

X-ray Photon Correlation Spectroscopy (XPCS)

at Synchrotron and FEL sources

Christian Gutt

Department of Physics, University of Siegen, Germany

gutt@physik.uni-siegen.de

Outline

- How to measure dynamics in condensed matter systems
- Coherence
- X-ray speckle patterns
- How to exploit X-ray intensity fluctuations
- Examples for slow dynamics
- XPCS at FEL sources

Wiley Classics Library

GOODMAN

Statistical Optics

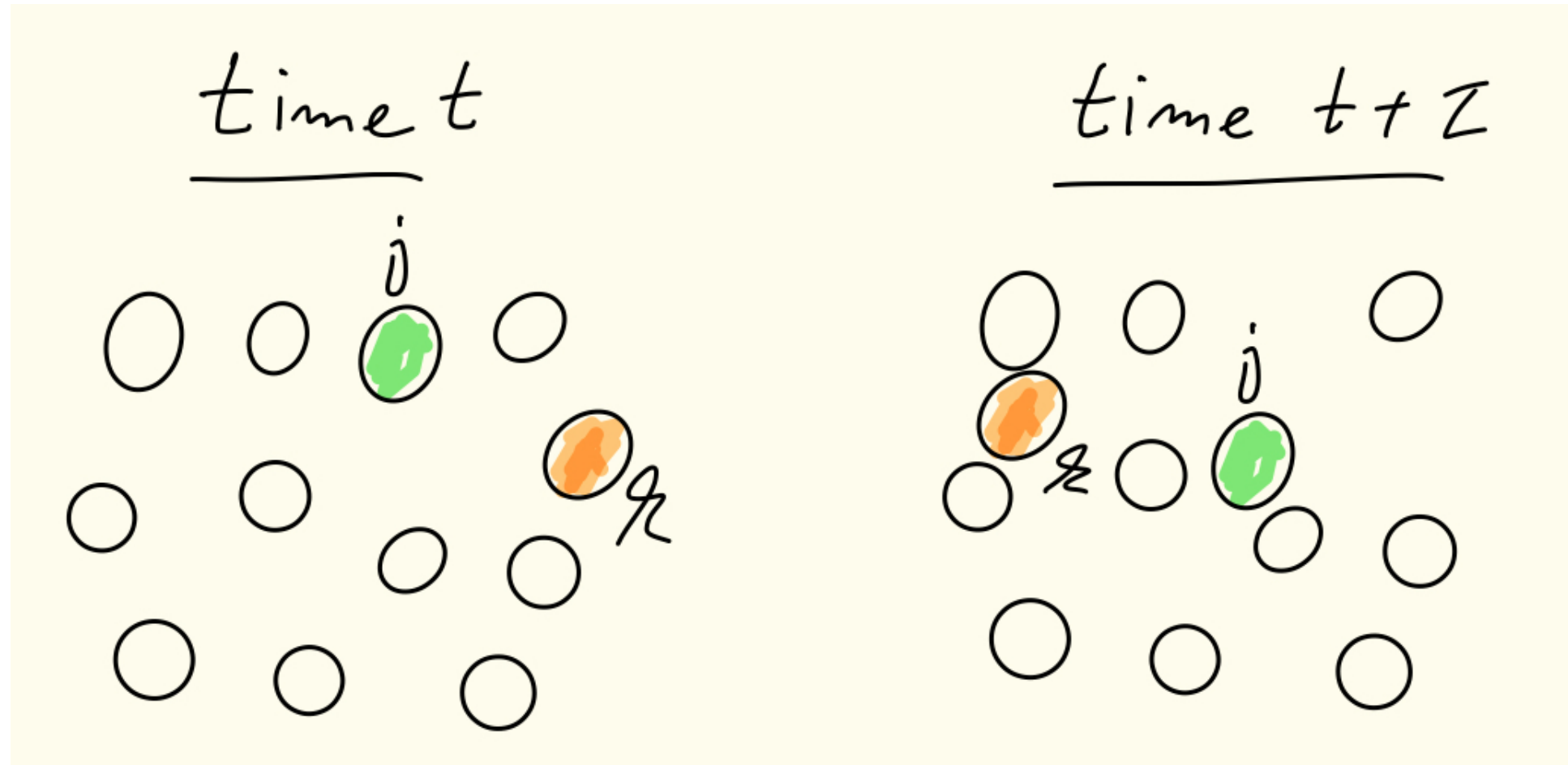


**SPECKLE
PHENOMENA
IN OPTICS**

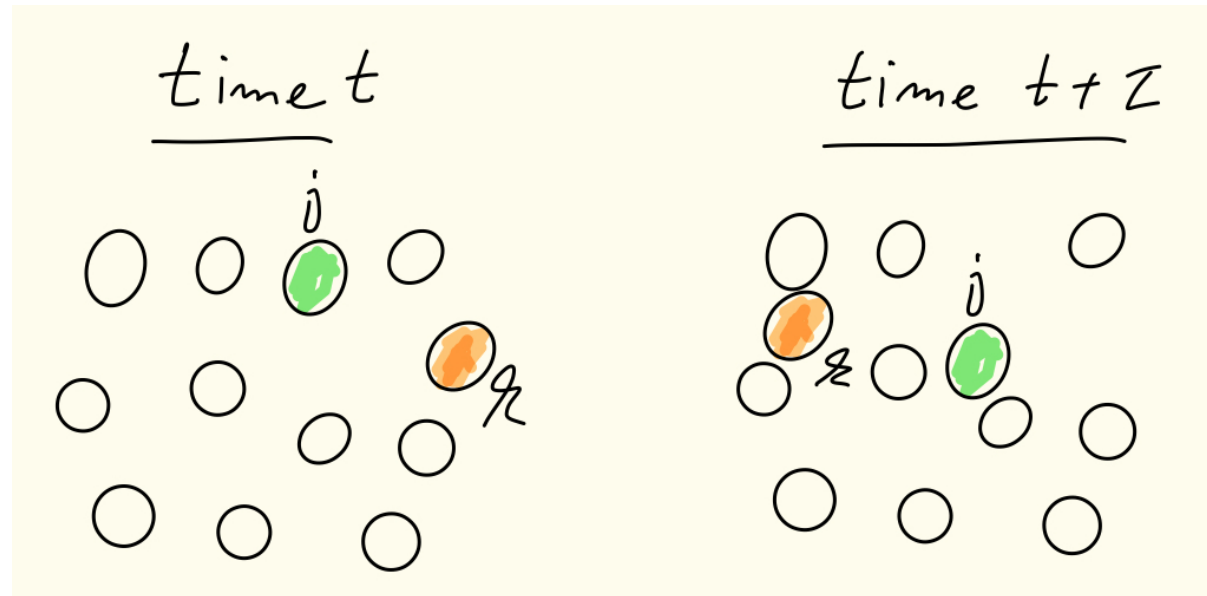
THEORY AND APPLICATIONS

JOSEPH W. GOODMAN

How to measure dynamics in condensed matter systems



How to measure dynamics in condensed matter systems



$$F(Q, \tau) = \frac{1}{N} \sum_j \sum_k \exp(iQ(r_j(t) - r_k(t + \tau)))$$

Time domain

intermediate scattering function

$$S(Q, \omega) = \int F(Q, \tau) \exp(i\omega\tau) d\tau$$

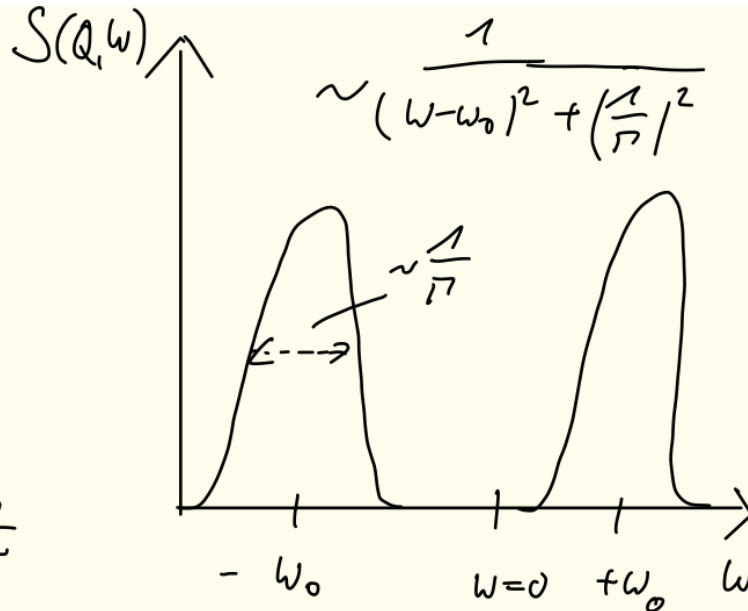
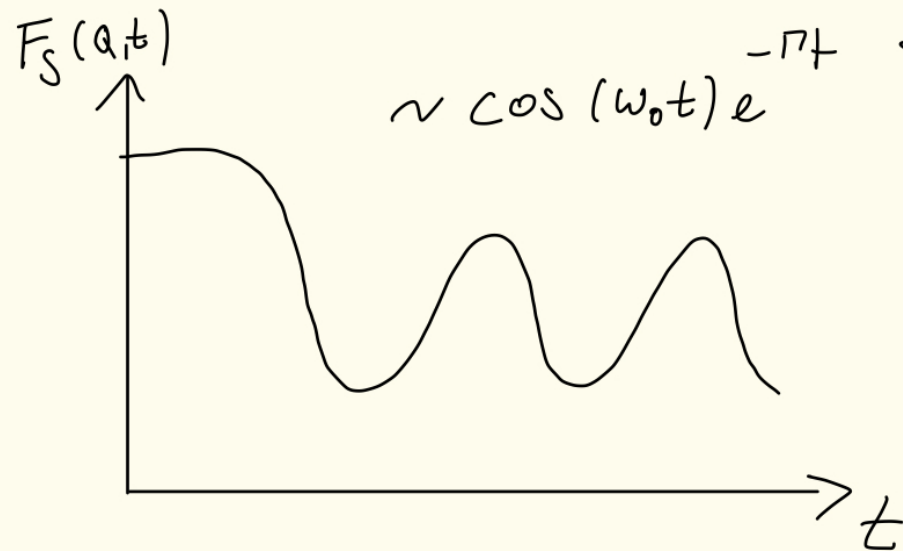
Frequency domain

dynamic structure factor

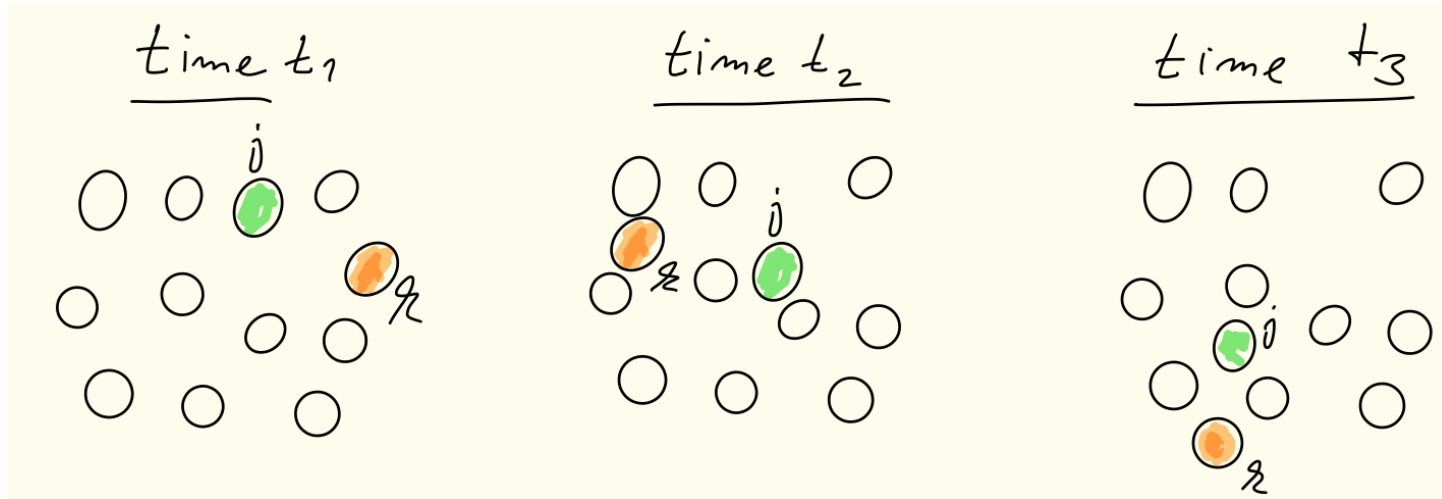
Elastic processes – waves, phonons...



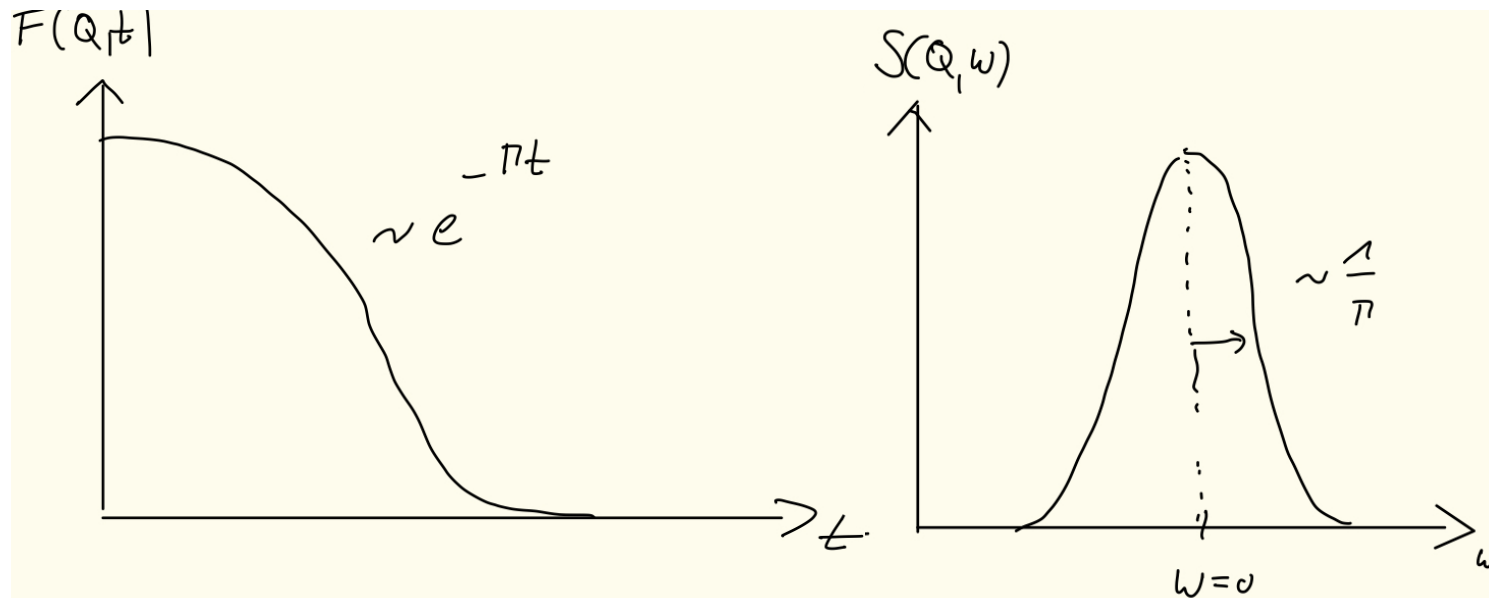
Restoring force – the system goes back to its previous configuration

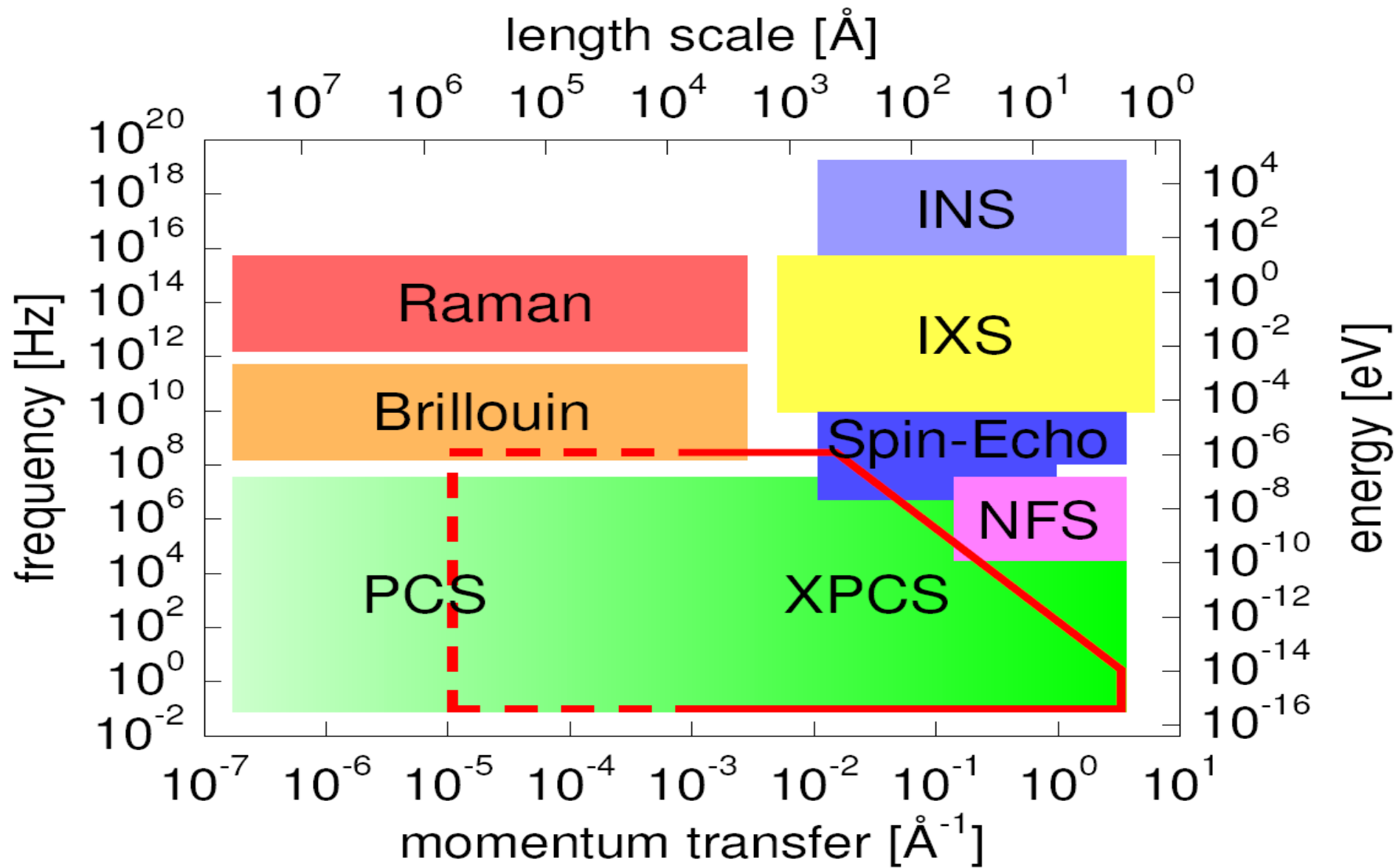


Relaxational processes – diffusion, viscosity...



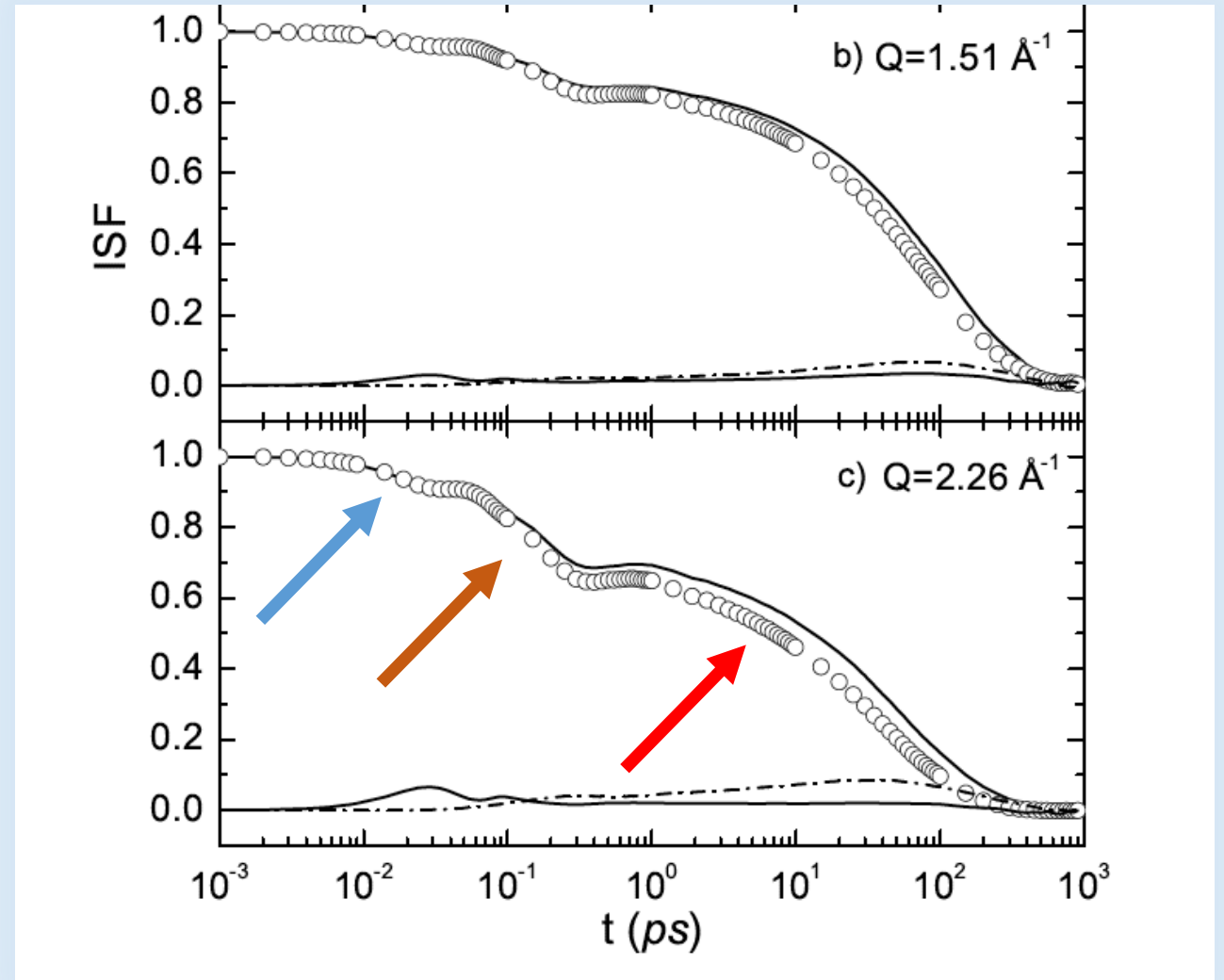
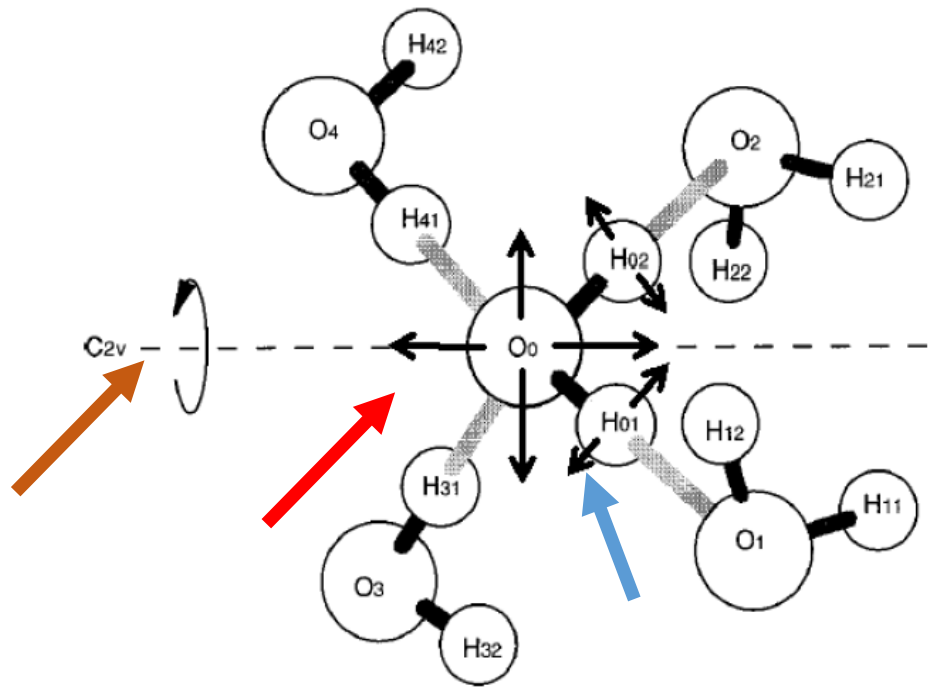
**No restoring force –
the system evolves with time
and does not come back**



