

# Higher Derivative Supergravity and Moduli Stabilization

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## Motivation

- ⇨ higher derivative interactions exist in string theory and are expected generically in any effective QFT
- ⇨ in supersymmetric theories: correction to potential possible
- ⇨ in general: subleading correction to two-derivative action
- ⇨ unless: flat directions exist
  - ⇒ higher derivative corrections can be dominant contribution for  
**moduli stabilization and inflation**
- ⇨ for early work see: [Cecotti, Ferrara, Girardello]  
more recently: [Khoury, Koehn, Lehnert, Ovrut, ...]

## Example in $d = 4, N = 1$ global supersymmetry

$$\mathcal{L} = \mathcal{L}_{(0)} + \mathcal{L}_{(1)} , \quad \Phi^i = A^i + \theta\chi^i + \theta^2 F^i$$

$$\mathcal{L}_{(0)} = \int d^4\theta K(\Phi, \Phi^\dagger) + \int d^2\theta W(\Phi) + \text{h.c.}$$

$$\mathcal{L}_{(1)} = \frac{1}{16} \int d^4\theta \mathbf{T}_{ij\bar{k}\bar{l}}(\Phi, \Phi^\dagger) D^\alpha \Phi^i D_\alpha \Phi^j \bar{D}_{\dot{\alpha}} \Phi^{\dagger\bar{k}} \bar{D}^{\dot{\alpha}} \Phi^{\dagger\bar{l}}$$

$\mathcal{L}_{(1)}$ : special higher derivative operator in that

- induces corrections to potential
- does not introduce kinetic terms for  $F^i$

$\Rightarrow$  recent systematic analysis by [\[Ciupke\]](#)

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$$\begin{aligned} \mathcal{L}_{\text{bos}} = & -G_{i\bar{j}} \partial_\mu A^i \partial^\mu \bar{A}^{\bar{j}} + G_{i\bar{j}} F^i \bar{F}^{\bar{j}} + F^i W_i + \bar{F}^{\bar{i}} \bar{W}_{\bar{i}} + \\ & \mathbf{T}_{ij\bar{k}\bar{l}}(\mathbf{A}, \bar{\mathbf{A}}) [(\partial_\mu A^i \partial^\mu A^j)(\partial_\nu \bar{A}^{\bar{k}} \partial^\nu \bar{A}^{\bar{l}}) - 2F^i \bar{F}^{\bar{k}} (\partial_\mu A^j \partial^\mu \bar{A}^{\bar{l}}) + F^i F^j \bar{F}^{\bar{k}} \bar{F}^{\bar{l}}] \end{aligned}$$

Note:

- $G_{i\bar{j}} = \partial_i \partial_{\bar{j}} K$  standard Kähler metric
- $\mathcal{L}_{\text{bos}}$  quartic in  $F \Rightarrow$  e.o.m. cubic

$$G_{i\bar{k}} F^i + \bar{W}_{\bar{k}} + 2F^i (F^j \bar{F}^{\bar{l}} - \partial_\mu A^j \partial^\mu \bar{A}^{\bar{l}}) \mathbf{T}_{ij\bar{k}\bar{l}} = 0$$

## Cubic equation of motion for auxiliary $F$

$$G_{i\bar{k}} F^i + \bar{W}_{\bar{k}} + 2F^i (F^j \bar{F}^{\bar{l}} - \partial_\mu A^j \partial^\mu \bar{A}^{\bar{l}}) \mathbf{T}_{ij\bar{k}\bar{l}} = 0 \quad .$$

### Solutions:

- (always) one solution analytic in  $\mathbf{T}$
- all others: non-analytic in  $\mathbf{T}$  — diverge as  $\mathbf{T} \rightarrow 0$   
 $\Rightarrow$  discard as artefact of approximation/truncation of infinite series of higher-derivative terms

confirmed in exact solution of one-loop WZ-model

[Buchbinder, Kuzenko, Yarevskaya, Tyler]

### Aside:

kinetic terms for  $F$  can be discarded by same reasoning [Ciupke]

## Analytic solution

$$F^i = -G^{i\bar{l}} \bar{W}_{\bar{l}} + 2 \mathbf{T}^{\bar{k}l ij} \bar{W}_{\bar{k}} \bar{W}_{\bar{l}} W_j - 2 \mathbf{T}^{\bar{k} j \bar{i} l} (\partial_\mu A^j \partial^\mu \bar{A}^{\bar{l}}) \bar{W}_{\bar{k}} .$$

