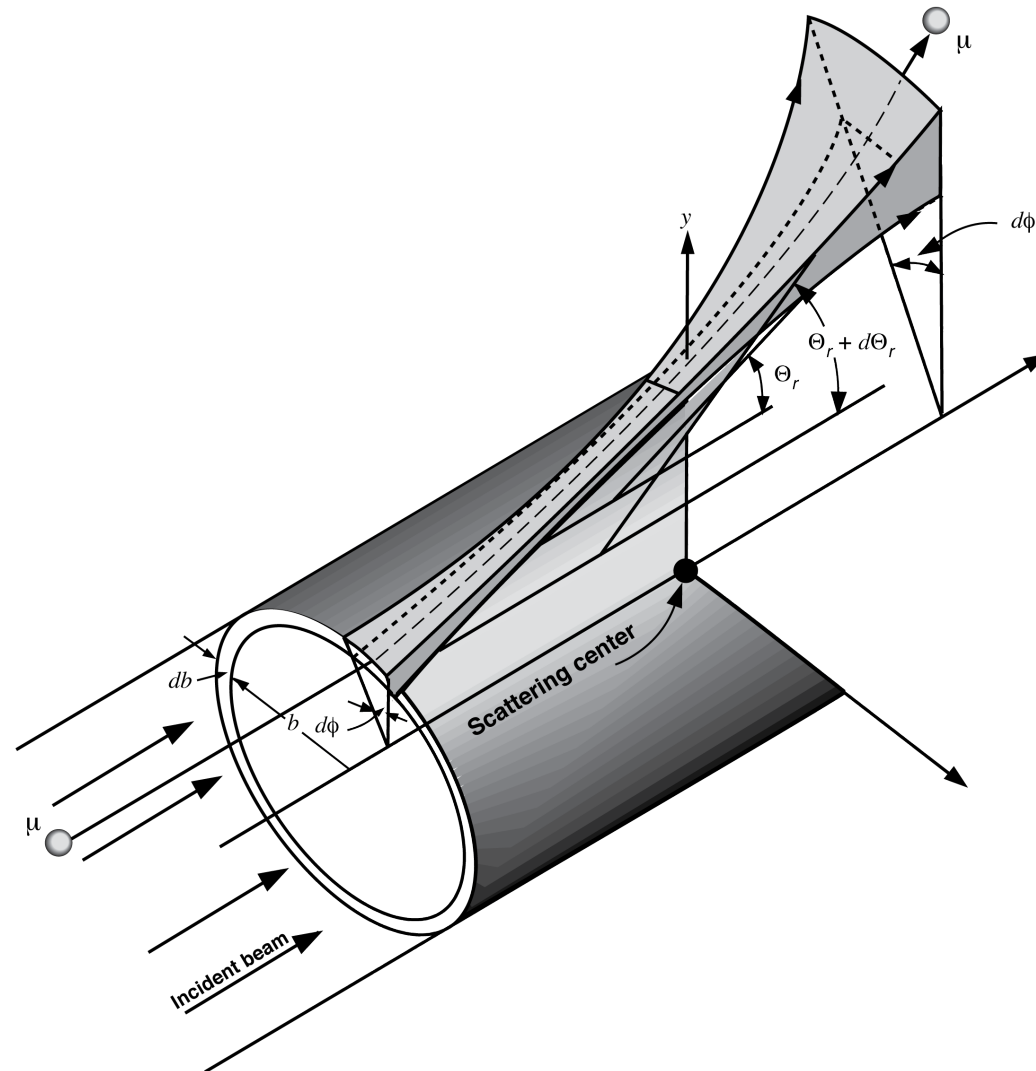
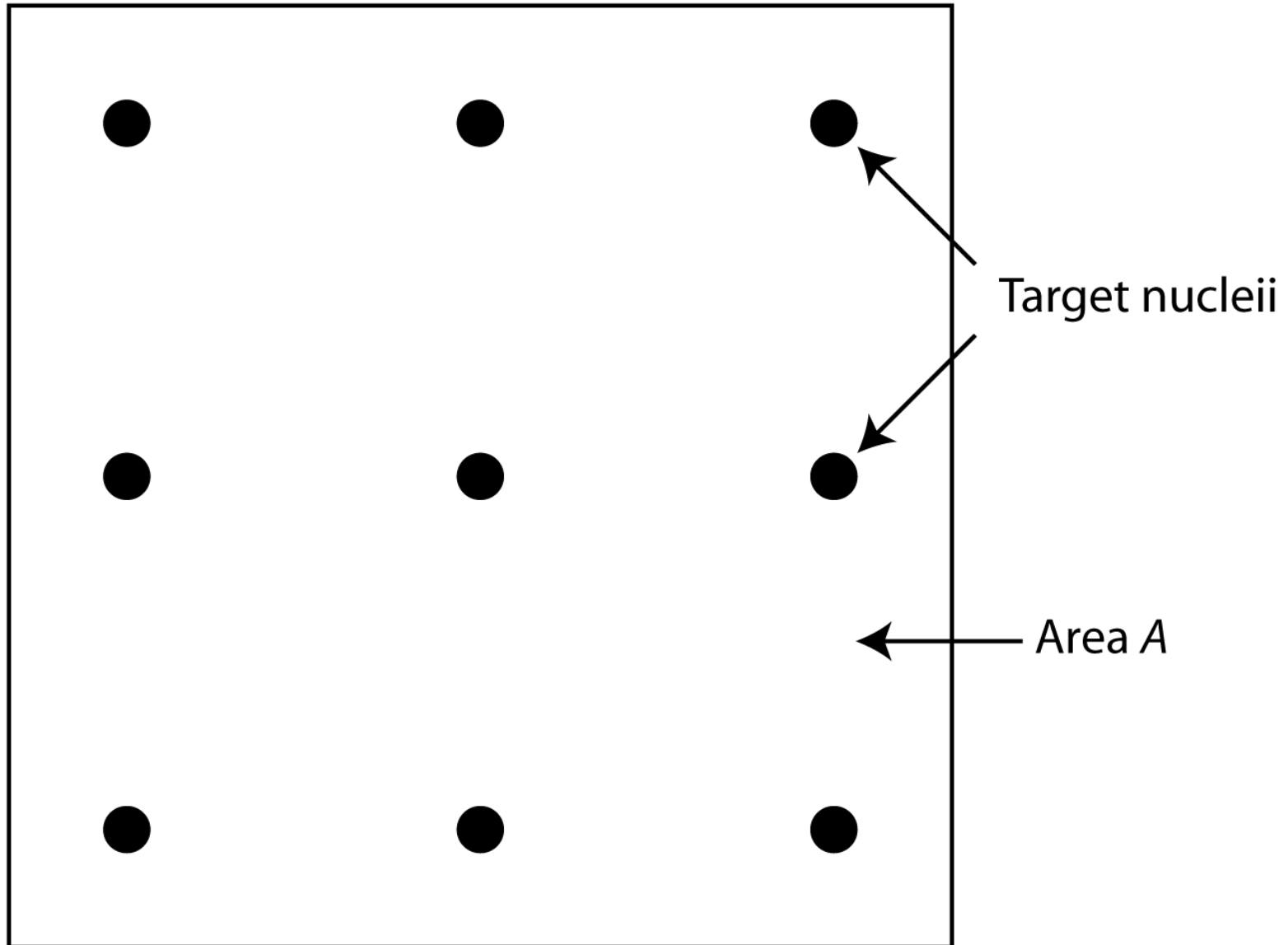


