

Chapter VI

Rectifying and amplifying stresses

I) Linear medium: dipole conservation

1. Mean-stress theorem
2. Dipole conservation in linear media
3. Relating the macroscopic stress to macroscopic work
4. Relating the macroscopic force dipole to macroscopic work

II) Nonlinear (bucklable) medium: amplification

1. Radial force balance
2. Double-check: the linear case
3. Force amplification

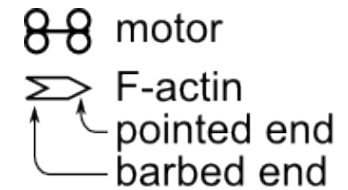
III) Discussion

Reference:

Ronceray & ML, *Soft Matter* **11**, 1597 (2015)

Ronceray, Broedersz & ML, *Proc. Natl. Acad. Sci. U.S.A.* **113**, 2827 (2016)

Actomyosin is an active, contractile material

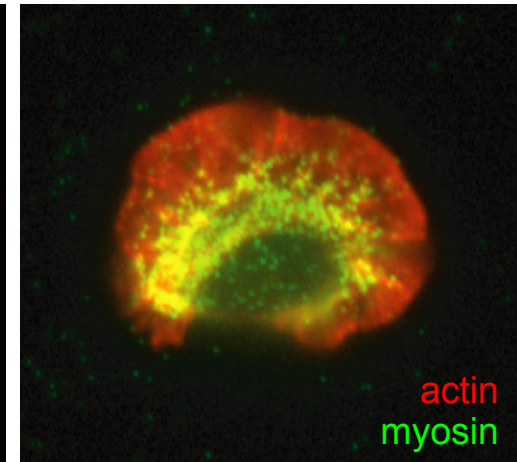
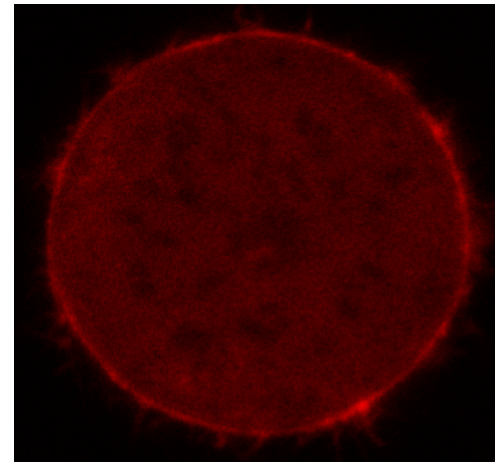
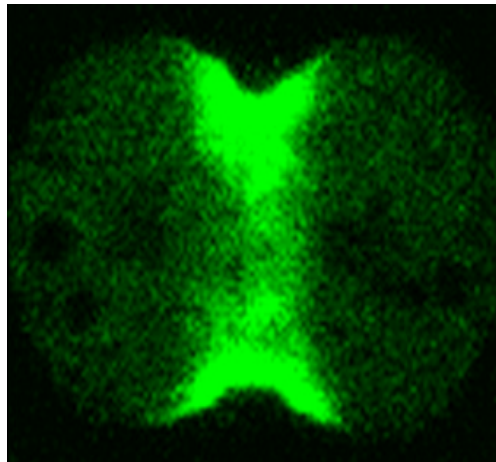
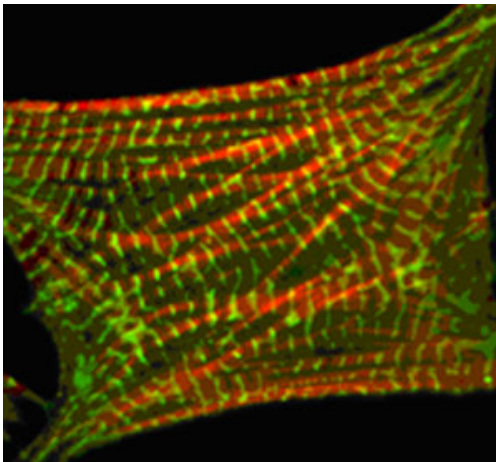


striated muscle

cytokinetic ring (1d)

cell cortex (2d)

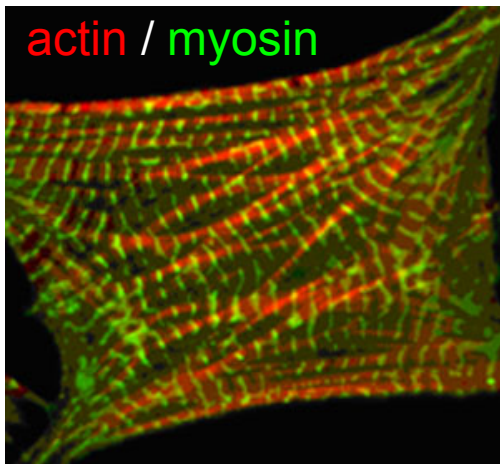
lamellar network (3d)



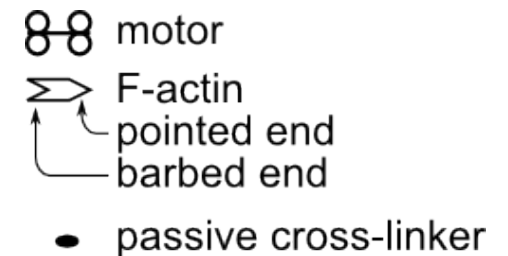
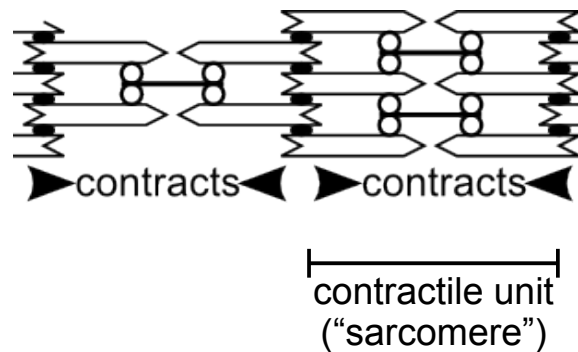
ordered

disordered

F-actin ordering accounts for striated muscle contractility



Contraction mechanism:



Universal contractility paradigm?

Myosin generates **contractile** forces, which the actin scaffold **transmits** over long distances.

Myosin motors have no intrinsic propensity for contraction



Why is disordered actomyosin contractile rather than extensile?

Linear networks cannot favor contraction over extension

$$\sigma_{\text{far field}} = -\rho \mathcal{D}_{\text{local}} + \Sigma$$

active stress

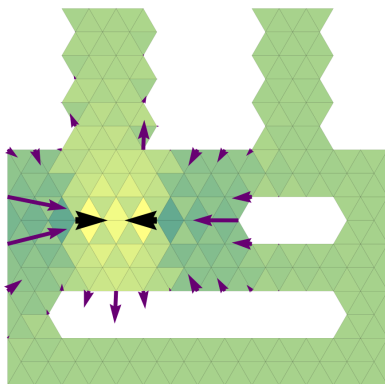
motor density

motor force dipole: $\mathcal{D} = \sum_i \mathbf{f}_i \cdot \mathbf{r}_i$

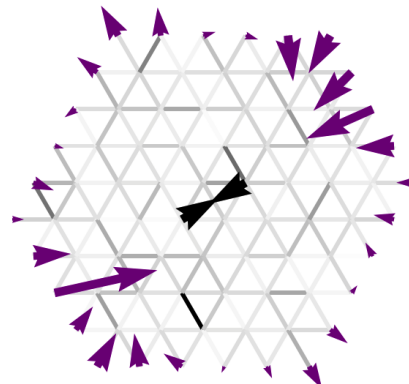
$$\Sigma = \int \text{Tr}(\sigma) dV$$

integrated "pressure"
Vanishes for homogeneous linear materials.

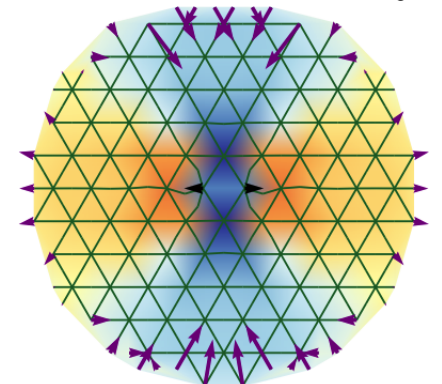
$\Sigma = 0$
even in strange geometries



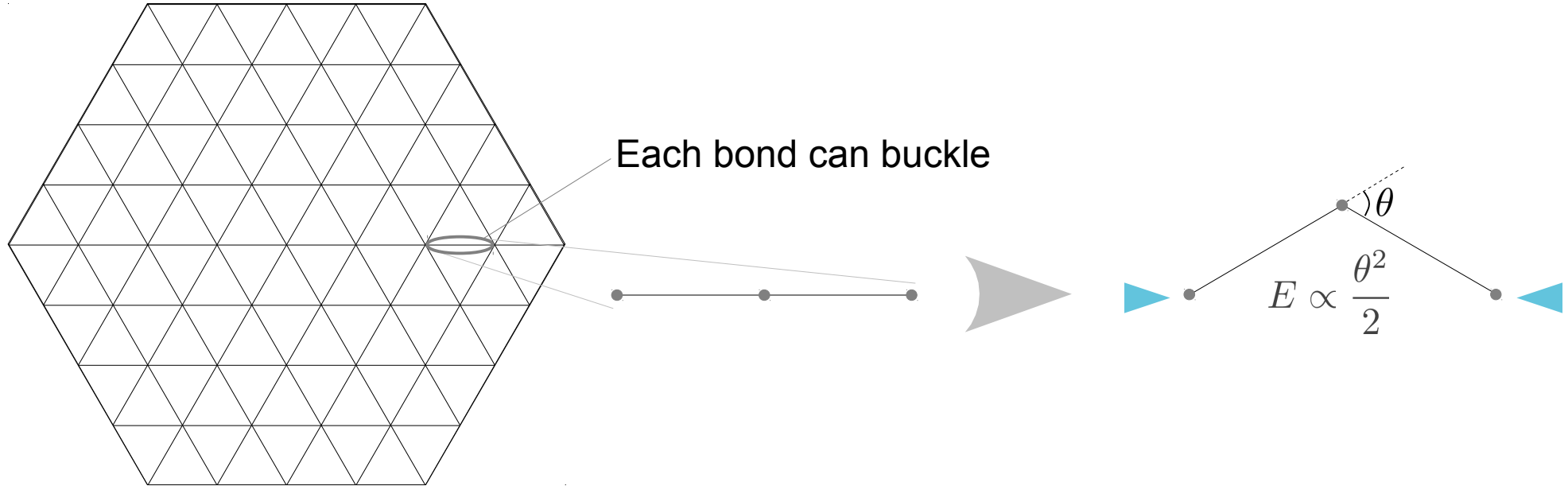
$\langle \Sigma \rangle = 0$
even in disordered materials



