

# Symplectic symmetries of $\text{AdS}_3$ phase space

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# Motivation

## Dynamics in 3d Einstein gravity

- No local dynamics in the bulk
- Chern-Simons formulation

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 0 \Leftrightarrow F = 0 = \bar{F}$$

- Many interesting solutions
- Boundary dynamics

## Surface Dynamics in 3d Einstein gravity

- Asymptotic symmetries
- Conserved charges uniquely label the solutions

### Extension into the bulk

- Boundary data  $\xrightarrow{\text{Einstein equations}}$  asymptotically AdS<sub>3</sub> solutions
- Asymptotic symmetries  $\xrightarrow{\text{gauge fixing}}$  symmetries everywhere in the bulk
- *symplectic symmetry*

# Symplectic geometry

- *Definition:*

*phase space* = a *manifold*  $\Gamma$  + a *symplectic form*  $\Omega$

- Properties of symplectic form  $\Omega$

**1** 2-form  $\Omega_{ab} = -\Omega_{ba}$

**2** Closed  $d\Omega = 0$

**3** Nondegenerate  $\det \Omega \neq 0$

- Symmetries  $\rightarrow$  charges

$$X^a = \Omega^{ab} \partial_b H_X$$

- Poisson bracket  $\rightarrow$  Algebra of charges

$$\{f, g\} = \Omega^{ab} \partial_a f \partial_b g$$

# Asymptotically AdS<sub>3</sub> Geometries

# Asymptotically AdS<sub>3</sub> geometries

- **Gauge Fixing.** Coordinates  $(z, x^+, x^-)$  in Fefferman-Graham gauge

$$ds^2 = \frac{1}{z^2} \left[ dz^2 + \gamma_{ij} dx^i dx^j \right]$$

$$\gamma_{ij} = g_{ij}^{(0)}(x^+, x^-) + g_{ij}^{sub}(z, x^+, x^-)$$

- **boundary conditions.** Brown-Henneaux

$$g^{(0)} = \eta = -dx^+ dx^- \quad \varphi \sim \varphi + 2\pi$$

## Asymptotically AdS<sub>3</sub> geometries

- Asymptotically AdS<sub>3</sub> geometries (*off-shell*)

$$z = 1/r \quad \rightarrow \quad ds^2 = \frac{dr^2}{r^2} + \left( r^2 \eta_{ij} + g_{ij}^{(sub)}(r, x^\pm) \right) dx^i dx^j$$

- Asymptotically AdS<sub>3</sub> geometries (*on-shell*) [Bañados(1998)]

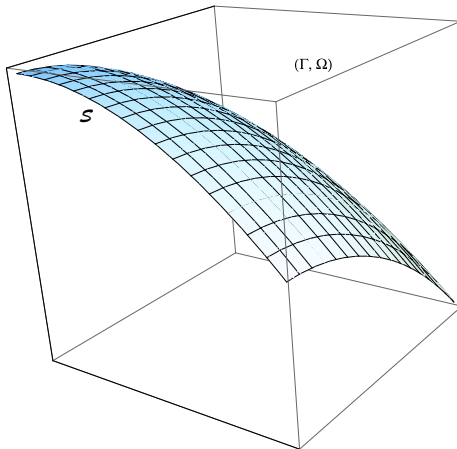
$$ds^2 = \ell^2 \frac{dr^2}{r^2} - \left( r dx^+ - \ell^2 \frac{L_-(x^-) dx^-}{r} \right) \left( r dx^- - \ell^2 \frac{L_+(x^+) dx^+}{r} \right)$$

# Phase Space of asymptotically AdS<sub>3</sub> Geometries



# Schematic picture of the phase space

- $\Gamma$ : The set of all asymptotically AdS<sub>3</sub> geometries (off-shell)
- $\mathcal{S}$ : The set of all asymptotically AdS<sub>3</sub> geometries (on-shell)



# Symplectic structure

[Lee, Wald(1990), Barnich, Brandt(2001)]

## Covariant phase space formulation

- Symplectic structure

$$\Omega(\delta_1\phi, \delta_2\phi) = \int_{\Sigma} \omega(\phi, \delta_1\phi, \delta_2\phi)$$