Part II

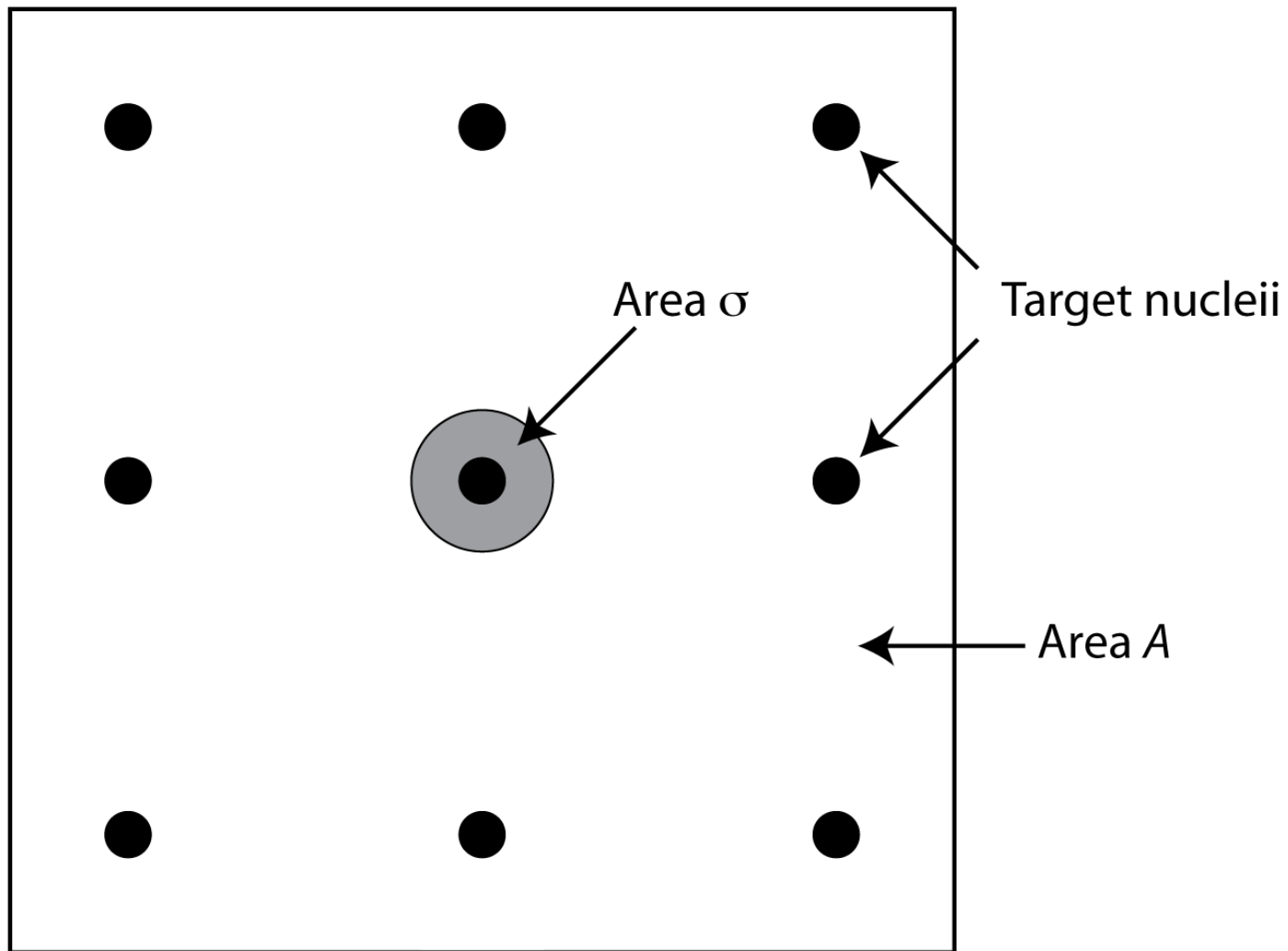
Introduction to Reaction and Scattering Cross-Sections



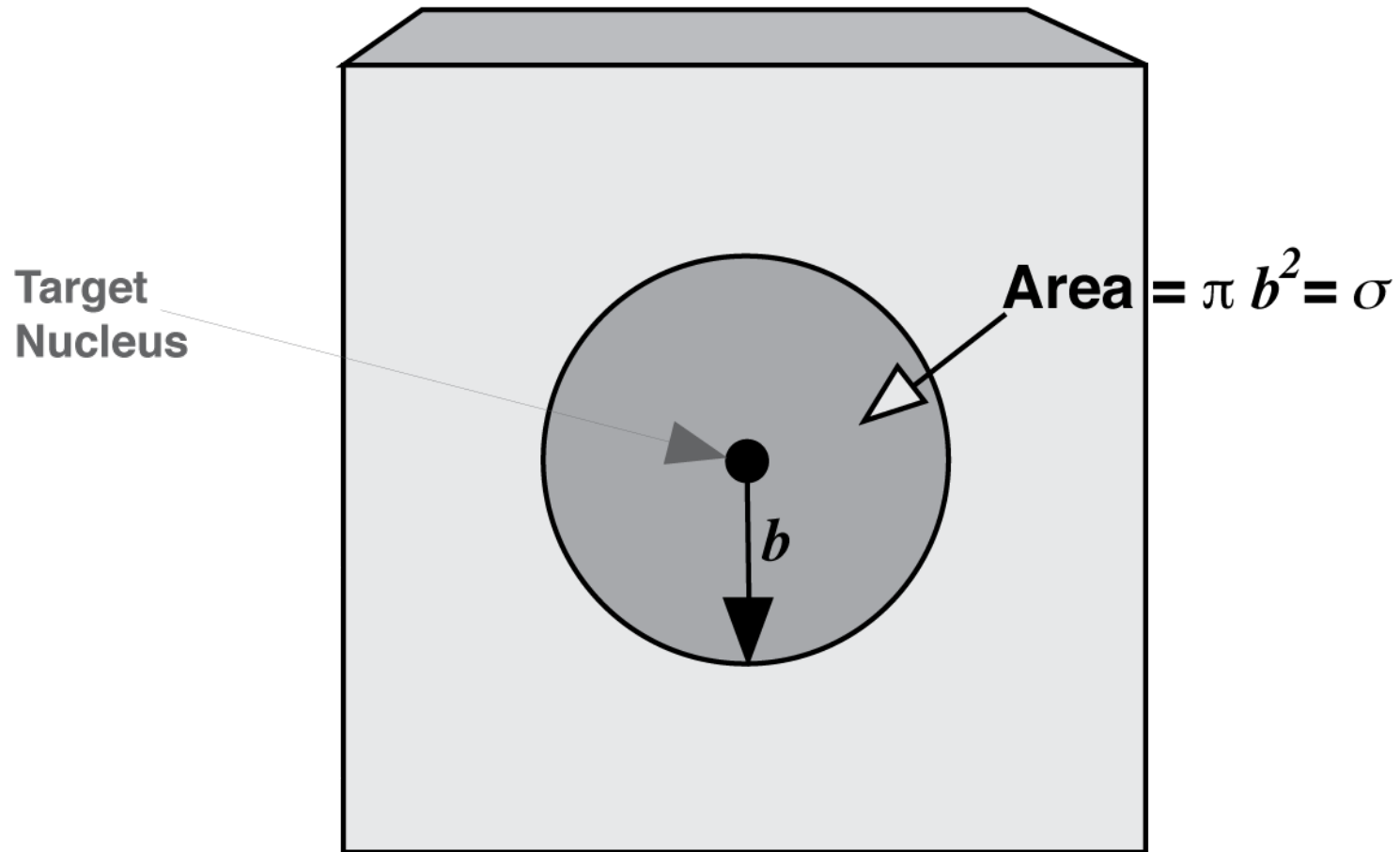
1. Reaction Cross-Section Introduction



The cross-section, σ , is the area that one atom presents to a beam of incident particles for initiating a specific “interaction”



The impact parameter, b

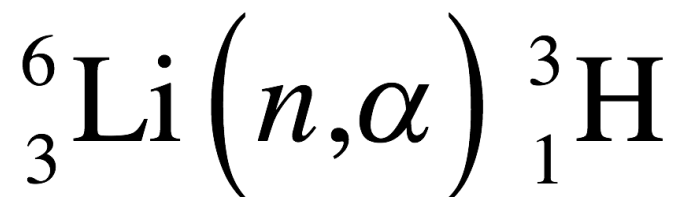
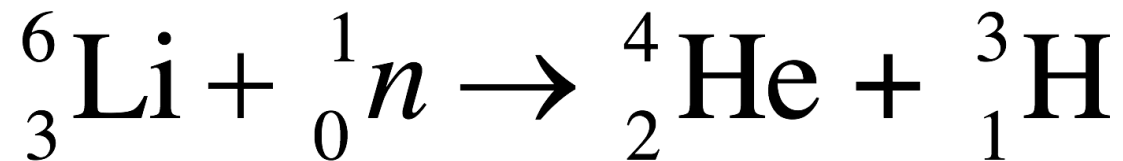


$$\sigma \left[\text{cm}^2 \right] = \pi (b \left[\text{cm} \right])^2$$

$$\sigma = \pi b^2$$

Transmutation reactions

	Be6 5.0E-21 s 2p, α 6.01972	Be7 ^{3/-} 53.28 d € γ 477.6 σ _p 3.9E4, 2E4 σ _α .14, .06 E .862	Be8 ~7E-17 s 2α .0461 8.0053051	Be9 ^{3/-} 100 σ _γ 8 mb, 4 mb 9.0121822	Be10 1.6E6 a β ⁻ .556 noγ σ _γ <1 mb E .5561
	Li5 ^{3/-} ~3E-22 s p, α 5.01254	Li6 ¹⁺ 7.5* σ _α 941, 423 σ _γ 39 mb, 17 mb 6.015122	Li7 ^{3/-} 92.5* σ _γ .045, .020 7.016003	Li8 ²⁺ 0.84 s β ⁻ 13 (2α) 1.57 E 16.004	Li9 ^{3/-} 177 ms β ⁻ 13.5, 11.0, ... (n) .3, ... (2α) .7, ... E 13.606
He3 ^{1/+} 0.000138 σ _p 5.33E3, 2.40E3 σ _γ .05 mb 3.01602930	He4 99.999862 σ _γ 0 4.00260323	He5 ^{3/-} 7.6E-22 s n, α 5.01222	He6 807 ms β ⁻ 3.510 noγ E 3.507	He7 ^{(3)/-} 3E-21 s n 7.02803	He8 119 ms β ⁻ 10, ... γ 980.7 (n) .61-3.0 t ω E 10.65
H2 ¹⁺ 0.015 σ _γ .52 mb, 23 mb 2.01410178	H3 ^{1/+} 12.3 a β ⁻ .0186 noγ σ _γ <6 μb E .01860				6



Li transmutation via thermal n capture

$$\sigma_{\alpha} = 941 \text{ [barns]} = 941 \cdot 10^{-24} \text{ [cm}^2\text{]}$$

$$b_{\alpha} = \sqrt{\frac{\sigma_{\alpha}}{\pi}} = 1.73 \cdot 10^{-11} \text{ [cm]} = 1.73 \cdot 10^{-13} \text{ [m]} = 173 \text{ [fm]}$$

How large is 173 fm?

