On the status of flavor anomalies

Diego Guadagnoli LAPTh Annecy (France)

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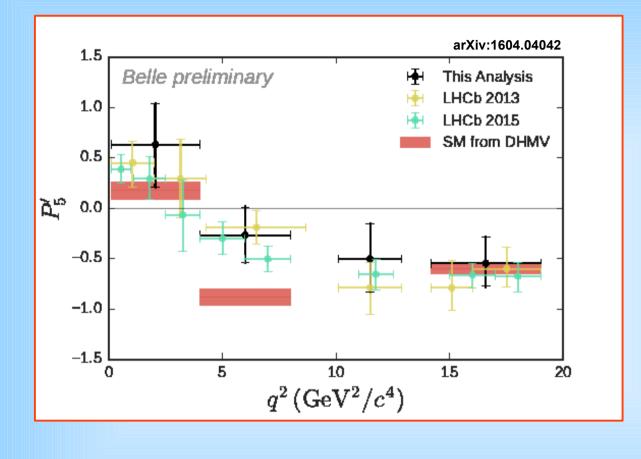
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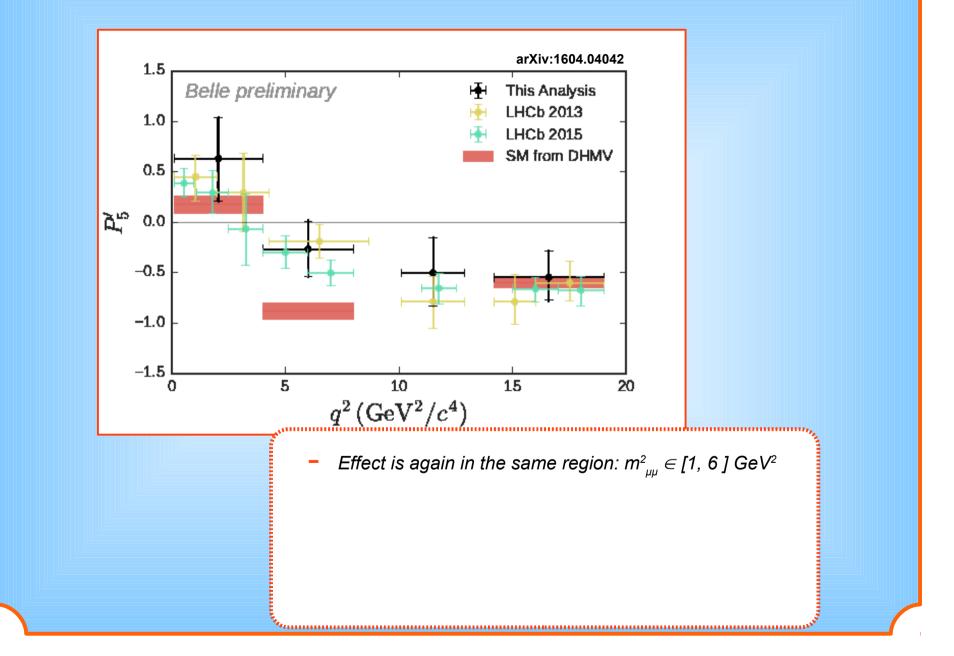
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- **B** \rightarrow **K**^{*} $\mu\mu$ angular analysis: discrepancy in one combination of the angular expansion coefficients, known as P'₅

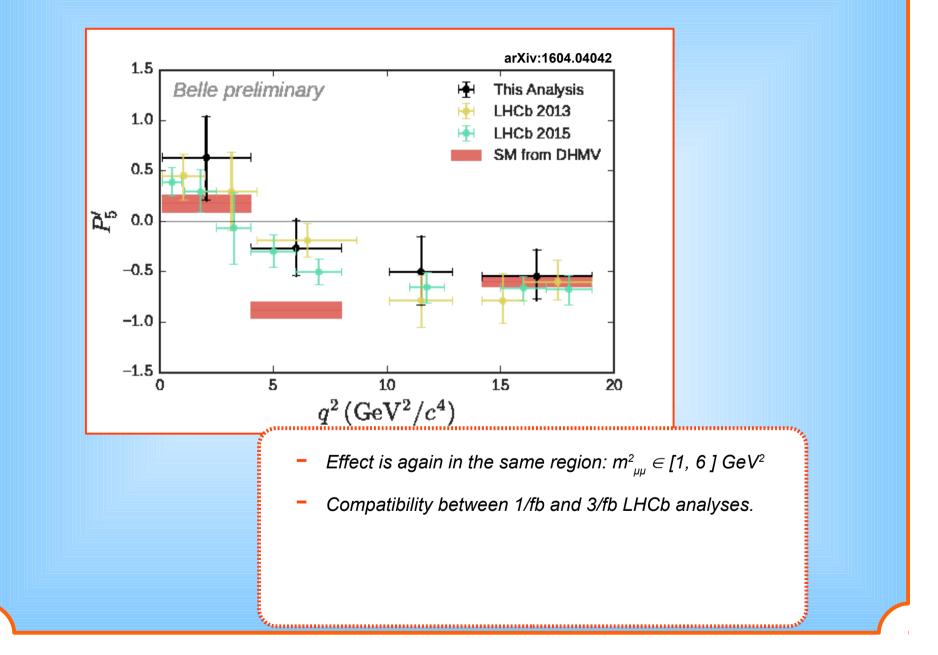
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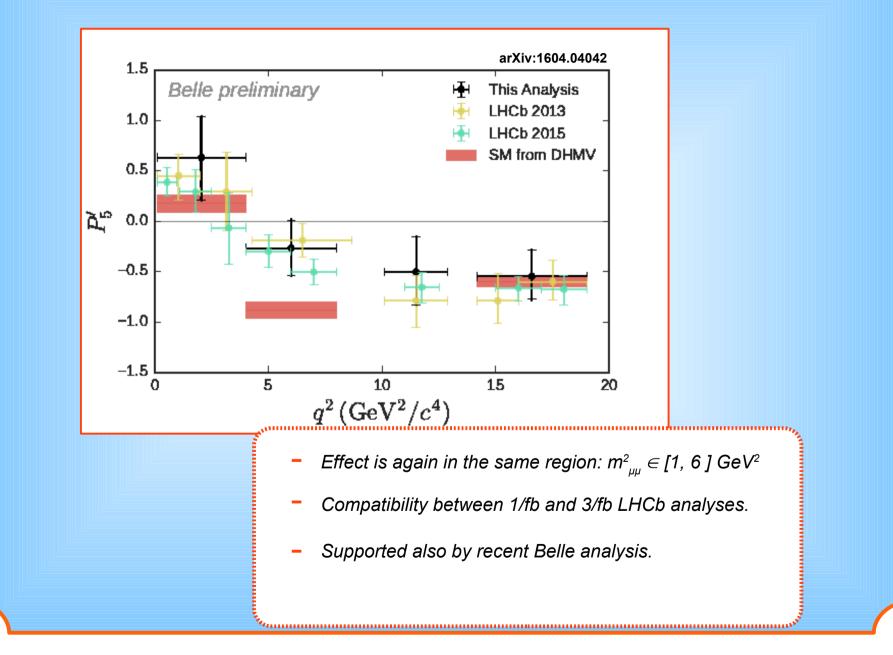
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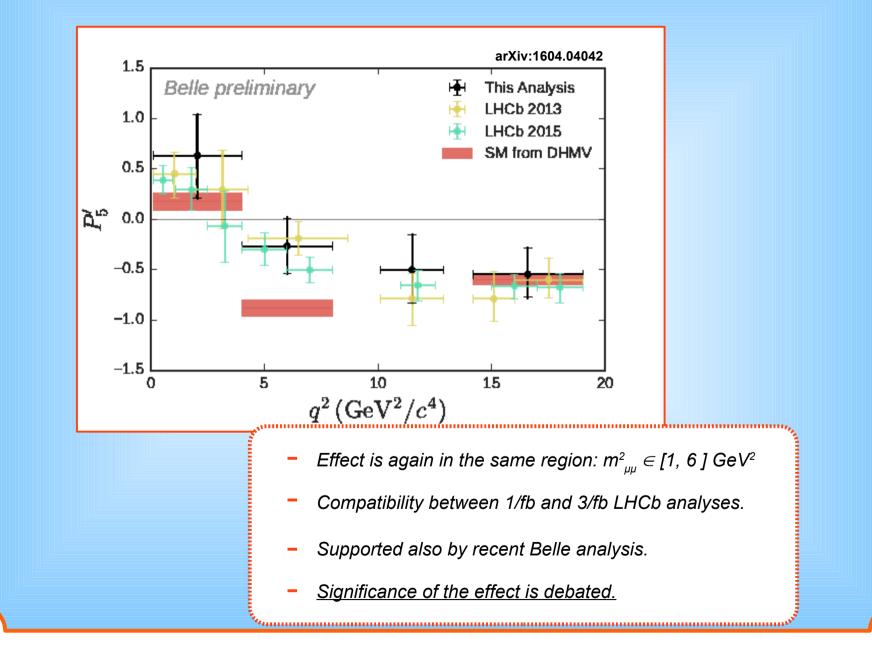
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LHCb and B factories measured several key $b \rightarrow s$ and $b \rightarrow c$ modes. Agreement with the SM is less than perfect.

