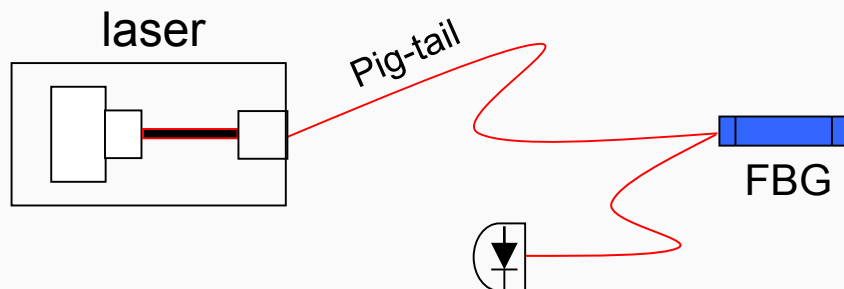
- static and dynamic point strain measurement;
- field and remote operation;
- immunity to EMI;
- embedding in materials;
- monitoring networks;
- low weight and cost;
- application as field sensors in several areas.

Conventional interrogation techniques:

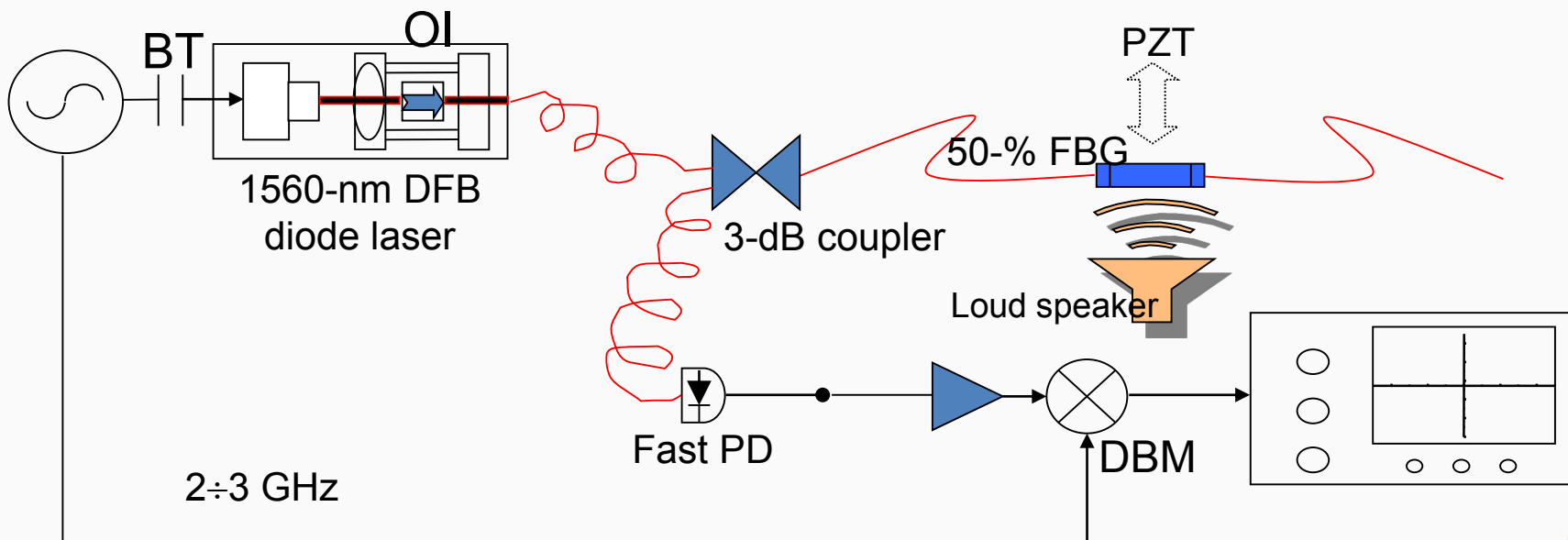
- broad-band sources combined to passive or active filters;
- limited sensitivity;
- frequency response.



Interrogation of FBGs using a laser source

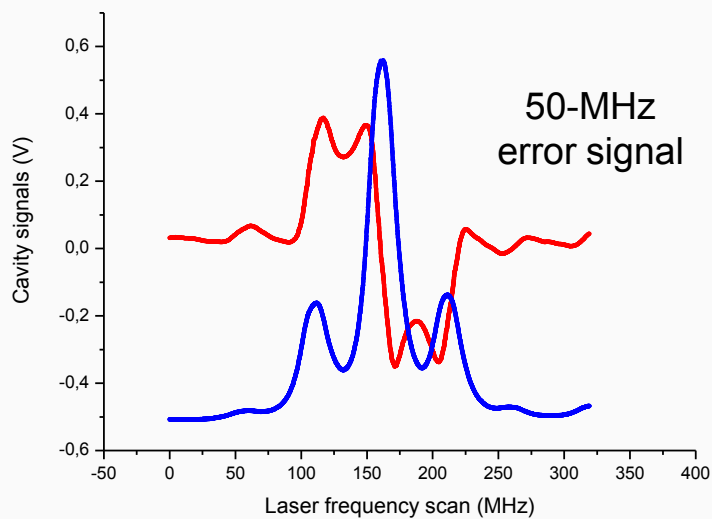
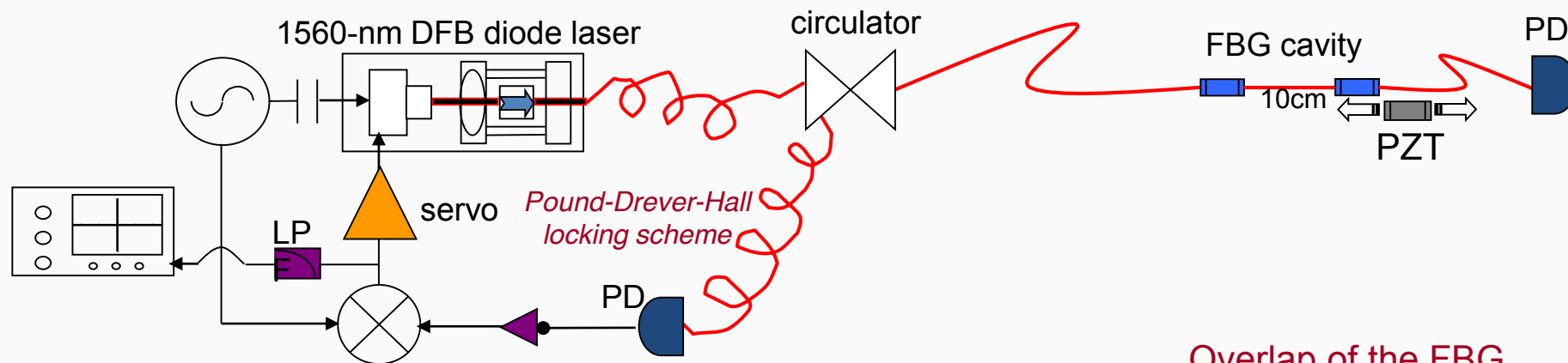


FM-based Bragg-grating set-up

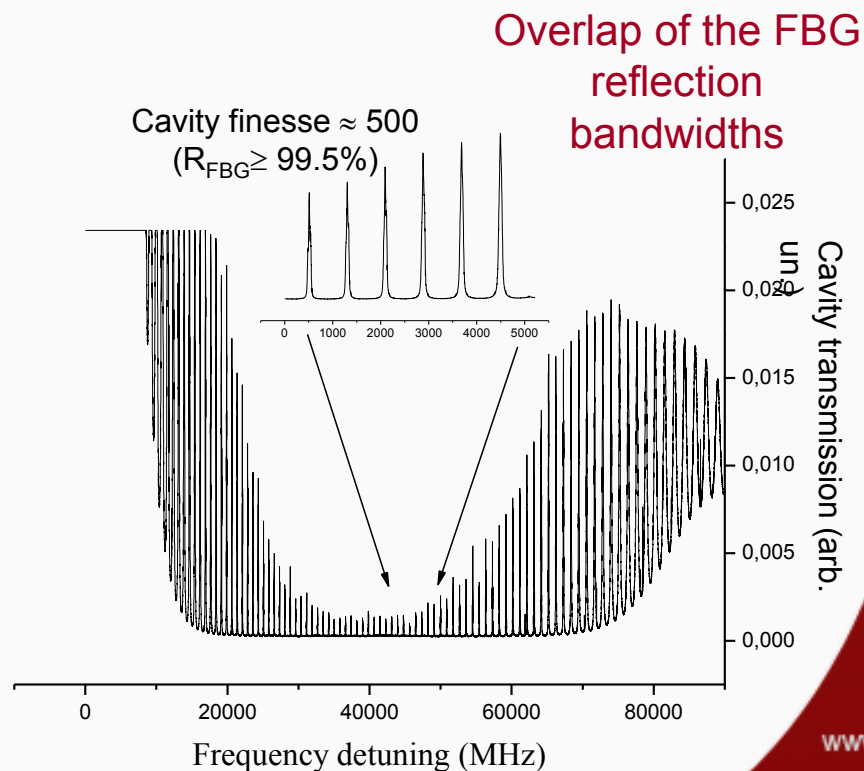




FBG-based Fabry-Pérot strain sensor



Strain monitor

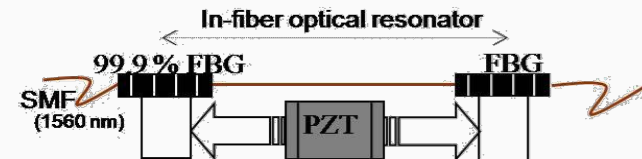


Fiber Optical Sensors



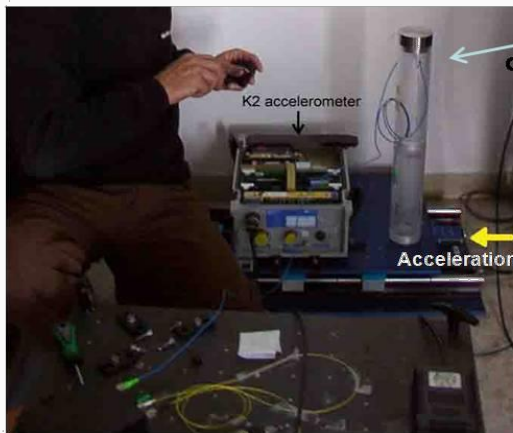
High-sensitivity strain sensing apparatus

•High-sensitivity laser-based strain and temperature sensing using “fiber Bragg grating” (FBG) resonators: interrogation of fiber optic sensors by highly-stabilized lasers



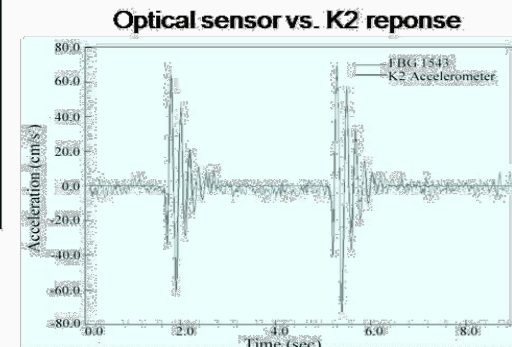
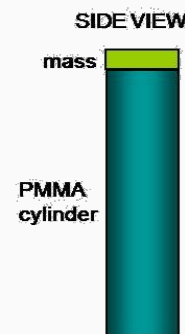
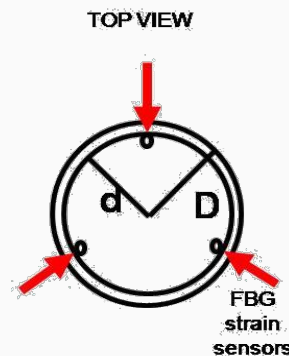
Optical-frequency combs for laser frequency stabilization

Strain meas towards the 10^{-14} level



FBG flexural beam accelerometer: comparison with a commercial K2 seismometer on a shaking table

•Seismometers based on FBG technology: spectroscopic interrogation by frequency-modulated telecom diode lasers





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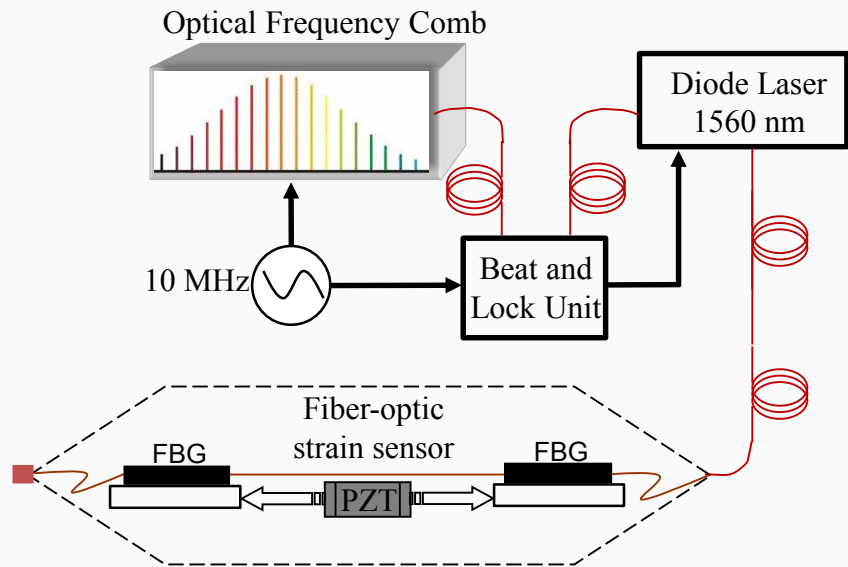
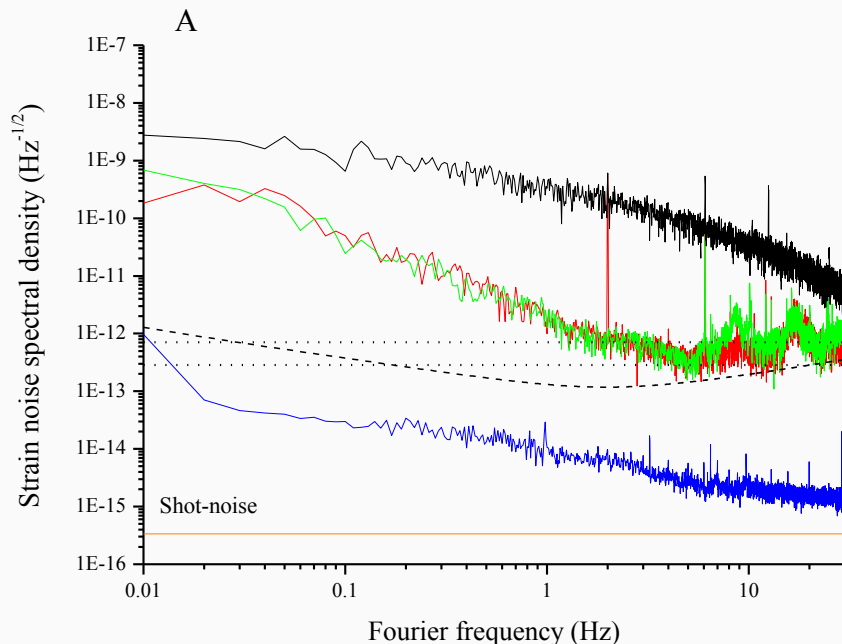
Comb-assisted strain sensing

Probing the Ultimate Limit of Fiber-Optic Strain Sensing

G. Gagliardi,^{1*} M. Salza,¹ S. Avino,¹ P. Ferraro,¹ P. De Natale²



SCIENCE VOL 330 19 NOVEMBER 2010



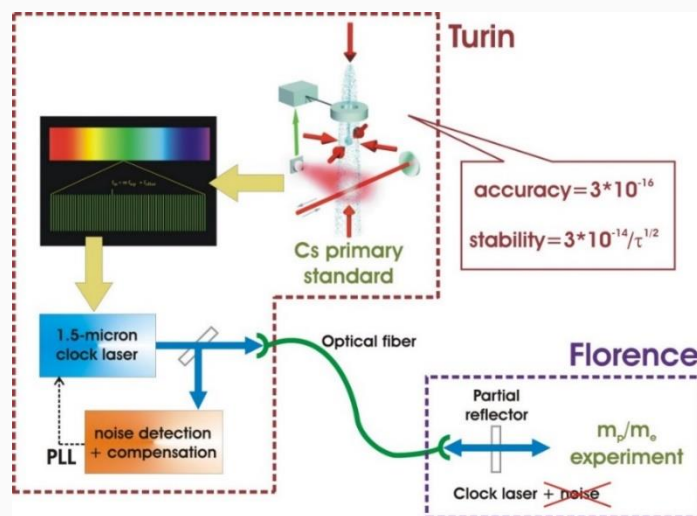
Thermodynamic noise limited strain resolution $\sim 900\text{-}350 \text{ f}\epsilon/\sqrt{\text{Hz}}$ in the infrasonic range



NATIONAL OPTICAL FIBER LINK

GPS limits the comb accuracy to 5×10^{-12}

- improved primary standard



Istituto Nazionale di
Ricerca Metrologica

A 1.5-micron laser, continuously referenced to the primary Cs fountain via an OFC, is disseminated through a long optical fiber. Part of the laser radiation is reflected back and used to cancel the phase noise accumulated during the fiber propagation

Accuracy = $3 \cdot 10^{-16}$

Stability = $3 \cdot 10^{-14} / \tau^{1/2}$

📖 D. Calonico et al., Appl. Phys. B **117**, 979 (2014)

Pushing the ultimate resolution in the spectroscopic frequency measurement down to 10^{-18} by stabilizing the frequency comb, to which the probe laser is referenced, against an optical atomic standard



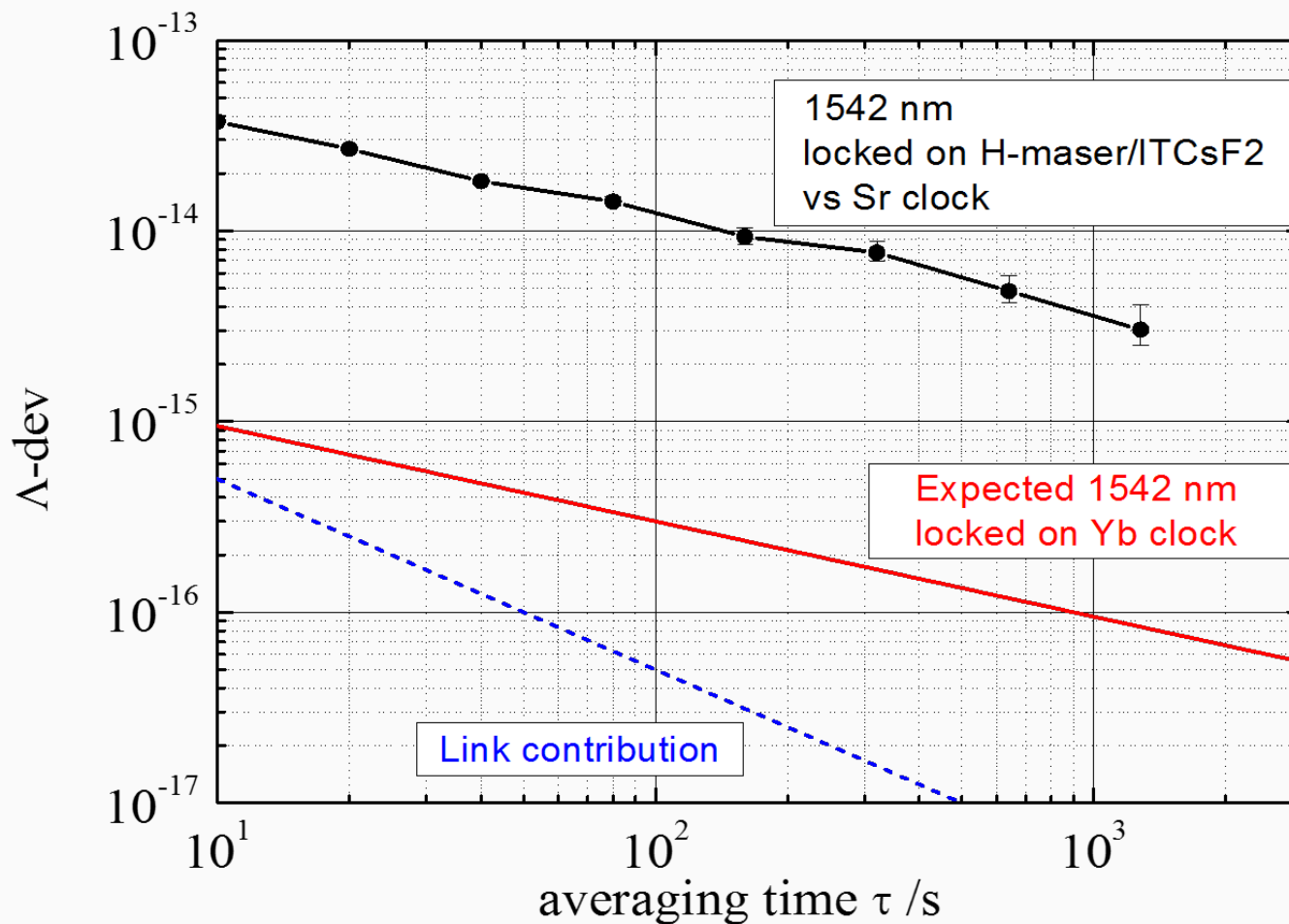
- Yb lattice clock under construction at INRIM

$$\frac{1}{\frac{\nu_{vib}(M)}{\nu_{el}(Yb)}} \frac{\partial \left[\frac{\nu_{vib}(M)}{\nu_{el}(Yb)} \right]}{\partial t} = -0.5 \frac{\dot{\beta}}{\beta} - N_{Yb} \frac{\dot{\alpha}}{\alpha}$$



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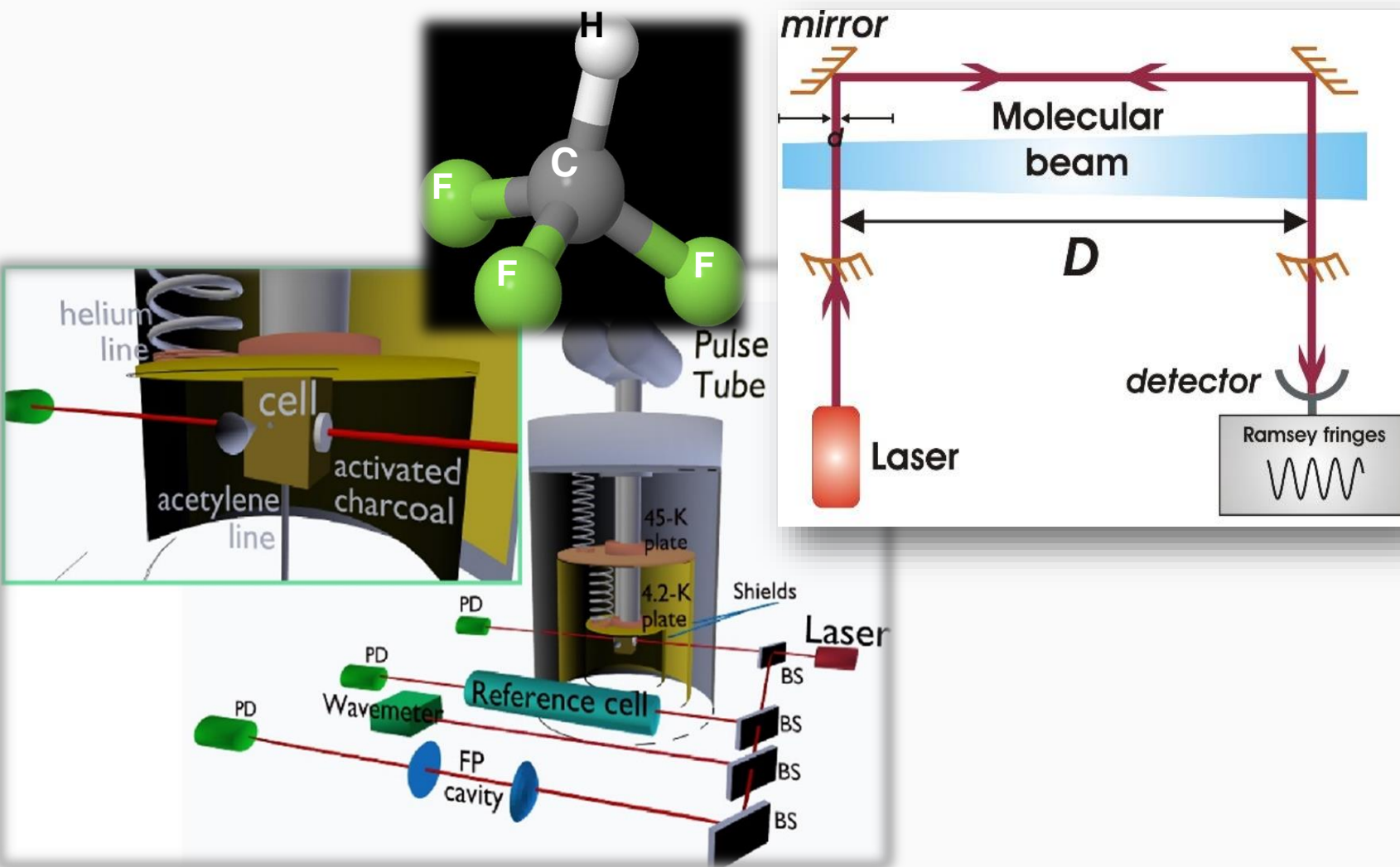
Stability transferred by the link (at 642 km distance, Turin-Florence, Italy)





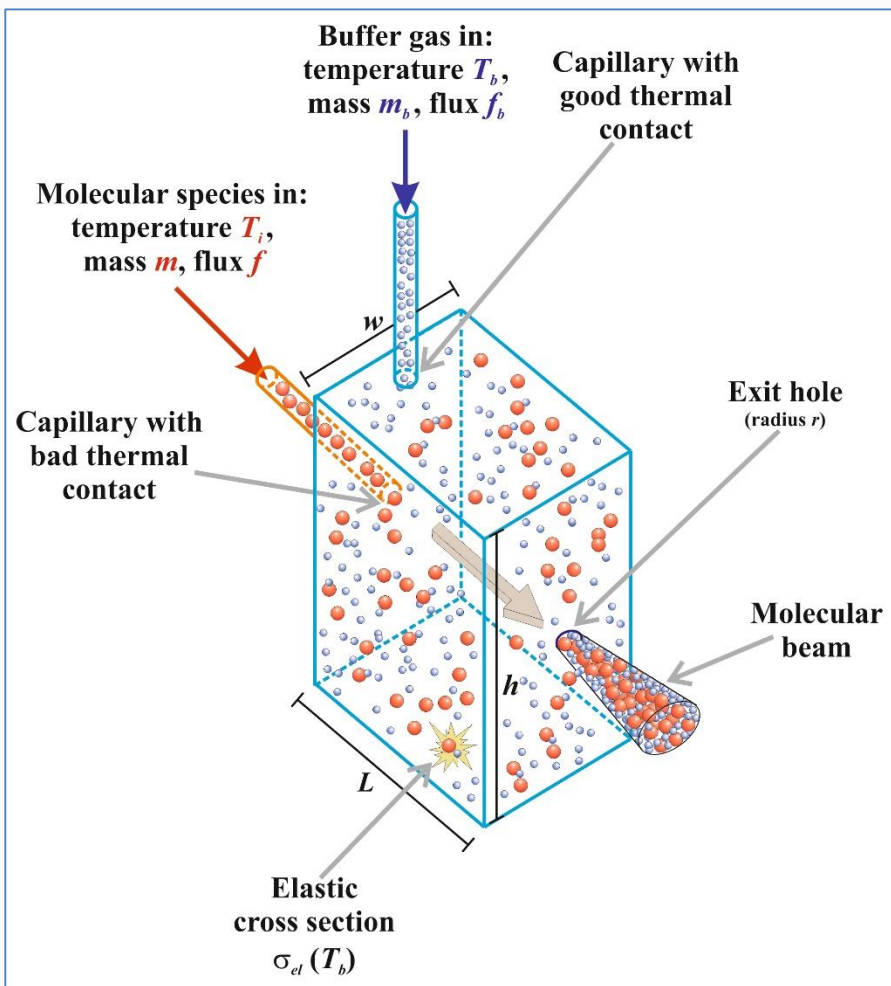
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Cold Molecules





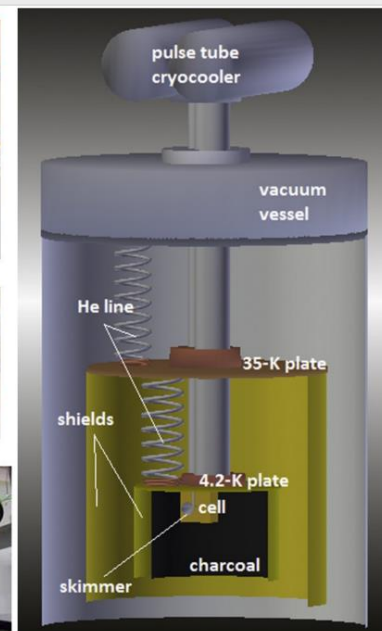
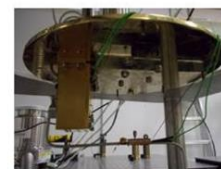
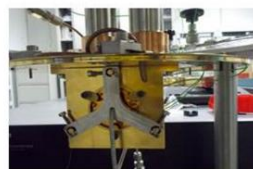
BUFFER GAS COOLING



Both translational and rotational (internal) degrees of freedom of the desired molecular species are cooled via collisions with a thermal bath of helium in a cryogenic cell

📖 S.E. Maxwell et al., Phys. Rev. Lett. **95**, 173201 (2005)

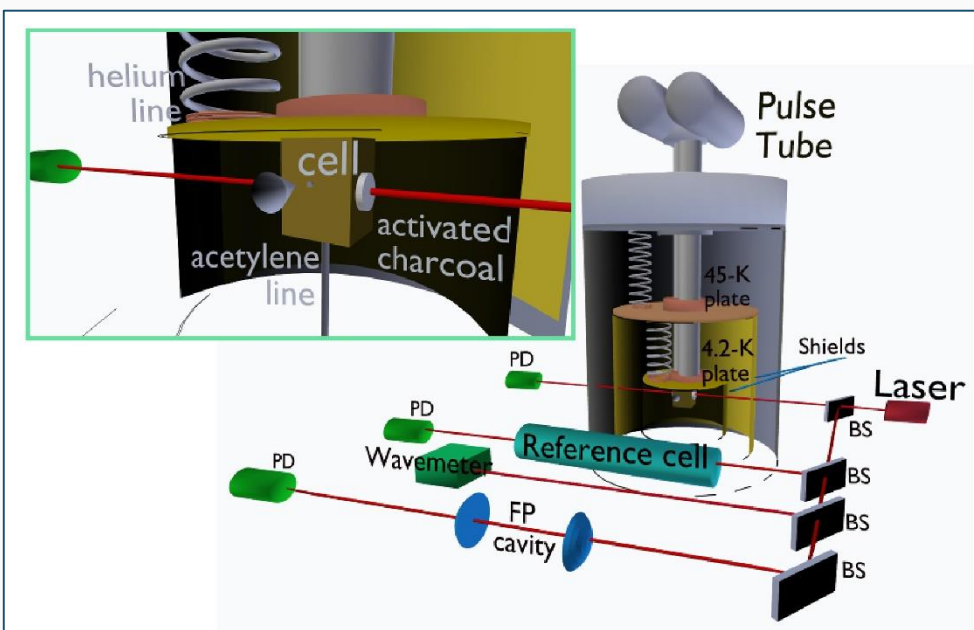
📖 N.R. Hutzler et al., Chem. Rev. **112**, 4803 (2012)





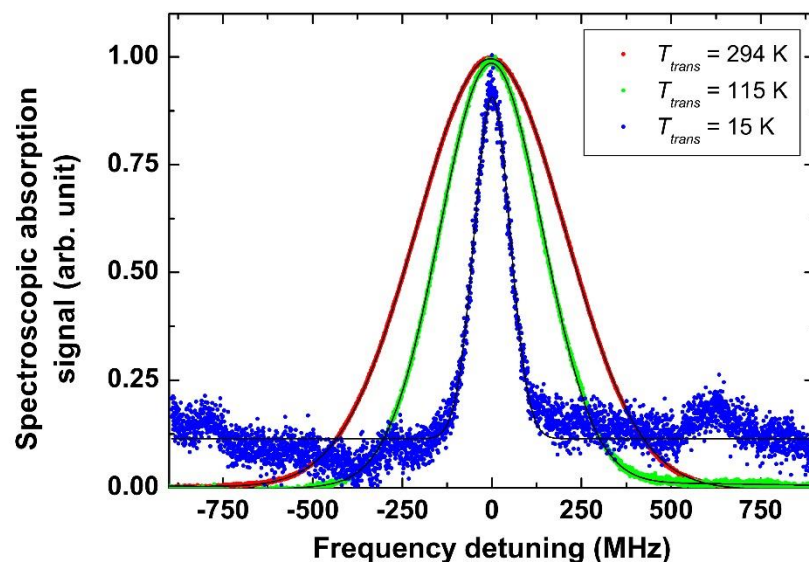
CHARACTERIZING THE COLLISIONAL COOLING PROCESS

Laser-absorption ro-vibrational spectroscopy of the $^{12}\text{C}_2\text{H}_2$ ($\nu_1 + \nu_3$) band is performed



Due to a non perfect thermal exchange between the copper pipe and the two PT plates, a temperature of about 15 K is measured on the He line just before the entrance into the BGC cell (at 4.2 K); to bridge this gap, an improved setup for better cooling of the He line is under construction

Translational temperatures are extracted by Doppler thermometry

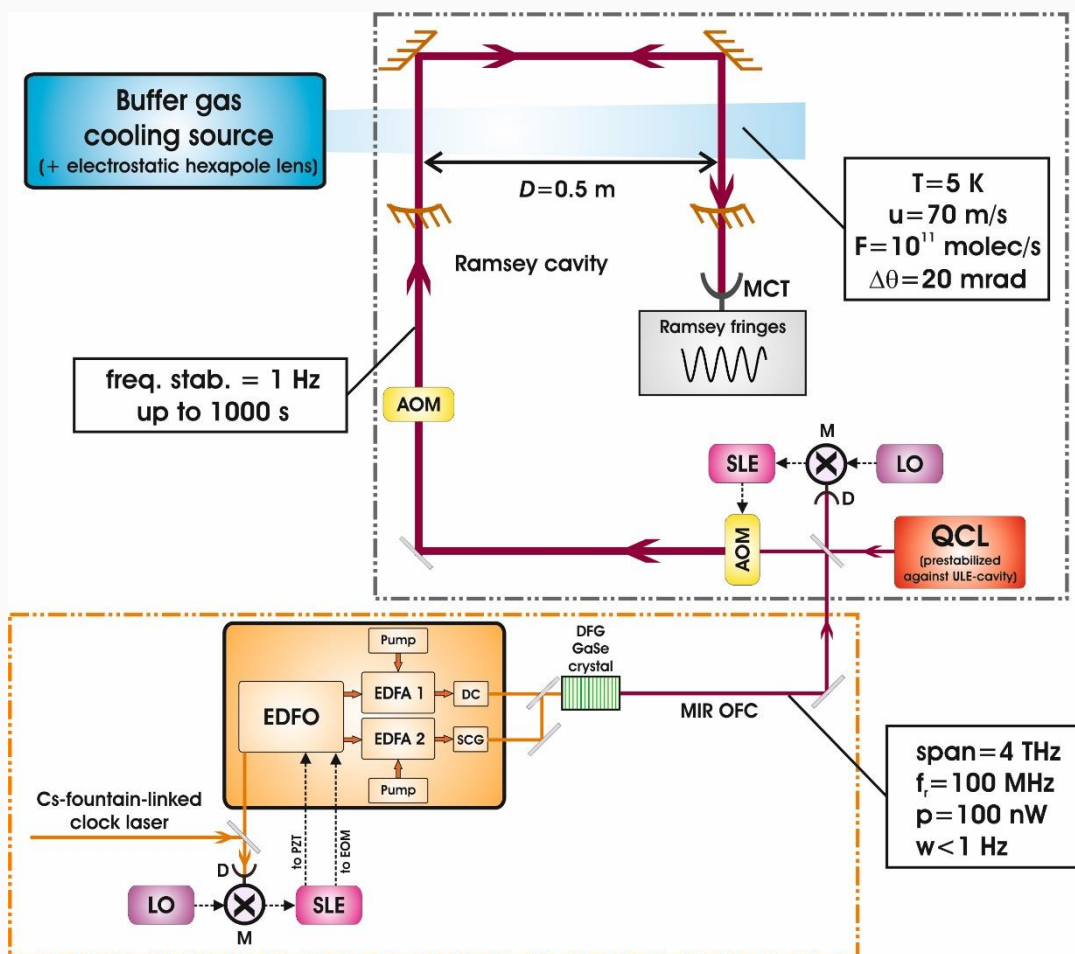


$$G(\nu) = G_0 \exp \left[-\frac{4 \ln 2 (\nu - \nu_0)^2}{\sigma_D^2} \right]$$

$$\sigma_D = \frac{\nu_0}{c} \sqrt{\frac{8 \ln 2 k_B T_{trans}}{m}}$$



MEASUREMENT OF TIME VARIATION OF CONSTANTS: A PROPOSAL



The cold fluoroform (CF_3H) beam is extracted from a buffer-gas-cooling source (and then collimated by means of an electrostatic hexapole lens). Employed in a 2-photon Ramsey-fringes interrogation, the probe source is a mid-infrared quantum cascade laser at 8.63 micron, phase-locked to a specially-developed optical frequency comb (OFC) that is ultimately referenced to the cesium primary standard via the National optical fiber link.