$\Sigma$  an arbitrary spacelike hypersurface

- The *symplectic current*  $\omega$  is obtained by the second variation of the Lagrangian (Lee-Wald 1990)

$$\begin{aligned} \delta\mathbf{L}[\Phi] &= \mathbf{E}[\Phi]\delta\Phi + d\Theta[\delta\Phi, \Phi] \\ \omega[\delta_1\Phi, \delta_2\Phi, \Phi] &= \delta_1\Theta[\delta_2\Phi, \Phi] - \delta_2\Theta[\delta_1\Phi, \Phi] \end{aligned}$$

- Alternative definition (Barnich-Brandt 2001)

# A strange property

[G.Compère, P.Mao, A.S and M.M.Sheikh-Jabbari, *JHEP* **1601**, 080 (2016).]

## Theorem

- Projection of symplectic form on the solution submanifold *vanishes*

$$\omega[\delta_1\Phi, \delta_2\Phi, \Phi] \Big|_{\mathcal{S}} = 0, \quad \forall \Phi \in \mathcal{S}, \quad \forall \delta_1\Phi, \delta_2\Phi \in T\mathcal{S}$$

- $\mathcal{S}$  is a *null* surface in  $\Gamma$  ( or an *isotropic* submanifold)
- This also happens in classical mechanics
- If  $\gamma$  is a first class constraint,

$$X^a = \Omega^{ab} \partial_b \gamma$$

is a null direction on  $\gamma = \text{const}$  surface

## Symmetries of AdS<sub>3</sub> phase space

# Symplectic symmetries of AdS<sub>3</sub> phase space

[G. Barnich, C. Troessaert (2010)]

[ G. Barnich, A. Gomberoff and H. A. Gonzalez, (2012)]

[G.Compère, P.Mao, A.S and M.M.Sheikh-Jabbari, (2015).]

- Diffeomorphism is the gauge symmetry of Einstein theory

$$g_{\mu\nu} \rightarrow g_{\mu\nu} + \mathcal{L}_\chi g_{\mu\nu}, \quad \forall \chi$$

- However, we have fixed the gauge  $g_{rr} = \frac{1}{r^2}$ ,  $g_{ri} = 0$
- Moreover, we have fixed a boundary condition  $g^{(0)} = -dx^+ dx^-$

## Residual gauge symmetries over solutions

$$\chi = -\frac{r}{2}(\epsilon'_+ + \epsilon'_-)\partial_r + \left(\epsilon_+ + \frac{\ell^2 r^2 \epsilon''_- + \ell^4 L_- \epsilon''_+}{2(r^4 - \ell^4 L_+ L_-)}\right)\partial_+ + \left(\epsilon_- + \frac{\ell^2 r^2 \epsilon''_+ + \ell^4 L_+ \epsilon''_-}{2(r^4 - \ell^4 L_+ L_-)}\right)\partial_-$$

- *Symplectic symmetry*  $\omega[\delta\Phi, \delta_\chi\Phi, \Phi] \Big|_S = 0$

# Symmetries of the phase space

$$\chi = -\frac{r}{2}(\epsilon'_+ + \epsilon'_-)\partial_r + \left(\epsilon_+ + \frac{\ell^2 r^2 \epsilon''_- + \ell^4 L_- \epsilon''_+}{2(r^4 - \ell^4 L_+ L_-)}\right)\partial_+ + \left(\epsilon_- + \frac{\ell^2 r^2 \epsilon''_+ + \ell^4 L_+ \epsilon''_-}{2(r^4 - \ell^4 L_+ L_-)}\right)\partial_-$$

- Extension of Brown-Henneaux asymptotic symmetries into the bulk

$$\chi^{(BH)} = -\frac{r}{2}(\epsilon'_+ + \epsilon'_-)\partial_r + \epsilon_+ \partial_+ + \epsilon_- \partial_-$$

- Transform solutions to solutions **everywhere**

$$g_{\mu\nu}(L) + \mathcal{L}_\chi g_{\mu\nu}(L) = g_{\mu\nu}(L + \delta_\epsilon L),$$

$$\delta_\epsilon L = \epsilon \partial L + 2L \partial \epsilon - \frac{1}{2} \partial^3 \epsilon$$

- $[\chi(\epsilon_1; L), \chi(\epsilon_2; L)]$  does not close !

