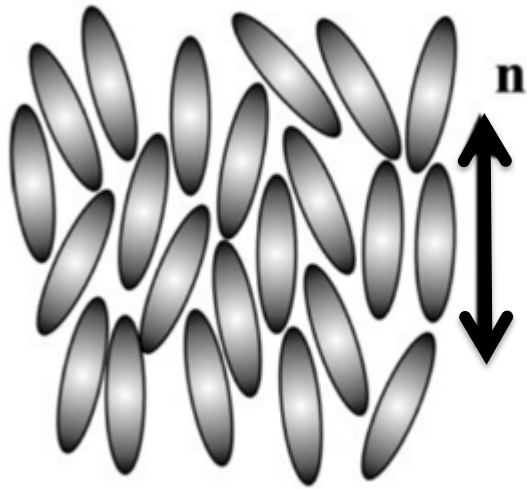


## **Orientational Phase Transitions and the Assembly of Viral Capsids**

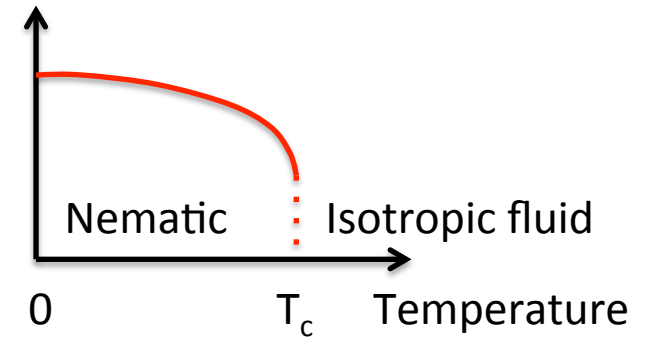
- i) Lorman-Rochal theory and the Landau theory of orientational ordering in liquid crystals and glasses.
  
- ii) Extensions of Lorman-Rochal theory.

# Orientational ordering in liquid crystals and glasses

- Nematic liquid crystals.



Order Parameter

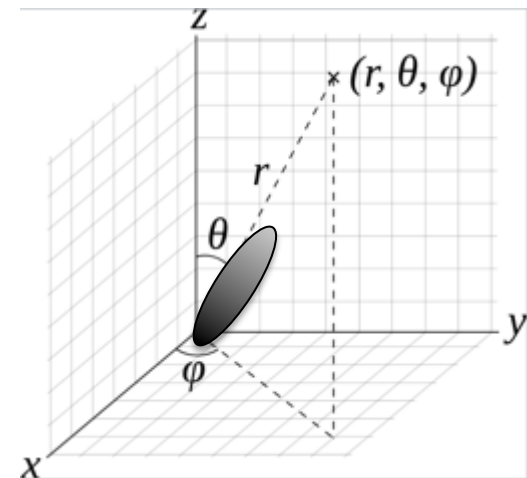


*First-order phase transition*

Order parameter

$$S = \langle P_2^0(\cos\theta) \rangle = \left\langle \frac{3 \cos^2 \theta - 1}{2} \right\rangle$$

Second Legendre polynomial  $P_2^0(x) = (3x^2 - 1)/2$  ?



- Biaxial Nematics (Freiser, 1970)

