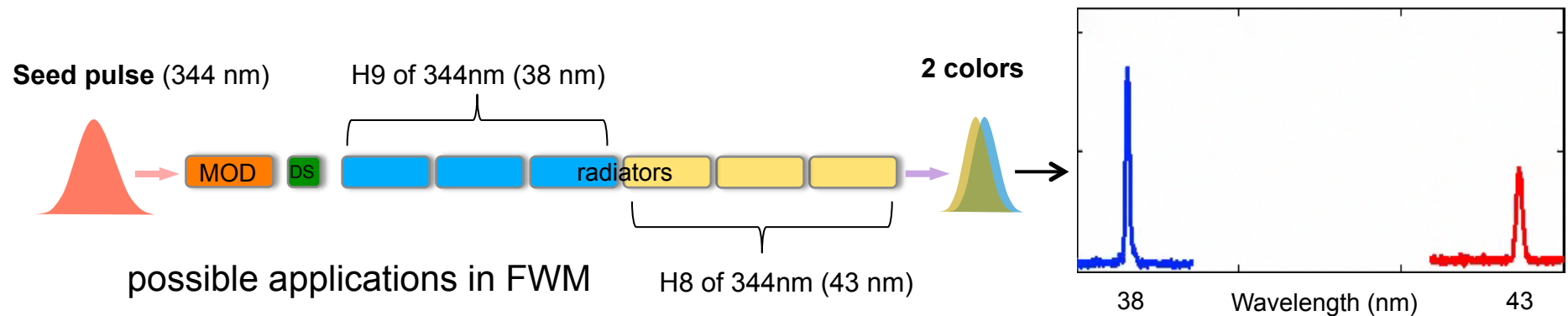
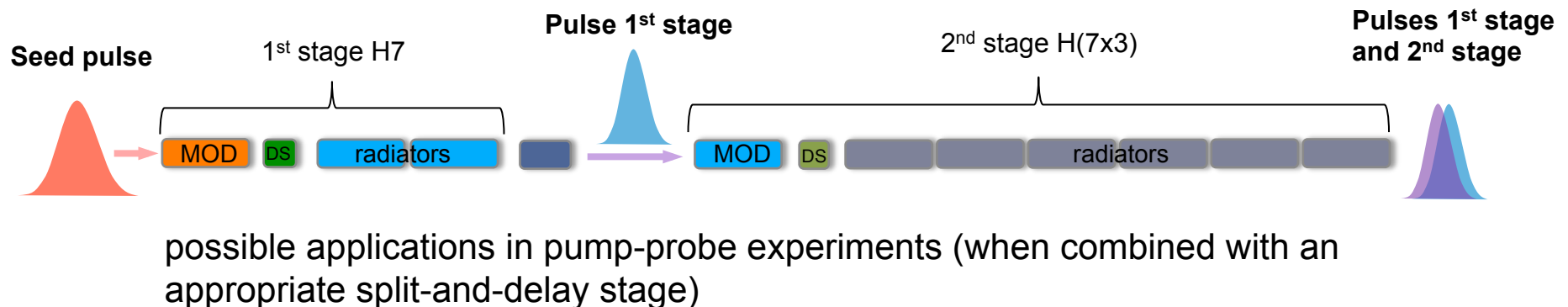


TWO COLOR OPERATION AT THE FERMI SEEDED FEL: ONE SEED PULSE

One seed pulse → two color “zero-delay” FEL pulses using the split undulator scheme

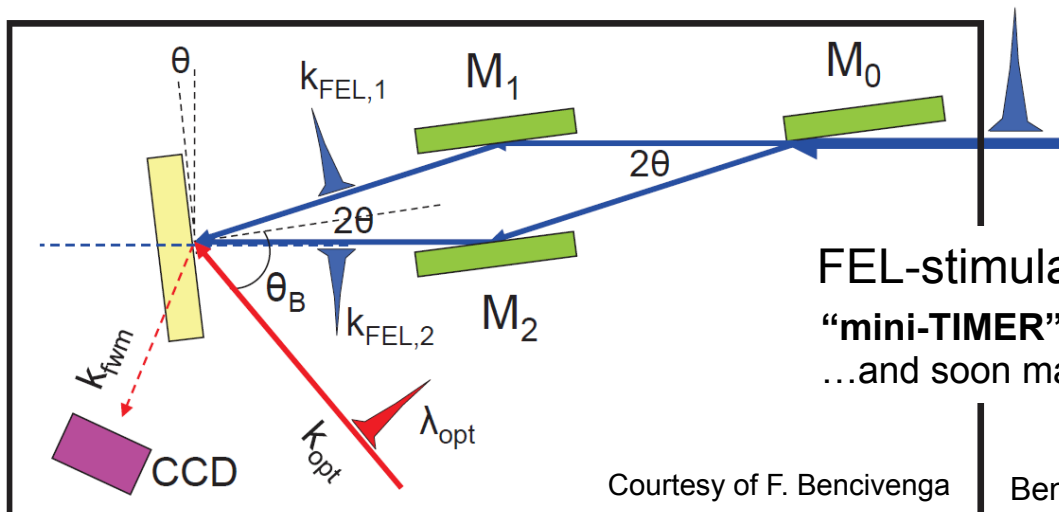
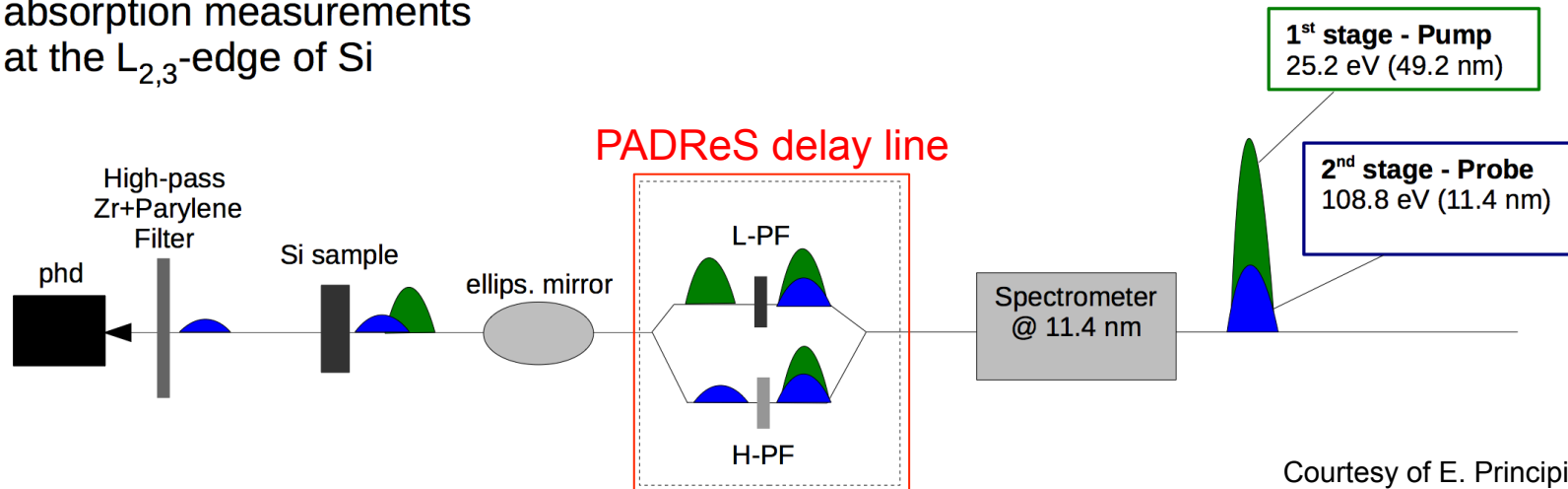


One seed pulse → two color “zero-delay” FEL pulses using the two-stage FEL2



SPLIT AND DELAY

EIS-TIMEX FEL/pump-FEL/probe
absorption measurements
at the $L_{2,3}$ -edge of Si



FEL-stimulated transient grating

“mini-TIMER”@DiProl

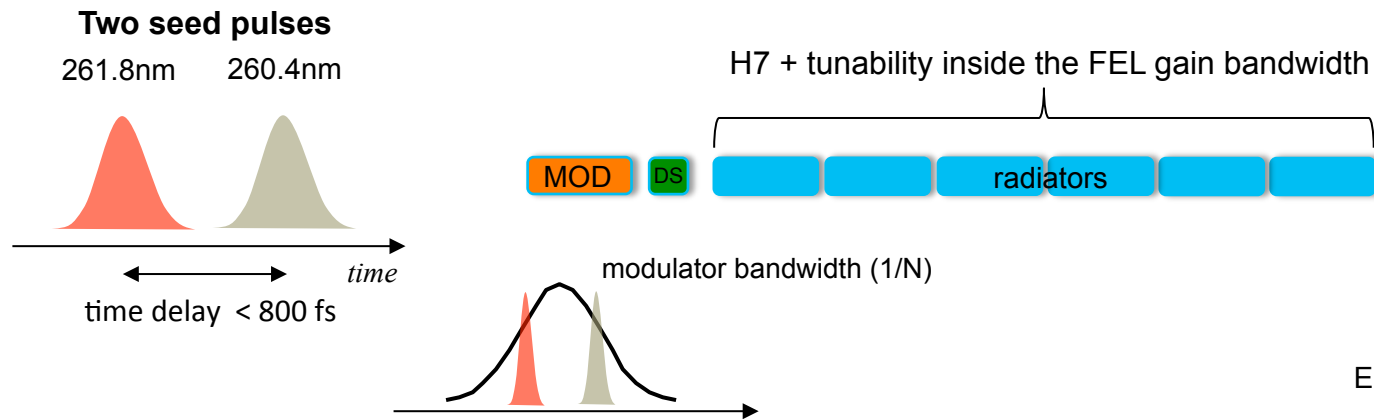
...and soon maxi-TIMER @EIS-TIMER beamline