Inserted into  $\mathcal{L}_{\text{bos}}$  (at order  $\mathcal{O}(\mathbf{T})$ )

$$\begin{aligned} \mathcal{L}_{\text{bos}} = & - \left( G_{i\bar{k}} + 2 \mathbf{T}^{\bar{l} j \bar{i} k} W_j \bar{W}_{\bar{l}} \right) \partial_\mu A^i \partial^\mu \bar{A}^{\bar{k}} \\ & + \mathbf{T}_{ij\bar{k}\bar{l}} (\partial_\mu A^i \partial^\mu A^j) (\partial_\mu \bar{A}^{\bar{k}} \partial^\mu \bar{A}^{\bar{l}}) - V(A, \bar{A}) \end{aligned}$$

$$V = G^{i\bar{j}} W_i \bar{W}_{\bar{j}} - \mathbf{T}^{ij\bar{k}\bar{l}} W_i W_j \bar{W}_{\bar{k}} \bar{W}_{\bar{l}} .$$

### Note:

- no ghosts
- non-Kähler correction to metric & correction to  $V$
- supersymmetric minima:  $F^i = 0 = W_i \Rightarrow \langle A^i \rangle$  unchanged
- non-supersymmetric minima:  $F^i \neq 0 \Rightarrow \langle A^i \rangle$  can shift

## Generalization to supergravity

$$\begin{aligned} \mathcal{L}_{\text{bos}} = & -\frac{1}{2}\mathcal{R} - (G_{i\bar{k}} + 2e^K \mathbf{T}^{\bar{l}j}_{i\bar{k}} D_j W \bar{D}_{\bar{l}} \bar{W}) \partial_\mu A^i \partial^\mu \bar{A}^{\bar{k}} \\ & + \mathbf{T}_{ij\bar{k}\bar{l}} (\partial_\mu A^i \partial^\mu A^j) (\partial_\nu \bar{A}^{\bar{k}} \partial^\nu \bar{A}^{\bar{l}}) - V(A, \bar{A}) \end{aligned}$$

where

$$V = e^K (G^{i\bar{j}} D_i W \bar{D}_{\bar{j}} \bar{W} - 3|W|^2) - e^{2K} \mathbf{T}^{\bar{i}jkl} \bar{D}_{\bar{i}} \bar{W} \bar{D}_{\bar{j}} \bar{W} D_k W D_l W$$

no-scale example:

$$K = -p \ln(A + \bar{A}), \quad W = \text{const.}, \quad T < 0$$

minimum of  $V$  at

$$\langle A + \bar{A} \rangle = \left( -\frac{27}{2} p \mathbf{T} |W_0|^2 \right)^{1/3}$$

## Computation of $\mathbf{T}$ in CY orientifold compactifications of type IIB

### strategy:

compactify  $d = 10$ ,  $R^4$ -term and read off  $\mathbf{T}$  from  $\mathbf{T}(\partial A)^4$  coupling

### leading order result:

$$\mathbf{T}_{ijk\bar{l}} = \lambda(\alpha')^3 g_s^{-3/2} (\mathbf{\Pi}_m \mathbf{t}^m) K_{(0),i} K_{(0),j} K_{(0),\bar{k}} K_{(0),\bar{l}} ,$$

where

$$\mathbf{\Pi}_m \mathbf{t}^m = \int_{CY} c_2 \wedge J(t)$$

$\lambda =$  (uncomputed) combinatorial number,  $c_2 =$  2nd Chern class,

$\mathbf{t}^m =$  Kähler moduli of CY,  $\mathbf{\Pi}_m =$  topological numbers ,

$J(t) =$  Kähler form of CY



## Minimization of $V$

⇨ correction to  $V$  implied:

$$V = a(\chi) \mathcal{V}^{-3} |W_0|^2 + b(\lambda) \mathcal{V}^{-4} |W_0|^4 (\prod_m t^m),$$

$W_0$  = flux superpotential,  $\mathcal{V}$  = volume of CY,  
 $a, b$  numerical factors

⇨ minimization:

for  $\lambda < 0$ ,  $\chi > 0$ :

- $V$  has non-supersymmetric AdS minimum fixing all  $\langle t^m \rangle$  without non-perturbative  $W$ !
- $m_{\frac{3}{2}} \sim |W_0|^{-2}$ ,  $\mathcal{V} \sim |W_0|^3$   
 $\Rightarrow m_{\frac{3}{2}}$  small and  $\mathcal{V}$  large for large  $|W_0|$
- “orthogonal” to LVS as  $\chi > 0$

## Conclusion/Outlook

⇒ application to inflation

[Broy,Ciupke,Pedro,Westphal; Cicoli,Muia,Pedro,...]

⇒ application to moduli stabilization in LVS

[Cicoli,Ciupke,de Alwis,Muia, to appear]

⇒ similar analysis for DBI-action

[Bielleman,Ibanez,Pedro,Valenzuela,Wieck]

**!!Happy Birthday Fernando!!**

and thank you for your friendship and  
your genuine interest in my work  
throughout the past 30 years