Two order parameters: S and T

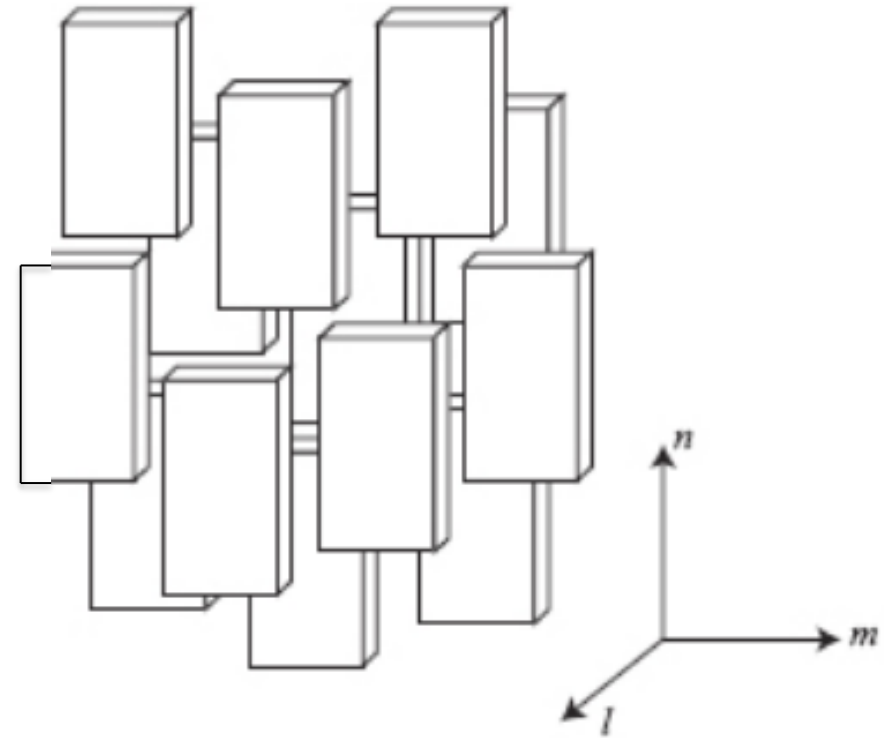
$S = T = 0$  isotropic



$S \neq 0, T = 0$  uniaxial nematic



$S \neq 0, T \neq 0$  biaxial nematic

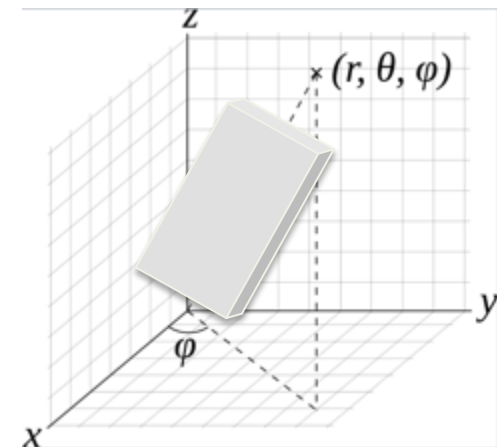


Order parameters

$$S = \langle P_2^0(\cos\theta) \rangle$$

$$T = \langle P_2^2(\cos\theta)\exp(i\varphi) \rangle = \langle P_2^{-2}(\cos\theta)\exp(-i\varphi) \rangle$$

$P_2^2(x) = 3(1-x^2)$  Associated Legendre polynomials ??



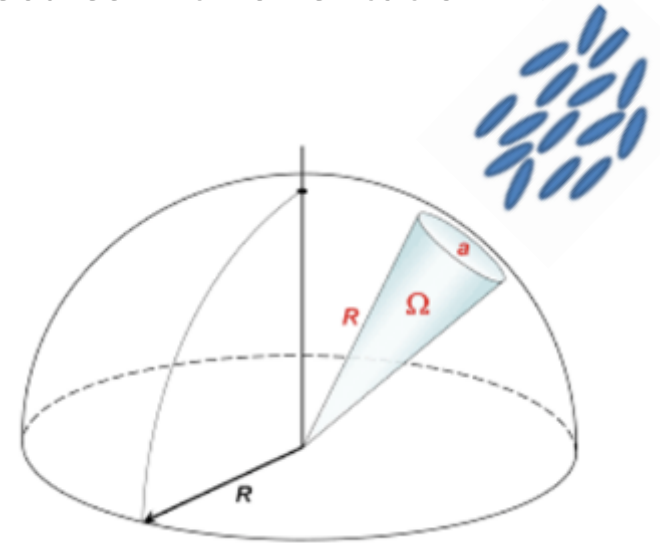
- General case: measure the density  $\rho(\Omega)$  of molecules with orientation  $\Omega$ .
- “Spherical Fourier decomposition”.