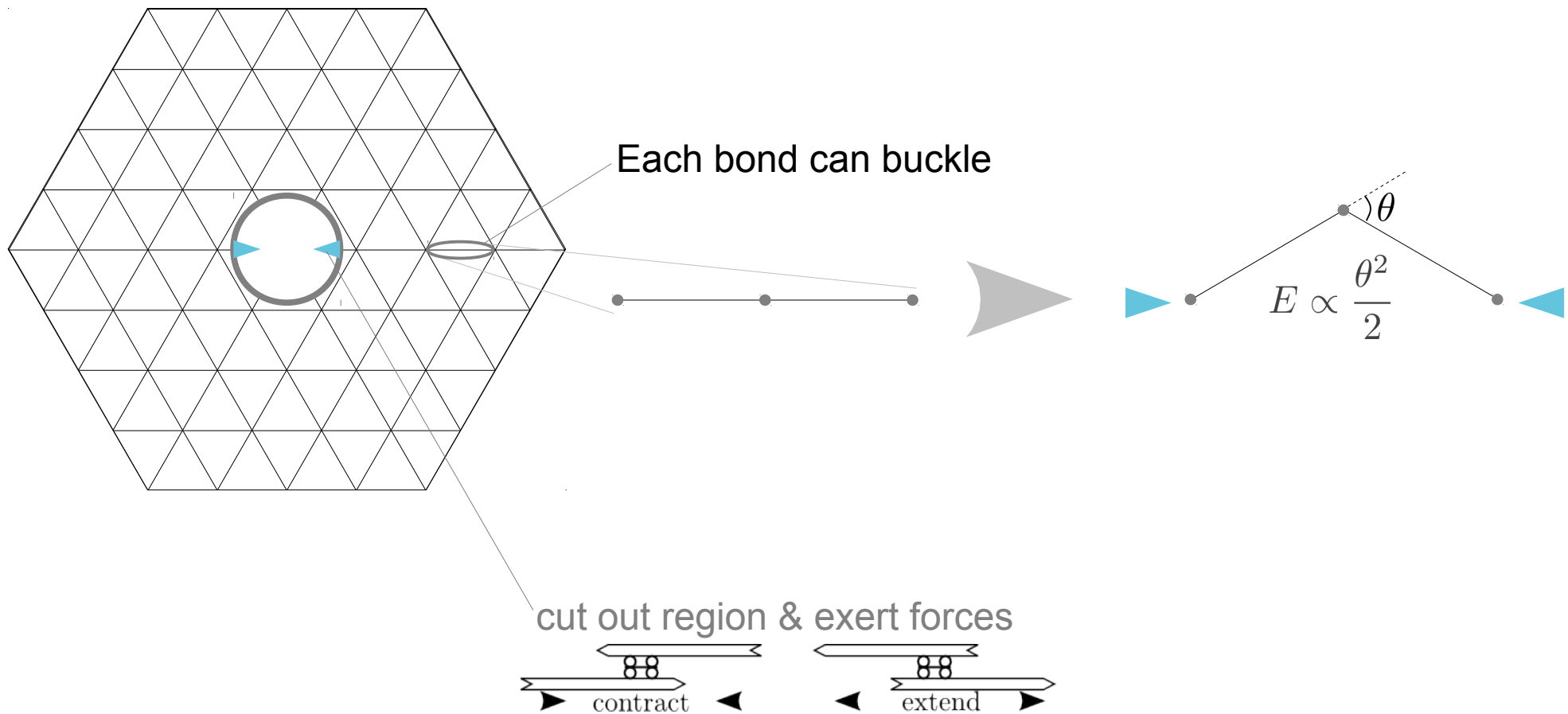
$\Sigma \neq 0$
in the presence of
nonlinear elasticity



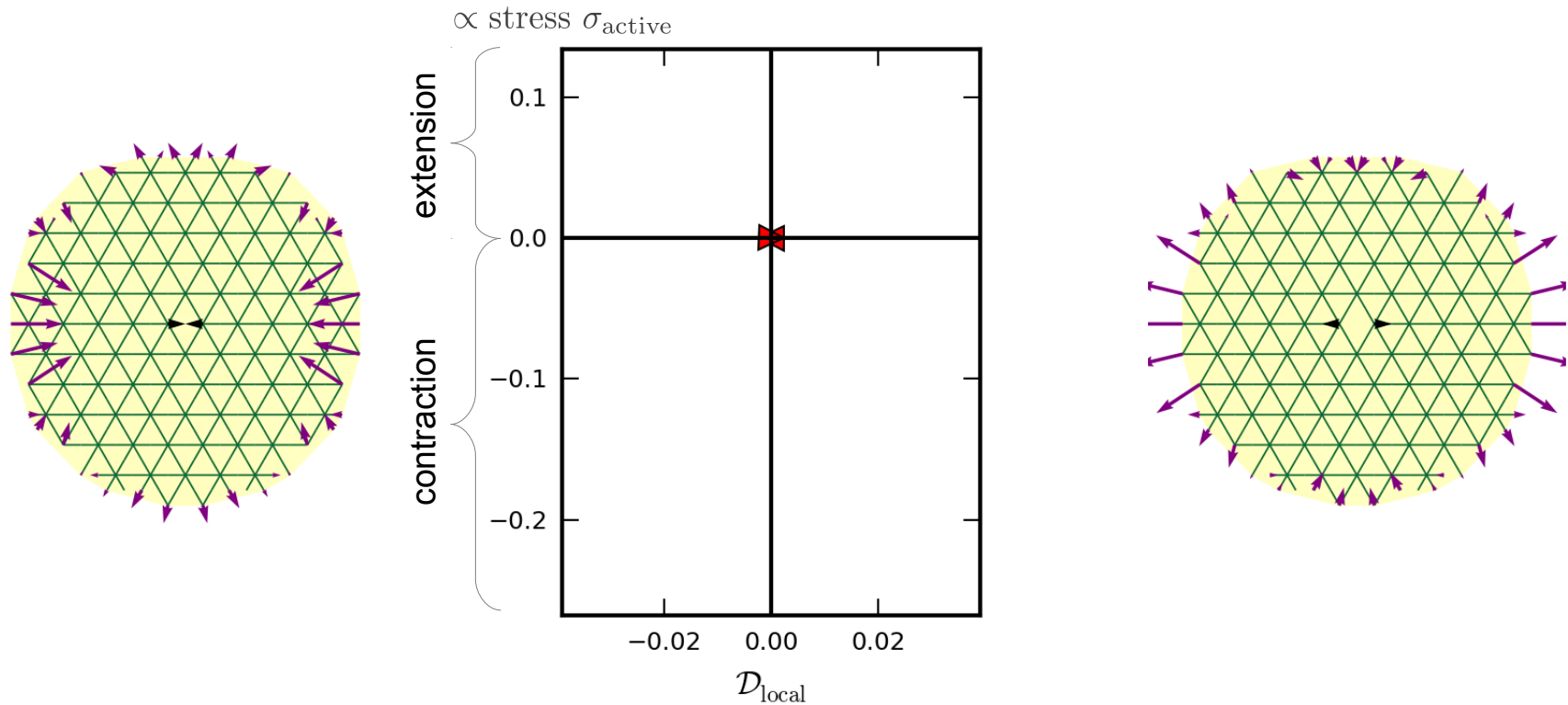
We consider a network with filaments susceptible to buckling



We consider a network with filaments susceptible to buckling



Forces below the buckling threshold yield linear force transmission

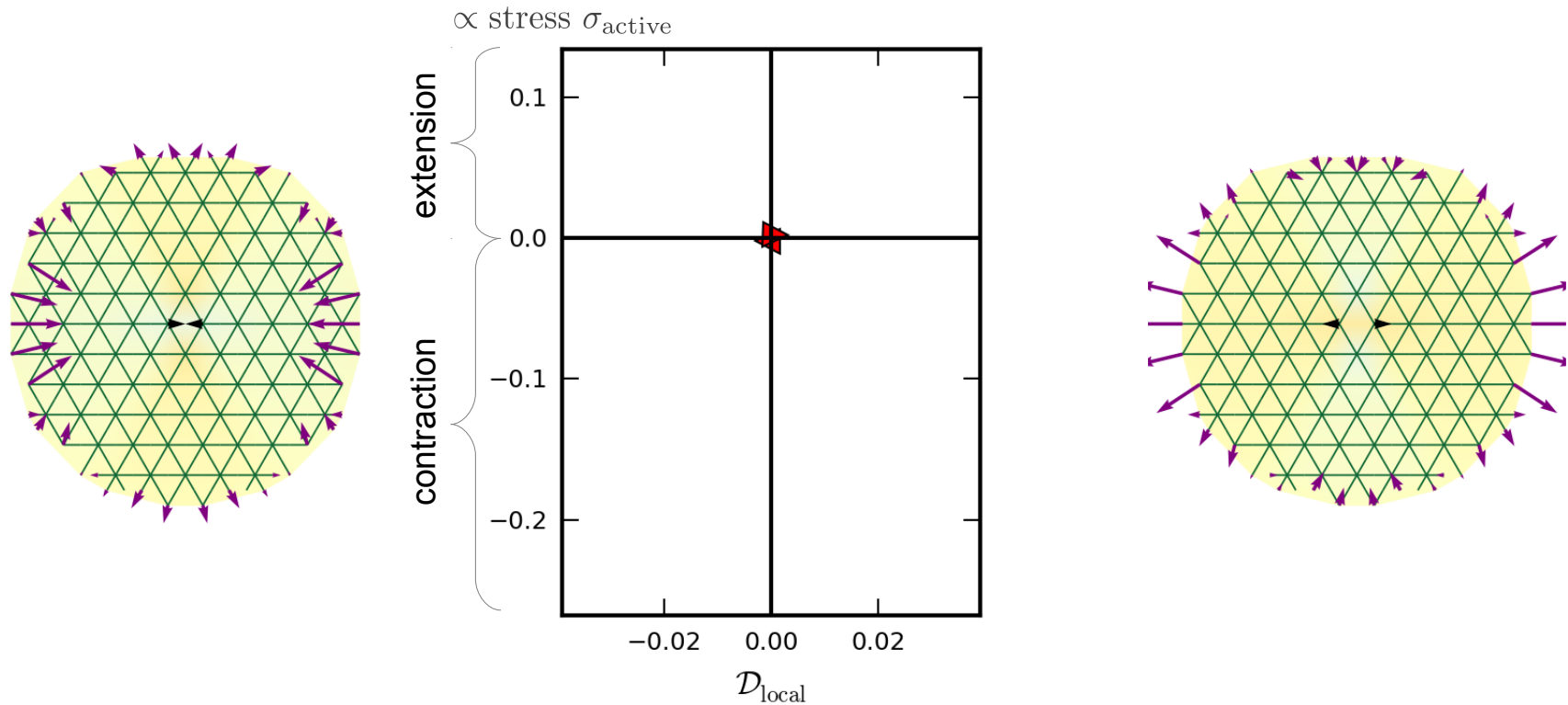


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Forces below the buckling threshold yield linear force transmission