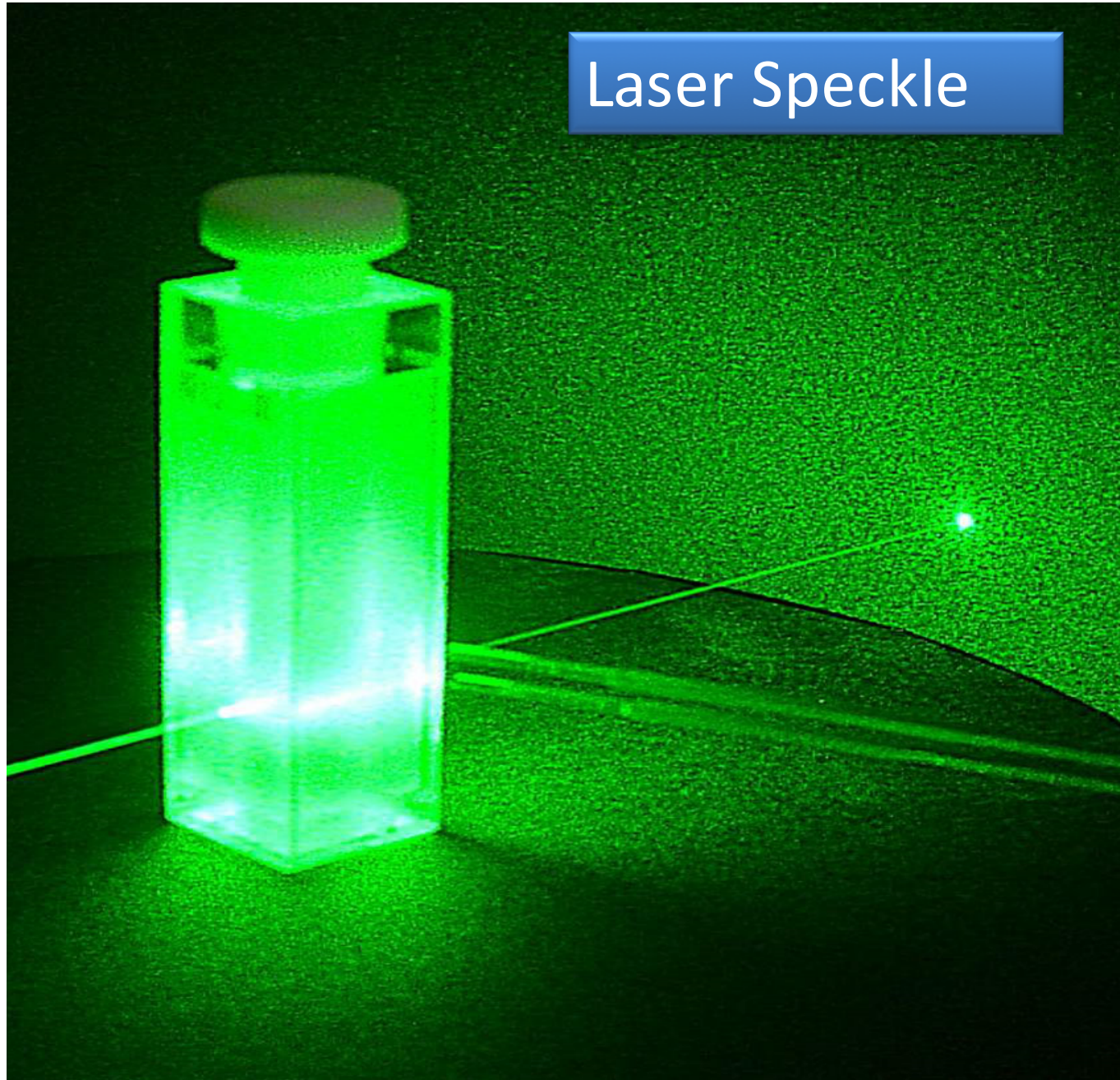


An example – molecular dynamics simulation of liquid water

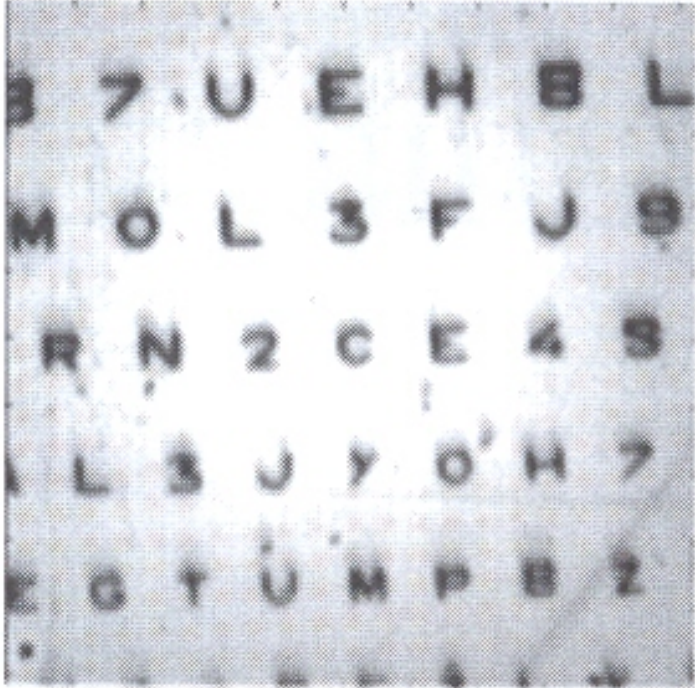
Intermediate scattering function is complex (many correlation processes) and spans many orders of magnitude
-> **experiments in the time domain**



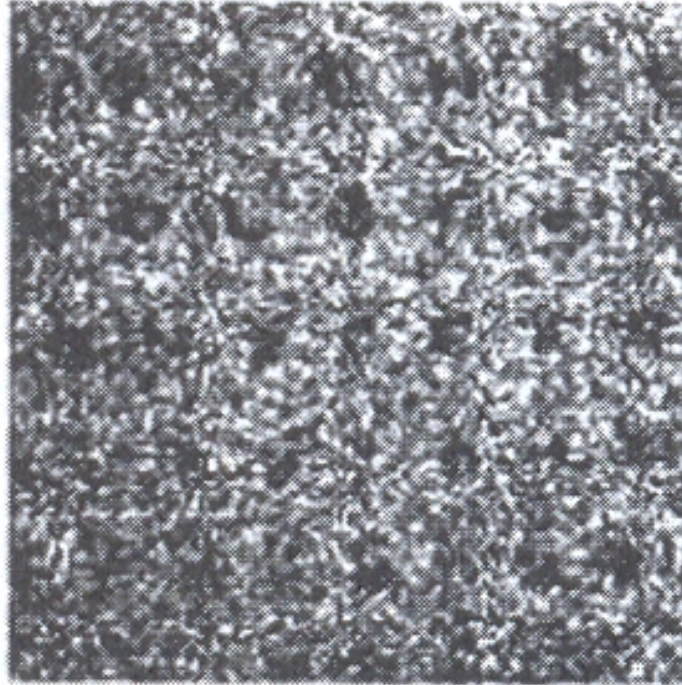
Laser Speckle



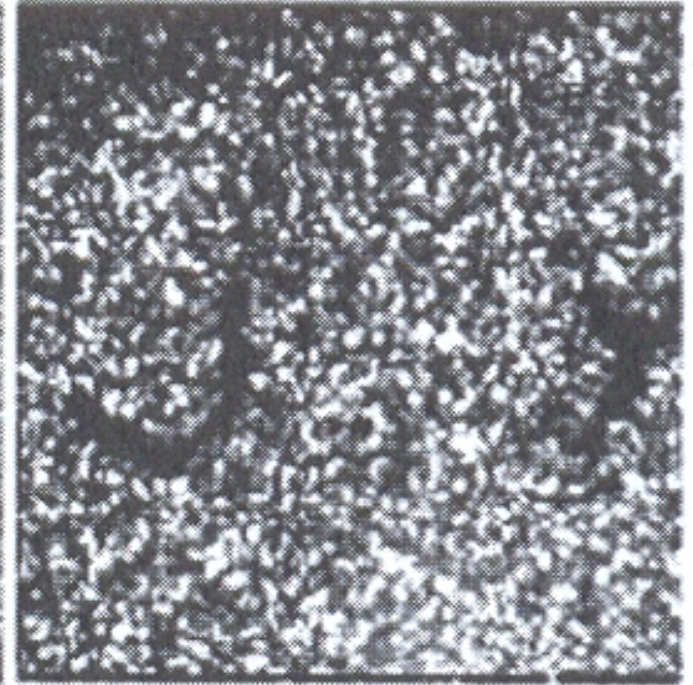
Optical Speckles



Incoherent light



Coherent light

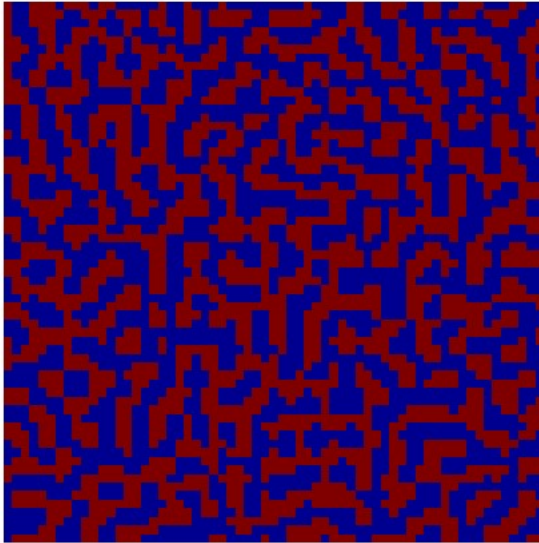


Close up

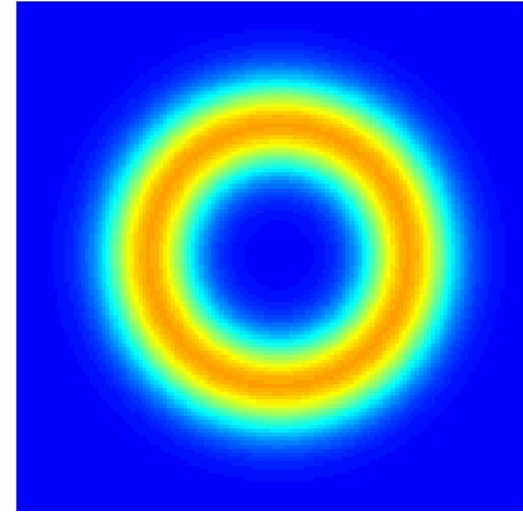
VLC movie

Coherent scattering from disorder: Speckle

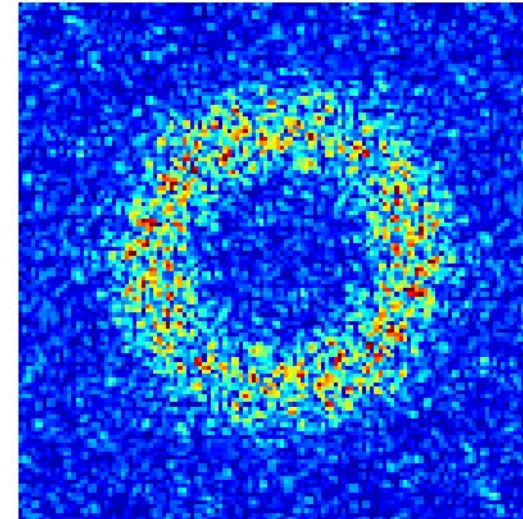
sample with disorder
(e.g. domains)



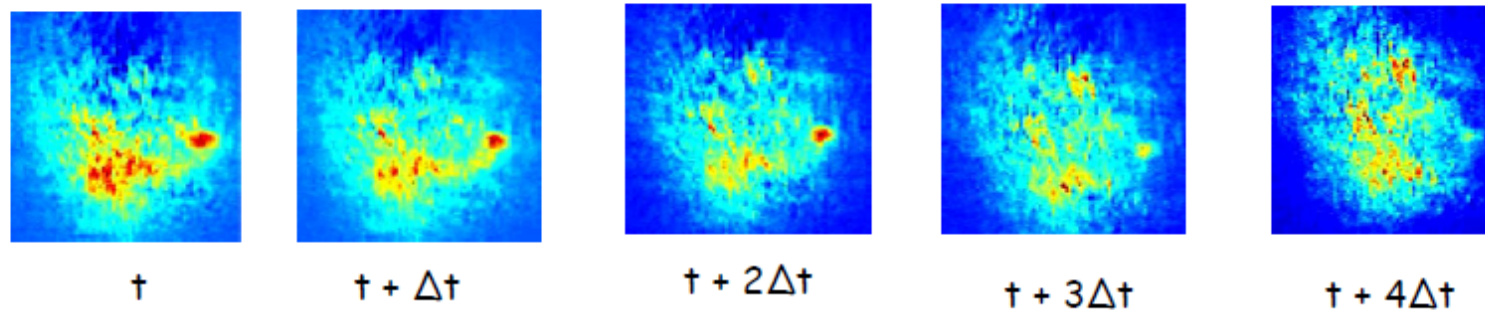
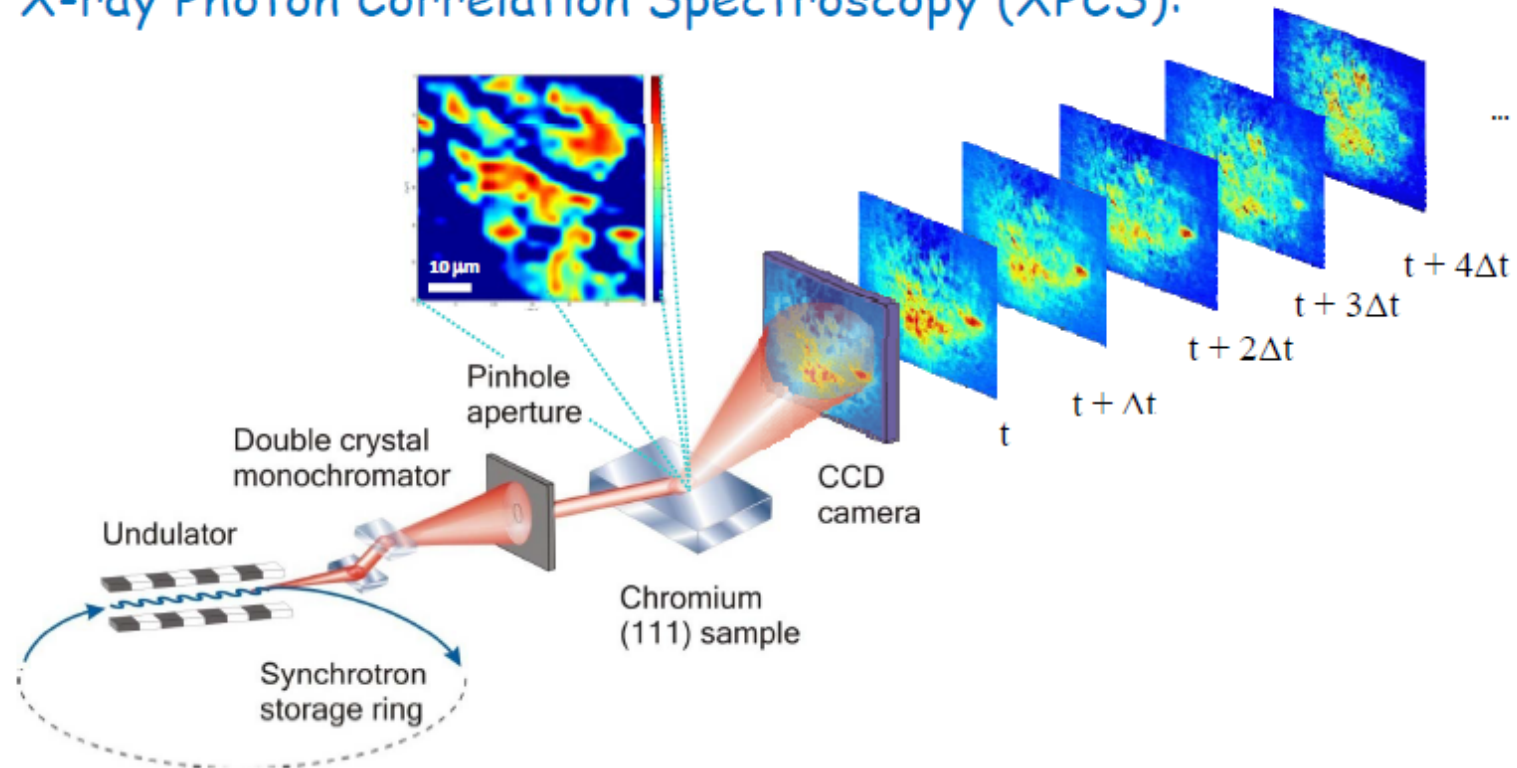
- ***Incoherent Beam: Diffuse Scattering***
 - Measures averages



- ***Coherent Beam: Speckle***
 - Speckle depends on **exact** arrangement
 - Speckel statistics encodes **coherence properties**



X-ray Photon Correlation Spectroscopy (XPCS):



XPCS – Theory

$$\langle I(t)I(t+\tau) \rangle = \langle E(t)E^*(t)E(t+\tau)E^*(t+\tau) \rangle$$

Gaussian momentum theorem

$$= \underbrace{\langle E(t)E^*(t) \rangle}_{\langle I(t) \rangle} \underbrace{\langle E(t+\tau)E^*(t+\tau) \rangle}_{\langle I(t) \rangle} + \underbrace{\left| \langle E(t)E^*(t+\tau) \rangle \right|^2}_{g_1(\tau)}$$

$$\frac{\langle I(t)I(t+\tau) \rangle}{\langle I(t) \rangle^2} = 1 + |g_1(\tau)|^2$$

XPCS Theory

$$E(t) = A \sum_{j=1}^N b_j \exp(iqr_j(t))$$

$$g_1(q, \tau) = A^2 \sum_{j,k=1}^N b_k b_j \exp(iq(r_j(t) - r_k(t + \tau)))$$

Time dependent density correlation function

Experiment

$$\frac{\langle I(t)I(t+\tau) \rangle}{\langle I(t) \rangle^2} = 1 + \beta |g_1(\tau)|^2$$