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- **3** $B \rightarrow K^* \mu \mu$ angular analysis: discrepancy in P'_5 Again same region $m^2_{\mu\mu} \in [1, 6]$ GeV² Compatibility between 1/fb and 3/fb LHCb analyses. Supported also by recent Belle analysis. Significance of the effect is debated.

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There seems to be BSM LFNU and the effect is in $\mu\mu$, not ee

Recap of flavor anomalies: $\textbf{b} \rightarrow \textbf{c}$

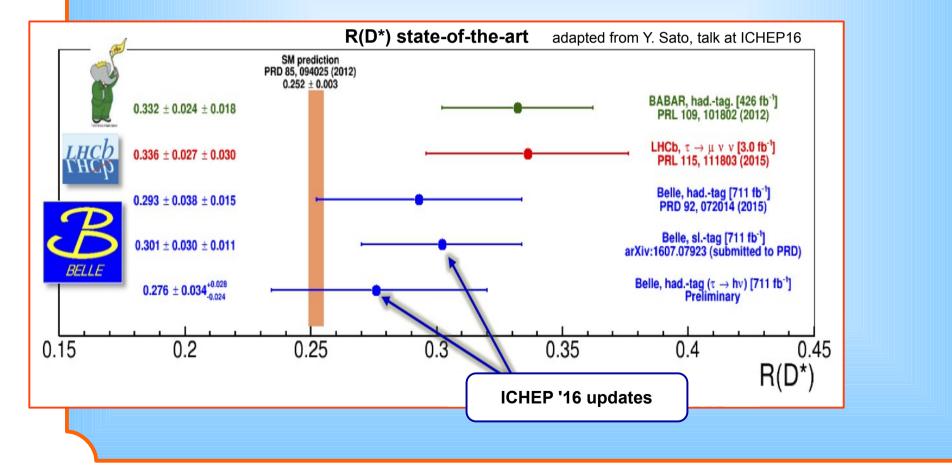
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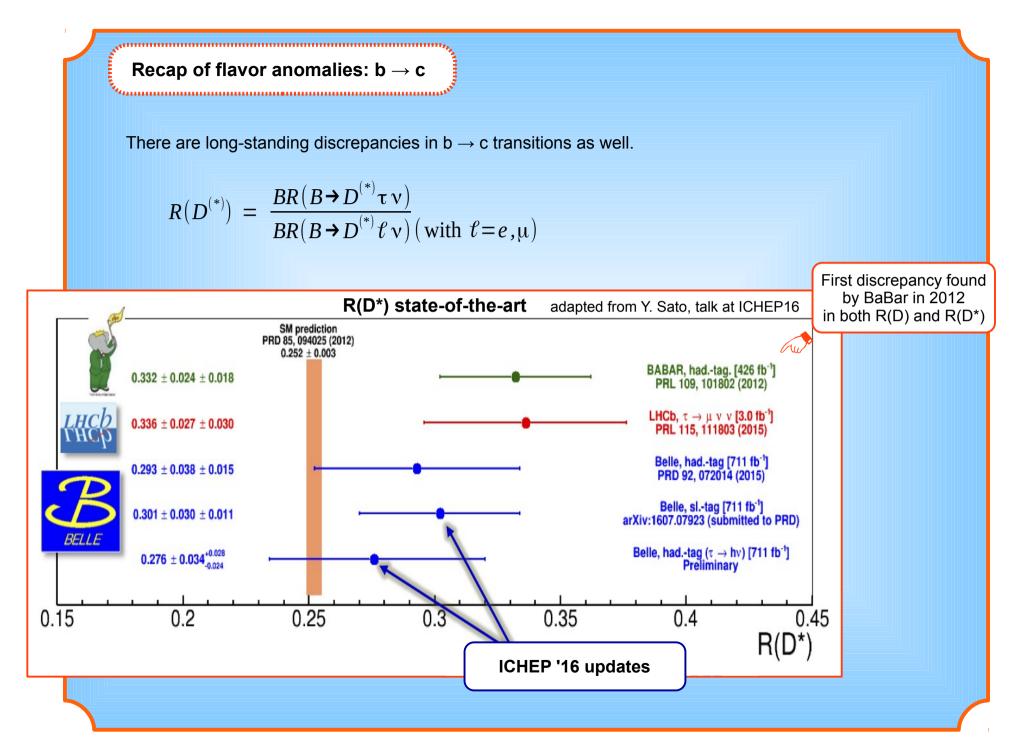
There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