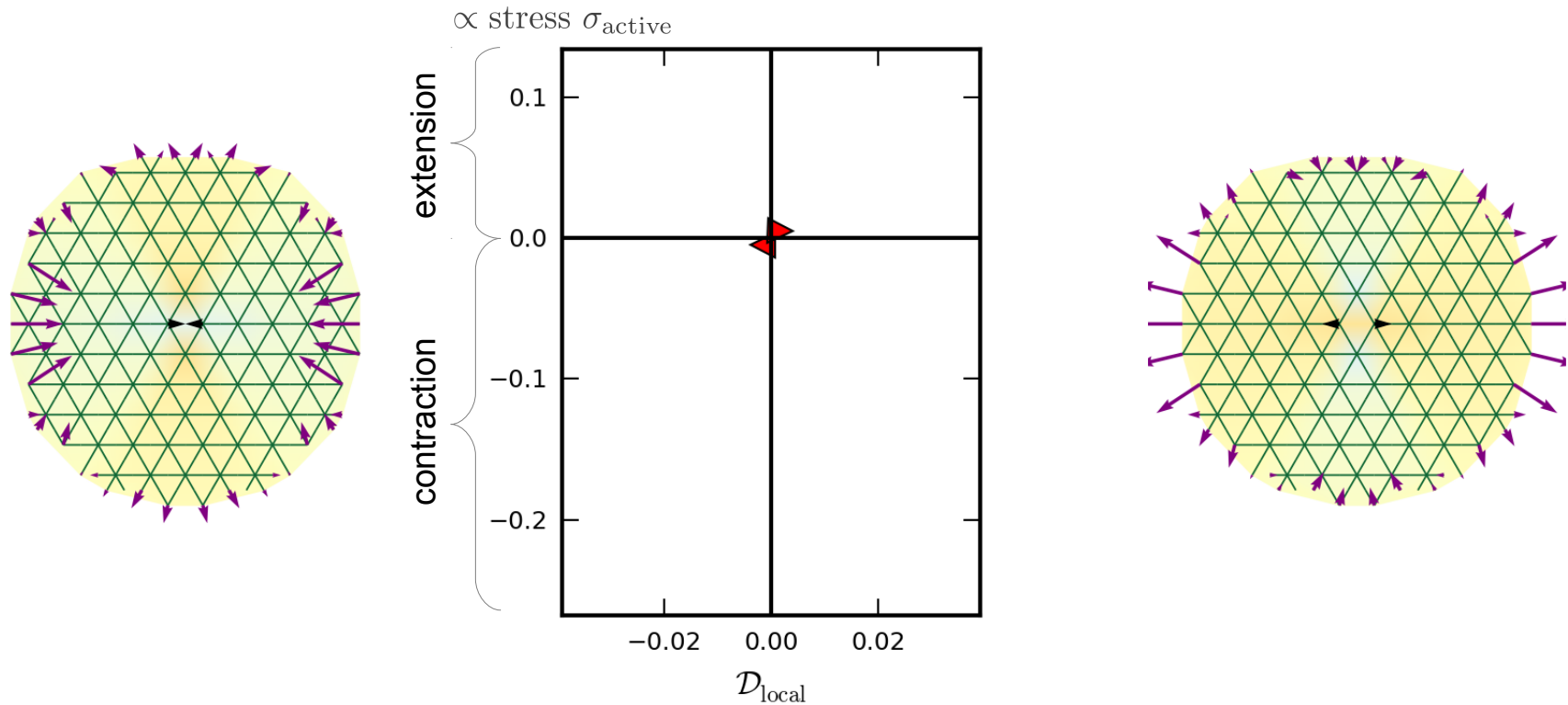


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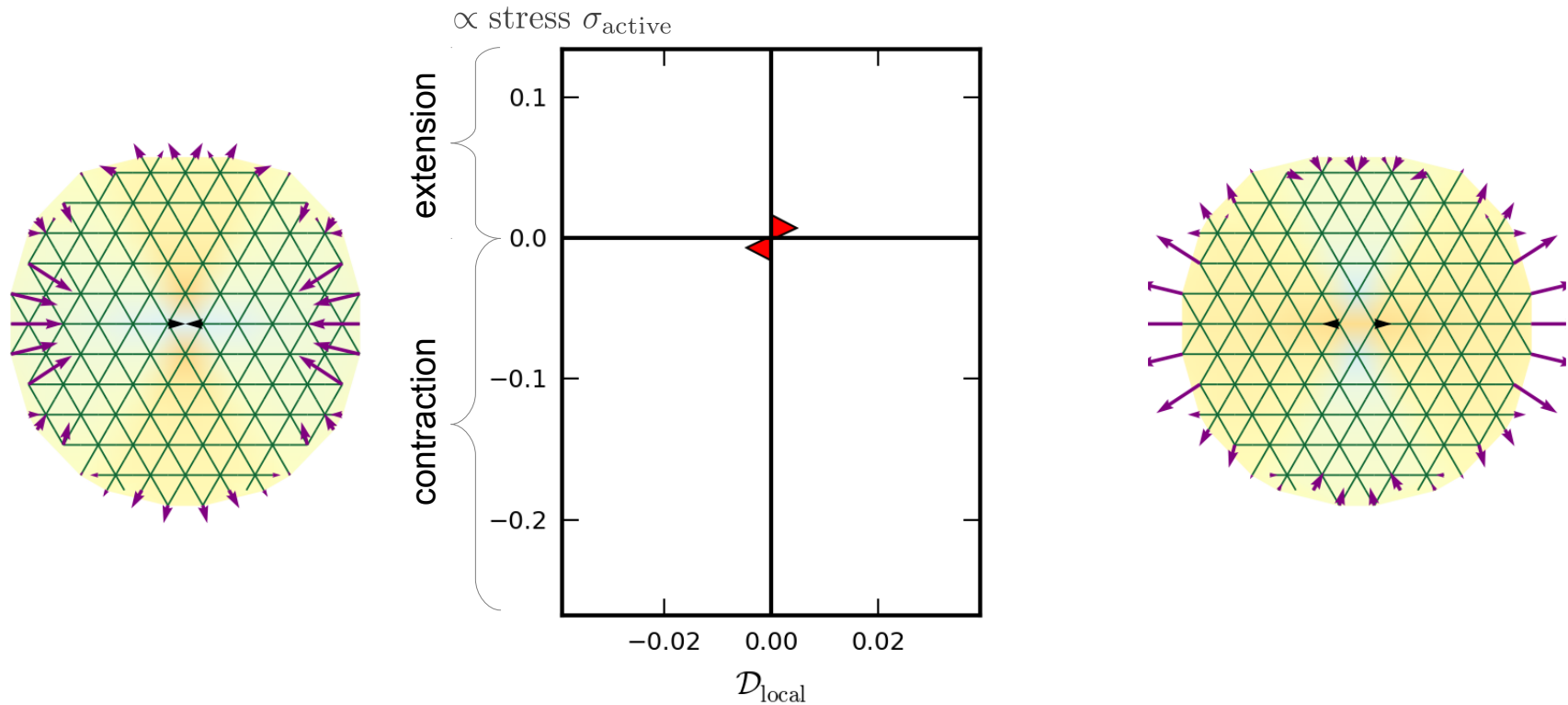


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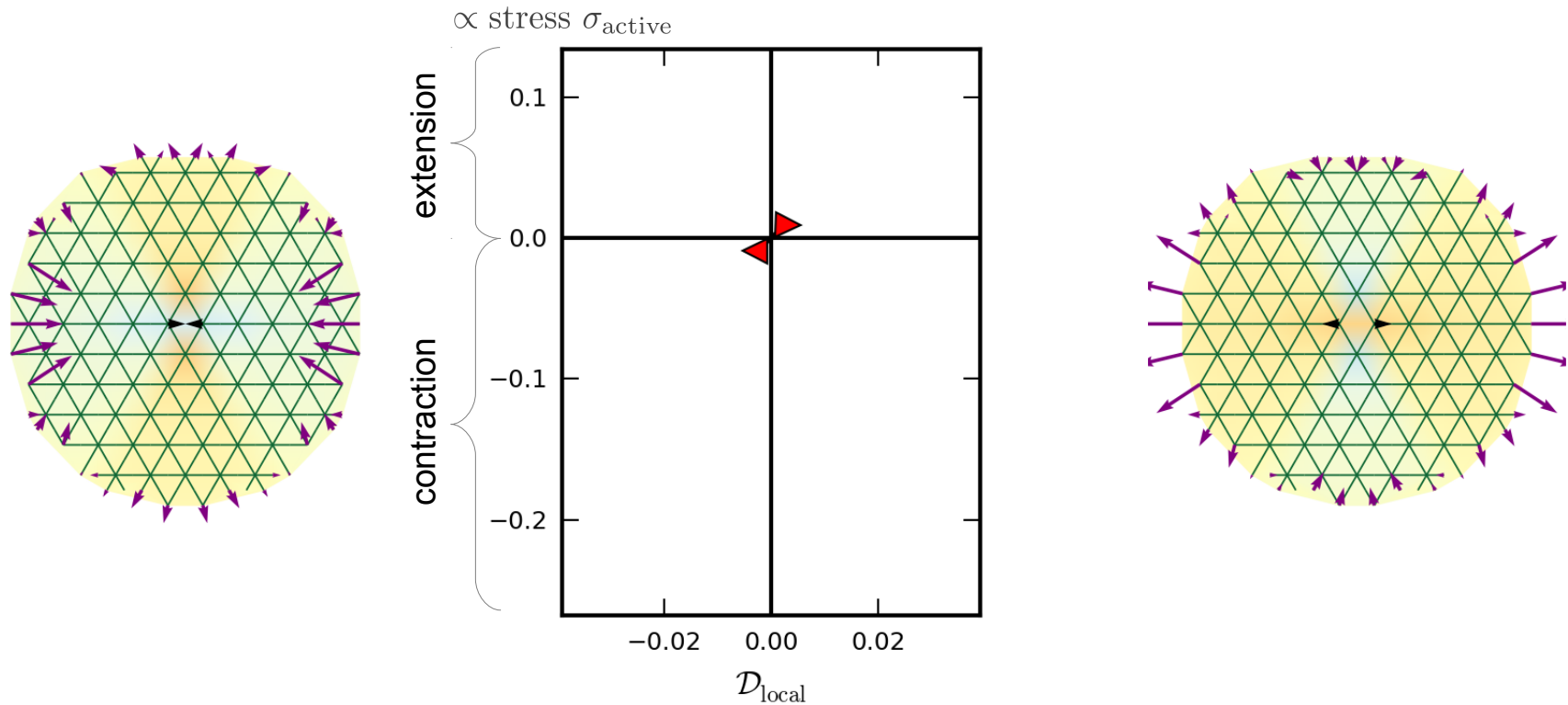


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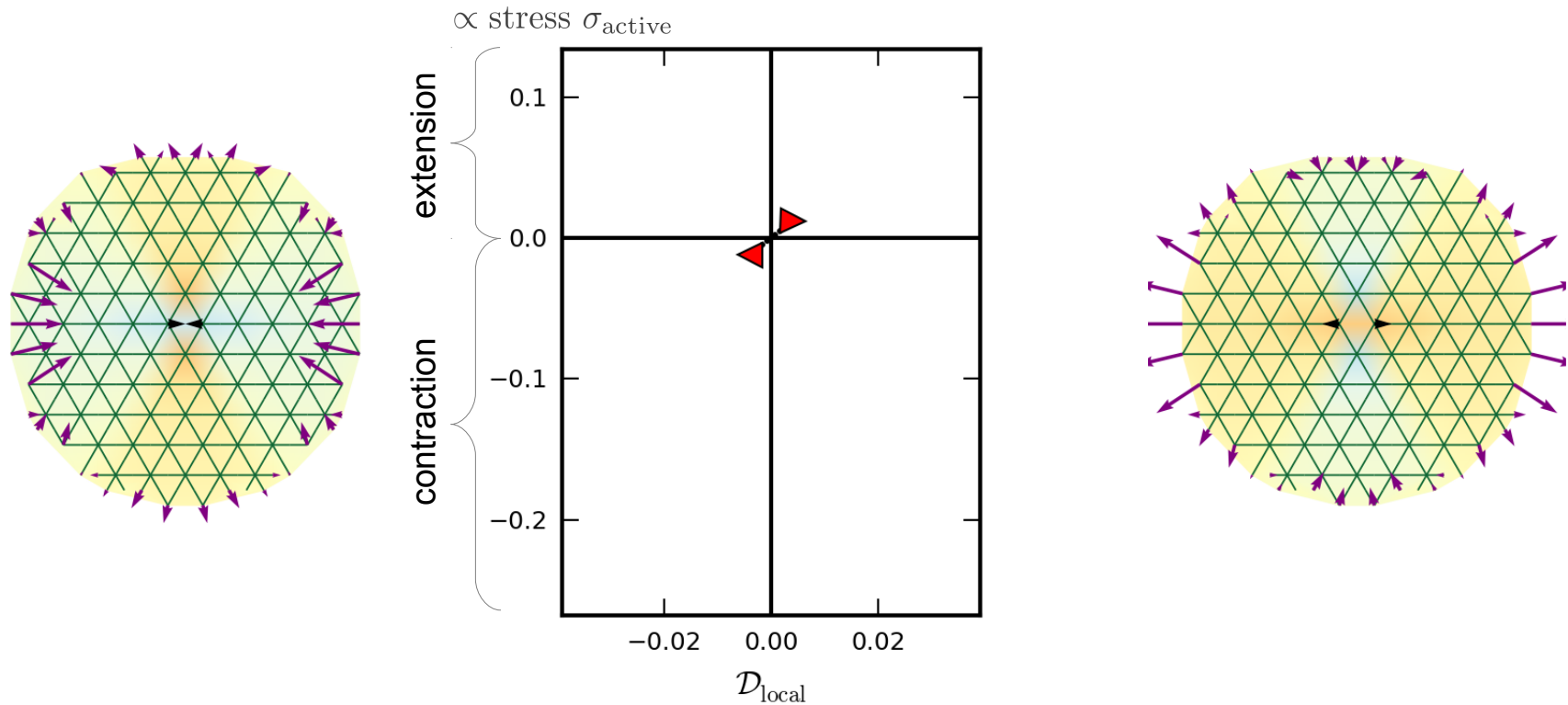