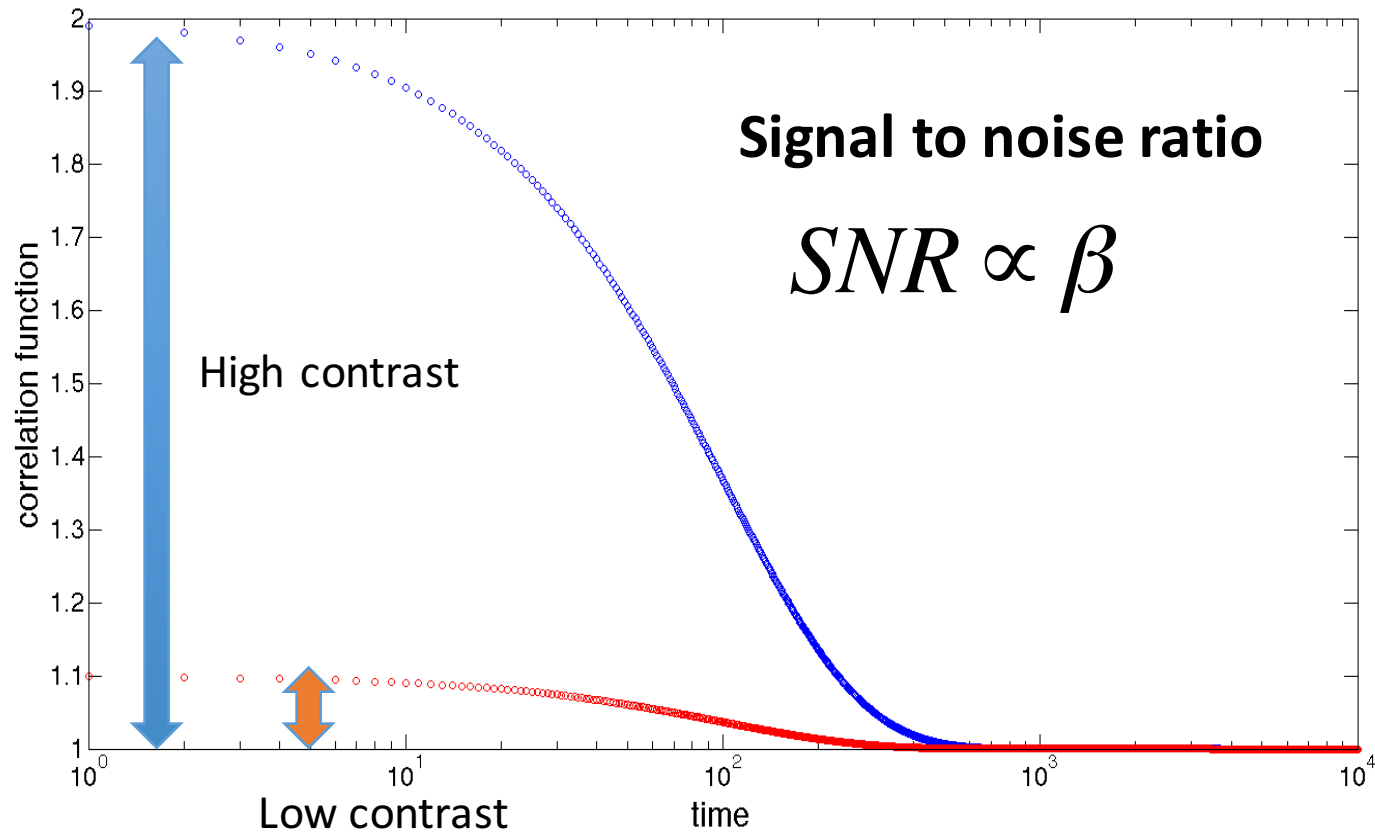
Speckle contrast < 1

Speckle blurring leads to small contrast

Partial coherence of the
x-ray source

Detector pixels P larger
than speckle size S

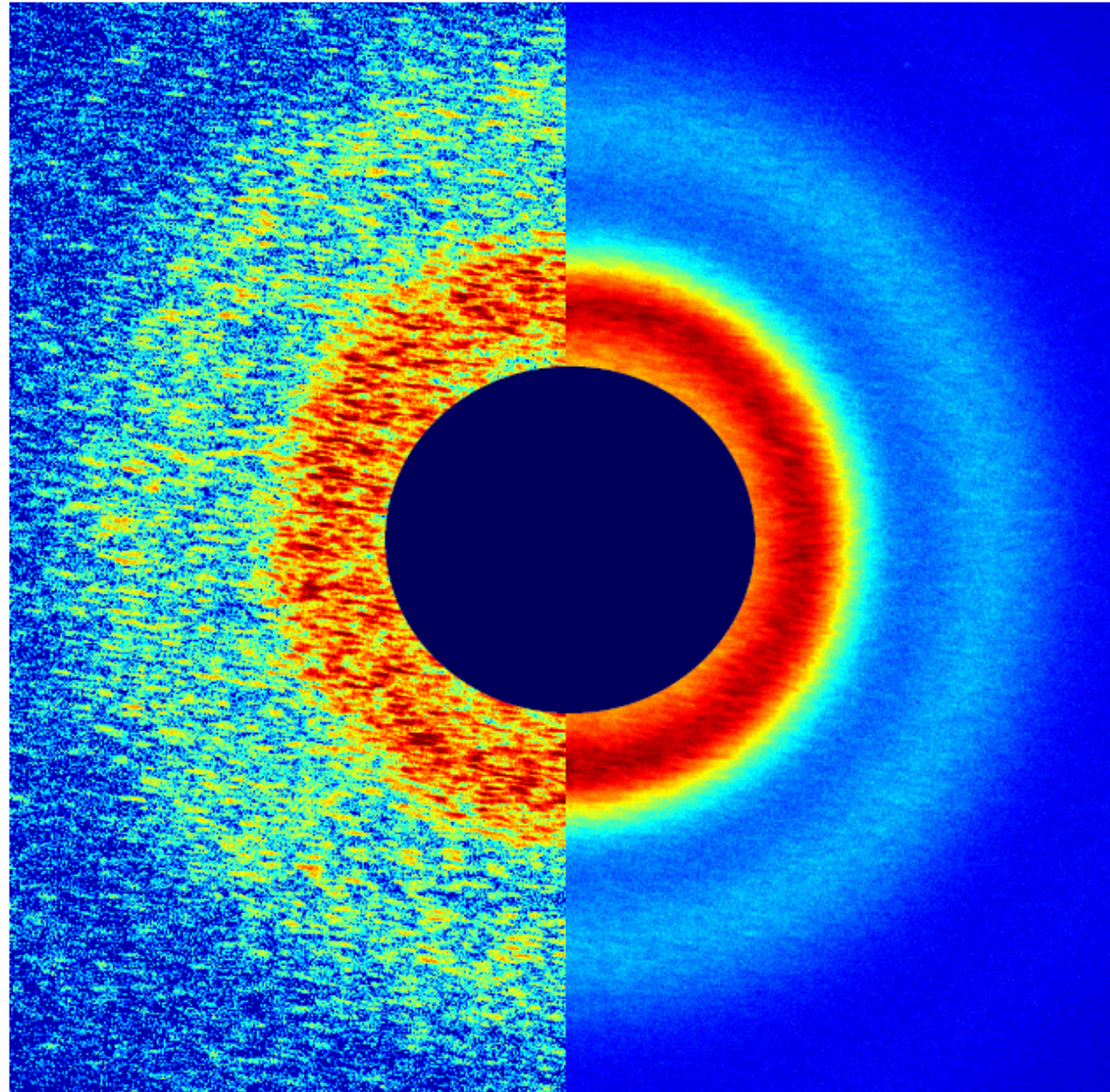
$$S \approx \frac{\lambda}{D} \times L$$

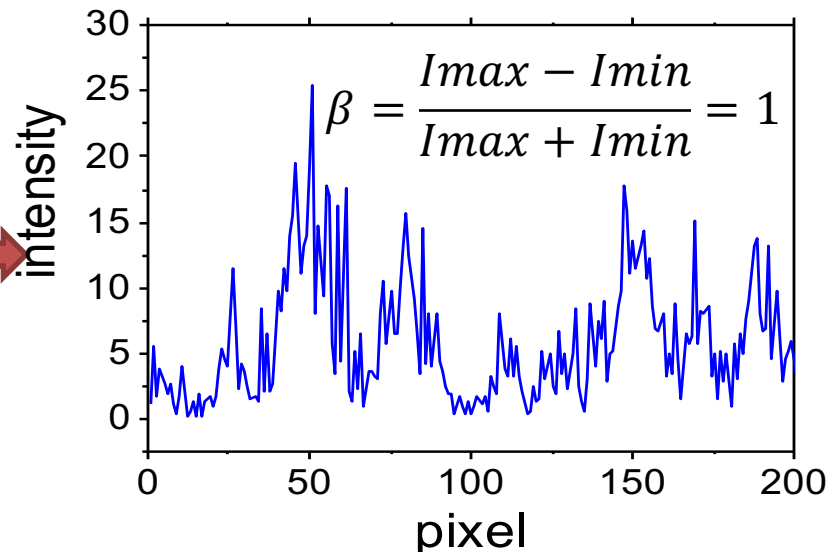
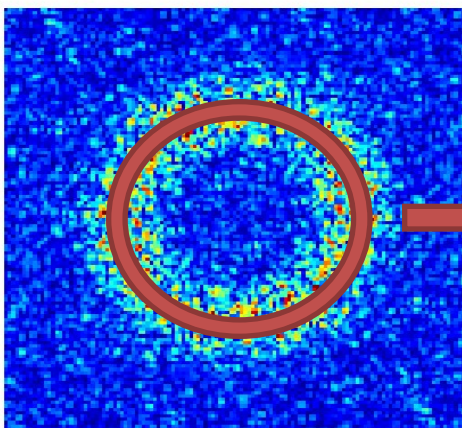
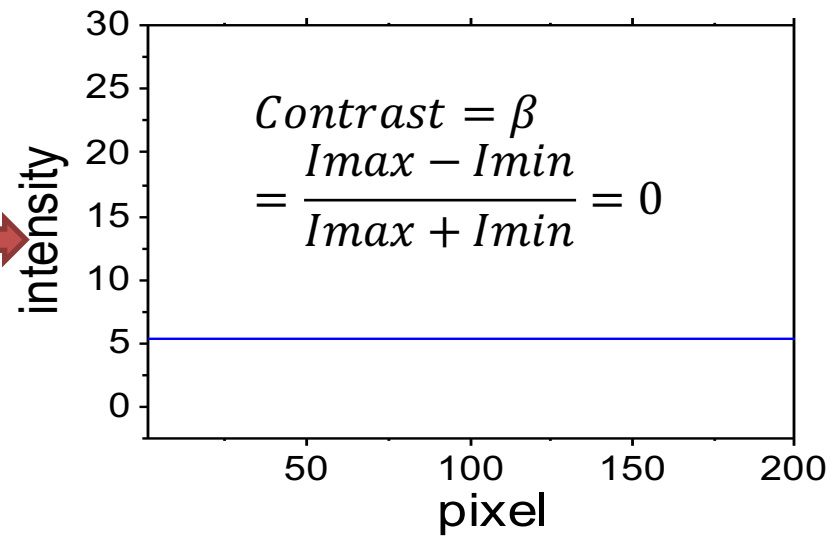
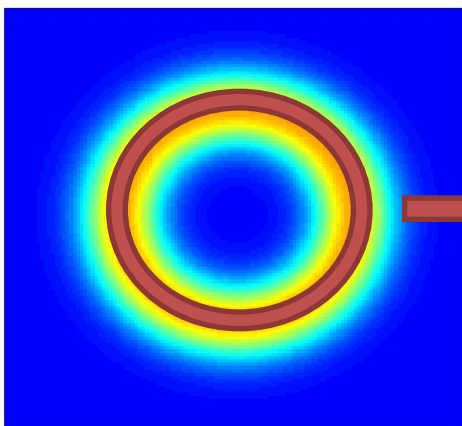


$$\frac{\langle I(t)I(t + \tau) \rangle}{\langle I(t) \rangle^2} = 1 + \beta |g_1(\tau)|^2$$

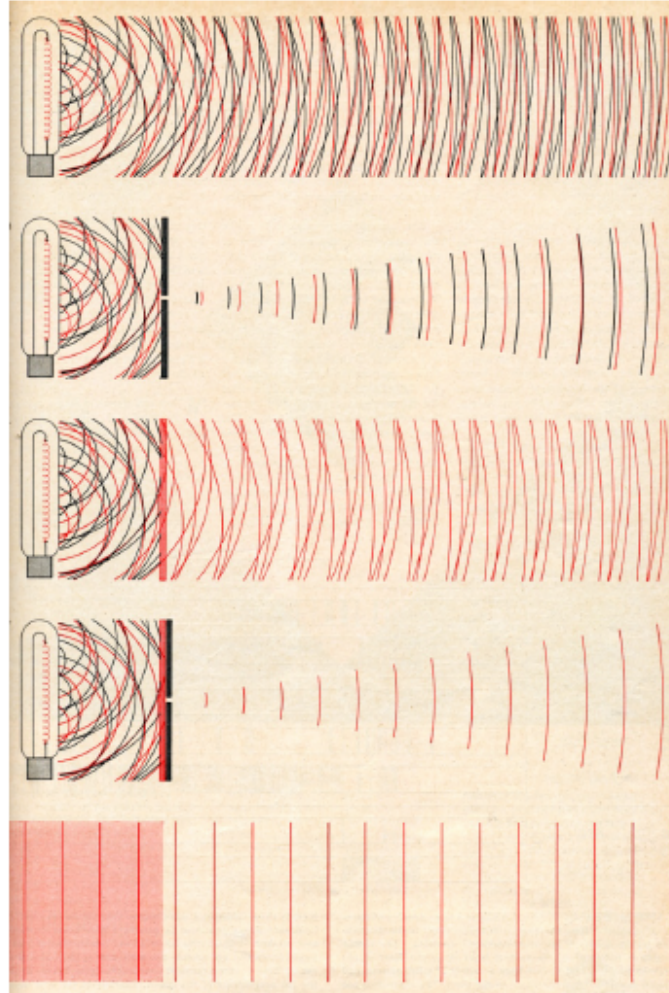
High coherence

Low coherence





Coherence

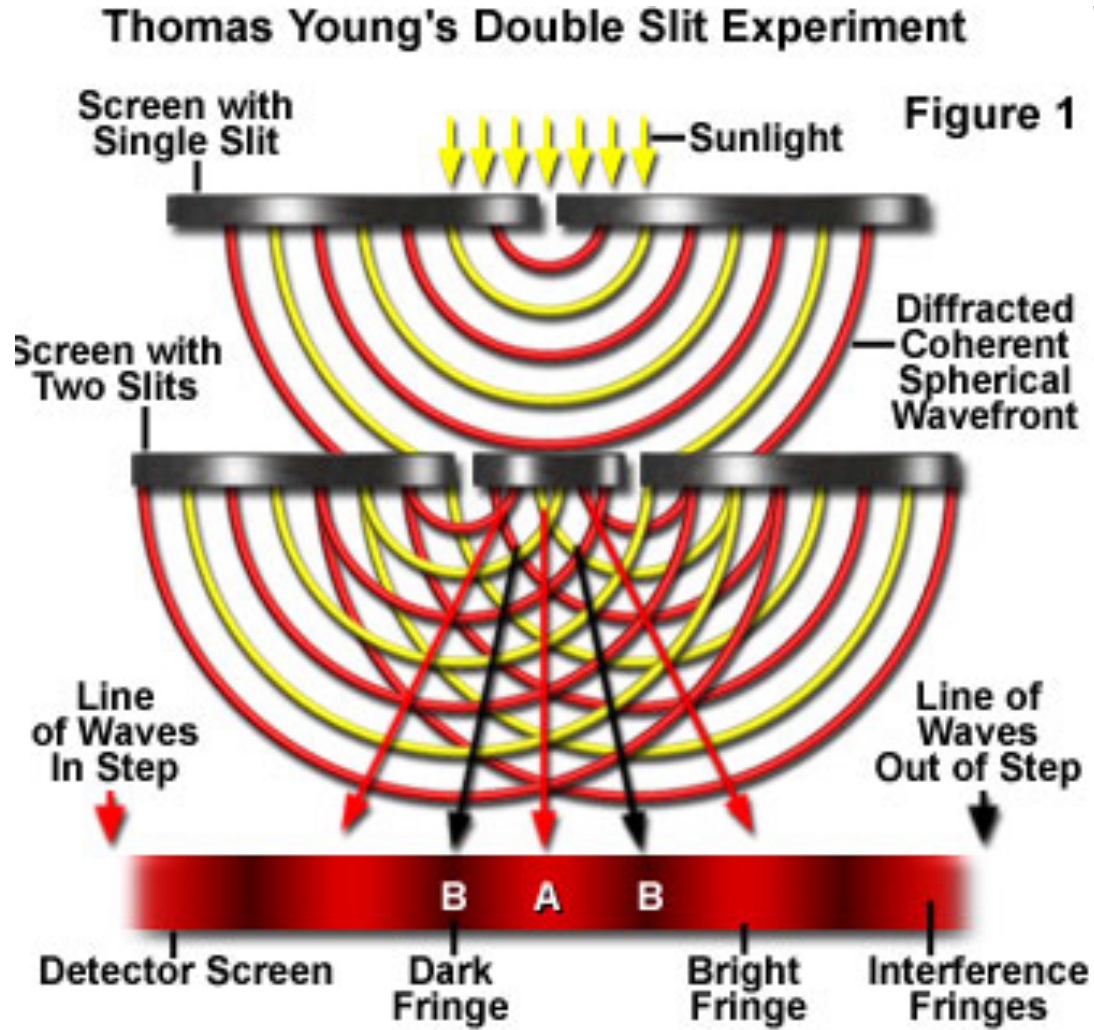


Spatial coherence

Temporal coherence

A. L. Schawlow "Laser Light"
Scientific American, 219 (3), p. 120, (1968)

Young's Double Slit Experiment



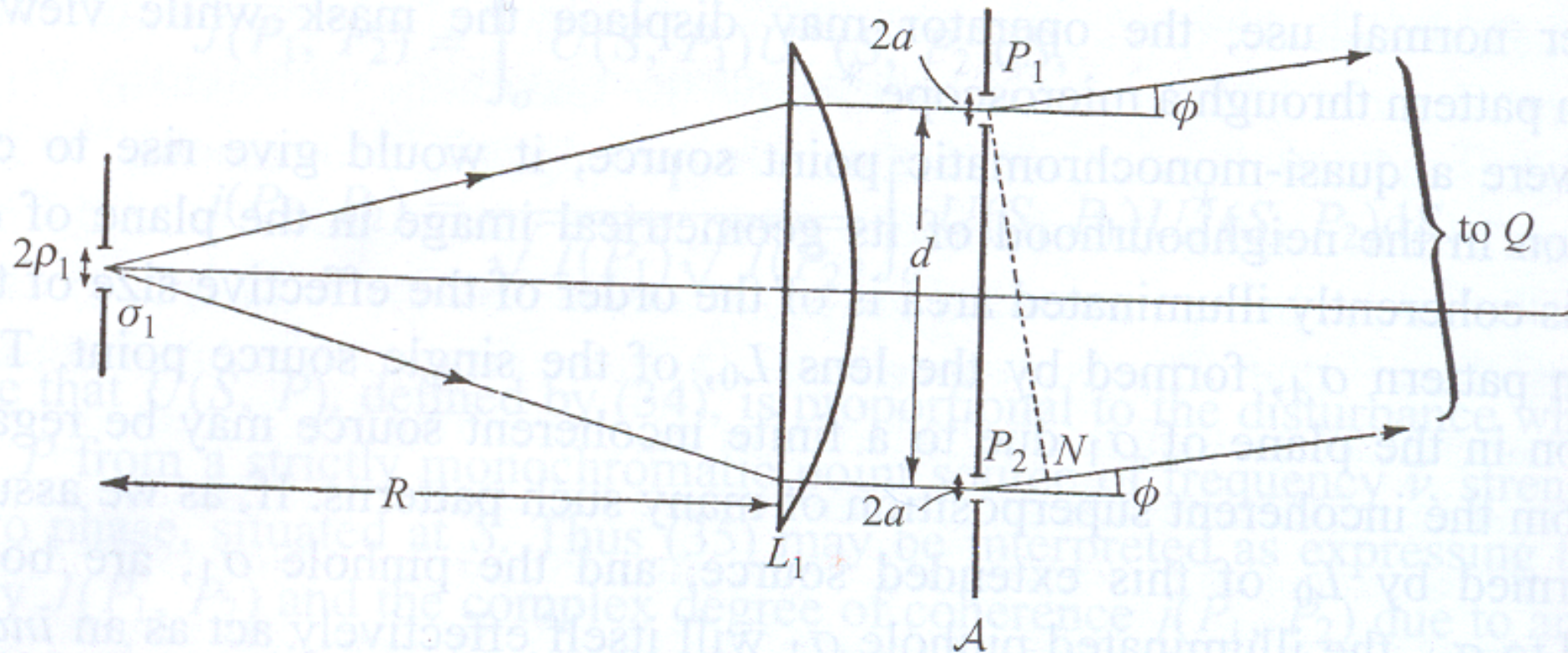
Thomas Young

Thomas Young, 1773-1829

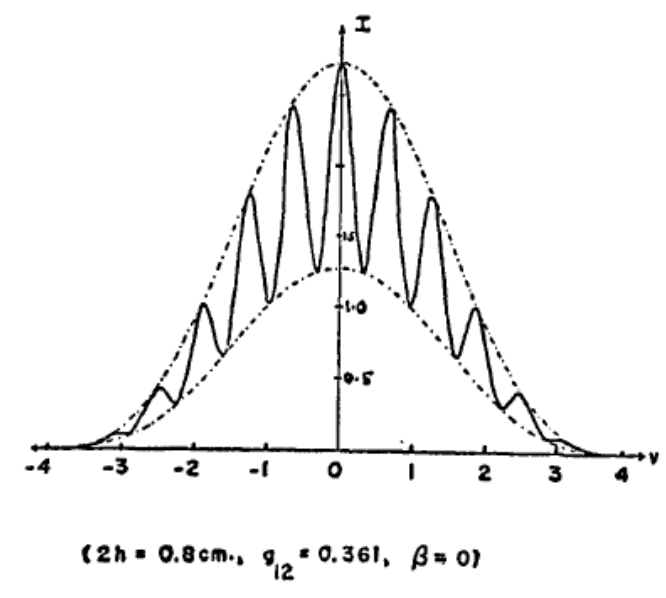
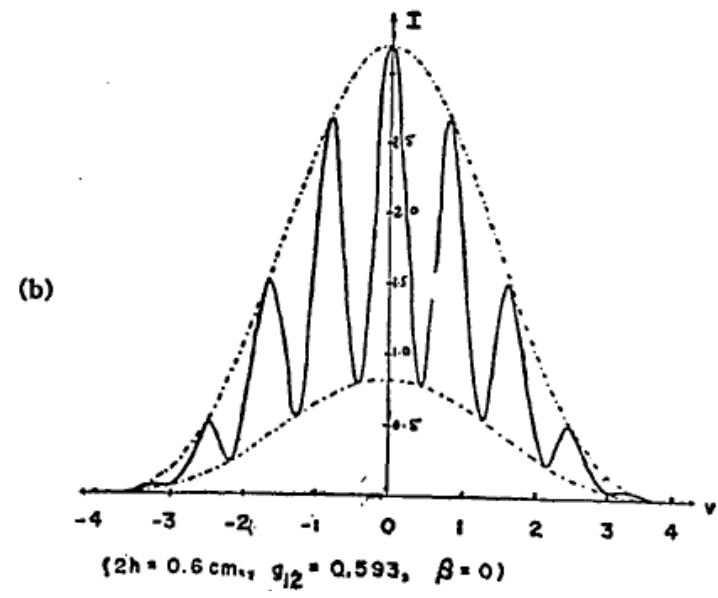
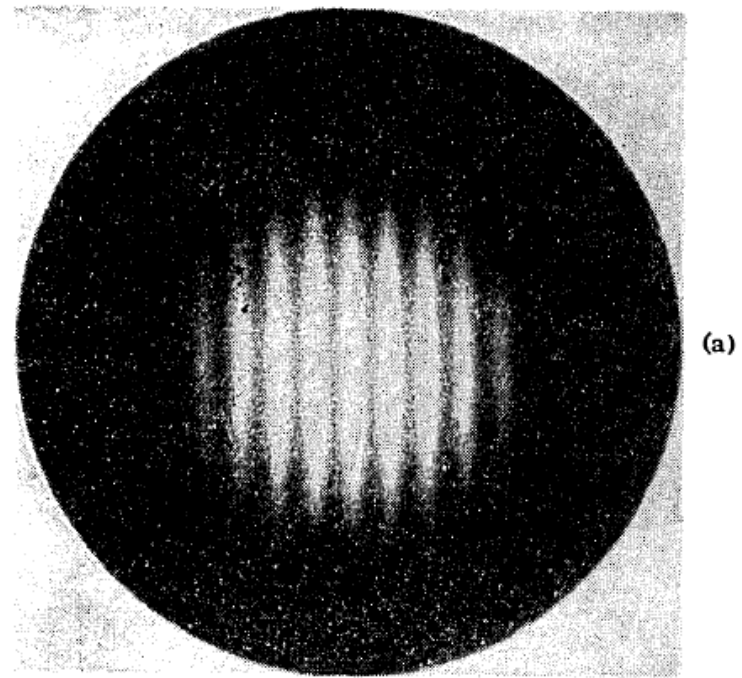
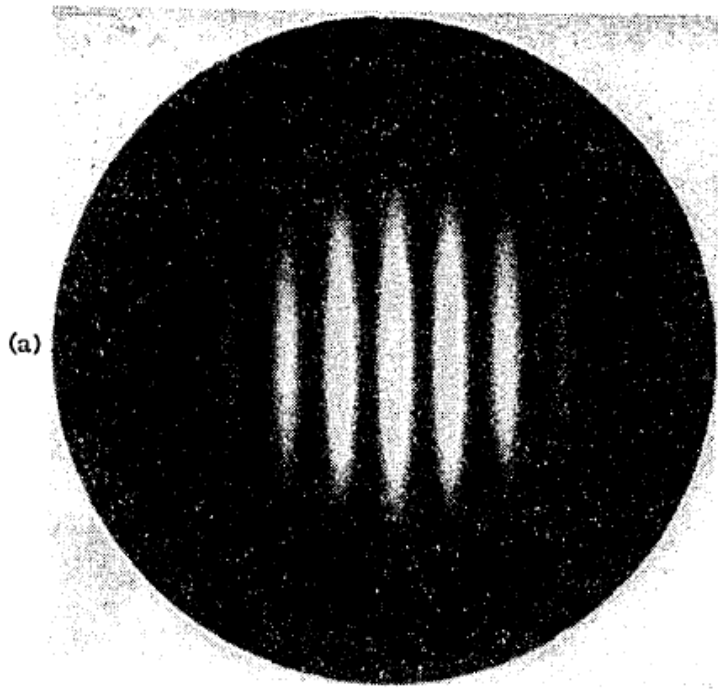
- Light is a wave
- Visibility (coherence)

$$V = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$

Spatial coherence in Young's Double-Slit experiment

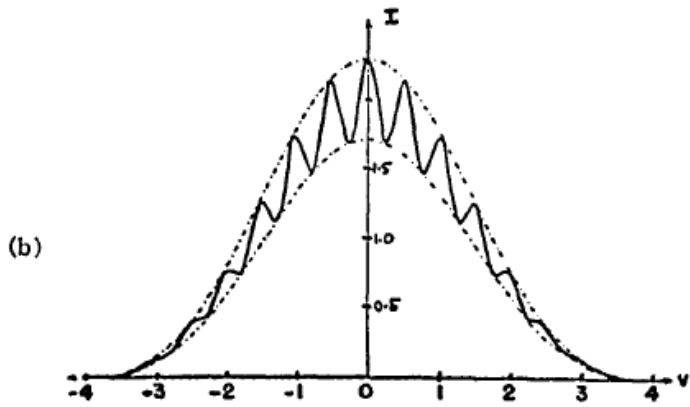
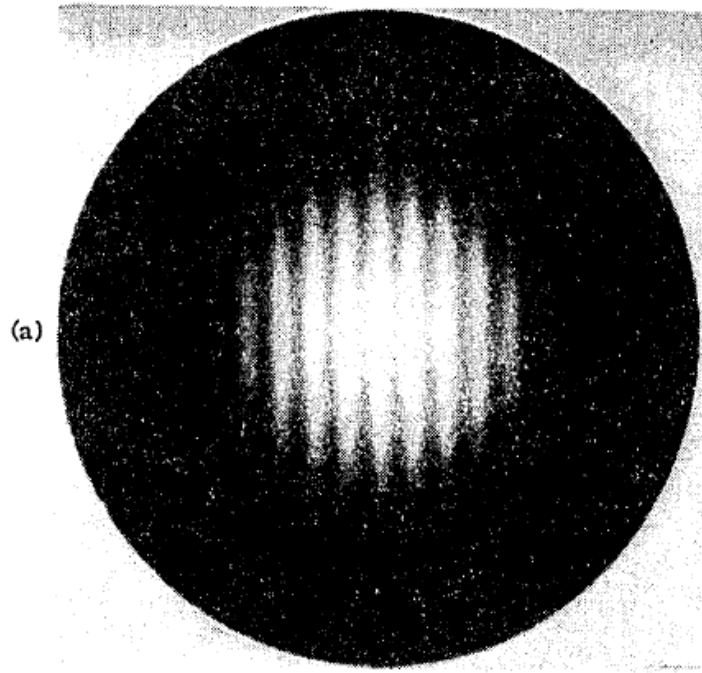