$$\mathcal{P} = \frac{u}{2D} = 35 \text{ Hz} \quad \text{SNR} = 25, N_{\text{avg}} = 100$$

$$\mathcal{R} \sim \frac{\mathcal{P}}{\text{SNR} \cdot \sqrt{N_{\text{avg}}}} = 140 \text{ mHz}$$

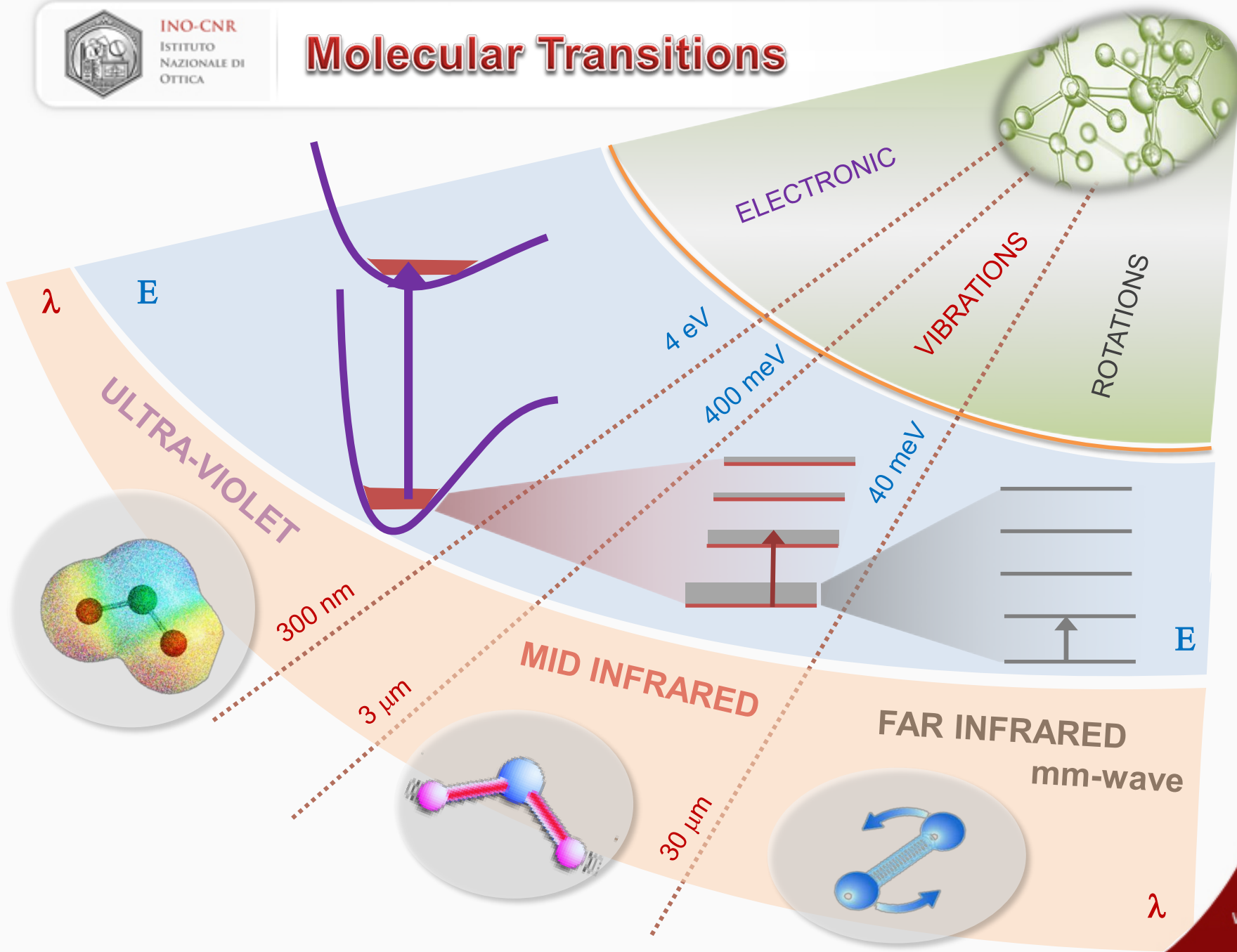


$$4 \cdot 10^{-15}$$



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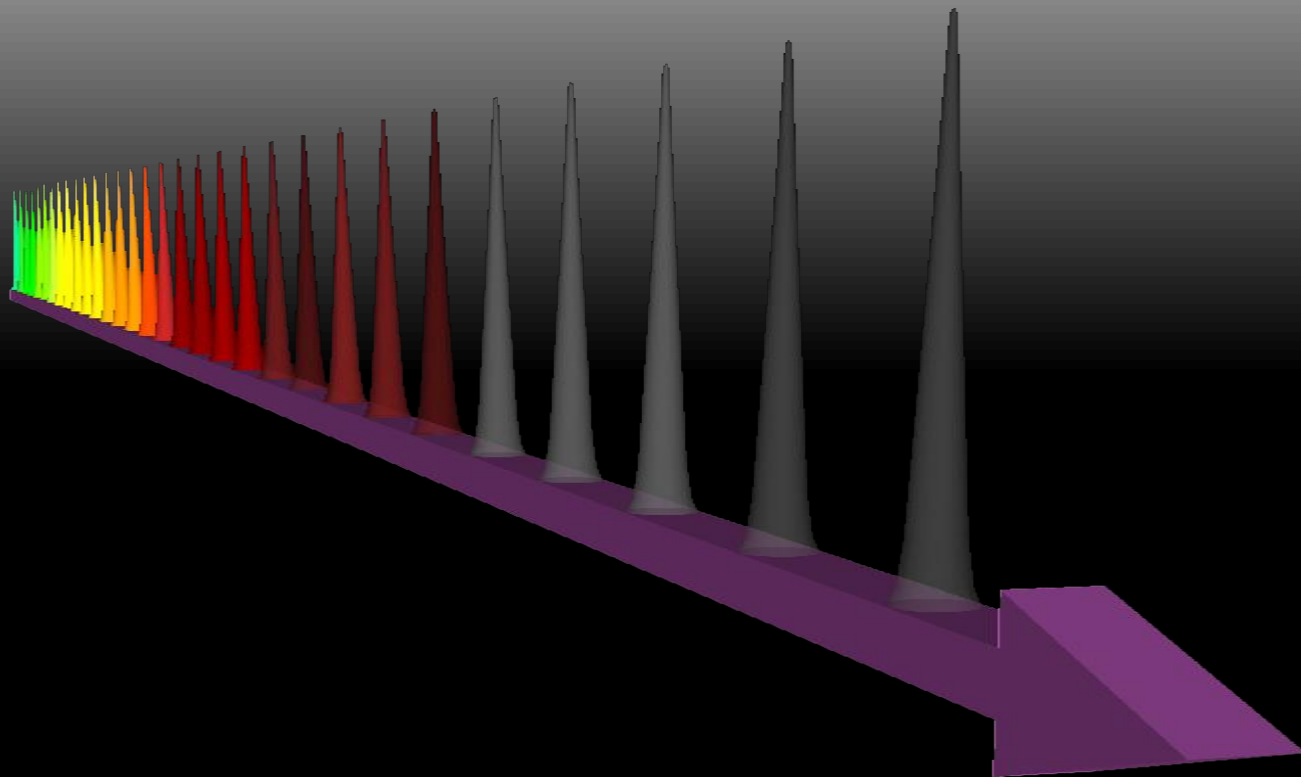
Molecular Transitions





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Moving deep into the Infrared



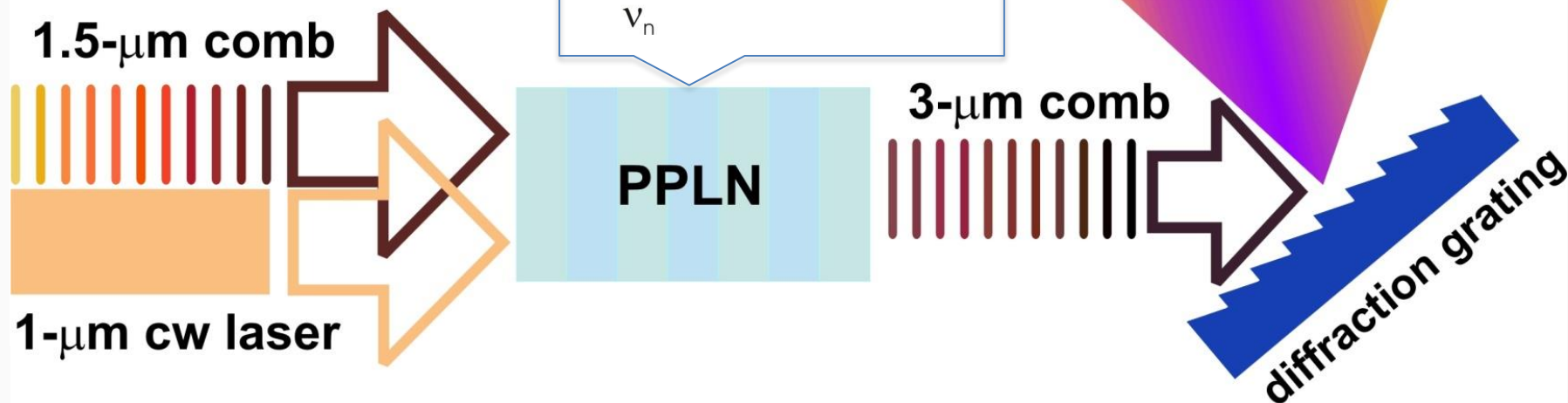


EXTENDING THE COMB TO THE MIR

only vis/NIR combs are commercially available ☐ linking direct IR sources to OFCs

Span: 180 nm (5 THz)
Power per mode: 100 pW

by tuning the pump source and
adjusting the QPM conditions, a
comb is created from from 2.9 to
3.5 μm in 180-nm-wide spans



P. Maddaloni et al., New Journal of Physics **8**, 262 (2006)



MIR-comb generation

OPTICS LETTERS / Vol. 37, No. 12 / June 15, 2012

Widely-tunable mid-infrared frequency comb source based on difference frequency generation

Axel Ruehl,¹ Alessio Gambetta,² Ingmar Hartl,³ Martin E. Fermann,³ Kjeld S. E. Eikema,¹ and Marco Marangoni^{2,*}

¹LaserLaB Amsterdam, VU University Amsterdam, de Boelelaan 1081, 1081 HV Amsterdam, The Netherlands

²Physics Department of Politecnico di Milano – Polo di Lecco, via Ghislanzoni 24, 23900 Lecco, Italy

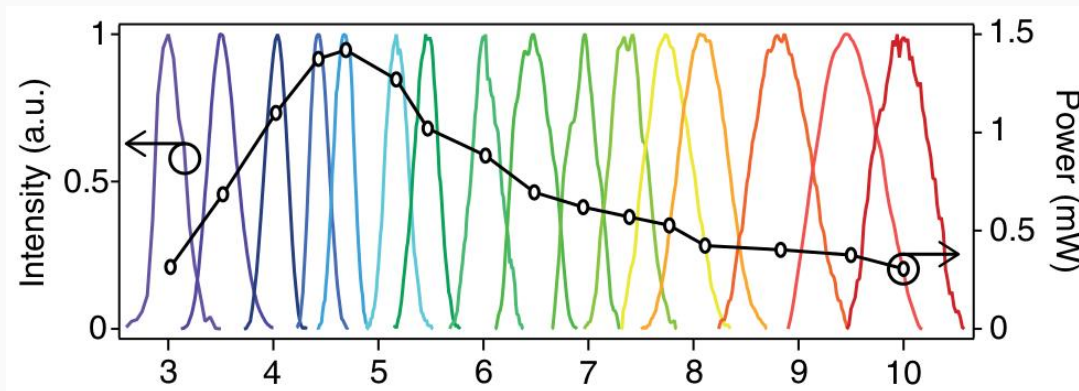
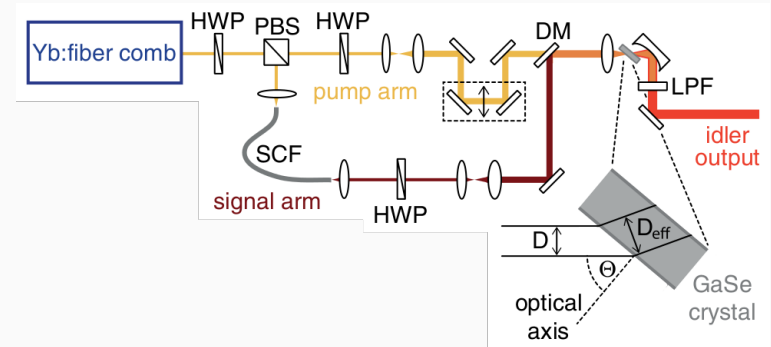
³IMRA America Inc., 1044 Woodridge Avenue, Ann Arbor, Michigan 48105-9774, USA

*Corresponding author: marco.marangoni@polimi.it

Received March 12, 2012; accepted April 23, 2012;
posted April 25, 2012 (Doc. ID 164602); published June 6, 2012

We report on a mid-IR frequency comb source of unprecedented tunability covering the entire 3–10 μm molecular fingerprint region. The system is based on difference frequency generation in a GaSe crystal pumped by a 151 MHz Yb:fiber frequency comb. The process was seeded with Raman-shifted solitons generated in a highly nonlinear suspended-core fiber with the same source. Average powers up to 1.5 mW were achieved at the 4.7 μm wavelength. © 2012 Optical Society of America

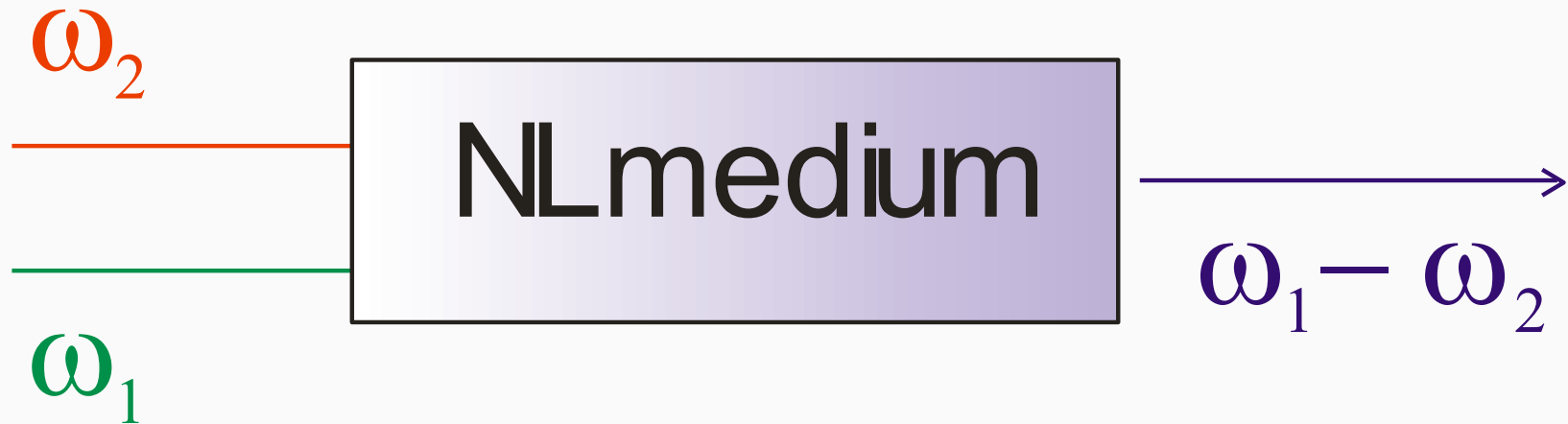
OCIS codes: 140.3070, 190.4410, 190.7110, 300.6340.



With a **GaSe** non-linear crystal they obtain a very **wide tunability** range (3-10 μm) and an overall optical power up to 1.5 mW.



Difference Frequency Generation - DFG



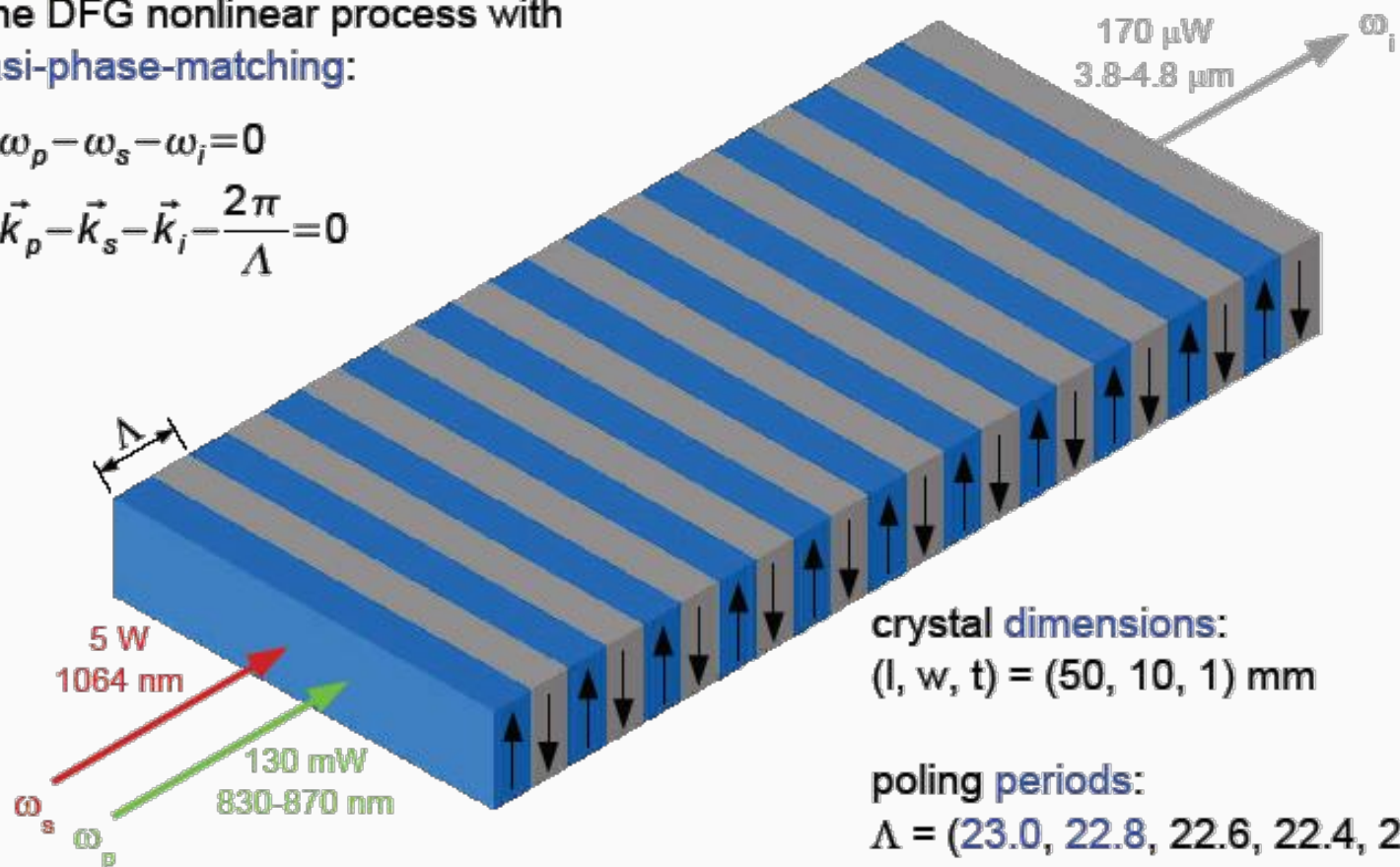


The Source: DFG in PP-Crystals

Energy and momentum conservation
in the DFG nonlinear process with
quasi-phase-matching:

$$\omega_p - \omega_s - \omega_i = 0$$

$$\vec{k}_p - \vec{k}_s - \vec{k}_i - \frac{2\pi}{\Lambda} = 0$$

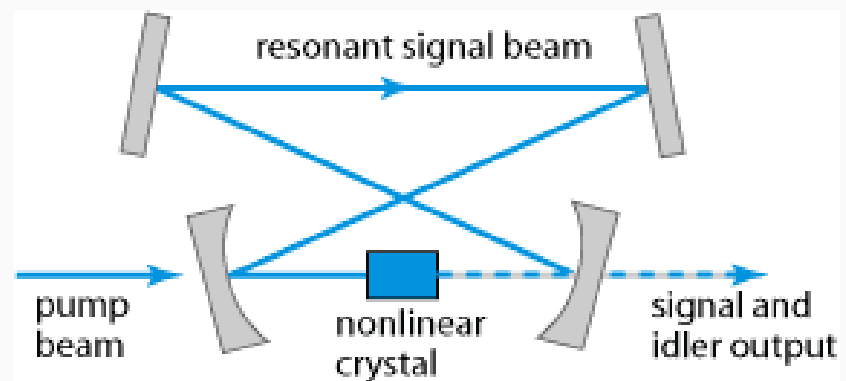
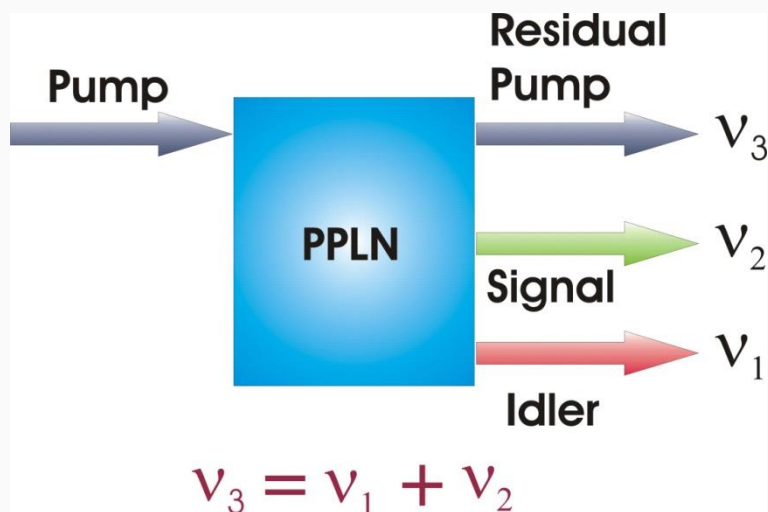
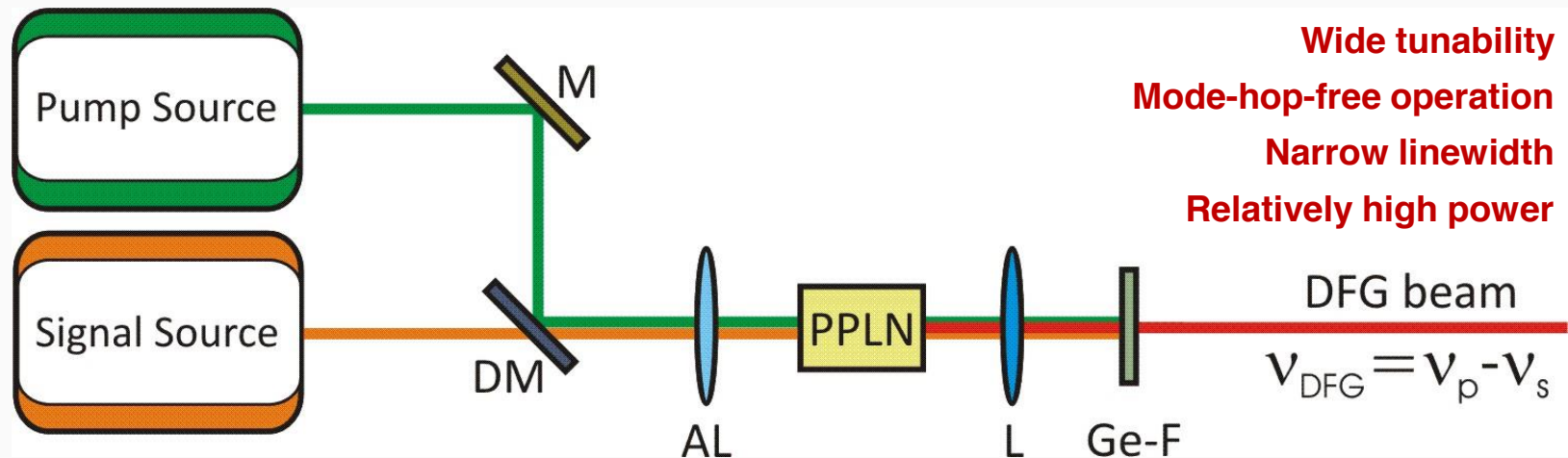


crystal dimensions:
(l, w, t) = (50, 10, 1) mm

poling periods:
 $\Lambda = (23.0, 22.8, 22.6, 22.4, 22.2,$
 $22.0, 21.8, 21.6, 21.4) \mu\text{m}$



Optical Frequency-Down Conversion

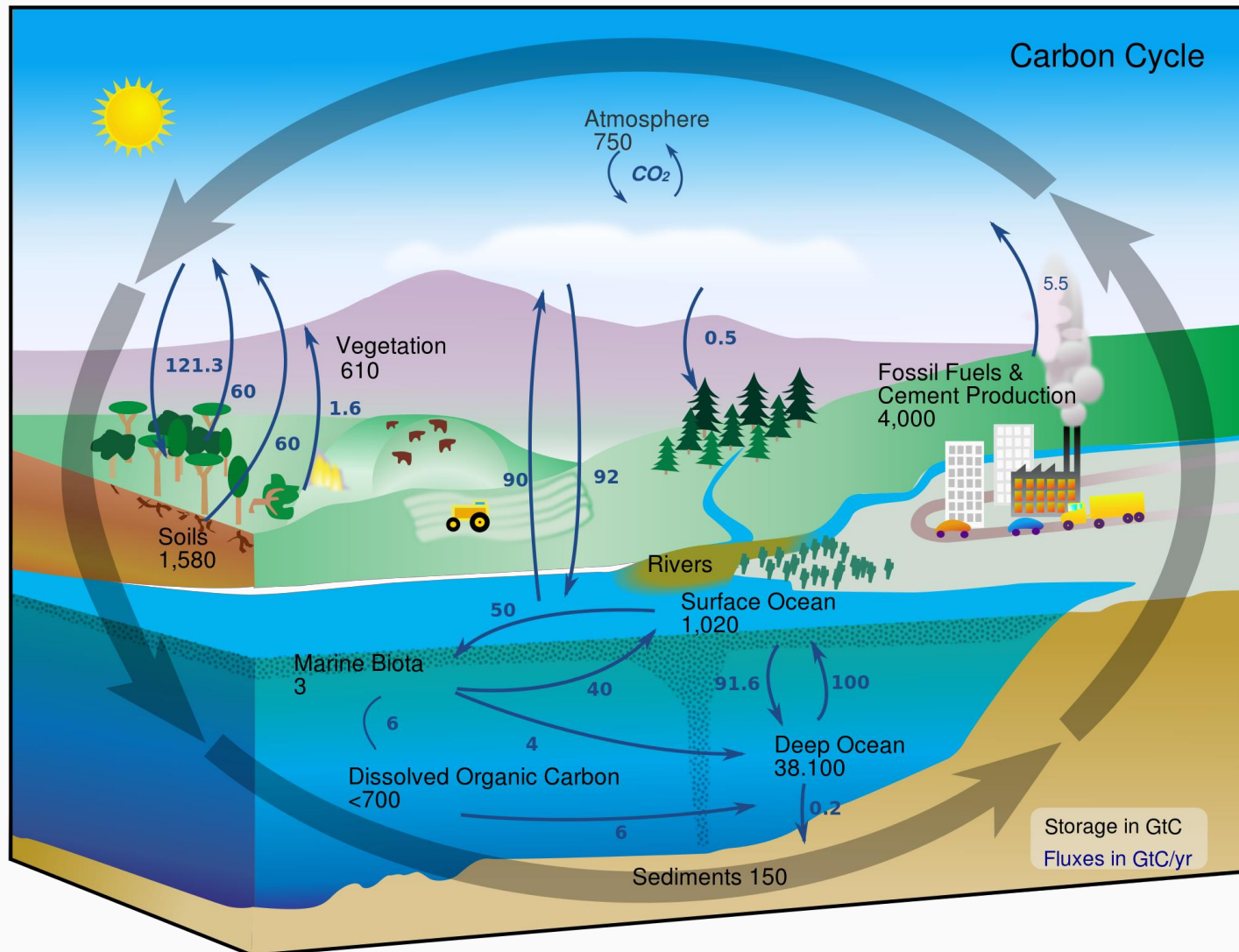


Much higher power, but cavity is needed to be frequency-stable and with single-mode behaviour



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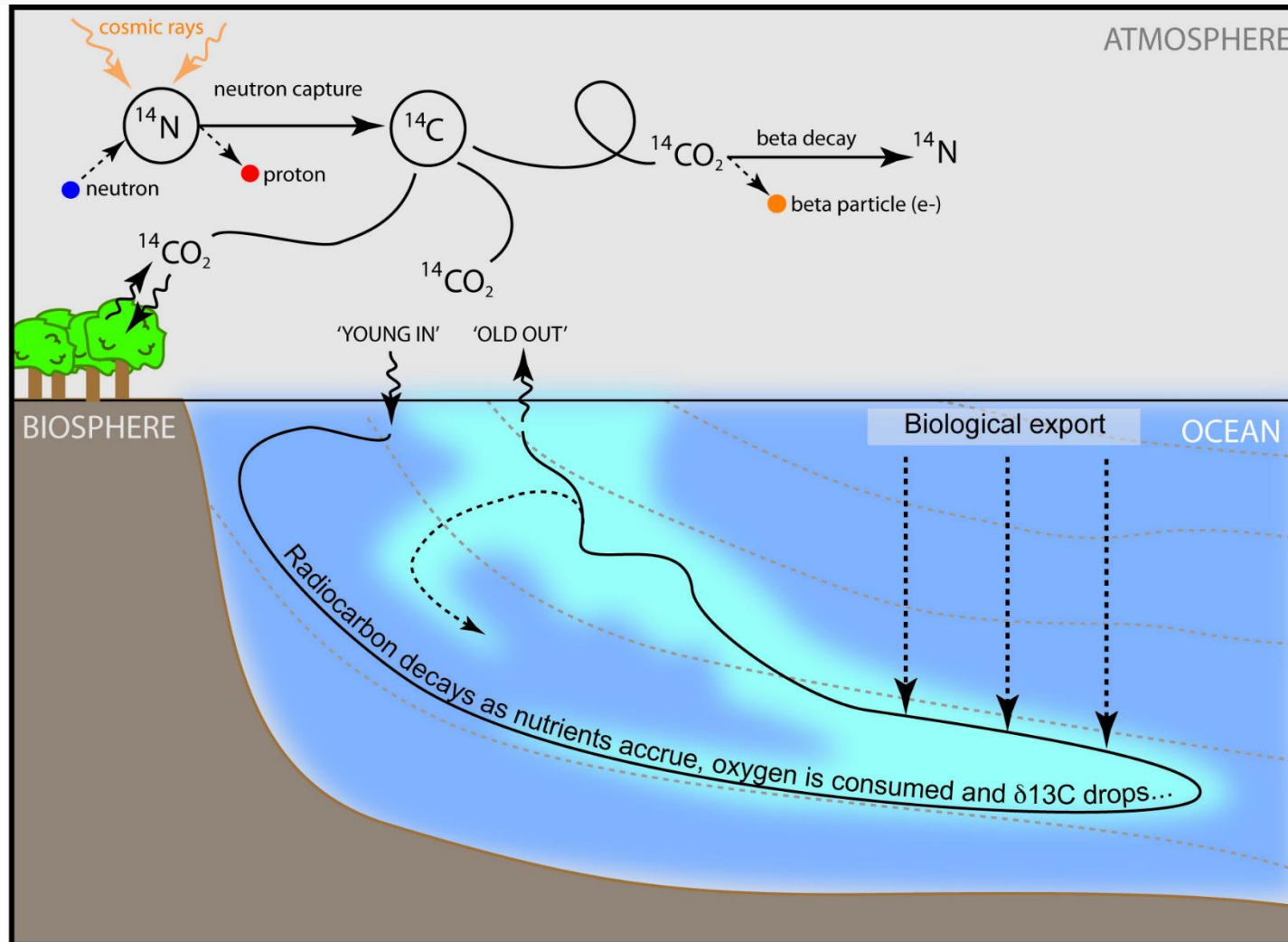
Carbon cycle...





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...and radiocarbon cycle





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The New York Times

ASIA PACIFIC | NEWS ANALYSIS

179 COMMENTS

Deal on Carbon Emissions by Obama and Xi Jinping Raises Hopes for Upcoming Paris Climate Talks

By CORAL DAVENPORT NOV. 12, 2014





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ENVIRONMENT

Energy and T for Managing

Aristides A. N. Patrinos^{1*} and Rich

Despite some uncertain
scientific and politic
is that the level of glo
of greenhouse gases (GHGs)
to atmospheric concentration
between 450 and 500 parts per
(1) to avoid serious, if not catast
on life and property. Achiev
poses some formidable challer
inertia in the climate system (C
for generations), as well as in (C
capital investment. Furthermor
nomic sector and country en
these challenges, a broad range
be required.

Despite some uncertainties, today's scientific and political consensus is that the level of global emissions of greenhouse gases (GHGs) needs to lead to atmospheric concentrations somewhere between 450 and 500 parts per million (ppm) (1) to avoid serious, if not catastrophic, effects on life and property. Achieving this goal poses some formidable challenges. There is inertia in the climate system (GHGs survive for generations), as well as in GHG-emitting capital investment. Furthermore, every economic sector and country emits. To meet these challenges, a broad range of actions will be required.