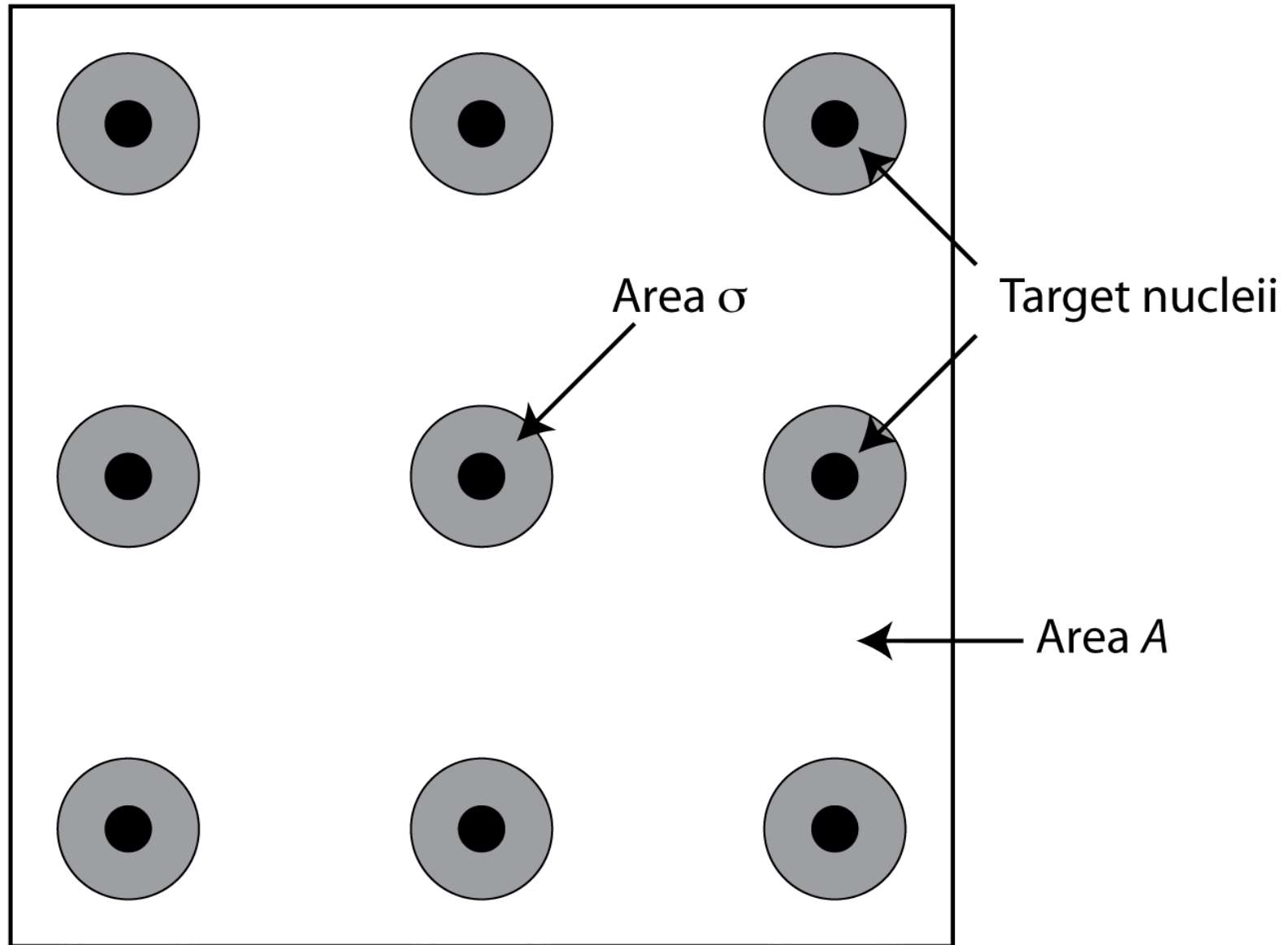
$$R = 1.25 \cdot 10^{-5} M_T^{1/3} [\text{\AA}]$$

$$M_T(^6\text{Li}) = 6.01522 [\text{amu}]$$

$$R = 2.27332 \cdot 10^{-5} [\text{\AA}]$$

$$= 2.27332 \cdot 10^{-5} [\text{\AA}] = 2.27 [\text{fm}]$$

Reaction Cross-Section(s). Relationship to the Probability of Interaction



Fluence

$$\Phi \left[\frac{\text{interactions}}{\text{cm}^2} \right] = \Phi_0 \left[\frac{\text{projectile particles}}{\text{cm}^2} \right] \hat{\rho} \left[\frac{\text{target atoms}}{\text{cm}^2} \right] \sigma \left[\frac{\text{cm}^2}{\text{target atom}} \right]$$

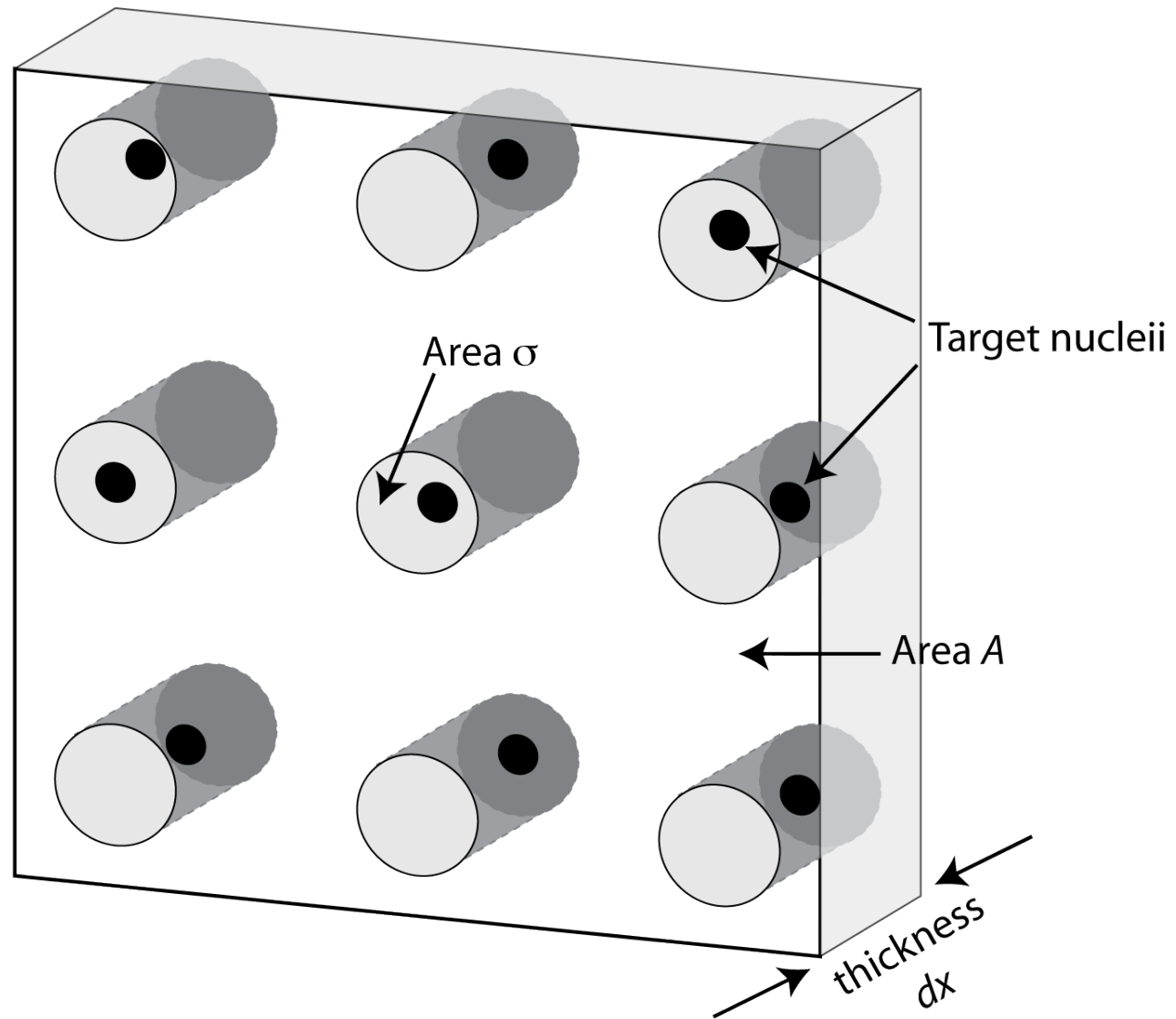
$$\Phi = \Phi_0 \hat{\rho} \sigma$$

Probability vs. Cross-Section

$$\begin{aligned} \textit{Probabilty of Interaction} &= \frac{\Phi \left[\frac{\text{interactions}}{\text{cm}^2} \right]}{\Phi_0 \left[\frac{\text{projectile particles}}{\text{cm}^2} \right]} \\ &= \hat{\rho} \left[\frac{\text{target atoms}}{\text{cm}^2} \right] \sigma \left[\frac{\text{cm}^2}{\text{target atom}} \right] \end{aligned}$$