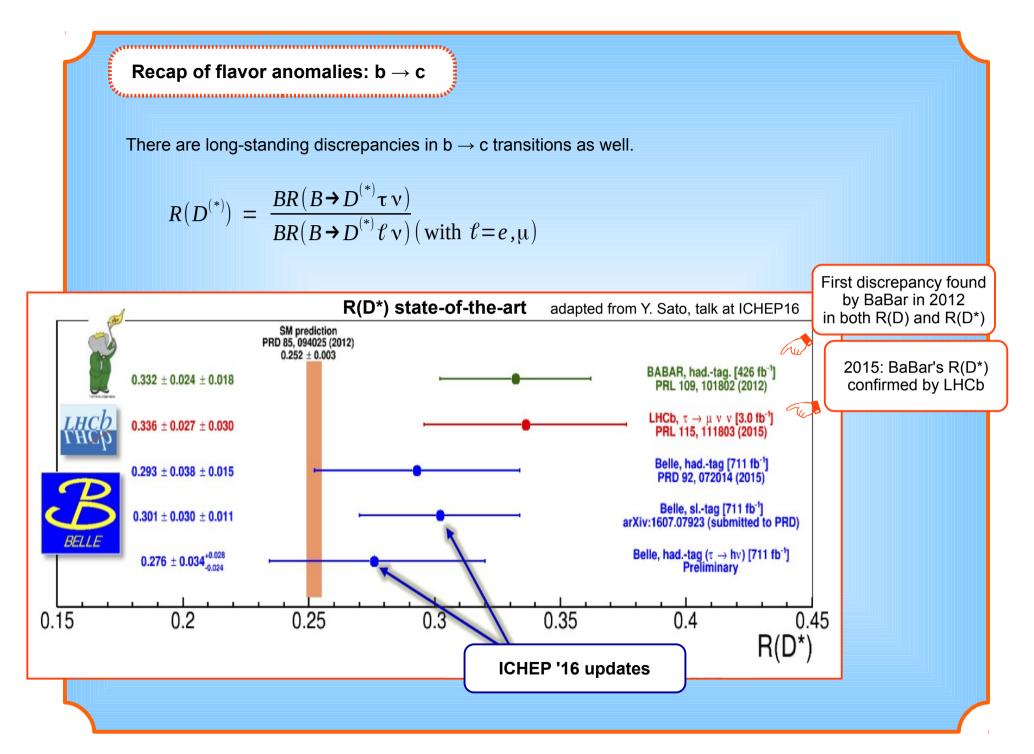
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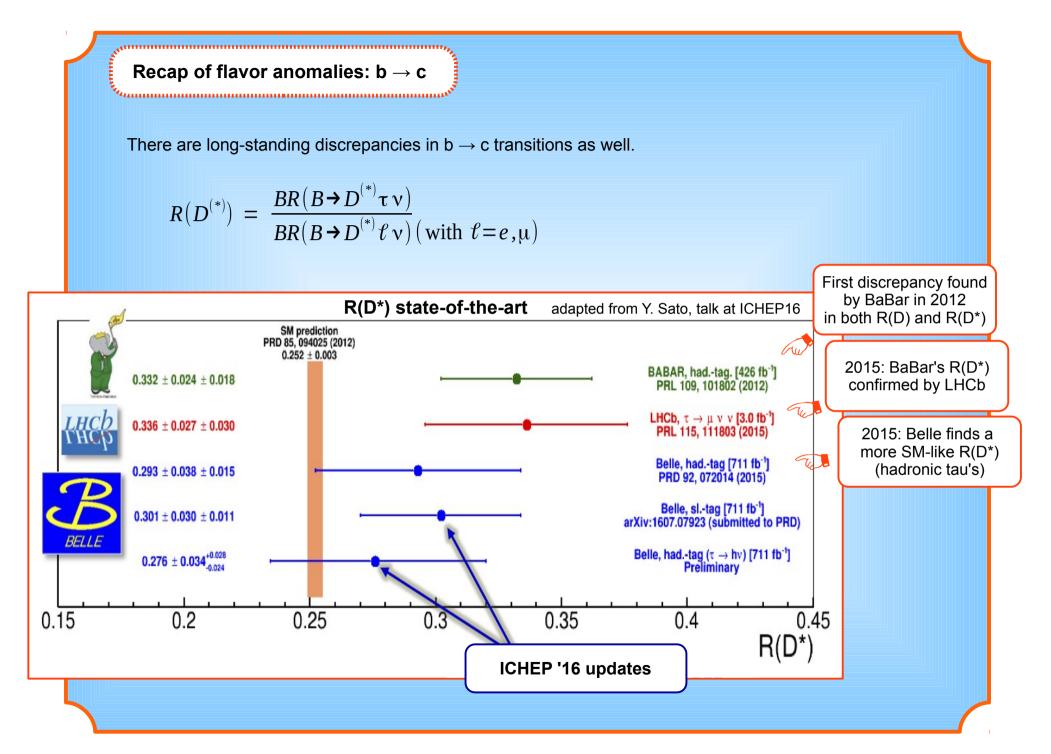
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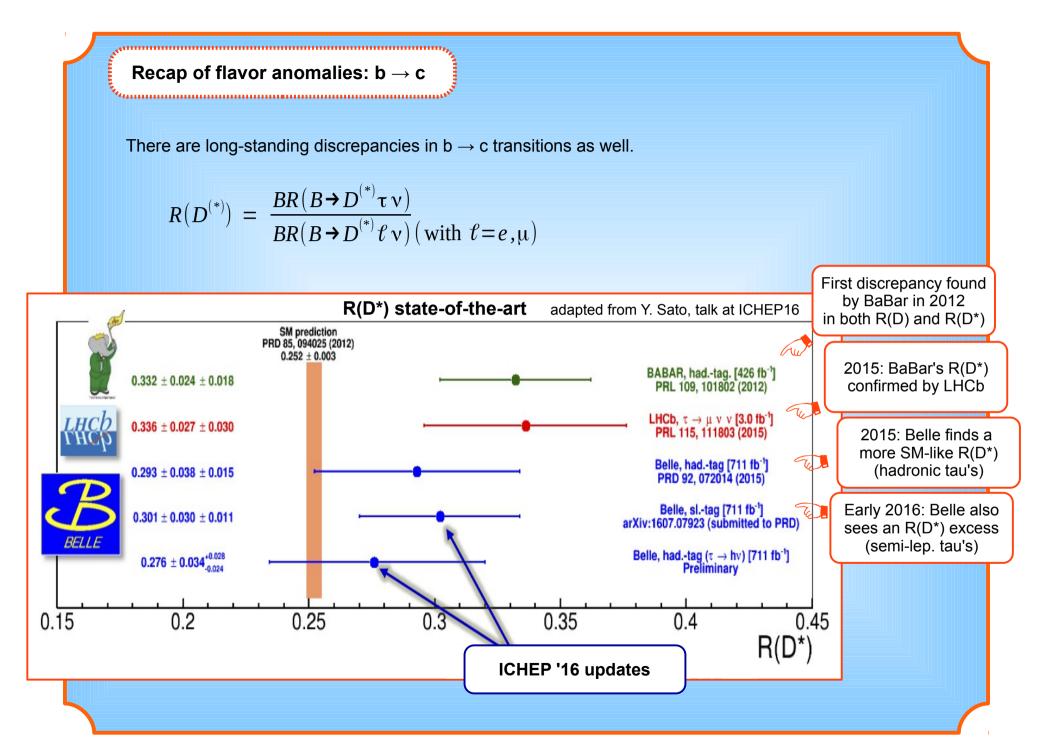
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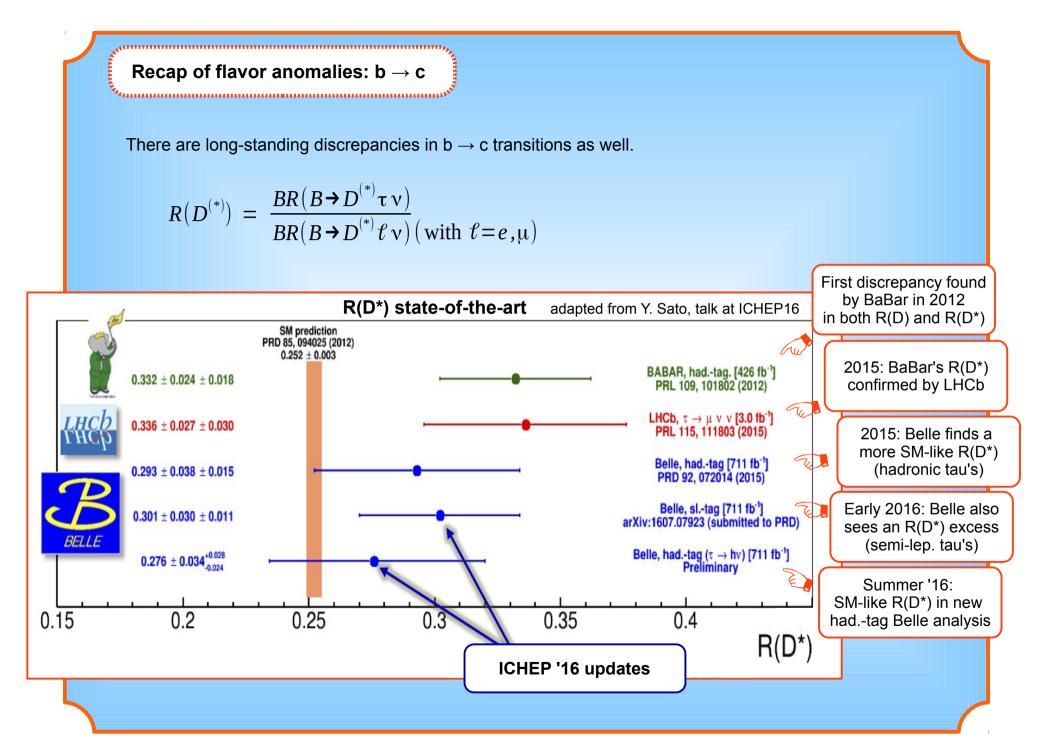