## Algebra of residual symmetries

Adjusted bracket [Barnich and C. Troessaert (2010)]

The adjusted (modified) bracket  $[\cdot, \cdot]_*$

$$[\chi(\epsilon_1; L), \chi(\epsilon_2; L)]_* \equiv [\chi(\epsilon_1; L), \chi(\epsilon_2; L)]_{L.B} - \left( \delta_{\epsilon_1}^L \chi(\epsilon_2; L) - \delta_{\epsilon_2}^L \chi(\epsilon_1; L) \right),$$

- Fourier expansion of  $\epsilon_+ = e^{inx^+} \rightarrow \chi_n, \quad \epsilon_- = e^{inx^-} \rightarrow \bar{\chi}_n$

$$[\chi_m, \chi_n]_* = (m - n) \chi_{m+n},$$

$$[\bar{\chi}_m, \bar{\chi}_n]_* = (m - n) \bar{\chi}_{m+n}$$

- 2 copies of Witt algebra

# Symmetries and conserved charges

**Theorem** [Wald (1993), Barnich, Brandt(2001), Compère (2007)]

Symplectic form associate a *charge* to a symmetry

$$\omega[\delta\Phi, \delta_\chi\Phi, \Phi] = d\mathbf{k}_\chi[\delta\Phi, \Phi]$$

- Charge variation

$$\delta Q_\chi \equiv \oint_{\partial\Sigma} \mathbf{k}_\chi[\delta\Phi, \Phi].$$

- Definition of charge (assuming *integrability* )

$$Q_\chi \equiv \int_{\phi_0}^{\phi} \delta Q_\chi$$



# Charges and their algebra

[G. Compre, L. Donnay, P. H. Lambert and W. Schulgin, (2015)]

[G.Compère, P.Mao, A.S and M.M.Sheikh-Jabbari, (2016).]

## Conserved charges everywhere

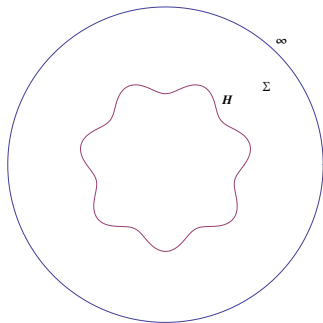
- Charges corresponding to *symplectic symmetries* can be computed **at any closed surface**

$$\omega[\delta\Phi, \delta_\chi\Phi, \Phi] = d\mathbf{k}_\chi[\delta\Phi, \Phi]$$

$$\Downarrow$$

$$\delta Q_\chi \Big|_\infty - \delta Q_\chi \Big|_H = \int_\Sigma \omega(\delta\Phi, \delta_\chi\Phi, \Phi)$$

$$= 0$$



# Algebra of charges

## Theorem

- Algebra of charges is a central extension of the algebra of symplectic symmetries

$$Q_X[g] = \frac{\ell}{8\pi G} \int_0^{2\pi} d\phi (\epsilon_+(x^+)L_+(x^+) + \epsilon_-(x^-)L_-(x^-))$$

- Fourier expansion

$$\begin{aligned} \{Q_m, Q_n\} &= (m-n)Q_{m+n} + \frac{c}{12}m^3\delta_{m+n,0}, \\ \{\bar{Q}_m, \bar{Q}_n\} &= (m-n)\bar{Q}_{m+n} + \frac{c}{12}m^3\delta_{m+n,0}, \end{aligned}$$

where  $c = \frac{3\ell}{2G}$ . This is the symmetry algebra of a 2 dimensional CFT

## Symmetries at the horizon

Symplectic symmetries *induce* a symmetry algebra near the *horizon* of black hole.

- Extremal black holes in 3 dimensions

[G.Compère, P.Mao, A.S and M.M.Sheikh-Jabbari, (2016).]

- Non-extremal black holes

[G. Compère and J. Long, (2016)]

[ H. Afshar, S. Detournay, D. Grumiller and B. Oblak, (2015)]

[L. Donnay, G. Giribet, H. A. Gonzalez and M. Pino, (2015)]

## Results

- Asymptotically AdS<sub>3</sub> geometries form a phase space.
- The set of solutions form a null surface in the phase space
- Symplectic symmetries are the extension of asymptotic symmetries inside the bulk
- Charges can be computed over *any* closed surface which is a smooth deformation of the circle at infinity

**Thank you for your attention**