Courtesy of F. Bencivenga

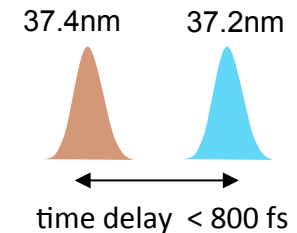
Bencivenga *et al.*, *Nature* 2015

TWO COLOR OPERATION AT THE FERMI SEEDED FEL: TWO SEED PULSES

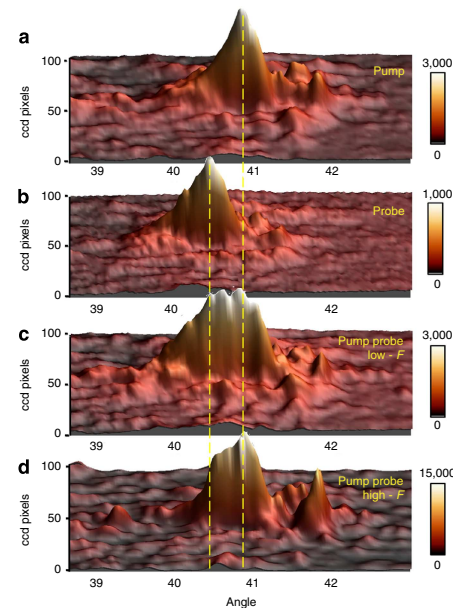
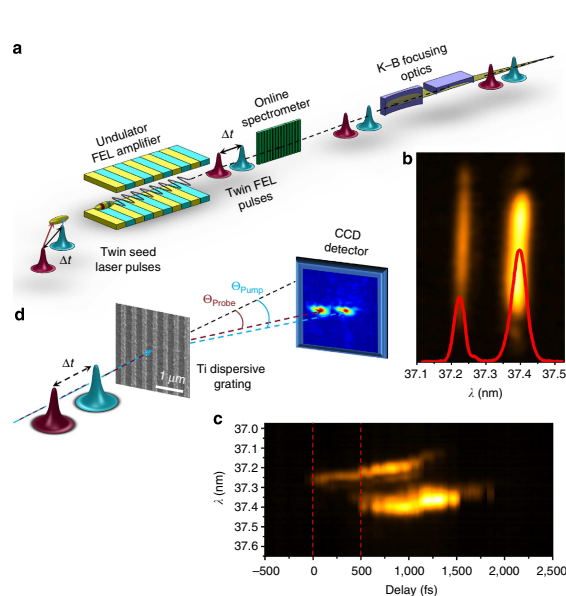
Two color seed pulses → two color FEL pulses inside the FEL bandwidth



**Two FEL pulses
(max wavelength
separation below 1%)**



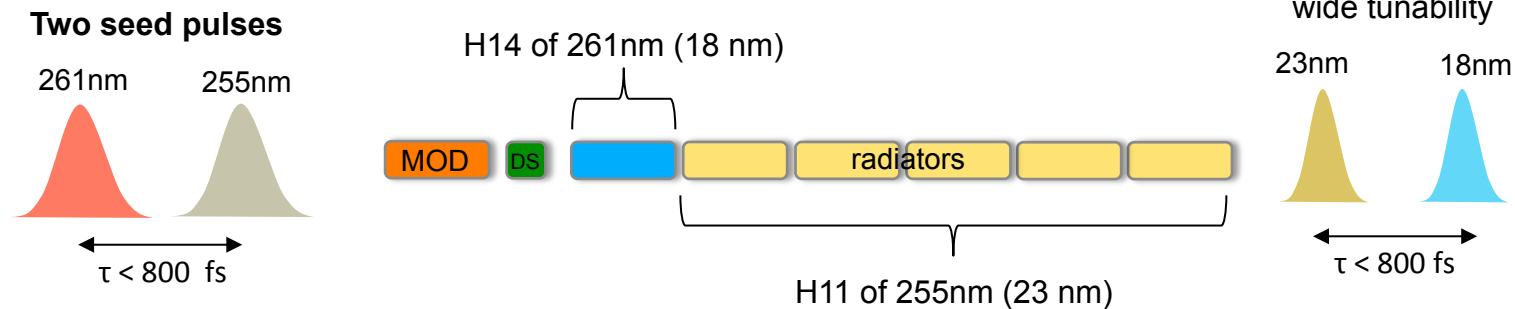
E. Allaria *et al.*, *Nat. Commun.*, 2013



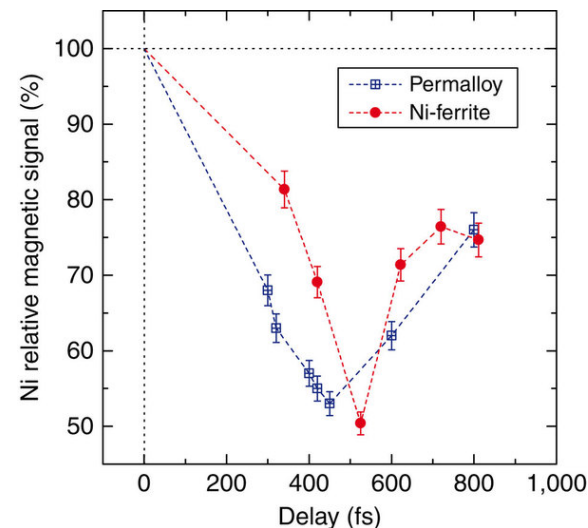
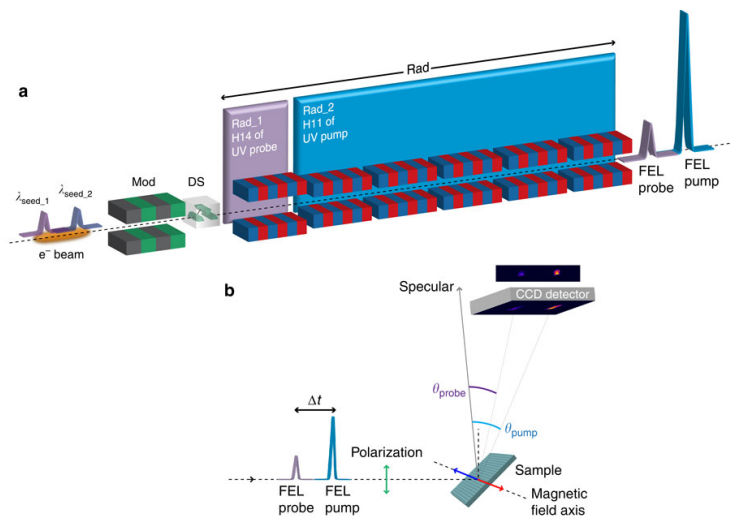
**probing structural
dynamics in solid
samples using
diffraction**