The Challenge

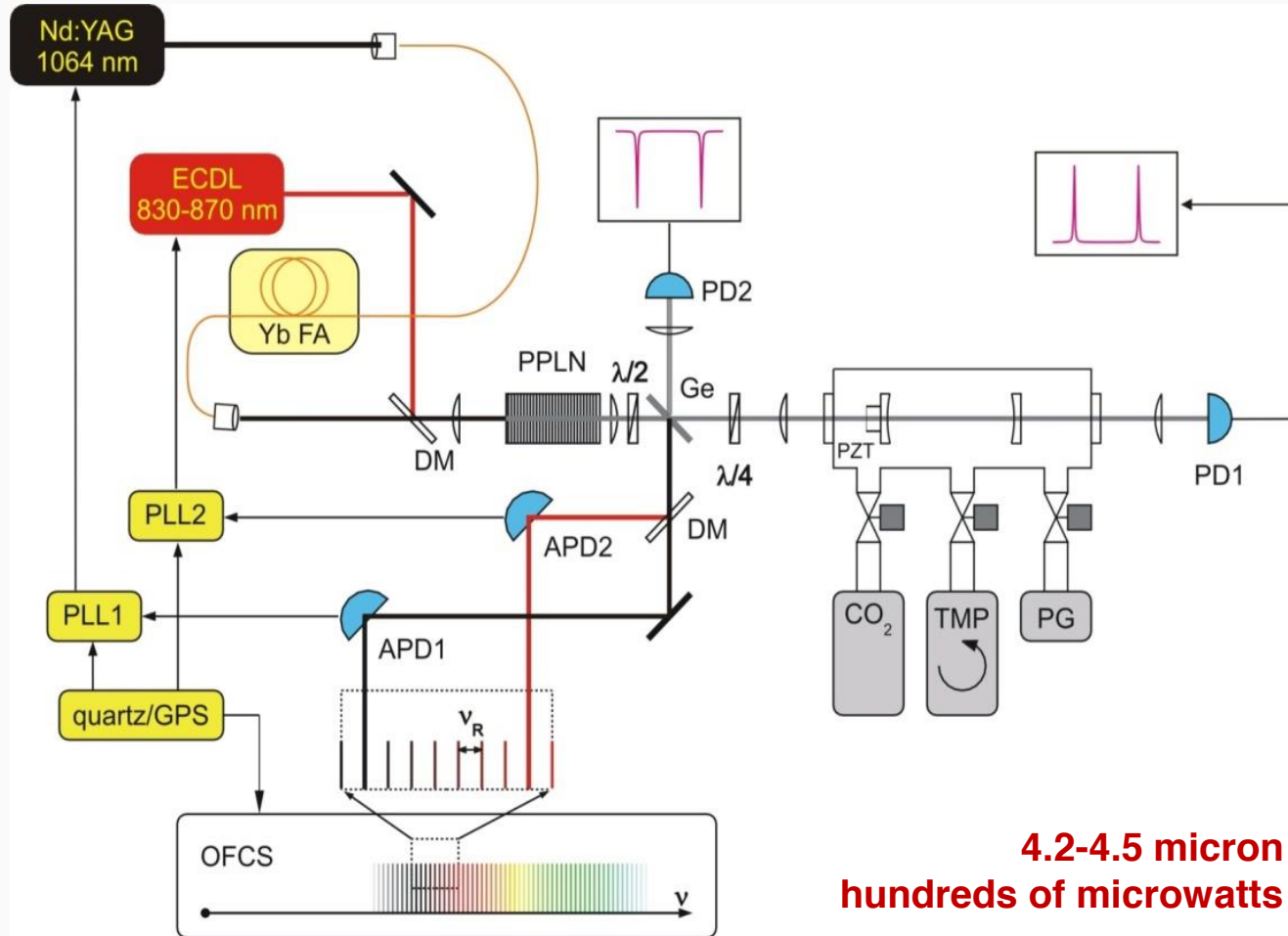
To meet these formidable challenges, sensors are needed with:

- High sensitivity, for trace gas sensing
- High resolution, for a high degree of discrimination
- Infrared spectral coverage, to match strong molecular bands
- Wide tunability, to interrogate a large number of molecules
- High precision, to reduce databases statistical uncertainties
- High accuracy, to reduce systematic effects getting high reproducibility



Laser-based Infrared Spectrometers

A “simple” Comb-Assisted DFG Setup



$$f_p = n_p f_r + f_0 + f_{beat}^p$$

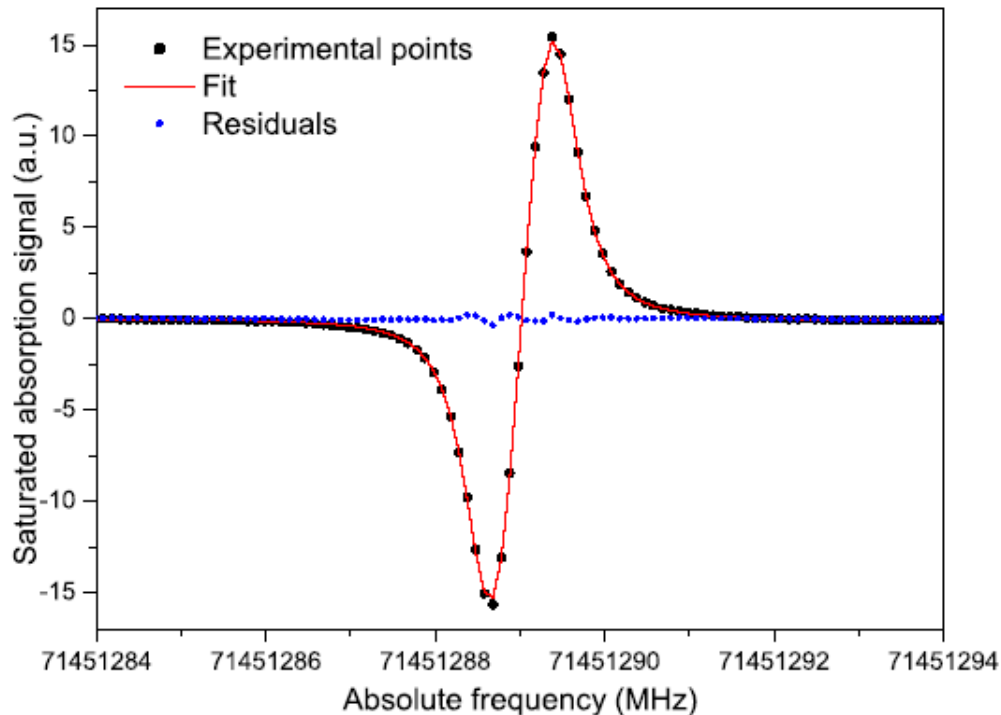
$$f_s = n_s f_r + f_0 + f_{beat}^s$$



$$f_{DFG} \equiv f_p - f_s = (n_p - n_s) f_r + (f_{beat}^p - f_{beat}^s)$$



Cavity-Enhanced Saturation Spectroscopy



CO₂ (00⁰1–00⁰0) R(56)
saturated-absorption
line at 2383.359 cm⁻¹

D. Mazzotti, et al., Opt. Lett. **30**, 997 (2005)

**uncertainty of 800 Hz (10⁻¹¹)
in the absolute frequency**

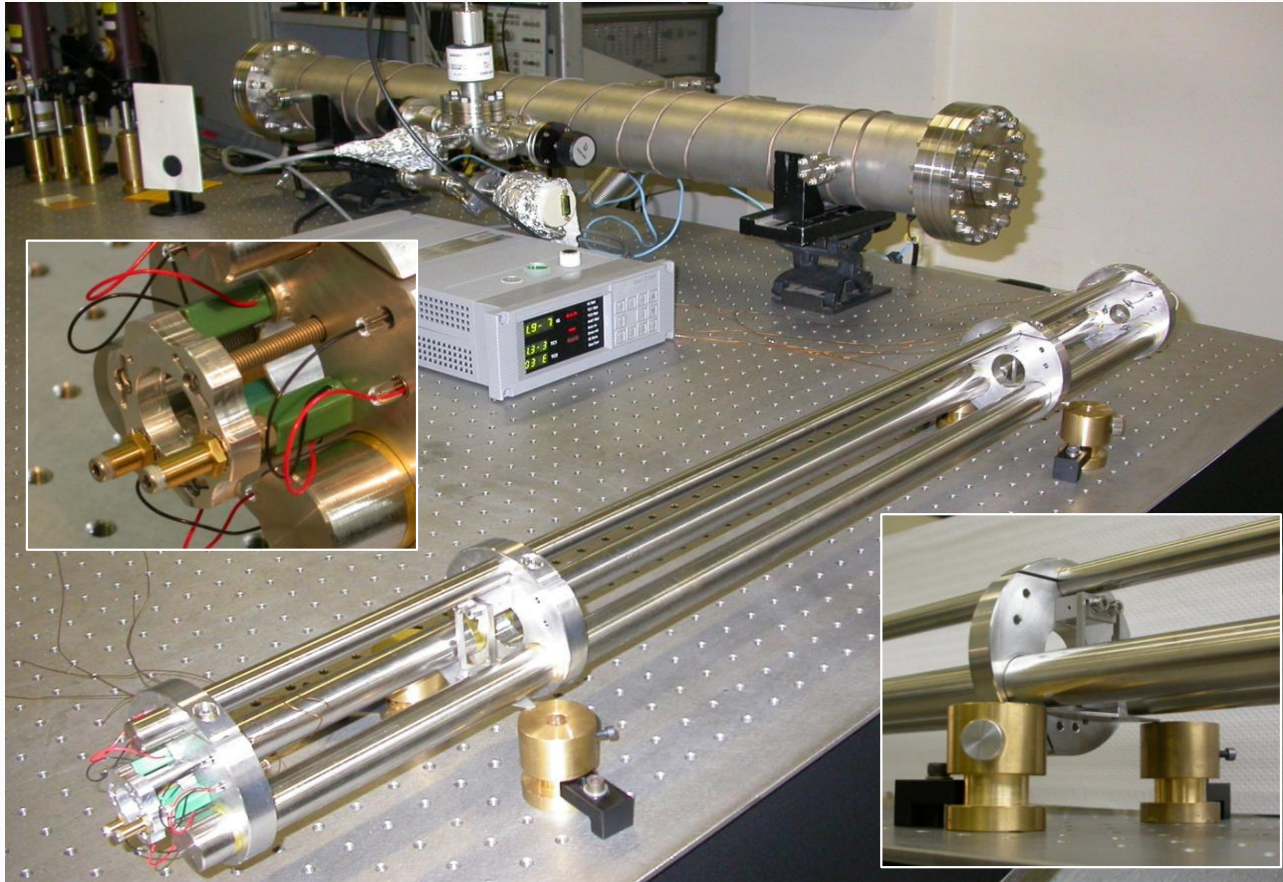
Source	Expression ^a	Actual Value
Natural lifetime	$k^3 \mu^2 / 3\pi \epsilon_0 \hbar$	0.2 kHz
Collisions	$c p P$	90 kHz
Transit time	$\sqrt{\ln 2 / \pi} \sqrt{k_B T / m} 1 / w$	400 kHz
IR jitter	$[1 - (\nu_s / \nu_p)] \Gamma_c$	<50 kHz
Power	$\sqrt{1 + (I / I_s)}$	$\times 1.4 [(I / I_s) = 1]$
Modulation		$\times 1.3 (\beta = 1)$

**Main limitation
comes from
transit time**



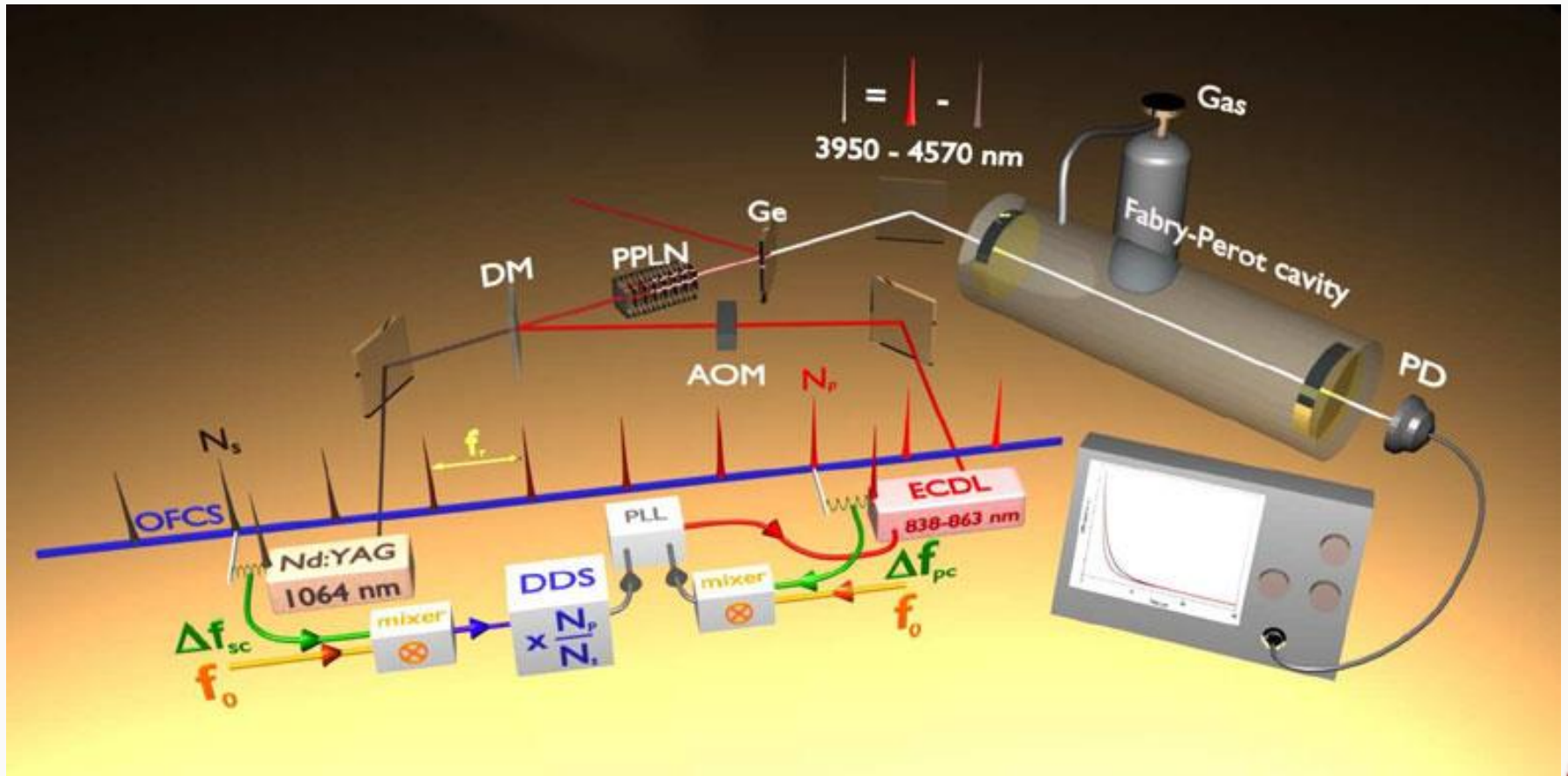
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Pushing up Sensitivity by Cavity Enhancement



$11000 < F < 25000$ @ $4.5 \mu\text{m}$
FSR = 150 MHz

Metrological spectroscopy: a smart DFG link to an Optical Frequency Comb Synthesizer

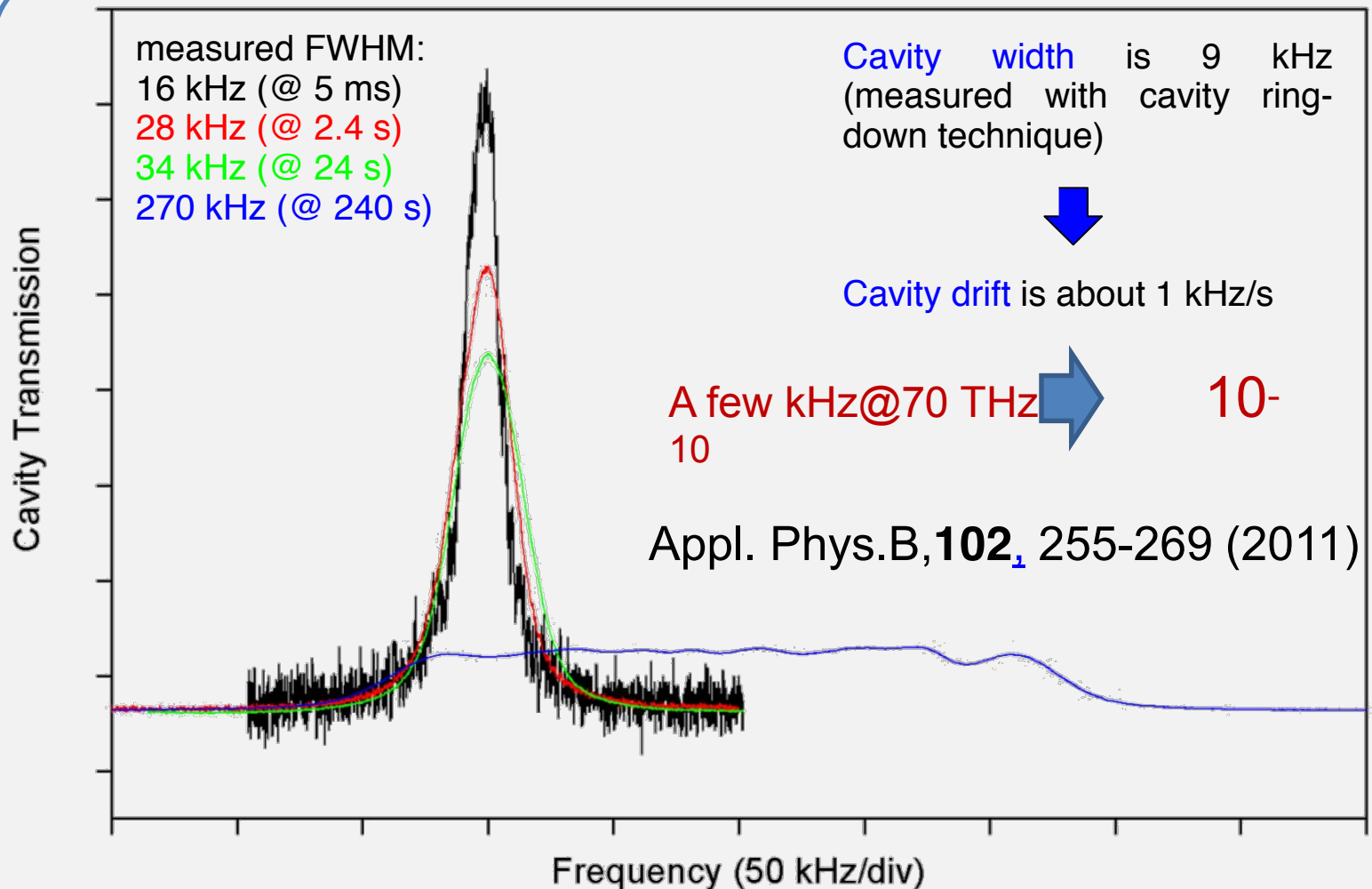


I. Galli et al., Optics Express 17, 9582 (2009)

original scheme proposed by Telle et al., Appl. Phys. B (2002))

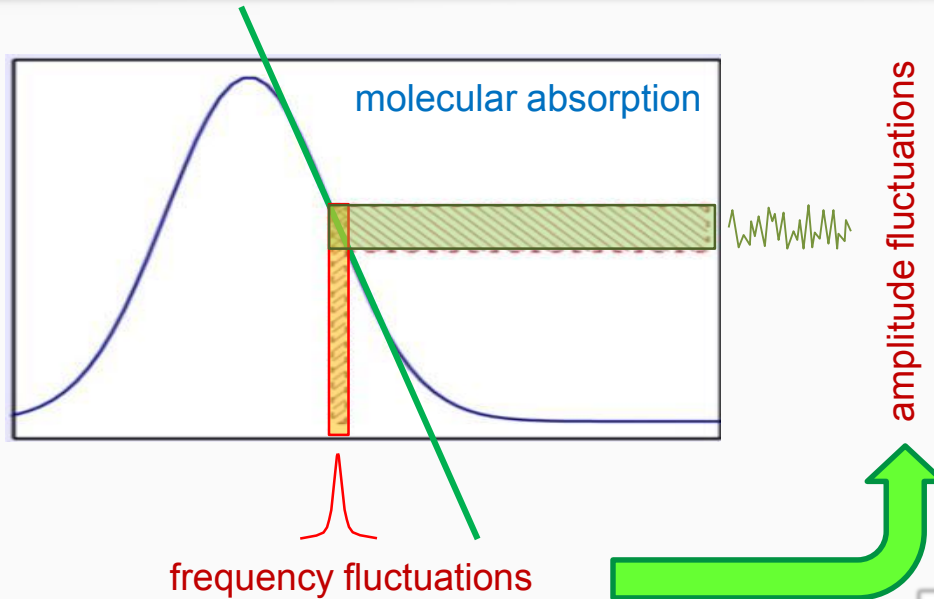


Coupling IR radiation into an enhancement cavity

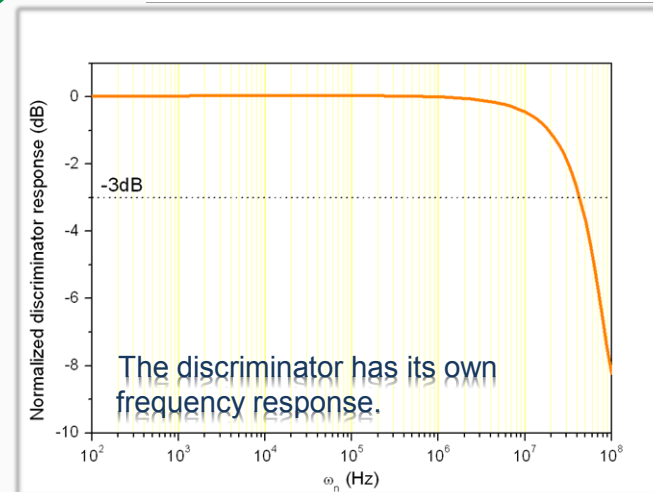
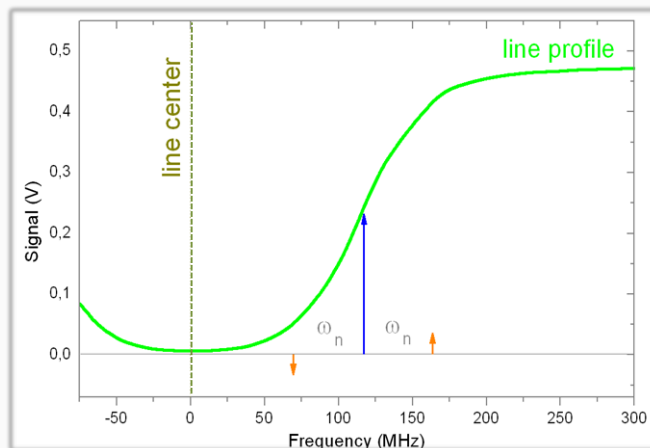
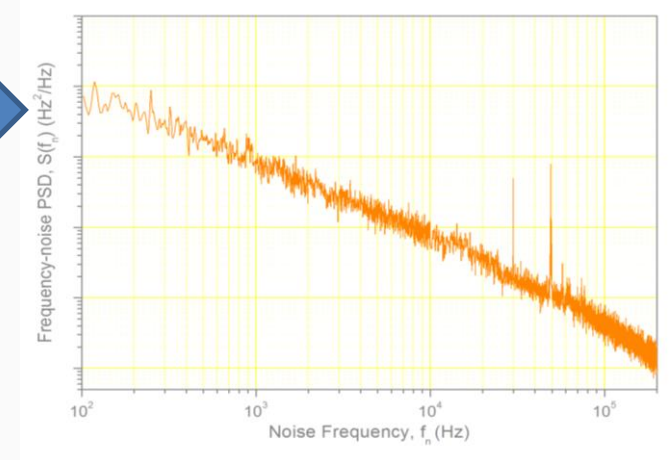




Measuring the Laser frequency-noise: principle

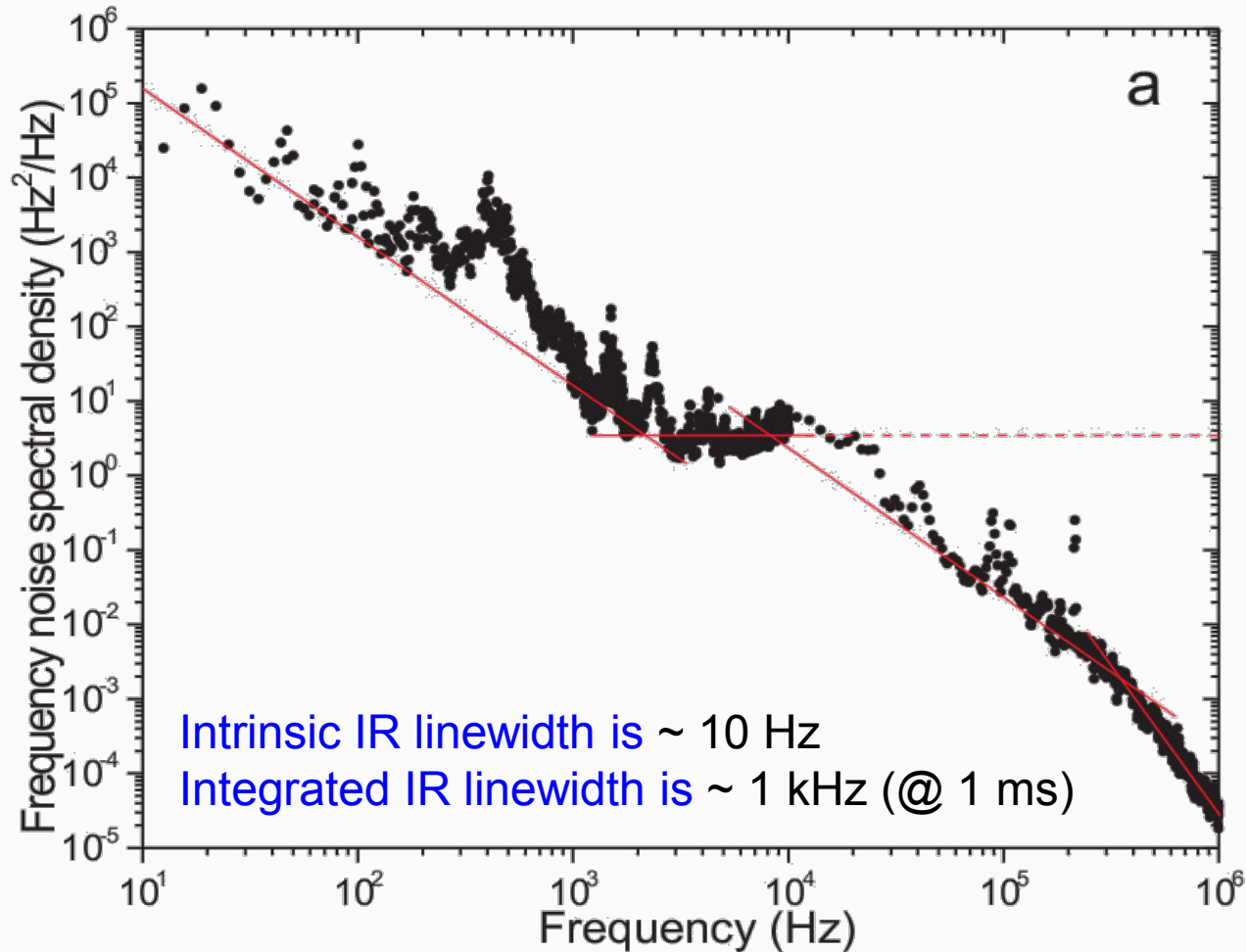


Frequency-noise
power spectral density



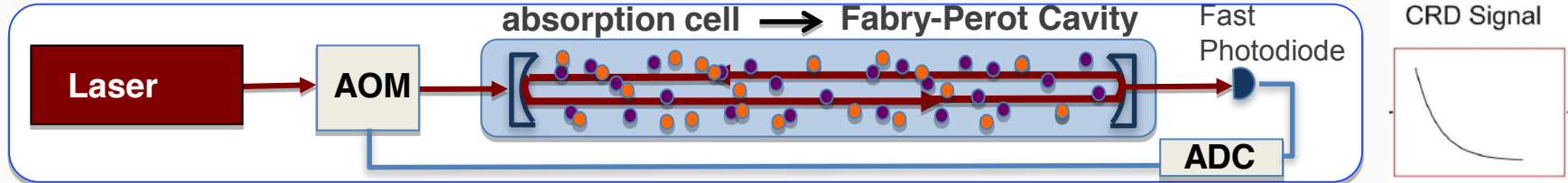


Narrowing the IR linewidth

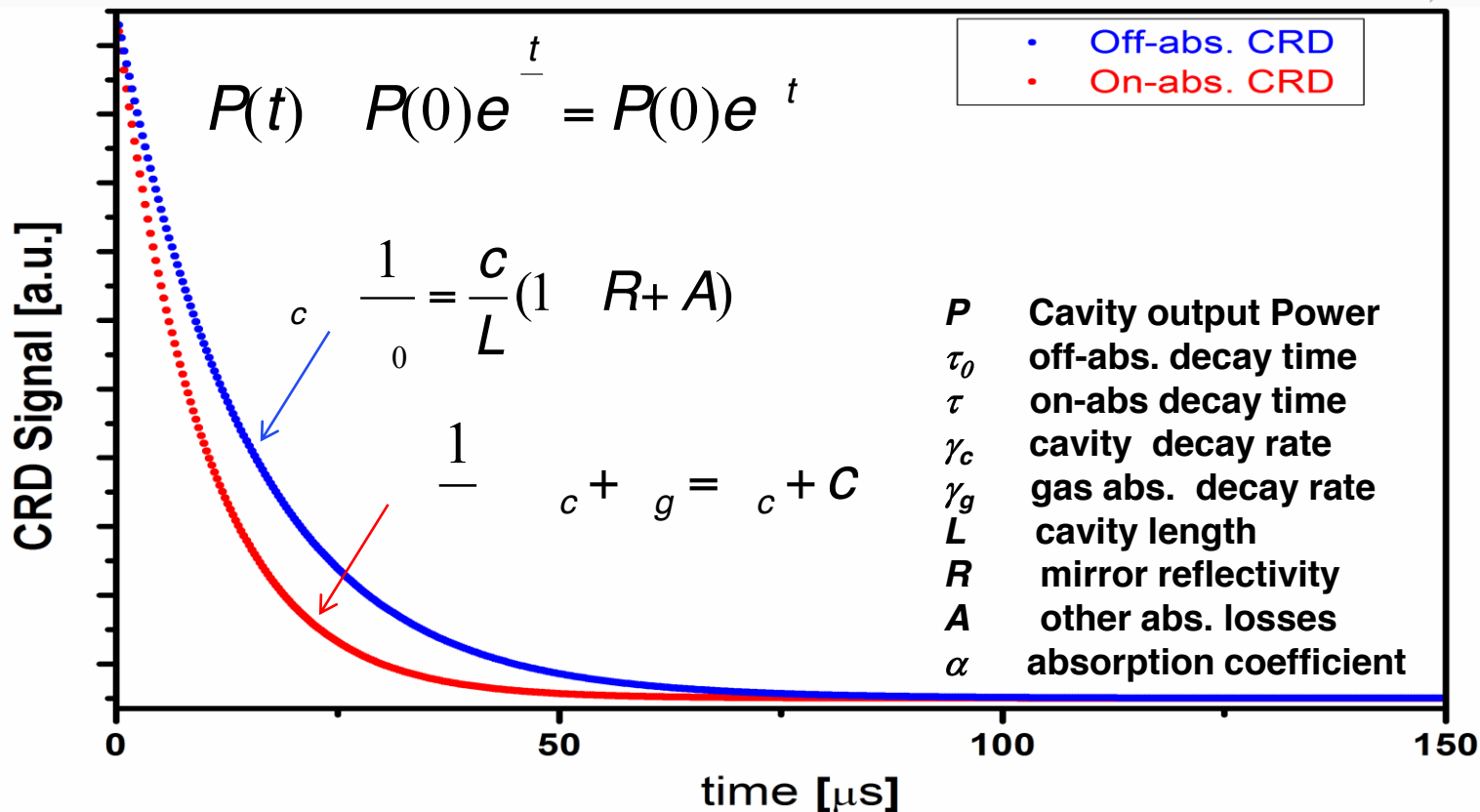


Ultra-stable, widely tunable and absolutely linked mid-IR coherent source
I. Galli et al., Optics Express **17**, 9582 (2009)

Linear Cavity Ringdown vs....

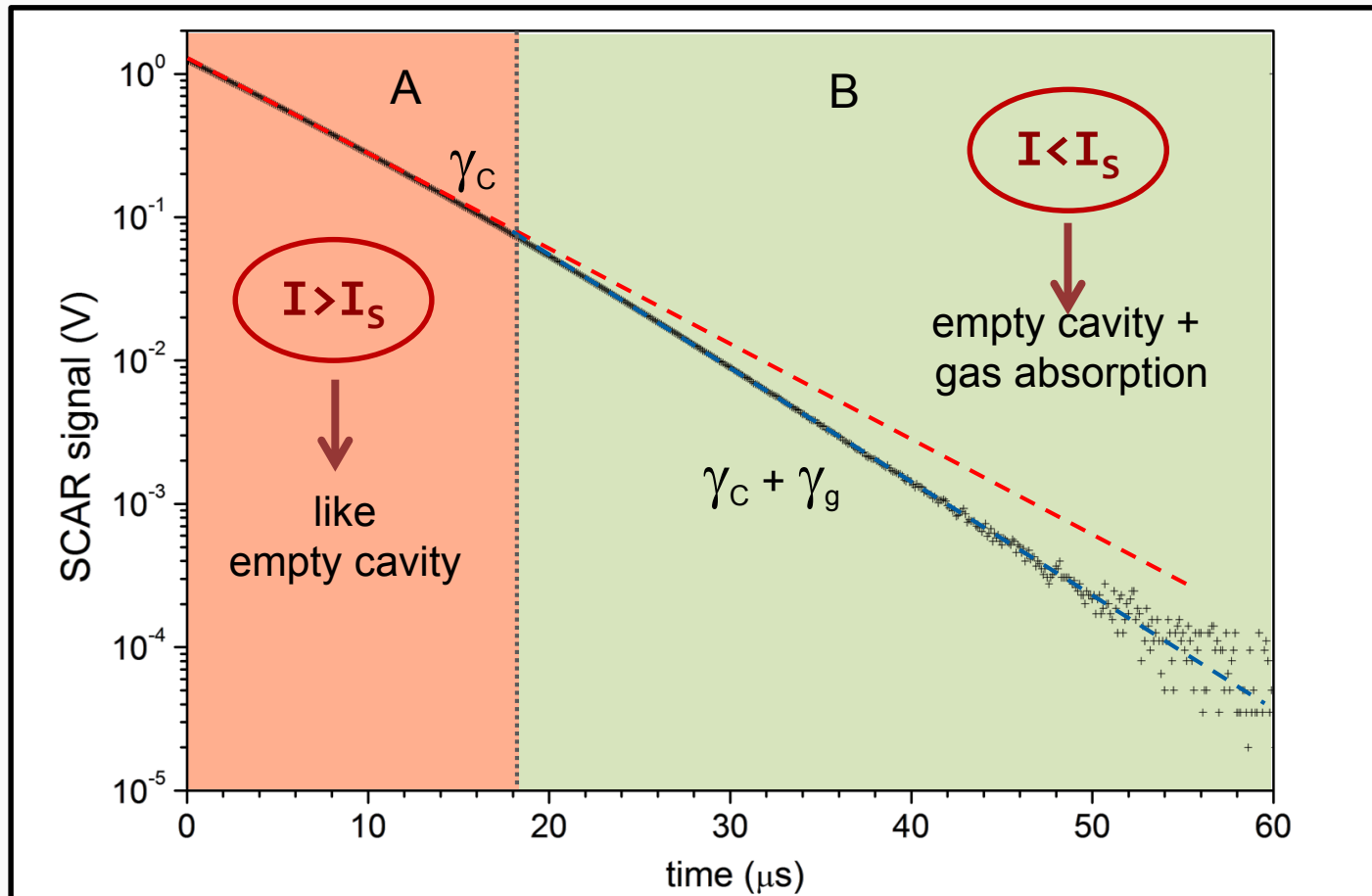


Measurement of the **Cavity Ring Down** time variations  **$C\alpha$**





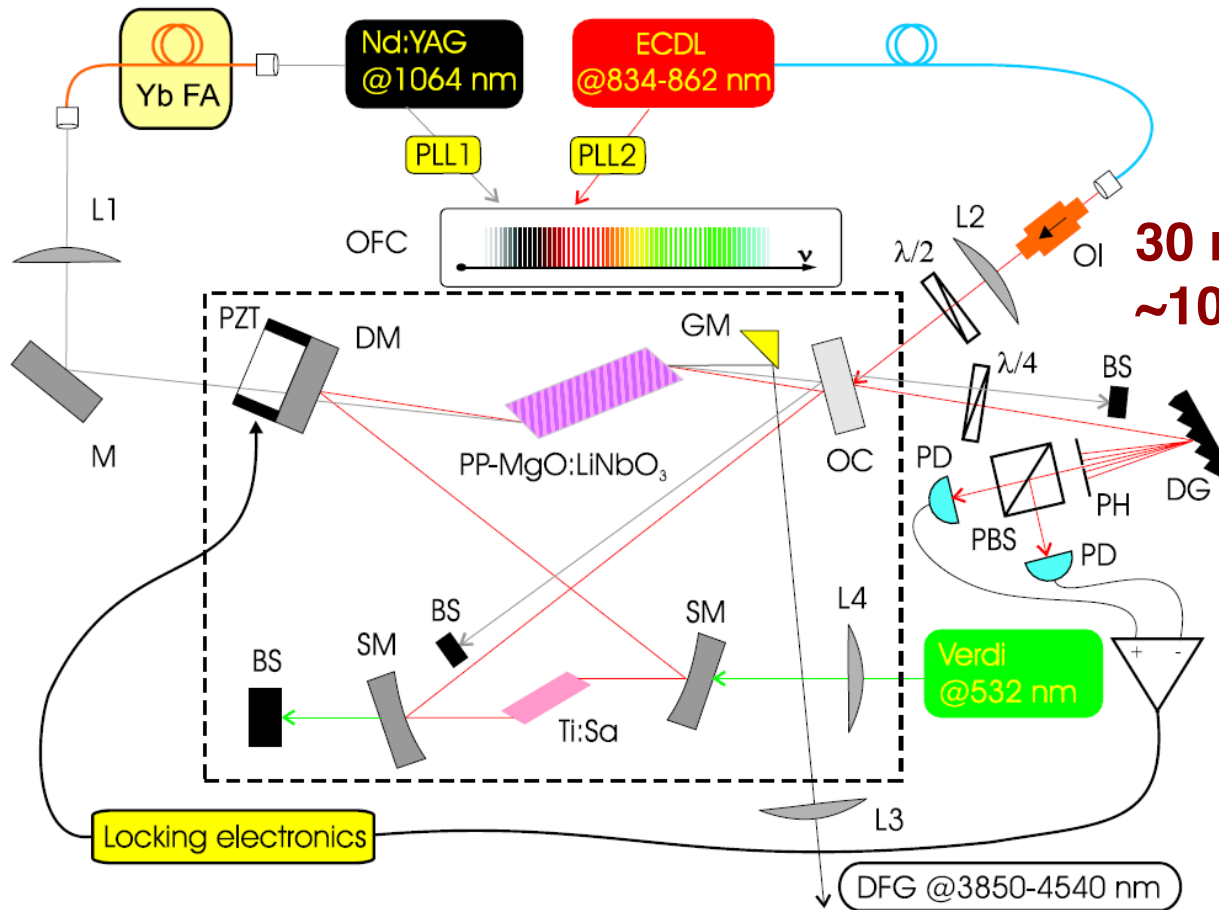
....Nonlinear Cavity Ringdown: Saturated absorption Cavity Ringdown-SCaR