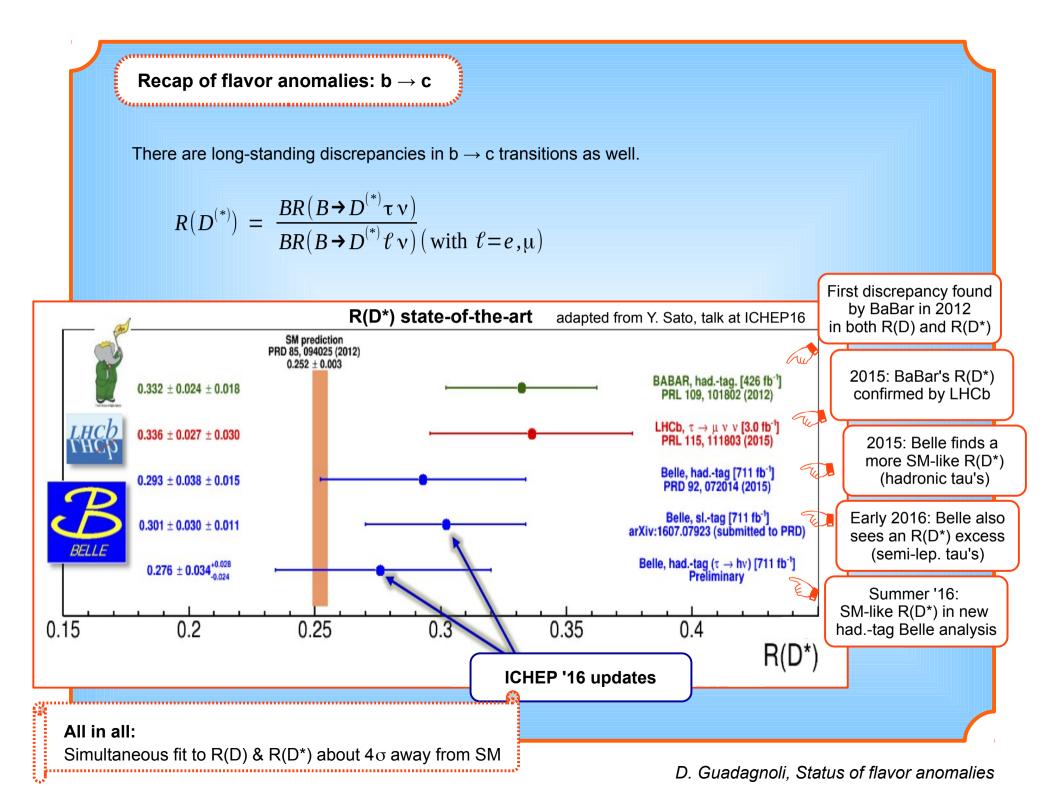


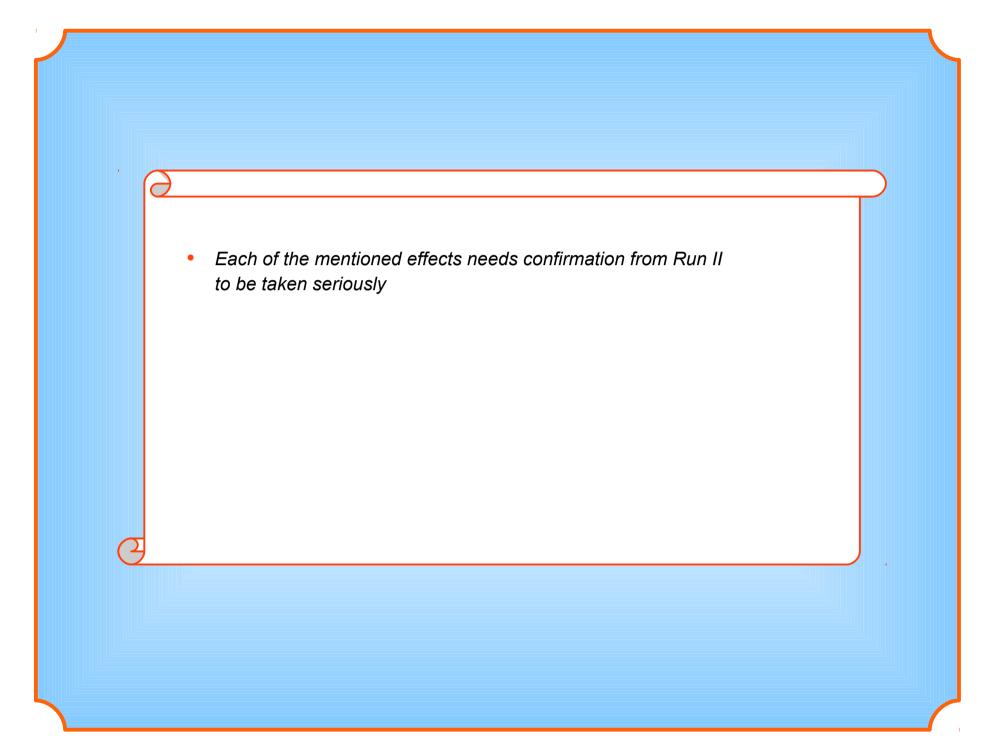












- Each of the mentioned effects needs confirmation from Run II to be taken seriously
- Yet, focusing for the moment on the $b \rightarrow s$ discrepancies
 - **Q1:** Can we (easily) make theoretical sense of data?
 - **Q2:** What are the most immediate signatures to expect ?

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Consider a new, LFNU interaction above the EWSB scale, e.g. with •

new vector bosons: $\overline{\ell} Z' \ell$ or leptoquarks: $\overline{\ell} \varphi q$

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 - (This basis doesn't yet even exist. We are above the EWSB scale.)
- Rotating q and t to the mass eigenbasis generates LFV interactions.

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• Physical LFV will appear in W couplings, but it's suppressed by powers of $(m_y / m_w)^2$

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So, BSM LFNU \implies BSM LFV (i.e. not suppressed by m_{y})

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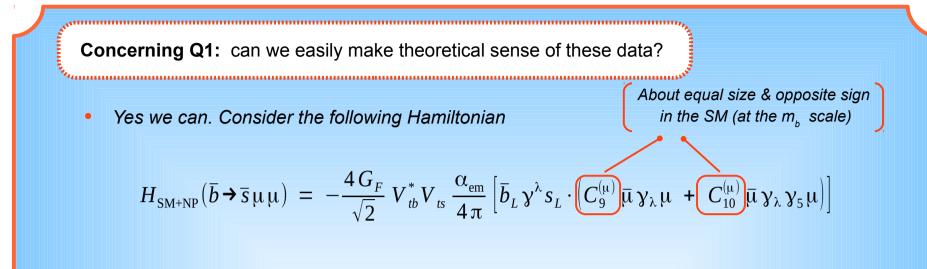
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• Yes we can. Consider the following Hamiltonian

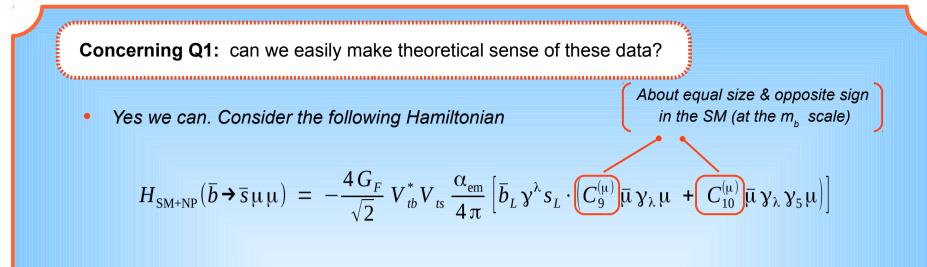
$$H_{\rm SM+NP}(\bar{b} \rightarrow \bar{s} \,\mu \,\mu) = -\frac{4 \,G_F}{\sqrt{2}} \,V_{tb}^* V_{ts} \,\frac{\alpha_{\rm em}}{4 \,\pi} \left[\bar{b}_L \,\gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \,\bar{\mu} \,\gamma_\lambda \mu + C_{10}^{(\mu)} \,\bar{\mu} \,\gamma_\lambda \gamma_5 \mu \right) \right]$$

Concerning Q1: can we easily make theoretical sense of these data:
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- Advocating the same $(V A) \times (V A)$ structure also for the corrections to $C_{9,10}^{SM}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_{κ} lower than 1
 - $B \rightarrow K \mu \mu \& B_s \rightarrow \mu \mu$ BR data below predictions
 - the P_5' anomaly in $B \rightarrow K^* \mu \mu$



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A fully quantitative test requires a global fit.

new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the χ^2 can be obtained either by a negative new physics contribution to C_9 (with $C_9^{\rm NP} \sim -30\% \times C_9^{\rm SM}$), or by new physics in the $SU(2)_L$ invariant direction $C_9^{\rm NP} = -C_{10}^{\rm NP}$, (with $C_9^{\rm NP} \sim -12\% \times C_9^{\rm SM}$). A positive NP contribution to C_{10} alone would also improve the fit, although to a lesser extent. [Altmannshofer, Straub, EPJC '15]

For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]

 As we saw before, all b → s data are explained at one stroke if:

- $C_{9}^{(t)} \approx -C_{10}^{(t)}$ (V - A) x (V - A) structure - $|C_{9,NP}^{(\mu)}| \gg |C_{9,NP}^{(e)}|$ LFNU

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- This pattern can be generated from a purely 3rd-generation interaction of the kind
 - $H_{\rm NP} = G \bar{b}'_{L} \gamma^{\lambda} b'_{L} \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$ with $G = 1/\Lambda_{\rm NP}^{2} \ll G_{F}$ expected e.g. in partial-compositeness frameworks

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- Note: primed fields
 - Fields are in the "gauge" basis (= primed)

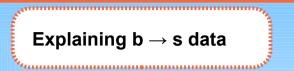
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$$mass \\ b'_{L} \equiv (d'_{L})_{3} = (U_{L}^{d})_{3i} (d_{L})_{i} \\ \tau'_{L} \equiv (\ell'_{L})_{3} = (U_{L}^{\ell})_{3i} (\ell_{L})_{i}$$

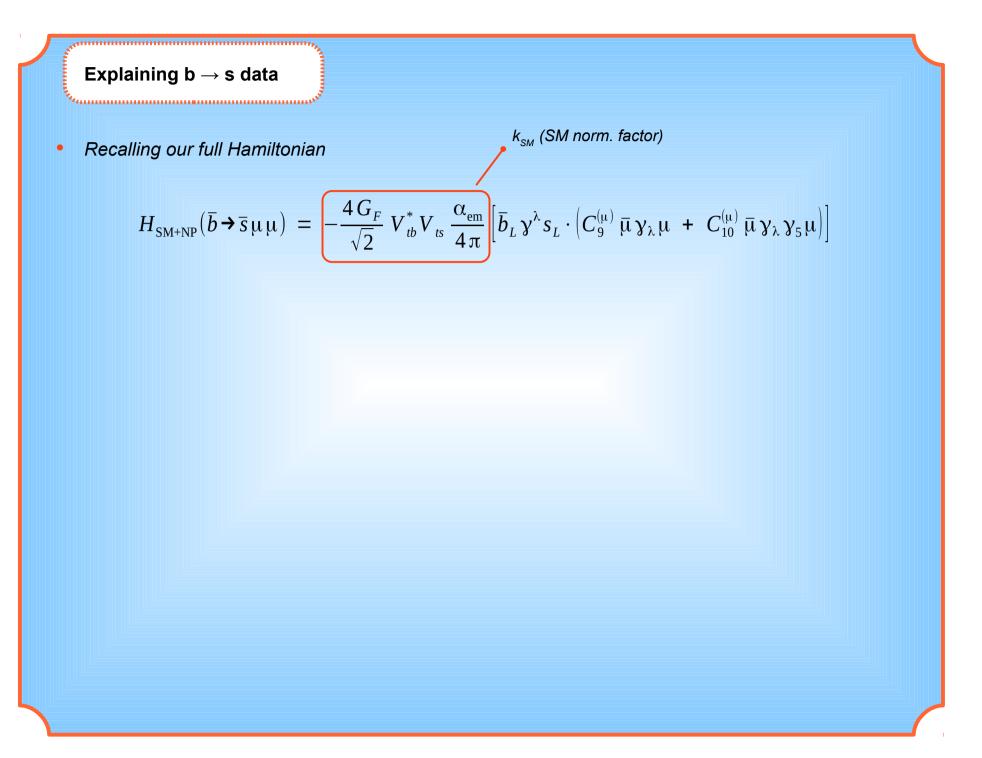
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 They need to be rotated to the mass eigenbasis
 This rotation induces <u>LFNU and LFV</u> effects
 b'_L = (d'_L)_3 = (U^d_L)_3 (d_L)_i (t_L)_i



