Probing high-scale physics using gravitational wave detectors

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BD and A. Mazumdar, Phys. Rev. D 93, 104001 (2016) [arXiv:1602.04203]

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A new window on the Universe



New Physics with GWs

"Hearing" things never seen before

- Three kinds of astrophysical sources:
 - Transient (e.g. compact binary inspirals)
 - Continuous (e.g. rapidly spinning neutron star)
 - Stochastic (e.g. superposition of unresolved sources)
- Stochastic signal also from cosmic events, e.g. inflation, cosmic strings, domain walls, phase transition.
- GWs can probe physics all the way up to the Planck epoch.

Phase Transition Basics



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First Order Phase Transition



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Bubble Nucleation





Energy Density



 In the envelope approximation [Caprini, Durrer, Servant (PRD '08); Huber, Konstandin (JCAP '08); Espinosa, Konstandin, No, Servant (JCAP '10); Weir '16],

$$rac{
ho_{
m GW}}{
ho_{
m tot}} \propto \kappa^2 v^3 \left(rac{lpha}{1+lpha}
ight)^2 \left(rac{H_*}{eta}
ight)^2.$$

In the strong first-order, thin-wall and vacuum-dominated limit:

$$\kappa \equiv \rho_v / \rho_{\rm vac} \to 1, \quad \alpha \equiv \rho_{\rm vac} / \rho_* \gg 1, \quad v \to 1.$$

- GW signal at T_* only depends on the nucleation rate $\beta/H_* \sim \log(m_{\text{Pl}}/T_*)$. [Kosowsky, Turner, Watkins (PRL '92); Kamionkowski, Kosowsky, Turner (PRD '94)]
- In realistic models with a given effective potential, typically $\beta/H_* \simeq 5/\epsilon \sim \mathcal{O}(100 1000)$. [Schwaller (PRL '15); Jaeckel, Khoze, Spannowsky '16]

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New Physics with GWs

GW Spectrum

$$\Omega_{\rm GW}(f)h^2 \equiv \frac{1}{
ho_c} \frac{d
ho_{\rm GW}}{d\log f} = \Omega_0 h^2 \, \frac{(p+q)\left(rac{f}{f_0}
ight)^p}{q+p\left(rac{f}{f_0}
ight)^{p+q}} \,,$$

with p = 2.8, q = 1.0, and the peak values are [Huber, Konstandin (JCAP '08)]

$$\begin{split} f_0 &\simeq (1.65 \times 10^{-7} \text{ Hz}) \left(\frac{0.62}{1.8 - 0.1 v + v^2} \right) \left(\frac{\beta}{H_*} \right) \left(\frac{T_*}{1 \text{ GeV}} \right) \left(\frac{g_*}{100} \right)^{1/6}, \\ \Omega_0 h^2 &\simeq (1.67 \times 10^{-5}) \kappa^2 \left(\frac{\alpha}{1 + \alpha} \right)^2 \left(\frac{0.11 v^3}{0.42 + v^2} \right) \left(\frac{H_*}{\beta} \right)^2 \left(\frac{100}{g_*} \right)^{1/3}, \end{split}$$

- $\Omega \propto f^{2.8}$ at low frequencies and f^{-1} at high frequencies.
- GW signal strength *decreases* with larger g_* .
- Need $T_* \sim 10^7 10^8$ GeV and $\beta/H_* \lesssim 100$ to be accessible at aLIGO.





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LIGO Sensitivity



BBH Background



Inflationary Spectrum (Blue-Tilted)



 Possible distinction between different stochastic GW signals using the frequency dependence:

$$\Omega_{\rm GW} \propto f^{2.8}$$
 (Phase Transition)
 $f^{2/3}$ (BBH)
 $f^{n_t(<0.36)}$ (Inflation)

• Feasible with the future worldwide GW network (LIGO+VIRGO+GEO+KAGRA+LIGO-India).

New Physics Scenarios

Toy model with two scalar fields (φ, χ):

$$V(\phi,\chi) = \frac{1}{4!}g^2(\phi^2 - v_*^2)^2 + \frac{1}{2}h\phi^2\chi^2.$$

• χ -field induces thermal corrections to the effective potential of ϕ :

$$V_T(\phi) = \frac{1}{24}h(T^2 - T_*^2)\phi^2 + \cdots$$
, where $T_* = \sqrt{\frac{2g}{h}}v_*$

- First-order phase transition for $T_*/v_* \leq 1$. [Jinno, Moroi, Nakayama (PLB '12)]
- Realistic examples: PQ axion, High-scale Supersymmetry.

Conclusion

- LIGO discovery has opened a new window on the Universe.
- Advanced LIGO design sensitivity can probe stochastic GW from cosmological phase transitions with $T_* \sim 10^7 10^8$ GeV.
- Distinct energy spectrum, as compared to other possible sources.
- Can be distinguished in future worldwide GW network.
- An unprecedented opportunity to constrain BSM physics at energy scales not directly accessible by laboratory experiments.