TWO COLOR OPERATION AT THE FERMI SEEDED FEL: TWO SEED PULSES

Two color seed pulses → split undulator scheme



E. Ferrari *et al.*, *Nat. Commun.*, 2016



**probing
demagnetization
dynamics in
magnetic
compounds**

EXOTIC TWO COLOR SCHEMES AND FULL SPECTRO-TEMPORAL SHAPING OF FEL PULSES AT FERMI

BUNCHING IN A SEEDED FEL

$$b_n(t) \sim J_n[-nBA(t)] \exp\{in[\phi_s(t) + \phi_e(t)]\}$$

Dispersive section strength points to B

Seed laser envelope points to $A(t)$

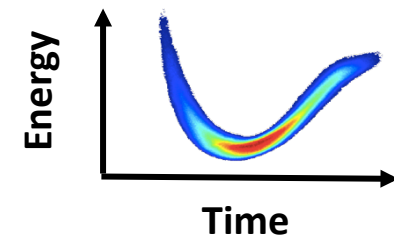
Bessel function points to J_n

Seed phase points to $\phi_s(t)$

Electron energy profile points to $\phi_e(t)$

Bunching envelope \longleftrightarrow FEL temporal profile

$\phi_s(t) + \phi_e(t)$ \longleftrightarrow FEL phase



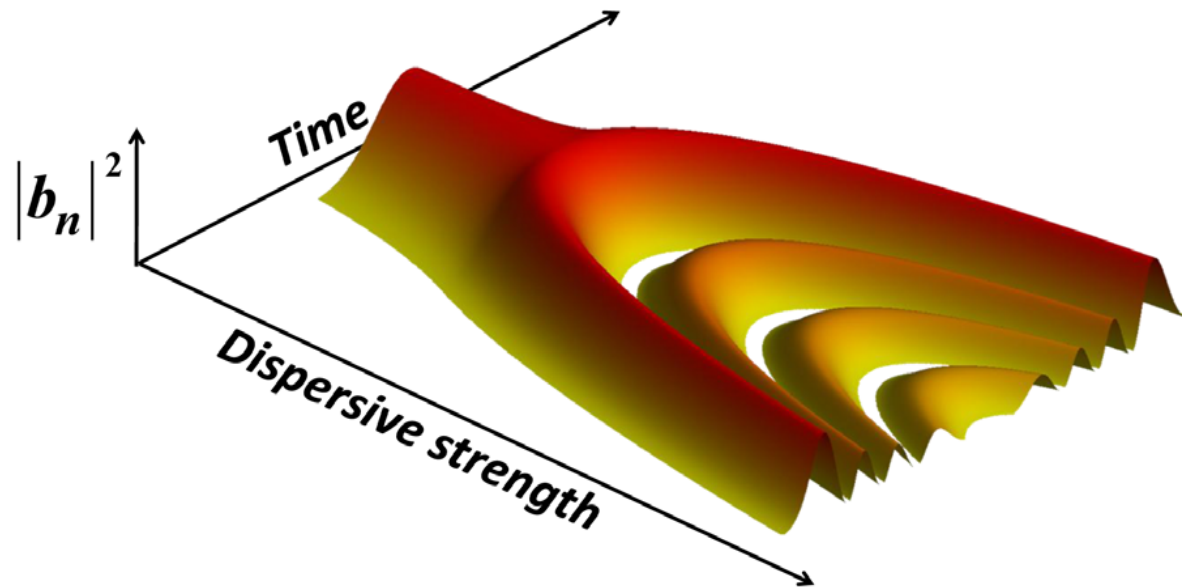
The FEL pulse can be shaped through the manipulation of the seed envelope $A(t)$ and phase $\phi_s(t)$