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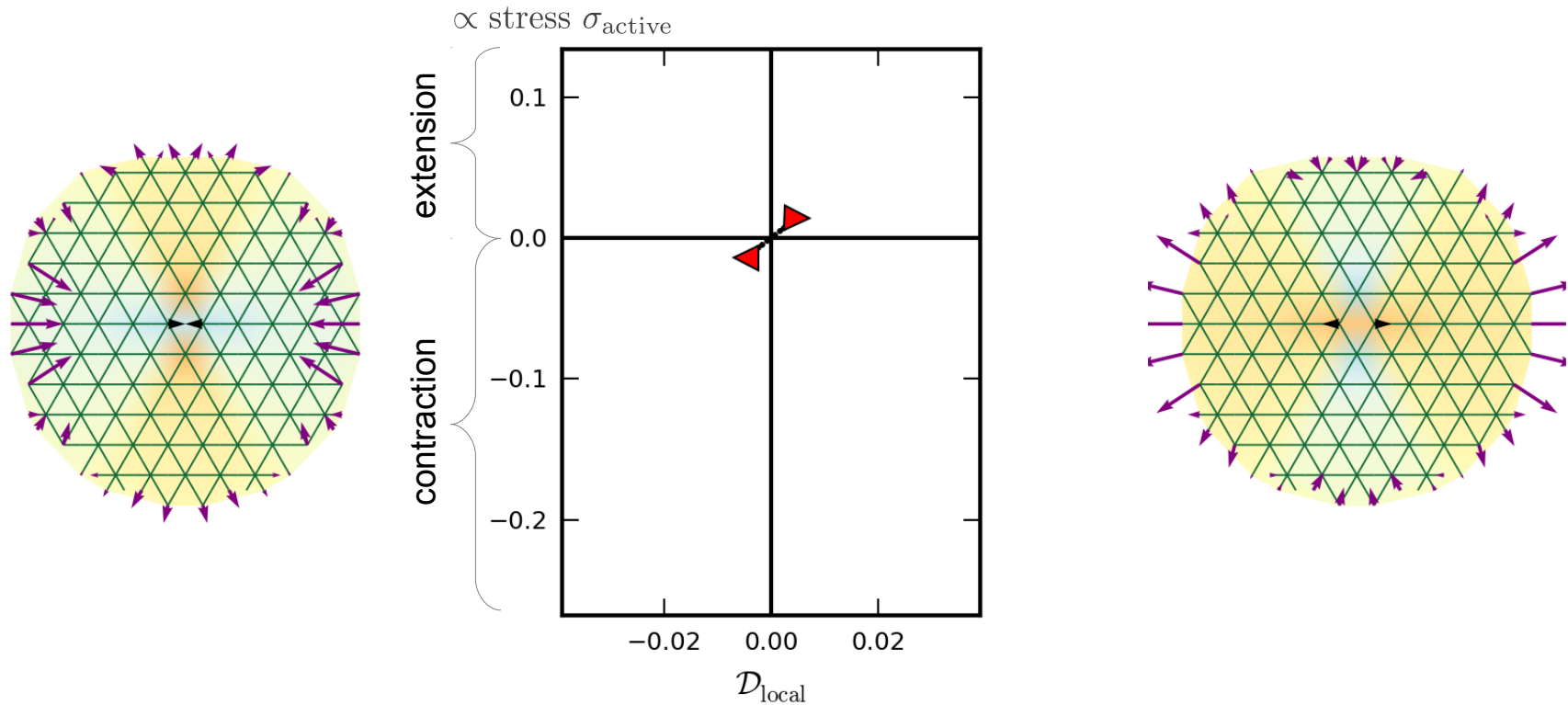


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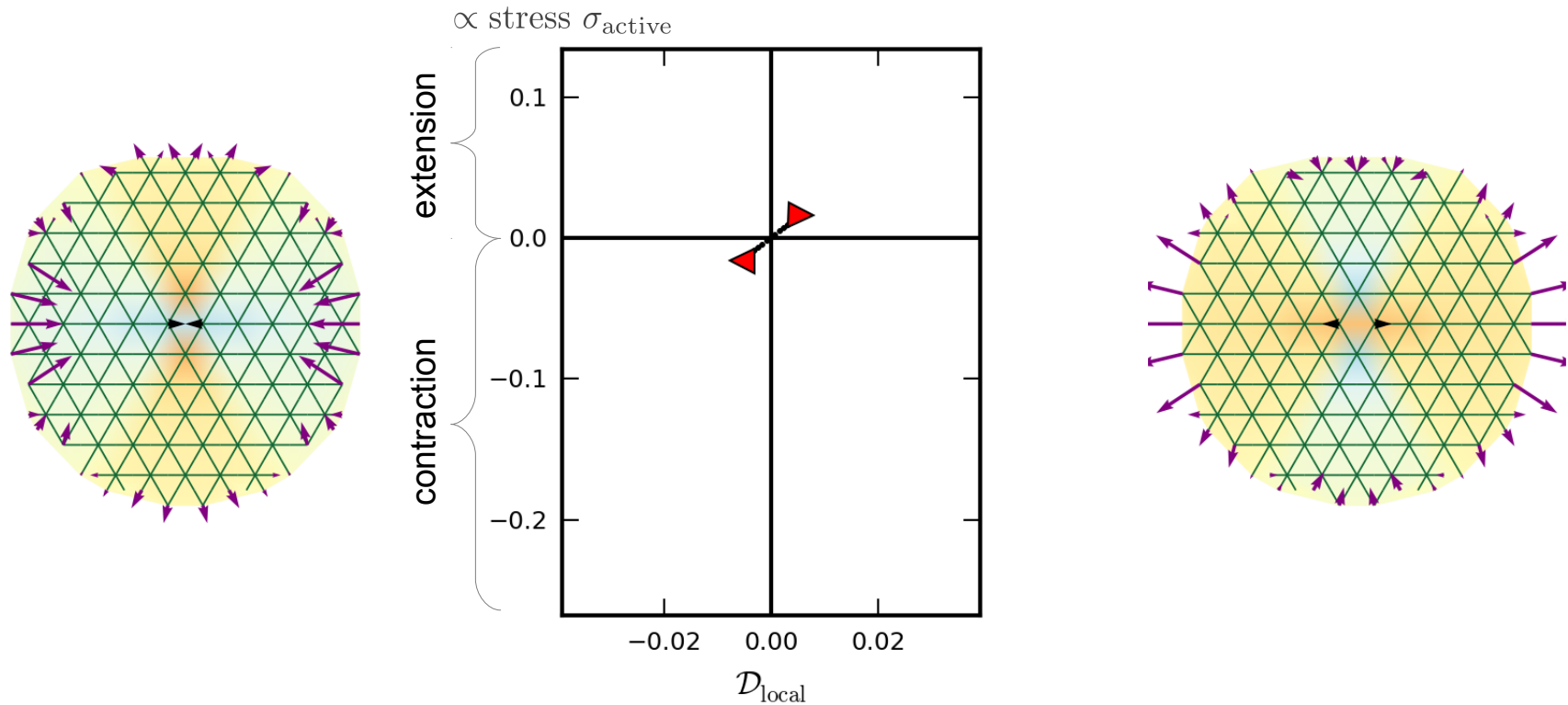


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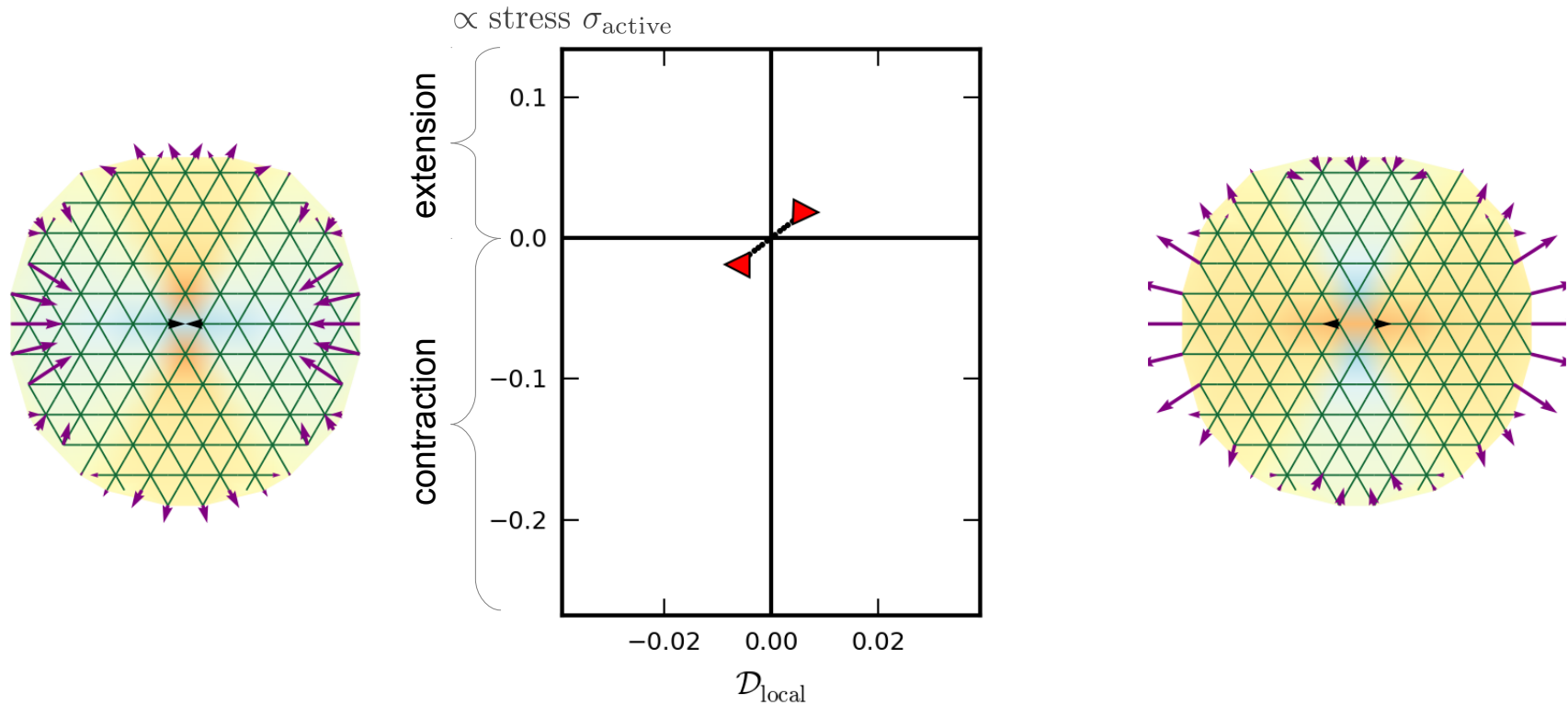


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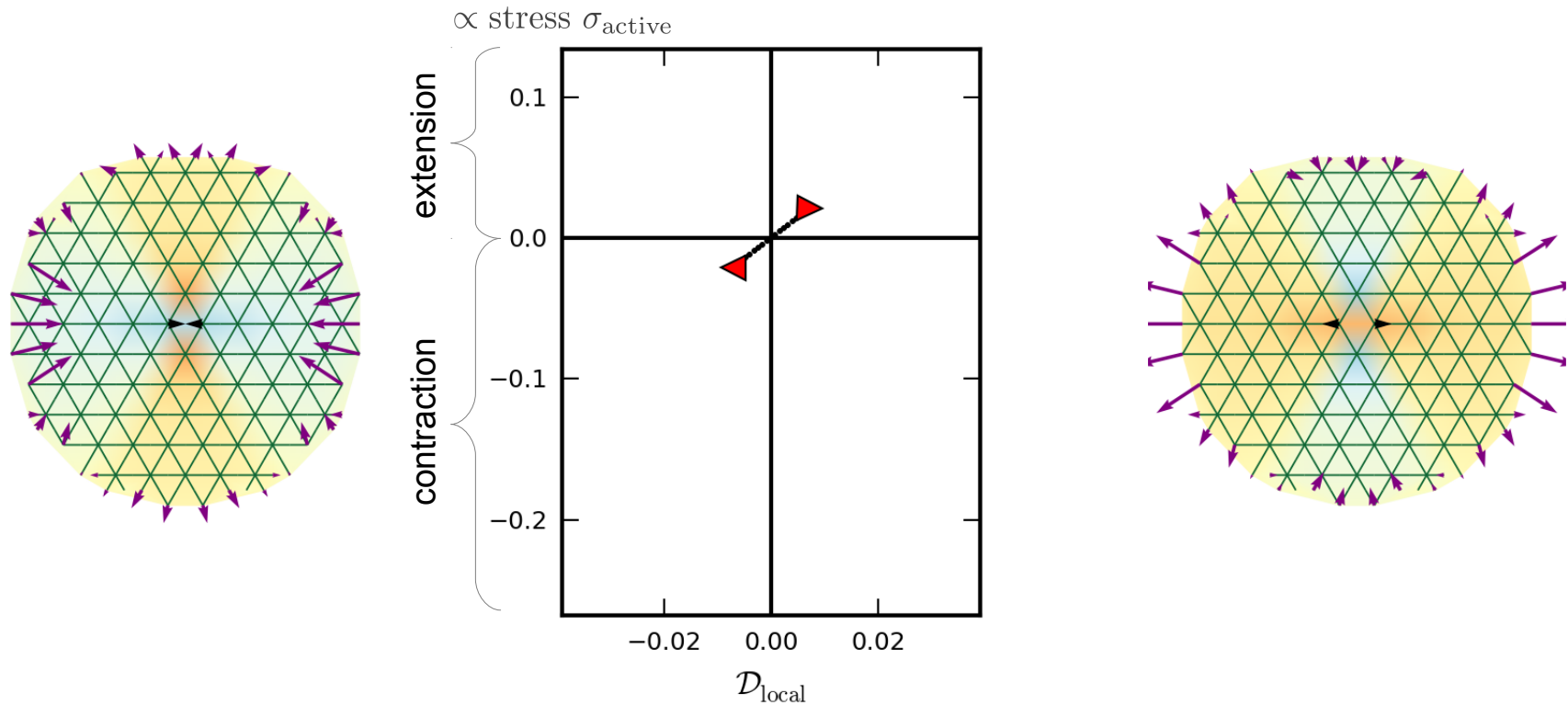
extension