• Recalling our full Hamiltonian

$$H_{\rm SM+NP}(\bar{b} \rightarrow \bar{s}\mu\mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\rm em}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$



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the shift to the $C_{_9}$ Wilson coeff. in the $\mu\mu$ -channel becomes

$$k_{\rm SM} C_9^{(\mu)} = k_{\rm SM} C_{9,\rm SM} + \frac{G}{2} (U_L^d)_{33}^* (U_L^d)_{32} |(U_L^\ell)_{32}|^2$$

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$$k_{\rm SM} C_9^{(e)} = k_{\rm SM} C_{9,\rm SM} + \frac{G}{2} (U_L^d)_{33}^* (U_L^d)_{32} |(U_L^t)_{31}|^2$$

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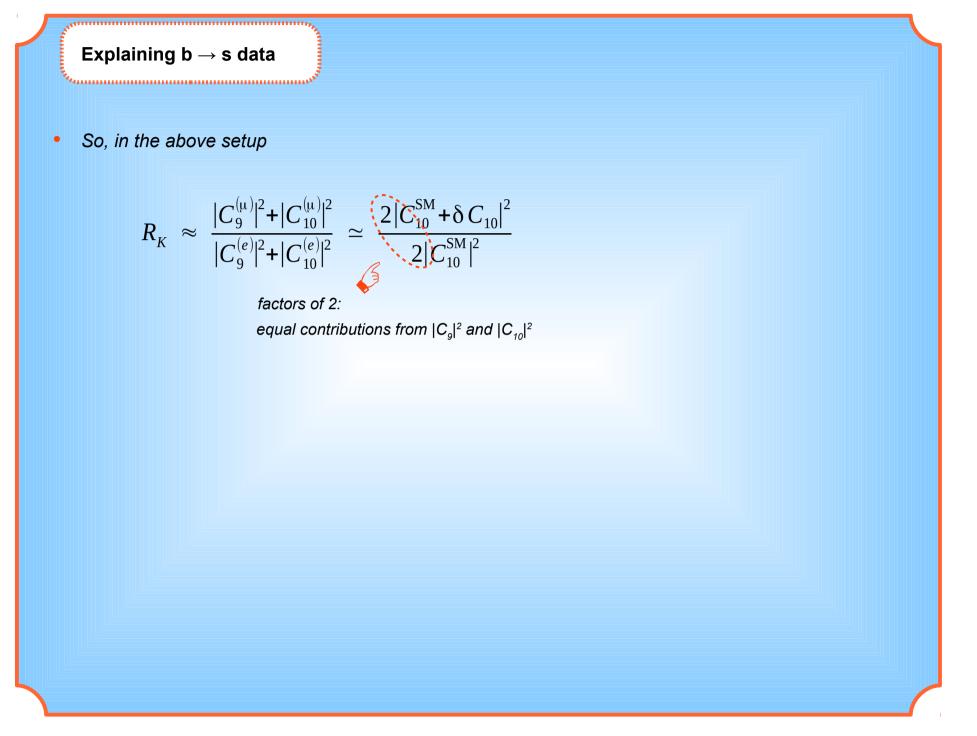
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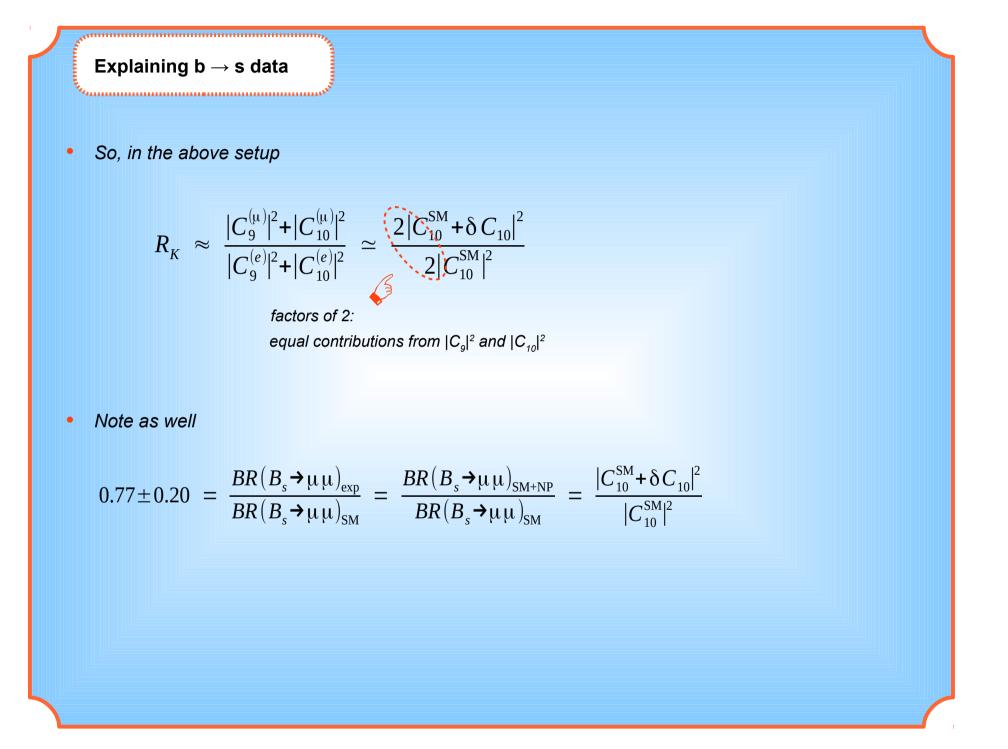
$$\left| U_{L,31}^{d} \right|^2 \ll \left| | U_{L,32}^{\ell} \right|^2$$