Born and Wolf, Optics



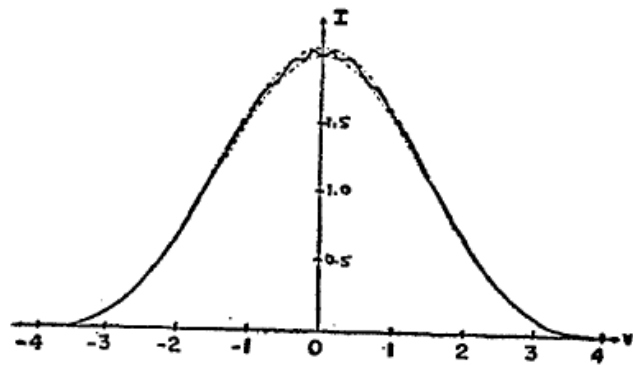
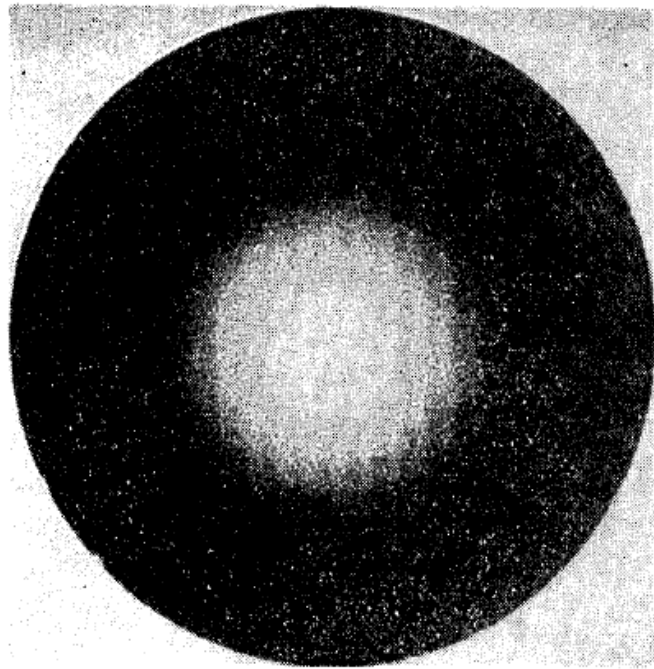
(b)

$$V = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$



$(2h = 1.0 \text{ cm.}, g_{12} = 0.146, \beta = 0)$

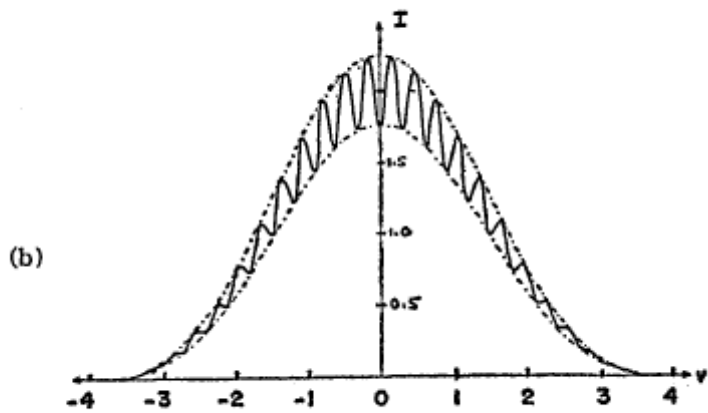
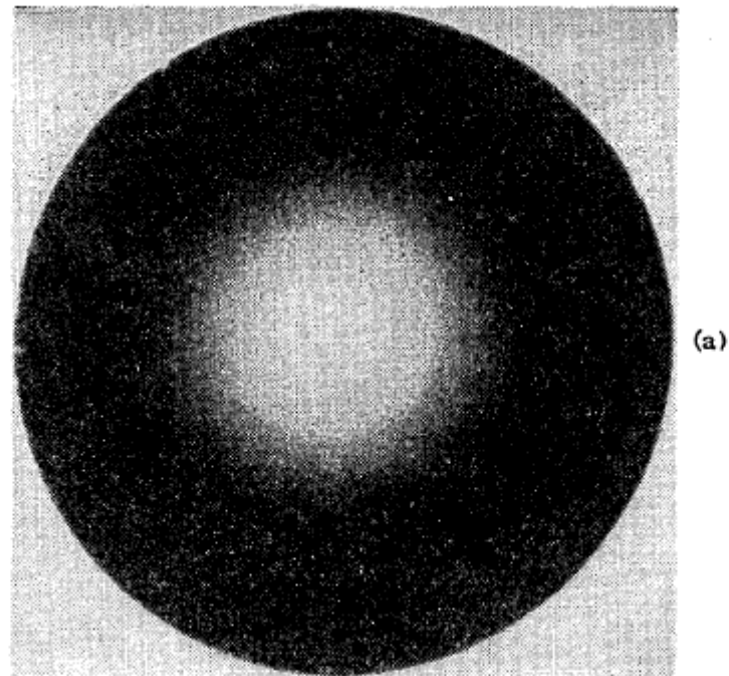
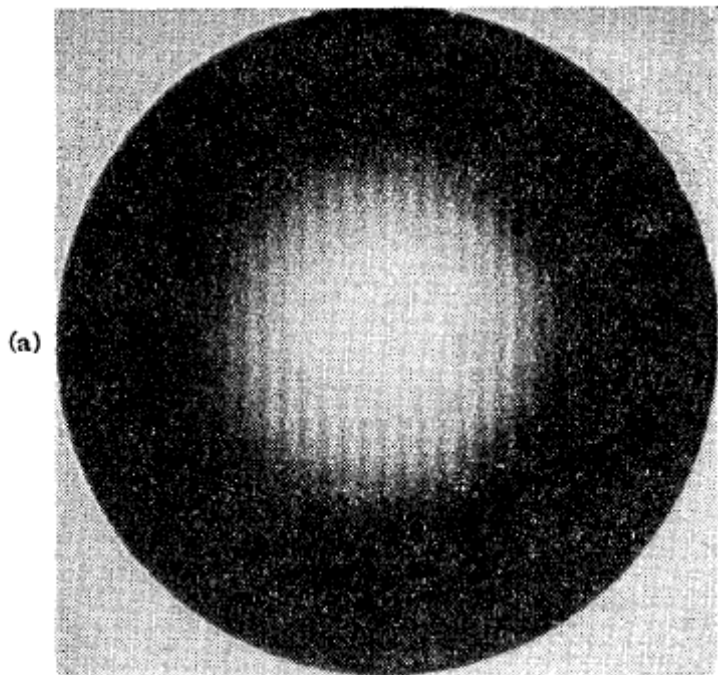
FIG. 4(C)



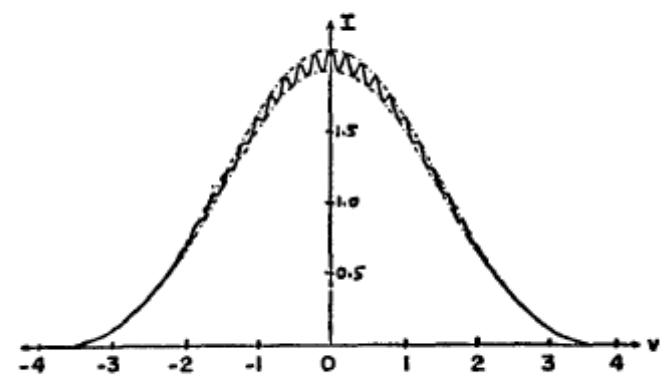
$(2h = 1.2 \text{ cm.}, g_{12} = 0.015, \beta = \pi)$

FIG. 4(D)

$$v = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$



$(2h=1.7\text{cm.}, g_{12}=0.123, \beta=\pi)$



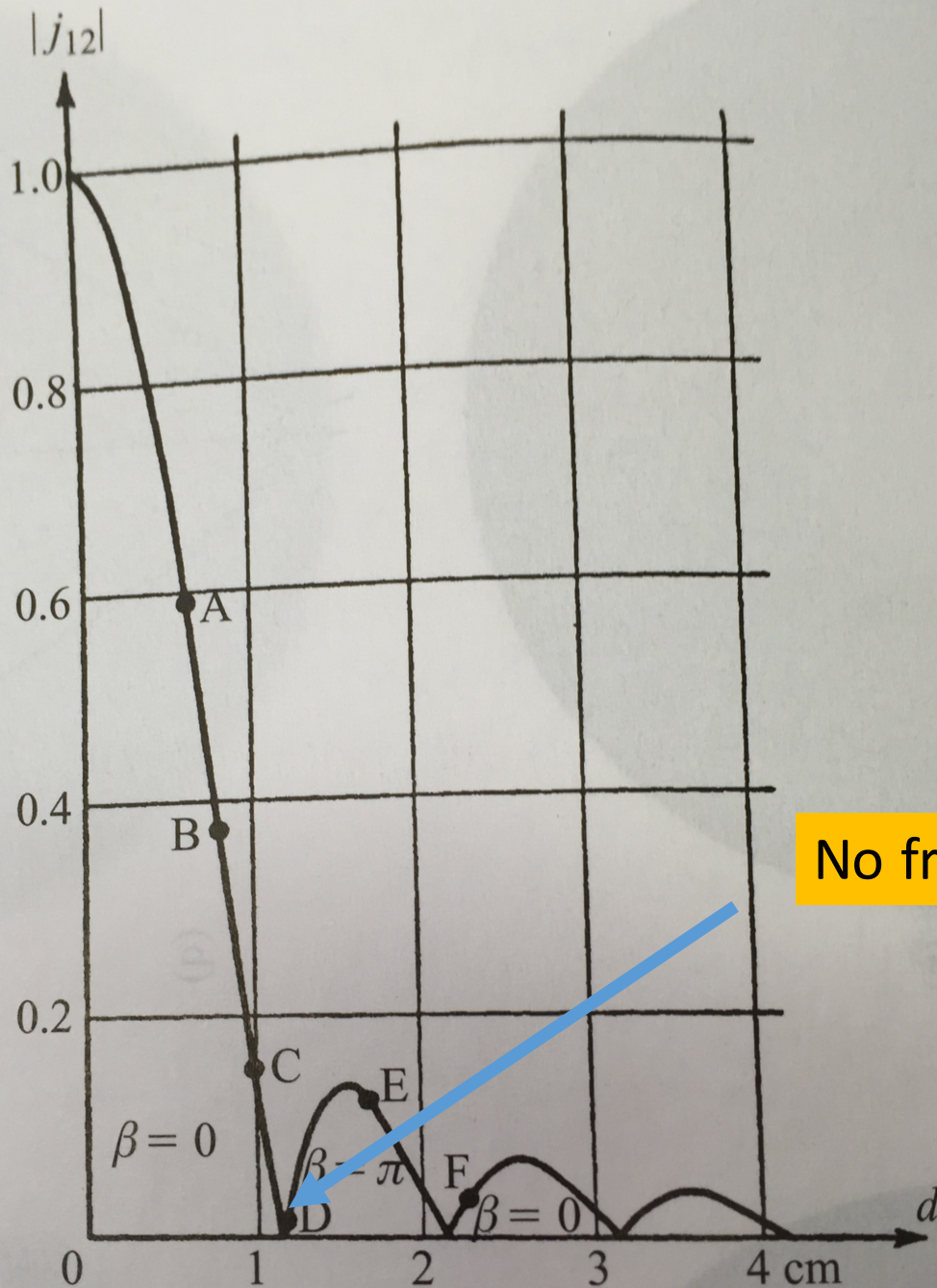
$(2h=2.3\text{cm.}, g_{12}=0.035, \beta=0)$

$$v = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$

(b)

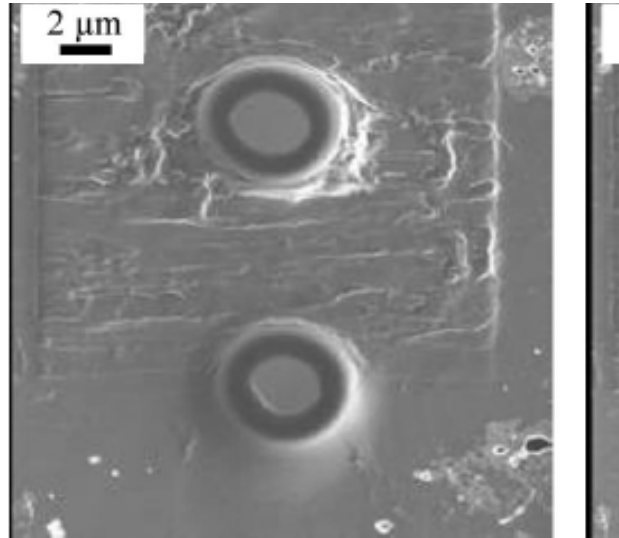
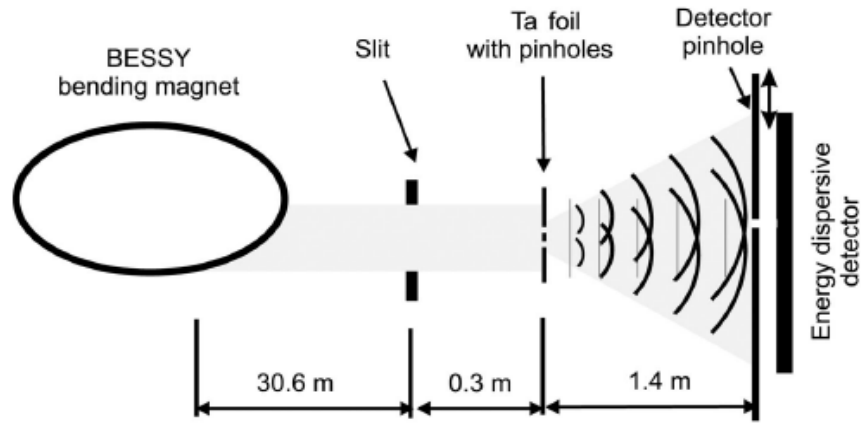
Fringe visibility as a function of distance between the pinholes

$$\Gamma(r_1, r_2, \tau) = \langle V^*(r_1, t) V(r_2, t + \tau) \rangle$$

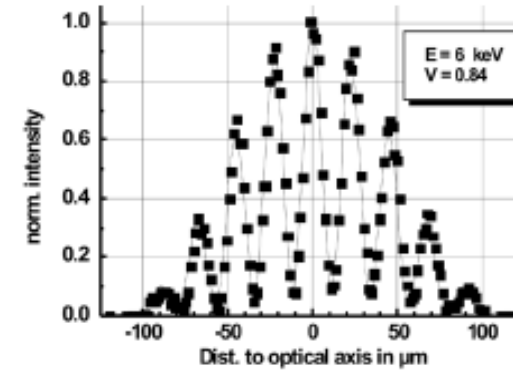


No fringes visibility: „coherence length exceeded“

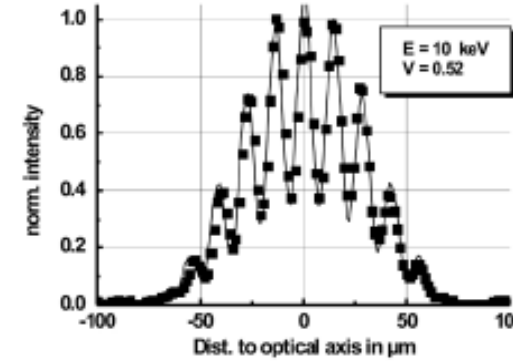
Young's experiment with X-rays



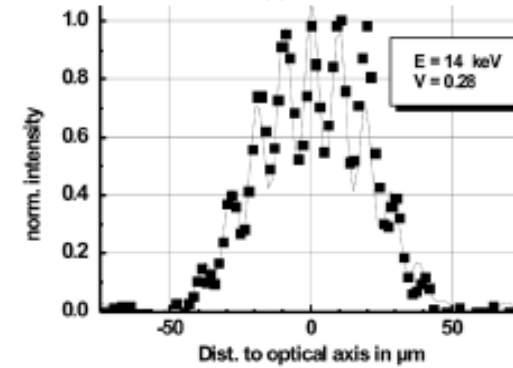
(a)



(a)



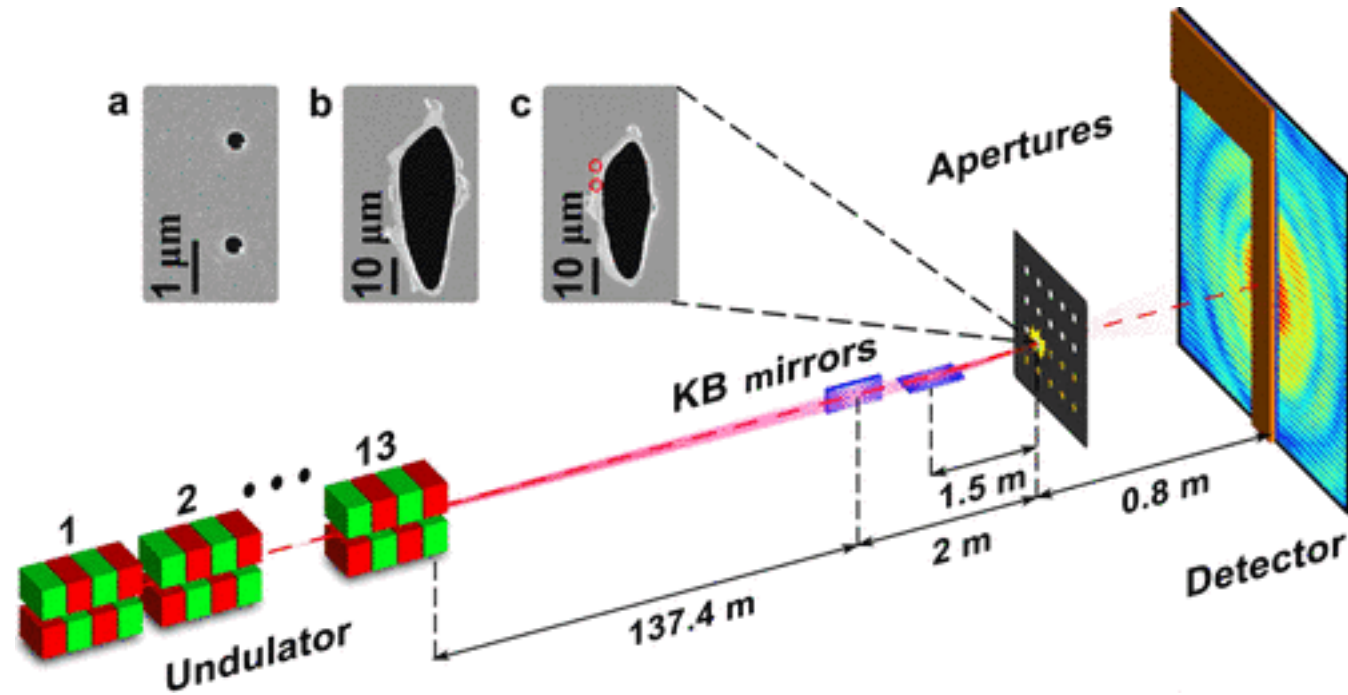
(b)