$$P = \frac{\Phi}{\Phi_0} = \hat{\rho} \sigma$$

3-D



$$\Phi \left[\frac{\text{interactions}}{\text{cm}^2} \right] = \Phi_0 \left[\frac{\text{projectile particles}}{\text{cm}^2} \right] \rho \left[\frac{\text{target atoms}}{\text{cm}^3} \right] \cdot dx \text{ [cm]} \sigma \left[\frac{\text{cm}^2}{\text{target atom}} \right]$$

$$\Phi = \Phi_0 (\rho \cdot dx) \sigma$$

3-D

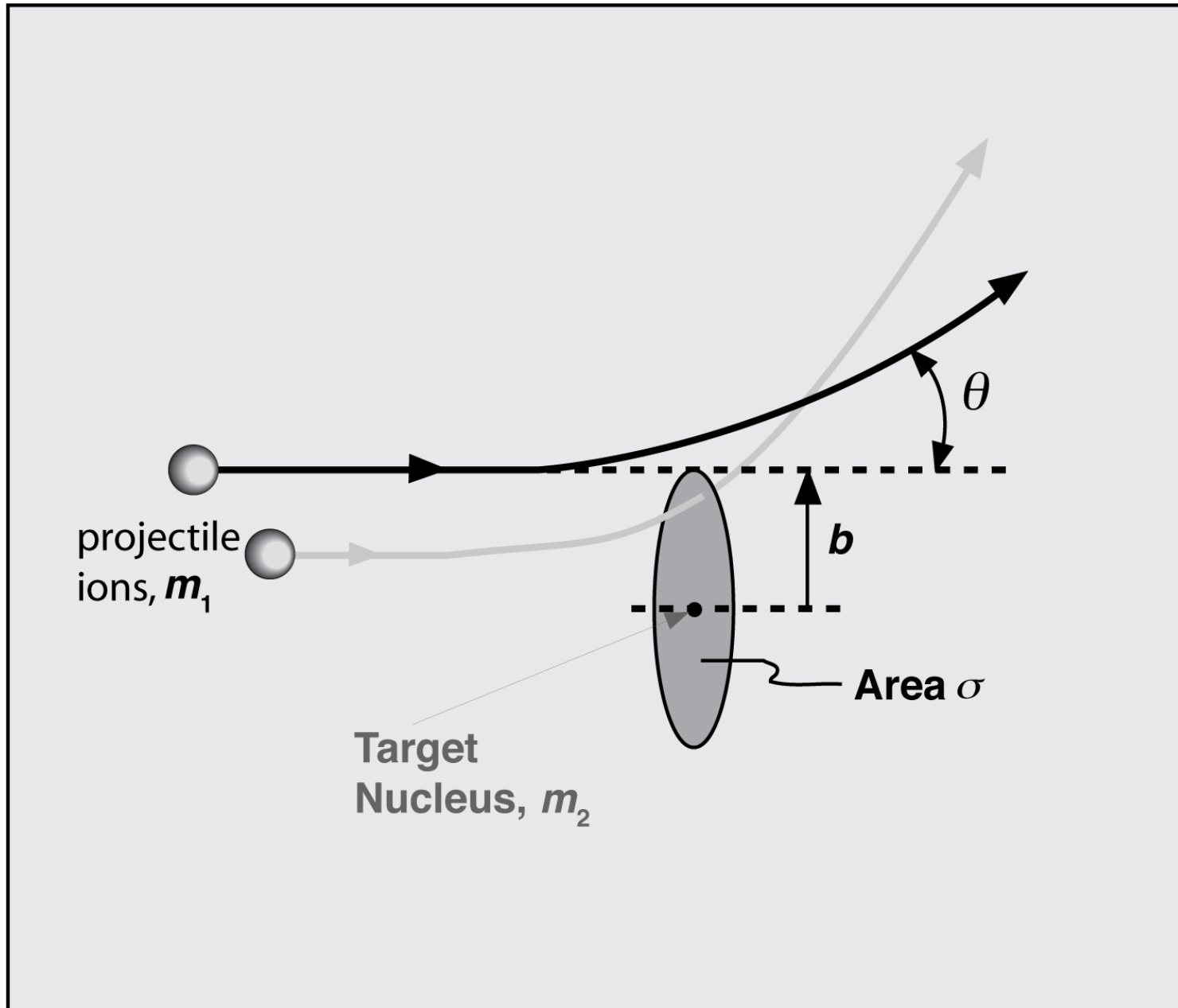
Probability of Interaction

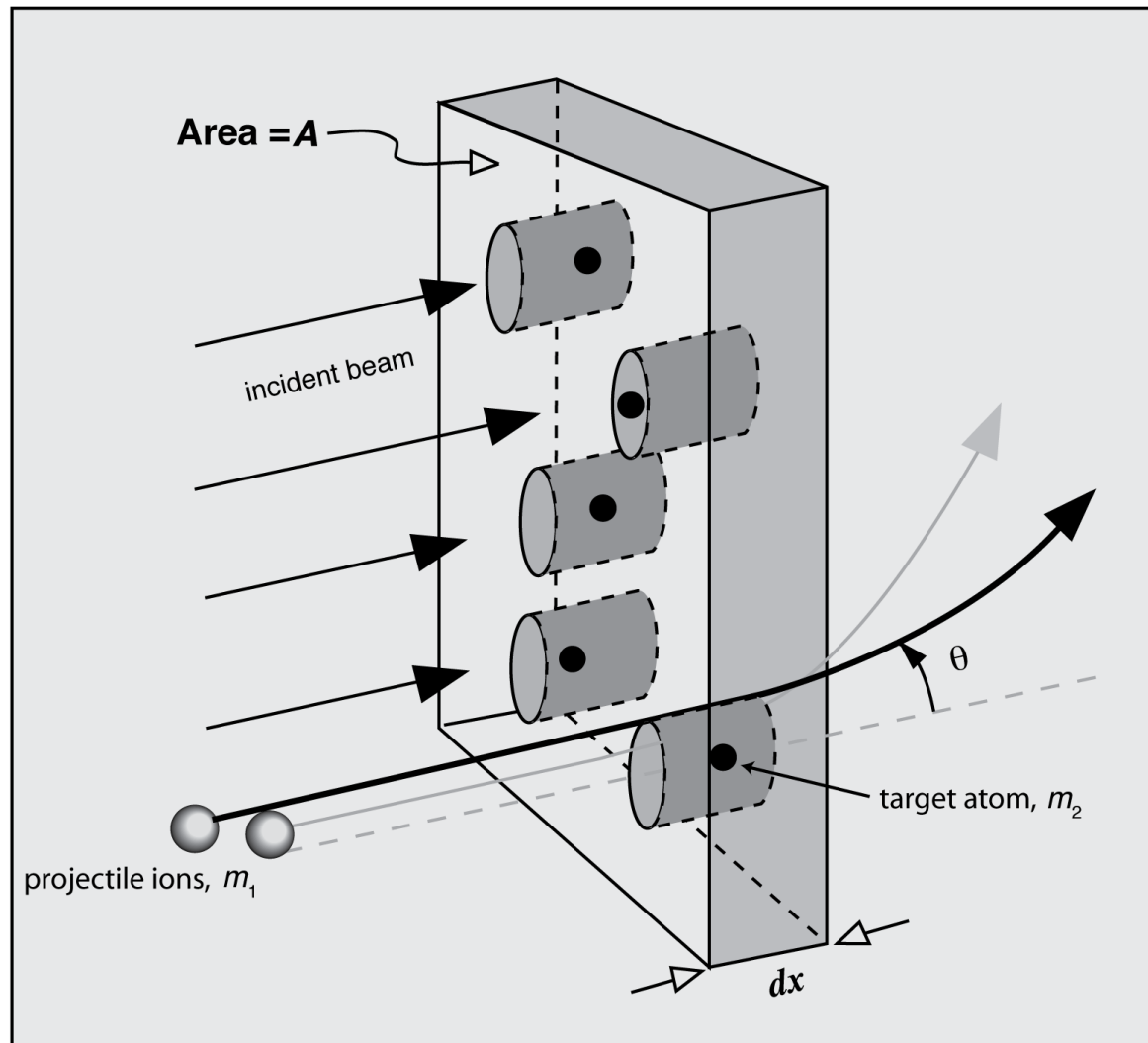
$$= \frac{\Phi \left[\frac{\text{interactions}}{\text{cm}^2} \right]}{\Phi_0 \left[\frac{\text{projectile particles}}{\text{cm}^2} \right]}$$