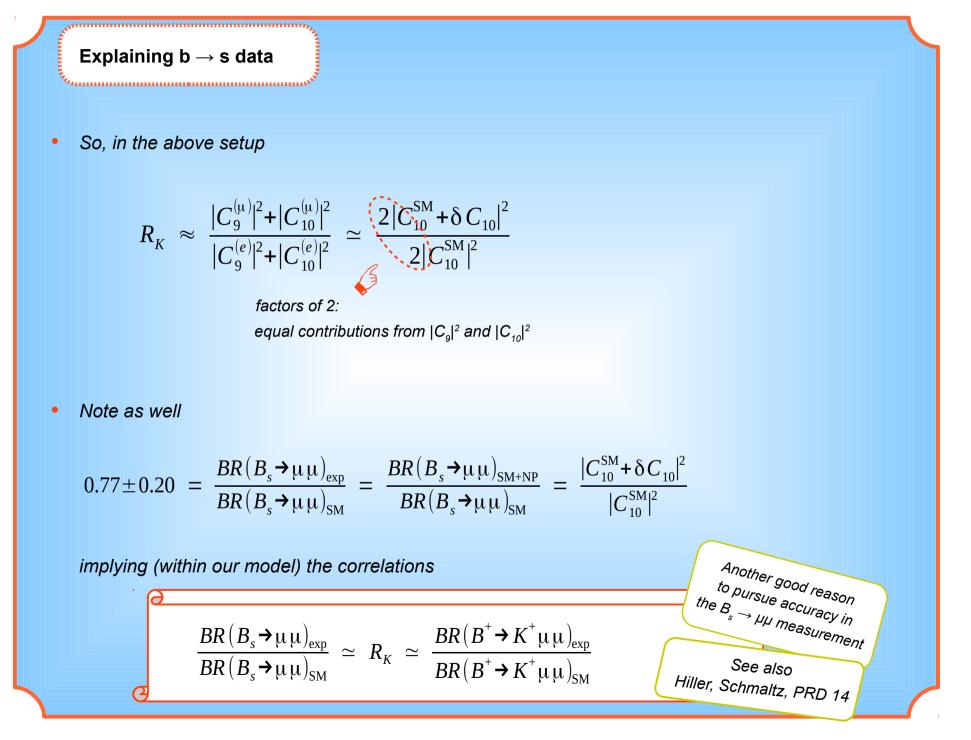


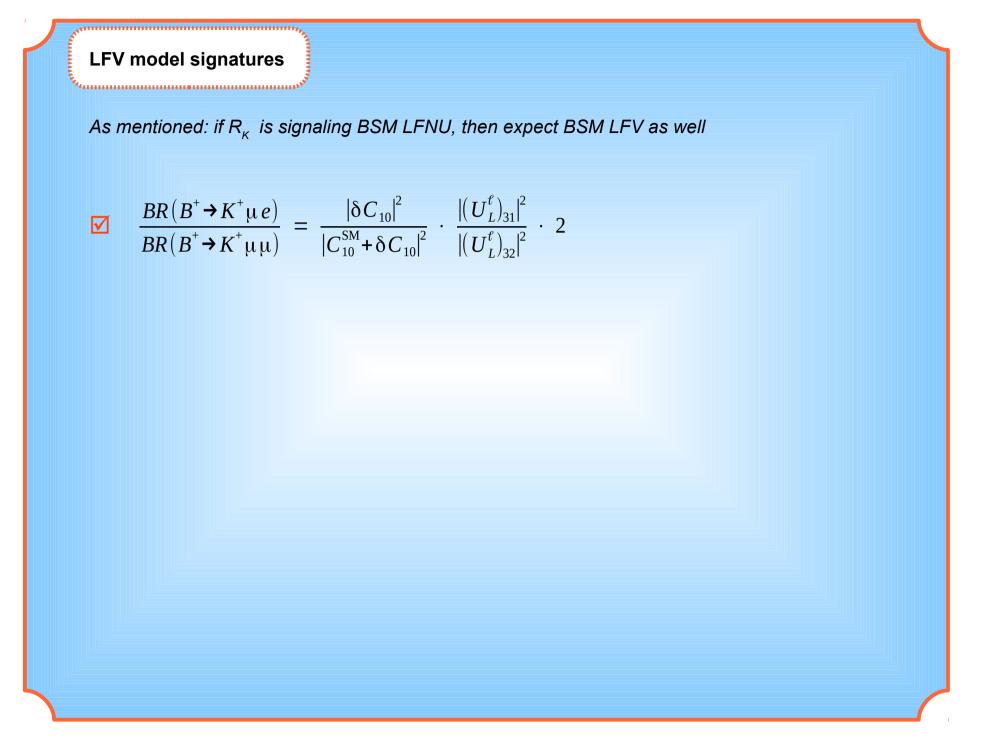
• So, in the above setup

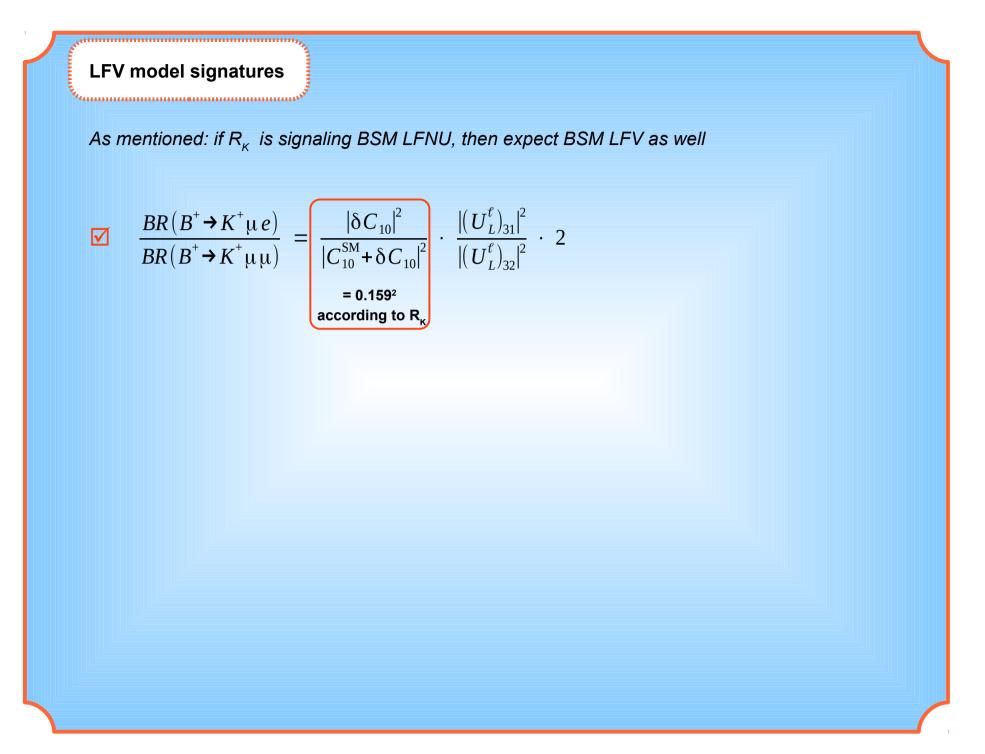
$$R_{K} \approx \frac{|C_{9}^{(\mu)}|^{2} + |C_{10}^{(\mu)}|^{2}}{|C_{9}^{(e)}|^{2} + |C_{10}^{(e)}|^{2}} \simeq \frac{2|C_{10}^{\text{SM}} + \delta C_{10}|^{2}}{2|C_{10}^{\text{SM}}|^{2}}$$