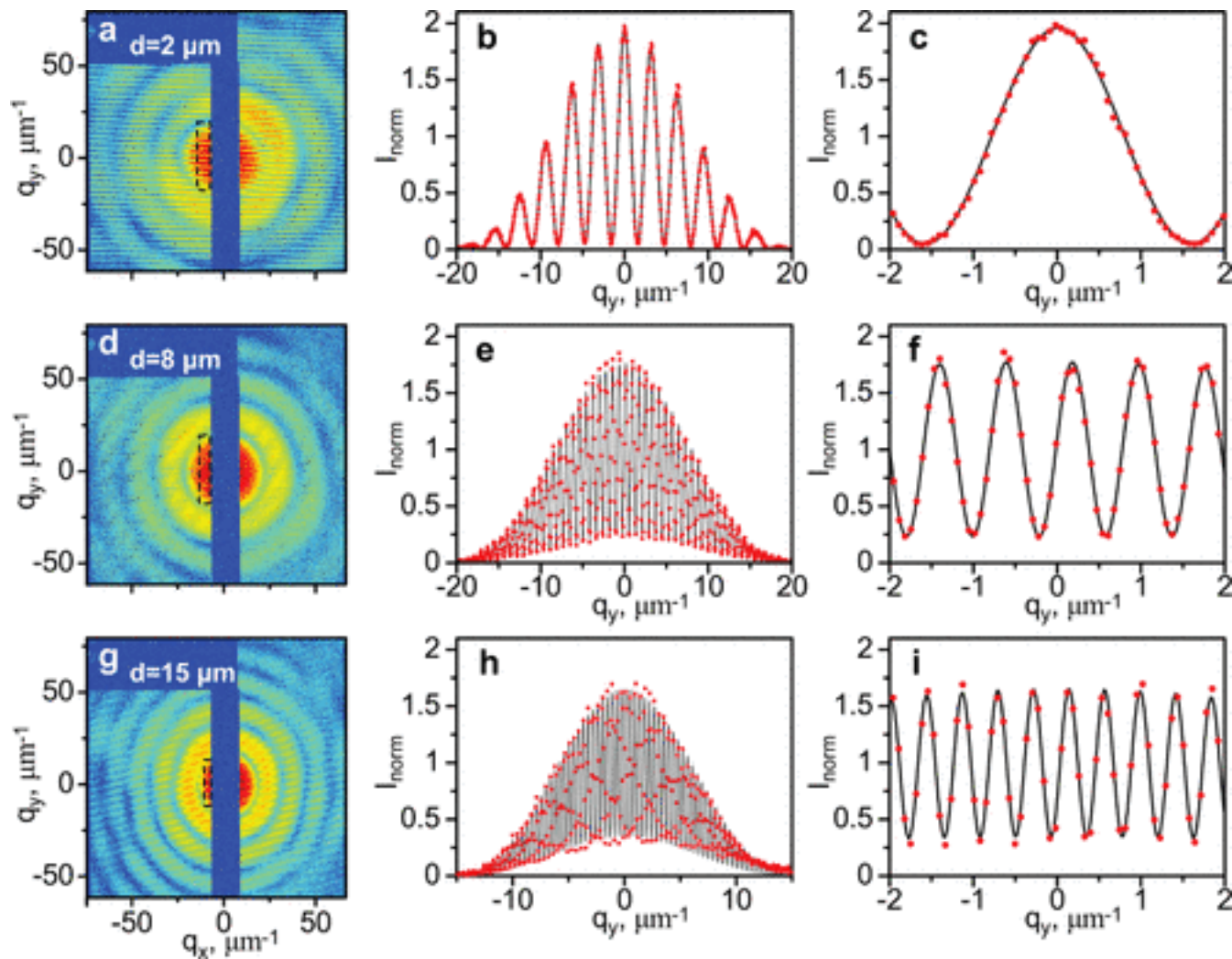
(c)

$$v = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$

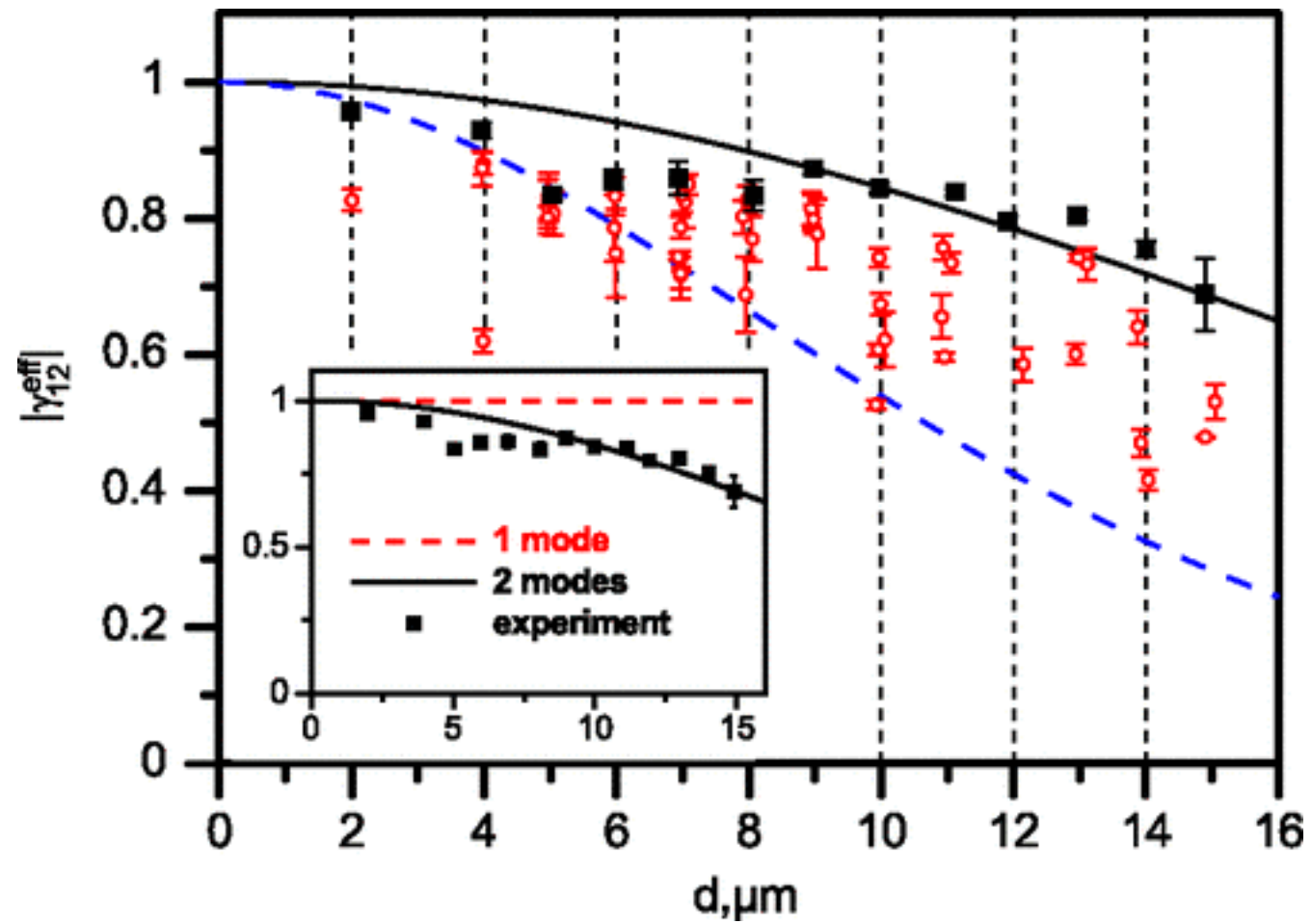
Young's experiment at an XFEL (here LCLS)



Vartanians et al. PRL 2012



$$v = \frac{I_{\text{max}} - I_{\text{min}}}{I_{\text{max}} + I_{\text{min}}}$$



Vartanians et al. PRL 2012

Coherent X-rays

third generation synchrotron sources

(ESRF, APS, Spring-8,...)

(and soon : Diamond, Soleil, SLS, Petra-III)



Transverse coherence lengths ξ_T

$$\xi_T = \frac{\lambda}{2} \cdot \frac{R}{s} = 155(V) \times 3.8(H) \mu\text{m}^2$$

Troika I Beamline parameters

$$\lambda = 1.55\text{\AA}, R=46\text{m},$$

source size : 23(V)×928(H) μm

Longitudinal coherence length ξ_L

$$\xi_L = \lambda \cdot \left(\frac{\lambda}{\Delta\lambda} \right) \approx 1\mu\text{m}$$

$$\lambda = 1.55\text{\AA}, \text{ Si}(111) \Delta\lambda/\lambda = 1.4 \cdot 10^{-4}$$

Selectioning coherent
part of the beam

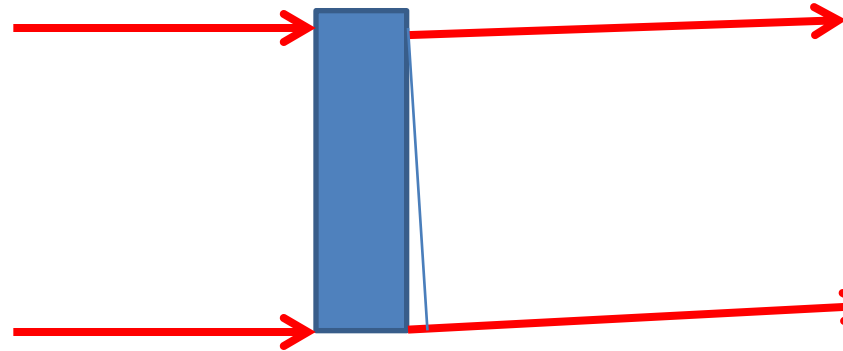
beam size $\approx \xi_T \times \xi_T$

Contrast (Visibility) $\beta(Q)$ of a speckle pattern is determined by the coherence properties of the X-ray beam

$$\beta(Q) = \frac{1}{V^2 \langle |E|^2 \rangle^2} \int_V \int_V |\Gamma(\vec{r}_2 - \vec{r}_1, \vec{Q} \cdot (\vec{r}_2 - \vec{r}_1) / ck_0)|^2 d\vec{r}_1 d\vec{r}_2$$

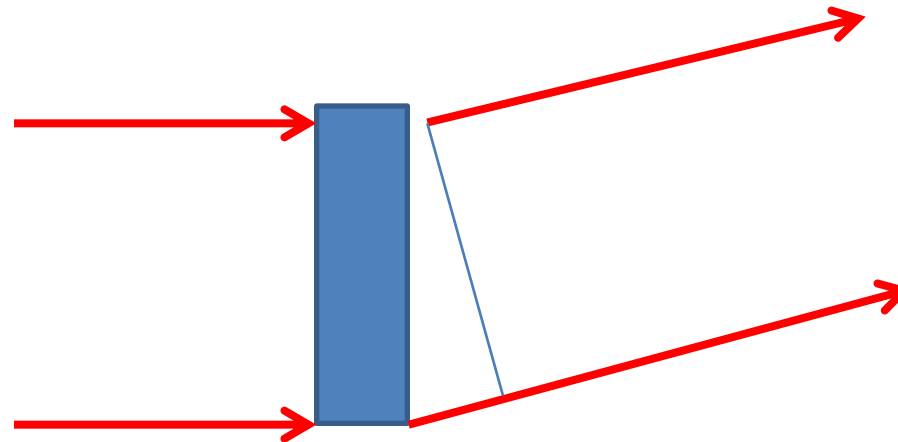
$\Gamma(r, \tau)$ mutual coherence function (MCF)

SAXS Q small
probing transverse coherence $\Gamma(r, 0)$

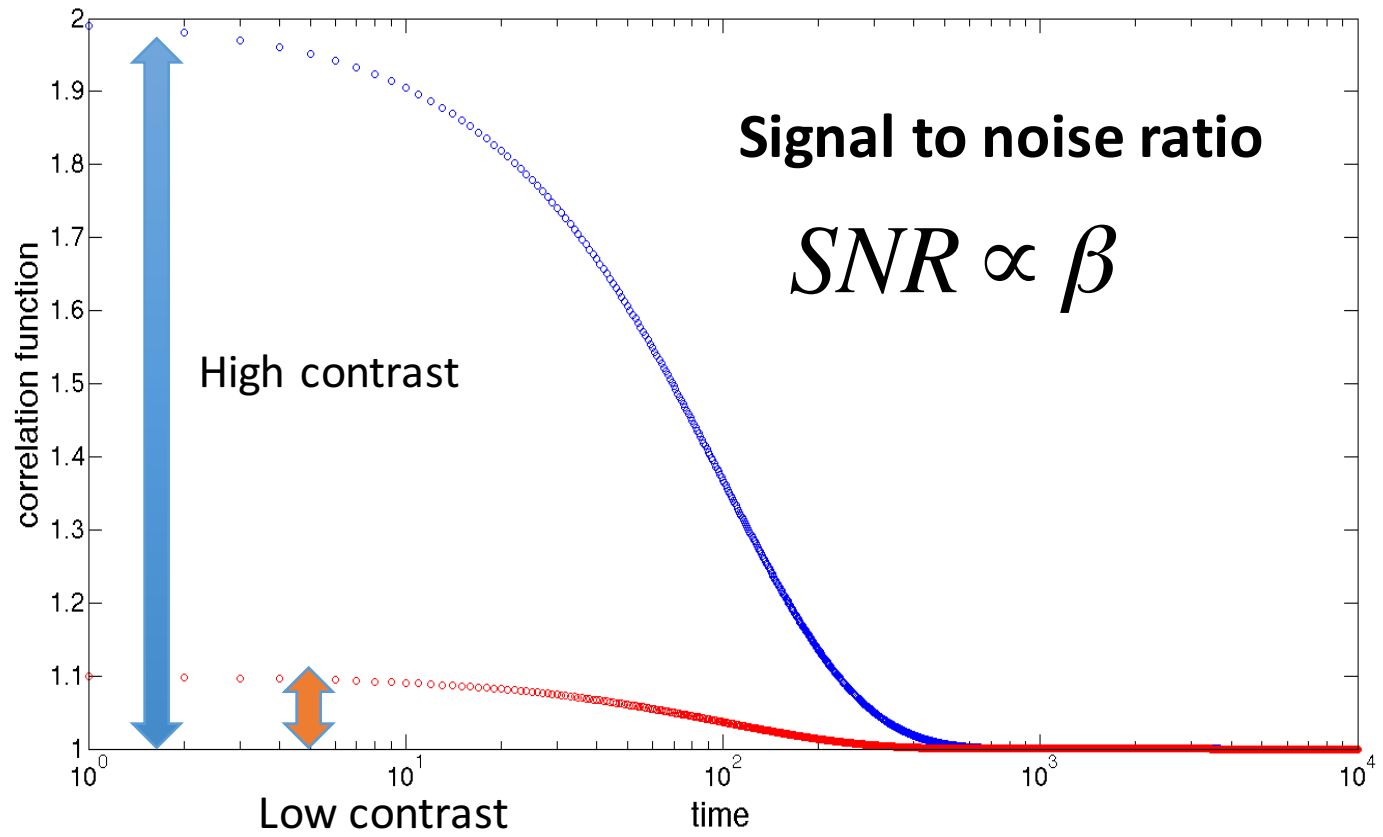


$$\Delta\tau = \frac{Q(r_2 - r_1)}{ck_0} \ll \tau_c$$

WAXS Q large
probing transverse AND temporal coherence $\Gamma(r, \Delta\tau)$

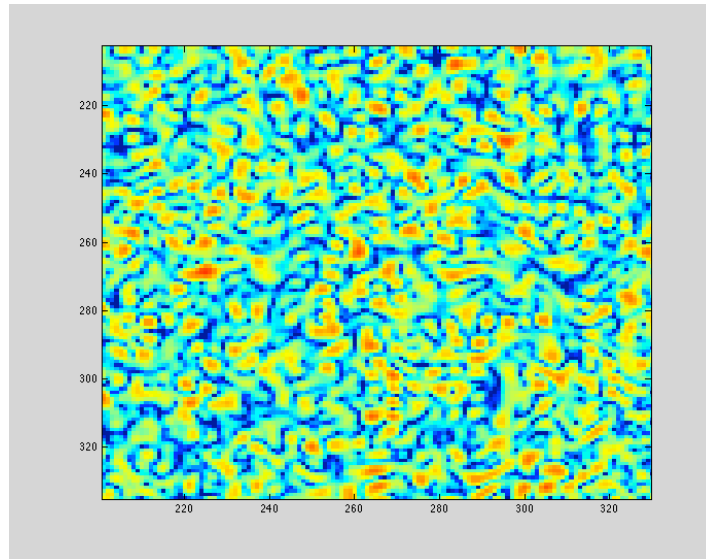


$$\Delta\tau = Q(r_2 - r_1) / ck_0 \sim \tau_c$$

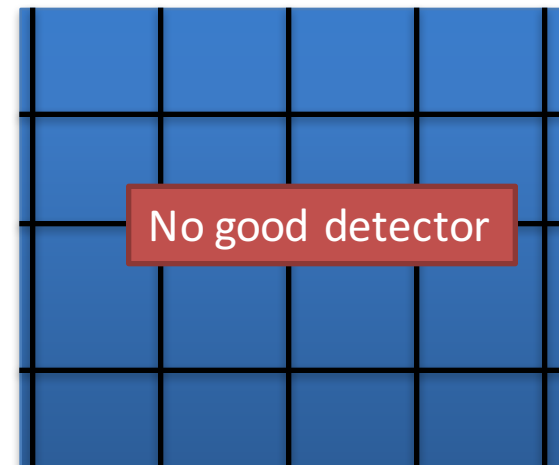
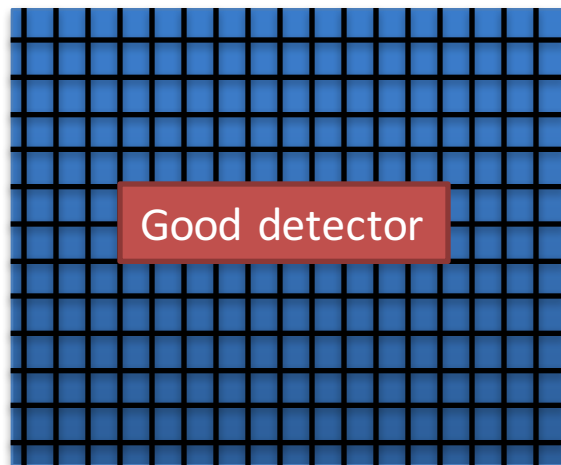
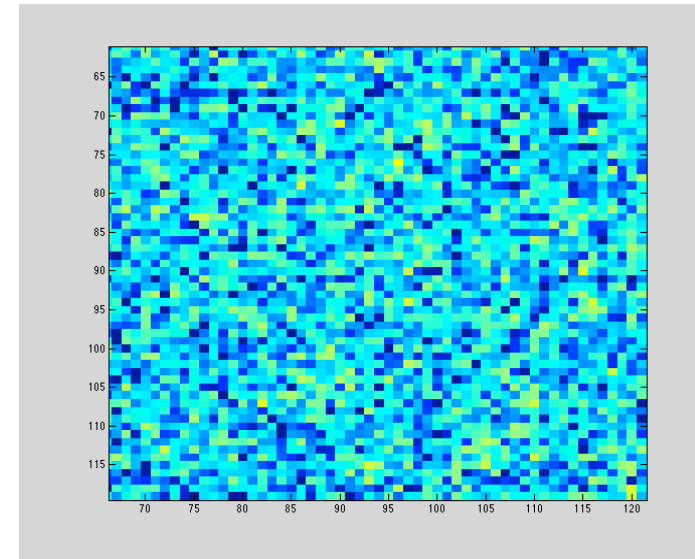


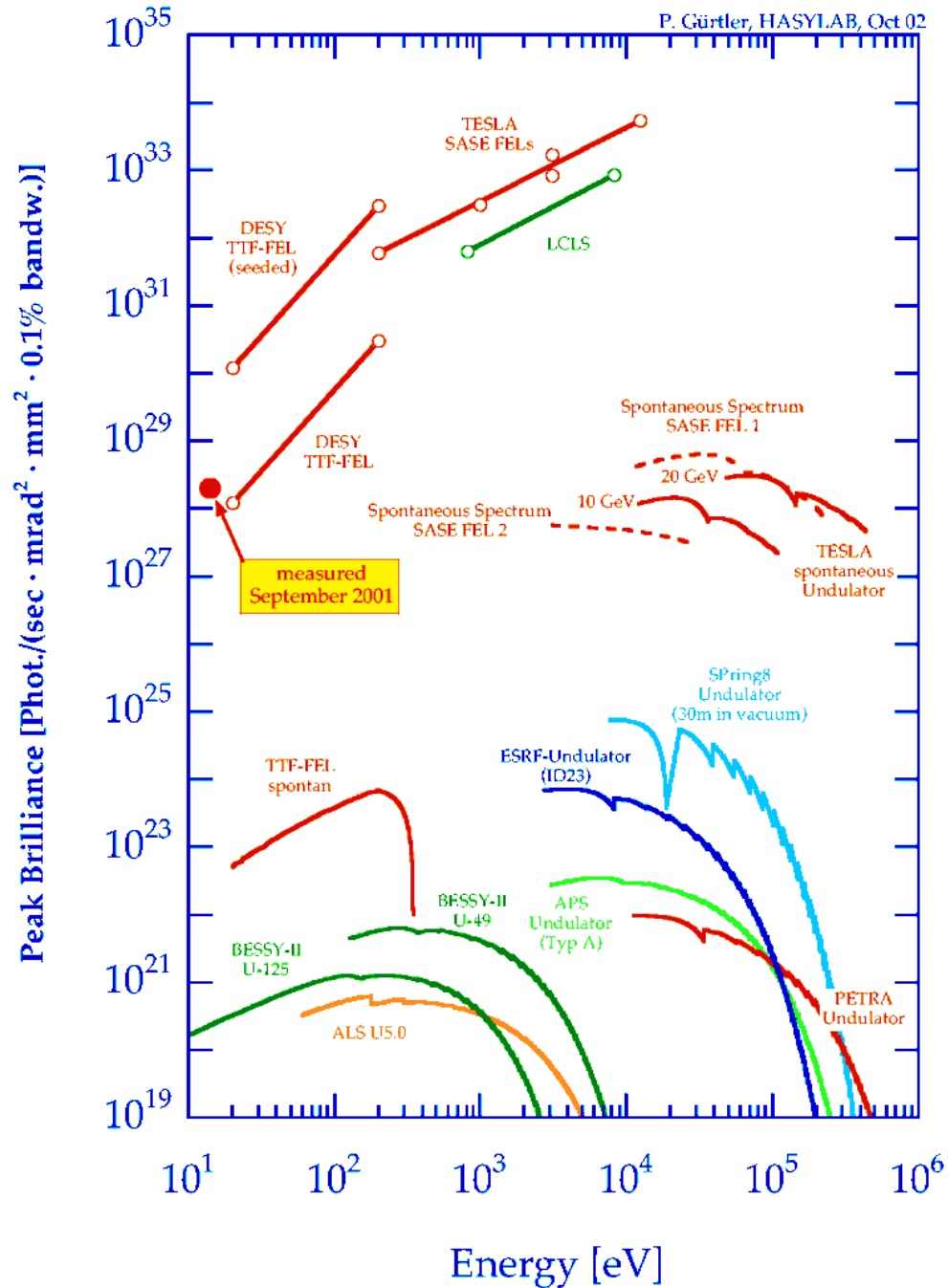
Speckle size needs to match pixel size of detector