Slope increase from region **A** to **B** is proportional to the absorption coefficient and hence to the **gas concentration**



Nonlinear absorption phenomena require power: Intra-cavity DFG



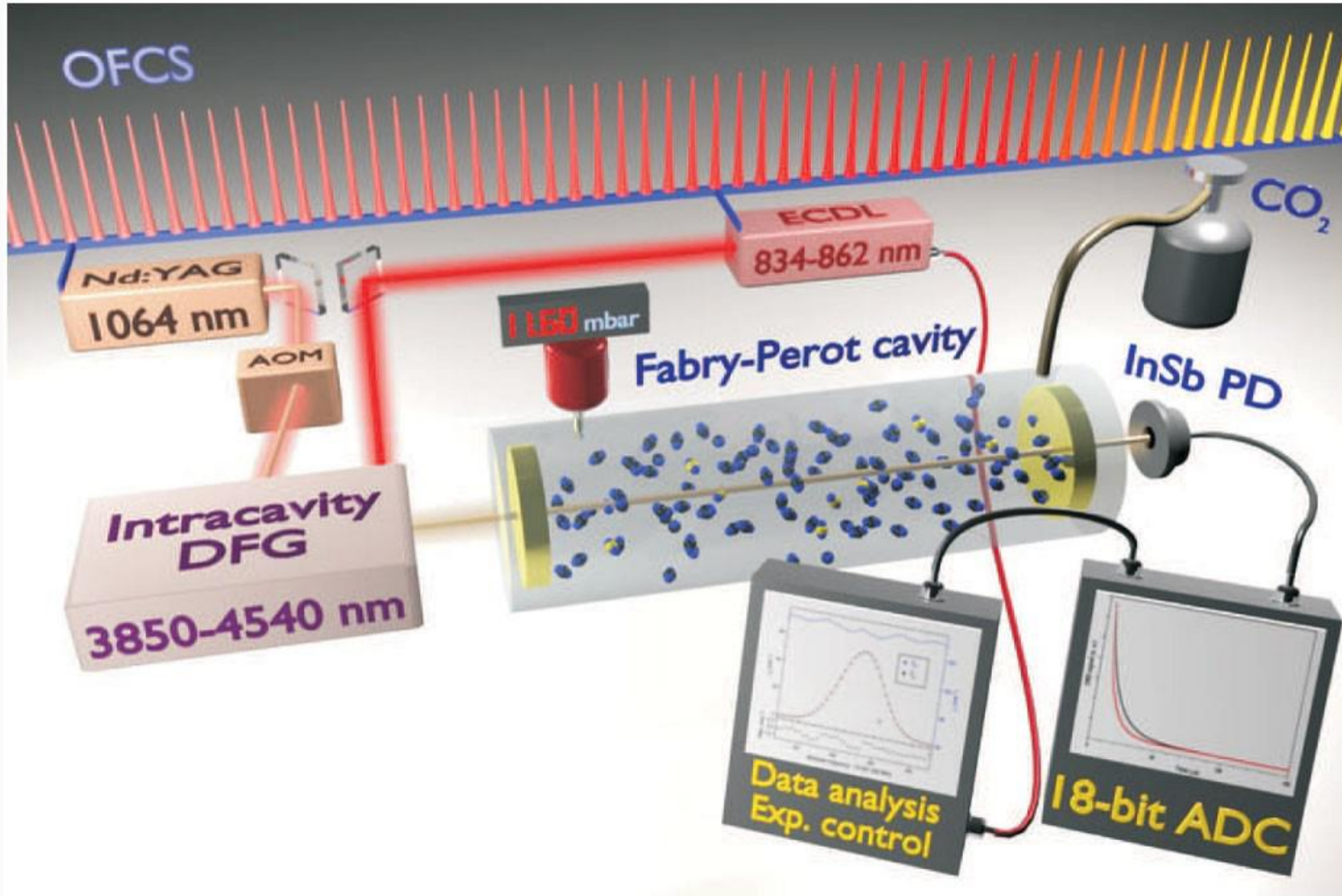
**30 mW cw power @ 4.5 μm
~10Hz linewidth**

Opt. Lett. 35, 3616 (2010)



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Combining Precision and Sensitivity: The SCAR set-up



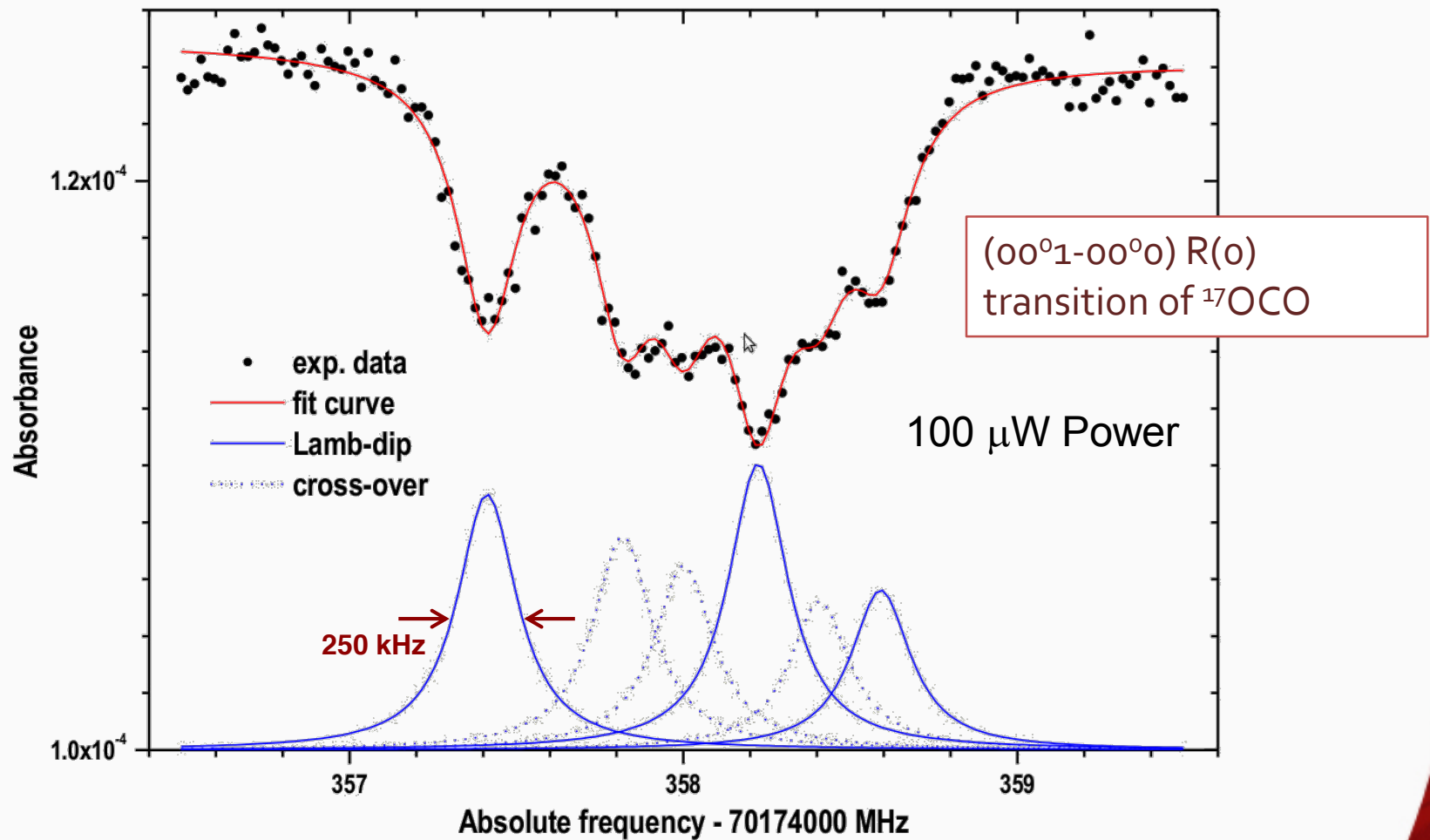


Hunting $^{14}\text{CO}_2$: Background and motivations

- ✓ Radiocarbon (^{14}C) is a key isotopic species for analysis in many different sectors worldwide. It is contained in $^{14}\text{CO}_2$ BUT it is a VERY elusive species (about 1 ppt concentration in natural abundance).
- ✓ To date, two main techniques are widely used to measure radiocarbon concentration, both with some disadvantages:
 - Liquid Scintillation Counting (LSC) has very long measurement times (hours-days), requires sample with complex preparation and big mass (1-10 g of C);
 - Accelerator Mass Spectrometry (AMS) requires large (7-200 m²), expensive (0.3-3 M€) and high-maintenance facilities, with slow turnaround time (1-4 weeks) and high costs (400-1000 €) for each measured sample.
- ✓ High-sensitivity mid-IR laser spectroscopy has recently approached the performance of these two techniques, while reducing many of their disadvantages.



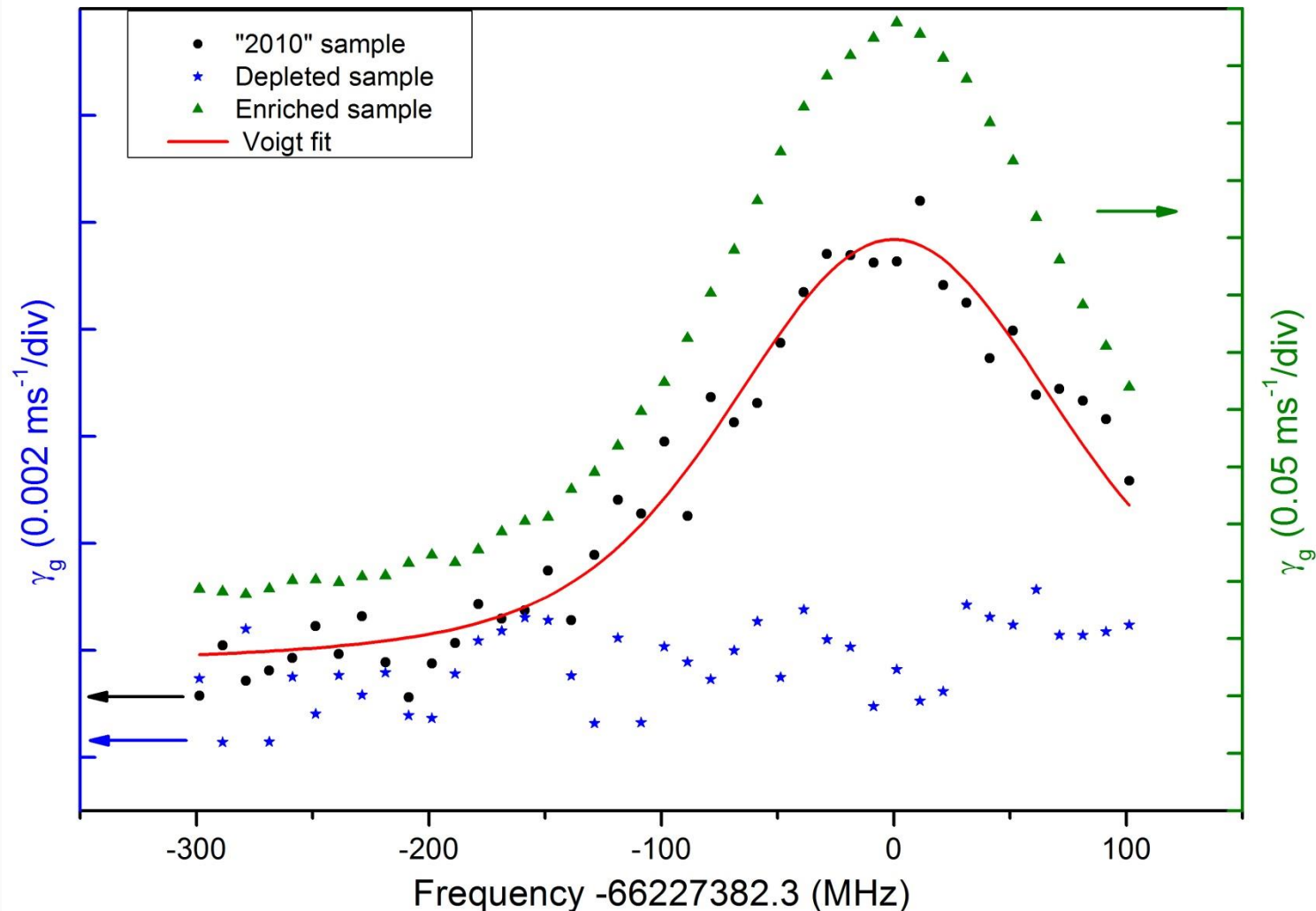
SCAR: the resolution



Phys. Rev. Lett. **104**, 110801 (2010)

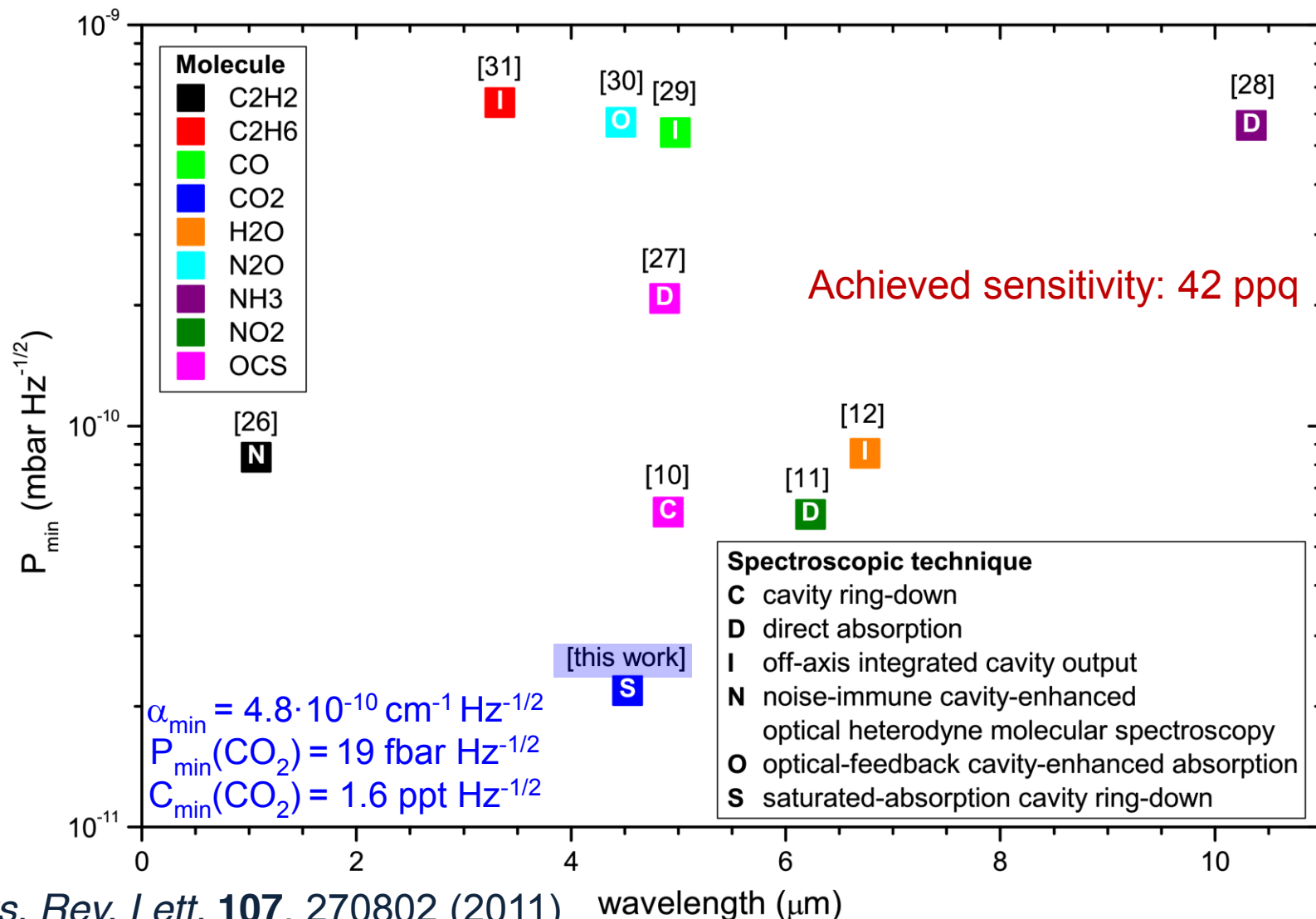


Natural abundance (1×10^{-12}) detection of $^{14}\text{C}^{16}\text{O}_2$



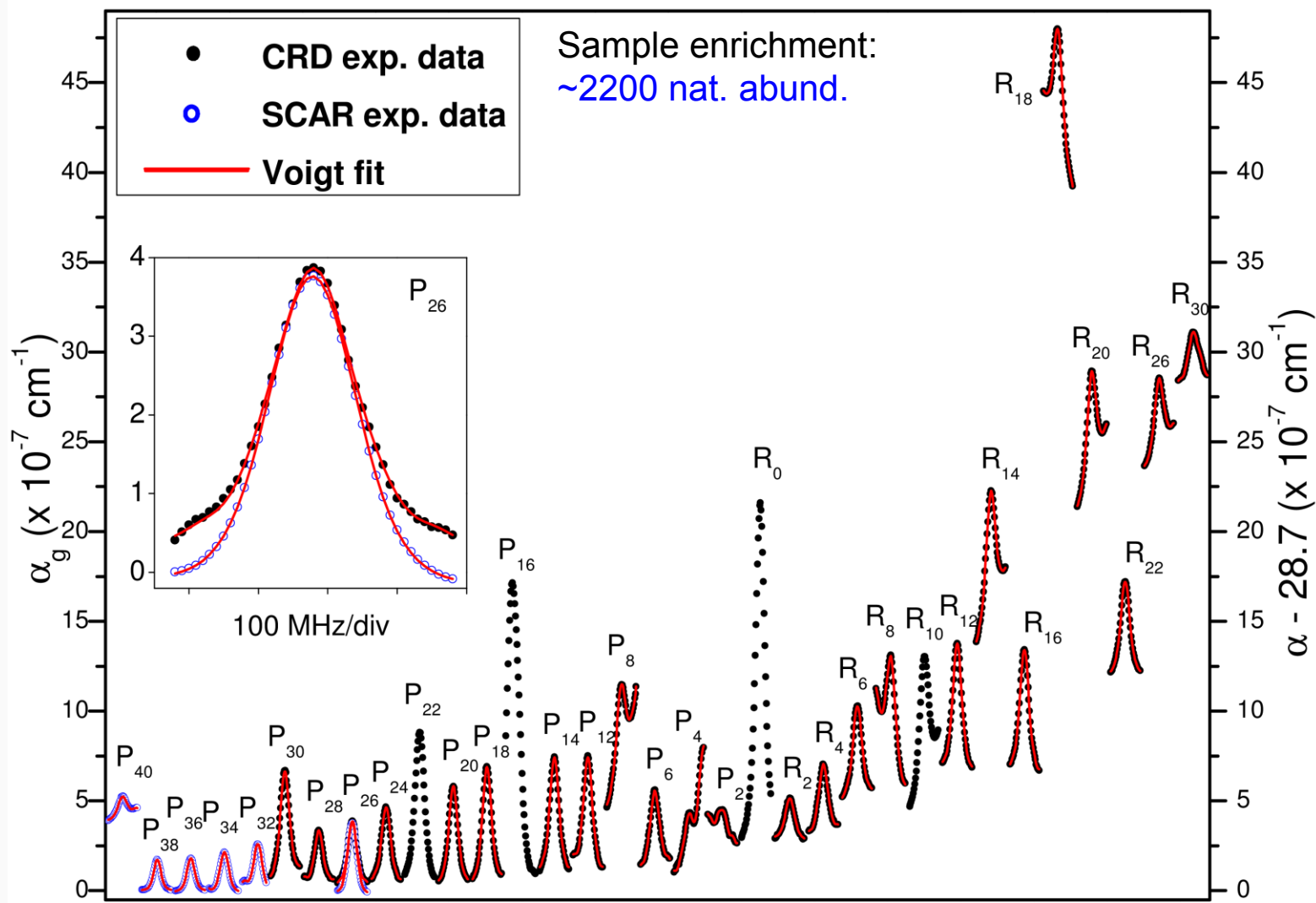


SCAR: the sensitivity





Spectroscopy of the $^{14}\text{C}^{16}\text{O}_2$ ν_3 band





Retrieved spectroscopic parameters

Table 1. Transitions frequencies, uncertainties and residuals from the fit of the lines measured for the ν_3 fundamental band of $^{14}\text{C}^{16}\text{O}_2$. All numbers are in MHz.

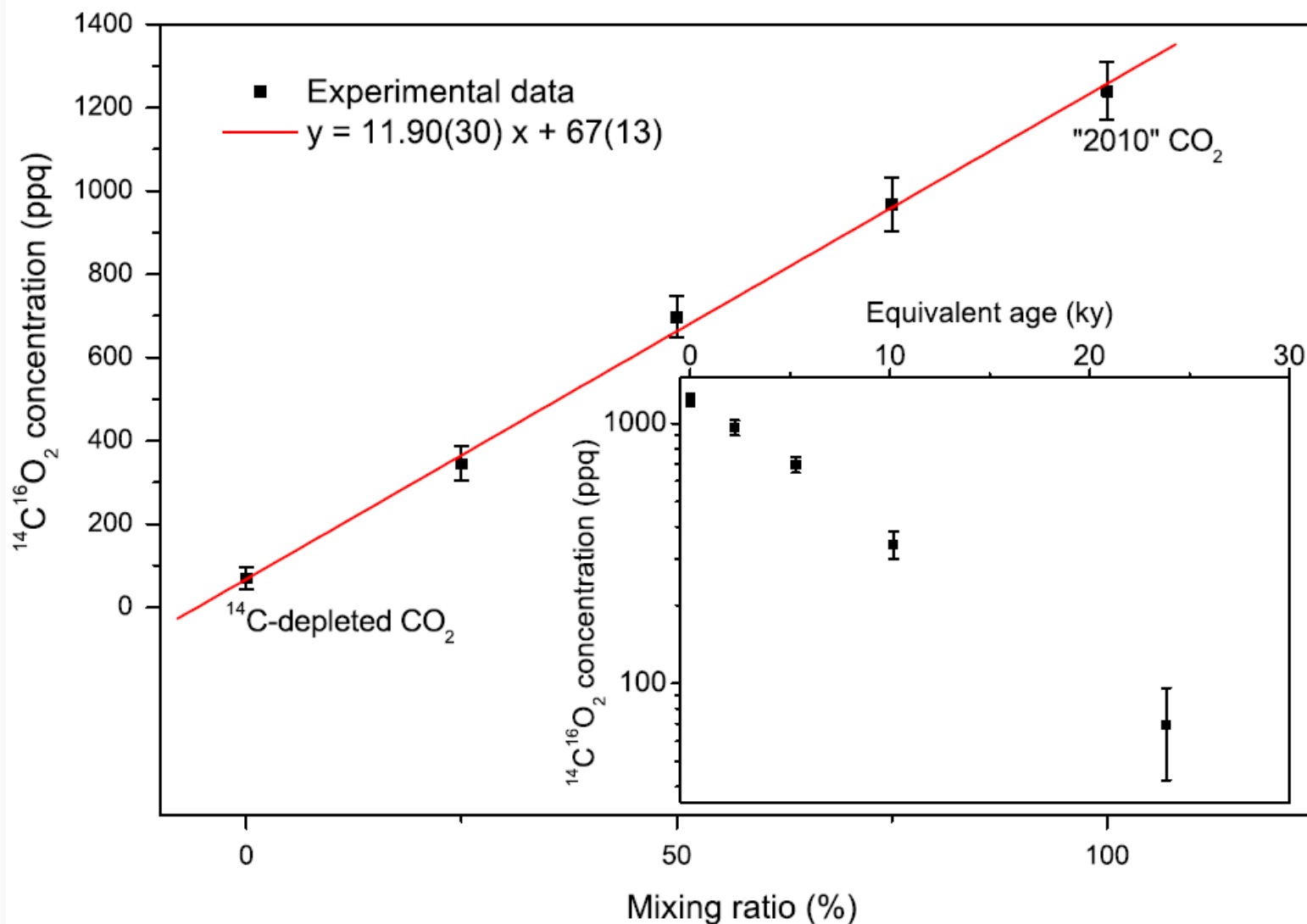
Transition	Obs. Freq.	Uncertainty ^a	Obs.-Calc. ^b
R(30)	67367614.181	4.795	1.759
R(26)	67294430.381	0.825	2.317 *
R(22)	67218457.563	0.875	0.047
R(20)	67179428.939	0.743	-0.260
R(18)	67139706.163	0.731	-0.314
R(16)	67099289.625	1.017	-0.450
R(14)	67058180.523	0.877	-0.197
R(12)	67016379.675	0.828	0.529
R(8)	66930706.117	1.178	3.834 *
R(6)	66886832.920	1.111	4.442 *
R(4)	66842264.775	1.112	-0.645
R(2)	66797014.526	2.003	0.666
P(2)	66680873.788	2.765	-3.952
P(4)	66633216.481	3.176	-5.929
P(6)	66584878.904	1.132	-3.113 *
P(8)	66535859.047	1.443	1.710
P(12)	66435758.242	0.996	-0.005
P(14)	66384685.331	2.458	-0.078
P(18)	66280498.427	0.801	1.310
P(20)	66227383.132	0.865	-0.128
P(24)	66119121.905	1.715	1.758
P(26)	66063972.706	0.755	0.194
P(28)	66008145.242	1.749	-3.339
P(32)	65894474.448	0.747	-0.680
P(34)	65836627.418	0.742	0.147
P(36)	65778106.638	0.729	0.186
P(38)	65718914.039	0.752	0.510
P(40)	65659047.695	1.795	-1.677

Table 2. Spectroscopic parameters (in cm^{-1}) obtained for the ν_3 fundamental band of $^{14}\text{C}^{16}\text{O}_2$.

Parameter	This work ^a	Ref. [6]
ν_0	2225.801399(17)	2225.80239(16)
B_0	0.390253082(58)	0.39025488(18)
$D_0 \times 10^{-7}$	1.33144(37)	1.3372(20)
B_1	0.38739025628 ^b	
$D_1 \times 10^{-7}$	1.327523 ^b	
$H_1 \times 10^{-14}$	-3.08 ^b	
$L_1 \times 10^{-17}$	1.244 ^b	
$\sigma_{\text{FIT}} \times 10^{-5}$	3.1 → 930 kHz	

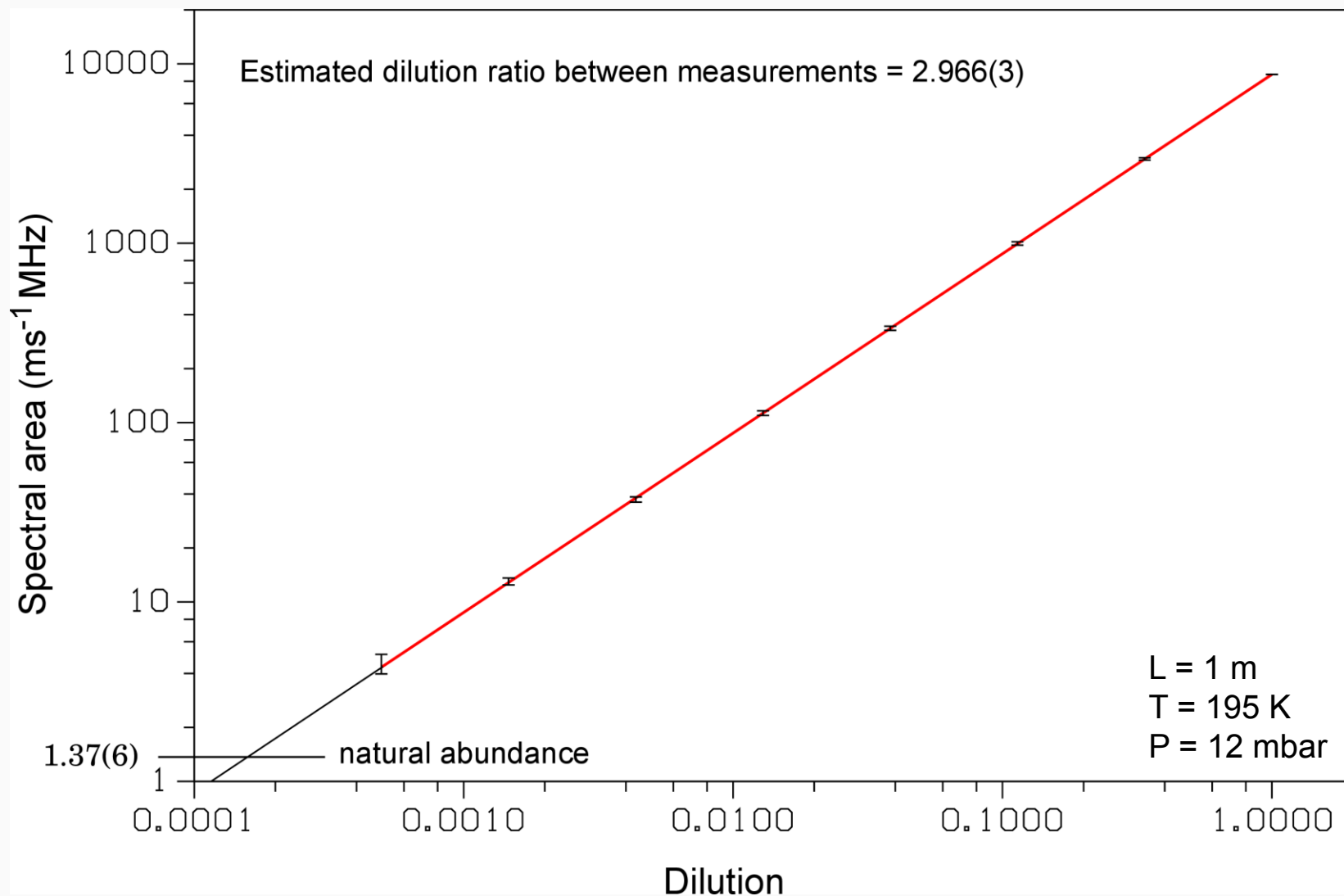


Detection linearity



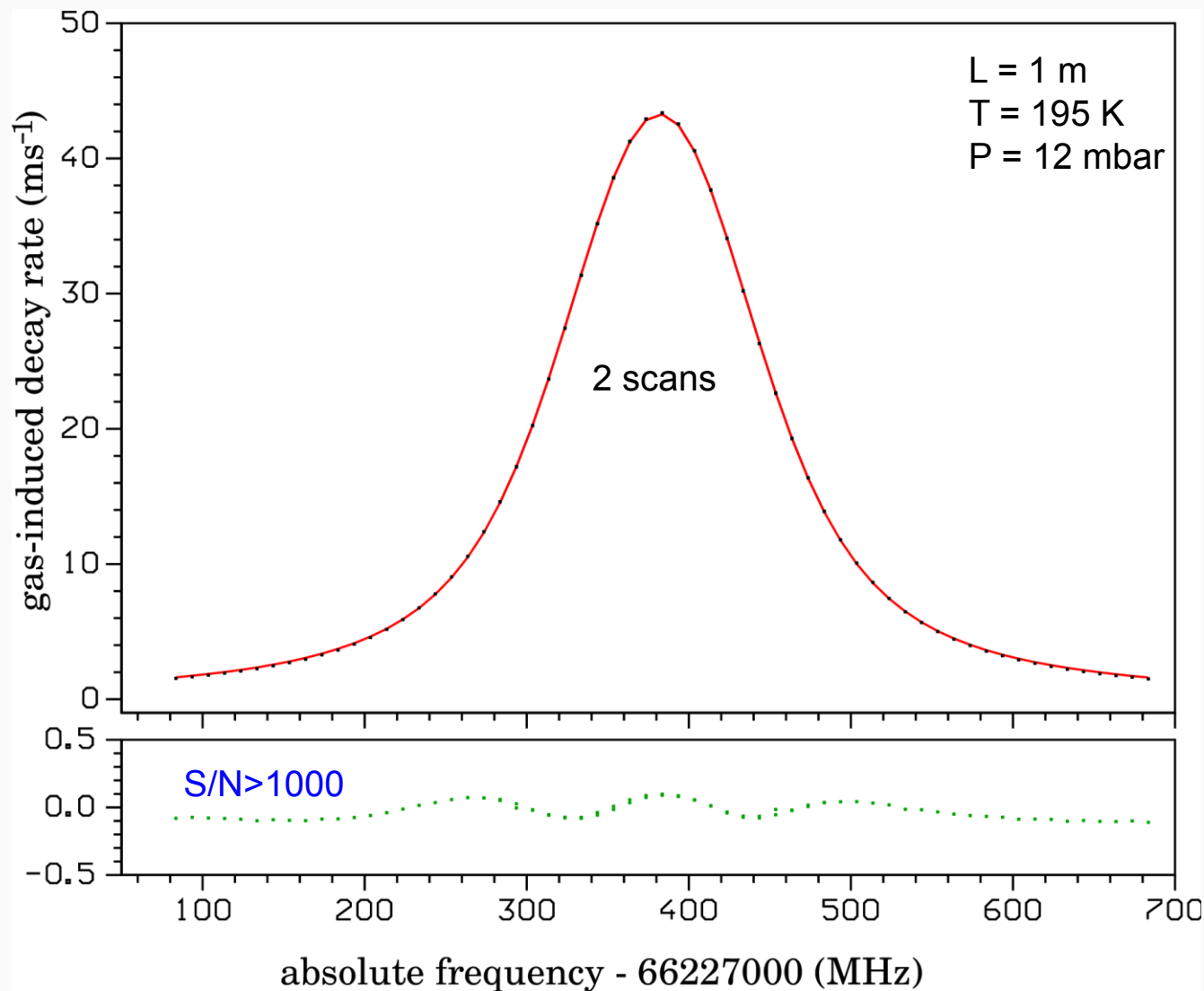


Detection linearity (1-6400 nat. abund.)







Voigt lineshape fitting (6400 nat. abund.)





Preliminary AMS-SCAR intercalibration

	AMS		SCAR	
	pMC	ppq	ms ⁻¹ MHz	ppq
modern sample 	106.68 ± 0.37	1253.3 ± 7.8	1.491 ± 0.026	1248 ± 22
fossil sample 	0.26 ± 0.05	3.1 ± 0.6	-0.004 ± 0.036	-4 ± 30

assuming no uncertainty
on linestrength!