FEL PULSE ENVELOPE

Bunching at the n th harmonic

$$b_n(t) \sim J_n[-nBA(t)] \exp\{in[\phi_s(t) + \phi_e(t)]\}$$

Amplitude profile

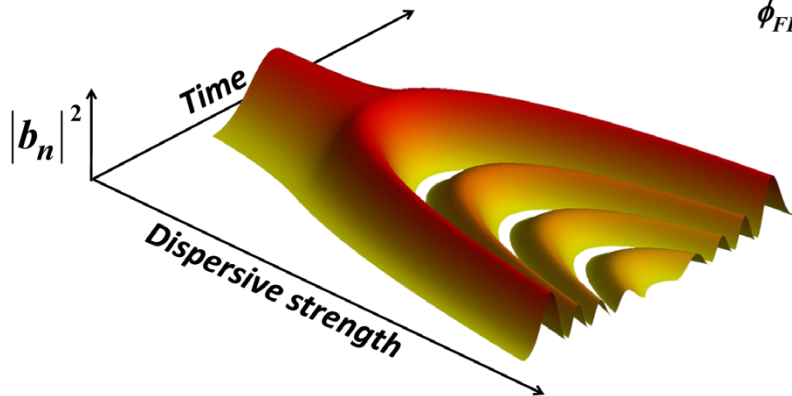


FEL SPECTRUM

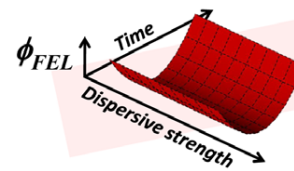
Bunching at the n th harmonic

$$b_n(t) \sim J_n[-nBA(t)] \exp\{in[\phi_s(t) + \phi_e(t)]\}$$

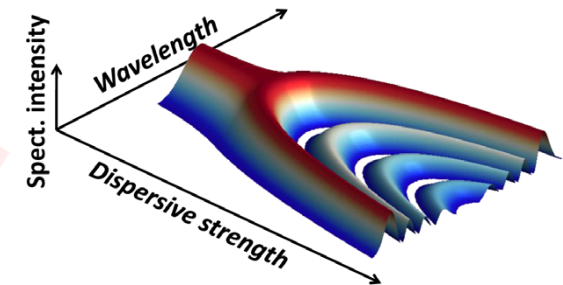
Amplitude profile



Phase profile

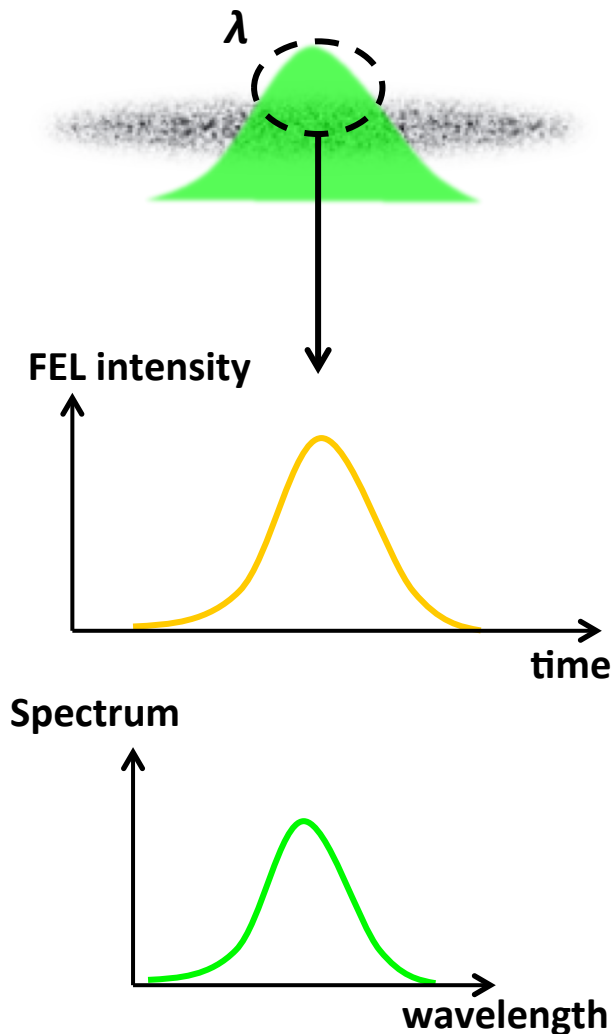


Chirped



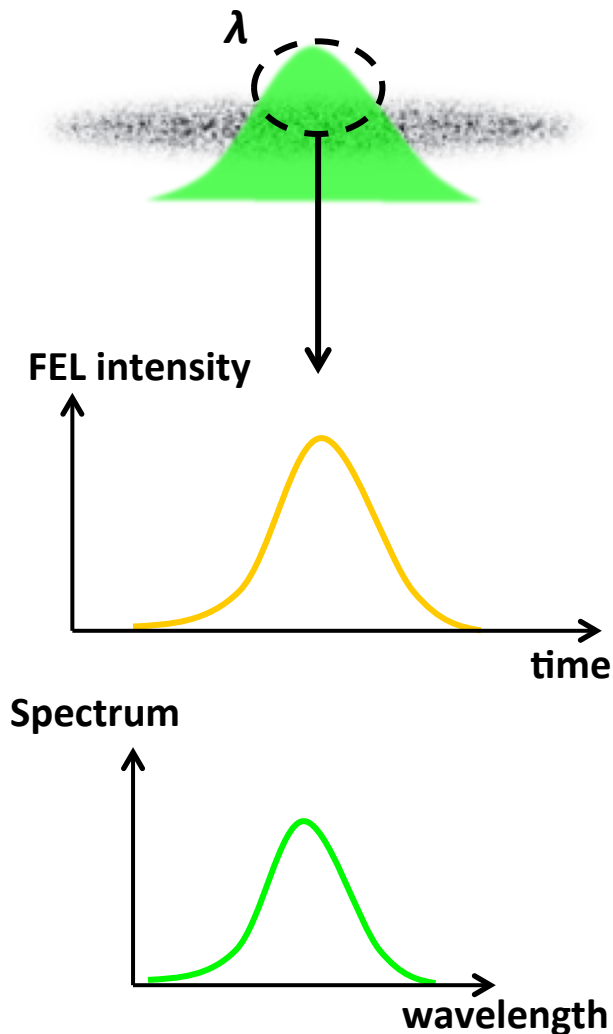
SEEDED FELs AS SELF-STANDING SOURCES FOR X-RAY PUMP – X-RAY PROBE EXPERIMENTS

moderate dispersive strength

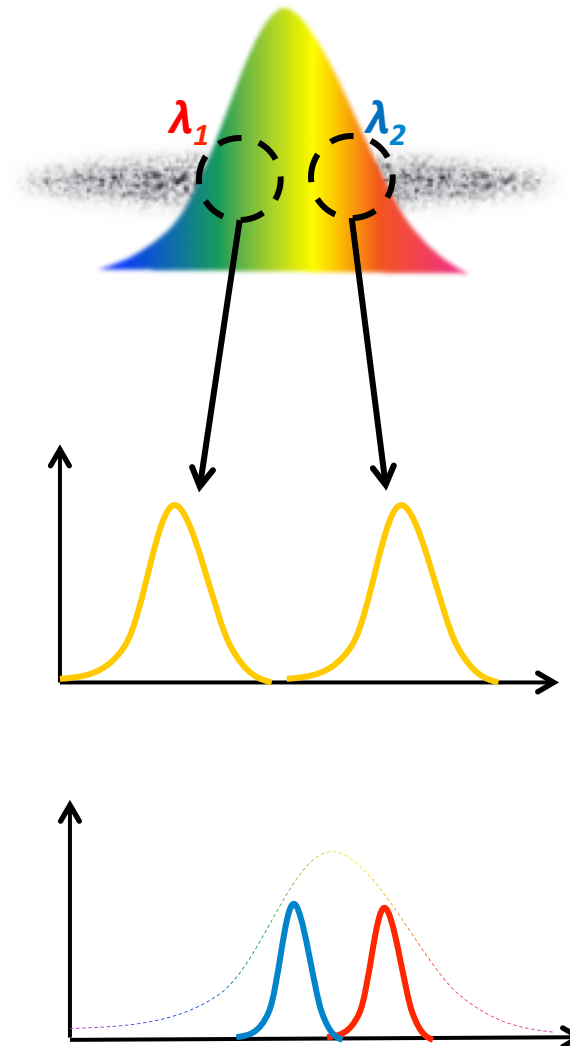


SEEDED FELs AS SELF-STANDING SOURCES FOR X-RAY PUMP – X-RAY PROBE EXPERIMENTS

moderate dispersive strength

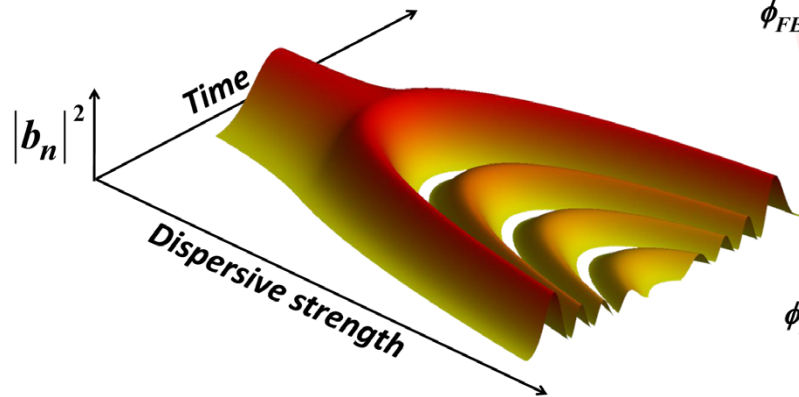


strong dispersive strength

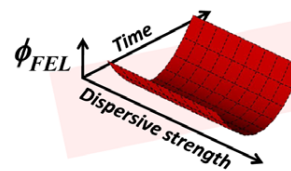


GENERATION OF TRANSFORM LIMITED FEL PULSES

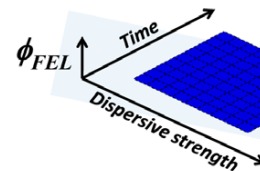
Amplitude profile



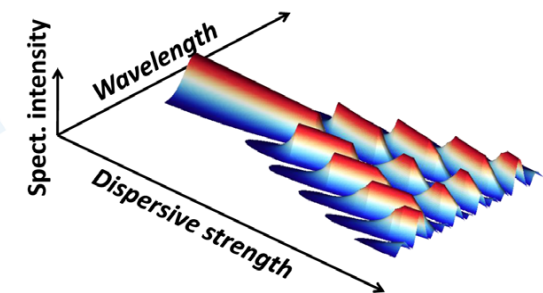
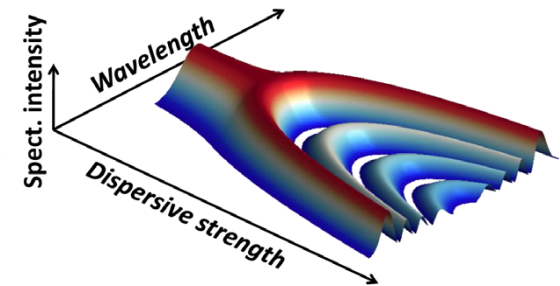
Phase profile