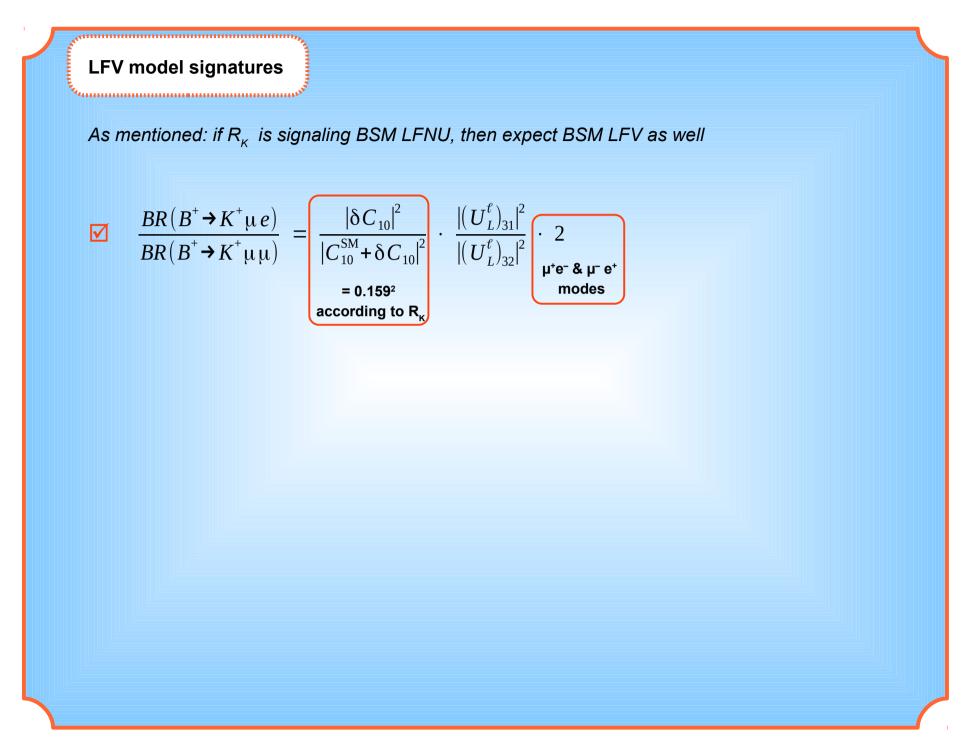


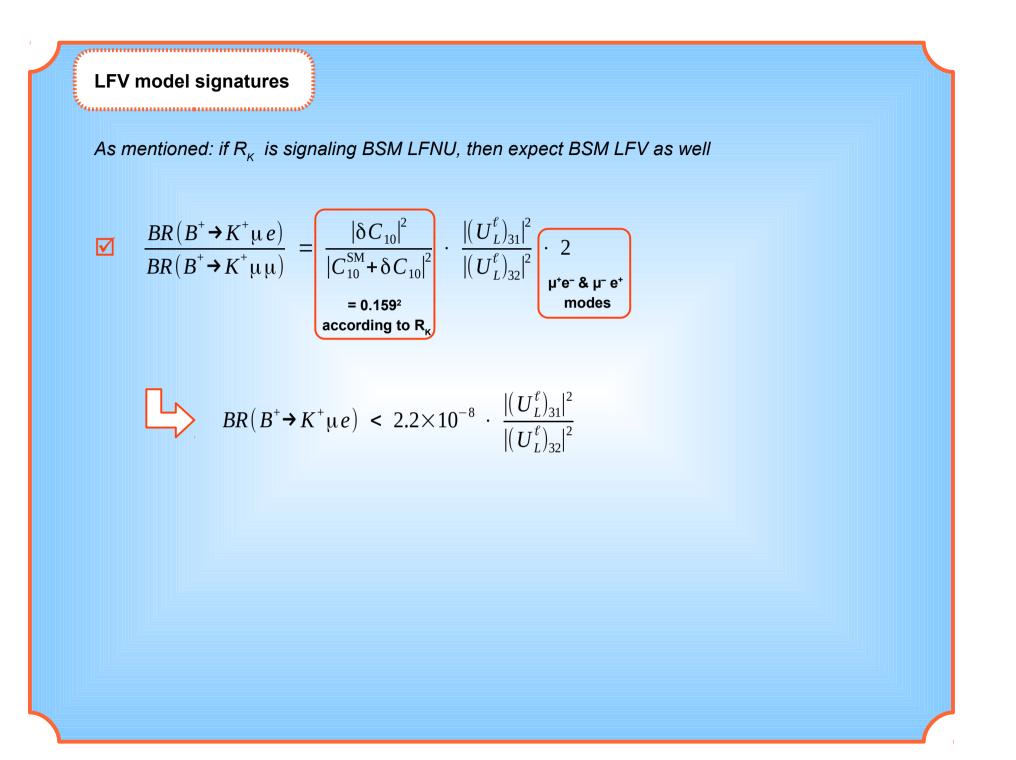


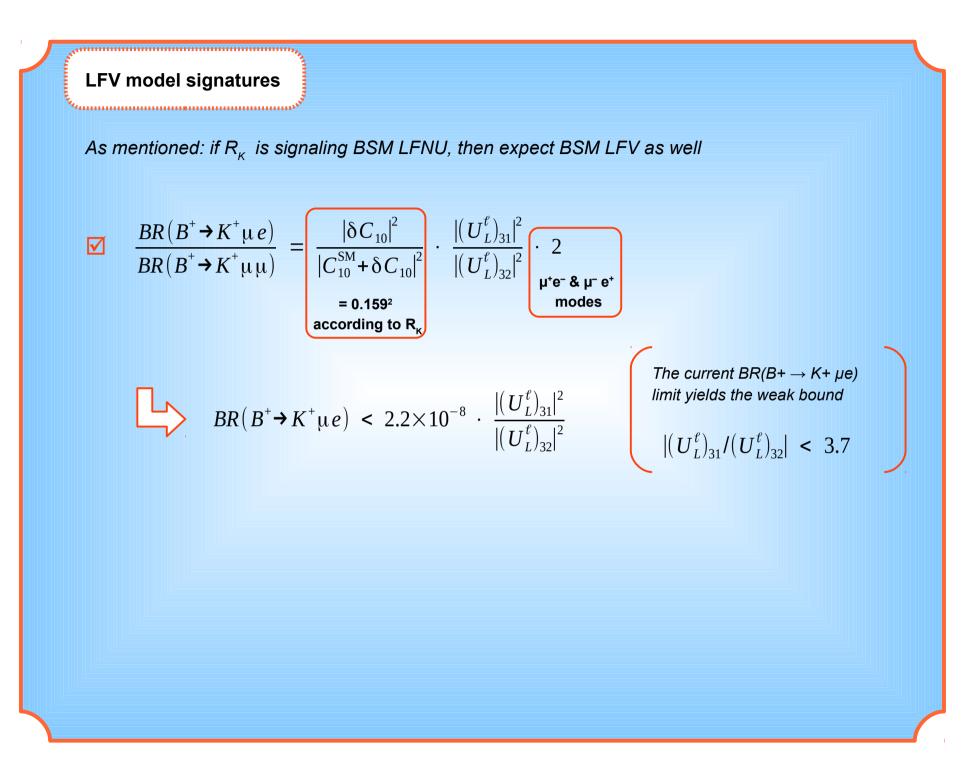


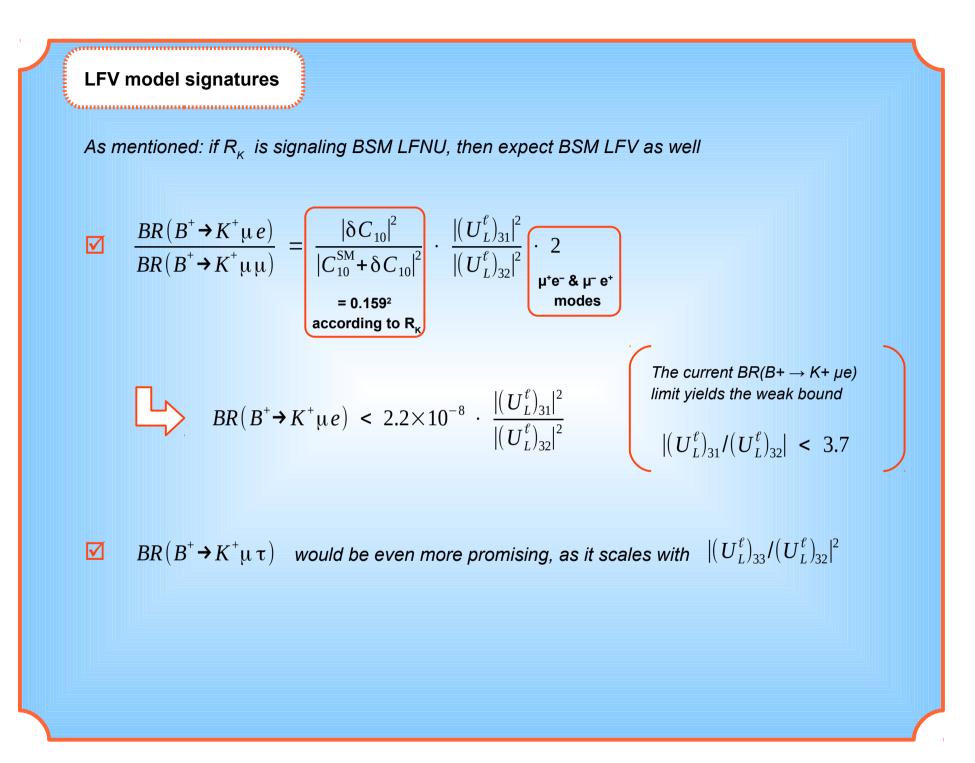


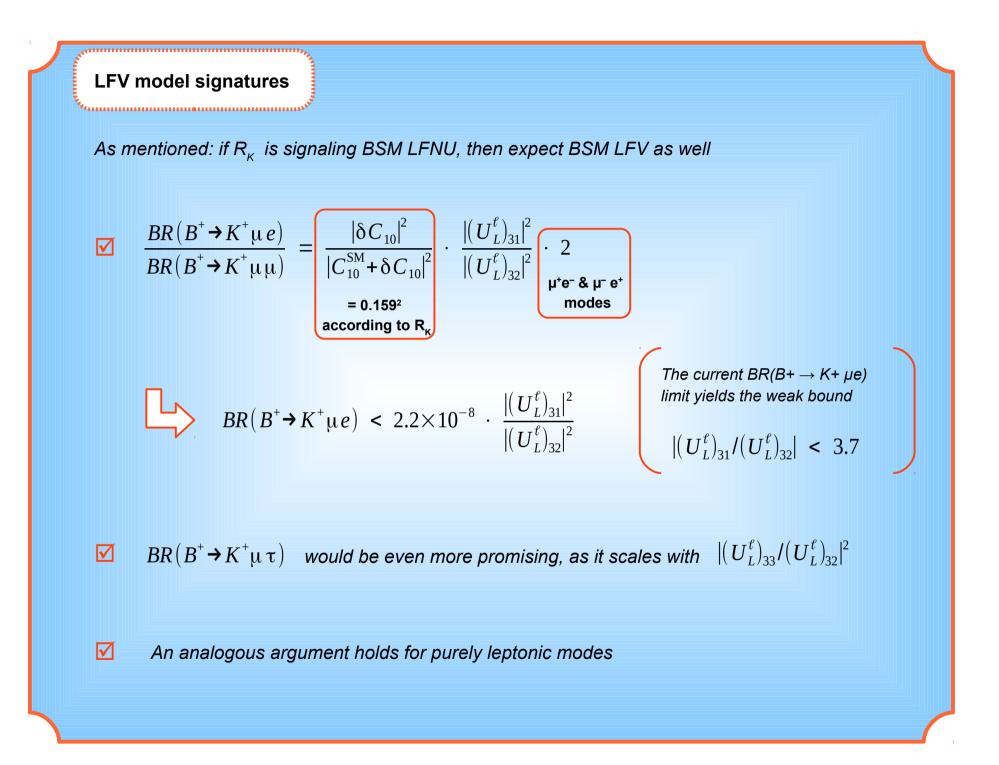


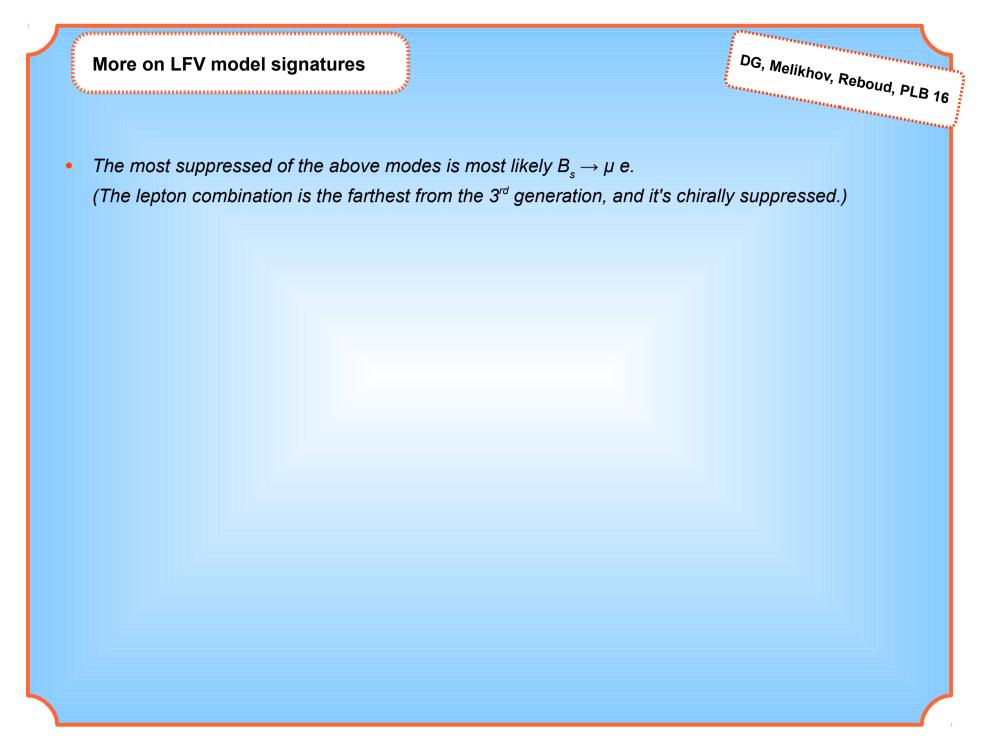




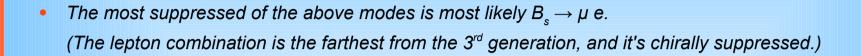










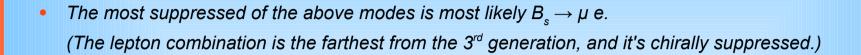


- What about $B_s \rightarrow \mu e \gamma$? •
 - γ = "hard" photon -

(hard = outside of the di-lepton Invariant-mass signal window)

DG, Melikhov, Reboud, PLB 16





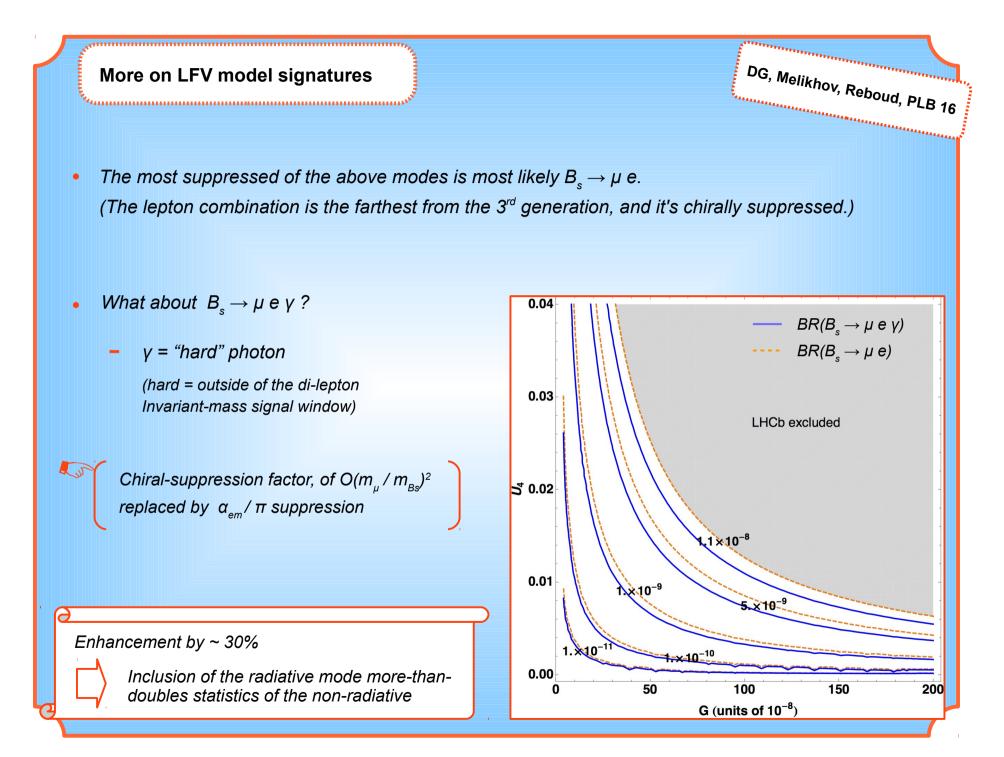
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Chiral-suppression factor, of $O(m_{\mu}^{}/m_{Bs}^{})^{2}$ replaced by $\alpha_{em}^{}/\pi$ suppression

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DG, Melikhov, Reboud, PLB 16



Being defined above the EWSB scale, • our assumed operator

 $G \bar{b}'_{L} \gamma^{\lambda} b'_{L} \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$

must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$



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must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$

Thus, the generated structures are all of:

$$t't'\nu'_{\tau}\nu'_{\tau}, \quad b'b'\nu'_{\tau}\nu'_{\tau},$$
$$t't'\tau'\tau', \quad b'b'\tau'\tau',$$

 $\left\{ \begin{array}{c} \bullet \quad \bar{Q}'_{L} \gamma^{\lambda} Q'_{L} \quad \bar{L}'_{L} \gamma_{\lambda} L'_{L} \\ \bullet \quad \bar{Q}'_{L}^{i} \gamma^{\lambda} Q'_{L}^{j} \quad \bar{L}'_{L}^{j} \gamma_{\lambda} L'_{L}^{i} \end{array} \right.$

SU(2)_L

inv.

[neutral-current int's only]

[also charged-current int's]

 Being defined above the EWSB scale, our assumed operator

$$G \ \bar{b}'_L \gamma^{\lambda} b'_L \ \bar{\tau}'_L \gamma_{\lambda} \tau'_L$$

must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$

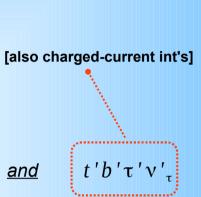
Thus, the generated structures are all of:

$$t't'v'_{\tau}v'_{\tau}, \quad b'b'v'_{\tau}v'_{\tau},$$
$$t't'\tau'\tau', \quad b'b'\tau'\tau'$$

 $\left\{ \begin{array}{c} \bullet \quad \bar{Q}'_{L} \gamma^{\lambda} Q'_{L} \quad \bar{L}'_{L} \gamma_{\lambda} L'_{L} \\ \bullet \quad \bar{Q}'^{i}_{L} \gamma^{\lambda} Q'^{j}_{L} \quad \bar{L}'^{j}_{L} \gamma_{\lambda} L'^{i}_{L} \end{array} \right.$