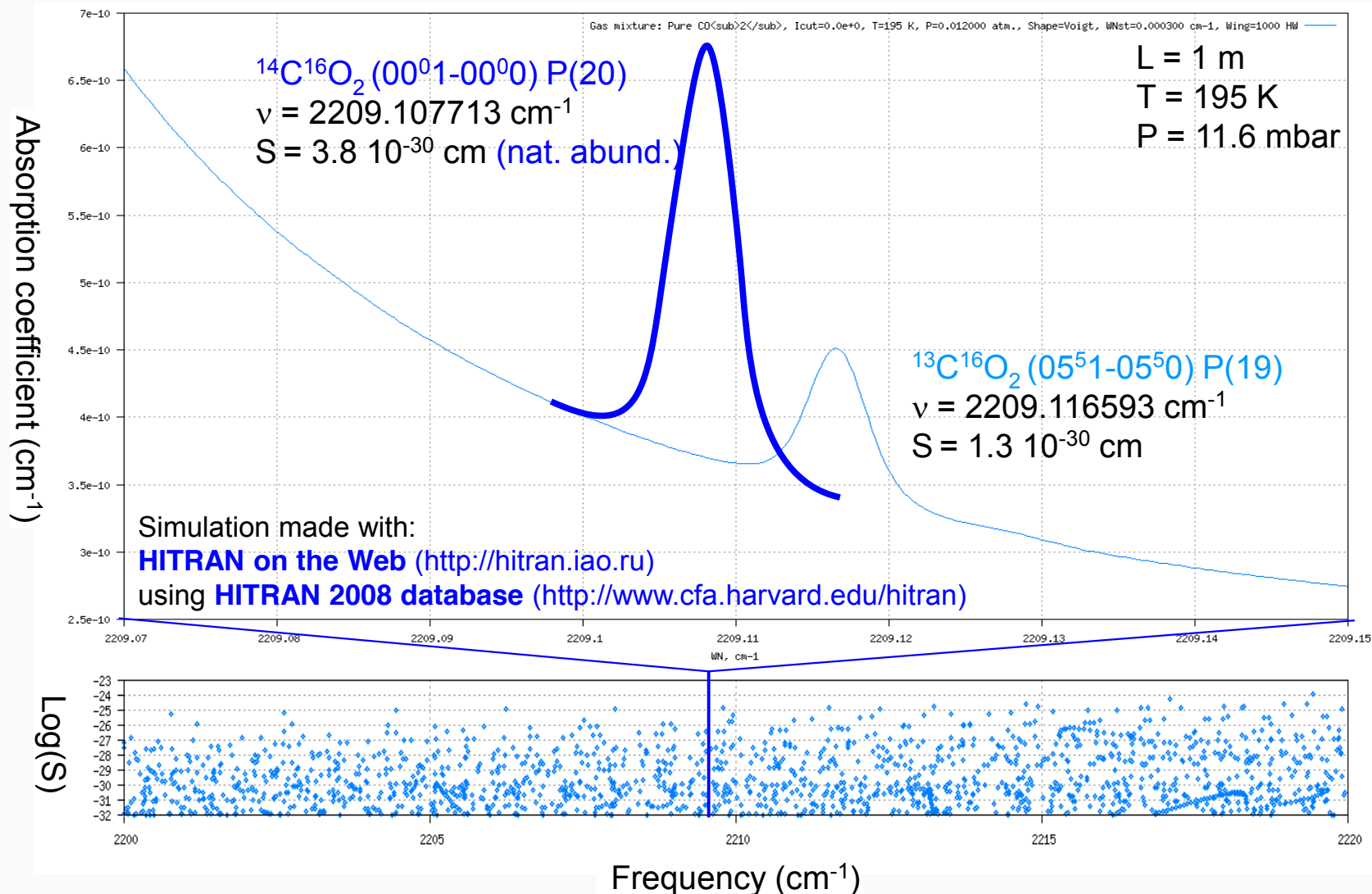


Comparison with competitor techniques

	LSC	AMS	SCAR
^{14}C detection method	β -decay count	^{14}C ion count	absorbed photons
C mass (mg)	~4000	<1	~70
Sample material	C_6H_6	C+Fe	CO_2
Measurement repeatability	restricted to same technique	restricted to same technique	allowed by any technique
Background and/or interferences	cosmic rays	^{13}CH $^{12}\text{CH}_2$	$^{13}\text{CO}_2$ (T>170 K) N_2O (>100 ppt) O_3 (>10 ppm)
Measurement time (h)	~17	~1	~3
Precision for modern samples (pMC)	~0.5	~0.3	~2
Limiting factor	Poisson statistics	Poisson statistics	optical S/N ratio
Footprint (m^2)	~1	~100	~2
Cost (k€)	~120	~3000	~300

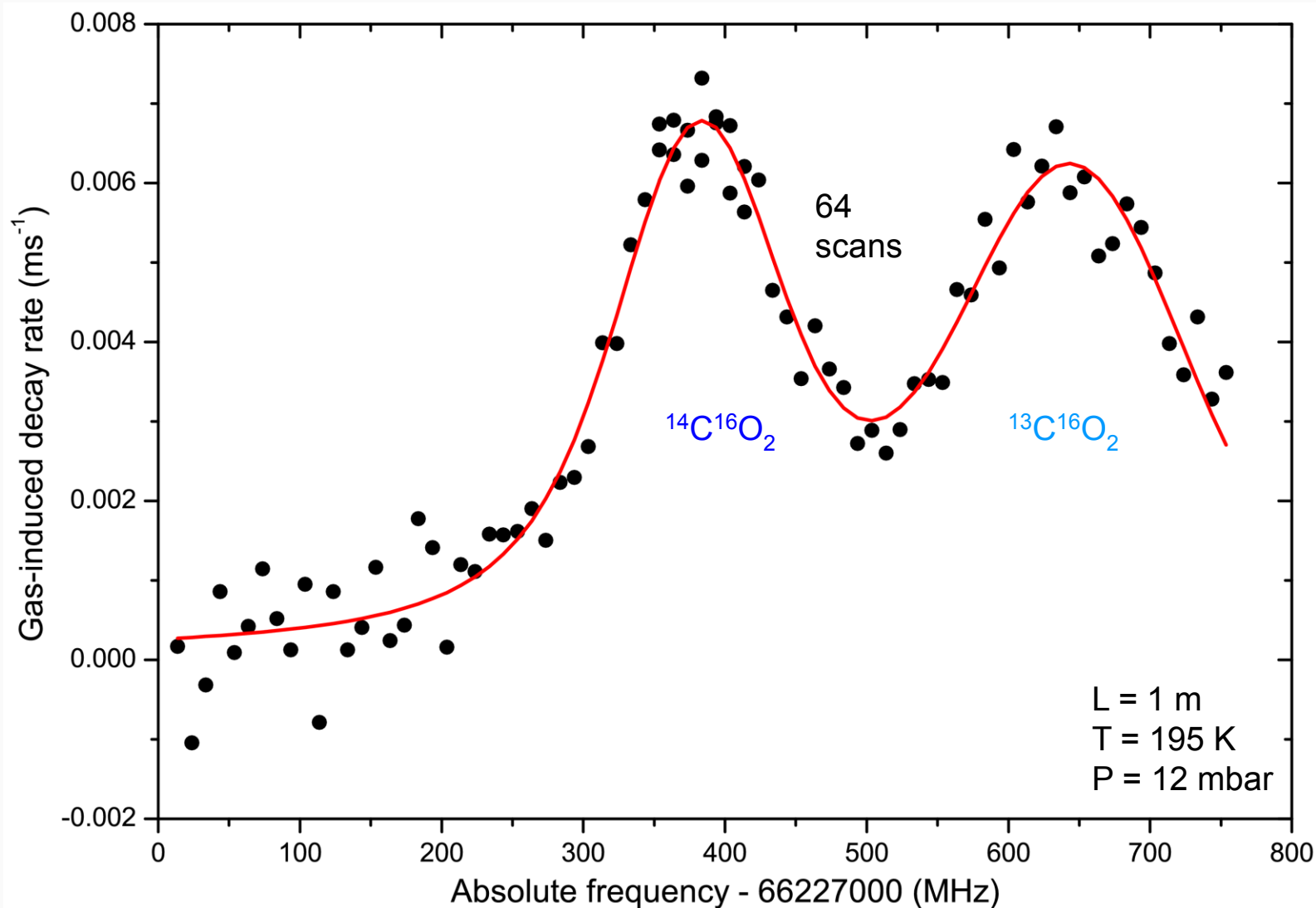


Target line for $^{14}\text{C}^{16}\text{O}_2$ detection



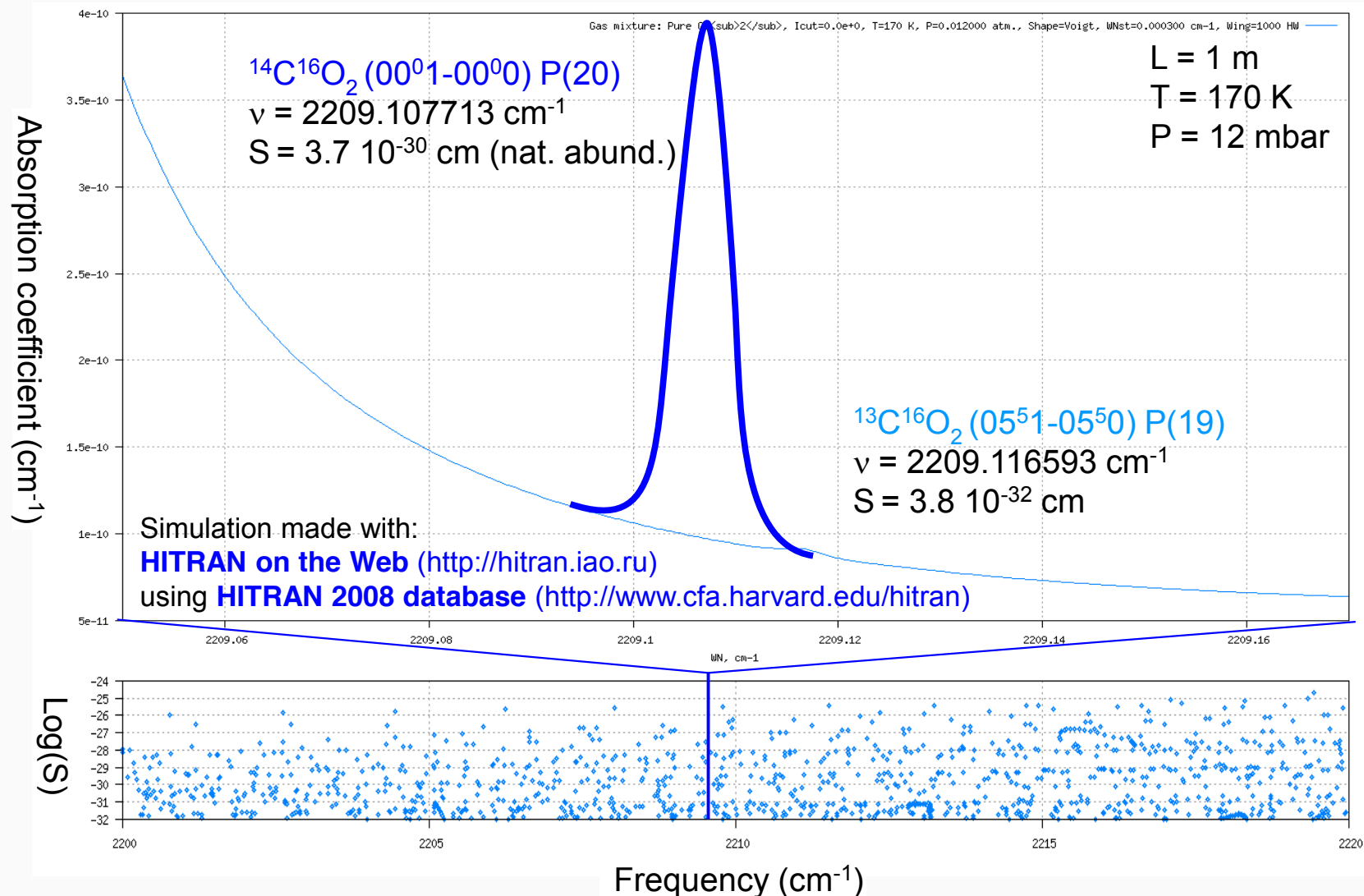


Interference from $^{13}\text{C}^{16}\text{O}_2$ at 195 K





Towards lower temperature





Relative measurement precision

$$\boxed{\frac{P_0}{P}} = \frac{1}{n_s} \frac{P_s}{P} \frac{T}{T_s} \frac{1}{S_{mn}} \int_{-\infty}^{+\infty} \alpha(\nu; P_i, T) d\nu$$

Pressure can be measured/stabilized
to better than **0.1%**

Temperature can be
measured/stabilized
to better than **0.3%**

Linestrength can be stabilized to better than **0.1%**

Spectral area can be measured at **2%**



Molecular Gas Sensing Below Parts Per Trillion: Radiocarbon-Dioxide Optical Detection

I. Galli, S. Bartalini, S. Borri, P. Cancio,* D. Mazzotti, P. De Natale, and G. Giusfredi

Istituto Nazionale di Ottica-CNR (INO-CNR)†, European Laboratory for Non-Linear Spectroscopy (LENS) Via N. Carrara 1,*

search and discovery

physics update

These items, with supplementary material, first appeared at <http://www.physicstoday.org>.

Radiocarbon (^{14}C) co-
absorption cavity ringdown
4.5 μm wavelength. The

attobar pre-
pressure ev-
compactne-
application
of other ve-
cover

Pr

OPN Video

Optics in 2012

POSTED: 30 November 2012

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OPTICS & PHOTONICS NEWS DECEMBER 2012

optics in 2012

SPECTROSCOPIC APPLICATIONS

by using saturated-
) molecules at the
are pushed down to
the lowest
range, the
 ^{14}C -tracing
e detection
R spectral

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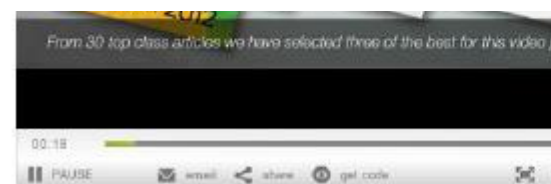
Science NOW

UP TO THE MINUTE NEWS FROM SCIENCE

ScienceShot: Sniffing Out the One in a Quadrillion

by Jon Cartwright on 9 December 2011, 11:41 AM | [2 Comments](#)

For
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This special issue of *Optics & Photonics News* highlights the research to have emerged over the past 12 months. Our panel of editors reviewed close to 80 submitted summaries from scientists all over the globe. They selected for publication the 30 stories that they felt most clearly communicated breakthroughs of interest to the optics community.

Science


AAAS

3.4.111

http://www.osa-opn.org/home/multimedia/optics_in_2012/

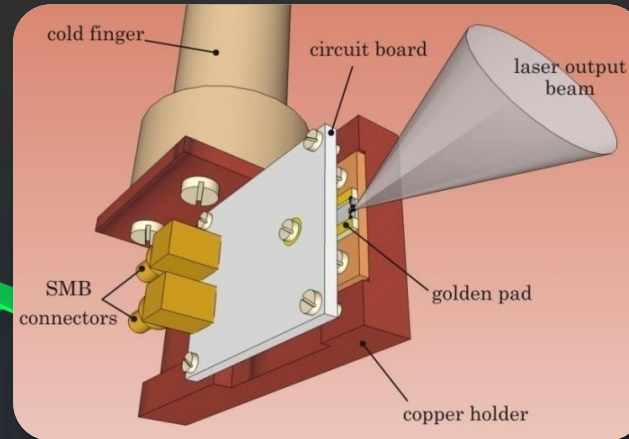
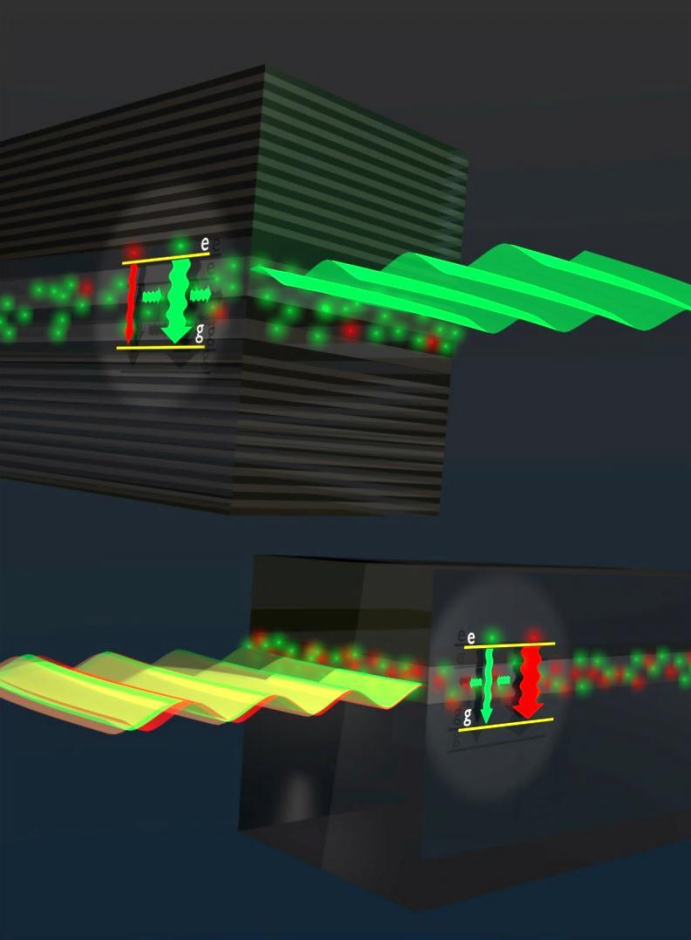


Perspectives: work in progress

	Present setup	 Future setup
Cell material	stainless steel	quartz
Volume (L)	~8	~0.7
Optical finesse	~11000	~15000
C mass (mg)	~70	~6
Laser source	intra-cavity DFG	tunable CW DFB-QCL1
Absolute freq. reference	OFCS referenced to quartz/Rb/GPS clock	$^{14}\text{C}^{16}\text{O}_2$ -locked/narrowed CW DFB-QCL2
Temperature (K)	195	170
Cooling method	dry ice	closed-cycle Stirling cryocooler
Footprint (m²)	~2	~1
Cost (k€)	~300	<100

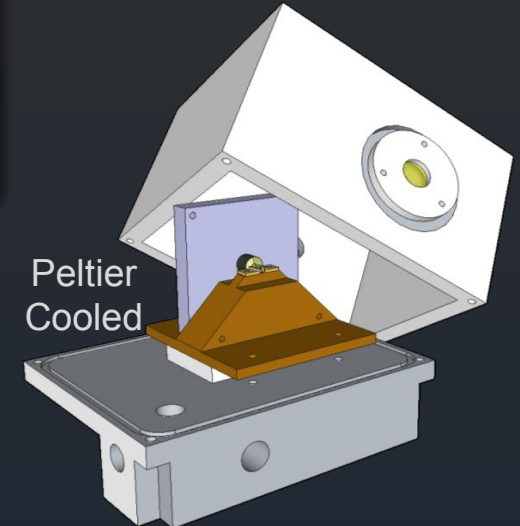


Quantum Cascade Lasers



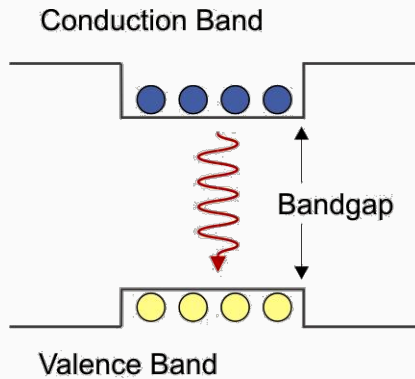
QCLs are emerging as the main sources in the mid and far infrared:

- ✓ Compactness
- ✓ Custom wavelength
- ✓ High power output
- ✓ Single mode emission

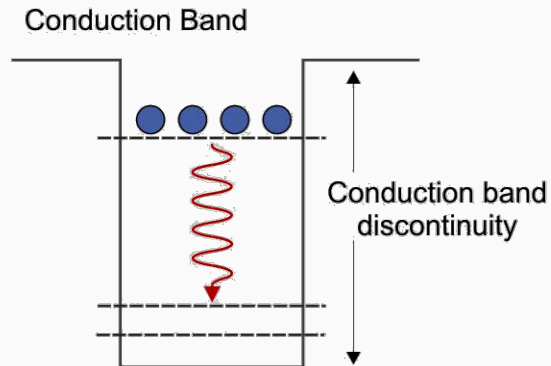




QCL vs. NIR Laser diodes



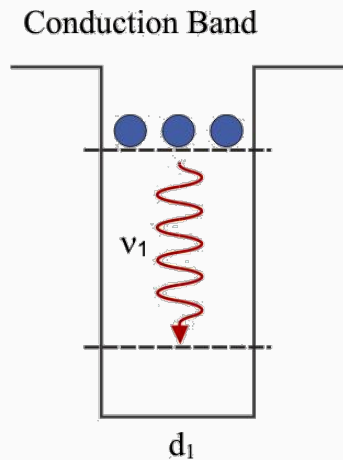
a)



b)

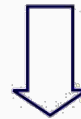
Diode laser: emission for electron-hole recombination (a).

QCL: only conduction band involved (b).

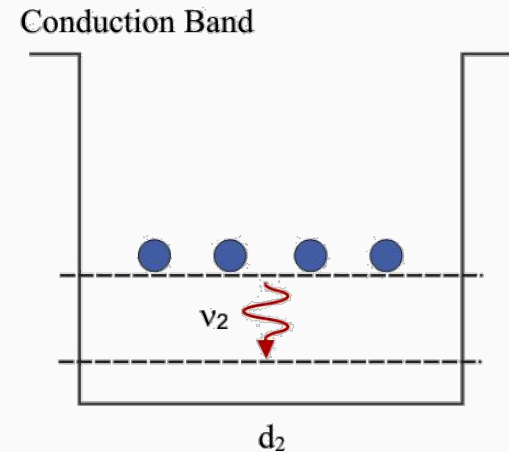


d_1

Same material system
Different quantum well widths



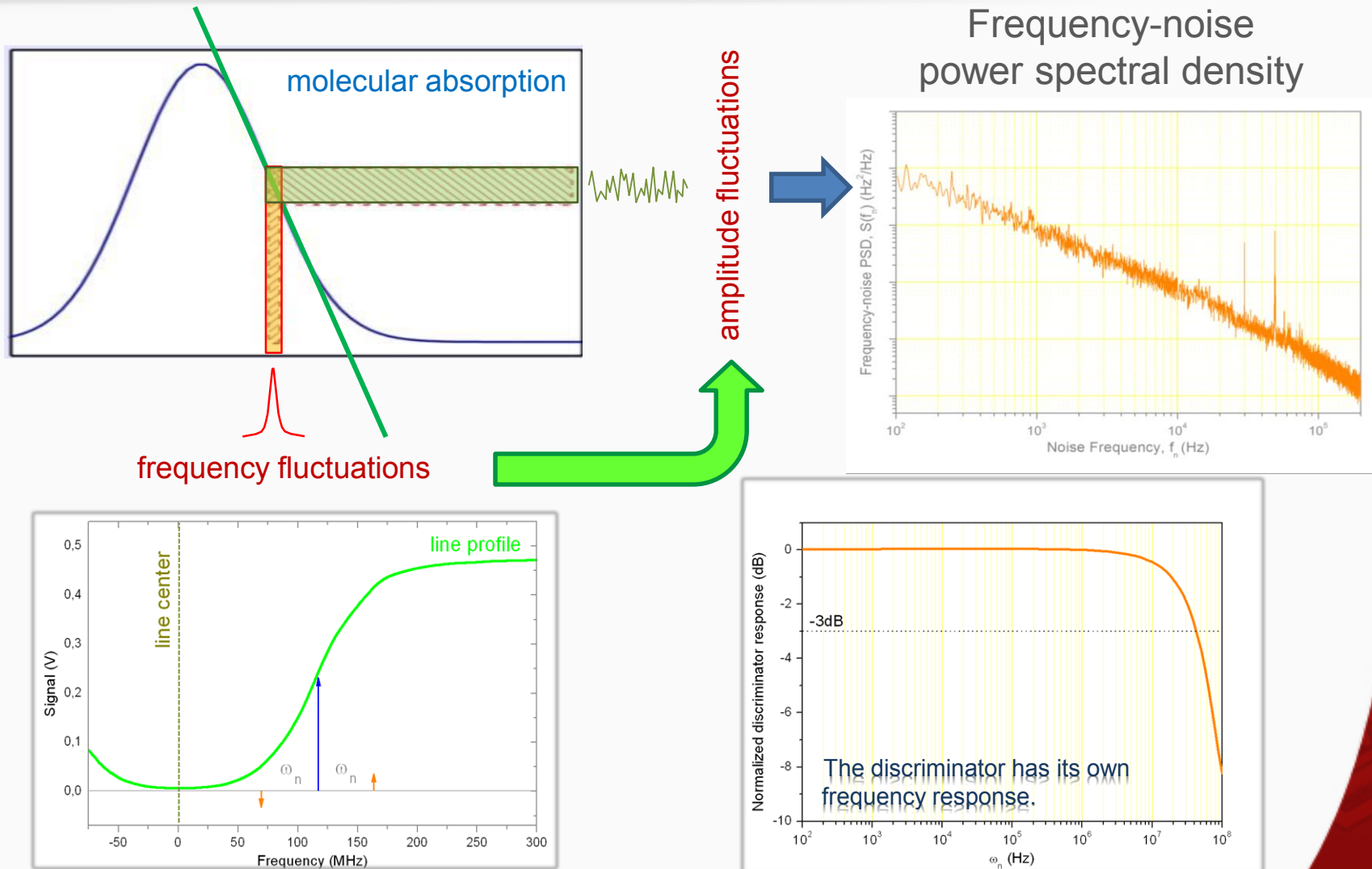
Different transition energies
 $d_1 < d_2 \rightarrow v_1 > v_2$



d_2

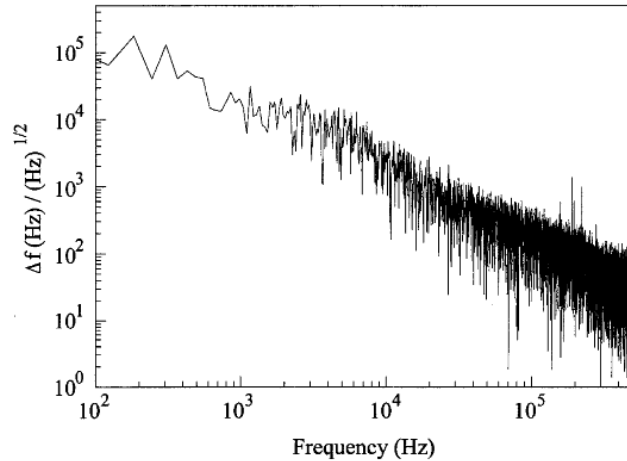


Measuring the QCL's frequency-noise: principle

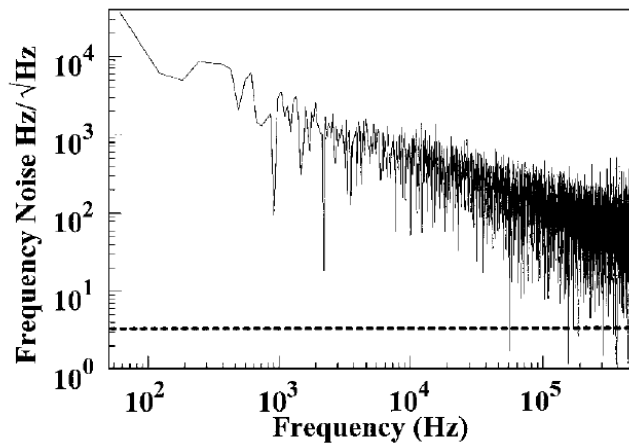




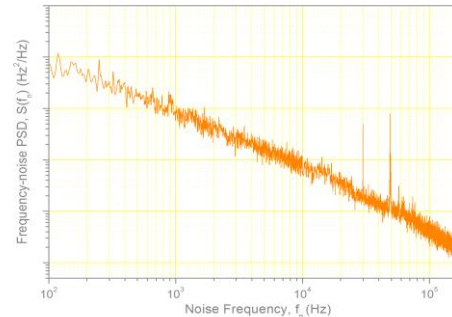
Observing QCL Intrinsic Linewidth



Williams *et al.*, *Opt. Lett.* **24**, 1844 (1999)



Myers *et al.*, *Opt. Lett.* **27**, 170 (2002)



“Observing the Intrinsic Linewidth of a Quantum-Cascade Laser: Beyond the Schawlow-Townes Limit”,

S. Bartalini *et al.*, *Phys. Rev. Lett.* **104**, 083904 (2010)

“The intrinsic limits of Quantum-Cascade Lasers”

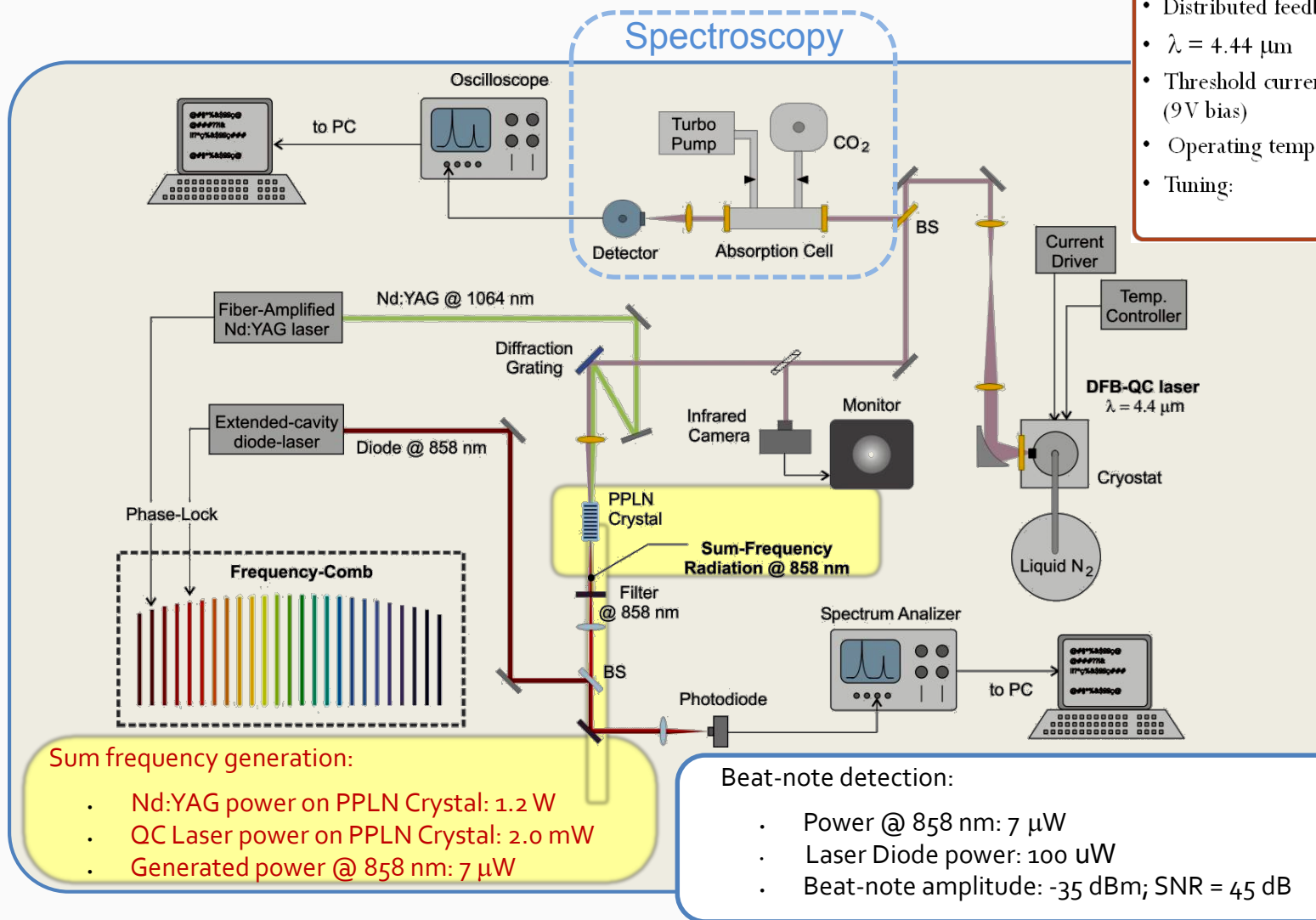
Physics Today, p.20, (May 2010)



Basic QCL-Comb set-up

QCL specs:

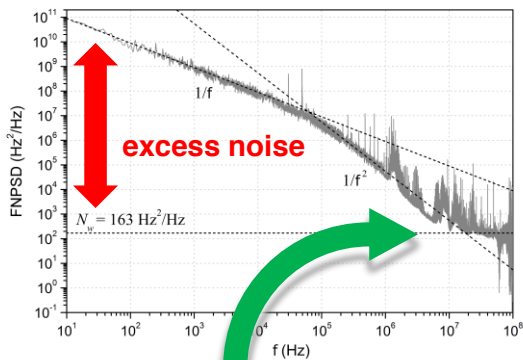
- Distributed feedback
- $\lambda = 4.44 \mu\text{m}$
- Threshold current: 240 mA (9V bias)
- Operating temperature: 84 K
- Tuning:
 $\sim 840 \text{ MHz}/\mu\text{A}$
 $\sim 2 \text{ GHz/K}$



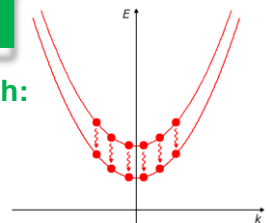


QCLs: noise suppression

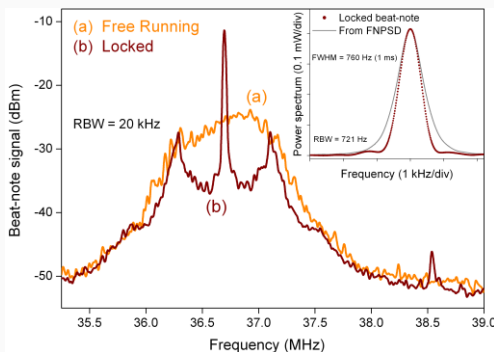
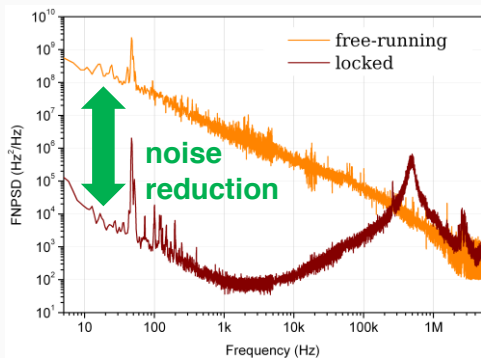
Bartalini et al., PRL 104, 083904 (2010)



Intrinsic linewidth:
inter-subband
transition limit



The results



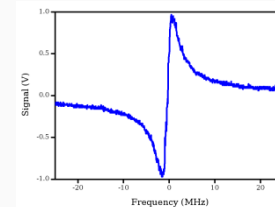
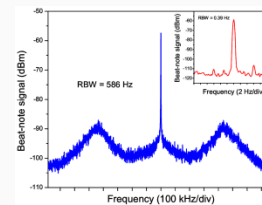
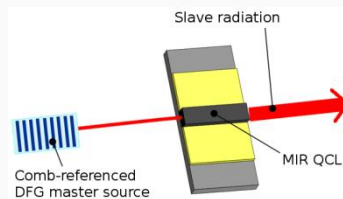
Three techniques:

Opt. Lett. **37**, 1011, (2012)

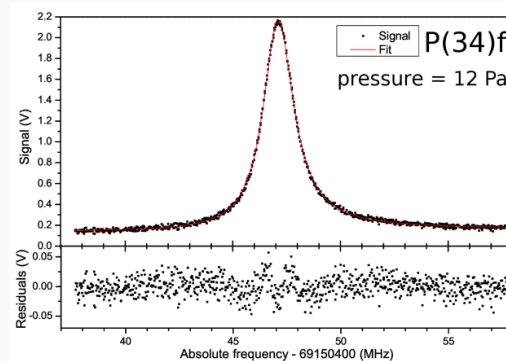
Opt. Expr. **37**, 4811 (2012)

	Optical injection locking	Phase lock	Frequency lock
Reference	Comb-referenced DFG source	Comb-referenced DFG source	Molecular absorption line
Linewidth	48 kHz	500 Hz	760 Hz

Appl. Phys. Lett. **102**, 121117 (2013)



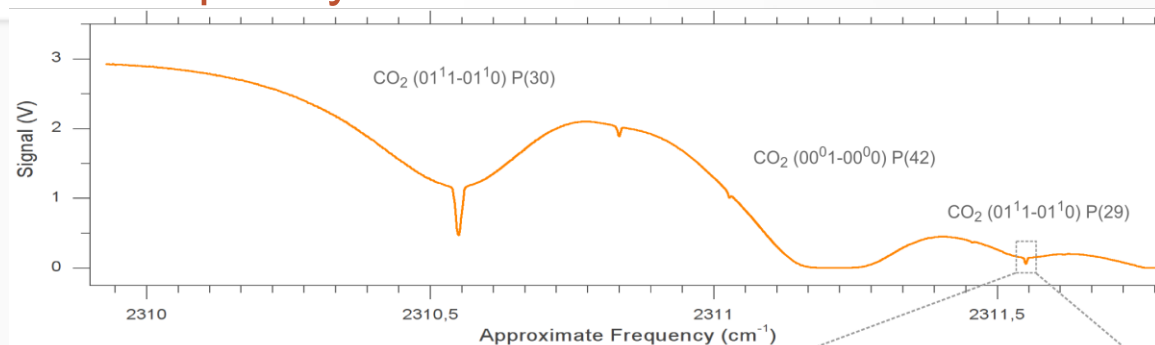
CO2 high resolution spectroscopy



Molecular Physics, **111**, 2041 (2013)