$$= \rho \left[\frac{\text{target atoms}}{\text{cm}^3} \right] \cdot dx \text{ [cm]} \quad \sigma \left[\frac{\text{cm}^2}{\text{target atom}} \right]$$

$$P = \frac{\Phi}{\Phi_0} = (\rho \cdot dx) \sigma$$

The Total Scattering Cross-Section, σ





$$\Phi \left[\frac{\text{projectile ions scattered to } \theta \leq \theta \leq \pi}{\text{cm}^2} \right] = \Phi_0 \left[\frac{\text{projectile ions}}{\text{cm}^2} \right] \rho \left[\frac{\text{target nuclei}}{\text{cm}^3} \right] \cdot dx [\text{cm}] \sigma \left[\frac{\text{cm}^2}{\text{target nucleus}} \right]$$

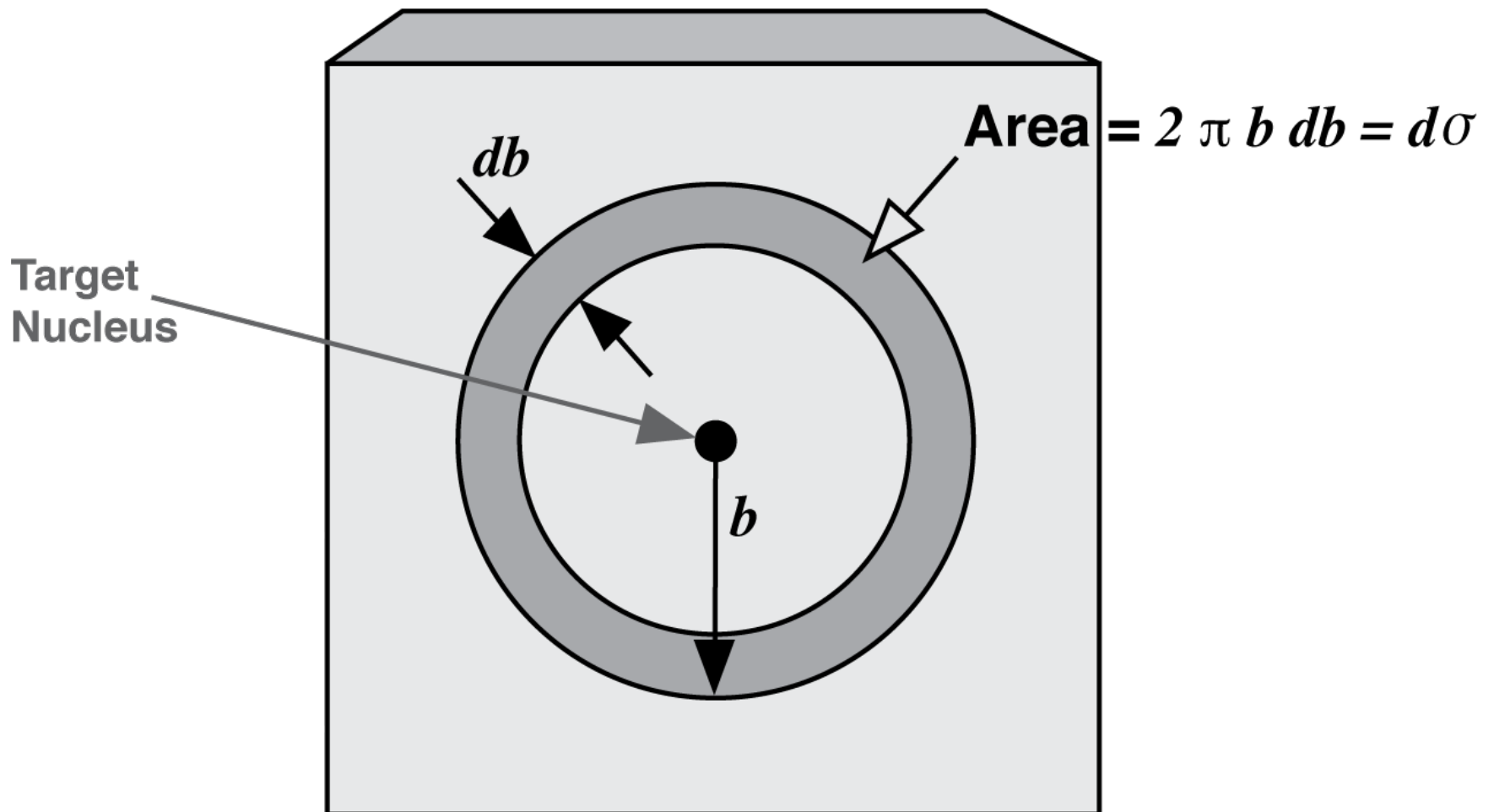
$$\Phi = \Phi_0 (\rho \cdot dx) \sigma$$

Probabilty of Scattering to $\theta \leq \theta \leq \pi$

$$= \frac{\Phi \left[\frac{\text{projectile ions scattered to } \theta \leq \theta \leq \pi}{\text{cm}^2} \right]}{\Phi_0 \left[\frac{\text{projectile ions}}{\text{cm}^2} \right]} = \rho \left[\frac{\text{target nuclei}}{\text{cm}^3} \right] \cdot dx \text{ [cm]} \sigma \left[\frac{\text{cm}^2}{\text{target nucleus}} \right]$$

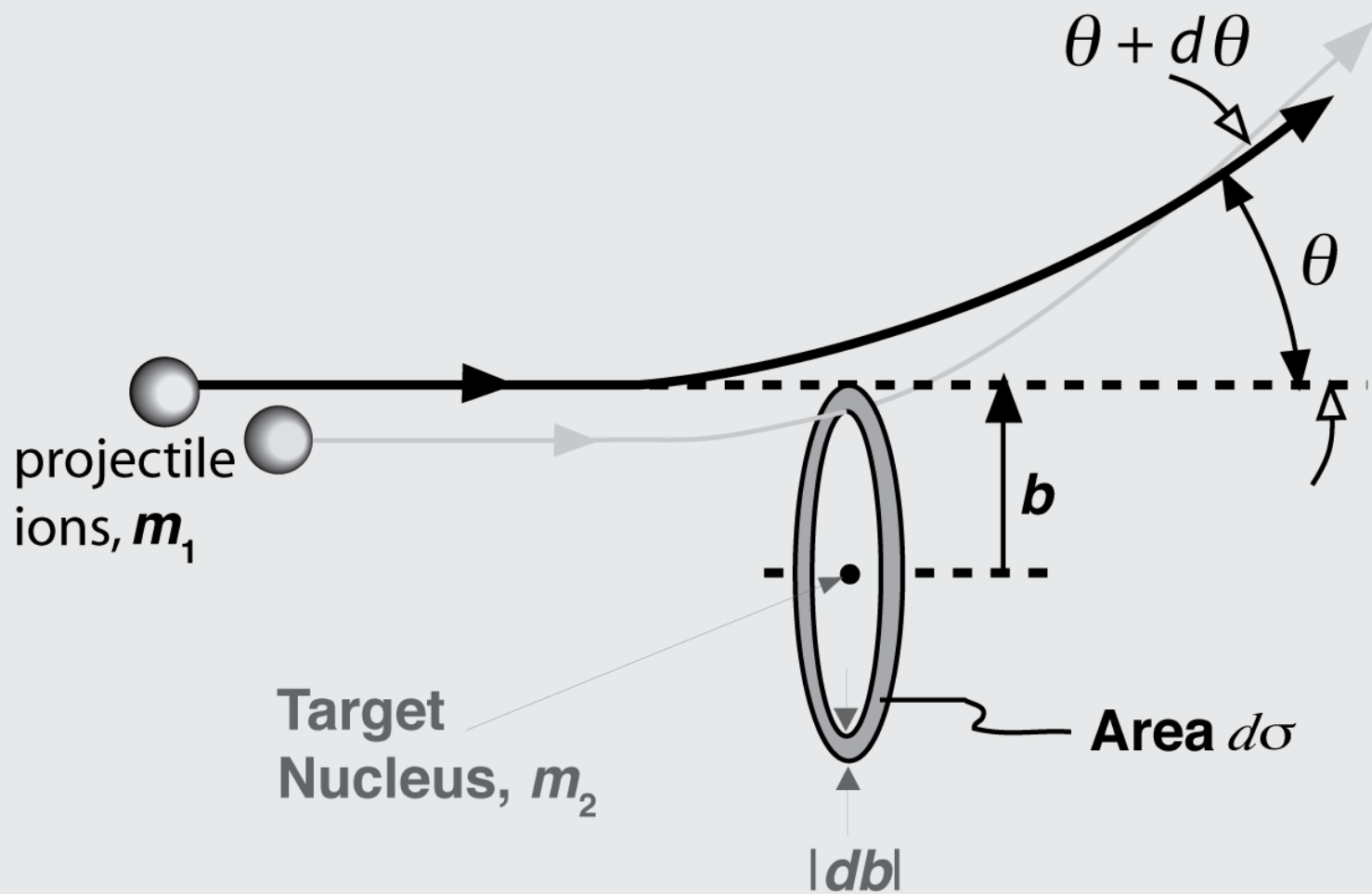
$$P = \frac{\Phi}{\Phi_0} = (\rho \cdot dx) \sigma$$