Large speckles



Small speckles





Brilliance of X-rays Sources

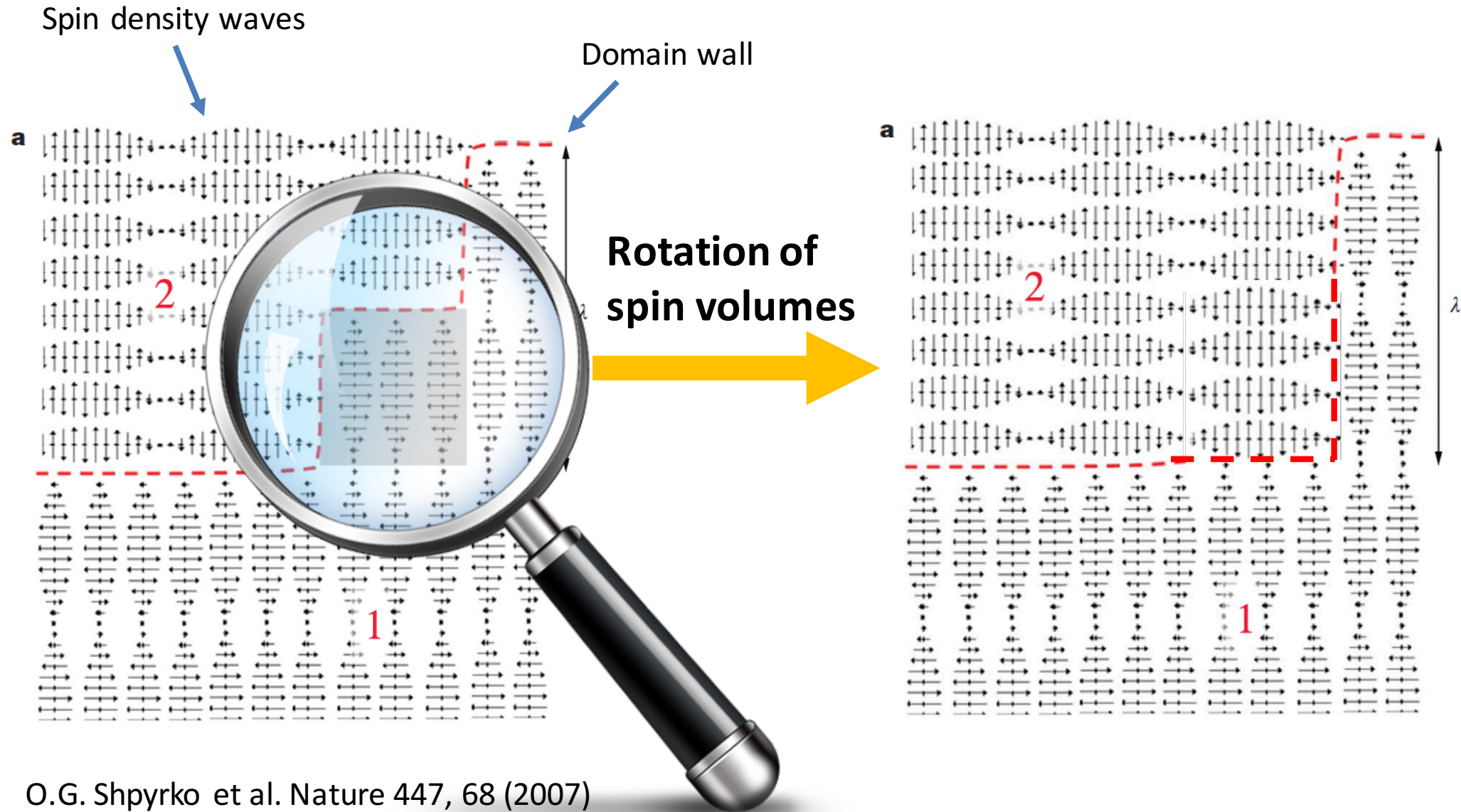
Coherent Flux: $F_0 = B \lambda^2 (\Delta\lambda/\lambda)$

(ESRF: ID10A $F_0 \sim 10^{10}$ ph/s)

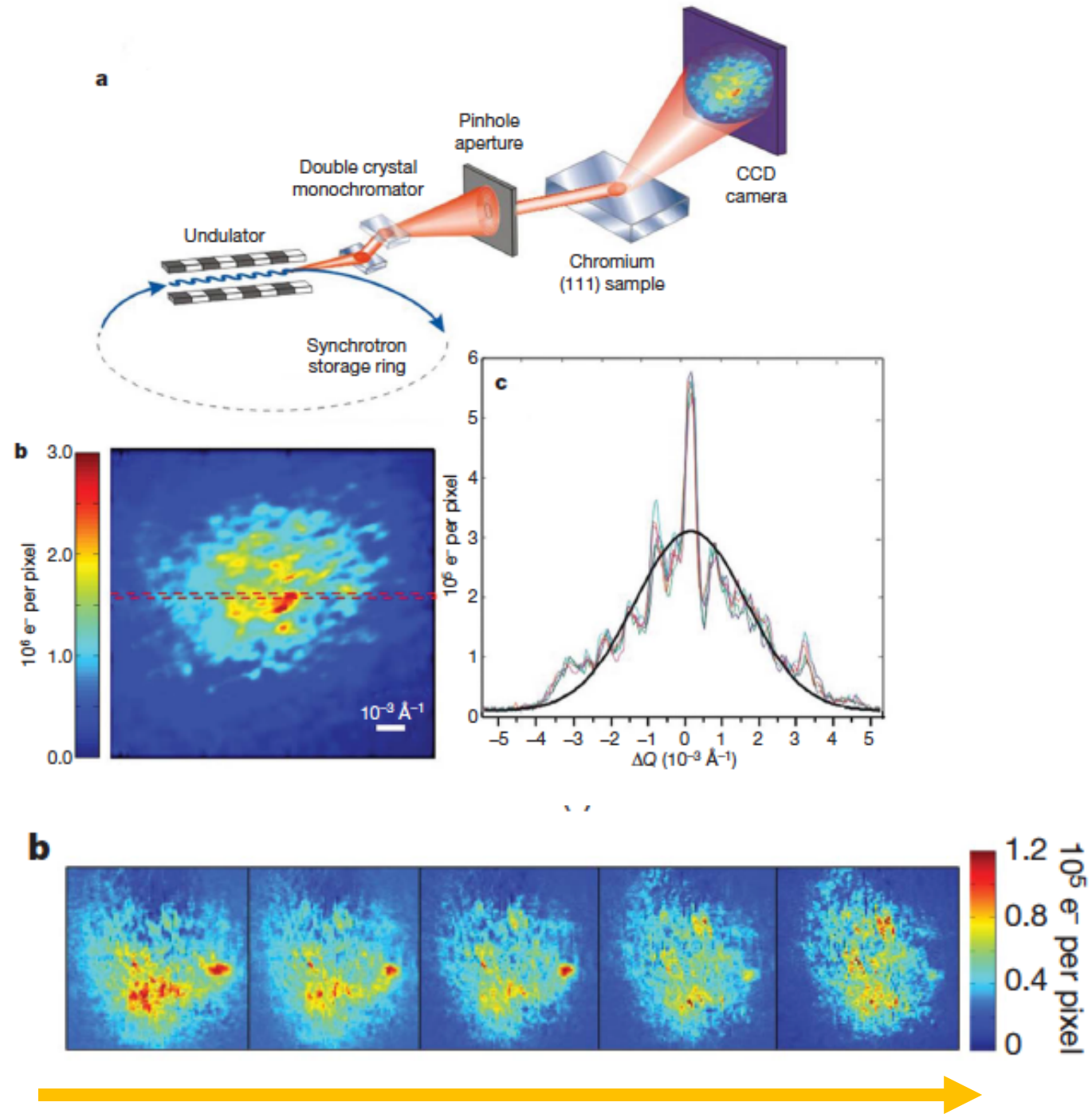
Examples

24	51.996
2672	1.6
1857	
Cr	
[Ar]3d ⁵ 4s	
7.19	2,3,6

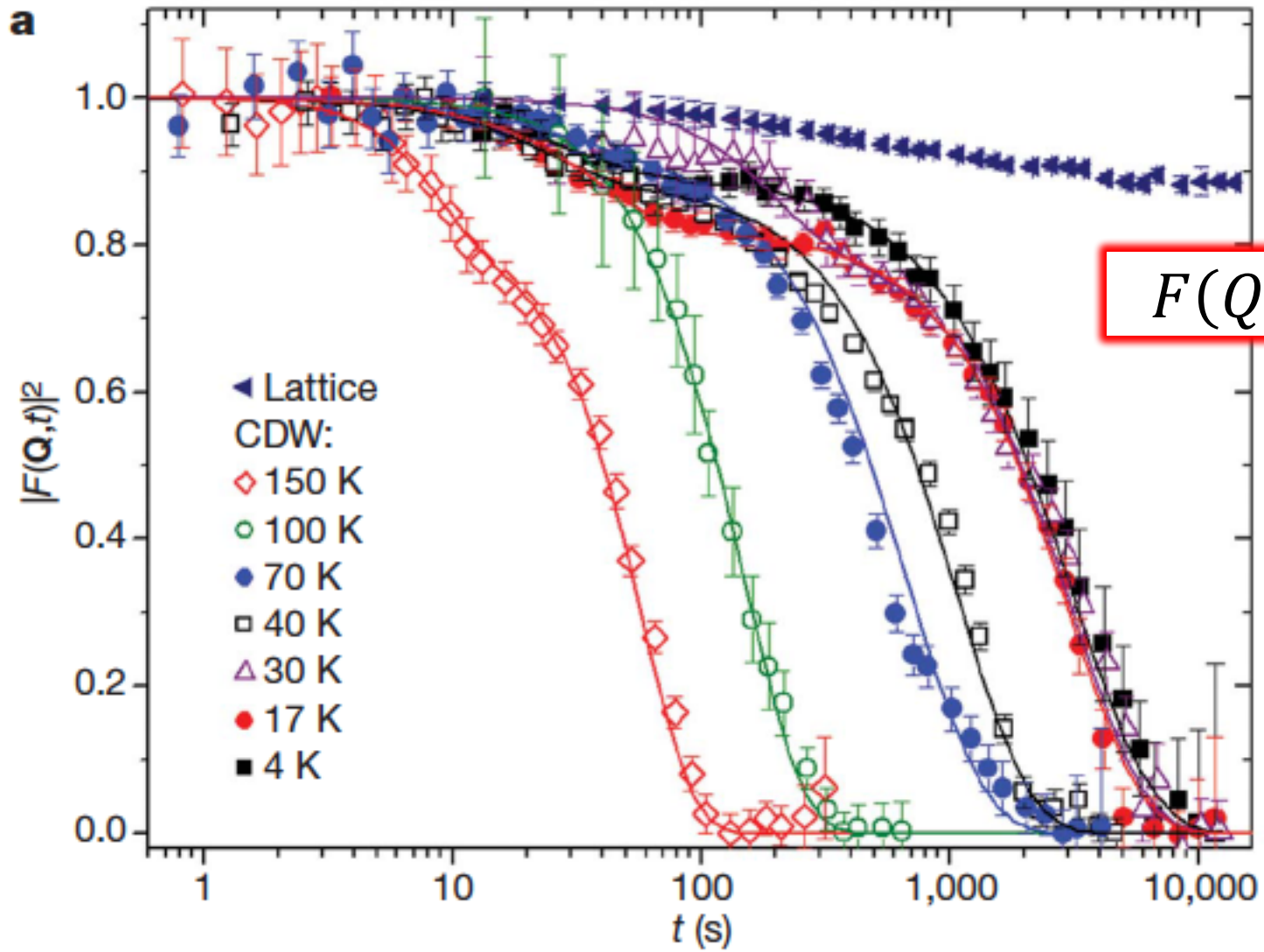
Antiferromagnetic domain fluctuations in Chromium



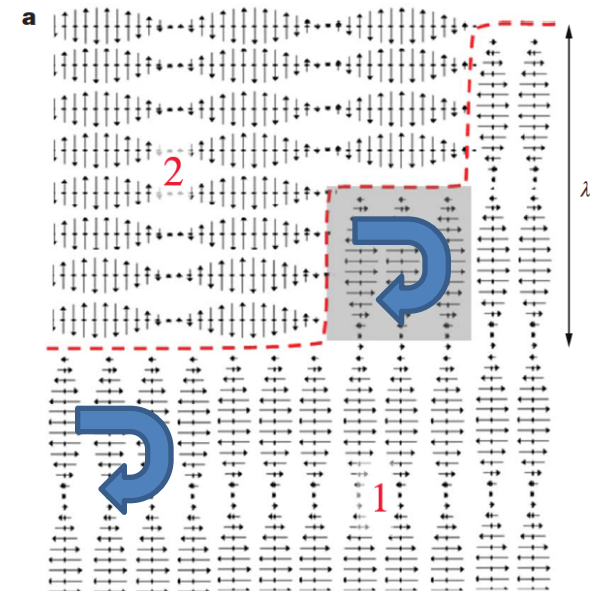
O.G. Shpyrko et al. Nature 447, 68 (2007)



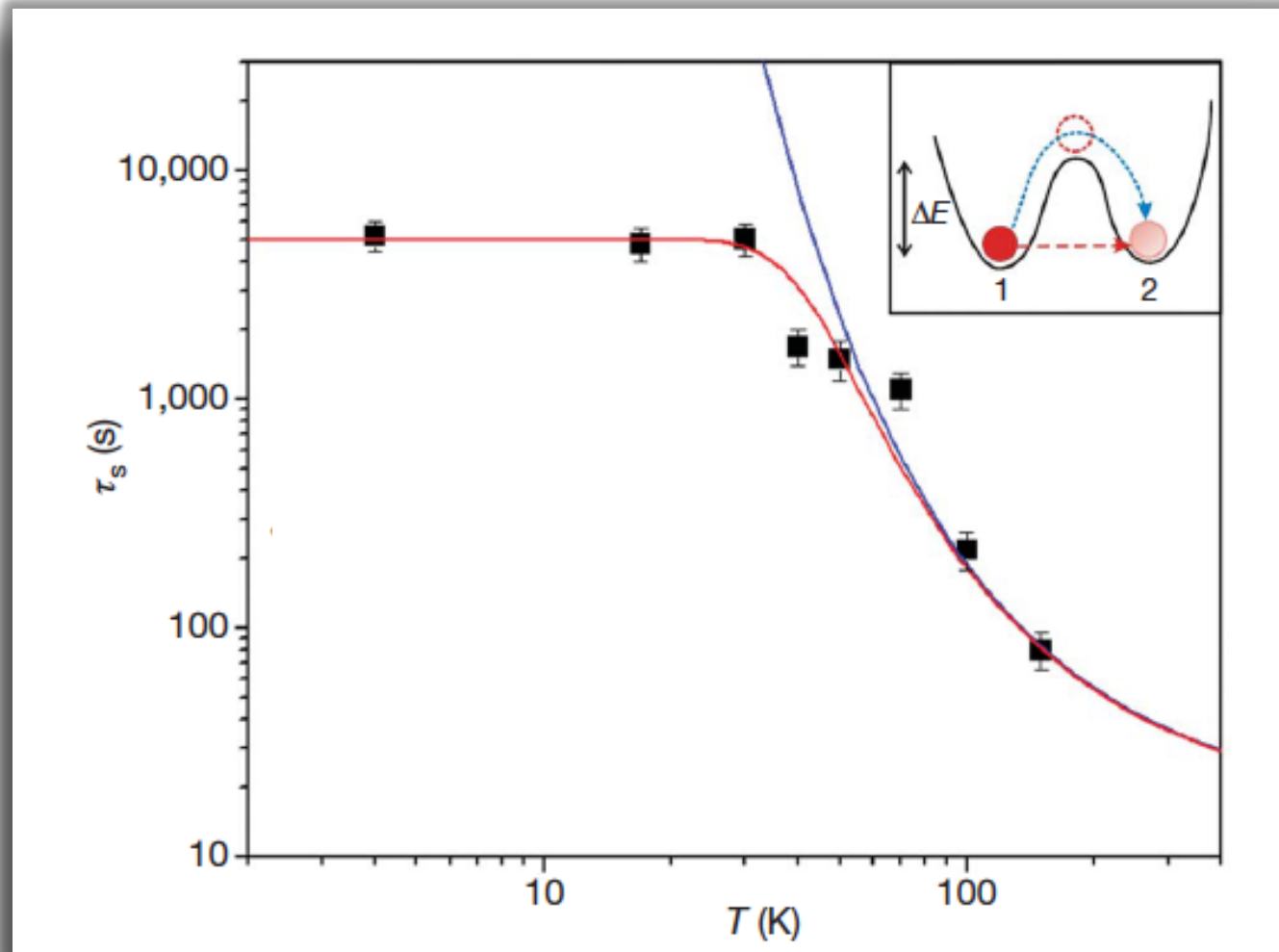
Correlation functions



$$F(Q, t) = \exp(-(t / \tau_s)^\beta)$$



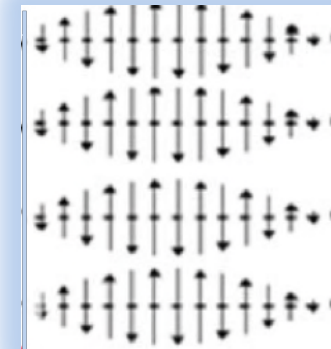
Quantum rotation of spin blocks



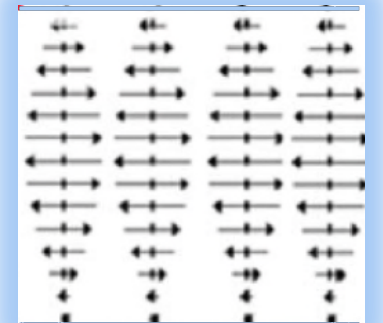
Blue line: Thermally activated jumps over an energy barrier

Red line: Quantum tunneling through an energy barrier

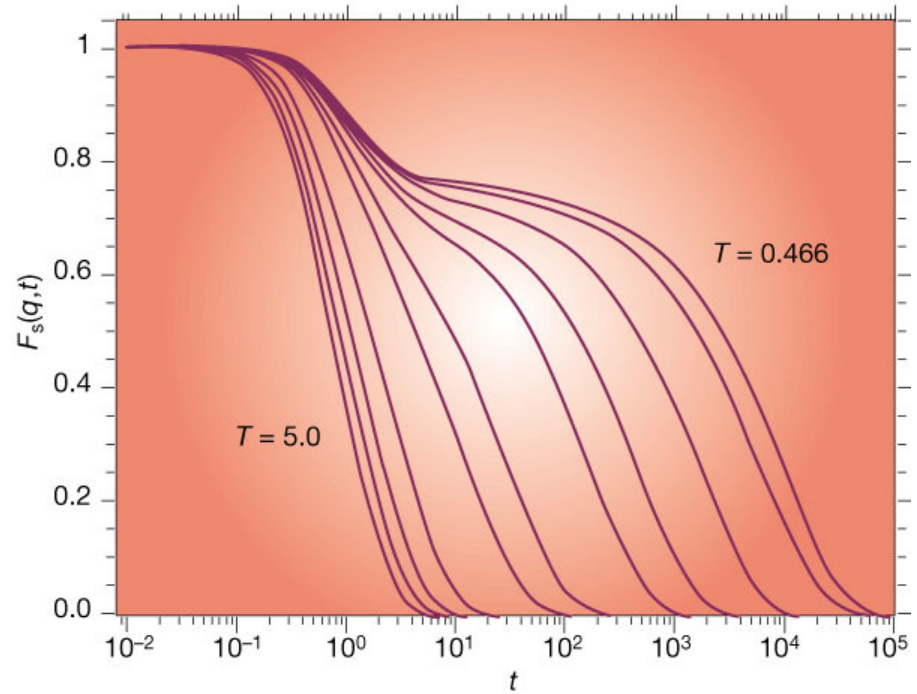
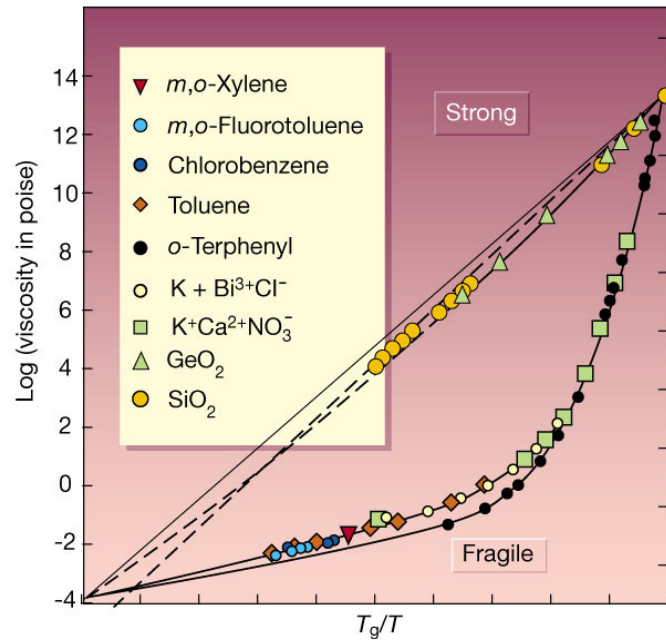
1



2



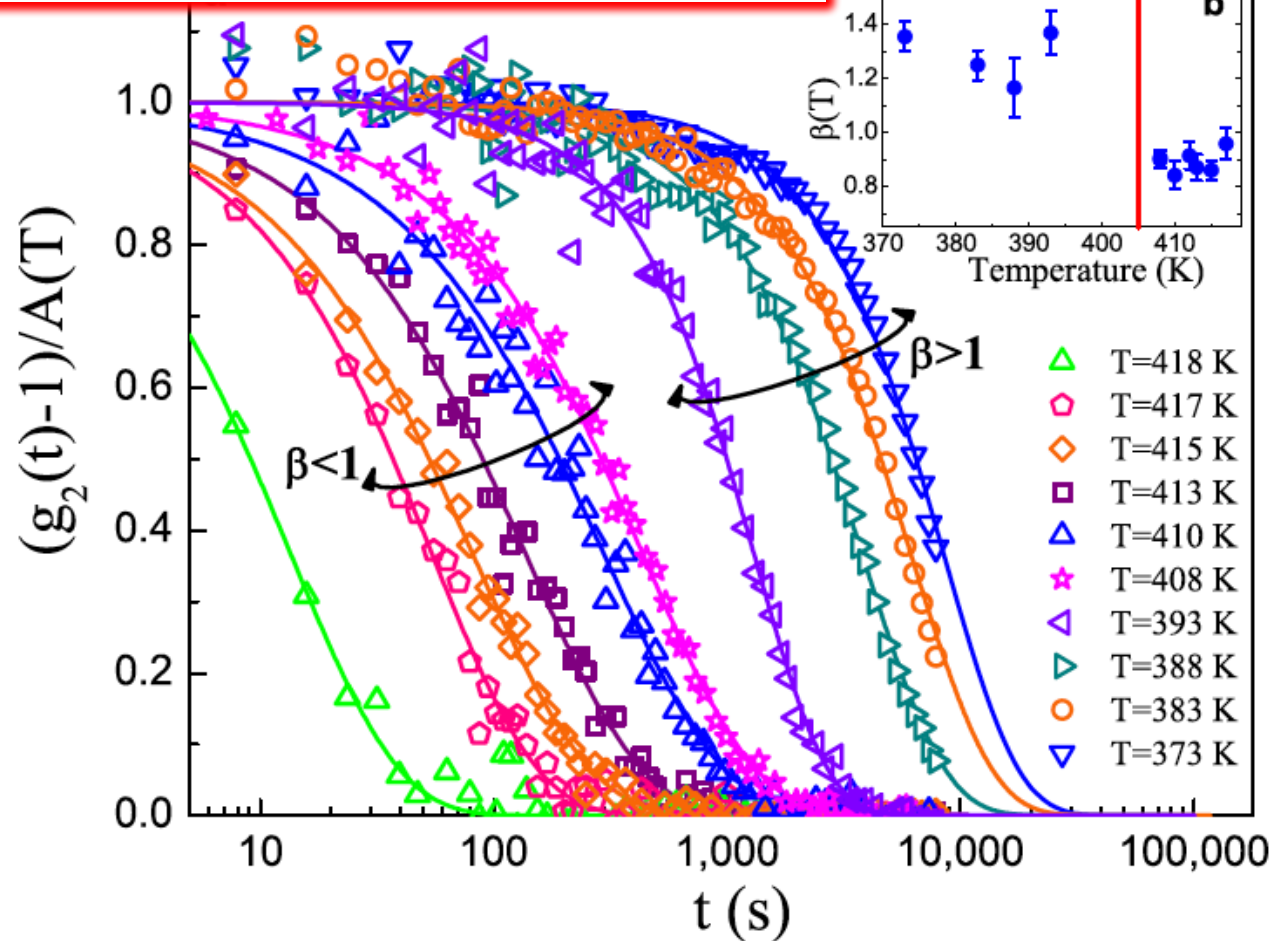
How Solid are Glasses ?