Accuracy improvement of absolute frequency measurements



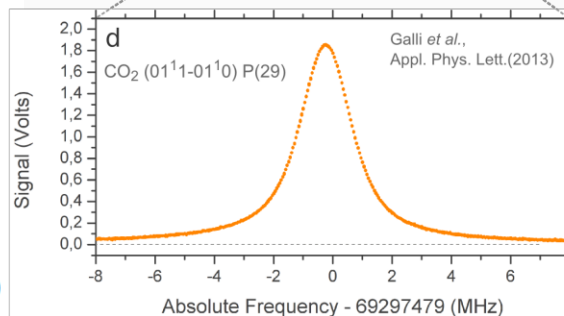
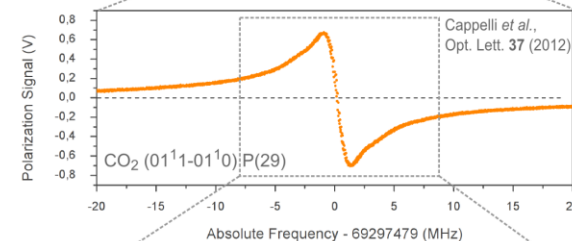
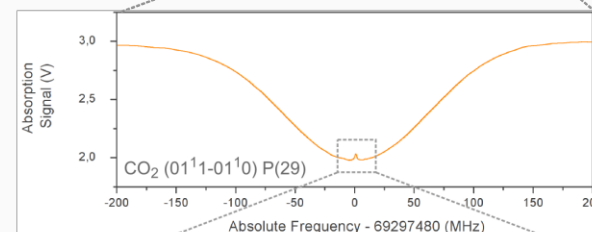
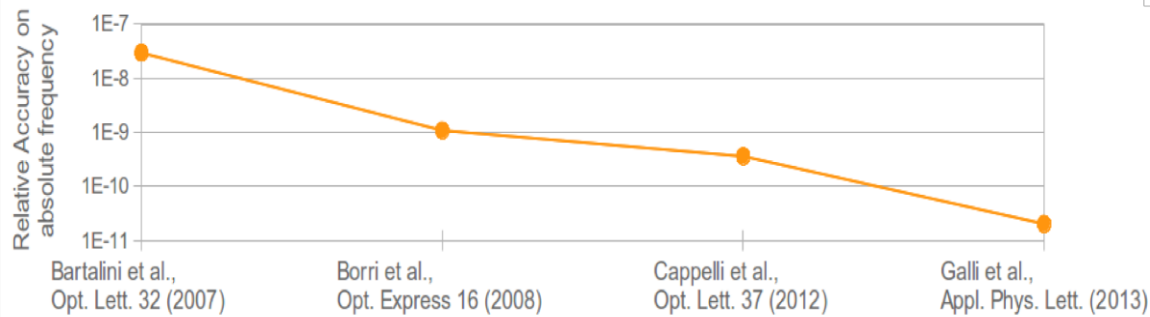
From

Doppler-limited direct-absorption spectroscopy
with a **free running** QCL

To

Sub-Doppler saturated-absorption spectroscopy
with a **comb-assisted** QCL

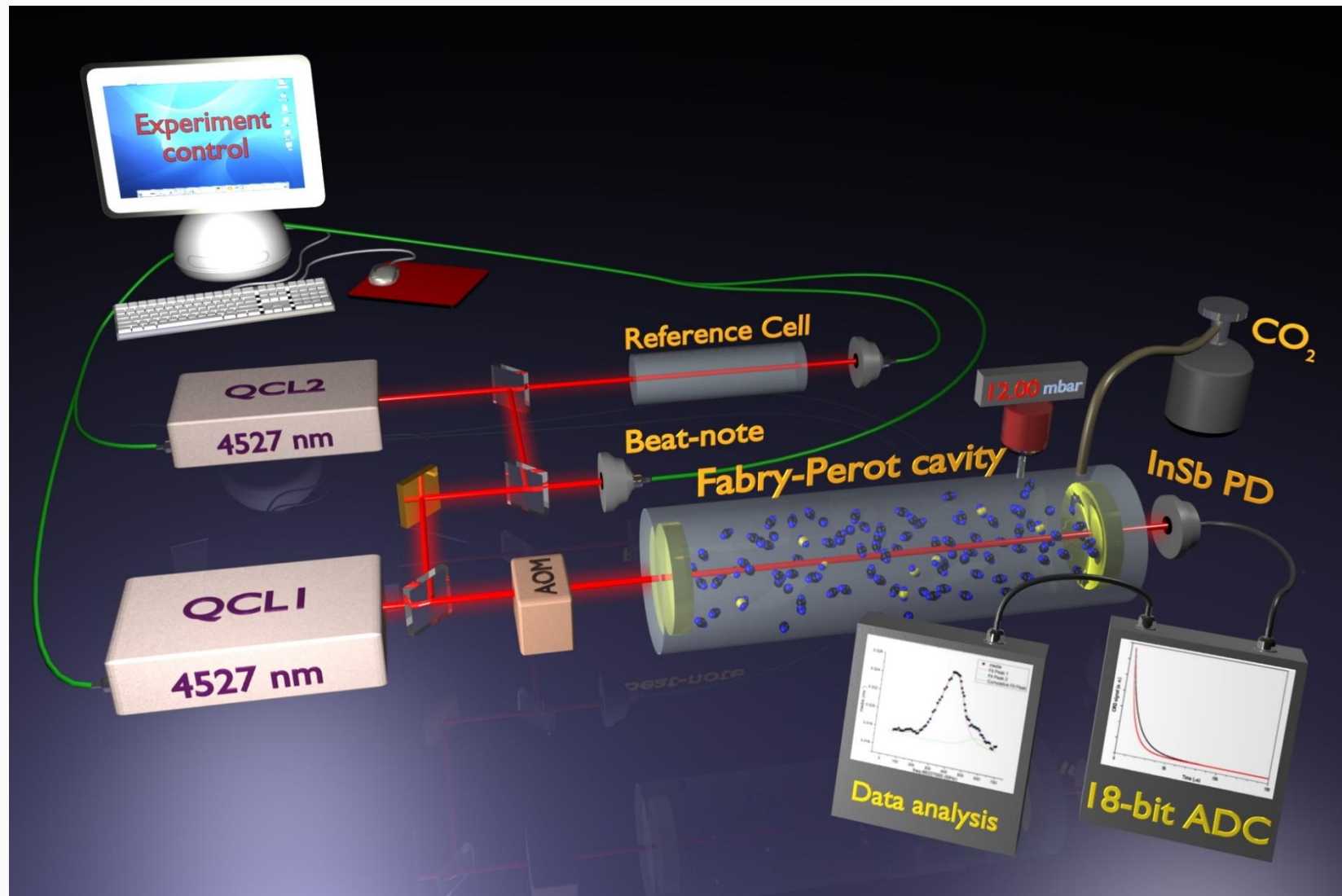
accuracy improved by
more than **3 orders of magnitude**





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SCAR2 setup

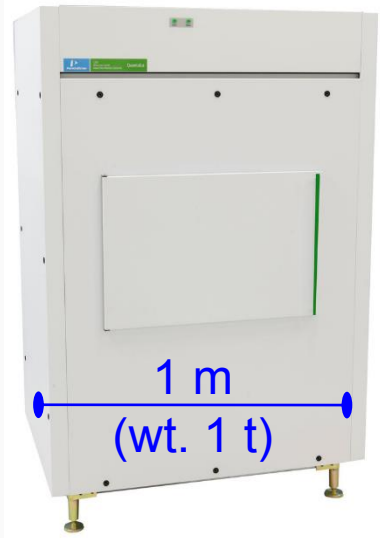




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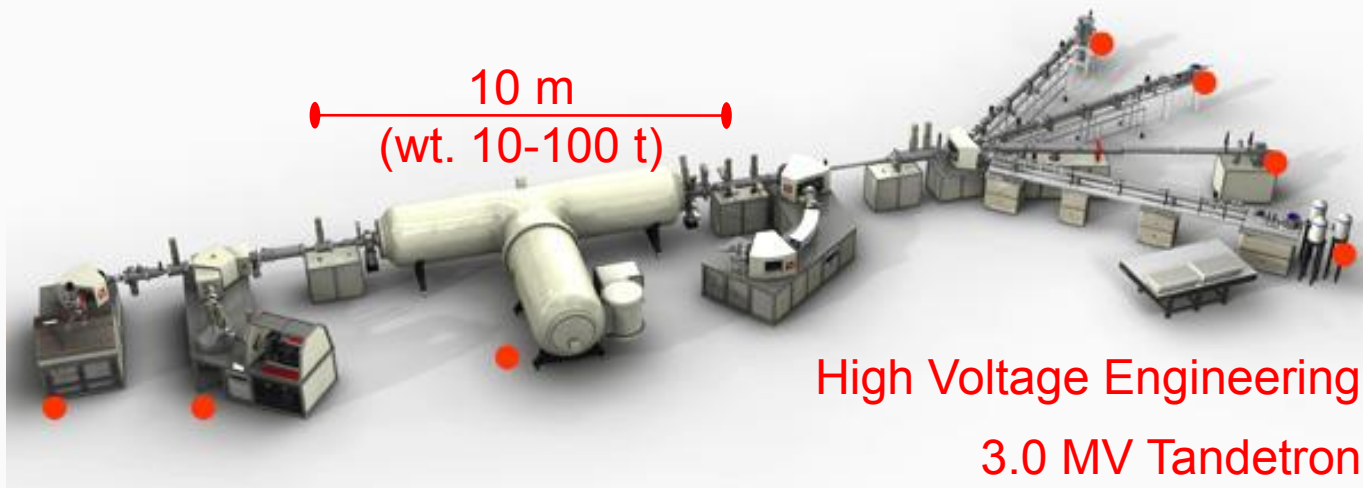
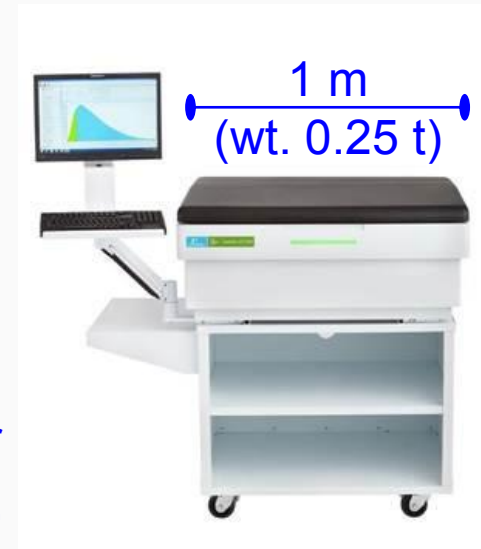
Examples of commercial apparatuses

LSC



PerkinElmer
Quantulus 1220

PerkinElmer
Quantulus GCT 6220



AMS

High Voltage Engineering
3.0 MV Tandatron




Saturated-absorption cavity ring-down (SCAR)

- ✓ Time-separated measurements in differential schemes (full/empty cavity, on/off-resonant gas) have **intrinsic technical fluctuations**.
- ✓ Recording CRD spectra in the **saturation-absorption regime** has a double advantage:
 - **sub-Doppler resolution** for free;
 - **gas absorption** can be **discriminated** against cavity losses in a **single ring-down event**.
- ✓ Due to a varying saturation parameter during the ring-down process, the decay curves are **not truly exponential**.
- ✓ An **accurate theoretical model** must be adopted to fit the experimental data, taking into account both the **transverse beam shape** and all homogeneous/inhomogeneous **broadening mechanisms**.

➡ *J. Opt. Soc. Am. B* **32**, 2223
(2015)



Technical upgrades

	SCAR1		SCAR2
Cell material	stainless steel		quartz
Volume (L)	~8		~0.7
Optical finesse	~11 000		~18 000
C mass (mg)	~70		~6
Laser source	intra-cavity DFG		tunable CW DFB-QCL1
Absolute freq. reference	OFCS referenced to quartz/Rb/GPS clock		N ₂ O-locked CW DFB-QCL2
Temperature (K)	195		170
Cooling method	dry ice		acoustic-Stirling cryocooler
Footprint (m²)	~2		~1
Cost (k€)	~300		~150



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Upcoming results for SCAR2

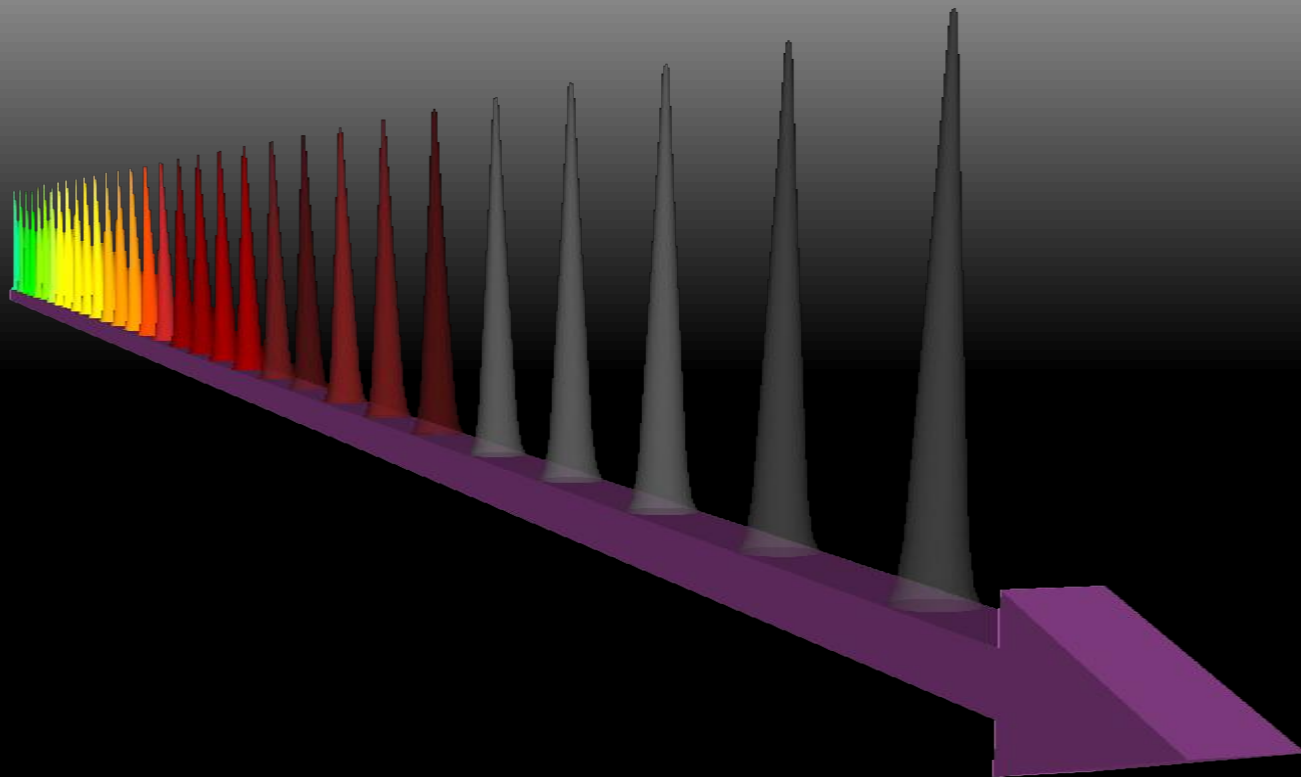
New results, demonstrating **few ppq** sensitivity of SCAR spectroscopy, will be published in:

Galli et al. Optica, (2016)



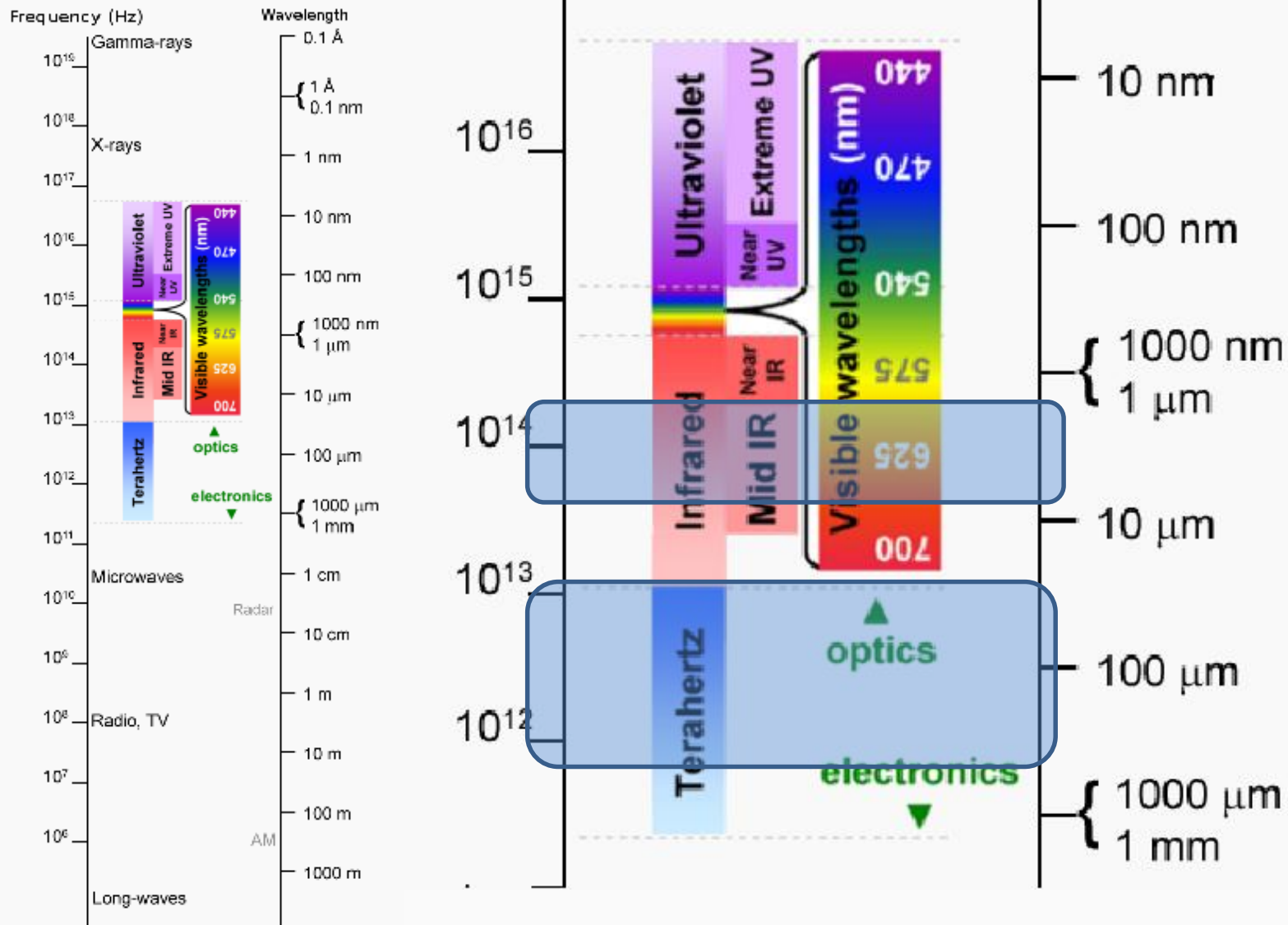
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OTTICA

Moving deeper into the (Far) Infrared





Electromagnetic Spectrum

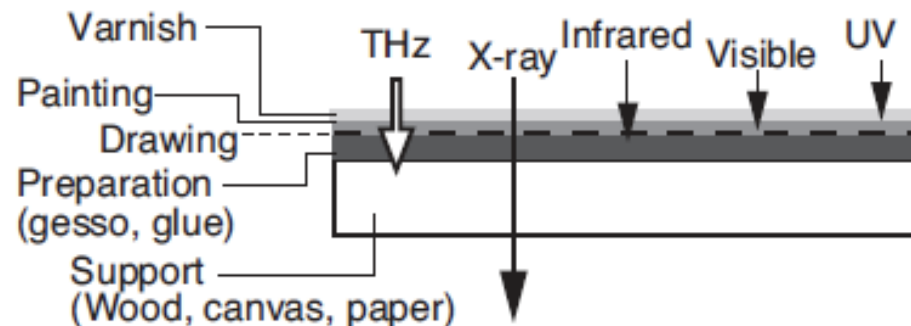


Properties & advantages

- Non-invasive, non-ionizing, non-contact and non-destructive examination method
- Can penetrate opaque materials
- Complementary to traditional methods such as X-ray, UV, VIS and IR imaging/spectroscopy
- Radiation frequency can tune the depth of penetration

State of the art

- Rapid instrumental progress and development in last decades
- No commercial spectral libraries in spite of a range of commercial instruments
- Currently a database with 500 spectra of art materials (pigments, dyes, binders)

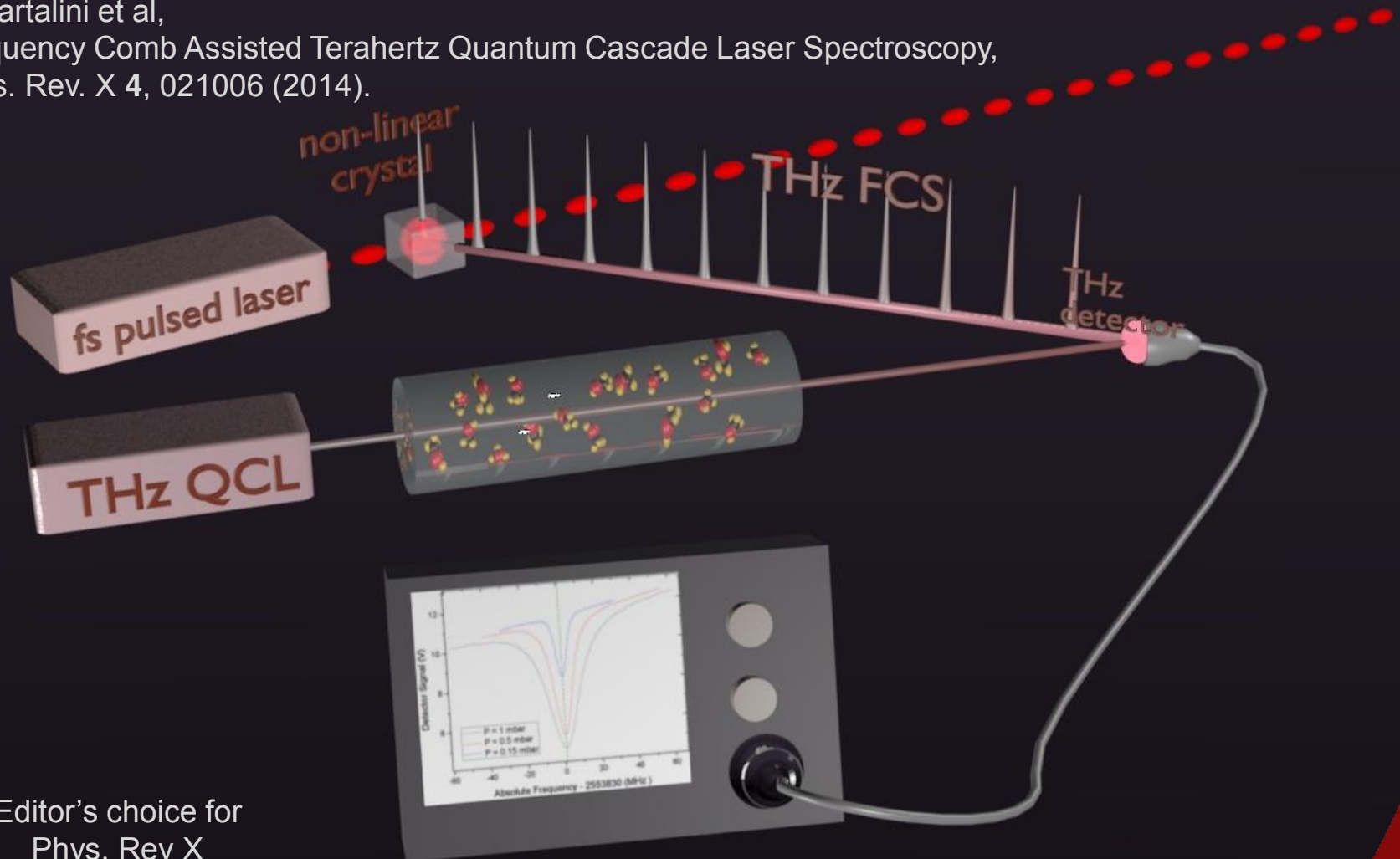




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Comb-assisted QCL-based THz Spectroscopy

S. Bartalini et al,
Frequency Comb Assisted Terahertz Quantum Cascade Laser Spectroscopy,
Phys. Rev. X 4, 021006 (2014).

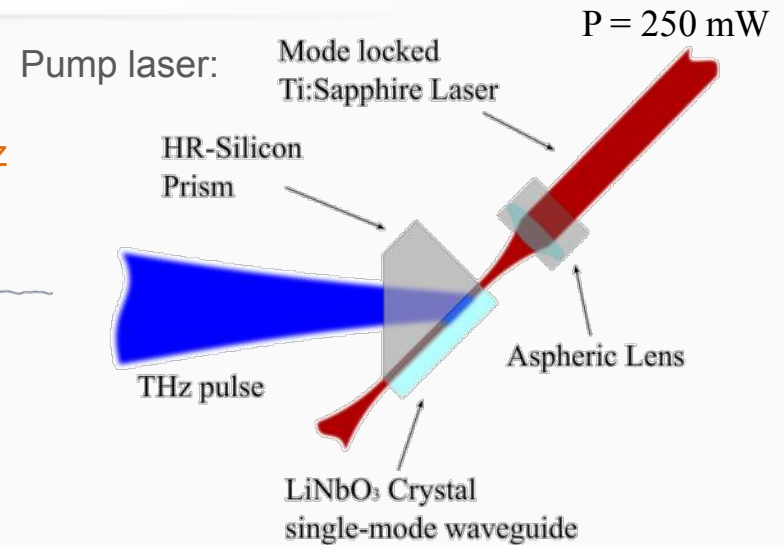
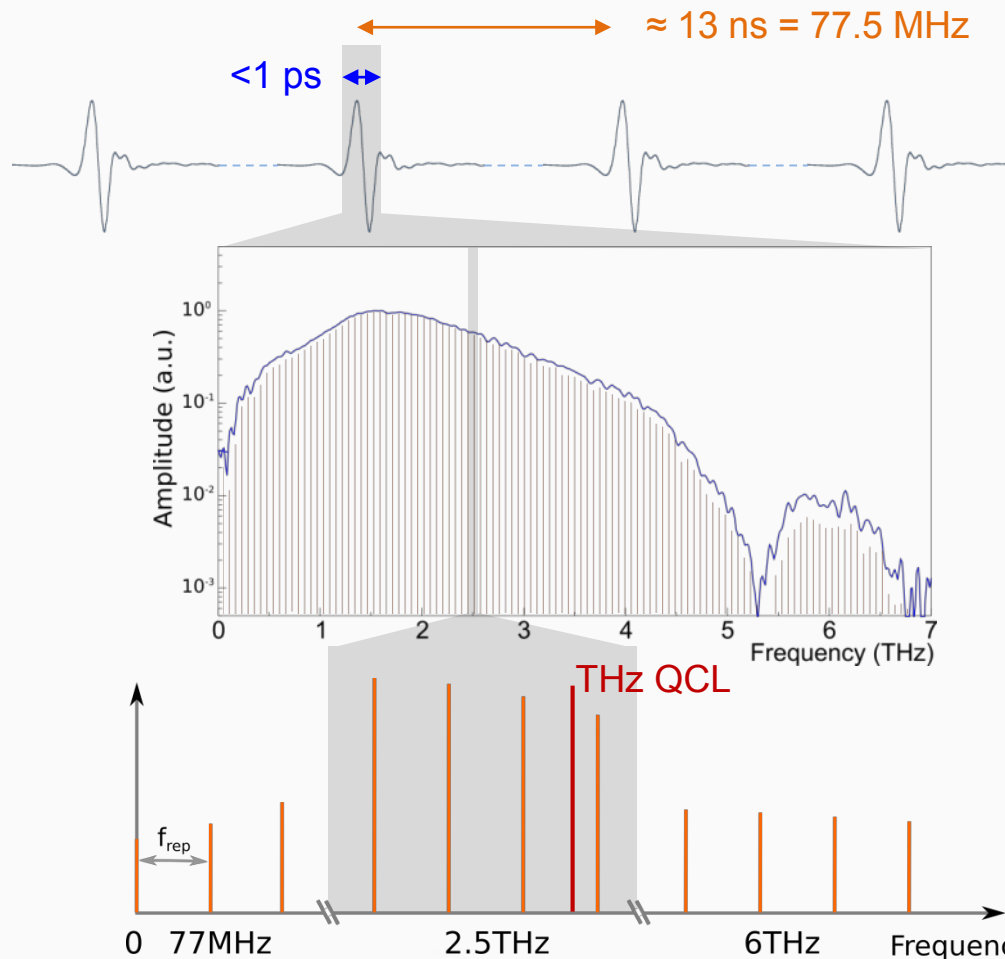


Editor's choice for
Phys. Rev X
Highlights



THz pulsed radiation: THz comb

Usually employed for THz TDS, but...

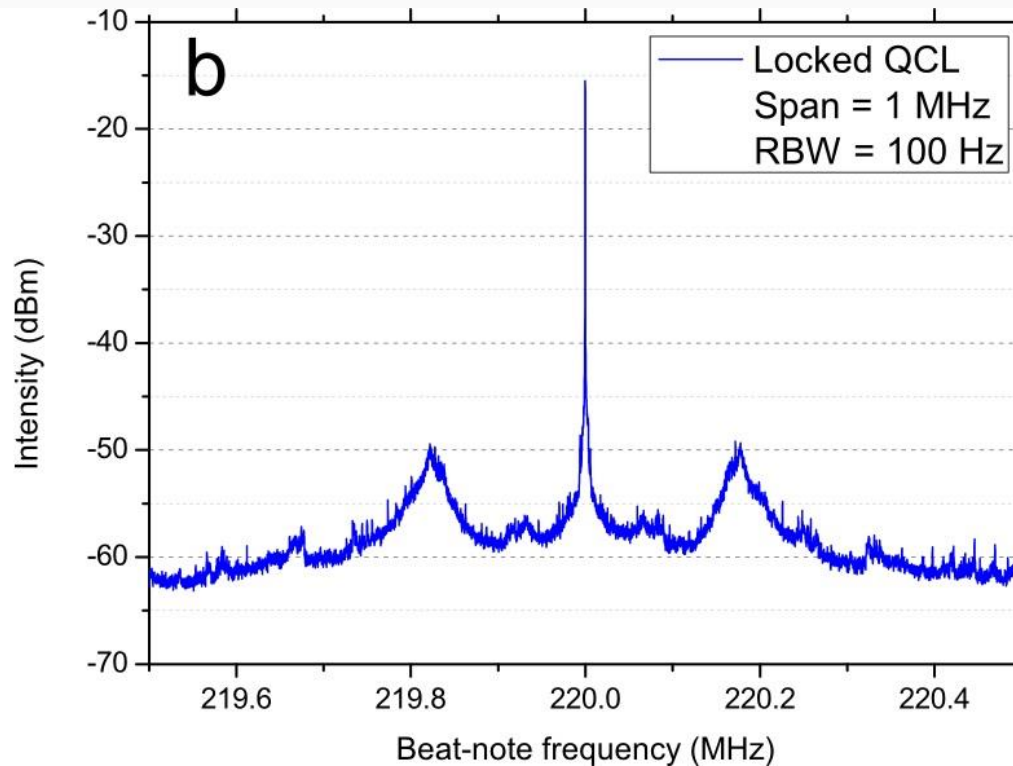


Total power = 1 μ W
Power per tooth = 30 pW

zero-CEO frequency
6-octaves-spanning COMB
(more than 60000 teeth)



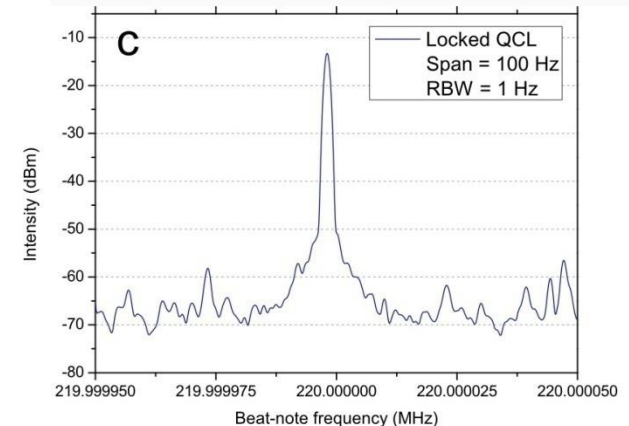
Phase-lock of a THz QCL to a THz COMB



About 70-80% of the QCL power
effectively phase-locked

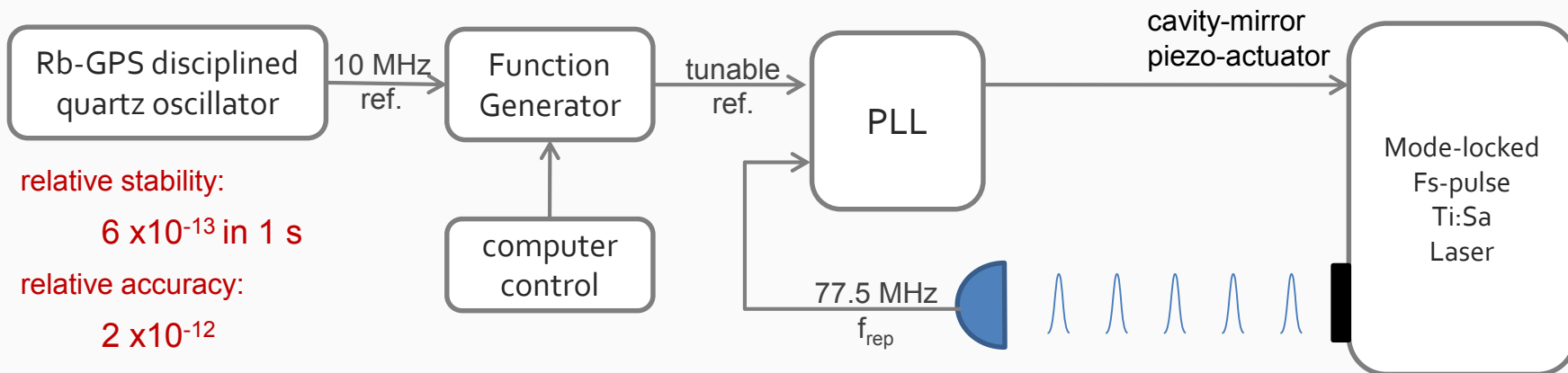
S/N = 50 dBm (1 Hz – RBW)

Electronic Bandwidth = 200 kHz





Stabilization of the COMB repetition rate

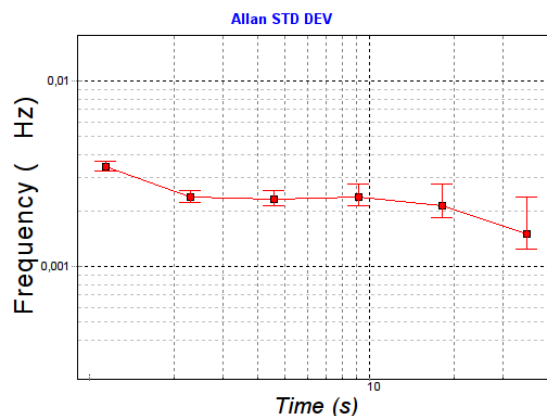
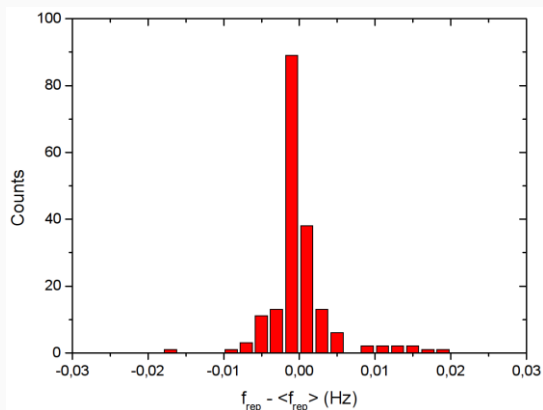


relative stability:

$$6 \times 10^{-13} \text{ in } 1 \text{ s}$$

relative accuracy:

$$2 \times 10^{-12}$$



For the THz comb:

- offset instabilities cancel out
- only f_{rep} instabilities contribute to the tooth linewidth