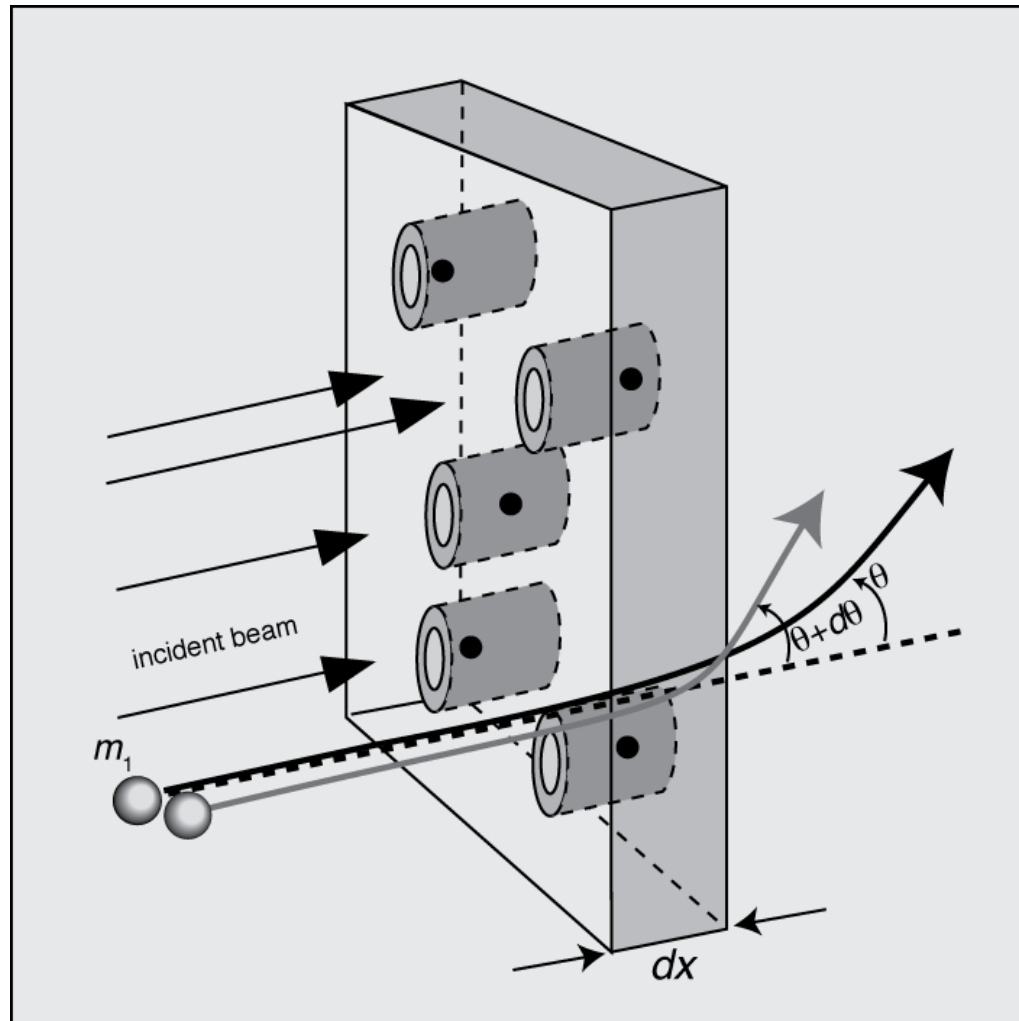
The Differential Scattering Cross-Section, $d\sigma/d\Omega$



$$d\sigma \left[\text{cm}^2 \right] = 2 \pi b \left[\text{cm} \right] db \left[\text{cm} \right]$$

$$d\sigma = 2 \pi b db$$





$$d\Phi \left[\frac{\text{projectile ions scattered to } \theta \leq \theta \leq \theta + d\theta}{\text{cm}^2} \right] = \Phi_0 \left[\frac{\text{projectile ions}}{\text{cm}^2} \right] \rho \left[\frac{\text{target nuclei}}{\text{cm}^3} \right] \cdot dx \text{ [cm]} \quad d\sigma \left[\frac{\text{cm}^2}{\text{target nucleus}} \right]$$

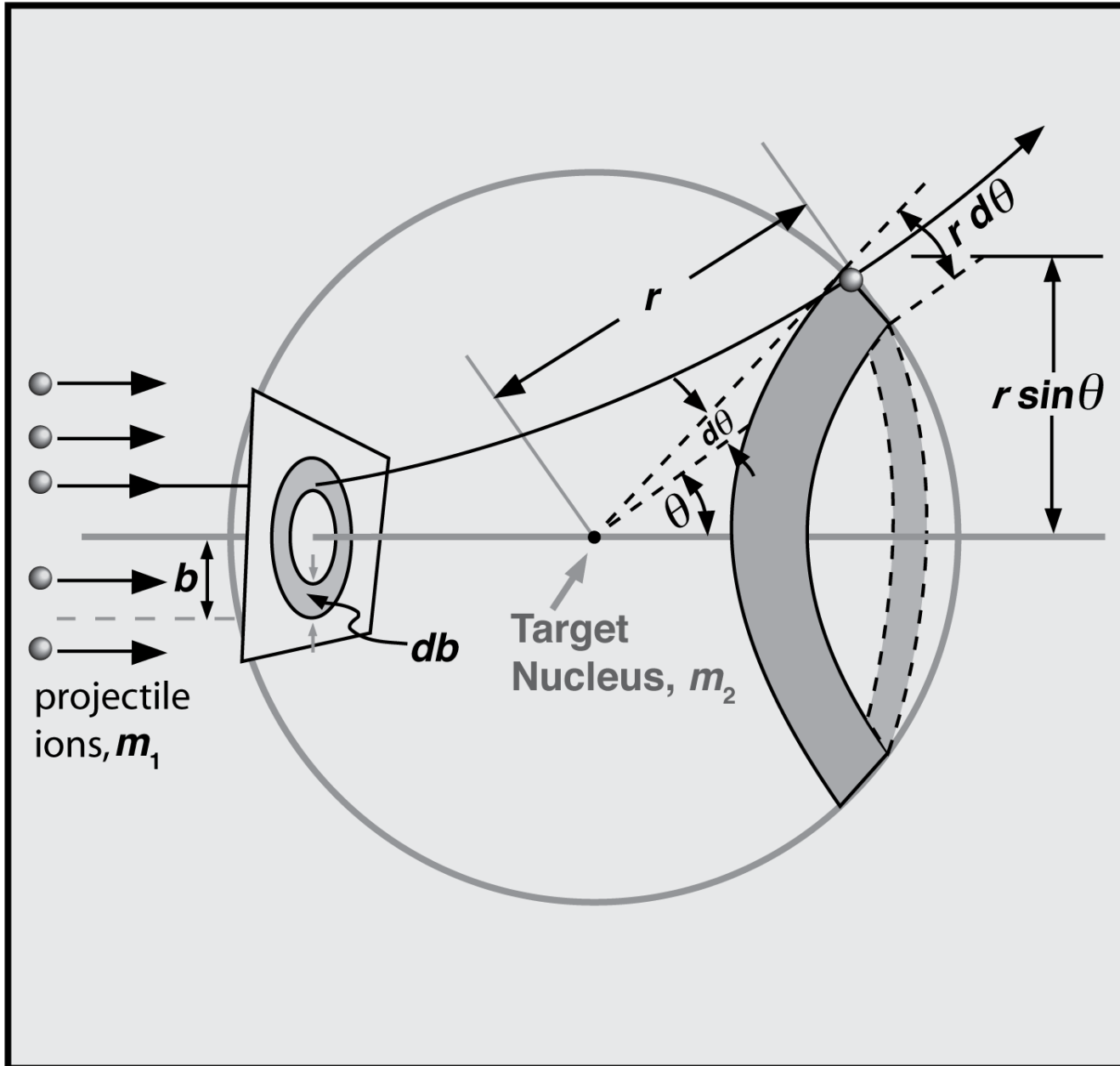
$$d\Phi = \Phi_0 (\rho \cdot dx) d\sigma$$