Chirped

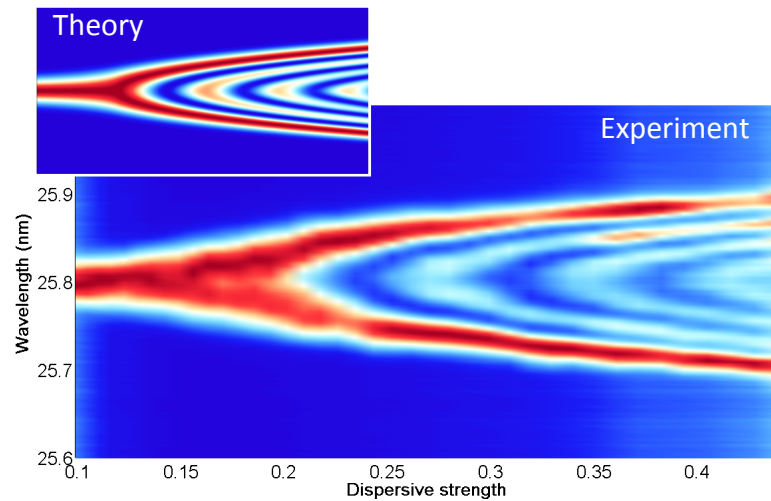


Fourier limit

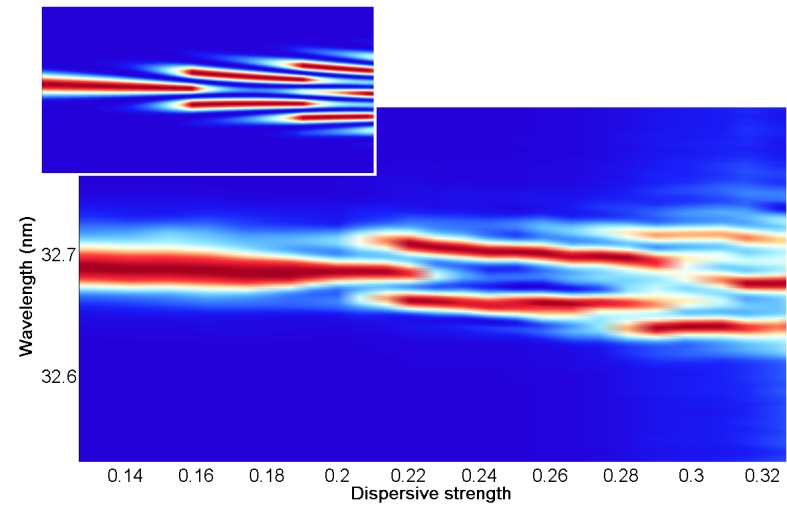


EXPERIMENTAL DEMONSTRATION

Strong chirp



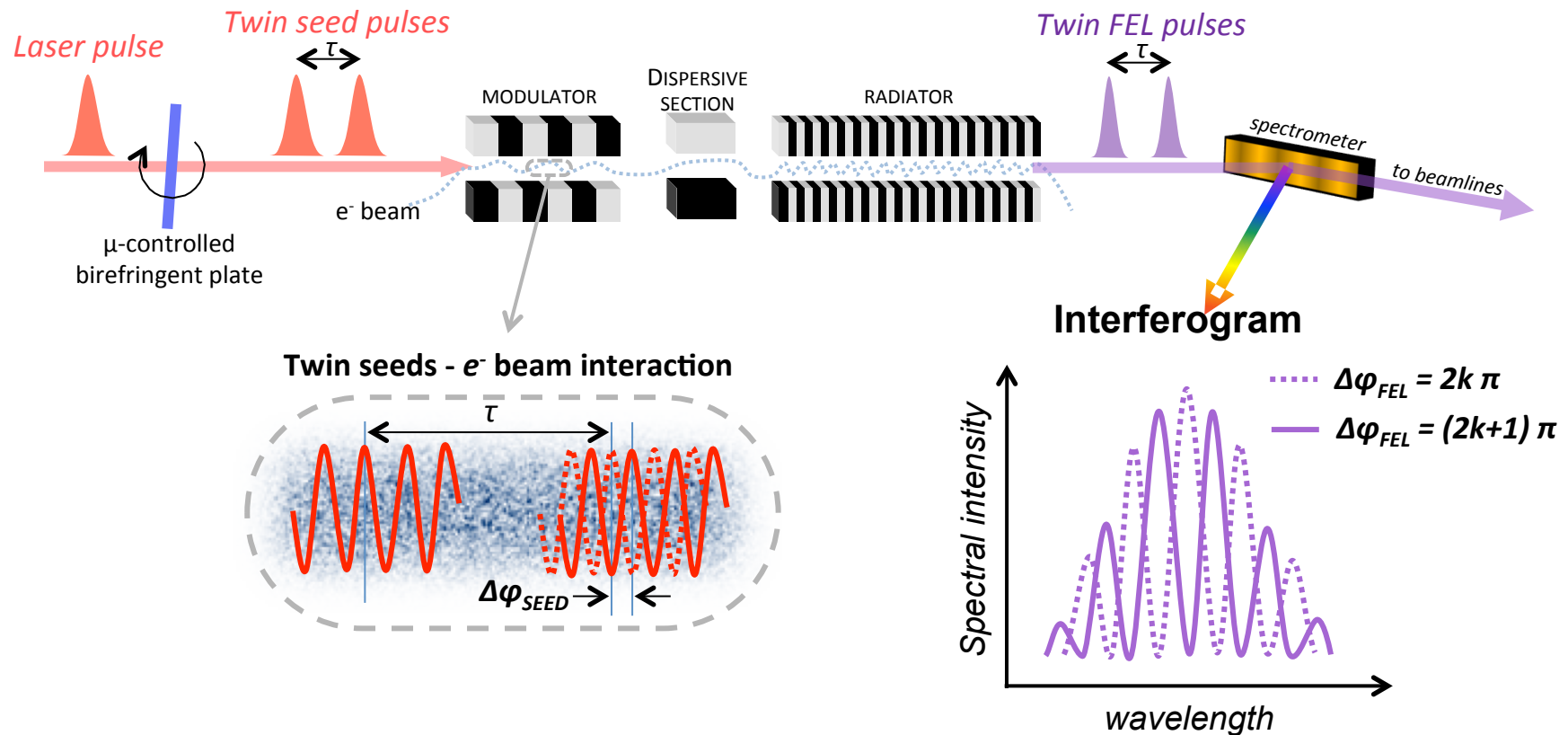
Compensated chirp



First demonstration of the possibility to generate a transform-limited FEL pulse

GENERATION OF TIME-DELAYED PHASE-LOCKED PULSES

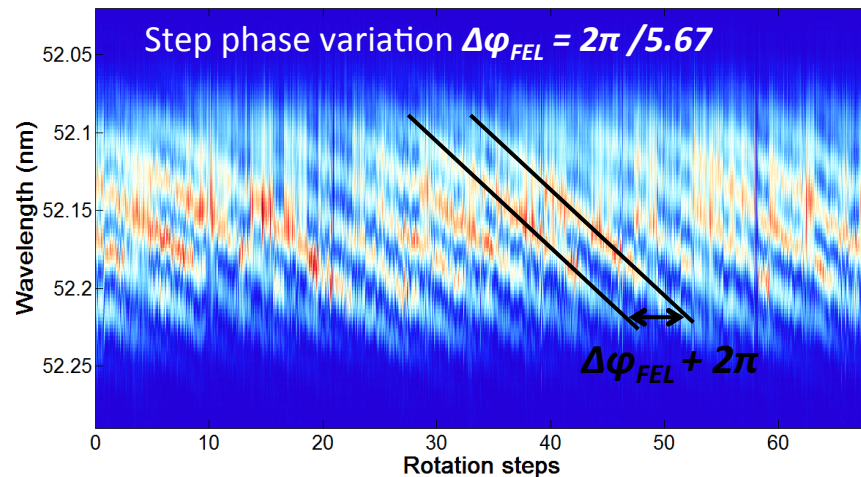
Two phase-locked seed pulses generate two phase-locked FEL pulses:



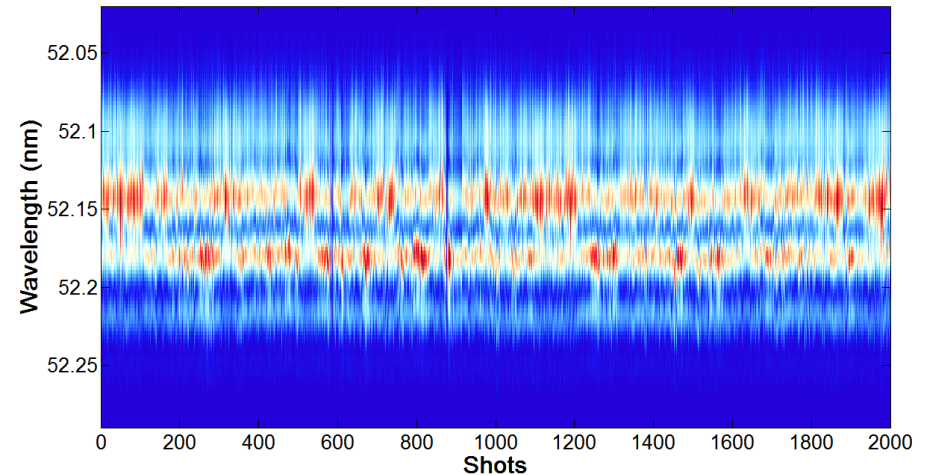
Control the phase difference between the carrier waves of the two time-delayed FEL pulses.

TIME-DELAYED PHASE-LOCKED PULSES: EXPERIMENTAL RESULTS

Interferograms vs. phase variation



Sequence of single-shot spectra



tuning of the twin-seed phase



precise control of the twin-FEL phase

phase stability: $\lambda_{FEL}/12$ RMS

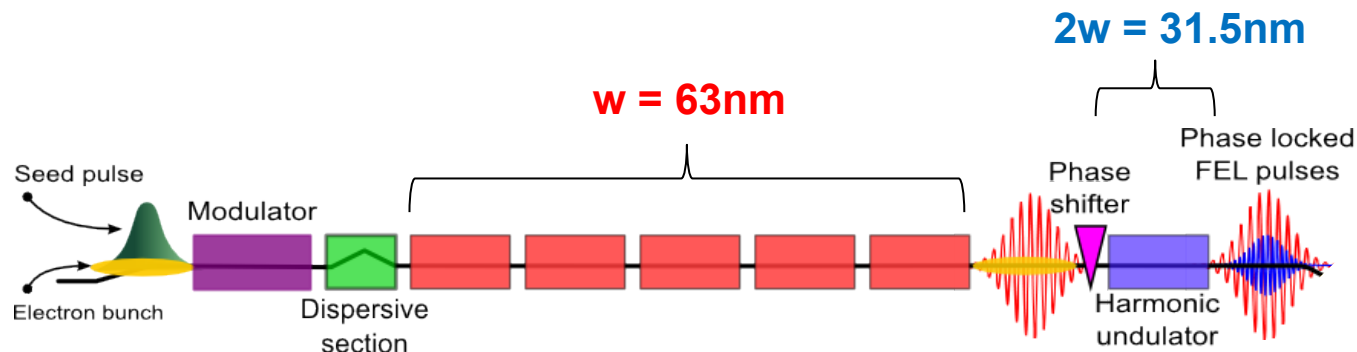


locking in phase better than **15 attoseconds**

Possible applications: nonlinear coherent transient interferometry and spectroscopy, spectral holography, quantum state holography, highly resolved spectroscopy, ...

“ZERO-DELAY” PHASE-LOCKED PULSES FOR COHERENT CONTROL

Phase locking between two harmonics of the seed, controlled by means of an electron phase shifter.

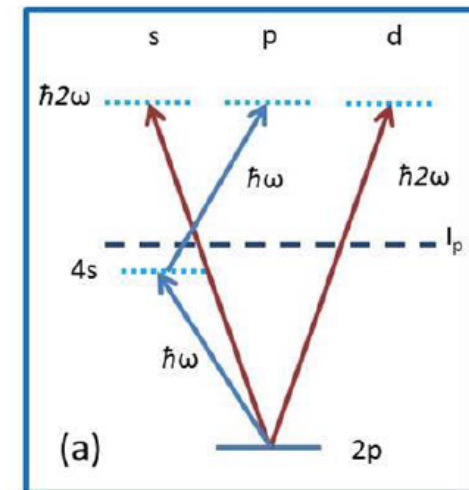


Proof-of-principle experiment:

Two-path quantum interference experiment (Brumer-Shapiro).

Interferences between 2 pathways for Ne ionization:

Ionization of Ne with 1 photon at $2w$ **vs.** 2 photons at w .