SU(2)_L

inv.



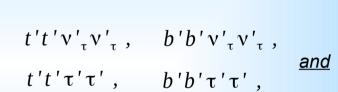
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Being defined above the EWSB scale, our assumed operator

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must actually be made invariant under $SU(3)_c \times SU(2)_1 \times U(1)_{\gamma}$

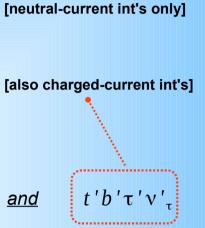
Thus, the generated structures are all of:



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SU(2)_L

inv.



See.

and a second second second

Bhattacharya, Datta, London,

Shivashankara, PLB 15



After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \ \tau \ v)$

i.e. it can explain deviations on R(D())*

Being defined above the EWSB scale, our assumed operator

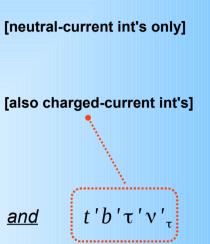
 $G \bar{b}'_{L} \gamma^{\lambda} b'_{L} \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$

must actually be made invariant under SU(3), x SU(2), x U(1),

Thus, the generated structures are all of:

 $\left\{ \bullet \ \bar{Q}'_L \gamma^{\lambda} Q'_L \bar{L}'_L \gamma_{\lambda} L'_L \right\}$ SU(2) • $\bar{Q}^{\prime i}_{L} \chi^{\lambda} Q^{\prime j}_{L} \bar{L}^{\prime j}_{L} \chi_{\lambda} L^{\prime i}_{L}$

 $t't'v'_{\tau}v'_{\tau}$, $b'b'v'_{\tau}v'_{\tau}$, $t't'\tau'\tau'$, $b'b'\tau'\tau'$,



and the second second

Bhattacharya, Datta, London,

Shivashankara, PLB 15

See

After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau v)$ i.e. it can explain deviations on R(D(*))

But this coin has a flip side.

Through RGE running, one gets also LFU-breaking effects in $\tau \rightarrow \ell v v$ (tested at per mil accuracy)

inv.

Such effects "strongly disfavour an explanation of the R(D(*)) anomaly model-independently"

D. Guadagnoli, Status of flavor anomalies

^{Feruglio,} Paradisi, Pattori, 2016





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 Their most convincing aspects are the following:
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- Early to draw conclusions. But Run II will provide a definite answer
- Timely to propose further tests. One promising direction is that of LFV. Plenty of channels, many of which largely untested.