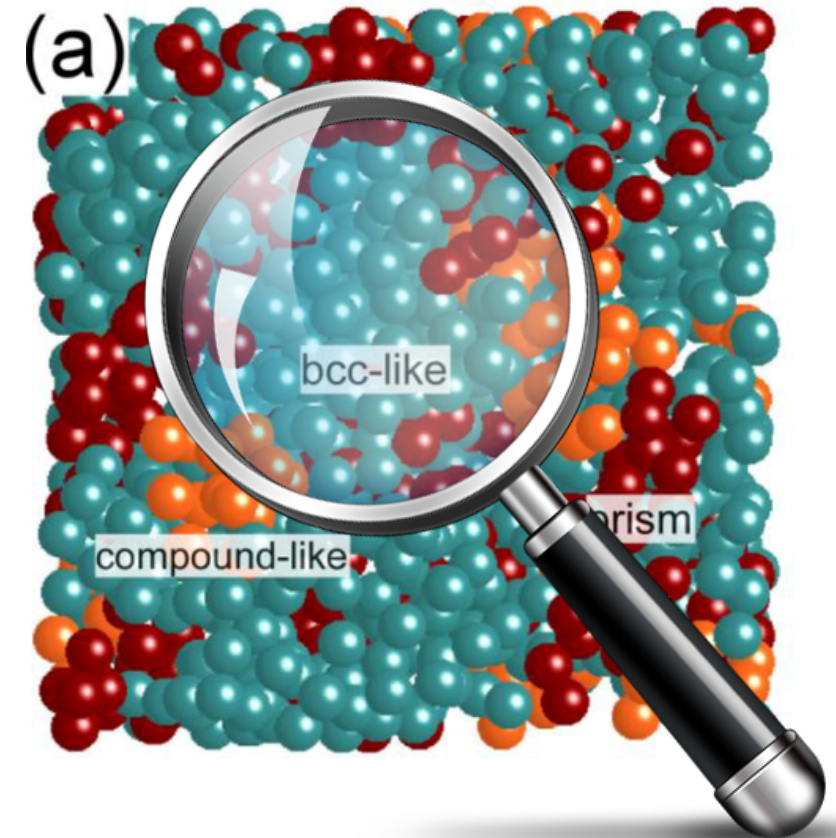
PABLO G. DEBENEDETTI AND FRANK H. STILLINGER,
Nature 410, 259 (2001)

Atomic dynamics in metallic glasses

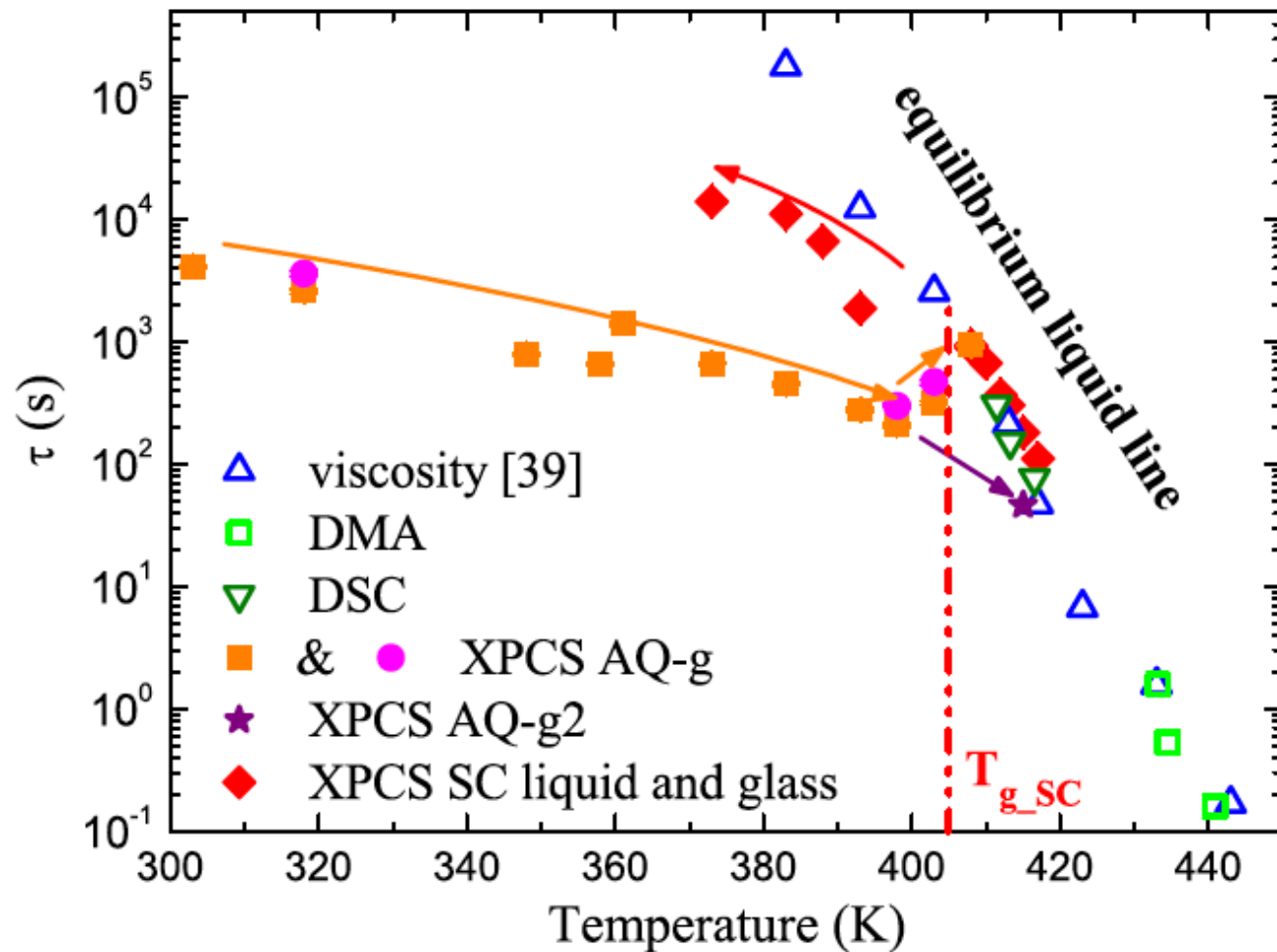
$$F(Q, t) = \exp(-(t / \tau_s)^\beta)$$



B. Ruta et al. Phys. Rev. Lett. 109, 165701 (2012)
B. Ruta et al. Nature Comm. 5, 3939 (2014)



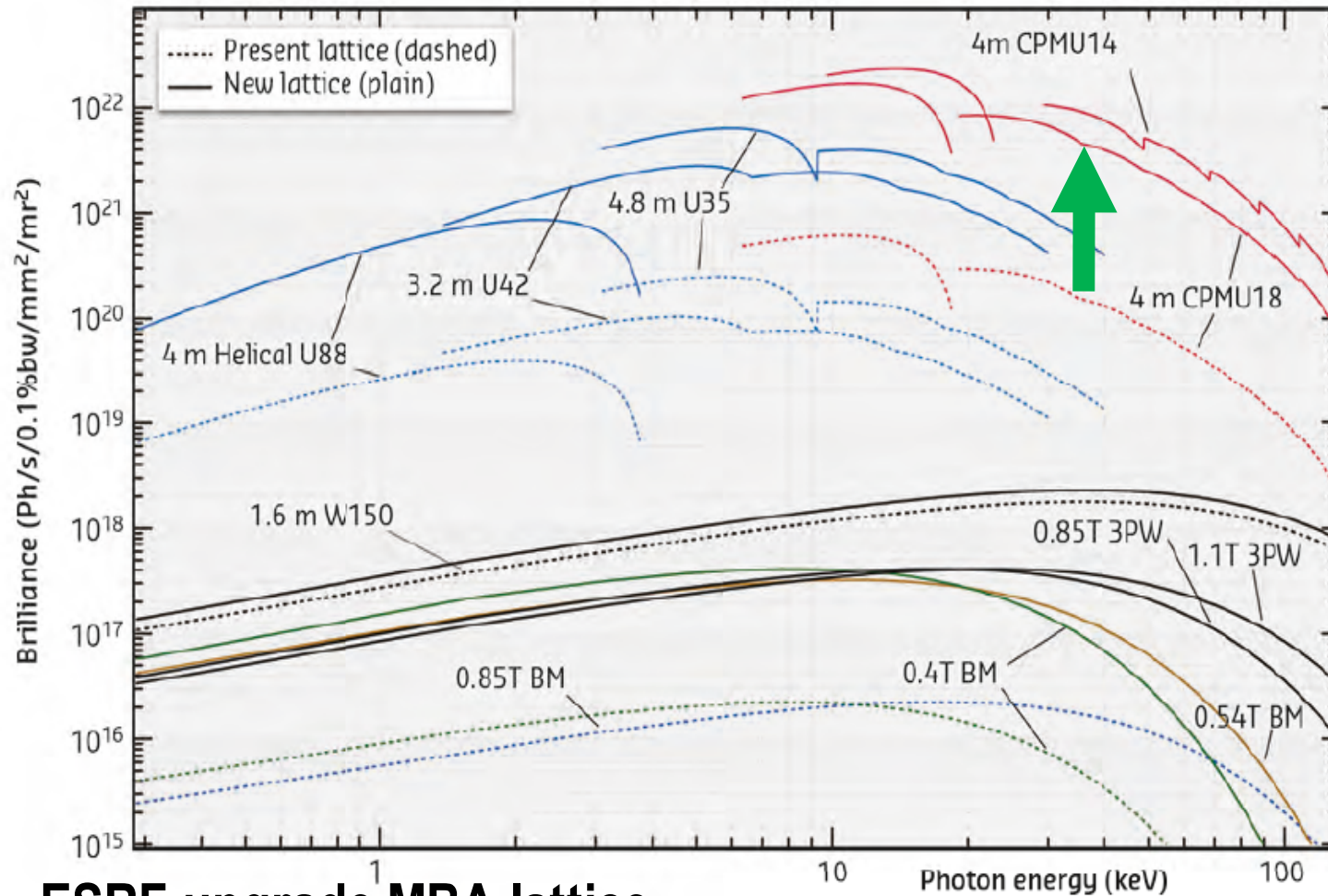
Reality check for glasses



- Fast relaxation dynamics exists below the glass transition temperature T_g .
- Glasses are not completely frozen in
- Stress dominates dynamics below T_g

B. Ruta et al. Phys. Rev. Lett. 109, 165701 (2012)
B. Ruta et al. Nature Comm. 5, 3939 (2014)

XPCS at diffraction limited storage rings (DLSR)



ESRF upgrade MBA lattice

Coherent Flux: $F_0 = B \lambda^2 (\Delta\lambda/\lambda)$

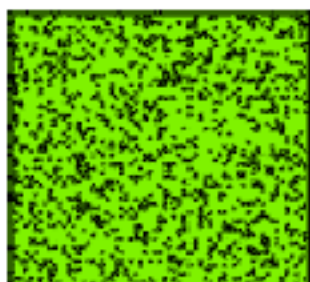
Increase of B by factor 50 - 100

up to 10.000 times faster time
scale accessible in XPCS

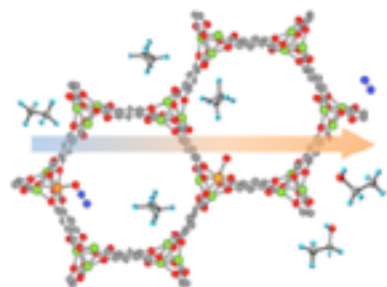
$$\tau \sim 1/B^2$$

unusual scaling because XPCS
correlates pairs of photons

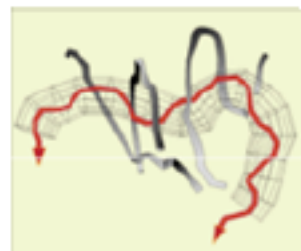
Nucleation kinetics
10.1126/science.1230915



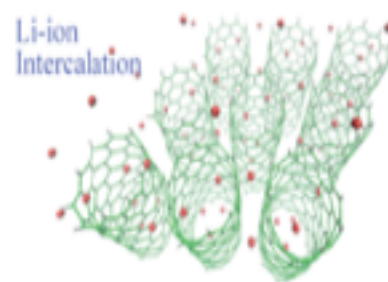
Selective Catalysis
10.1038/nchem.1956



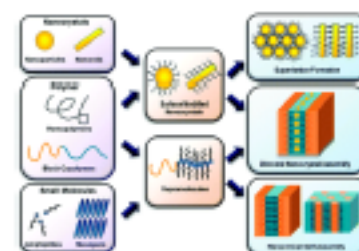
Polymer reptation



Intercalation kinetics,
10.1039/C0EE00473A



Self assembly



10^{-1}

10^{-3}

10^{-5}

10^{-7}

10^{-9}

10^{-11}

10^{-13}

10^{-15}

10^{-18}

Energy Scale
(eV)

10^{-12}

10^{-10}

10^{-8}

10^{-6}

10^{-4}

10^{-2}

10^0

10^2

Time Scale
(sec)

now/new

RIXS

FT-RIXS

Fast XPCS

XPCS

movies

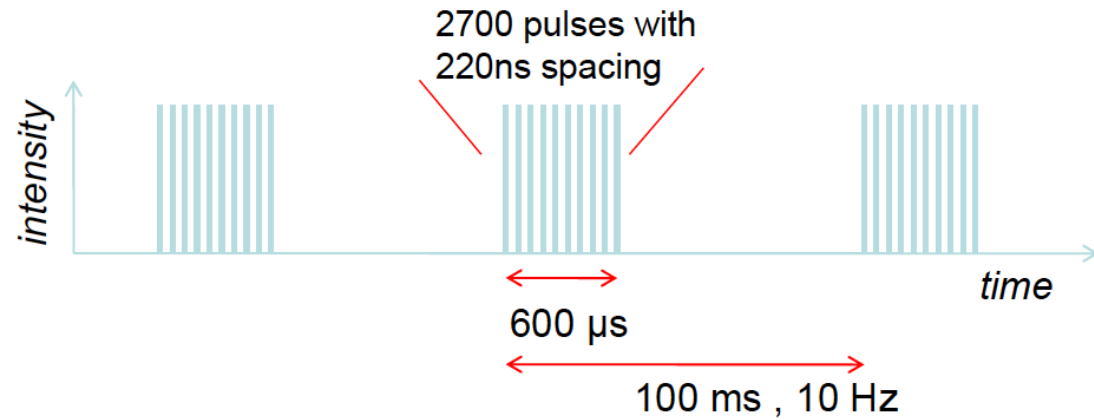
movies/XPCS

Problems that can be addressed at DLSR

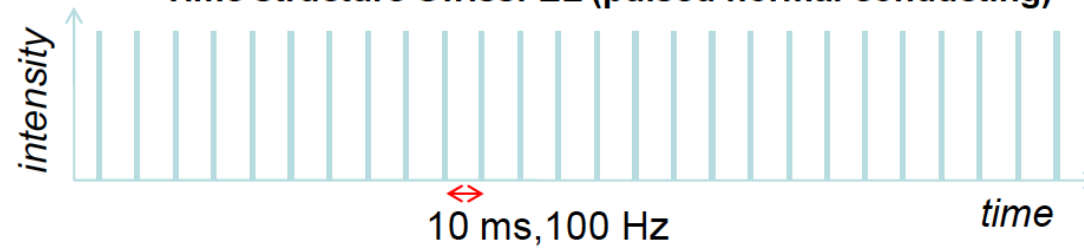
- Dynamics in the supercooled state
- Dynamics in confinement
- Domain fluctuations in hard condensed matter
- Protein diffusion in cells
- Kinetics of biomineralization processes
- Liquids under extreme conditions (e.g. pressure)
- Driven dynamics under external (B,E,T) fields
- Local structures and their relaxations
- ...

XPCS at XFELs

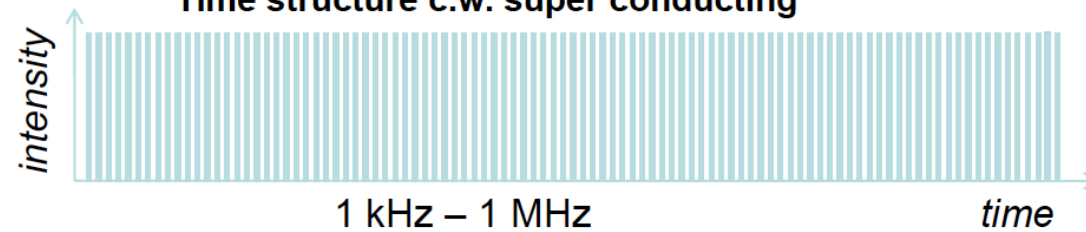
Time structure European XFEL (pulsed superconducting)



Time structure SwissFEL (pulsed normal conducting)