Forces below the buckling threshold yield linear force transmission



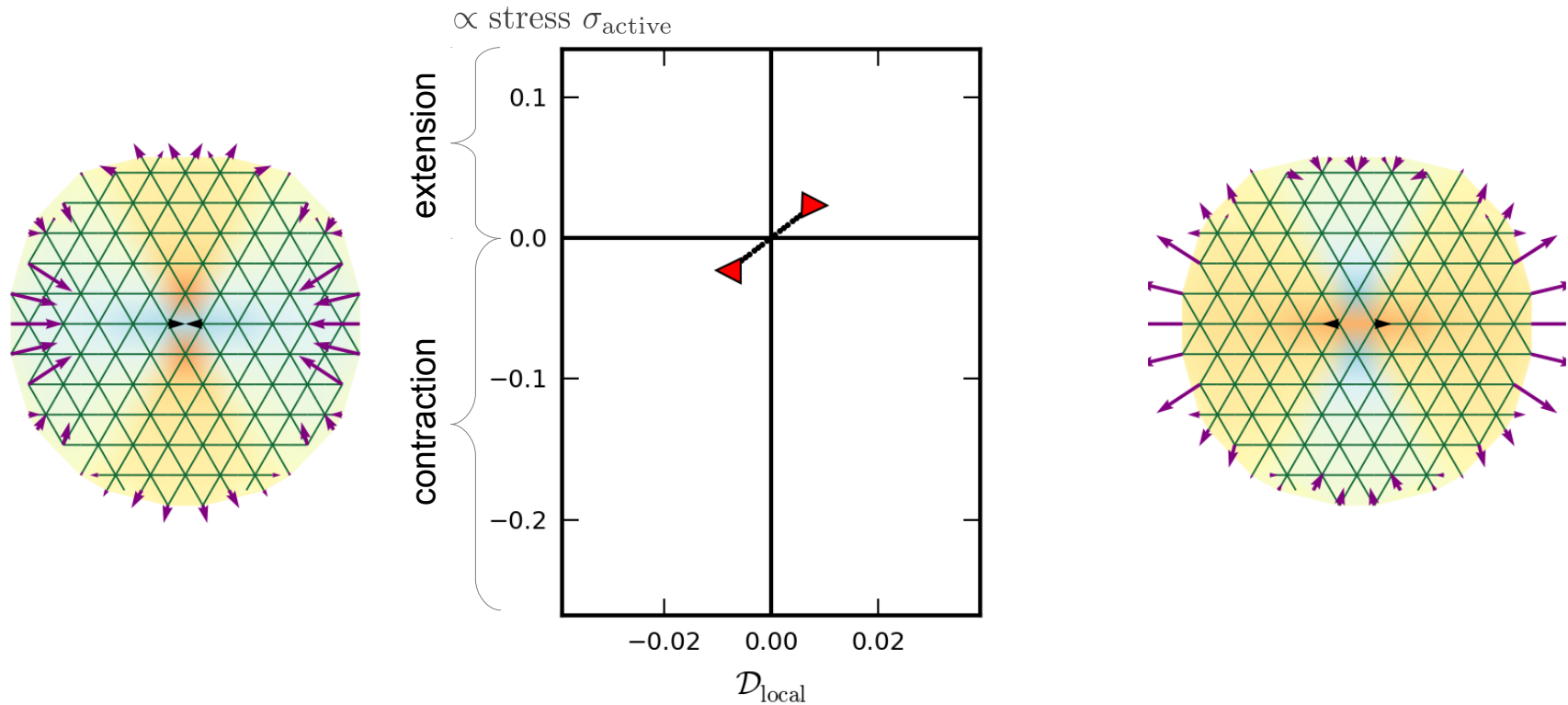
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Forces below the buckling threshold yield linear force transmission



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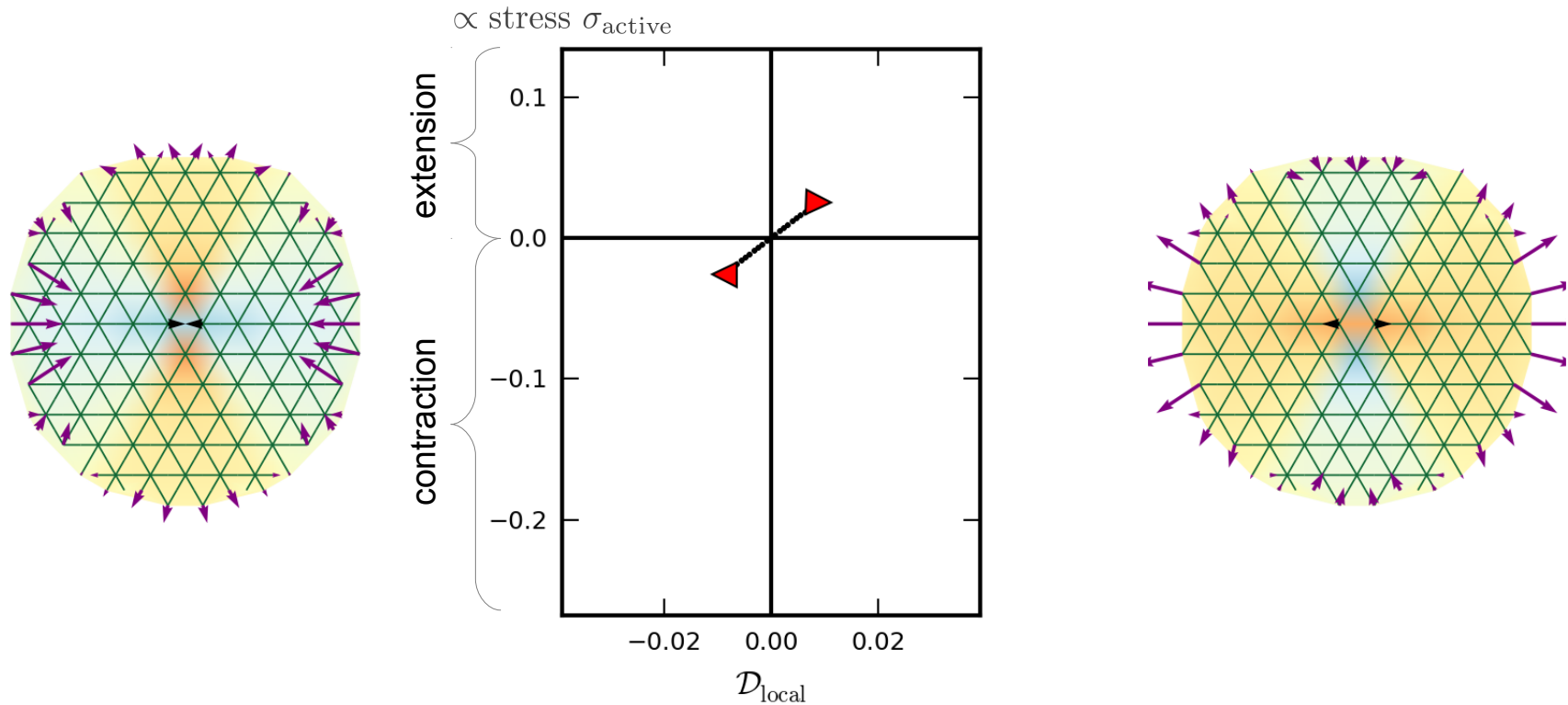


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Forces below the buckling threshold yield linear force transmission