Allan variance < 4mHz



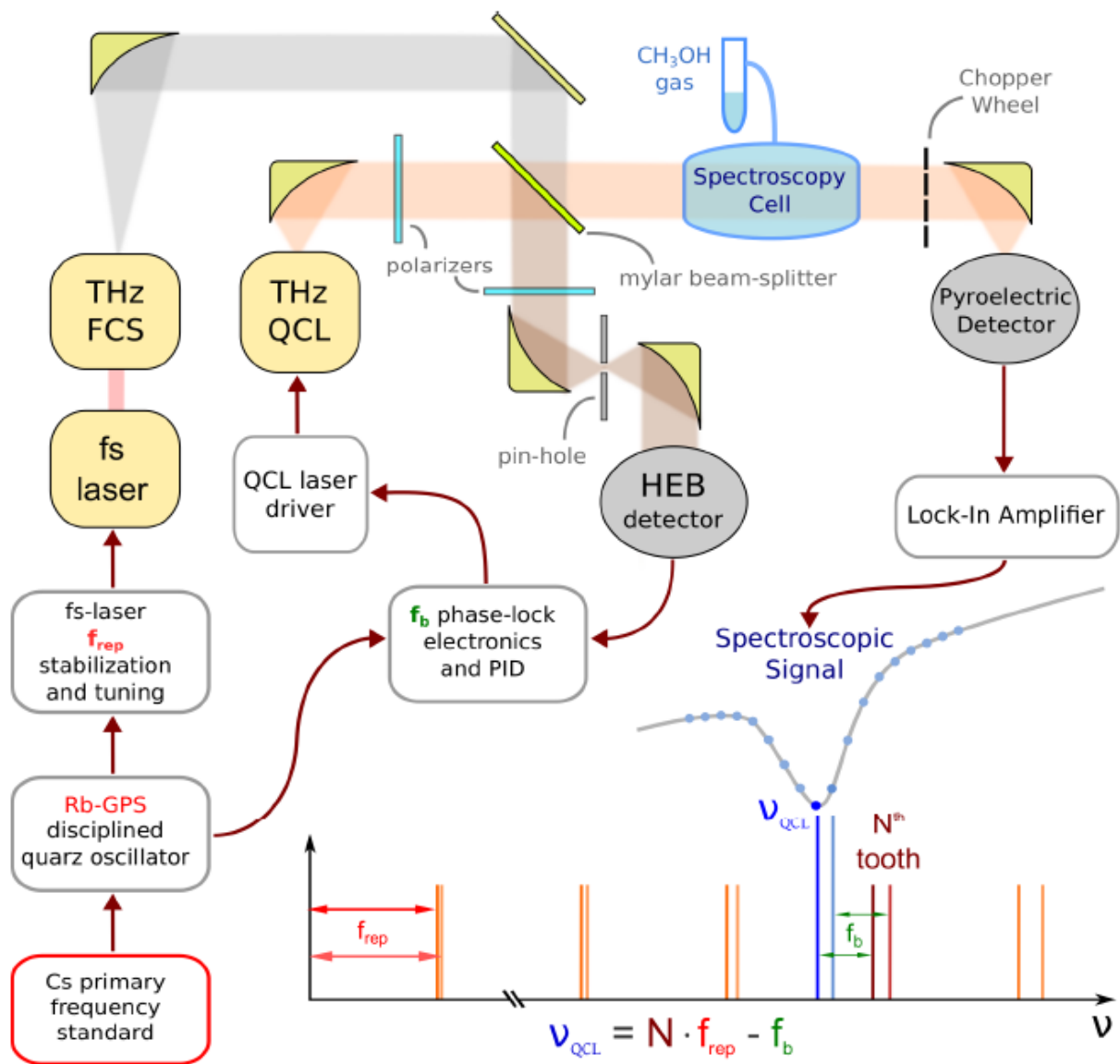
Measured f_{rep} relative stability better than 5×10^{-11}



~100Hz level @ 2.5 THz



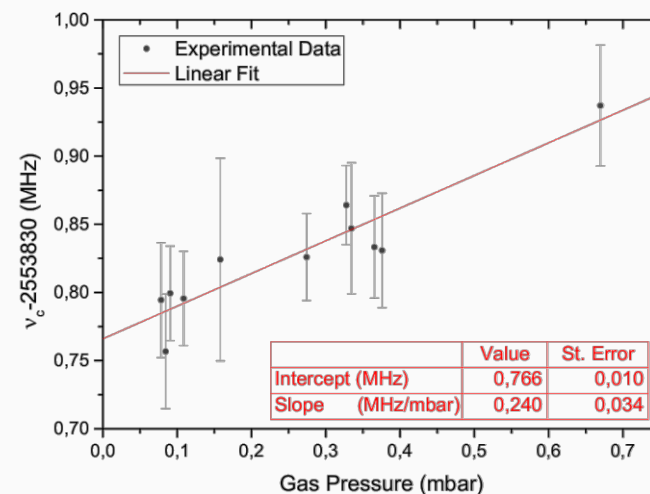
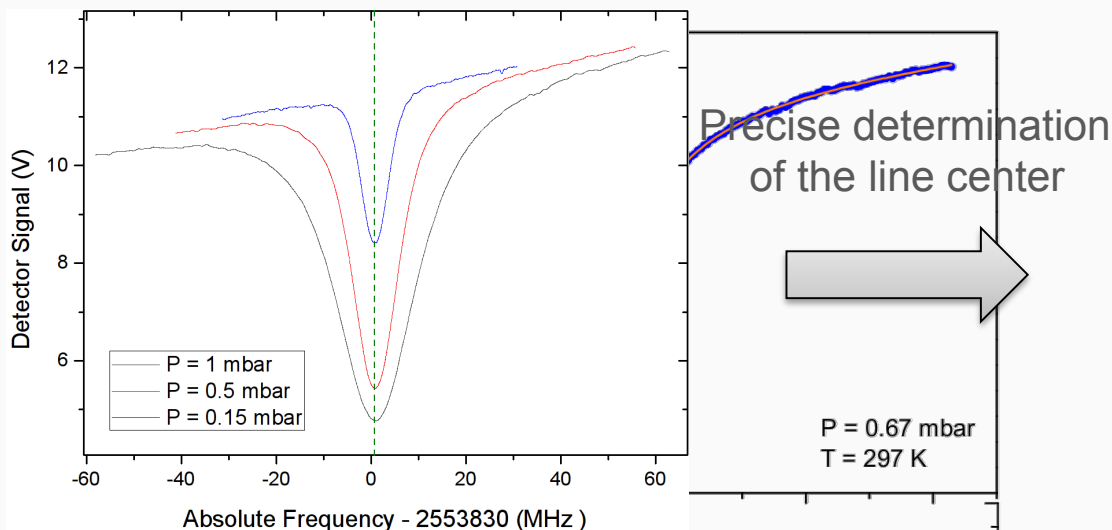
Spectroscopic Setup





THz-comb-assisted Spectroscopy

$$\nu_0 = 2.553\,830\,766\,(10)\,\text{THz}$$



Comparison with state-of-the-art alternative techniques:

Transition linewidth \sim MHz (FWHM)

S/N > 100



expected
error on ν_c :
 ~ 10 kHz

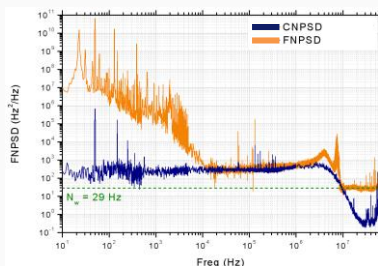
Doppler limited spectroscopy



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QCL based metrology

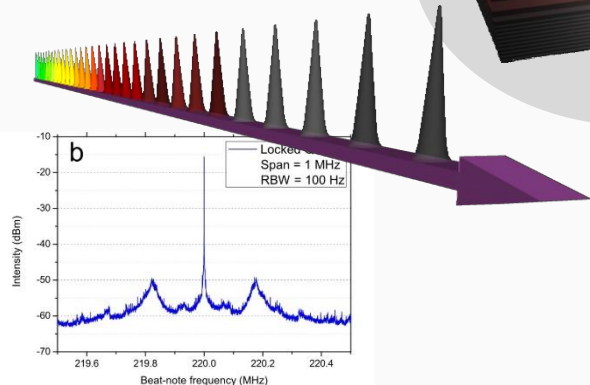
SOURCE: QCL



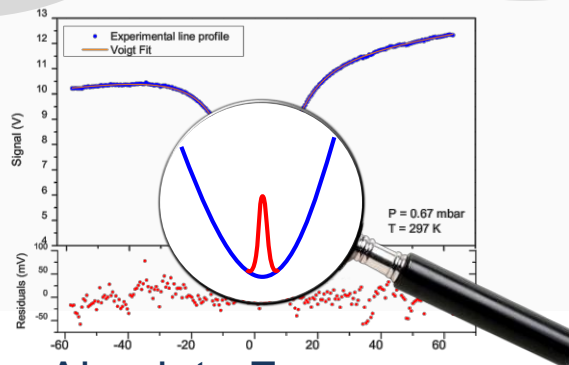
M.S. Vitiello *et al.*
Nat. Photon. **6**, 525 (2012).

THz QCLs are suitable for metrological-grade experiments

$$\Delta\nu = \pi N_W \sim 90 \text{ Hz}$$



L. Consolino *et al.*,
Nat. Comm. **3**, 1040 (2012).

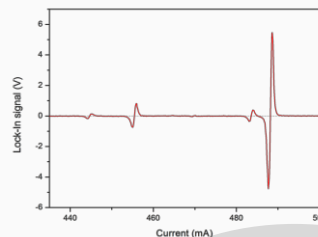


Absolute Frequency

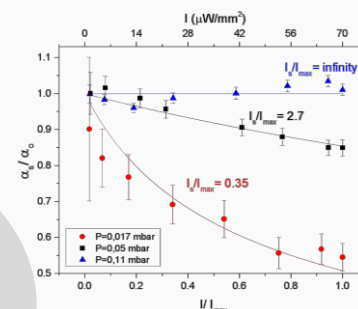
S. Bartalini *et al*, *Phys. Rev. X* **4**, 021006 (2014).

SPECTROSCOPY

High-sensitivity techniques

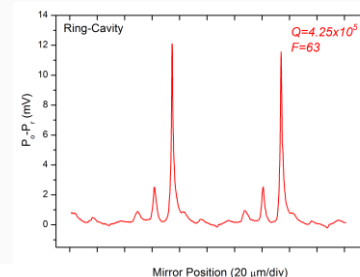


L. Consolino *et al.*,
Sensors **13**, 3331 (2013). **Saturation**



L. Consolino *et al.*
Appl. Phys. Lett.
106, 021108 (2015)

Cavity enhancement



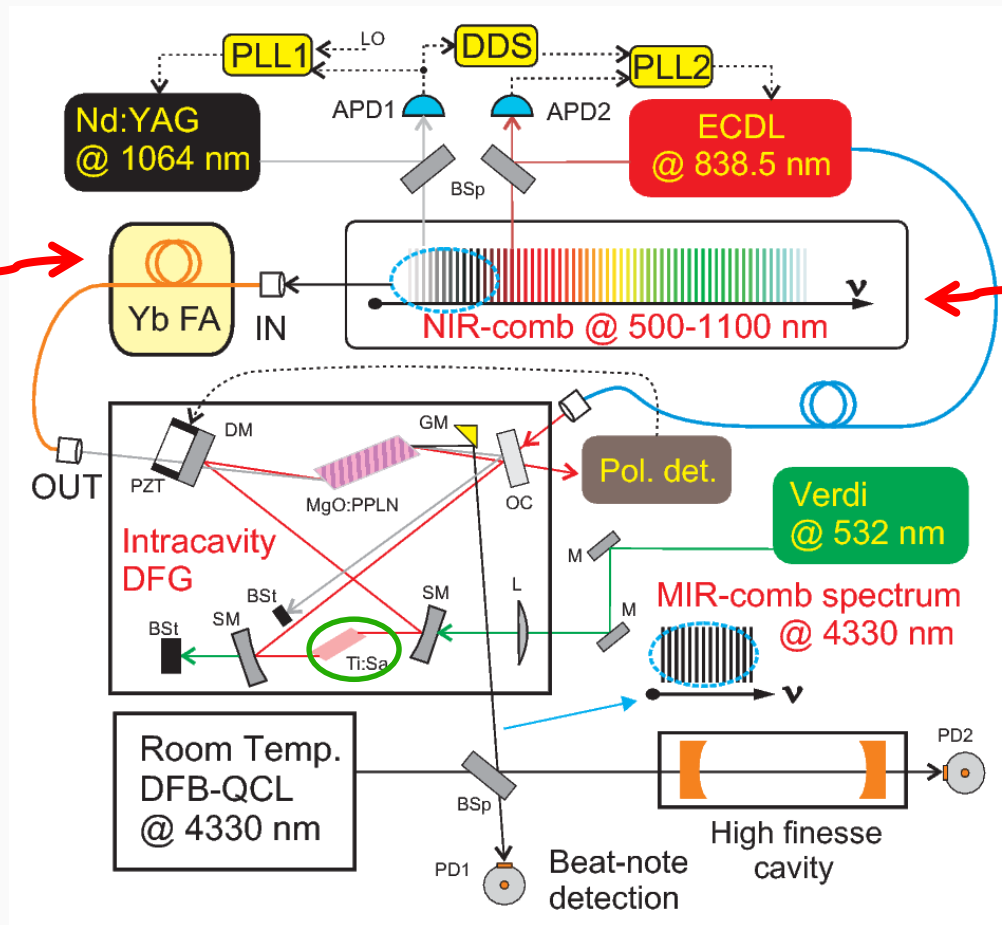
A. Campa *et al.*
Optics Express,
23, 3751 (2015).

IC-DFG-comb generation: the setup

First approach: **intracavity DFG generation**.

Yb FA output
(*signal*):
1032 – 1045 nm

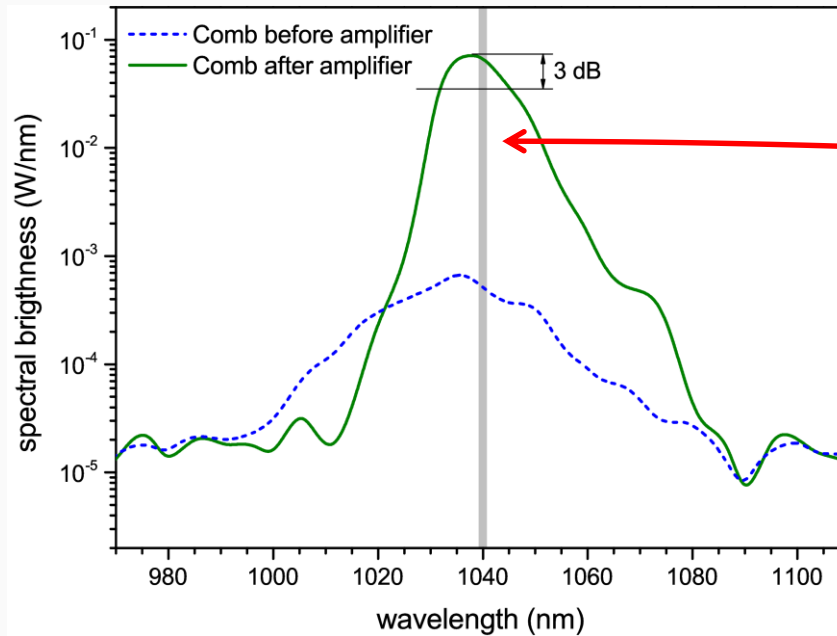
NIR-comb:
mode-locked Ti:Sa
with $f_r \approx 1$ GHz



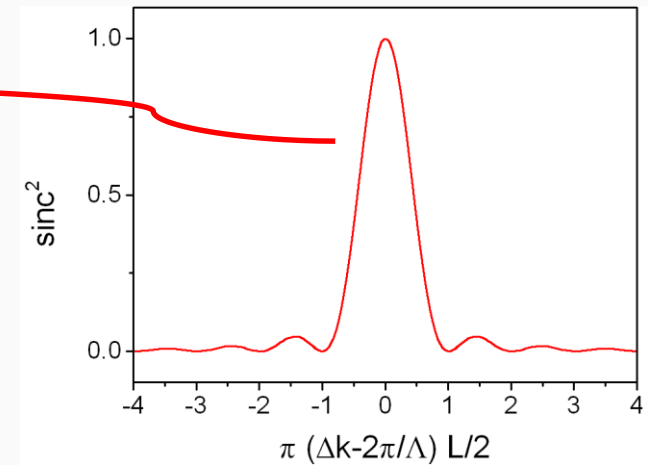
Galli et al., Opt. Express 21, 28877-28885 (2013).

IC-DFG-comb : the characteristics

NIR-comb spectrum



IC-DFG-comb envelope
(phase-matching bandwidth)



FWHM = 27 nm

	tunability	power
ECDL/Ti:Sa	838 – 863 nm	30 W
NIR-comb	1032 – 1045 nm *	5 W (ampl.)
IC-DFG-comb	4.2 – 5.0 μm	0.5 mW

* amplified NIR-comb
FWHM

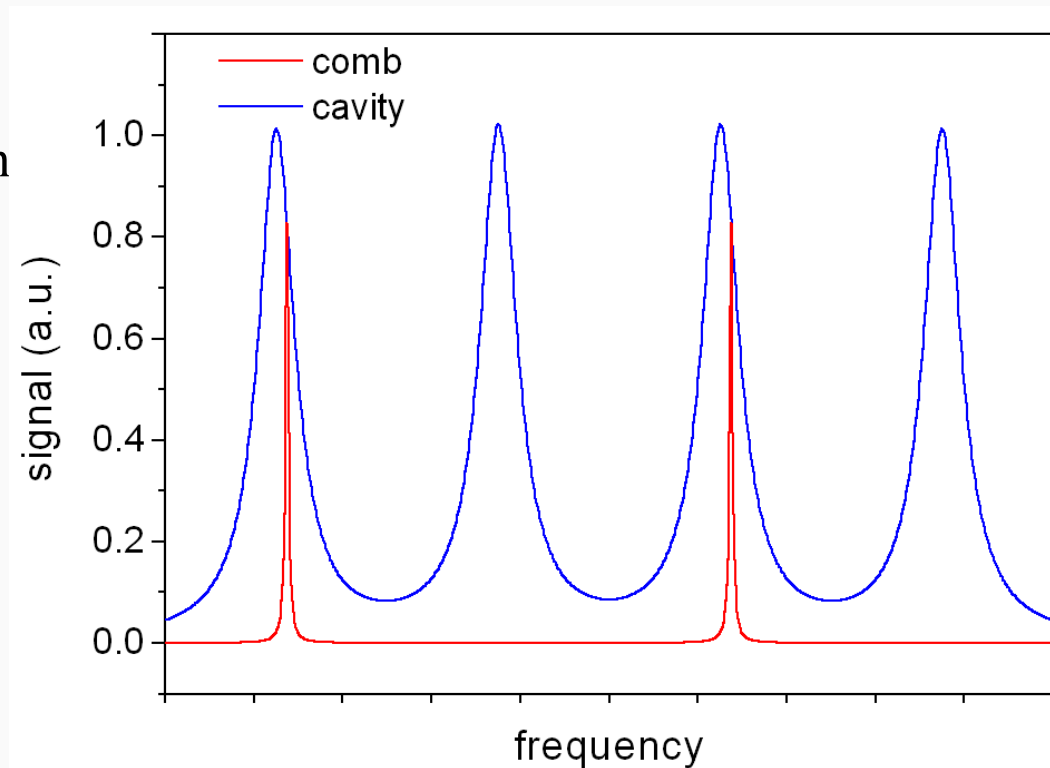


IC-DFG-comb : the high-finesse cavity

Characteristics:

Length = 1 m Finesse = 8000

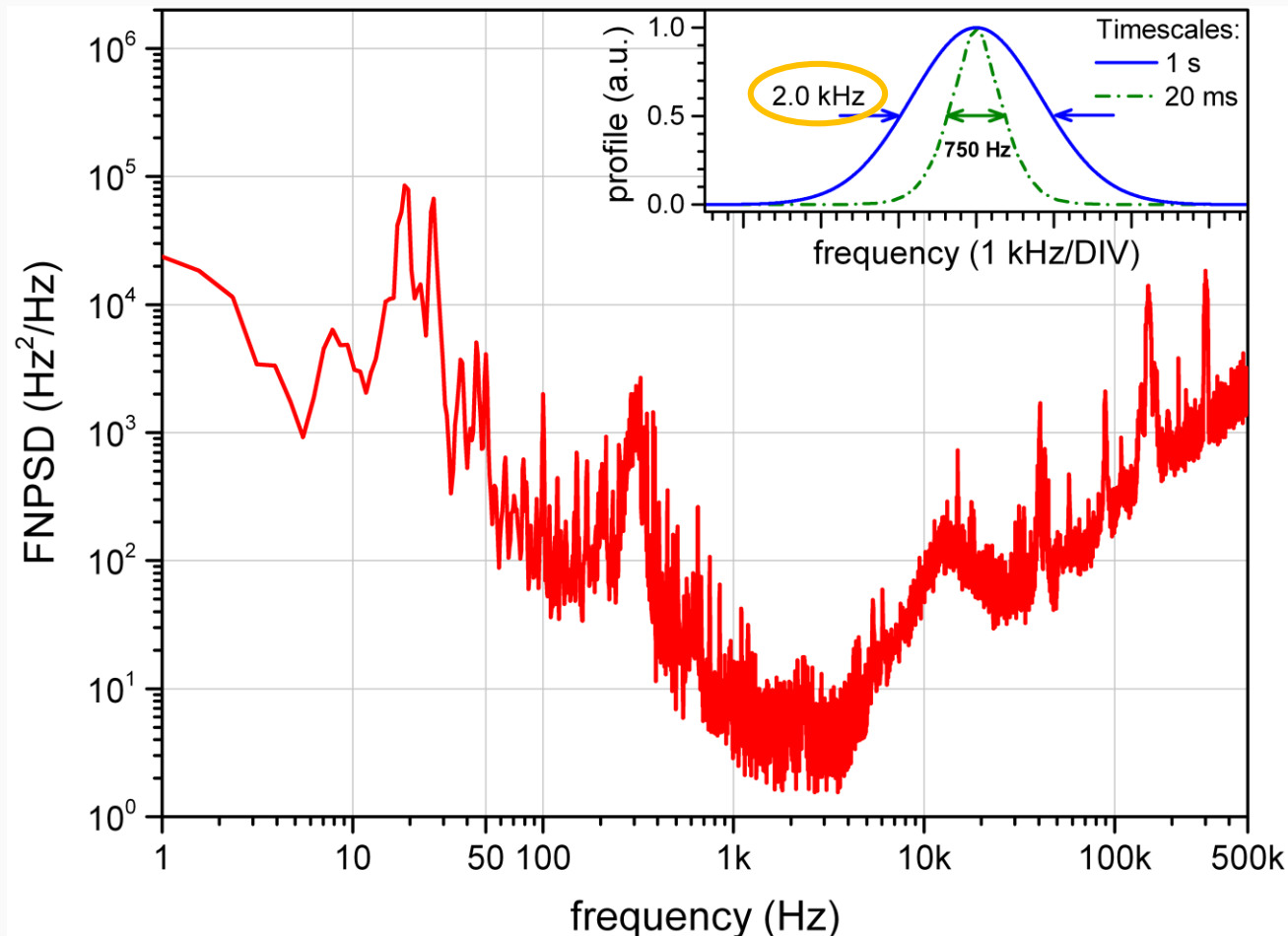
To **maximize** the transmission we have matched the repetition rate of the comb and the free spectral range of the cavity.





IC-DFG-comb : the teeth linewidth

The FNPSP of the comb in matching condition:



2 kHz linewidth
instead of
100 kHz
thanks to the **DDS**

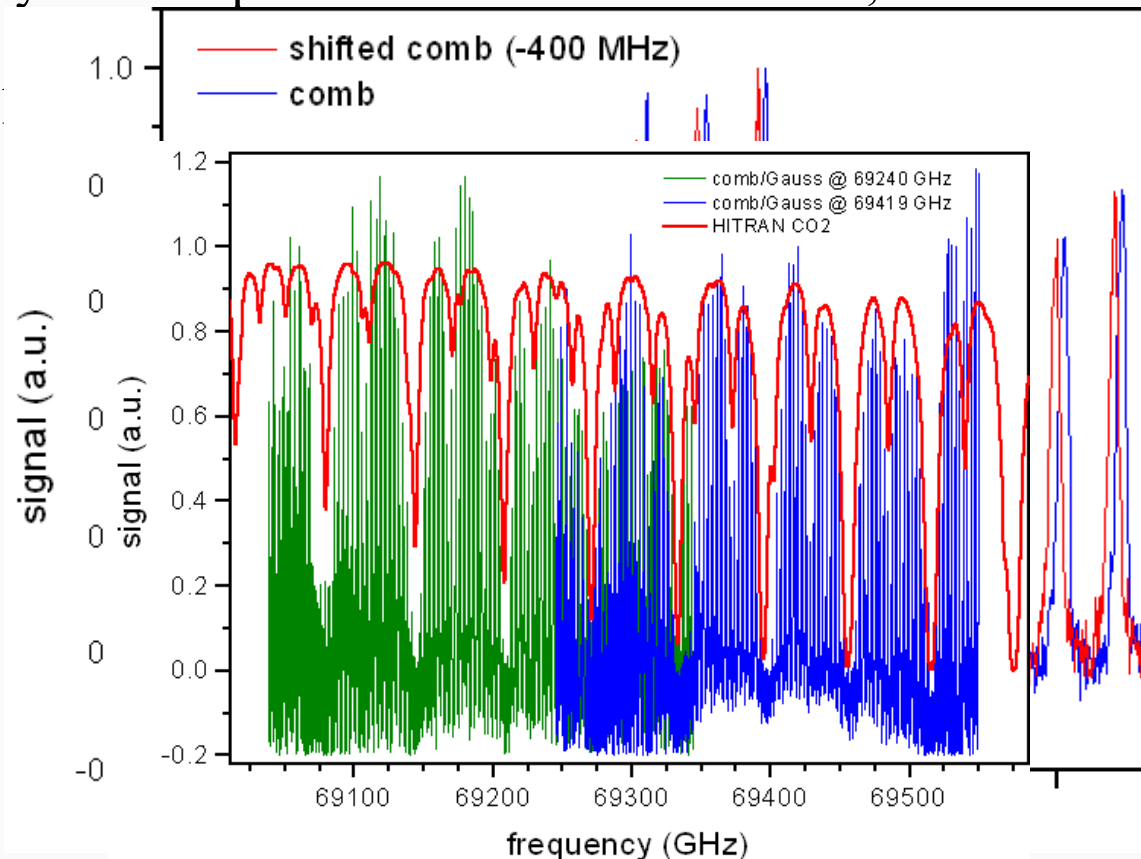


IC-DFG-comb : CO₂ spectroscopy

It is also possible to shift the comb:

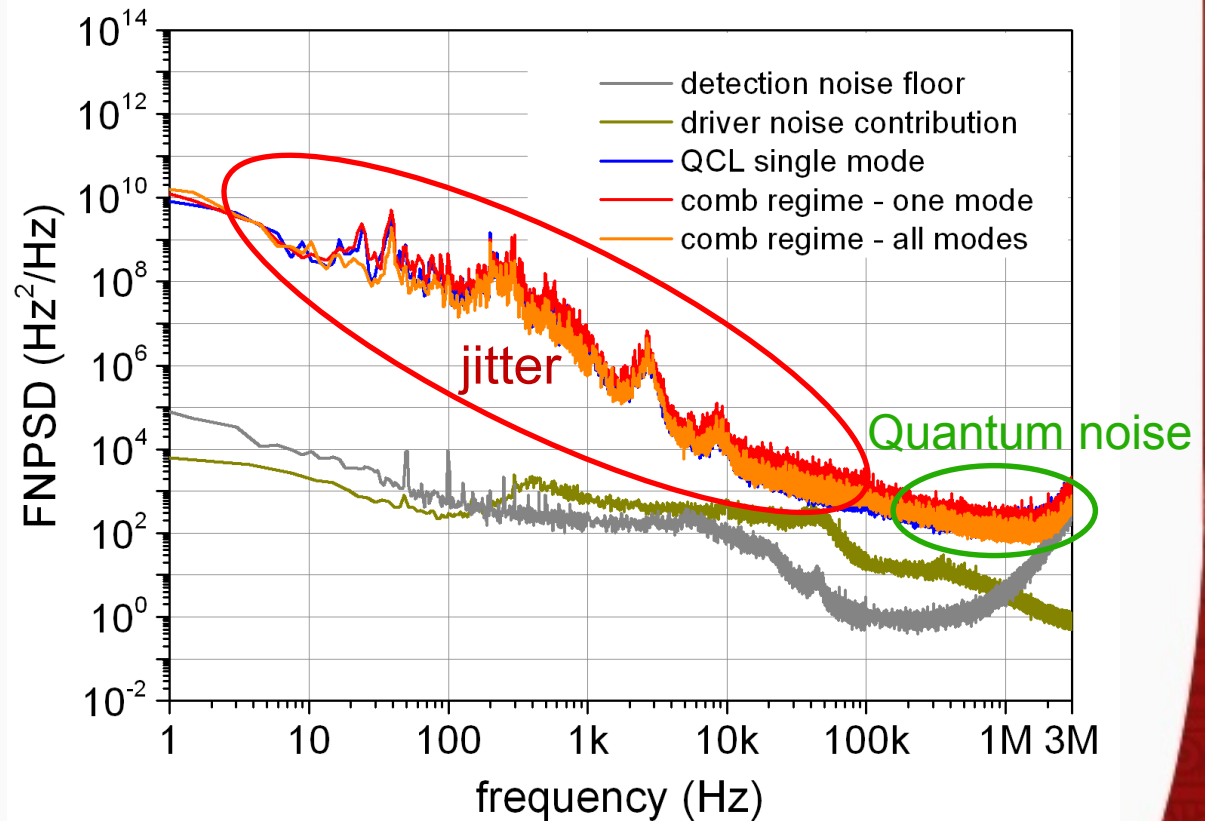
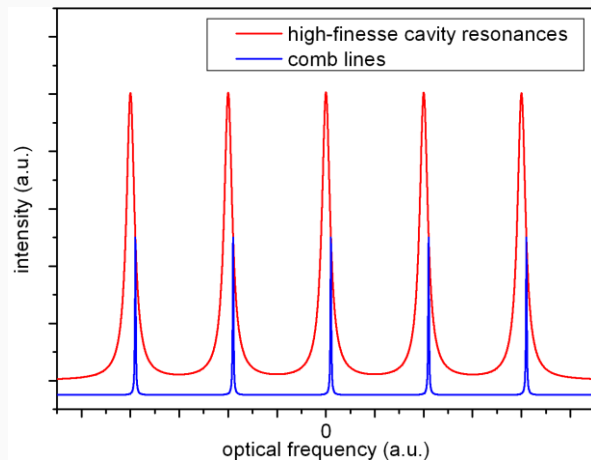
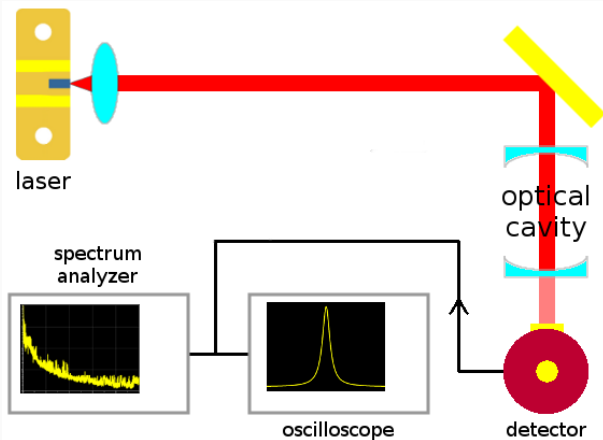
- **fine shift** in order to cover the frequencies between the teeth, either by tuning the pump frequency or the repetition rate of the NIR-comb;

- **wide shift** in temperature (c



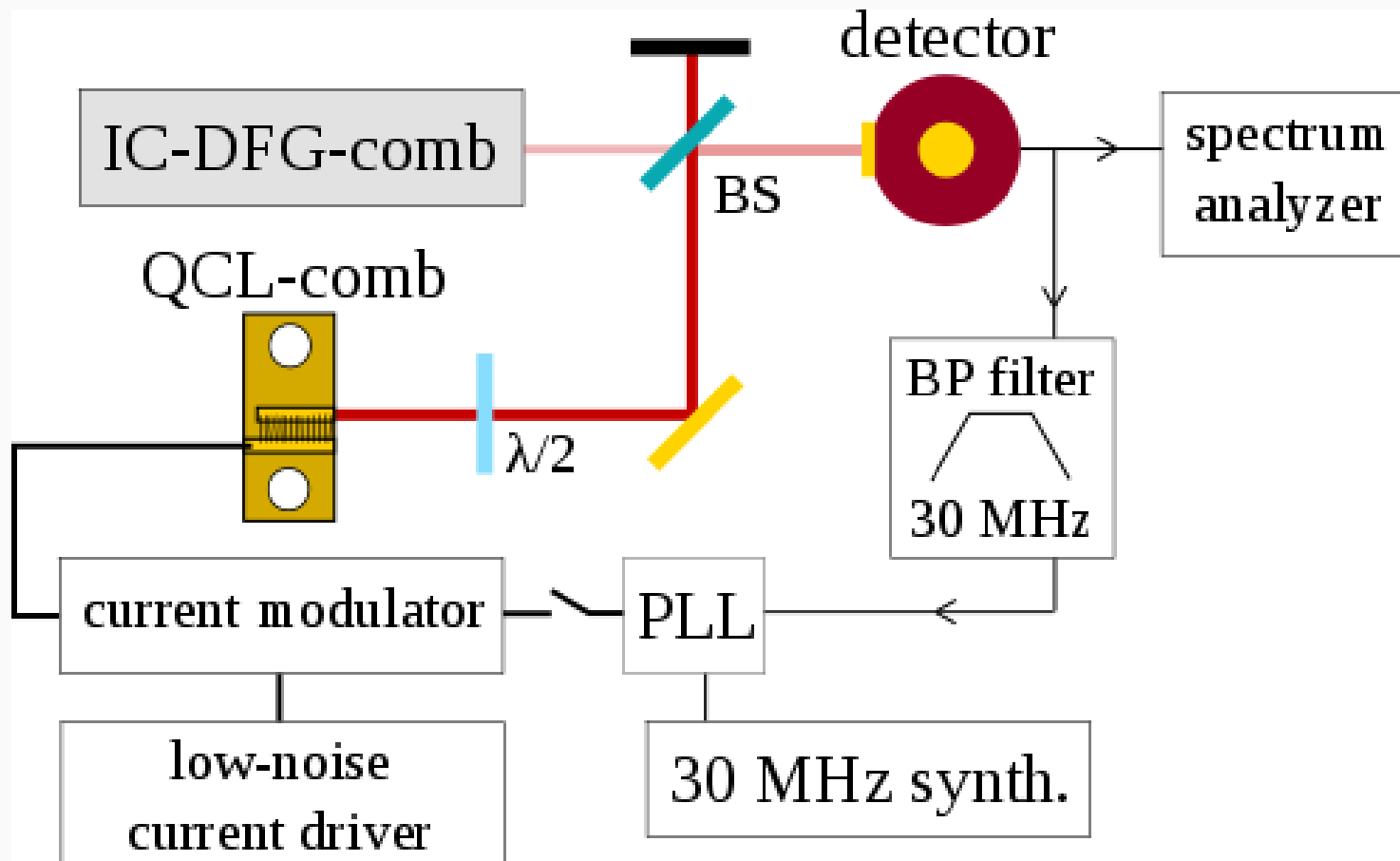
PLN crystal

QCL-comb: the frequency noise



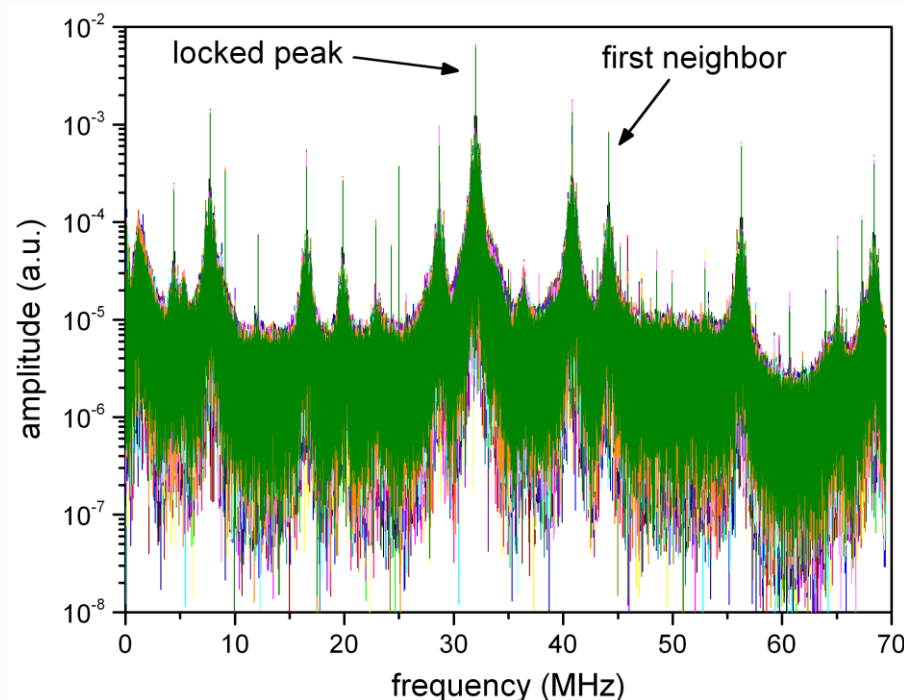
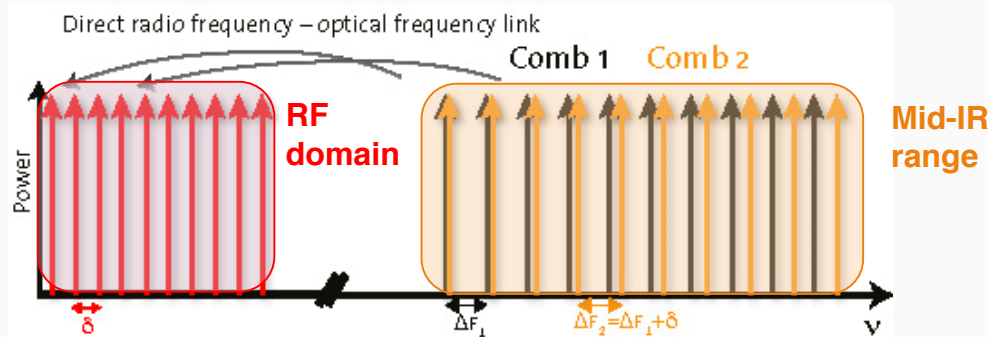
The modes are coherent among themselves,
but the overall emission has to be stabilized.