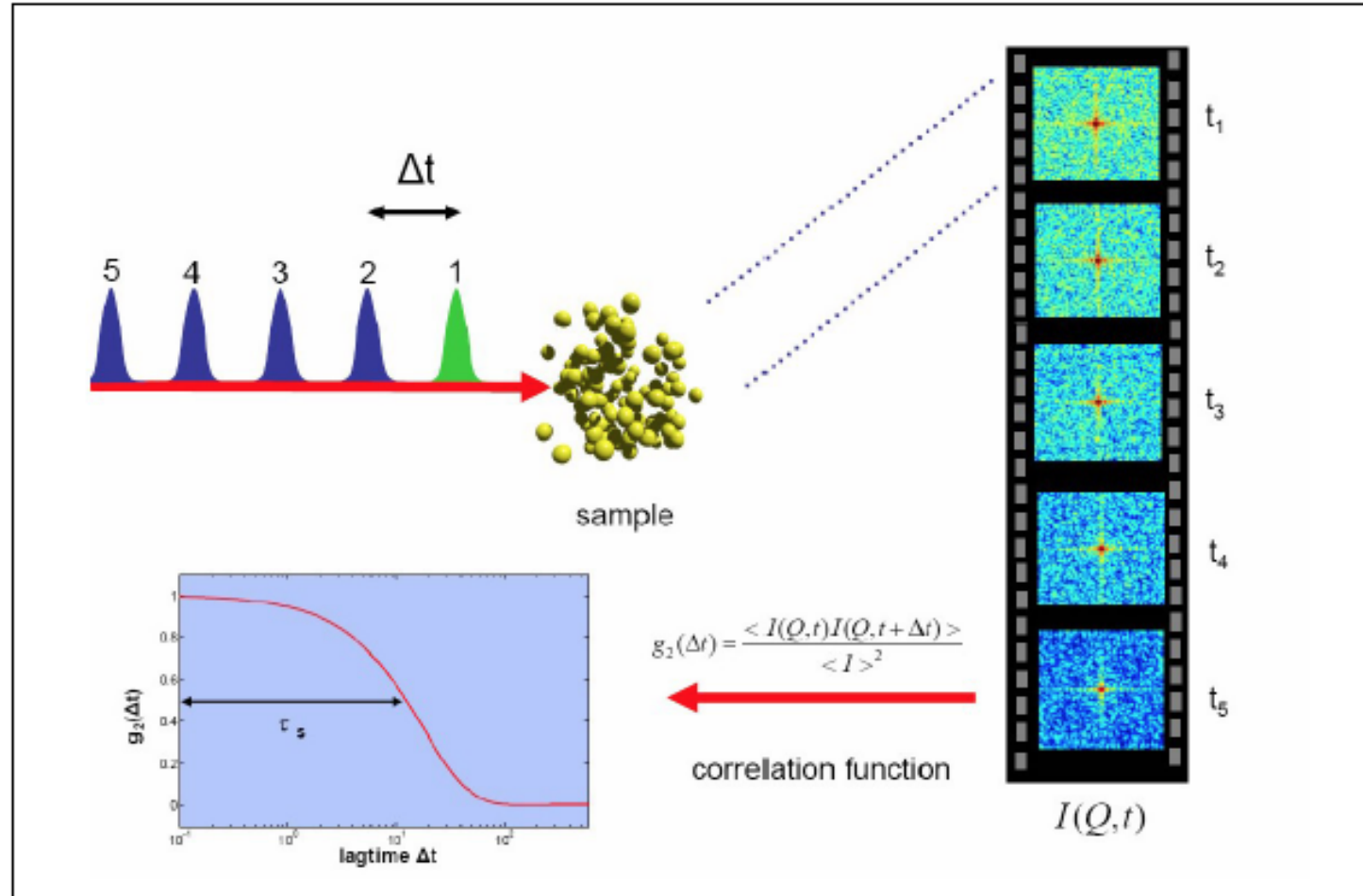


Time structure c.w. super conducting

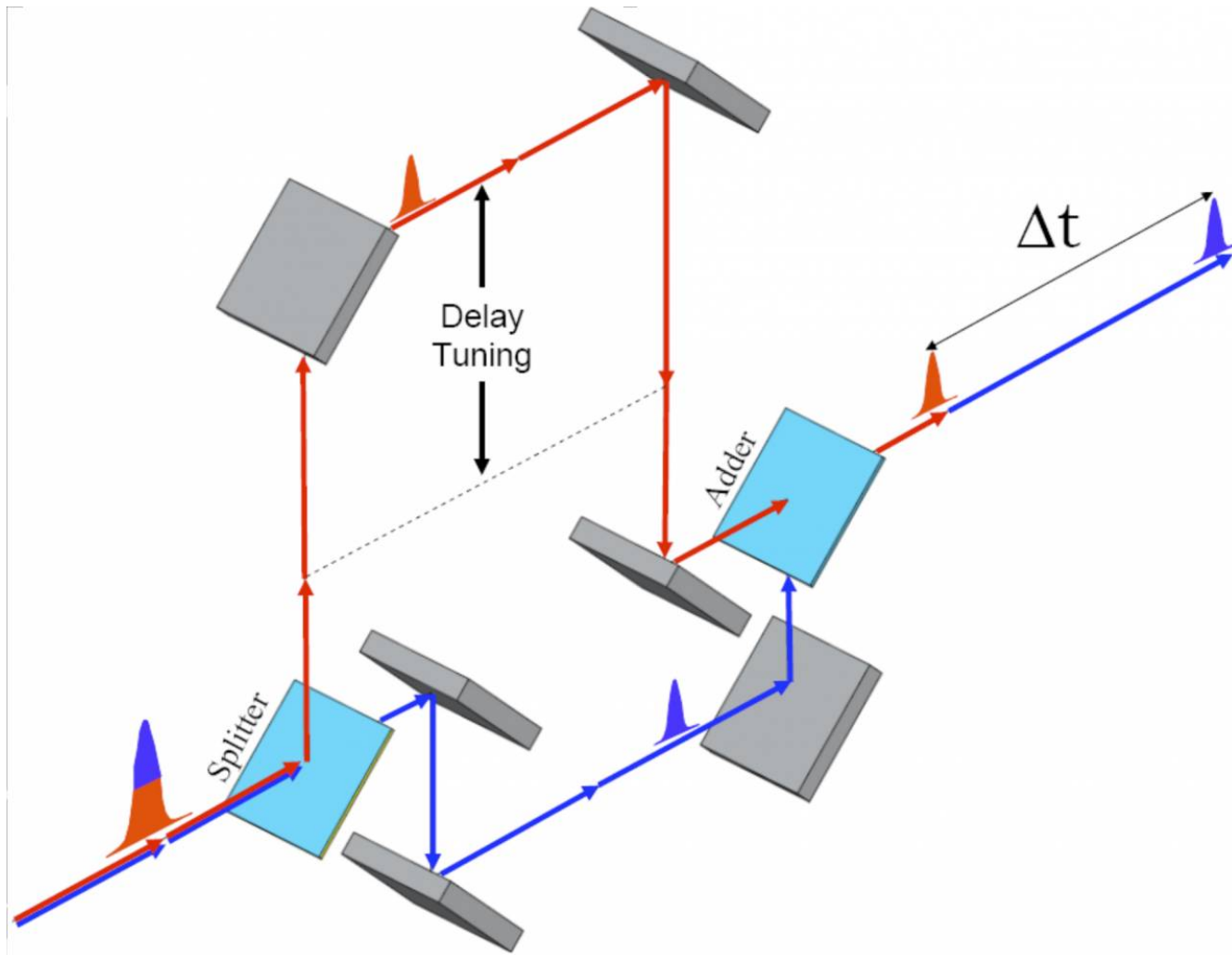


Serial mode

Temporal resolution depends on rep rate of the machine

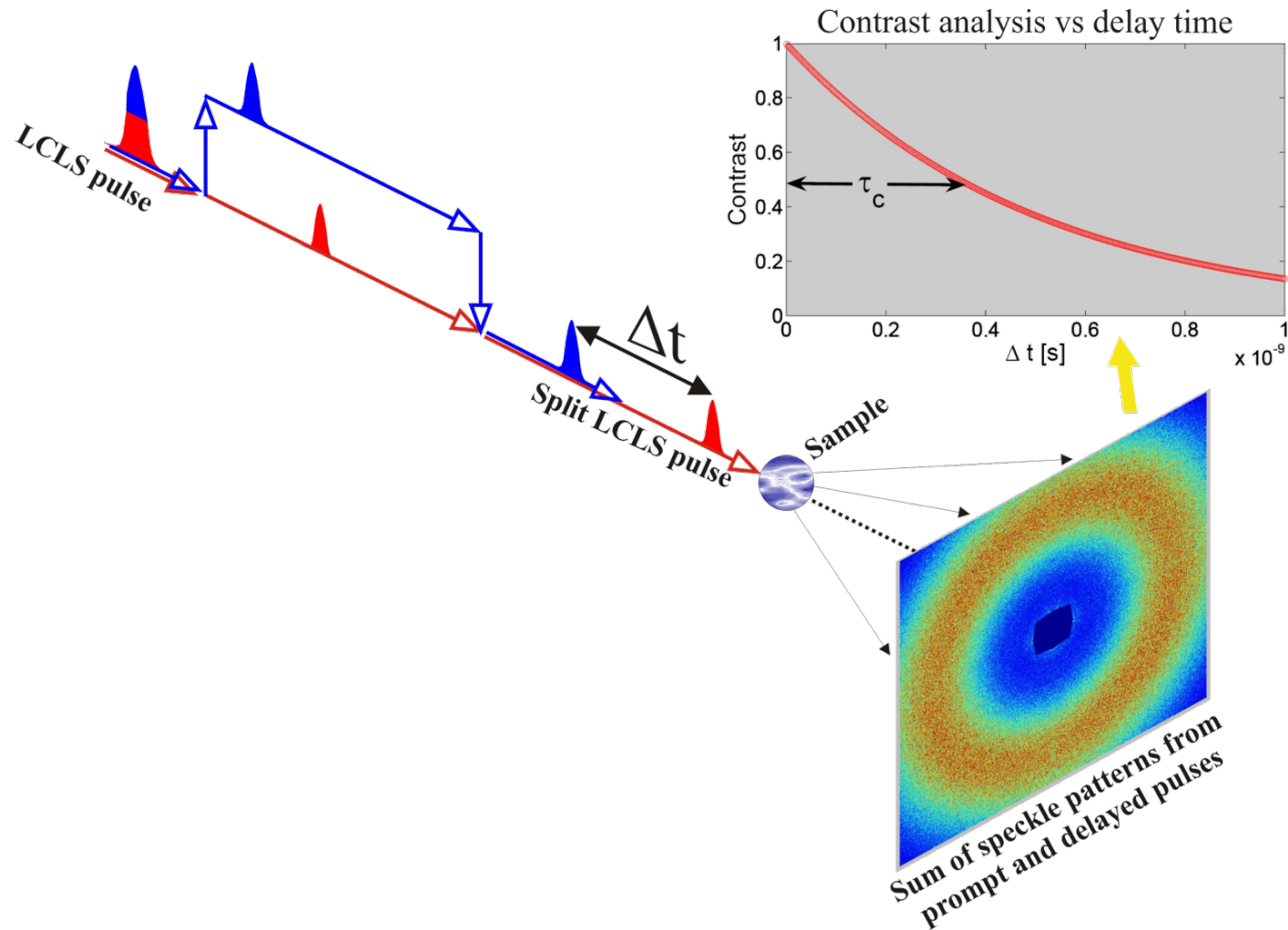


Ultrafast XPCS using a split and delay line



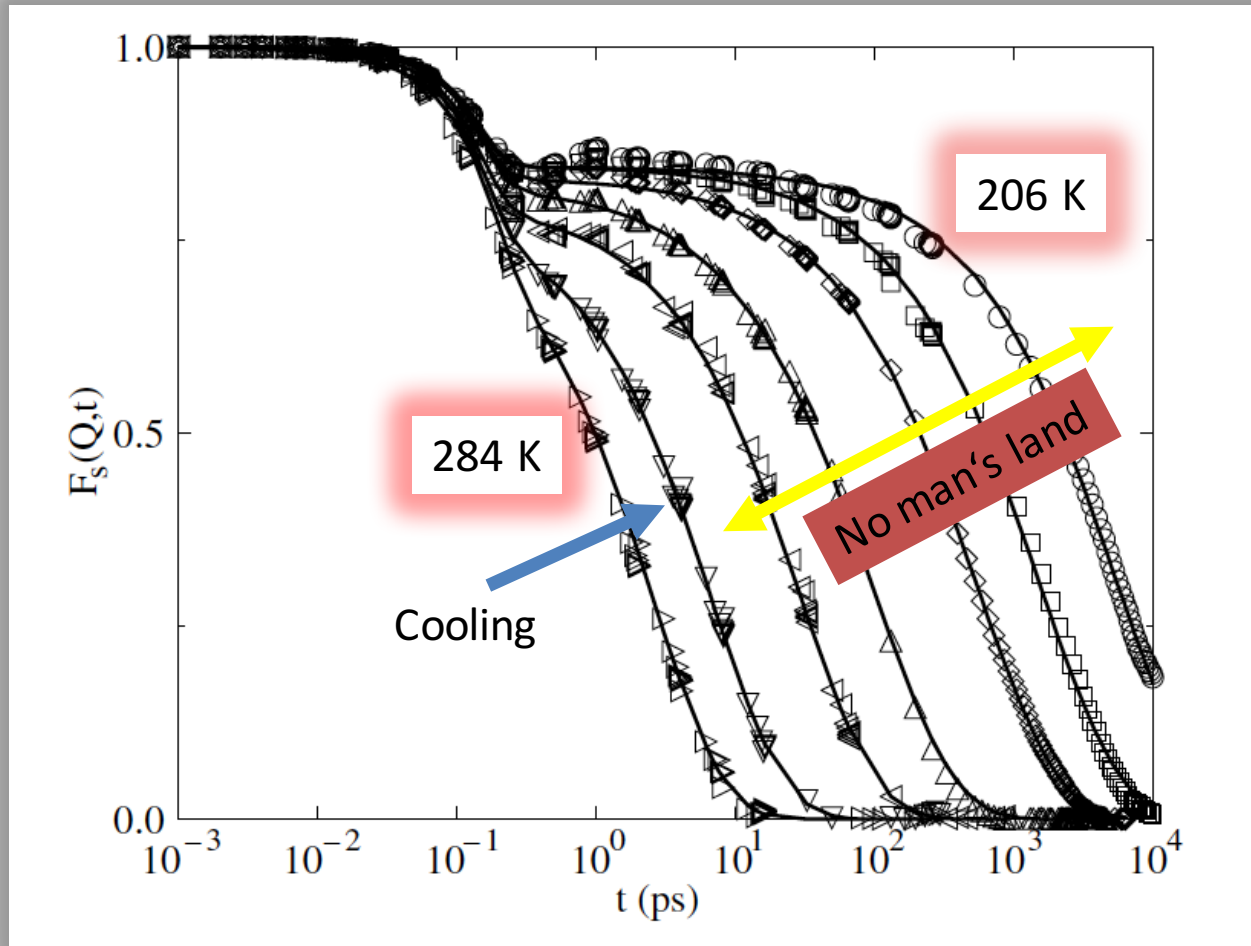
Delay times between 100 fs and 1 ns

Measure speckle contrast as a function of pulse separation

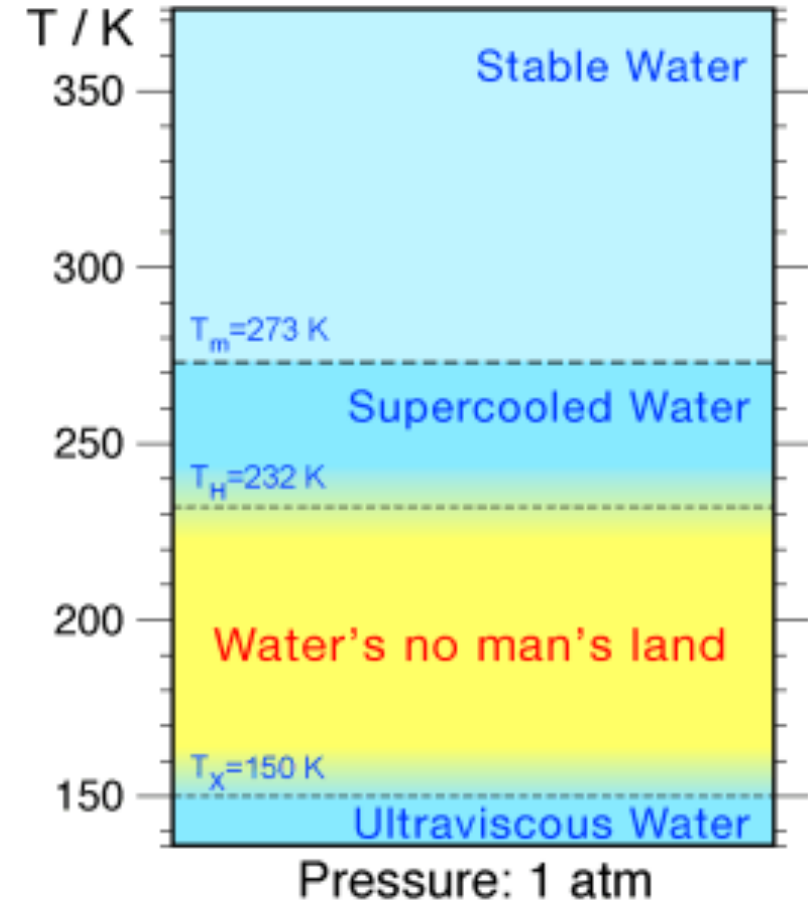


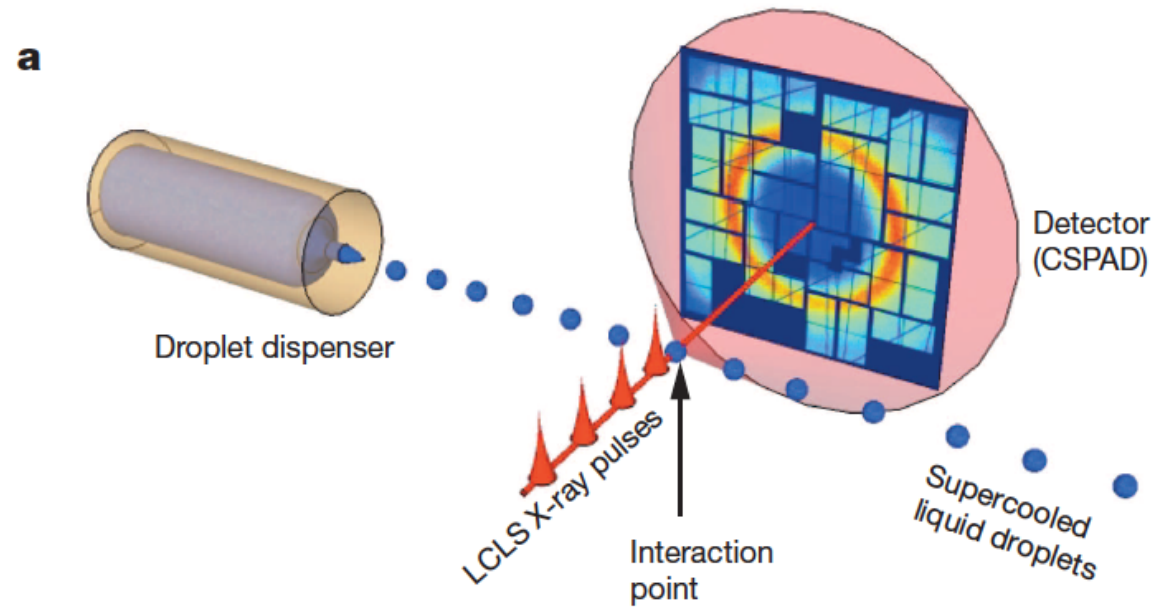
Ultrafast XPCS at XFEL - dynamics in extreme conditions

Calculated correlation function supercooled liquid water



Dynamics on time-scales ranging from 100 fs to 1000 ps

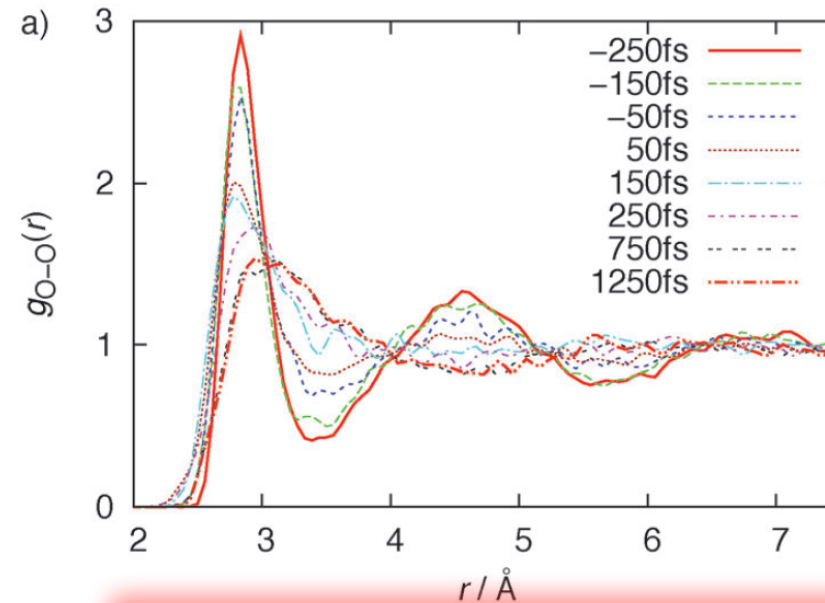
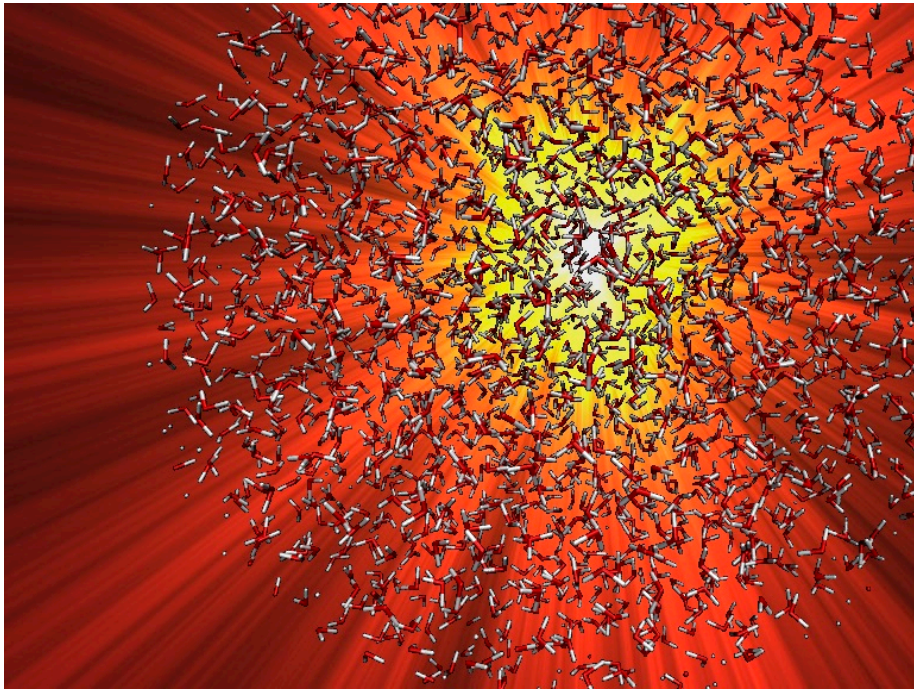




J.A. Sellberg et al. Nature 510, 381 (2014)

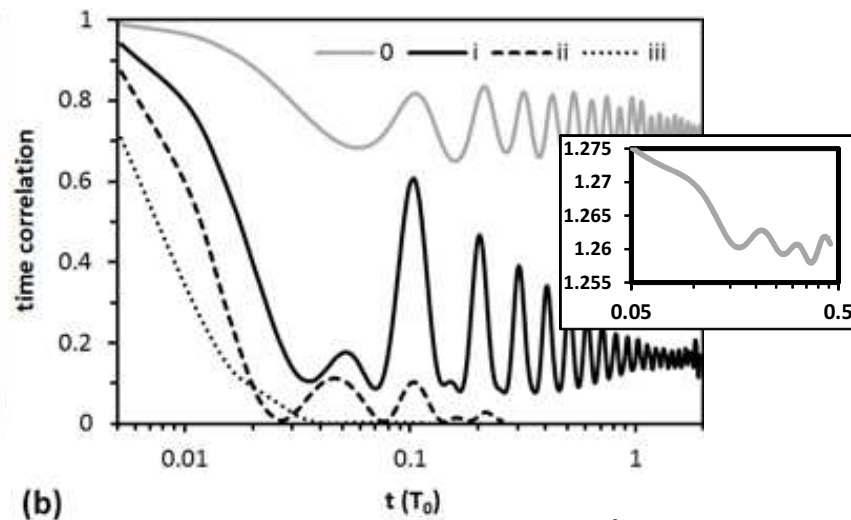
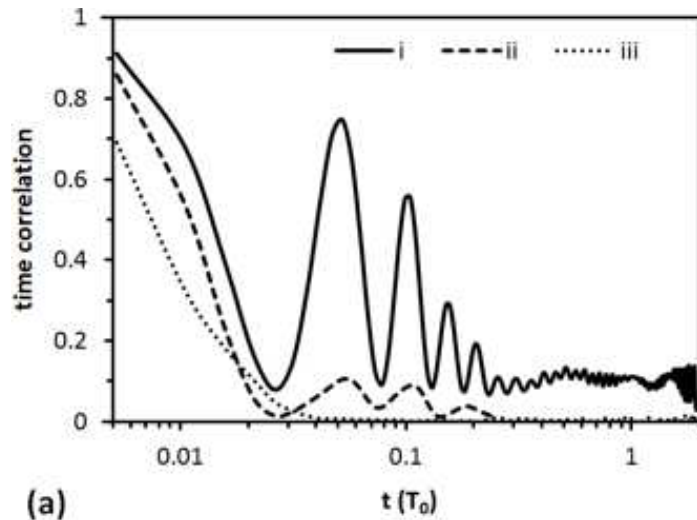
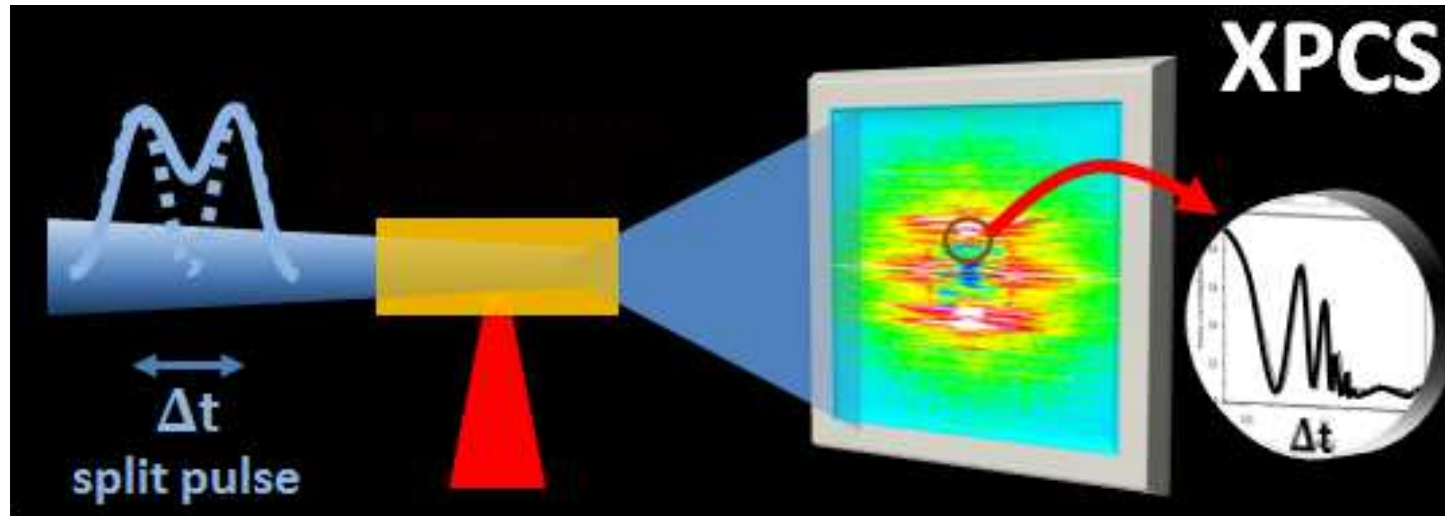
Ultrafast Energy Transfer to Liquid Water by Sub-Picosecond High-Intensity Terahertz Pulses: An Ab Initio Molecular Dynamics Study**

Pankaj Kr. Mishra, Oriol Vendrell,* and Robin Santra



Water at $T=1500 \text{ K}$, $p = 12 \text{ GPa}$ at least for a few ps

Pump-probe XPCS in Plasma Physics



The end