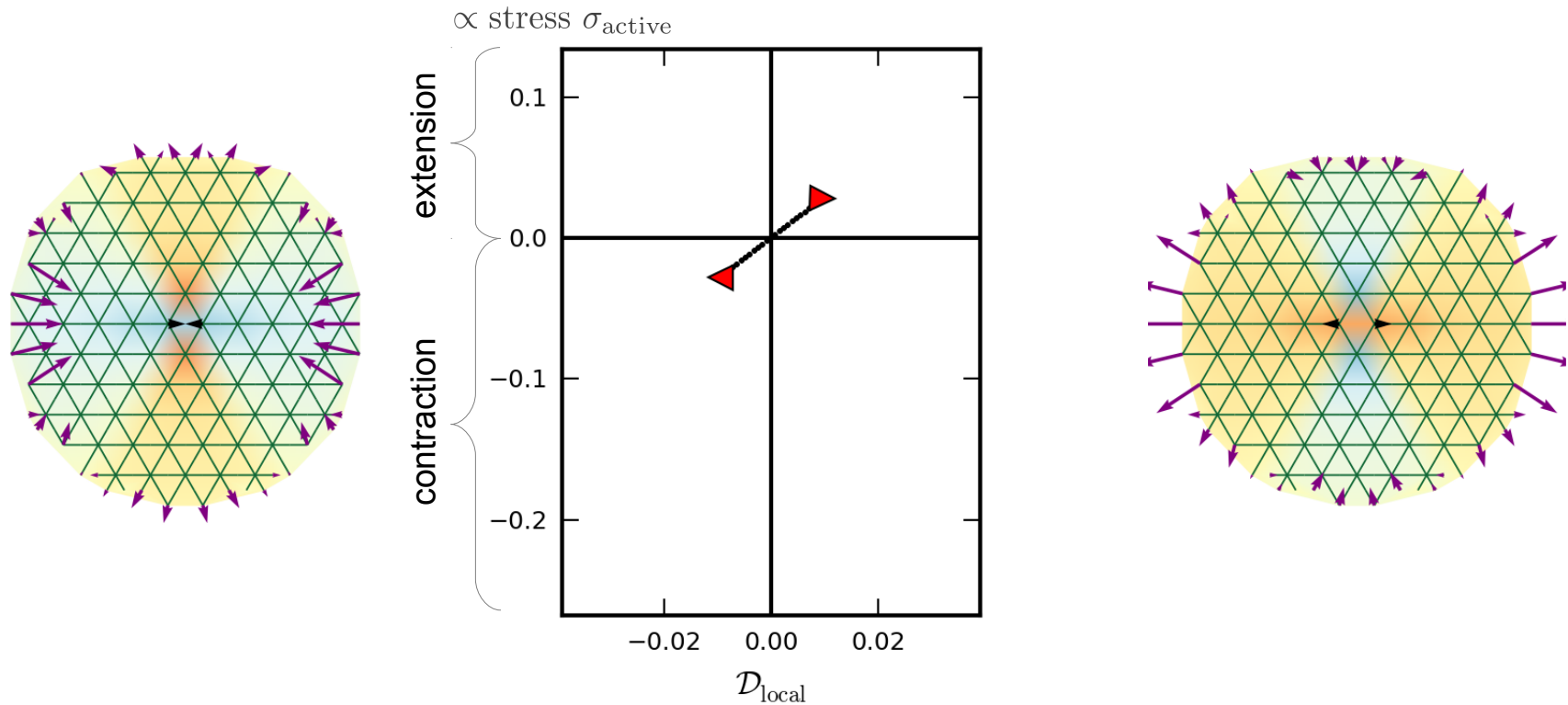


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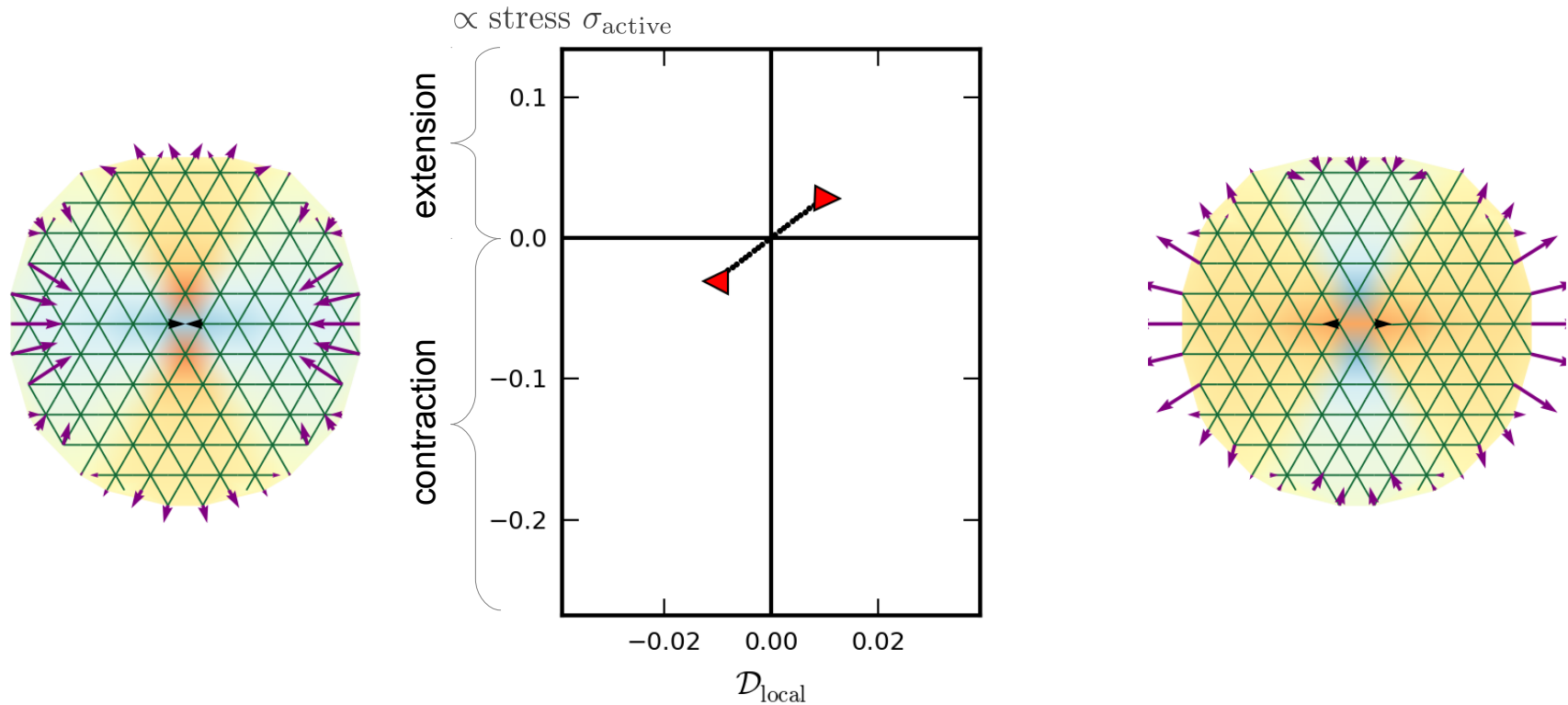
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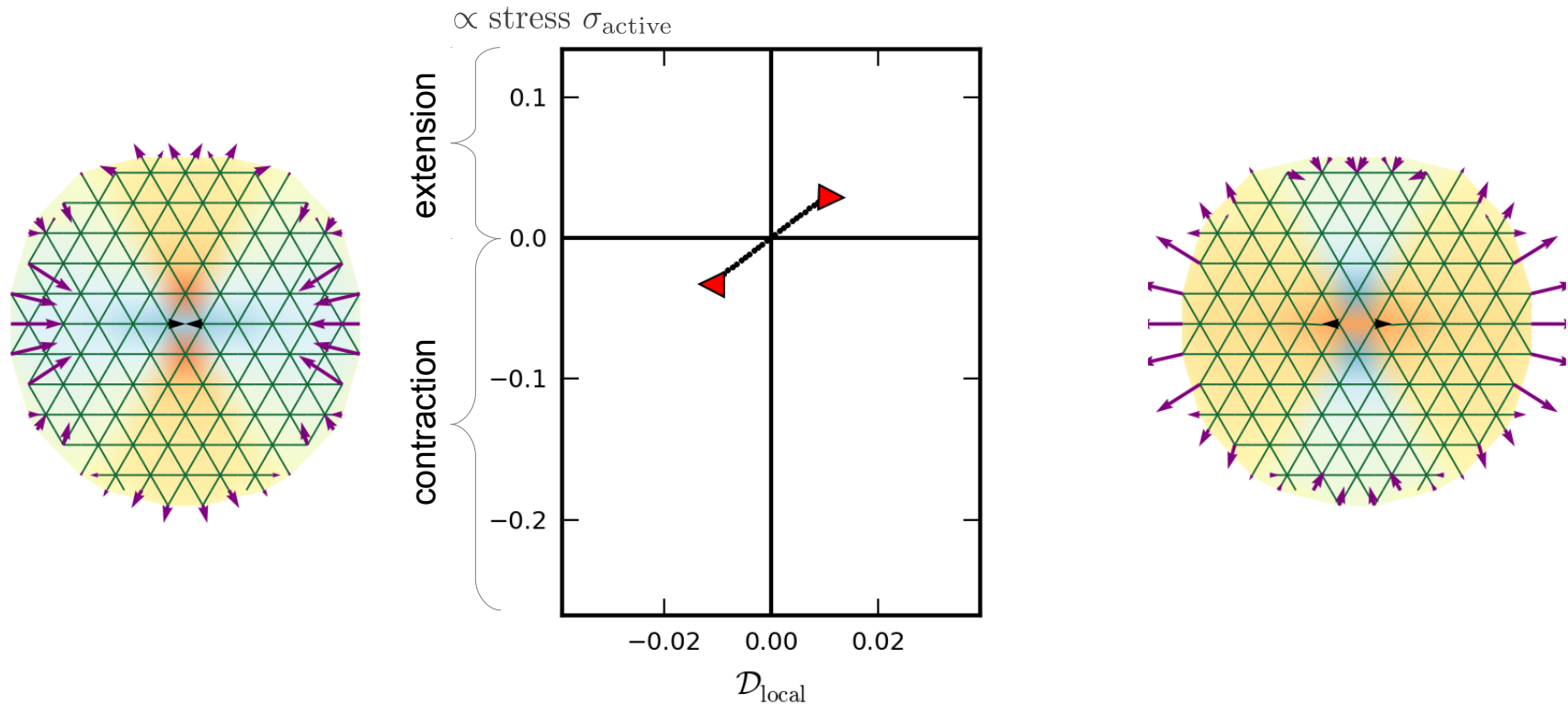
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Buckling rectifies forces of any sign towards contraction