QCL-comb: the phase-locking of a mode to an IC-DFG-comb mode



QCL provided by ETH Zürich, Switzerland.

The dual comb technique



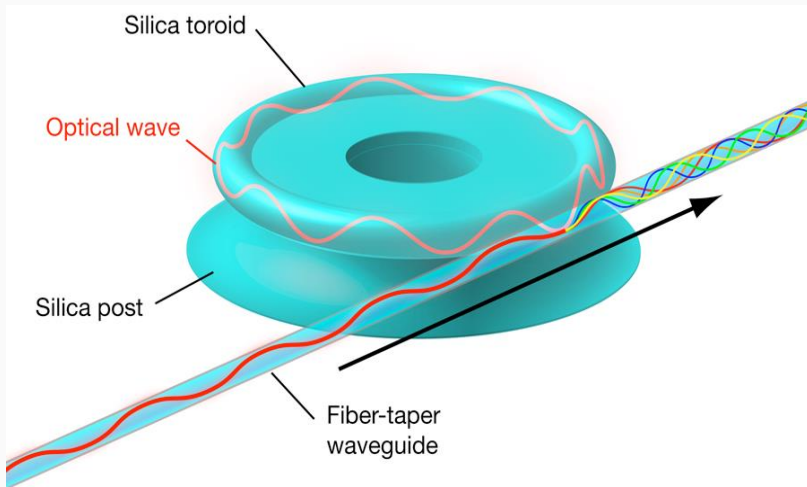
Dual comb technique:
down-conversion of the IR
spectra down to the RF.

← 20 frames 2-ms long

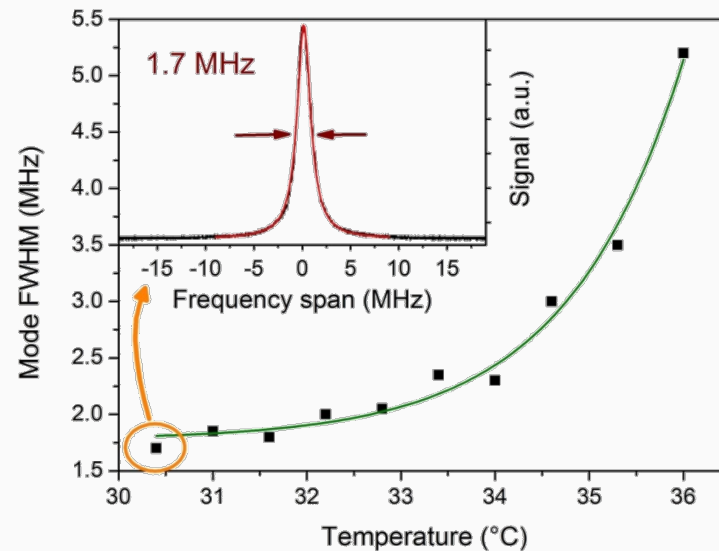
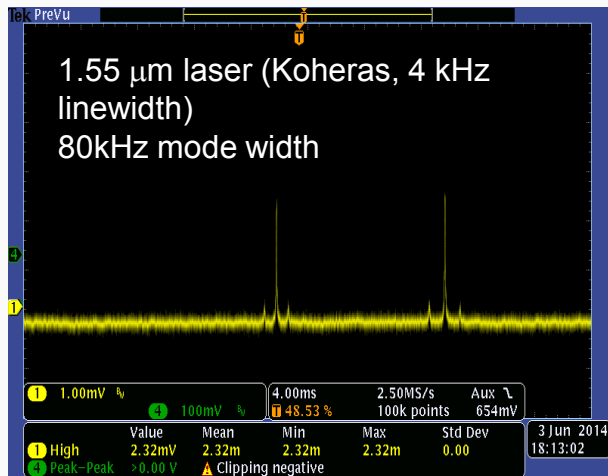
The locking effectively
narrows the peaks from 500
kHz to **less than 20 kHz** on
a 1-s time scale.



Whispering Gallery Mode micro-Resonators



- coupling through evanescent waves
- compact and robust devices
- transparent from near- to mid-IR (CaF₂, MgF₂)
- suitable for laser stabilization and comb generation

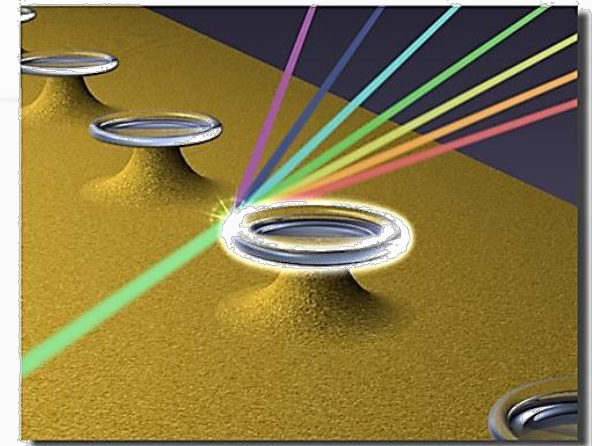
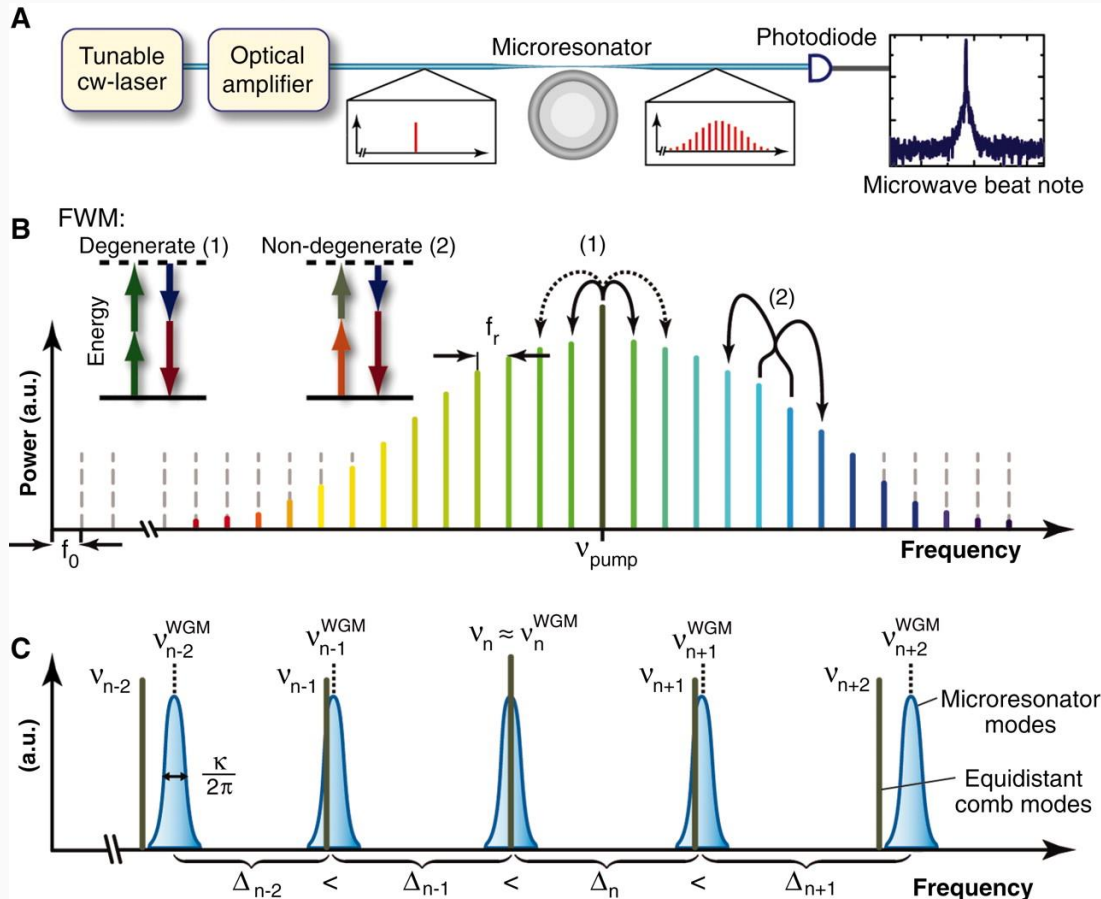


- modes as narrow as <10 kHz in the near-IR ($Q \sim 10^{10}$)

- ~ 1 MHz width in the mid-IR ($Q \sim 10^7 \div 10^8$)



Kerr-comb generation



**Kerr-comb generation
by four wave mixing**

from Kippenberg et al.,
Science **332**, 555-559
(2011)

**high Q
small
size**



**nonlinear optical phenomena can be
generated at very low coupled power
levels**



Kerr-comb generation using a silica microresonator

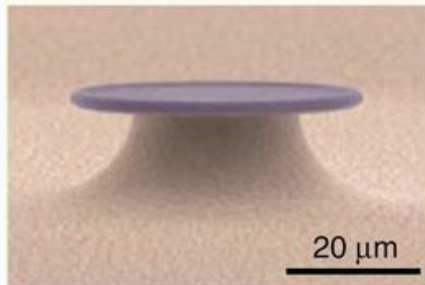
nature

Vol 450 | 20/27 December 2007 | doi:10.1038/nature06401

LETTERS

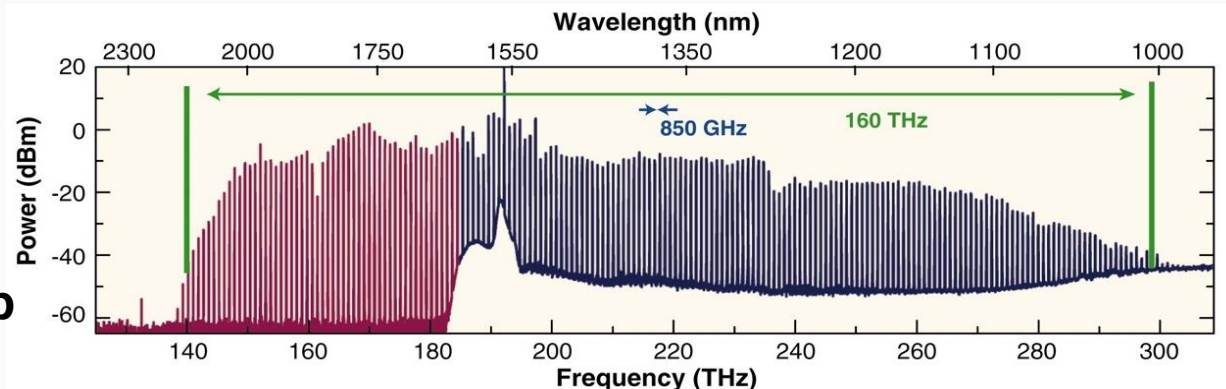
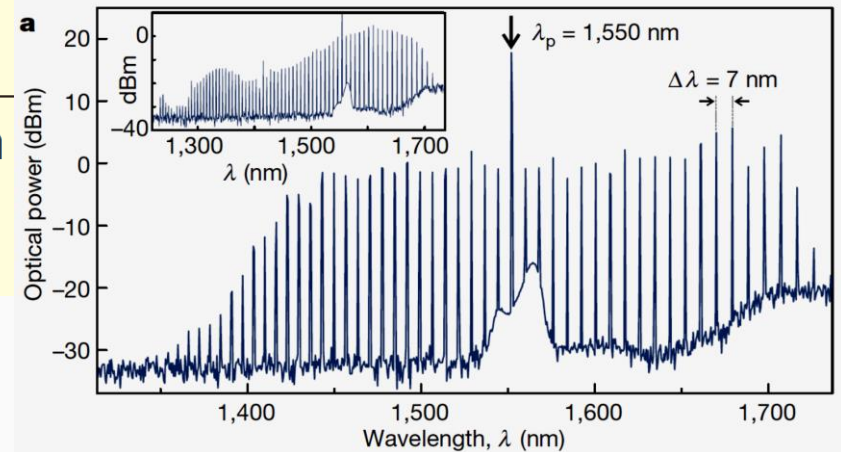
Optical frequency comb generation from a microresonator

P. Del'Haye¹, A. Schliesser¹, O. Arcizet¹, T. Wilken¹, R. Holzwarth¹ & T. J. Kippenberg¹



Silica

Octave-spanning comb



Del'Haye, Herr, Gavartin, Gorodetsky, Holzwarth, and Kippenberg, Phys. Rev. Lett. **107**, 063901 (2011)

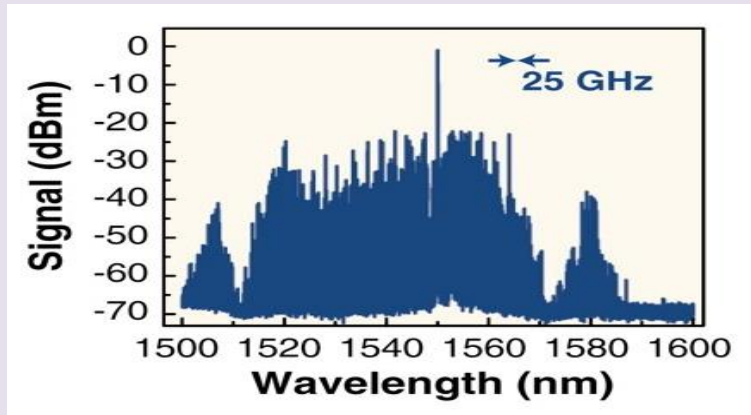


Kerr-comb generation

Tunable Comb with a Crystalline CaF₂ Whispering Gallery Mode Resonator



Crystalline



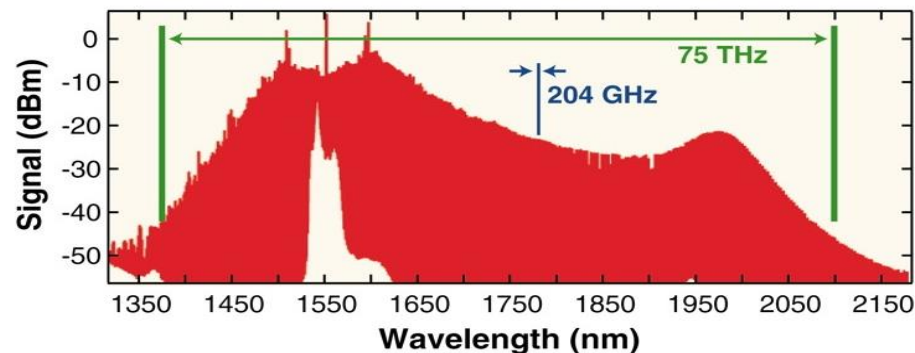
controllable tuning of the comb repetition frequency by changing the frequency of the pump laser

Savchenkov, Matsko, Ilchenko, Solomatine, Seidel, and Maleki, Phys. Rev. Lett. **101**, 093902 (2008)

Comb generated in integrated SiN resonator



Silicon nitride



Foster, Levy, Kuzucu, Saha, Lipson, and Gaeta, Opt. Express **19**, 14233-14239 (2011)



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Kerr-comb generation in the mid infrared



ARTICLE

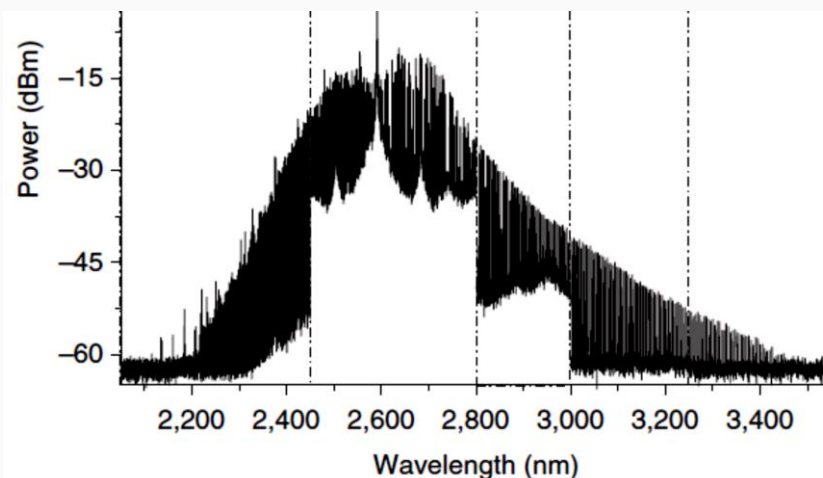
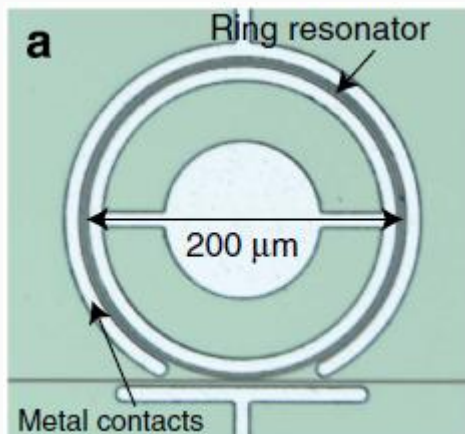
Received 10 Aug 2014 | Accepted 15 Jan 2015 | Published 24 Feb 2015

DOI: 10.1038/ncomms7299

Silicon-chip mid-infrared frequency comb generation

Austin G. Griffith¹, Ryan K.W. Lau², Jaime Cardenas¹, Yoshitomo Okawachi², Aseema Mohanty¹, Romy Fain¹, Yoon Ho Daniel Lee¹, Mengjie Yu², Christopher T. Phare¹, Carl B. Poitras¹, Alexander L. Gaeta^{2,3} & Michal Lipson^{1,3}

- SiN resonator pumped with an OPO (2.5 to 3.2 μm , 100-kHz linewidth)
- broadband frequency comb spanning from 2.1 to 3.5 μm





Kerr-comb generation in the mid infrared

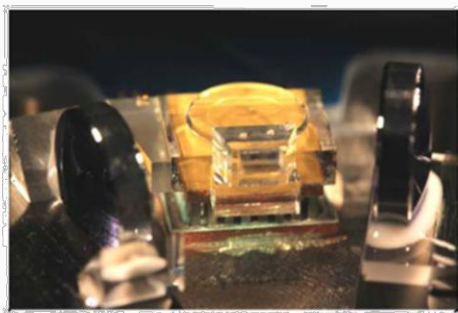
3468 Vol. 40, No. 15 / August 1 2015 / Optics Letters

Letter

Optics Letters

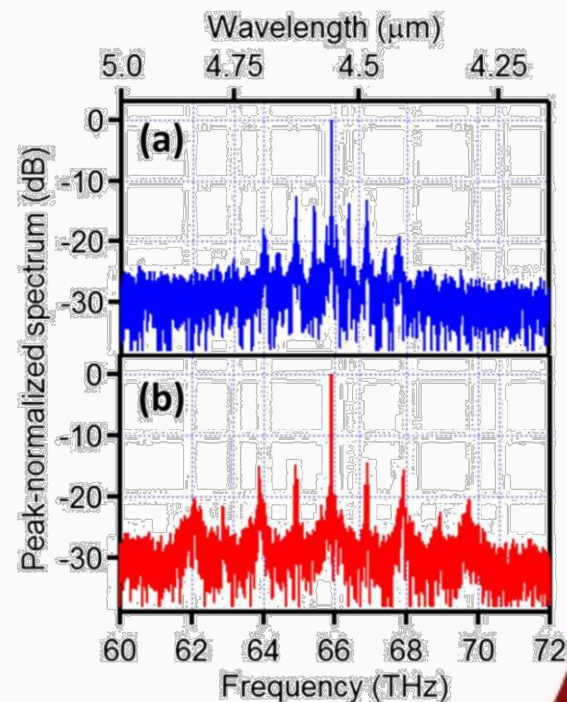
Generation of Kerr combs centered at 4.5 μm in crystalline microresonators pumped with quantum-cascade lasers

ANATOLIY A. SAVCHENKOV,¹ VLADIMIR S. ILCHENKO,¹ FABIO DI TEODORO,^{2,3} PAUL M. BELDEN,²
WILLIAM T. LOTSHAW,² ANDREY B. MATSKO,^{1,*} AND LUTE MALEKI¹

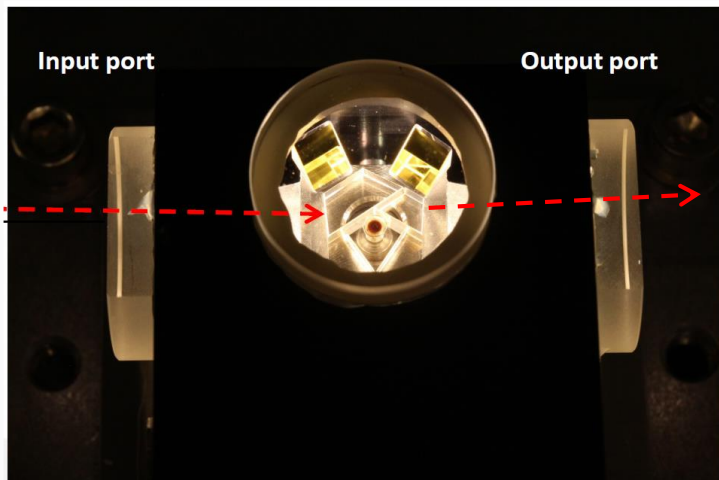


➤ high-finesse CaF₂ and MgF₂ resonators pumped with QCLs

➤ 4.5 μm half-octave spanning comb
➤ 60 kHz linewidth



CaF₂ micro-resonator for mid-IR laser stabilization



CaF₂ toroidal WGMR

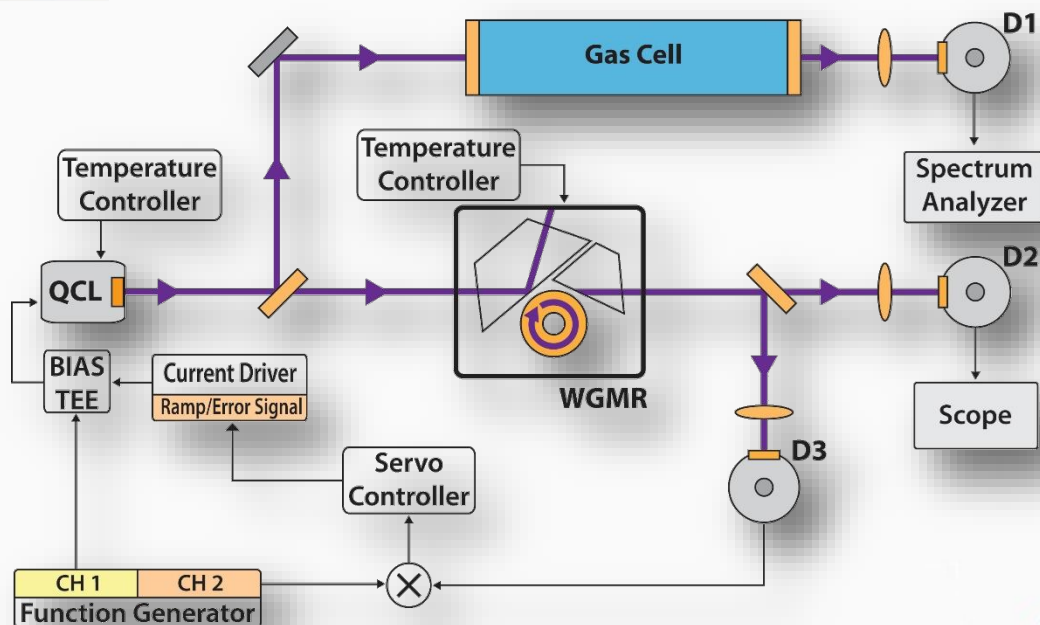
3.6 mm diameter --> FSR=18.9 GHz @ 4.3 μ m

Coupling: Al₂O₃ (sapphire) prisms

Q-factor $\sim 2 \times 10^9$ @ 1.5 μ m ($F=250k$)

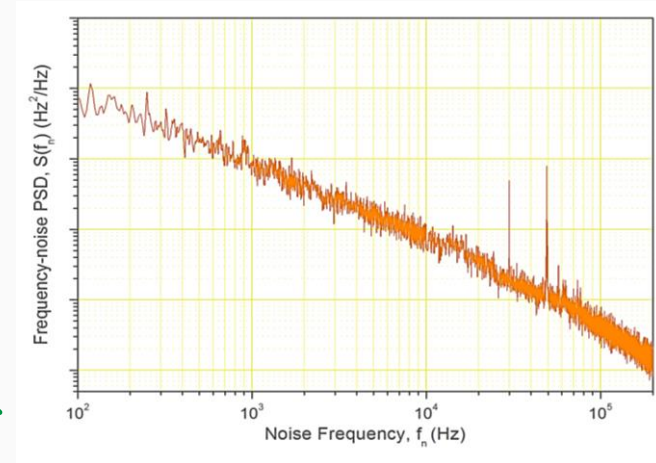
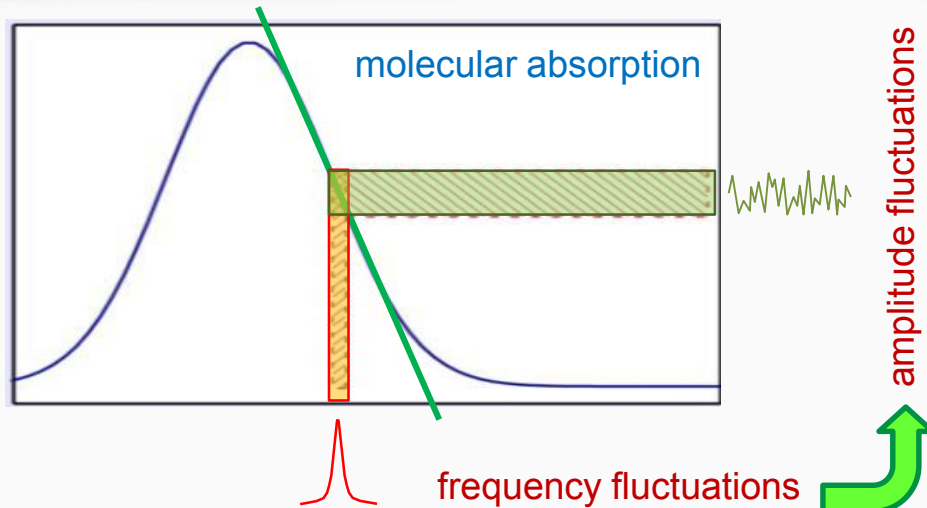
$\sim 4 \times 10^7$ @ 4.3 μ m ($F=11k$)

- DFB QCL @ 4.3 μ m (Hamamatsu)
700 mA thr., 10 mW
- low-noise current driver
- Reference cell with pure CO₂





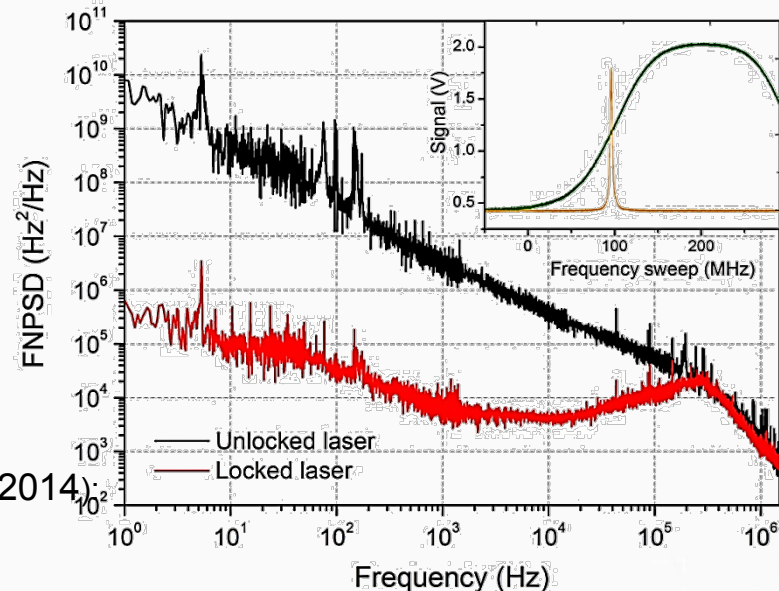
Measuring the laser frequency noise



Free-running QCL:
700 kHz FWHM (1 s timescale)

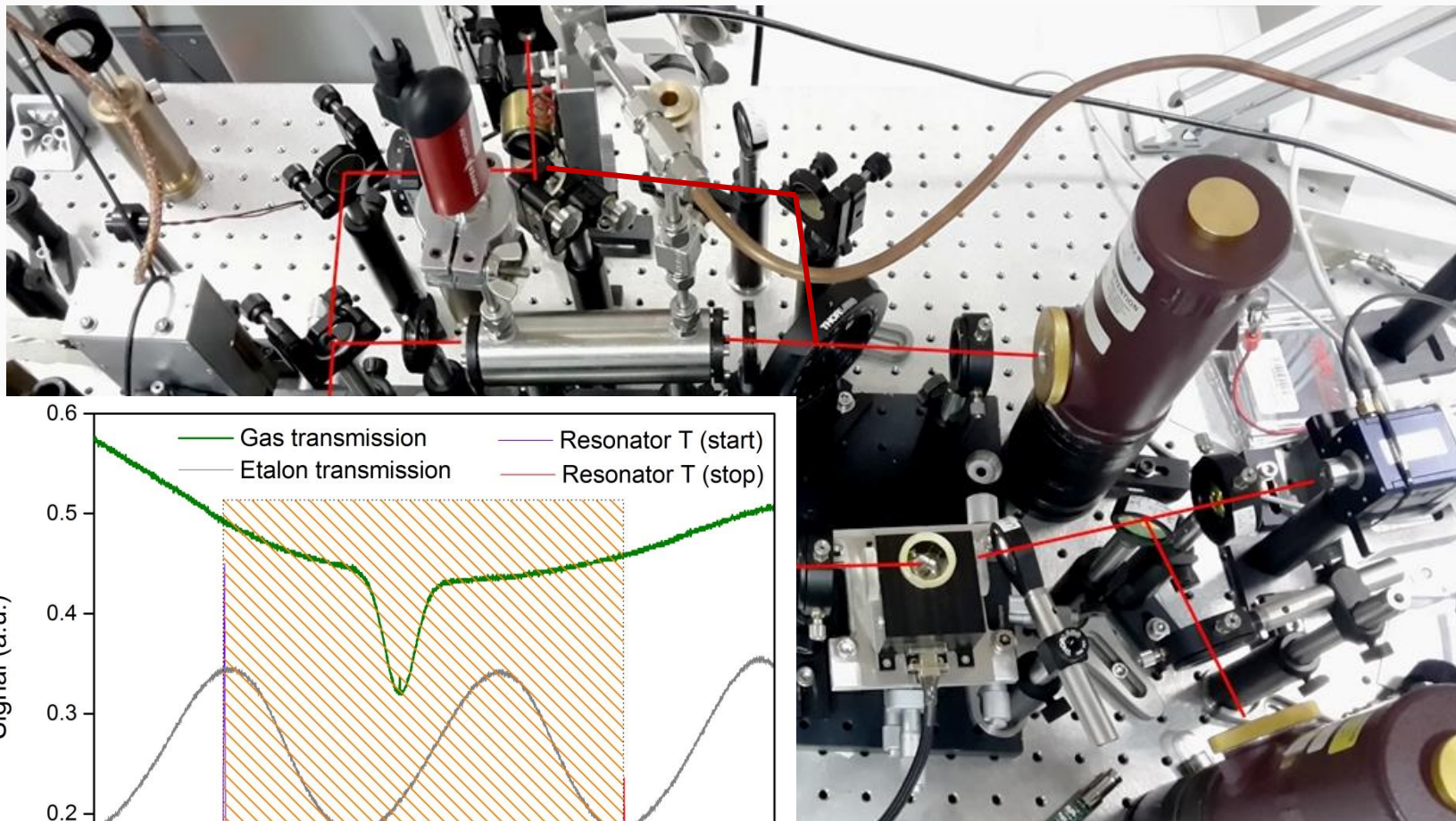
Locked QCL:
15 kHz FWHM (1 s timescale)
10 kHz (1 ms)

Locking to V-shaped cavity (Fasci et al. OL 39, 4946, 2014):
4 kHz (1 ms); 1 MHz (1 s)





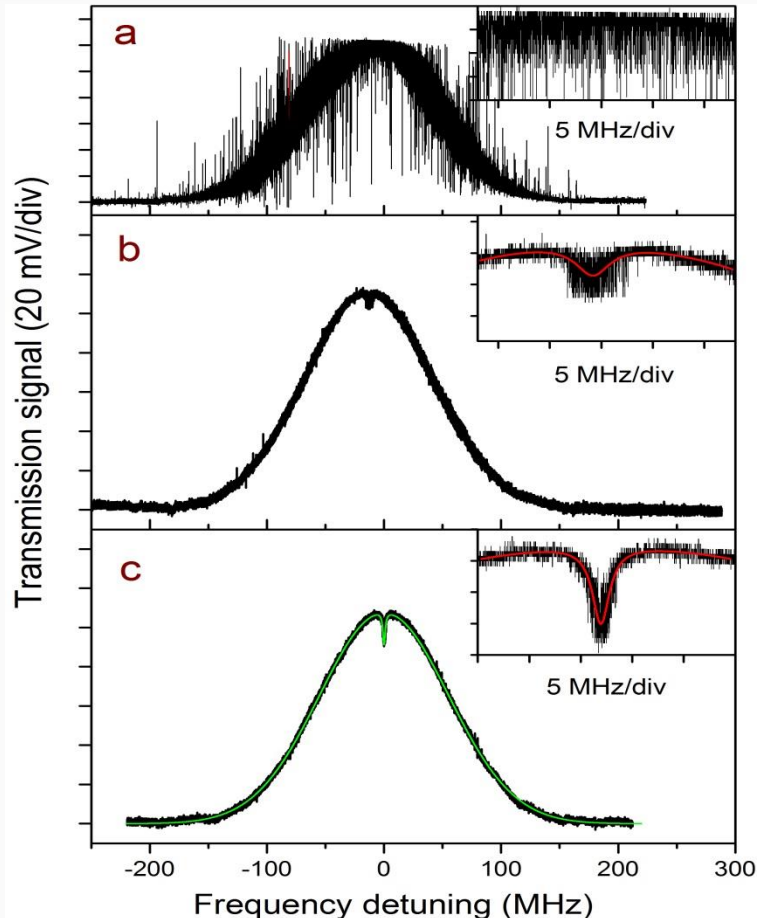
Doppler-free spectroscopy – electronic lock



More than 1.5 GHz tuning
by acting only on the resonator T



Doppler-free spectroscopy – electronic lock



Commercial vs home-made low-noise driver
free-running vs locked laser

- Pump power: about 5 mW
- 30 seconds span time

Improved precision on center
frequency determination by more
than a factor 3

WORK IN PROGRESS...

Borri et al., *Tunable microcavity-stabilized quantum cascade laser for mid-IR high-resolution spectroscopy and sensing*, Sensors 2016 (in press)

Physics

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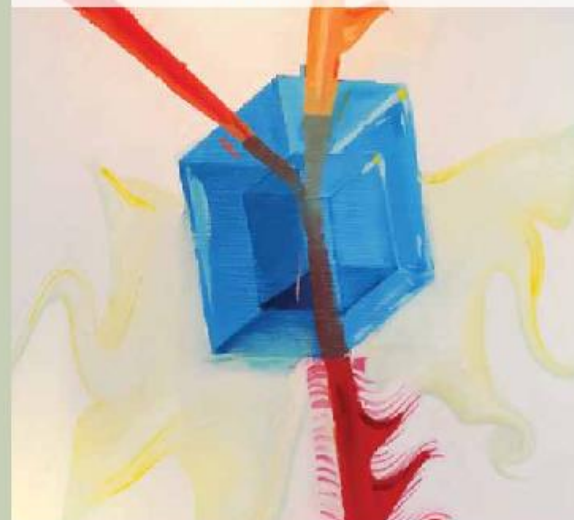


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A Handbook



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