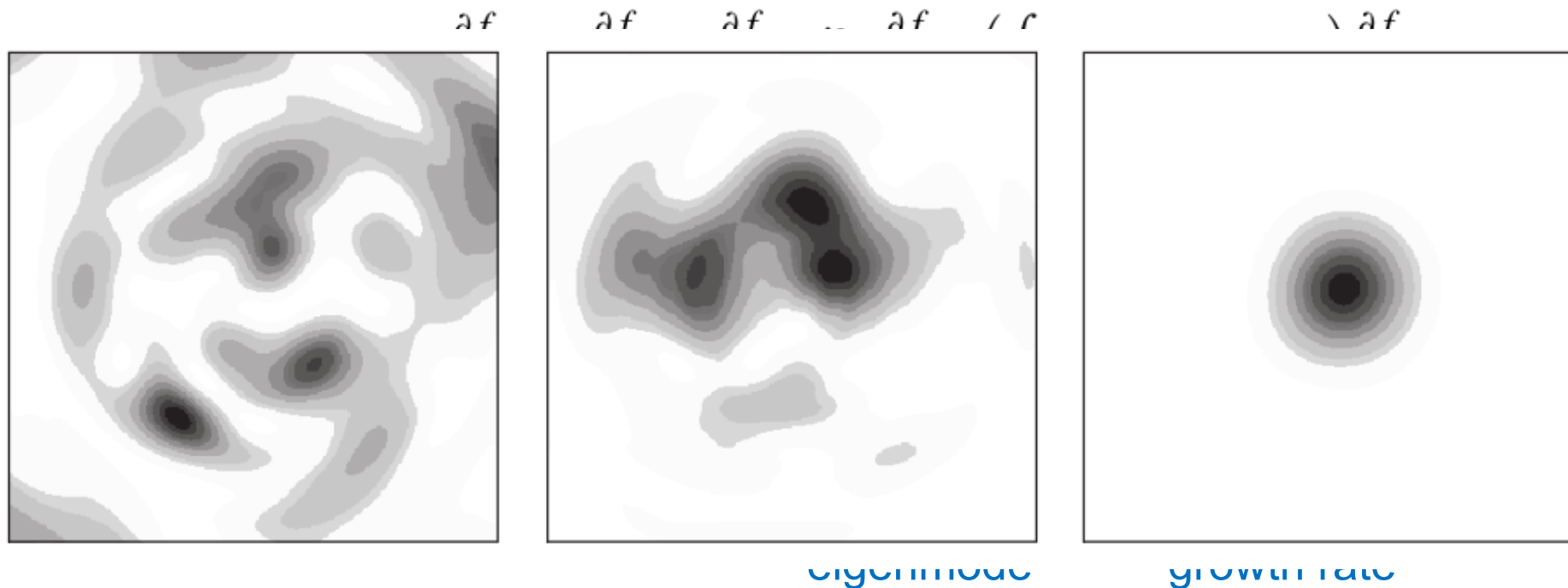


THERE ARE MANY MORE THINGS YOU
CAN DO IN THE TEMPORAL DOMAIN BY
USING AN EXTERNAL SEED TO TRIGGER
THE FEL EMISSION...

BUT LET'S SWITCH TO THE TRANSVERSE
PLANE...

FUL DESCRIPTION OF THE FEL RADIATION MECHANISM

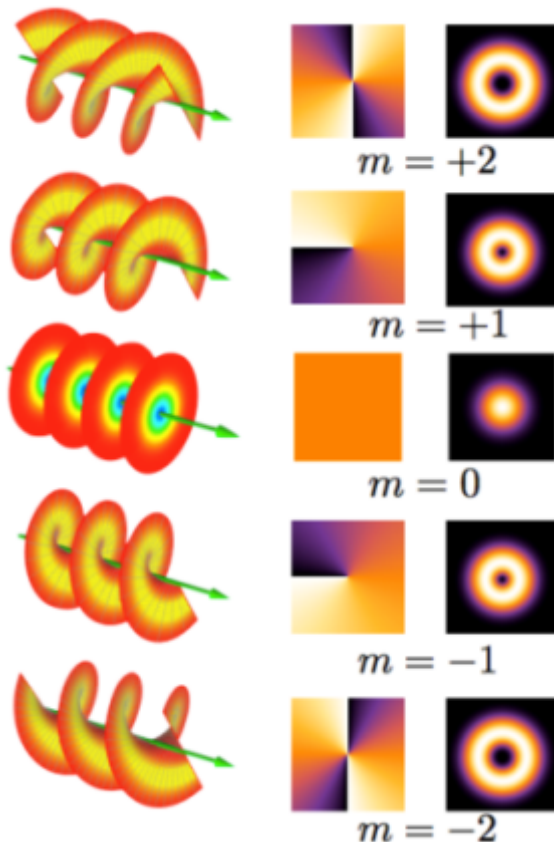
3D FEL theory:
$$\left(\frac{\partial}{\partial \hat{z}} + i\Delta \hat{v} + \frac{\hat{\nabla}_{\perp}^2}{2i} \right) a_{\nu}(\hat{\mathbf{x}}; \hat{z}) = - \int \frac{2\rho d\theta}{2\pi} e^{-i\nu\theta} \int d^2 \hat{p} \int d\hat{\eta} f(\theta, \hat{\eta}, \hat{\mathbf{x}}, \hat{\mathbf{p}}; \hat{s}),$$



At saturation the fundamental (TEM_{00}) mode, which has the highest growth rate, dominates.

ORBITAL ANGULAR MOMENTUM (OAM) OF LIGHT

Optical vortices, i.e., helically phased beams with a field dependence $E \propto e^{im\phi}$, carry orbital angular momentum*



Classically: $L = \epsilon_0 \int \mathbf{r} \times (\mathbf{E} \times \mathbf{B}) d\mathbf{r}$

Analogy with quantum mechanics:

$$L_z = -i\hbar \frac{\partial}{\partial \phi}$$

$$\text{If } \psi \propto e^{im\phi} \Rightarrow \langle L_z \rangle = \hbar m$$

SO, WHAT CAN WE DO WITH OPTICAL VORTICES?

Visible wavelengths:

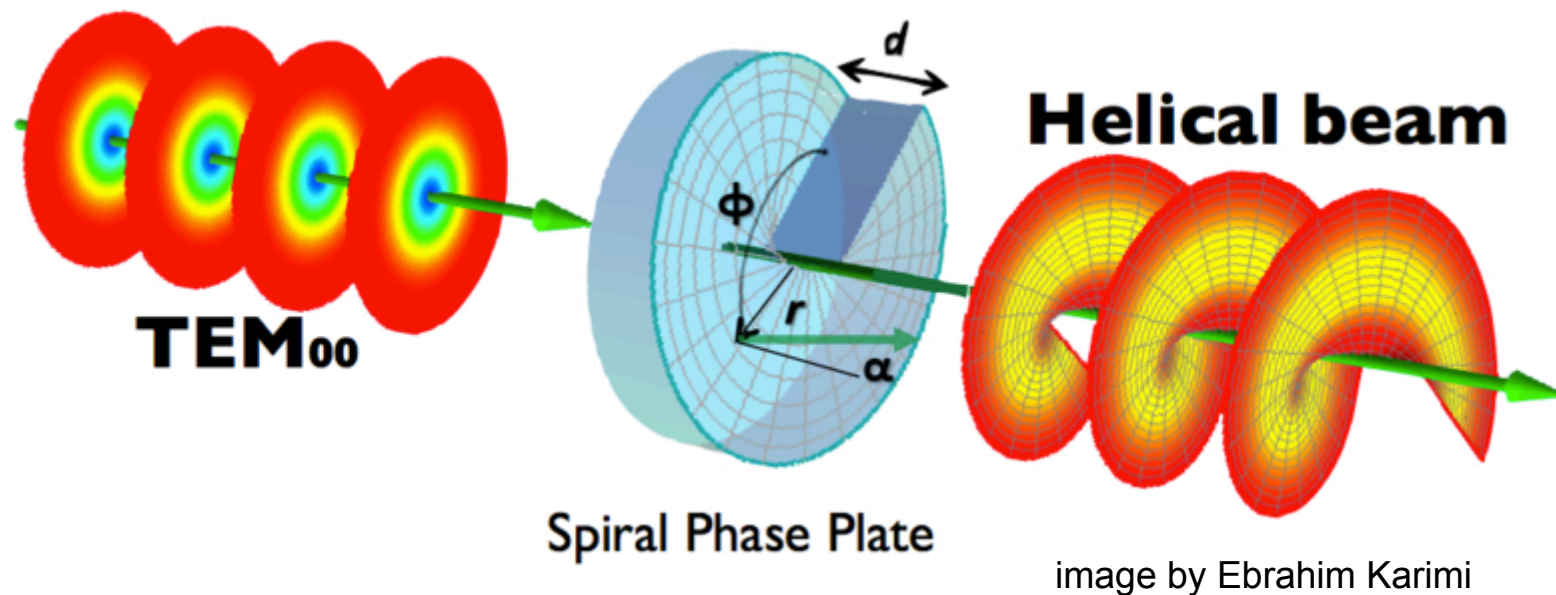
- H. He *et al.*, **Direct Observation of Transfer of Angular Momentum to Absorptive Particles from a Laser Beam with a Phase Singularity**, *Phys. Rev. Lett.* **75**, 826 (1995).
- M. P. J. Lavery *et al.*, **Detection of a Spinning Object Using Light's Orbital Angular Momentum**, *Science* **341**, 537 (2013).
- A. Jesacher *et al.*, **Shadow Effects in Spiral Phase Contrast Microscopy**, *Phys. Rev. Lett.* **94**, 233902 (2005).
- J. Wang *et al.*, **Terabit free-space data transmission employing orbital angular momentum multiplexing**, *Nature Photonics* **6**, 488 (2012).