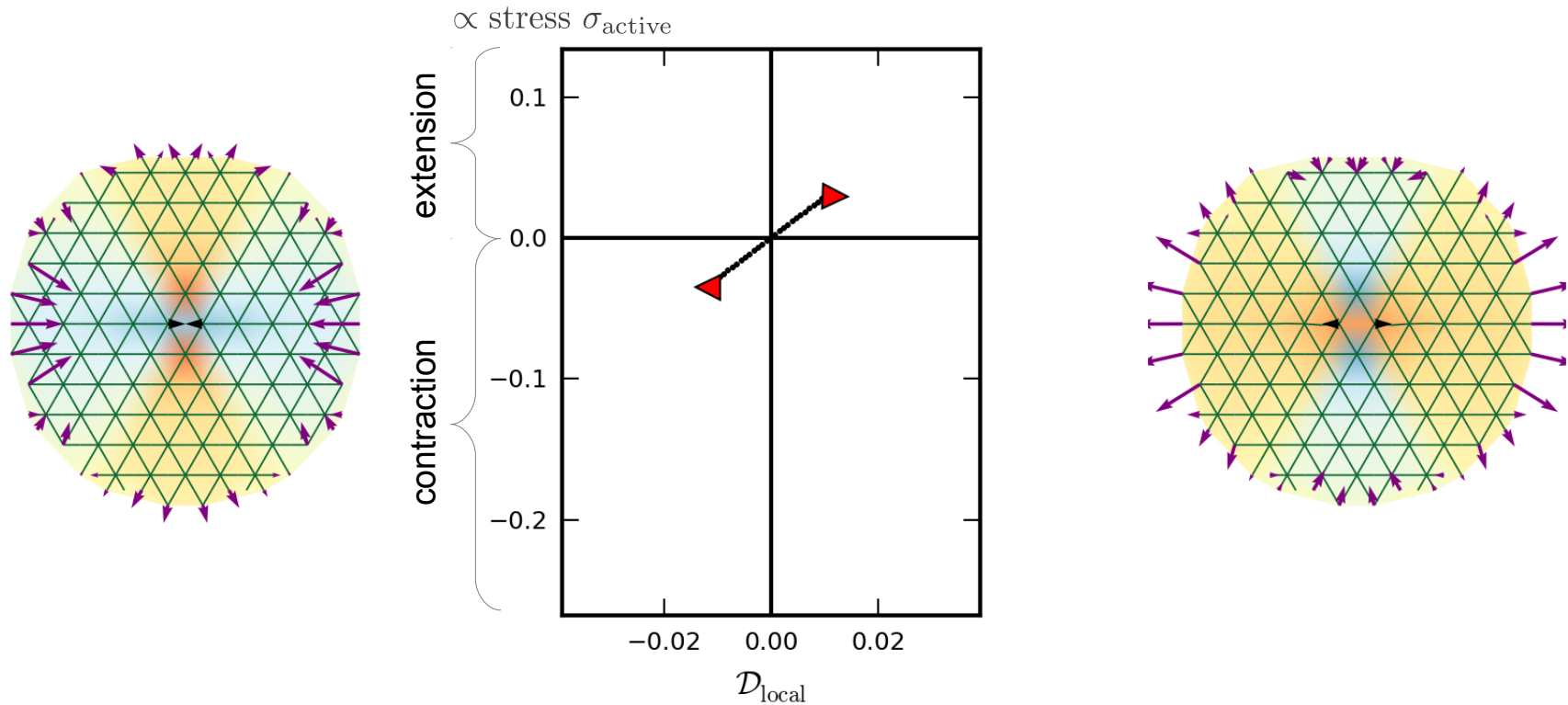
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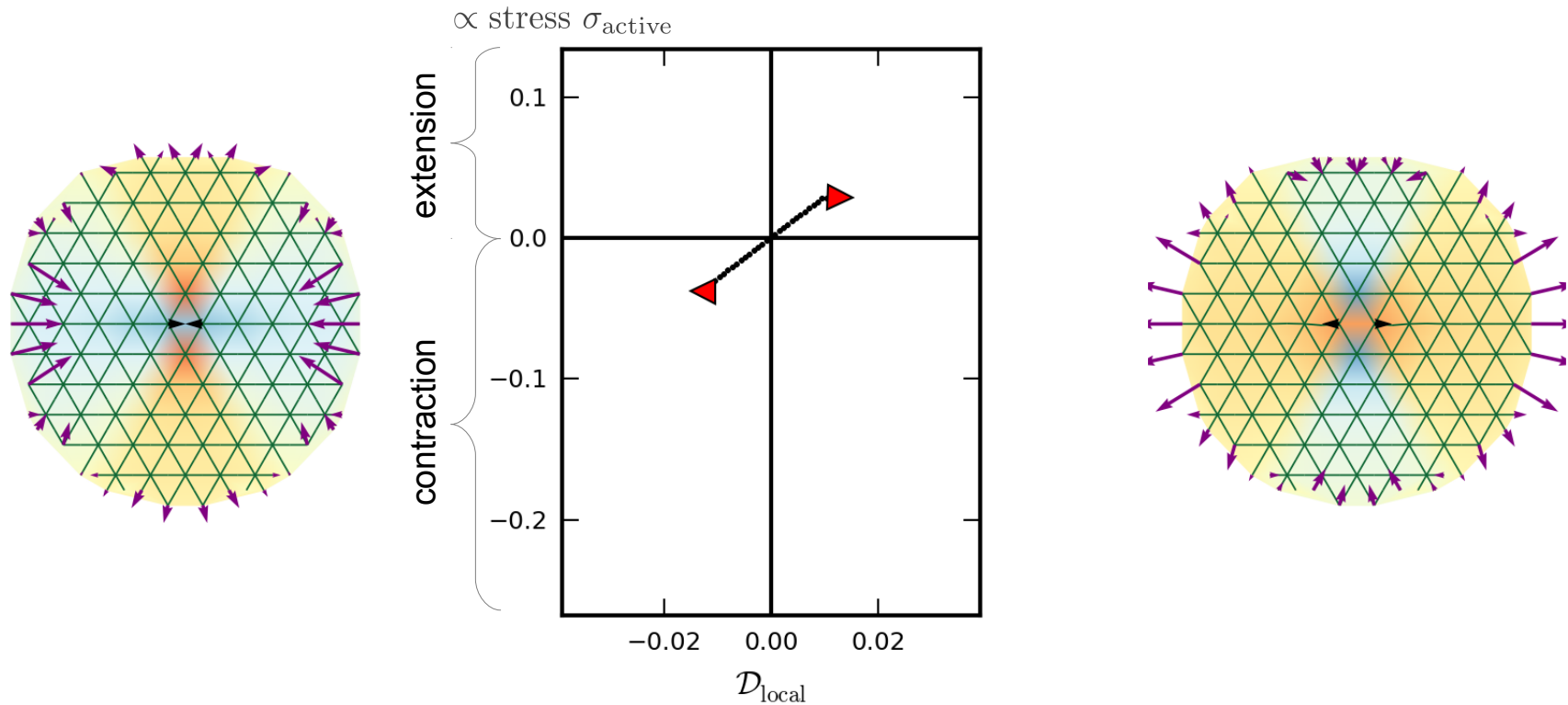
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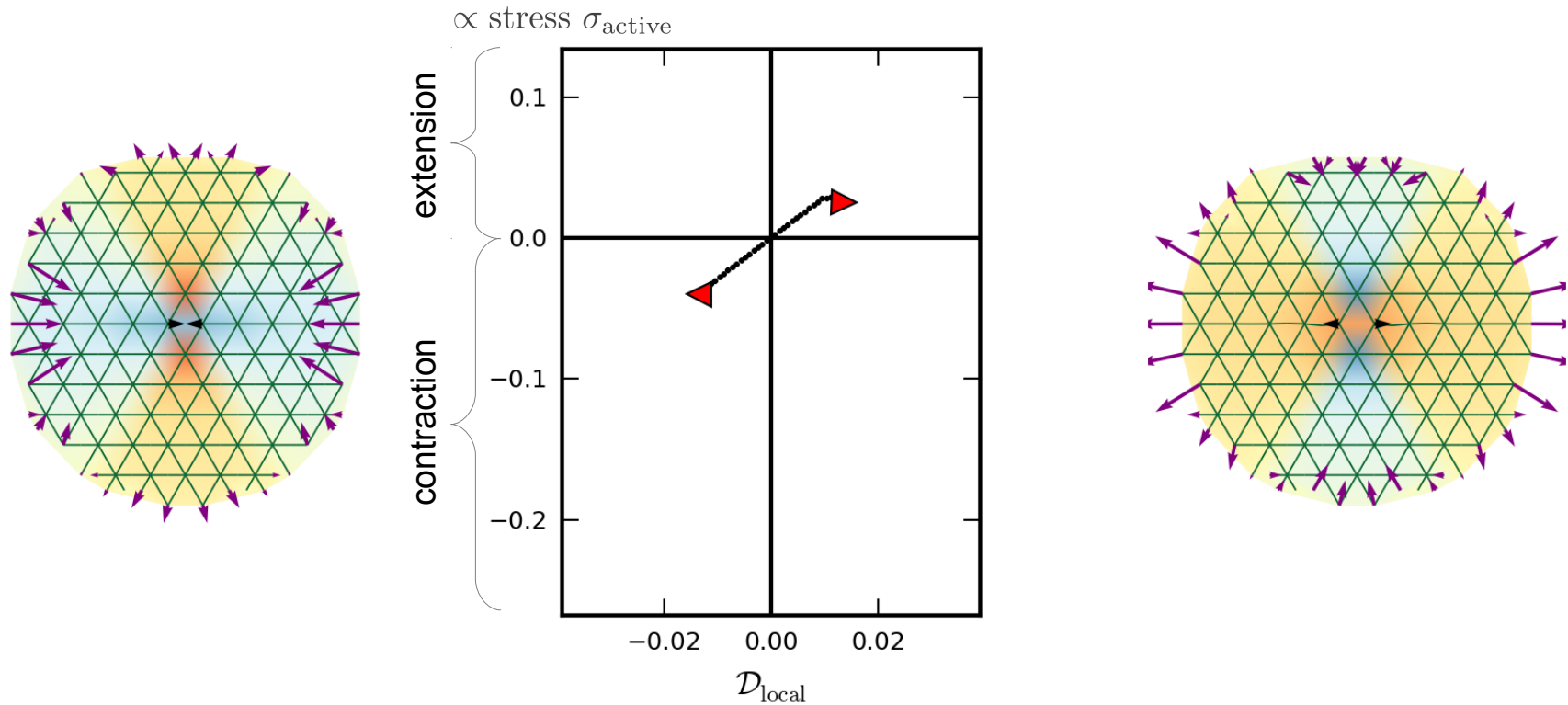
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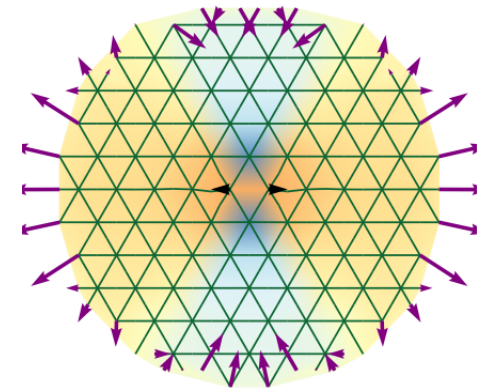
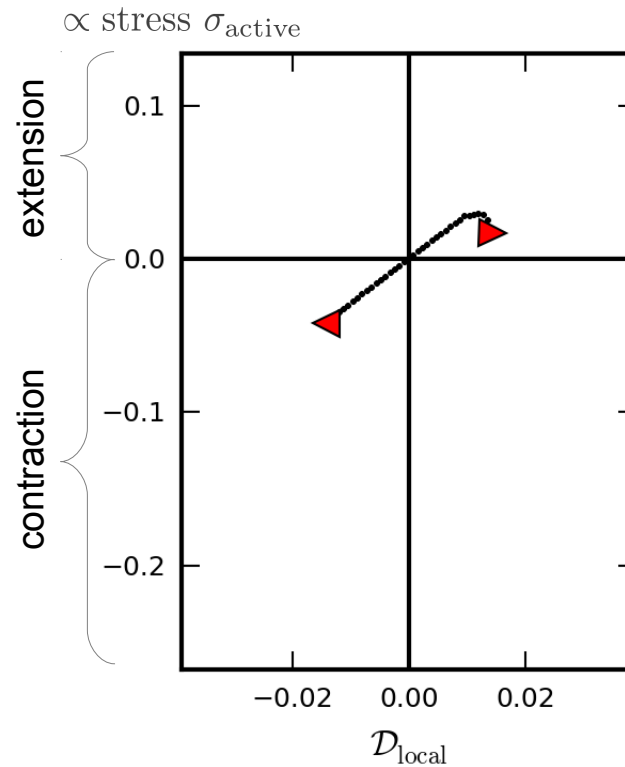
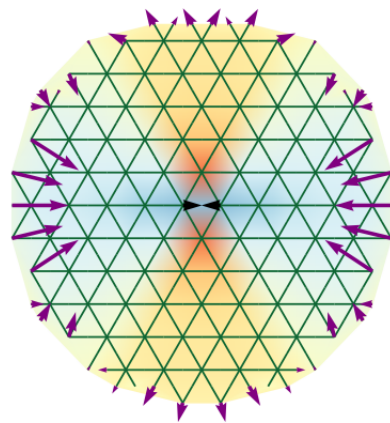
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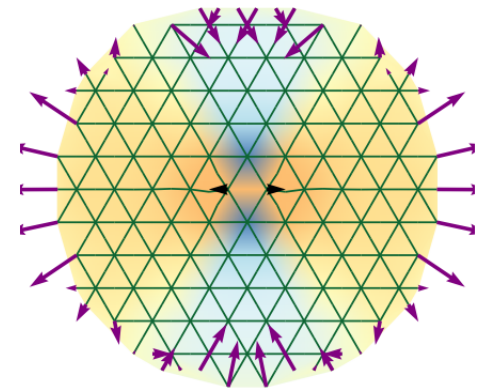
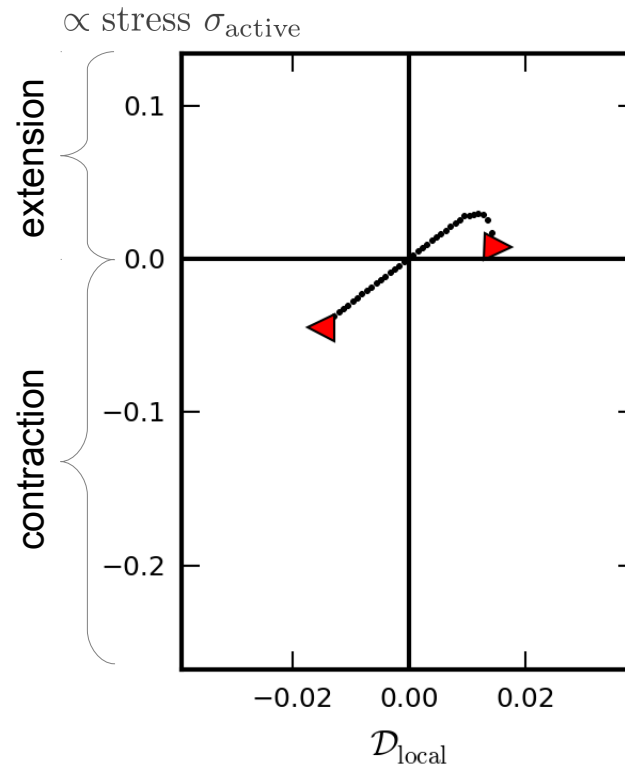
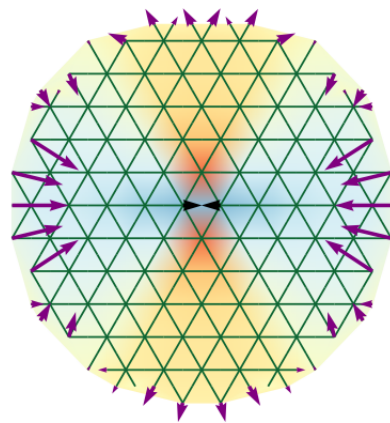
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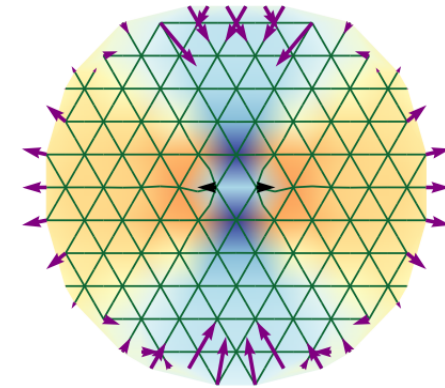
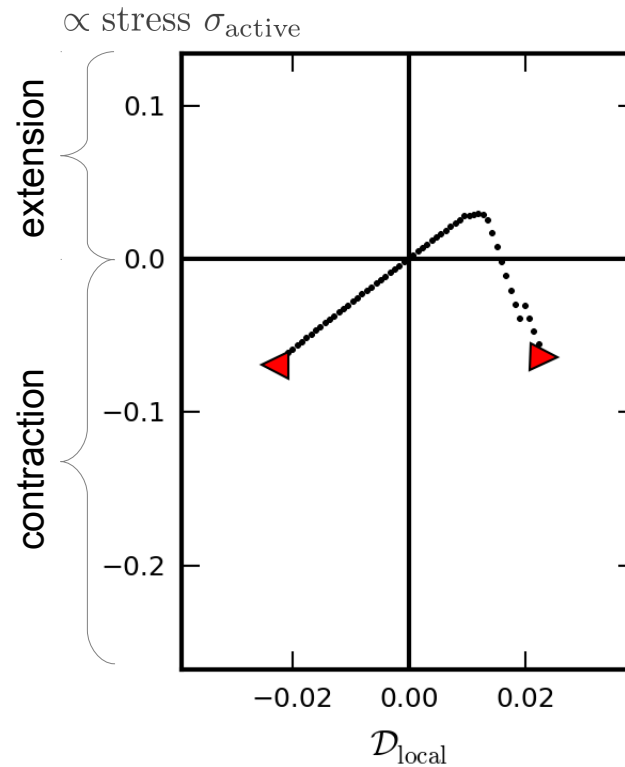
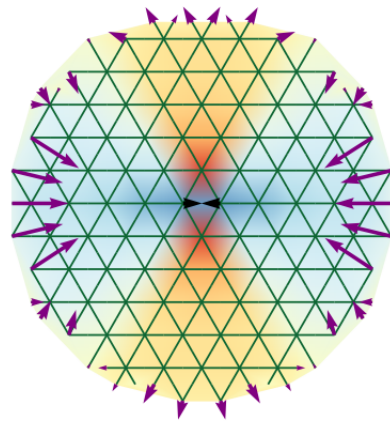
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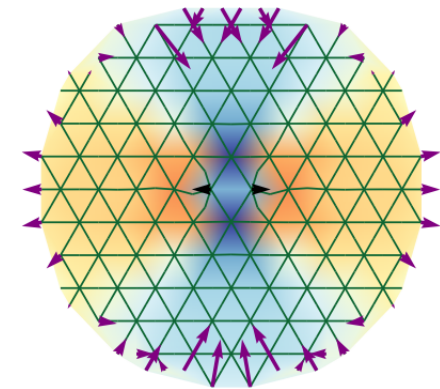
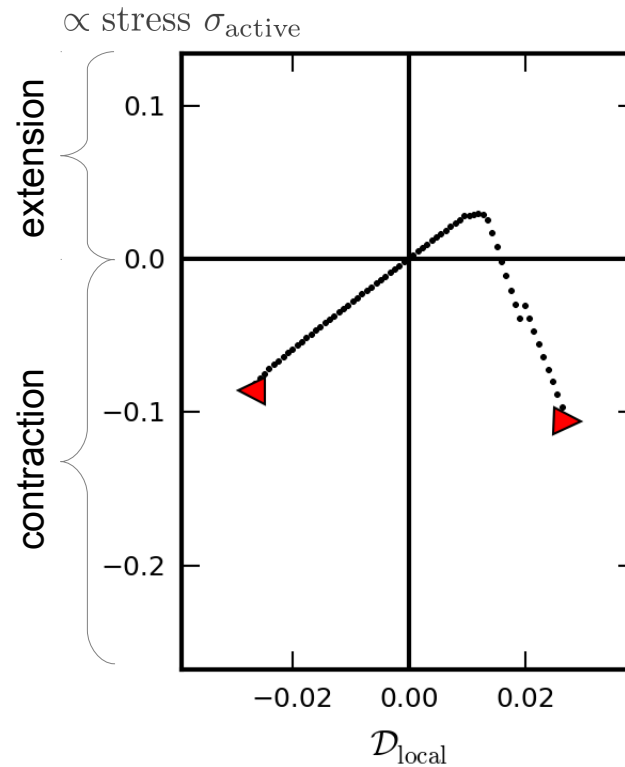
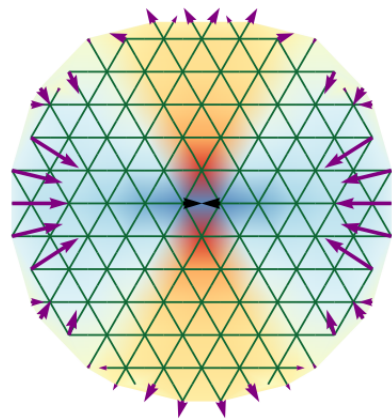
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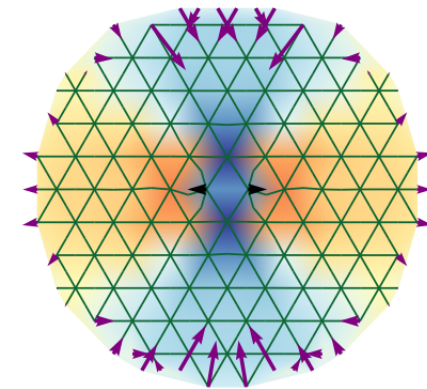
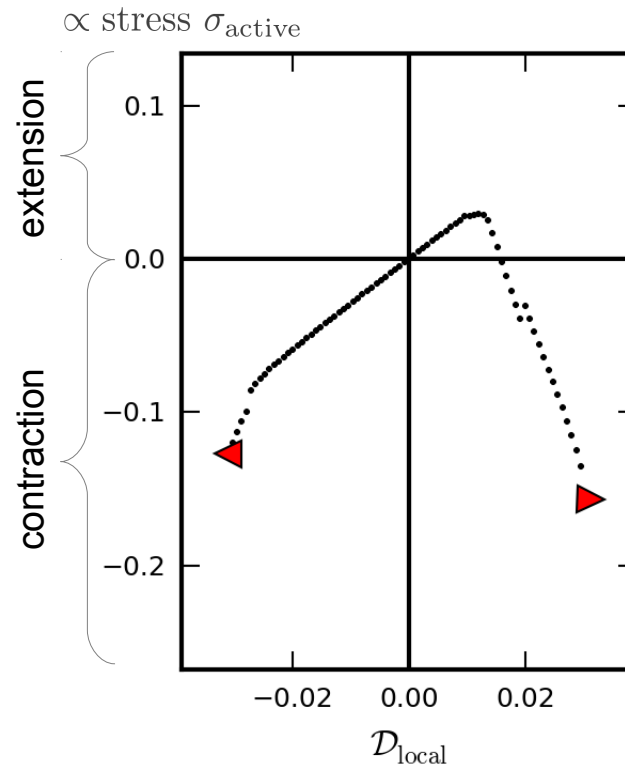
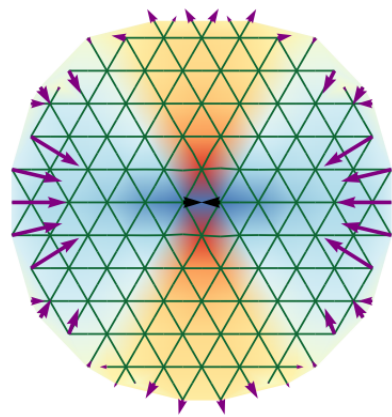
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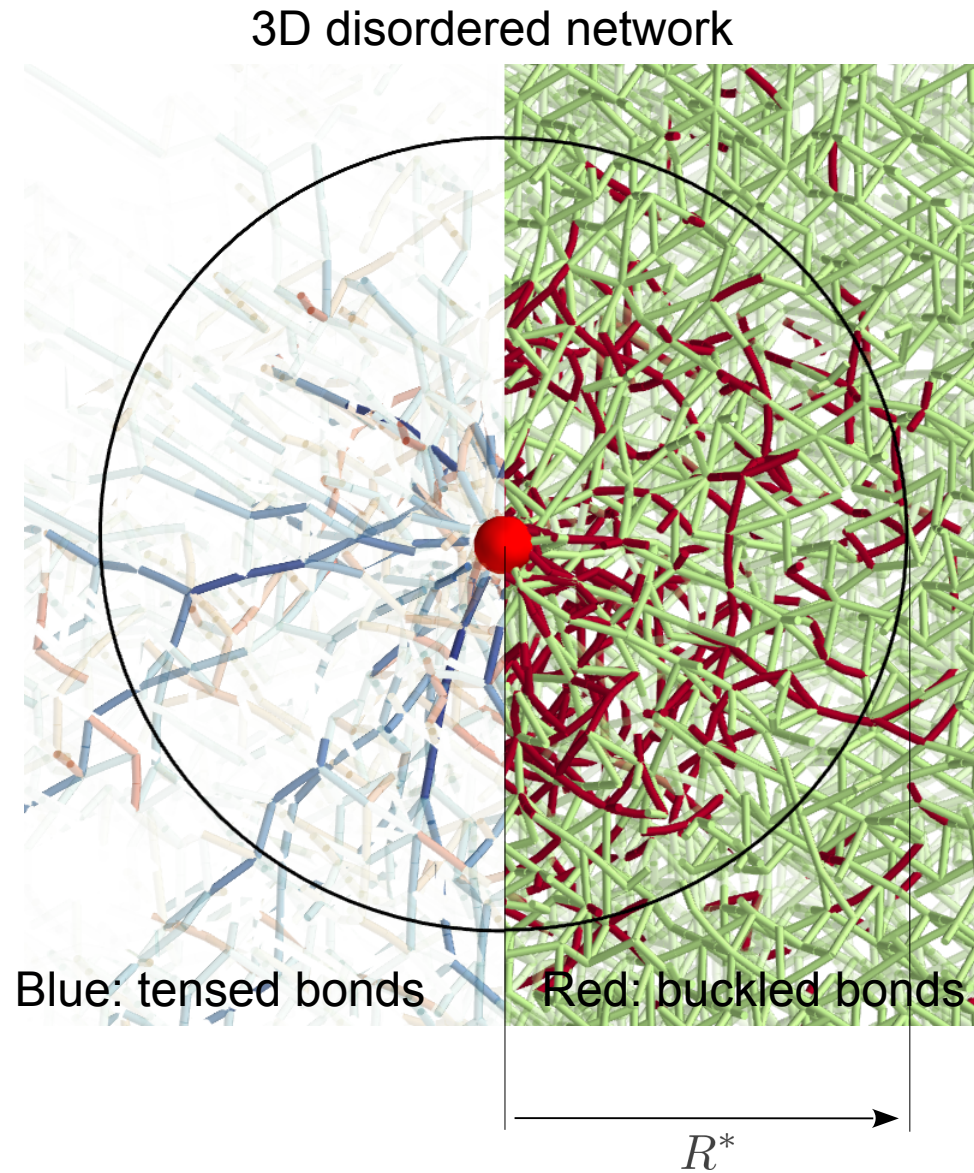
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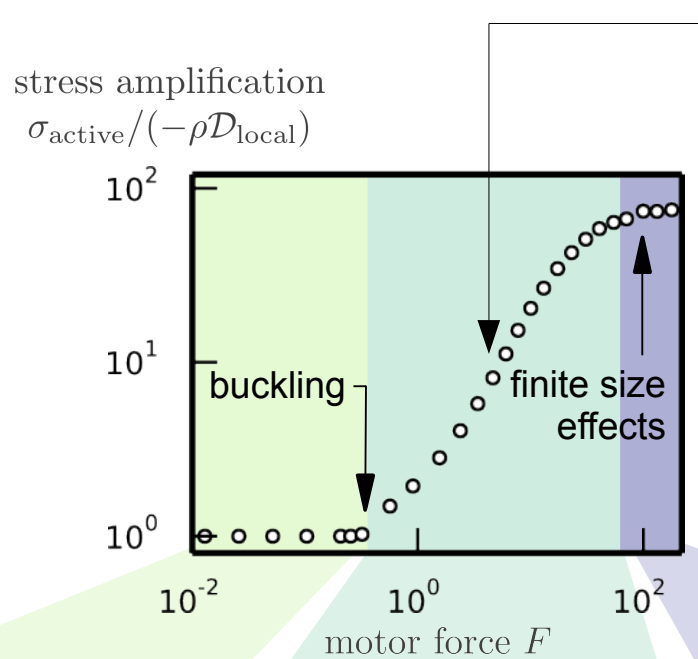