Differential Probability of Scattering to $\theta \leq \theta \leq \theta + d\theta$

$$= \frac{d\Phi \left[\frac{\text{projectile ions scattered to } \theta \leq \theta \leq \theta + d\theta}{\text{cm}^2} \right]}{\Phi_0 \left[\frac{\text{projectile ions}}{\text{cm}^2} \right]} = \rho \left[\frac{\text{target nuclei}}{\text{cm}^3} \right] \cdot dx \text{ [cm]} \quad d\sigma \left[\frac{\text{cm}^2}{\text{target nucleus}} \right]$$

$$P(\theta, \theta + d\theta) = \frac{d\Phi}{\Phi_0} = (\rho \cdot dx) d\sigma$$

Solid angle, $d\Omega$, associated with scattering in 3-D, to angles between θ and $\theta + d\theta$



$$\frac{d\Omega}{4\pi} = \frac{2\pi (r \sin \theta) (r d\theta)}{4\pi r^2}$$

$$d\Omega = 2\pi \sin \theta d\theta$$

Fractional ion fluence scattered through angle θ and within solid angle Ω

$$\begin{aligned}
 & \frac{d\Phi \left[\begin{array}{c} \text{projectile ions} \\ \text{scattered to} \\ \theta \leq \theta \leq \theta + d\theta \\ \hline \text{cm}^2 \end{array} \right]}{\Phi_0 \left[\begin{array}{c} \text{projectile ions} \\ \hline \text{cm}^2 \end{array} \right]} = \rho \left[\begin{array}{c} \text{target nuclei} \\ \hline \text{cm}^3 \end{array} \right] \cdot dx_{[\text{cm}]} \int_{\Omega} \frac{d\sigma}{d\Omega} \left[\begin{array}{c} \text{cm}^2 \\ \hline \text{target nucleus} \cdot \text{steradian} \end{array} \right] d\Omega_{[\text{steradians}]} \\
 P(\Omega(\theta)) &= \frac{d\Phi}{\Phi_0} = (\rho \cdot dx) \int_{\Omega} \frac{d\sigma}{d\Omega} d\Omega
 \end{aligned}$$

Rutherford backscattering (RBS) detector:

Fractional ion fluence scattered through angle θ and within solid angle Ω

$$P(\Omega(\theta)) = \frac{d\Phi}{\Phi_0} = (\rho \cdot dx) \frac{d\sigma}{d\Omega} \Omega$$

The Differential Energy Transfer Cross-Section, $d\sigma(E,T)/dT$

$$d\sigma(b) / db \rightarrow d\sigma(E, \theta) / d\Omega \rightarrow d\sigma(E, T) / dT$$

State (without proof) an important relationship:

$$\frac{d\sigma(E, T)}{dT} = \frac{4 \pi}{\Lambda E} \frac{d\sigma(E, \theta(T))}{d\Omega}$$

Cross-Sections for a Hard-Sphere Interaction Potential

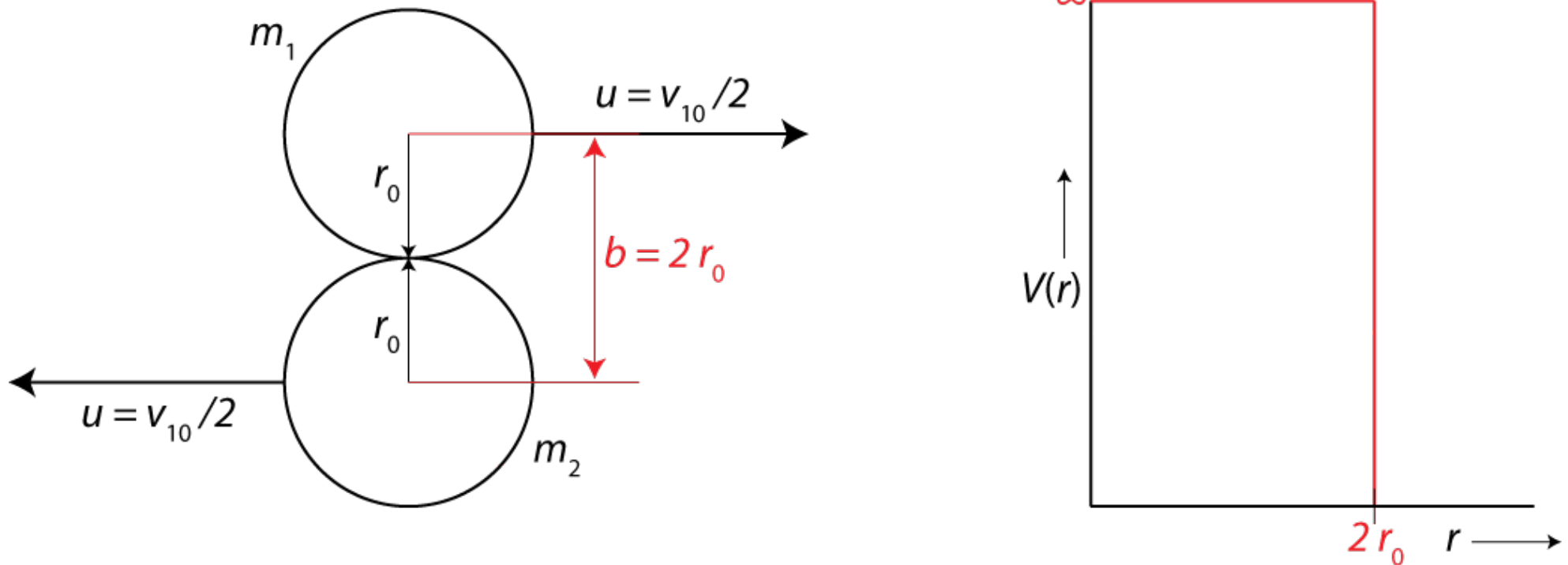


Figure 13.

$$\sigma = \pi b^2 \quad \rightarrow \quad d\sigma = 2\pi b db \quad \rightarrow \quad \boxed{\frac{d\sigma(b)}{db} = 2\pi b}$$

Hard sphere – scattering is isotropic in
the *CM* system

$$\frac{d\sigma\left(E, \theta(T)\right)}{\sigma(E)} = \frac{d\Omega(\theta)}{4\pi} = \frac{2\pi \sin\theta d\theta}{4\pi}$$

Hard-sphere

Differential energy transfer cross-section

$$\begin{aligned}\frac{d\sigma(E, T)}{dT} &= \frac{4 \pi}{\Lambda E} \frac{d\sigma(E, \theta(T))}{d\Omega} \\ &= \frac{4 \pi}{\Lambda E} \frac{\sigma(E)}{4 \pi}\end{aligned}$$