XUV and X-rays:

- M. van Veenendaal *et al.*, **Prediction of Strong Dichroism Induced by X Rays Carrying Orbital Momentum**, *Phys. Rev. Lett.* **98**, 157401 (2007).
- A. Picón *et al.*, **Photoionization with orbital angular momentum beams**, *Opt. Express* **18**, 3660 (2010).
- A. S. Rury *et al.*, **Examining resonant inelastic spontaneous scattering of classical Laguerre-Gauss beams from molecules**, *Phys. Rev. A* **87**, 043408 (2013).

HOW CAN WE GENERATE OAM?

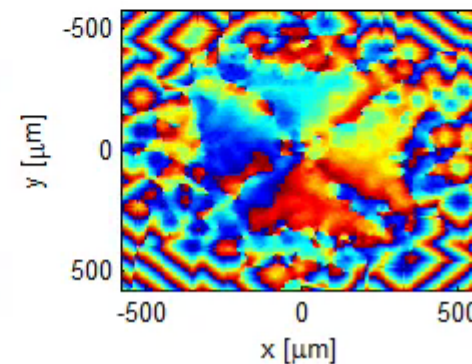
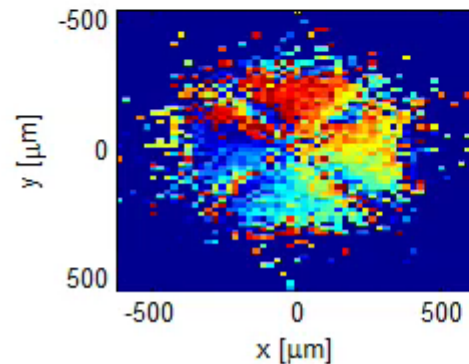
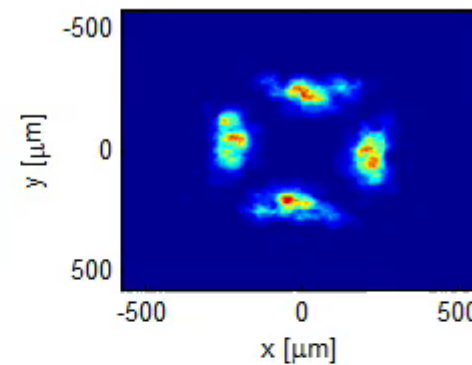
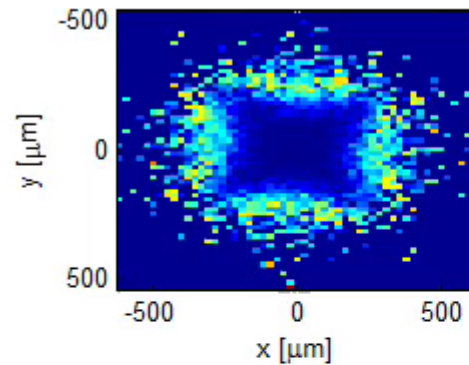
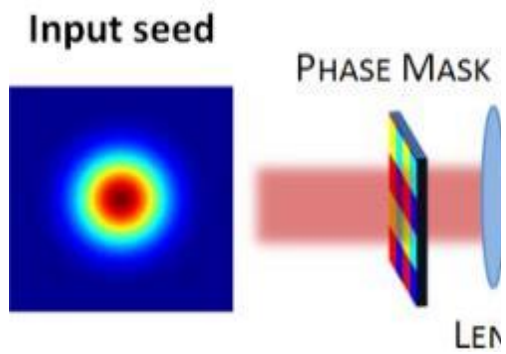
Using optical elements, e.g., a spiral phase plate:



Not practical at XUV and X-ray wavelengths and FEL intensities \Rightarrow *in situ* generation preferred

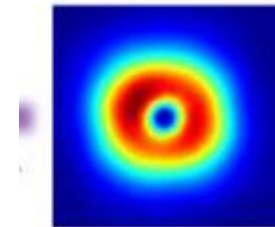
WHAT CAN WE DO AT FERMI?

Use a phase-mask

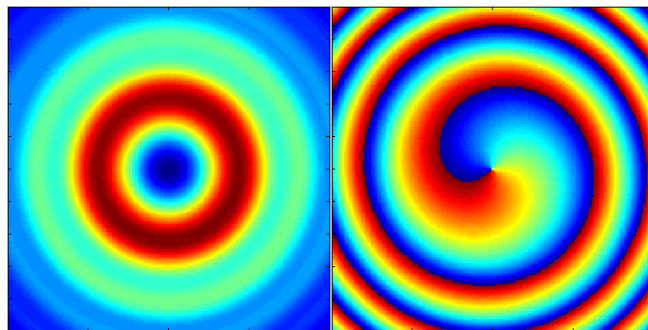


:

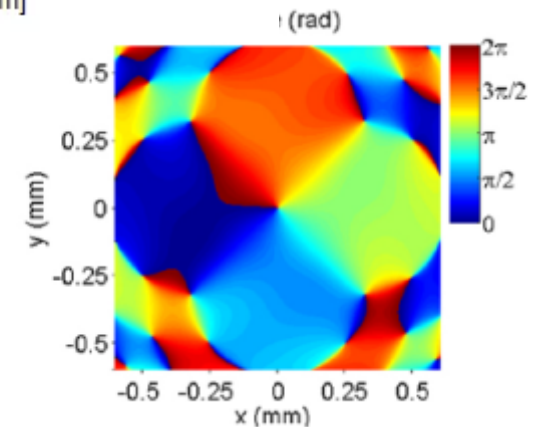
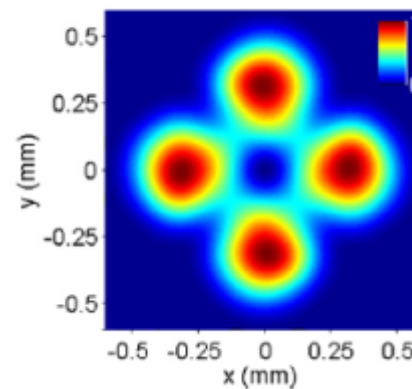
Output FEL



Use a spiral phase mask? It doesn't work

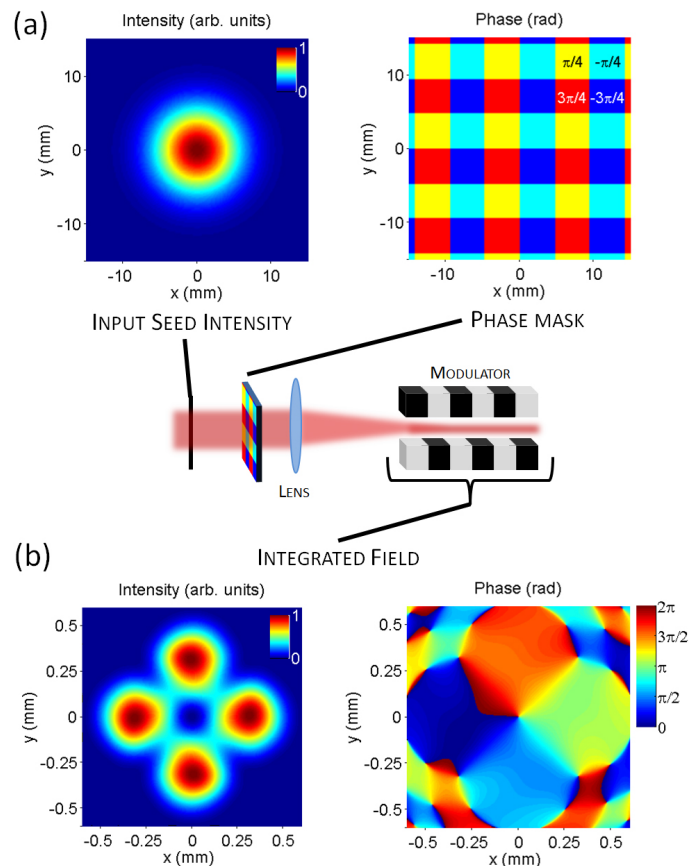


-like phase

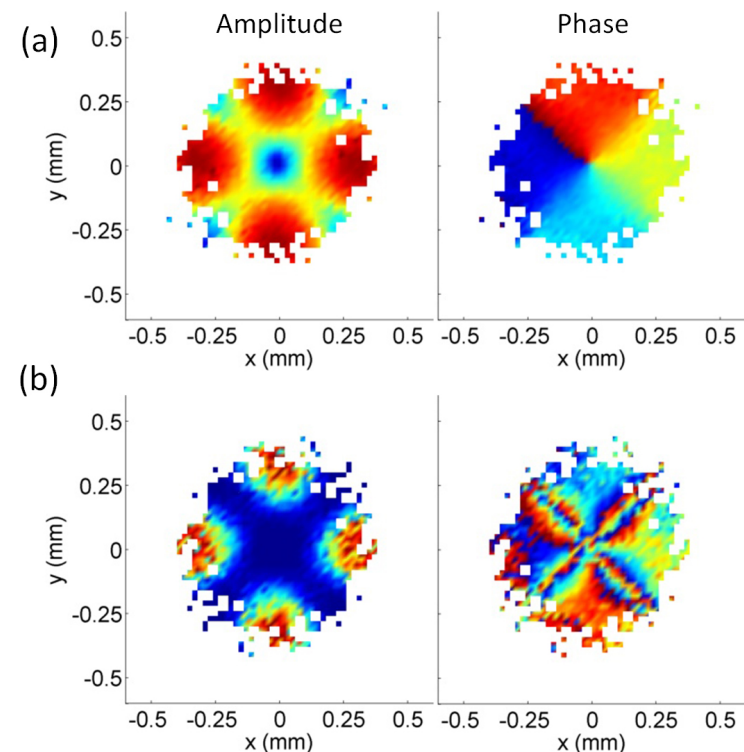


MICROBUNCHING CONSTRUCTION IN THE MODULATOR

Modification of the transverse seed profile using a phase mask:

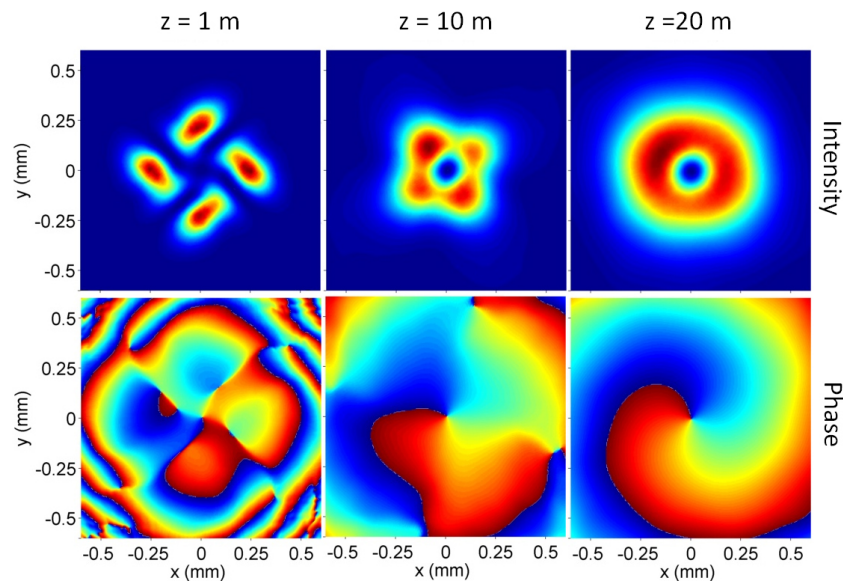


Formation of microbunching in the modulator: transverse profile at a) the fundamental (260 nm) and b) 7th harmonic (37 nm)

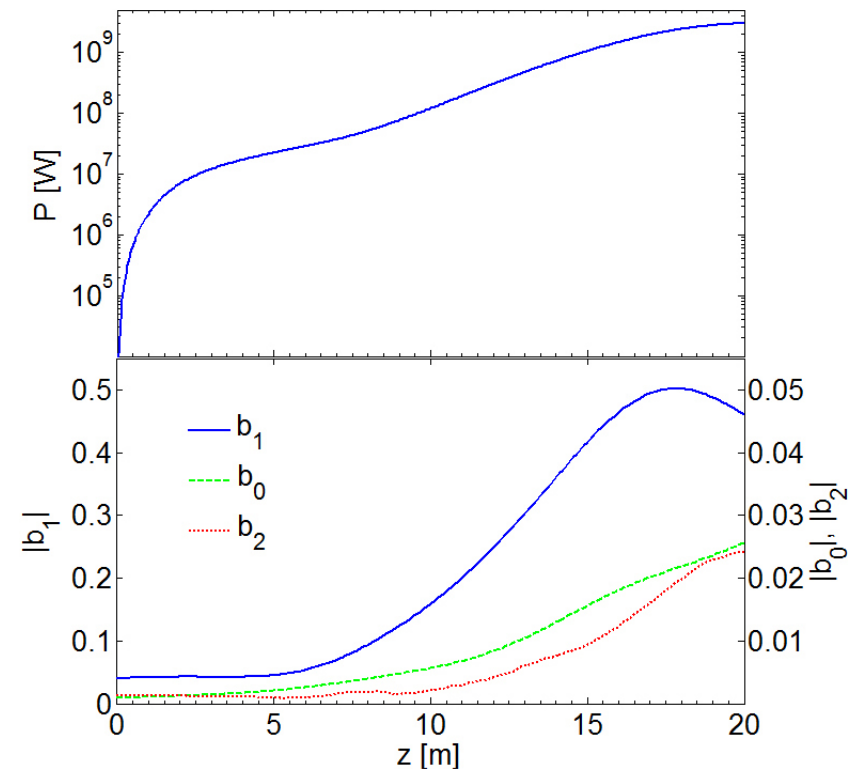


EVOLUTION OF THE RADIATION PROFILE

Transverse radiation profile in the undulator (7th harmonic)

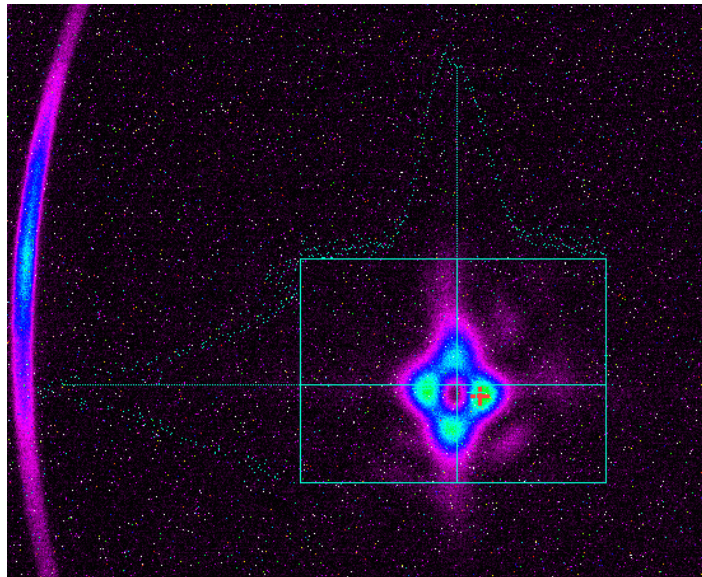


Evolution of the FEL power and bunching factors ($I_{\text{beam}} = 1$ kA, $P_{\text{seed}} = 1$ GW, normalized emittance = 5×10^{-6} m)

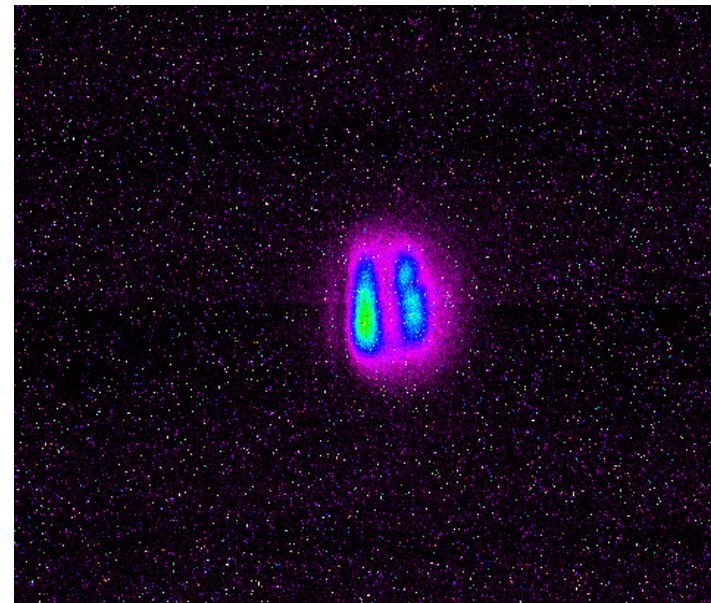


EXPERIMENTAL (ALMOST) DEMONSTRATION

seed transverse intensity profile



FEL intensity profile



Shaping FEL radiation in the transverse plane is much more difficult compared to shaping in the temporal domain!

WHAT HAVE WE JUST LEARNED?

- compared to synchrotrons FELs produce more powerful (orders of magnitude higher peak brilliance) and shorter (femtosecond) pulses with laser-like properties
- self seeding and HGHG improve FEL performance (spectral brightness, central wavelength and pulse energy stability)
- different schemes for SASE and seeded FELs can deliver two color pulses with tunable properties for pump-probe experiments
- HGHG offers full control over the spectro-temporal and spatial properties of FEL light

Acknowledgements:

Giovanni De Ninno, Elettra/UNG

David Gauthier, Elettra

FERMI comissioning team

Giorgio Margaritondo, EPFL