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Stress amplification is governed by an emergent buckling length scale



Large-scale buckling generates large stress amplification

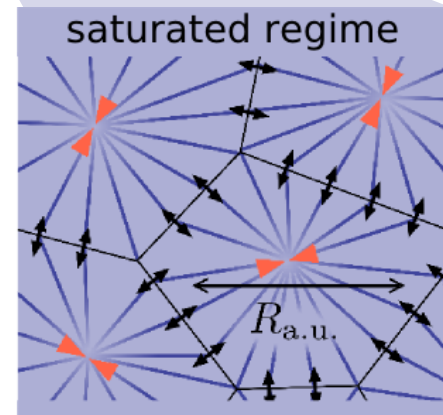
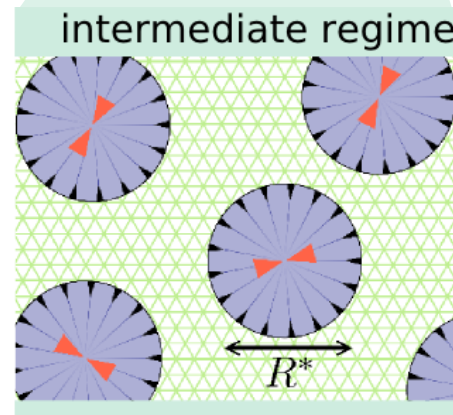
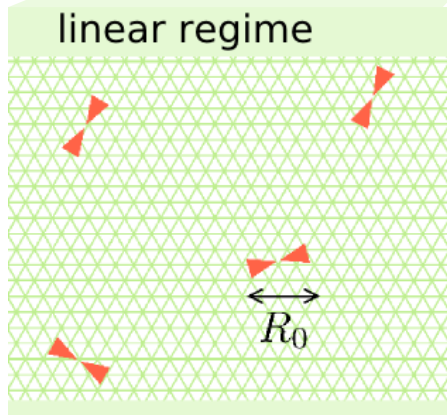


buckling zone controlled

$$\sigma_{\text{active}} = -\rho \mathcal{D}_{\text{eff}} \approx \rho F R^*$$

$$R^* \propto F^\alpha$$

	2D	3D
disorder regime		
stretching-dominated	$\alpha = 1$	$\alpha = 1/2$
bending-dominated	$\alpha \simeq 0.4$	$\alpha \simeq 0.25$



Experiments support all predicted amplification regimes