$$\boxed{\frac{d\sigma(E, T)}{dT} = \frac{\sigma(E)}{\Lambda E}}$$

Cross-Sections for a *Rutherford* Interaction Potential

$$V(r) = \frac{K}{r}$$

$$\begin{aligned} K &= (Z_1 e) (Z_2 e) \\ &= Z_1 Z_2 e^2 \end{aligned}$$

Cross-Sections for a *Rutherford* Interaction Potential

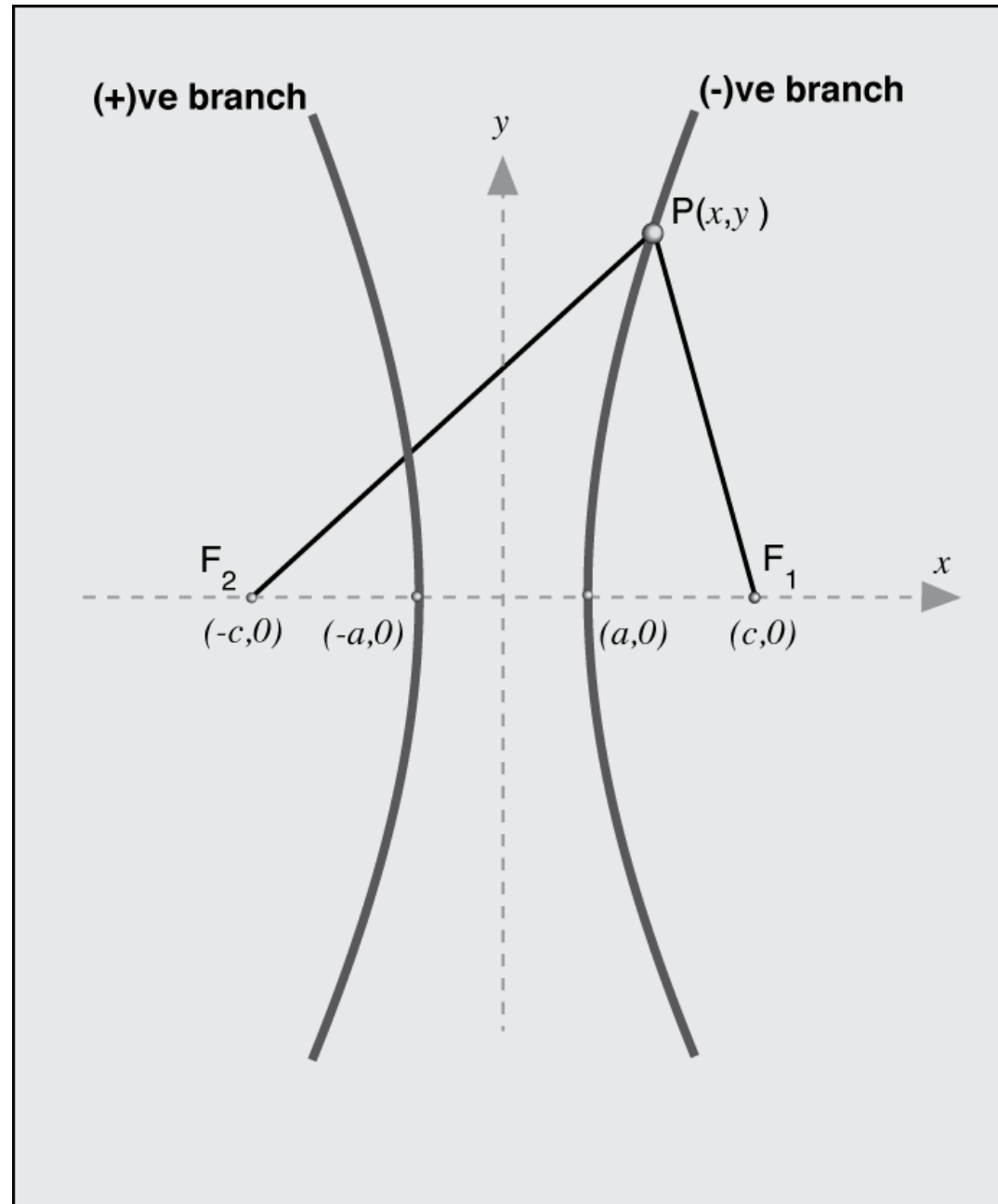
$$d\sigma(b) / db \rightarrow d\sigma(E, \theta) / d\Omega \rightarrow d\sigma(E, T) / dT$$

$$\frac{d\sigma(b)}{db} = 2 \pi b$$

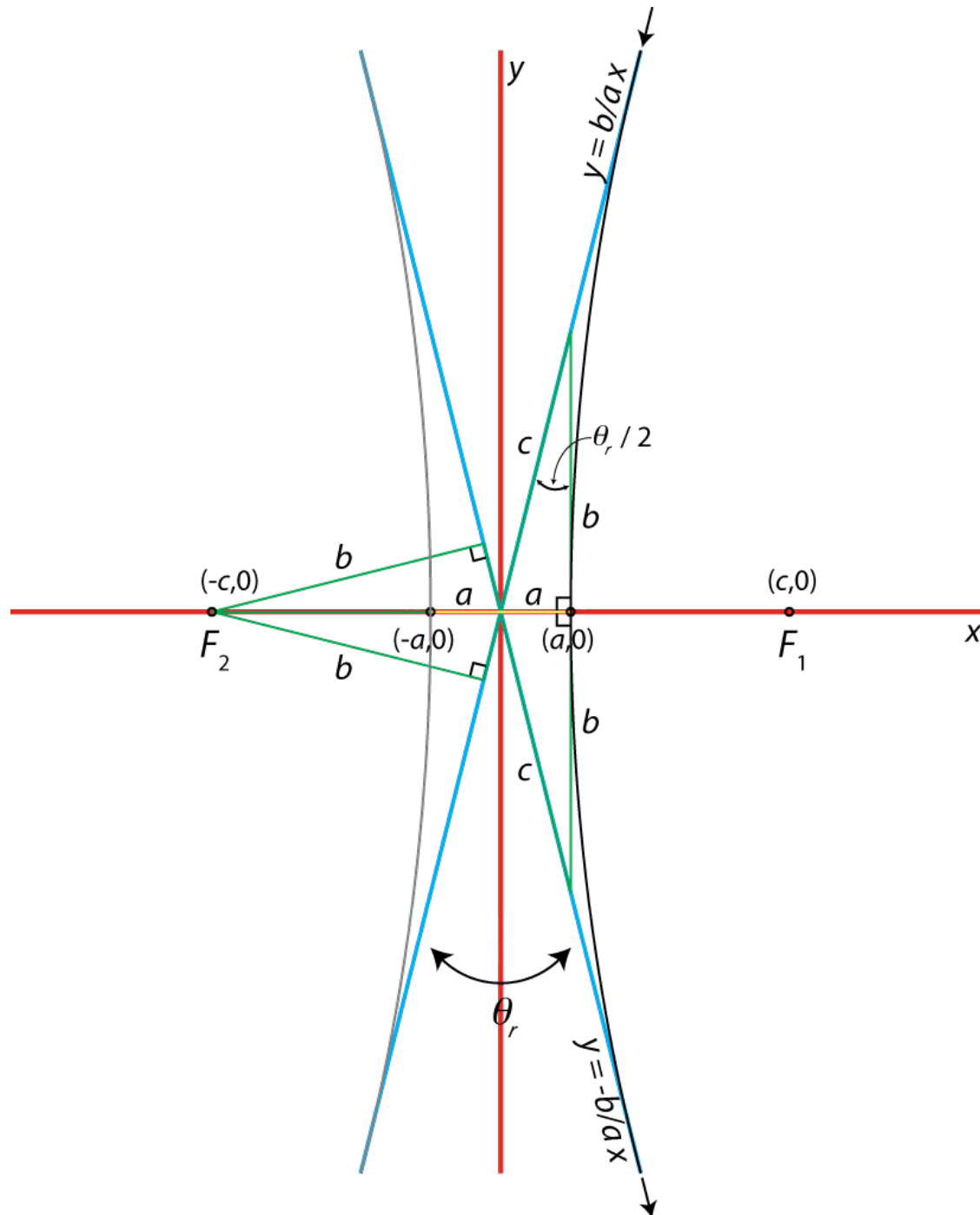
Next, we require an equation that provides a one-to-one mapping of

$$b \rightarrow \theta$$

Hyperbolic solution to the equation of motion



Hyperbolic solution to the equation of motion



$$\frac{b}{a} = \cot\left(\frac{\theta_r}{2}\right)$$

$$a = \frac{K \left(1 + \frac{m_1}{m_2} \right)}{2 E}$$

Rutherford cross-section for angular scattering

$$\frac{d\sigma(E, \theta)}{d\Omega} = \left(\frac{K \left(1 + \frac{m_1}{m_2} \right)}{4 E} \right)^2 \frac{1}{\sin^4 \left(\frac{\theta}{2} \right)}$$

Rutherford cross-section for energy transfer

$$d\sigma(E, \theta) / d\Omega \rightarrow d\sigma(E, T) / dT$$

$$\boxed{\frac{d\sigma(E, T)}{dT} = \frac{\pi K^2 x}{E T^2}}$$

where $x = m_1 / m_2$ and $K = Z_1 Z_2 e^2$

Rutherford scattering: *Differential probability* of energy transfer between T and $T+dT$

$$P(T, T + dT) = (\rho \, dx) \frac{d\sigma(E, T)}{dT} dT$$
$$= (\rho \, dx) \frac{\pi \, K^2 \, x}{E \, T^2} dT$$

$$\begin{aligned}
\sigma(E) &= \int_{T=0}^{T=\Lambda E} d\sigma(E, T) \\
&= \int_{T=0}^{T=\Lambda E} \frac{\pi K^2 x}{E} \frac{dT}{T^2} \\
&= \frac{\pi K^2 x}{E} \int_{T=0}^{T=\Lambda E} \frac{dT}{T^2} \\
&= \frac{\pi K^2 x}{E} \int_{T=0}^{T=\Lambda E} \frac{dT}{T^2} \\
&= - \frac{\pi K^2 x}{E} \frac{1}{T} \bigg|_{T=0}^{T=\Lambda E} \\
&= \frac{\pi K^2 x}{E} \left(\frac{1}{0} - \frac{1}{\Lambda E} \right)
\end{aligned}$$

The solution:
The *Screened Coulomb Potential*

$$V(r) = \frac{K}{r} e^{-\frac{r}{\hat{a}}}$$