$$\rho(\Omega) = \sum_{L=0}^{\infty} \sum_{M=-L}^{+L} Q_{L,M} Y_{L,M}(\Omega)$$

- Spherical Harmonics

$$Y_{L,M}(\theta, \varphi) \propto P_L^M(\cos\theta) \exp(iM\varphi)$$

- Expansion coefficients  $Q_{L,M}$  *Orientational order parameter set*



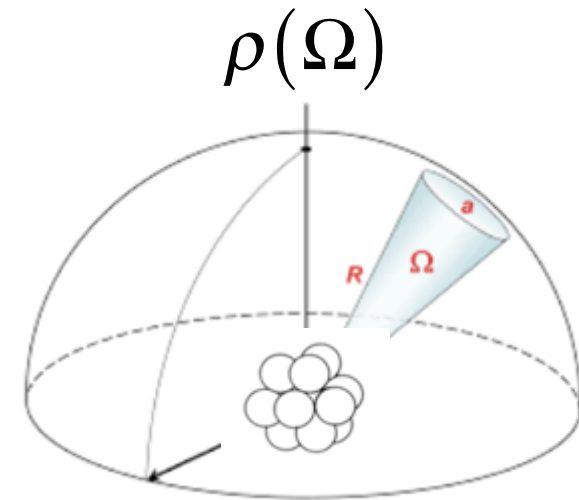
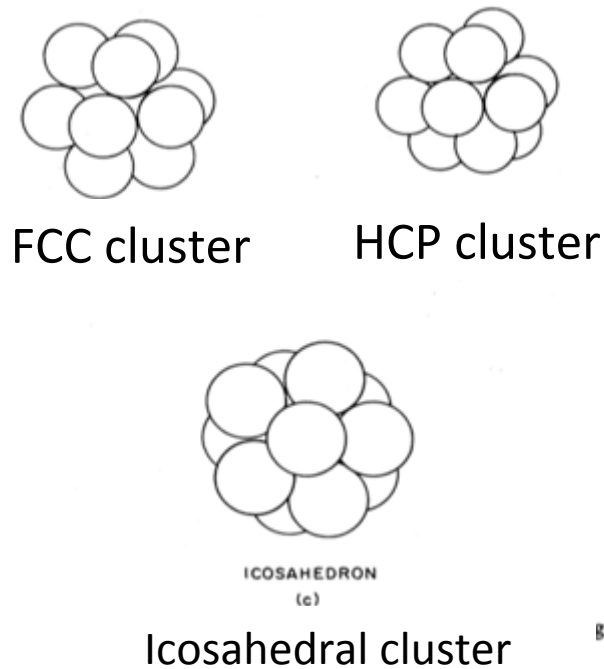
### Infinite hierarchy of Orientational Transitions

- Biaxial nematics:  $L = 2, M = 0, \pm 2$
- Cholesterics, blue phases:  $L = 2, M = 0, \pm 1, \pm 2$       Hornreich & Shtrikman (1980, 1981).
- “Cubatic” liquid crystal:  $L = 4$       Nelson & Toner (1981)

- Icosahedral Glasses/Quasicrystals:  $L=6$

(Steinhardt, Nelson, Ronchetti 1983)

- $\rho(\Omega)$ : surface density modulation of **cluster** of atoms in liquids near melting.  
(no positional order)



- What is the best choice for the symmetry of the clusters?

- **Rules of Landau Theory**



Lev  
Landau

Rule # 1: Free energy = Sum of the **scalar invariants** of the order parameters  $Q_{L,m}$  (under group SO(3) of rotations).

Rule # 2: Use one "irreducible representation". **Pick one single L.** (but which L?)

Rule # 3: **Minimize** free energy with respect to  $Q_{L,M}$ .

$$F = r \sum_{M=-L}^{+L} |Q_{LM}|^2 + w \sum_{\substack{-L \leq M_{1,2,3} \leq L \\ M_1 + M_2 + M_3 = 0}} \left( \begin{matrix} L & L & L \\ M_1 & M_2 & M_3 \end{matrix} \right) Q_{LM_1} Q_{LM_2} Q_{LM_3} + u \{Q^4\} + \dots$$

quadratic invariant
cubic invariant
quartic invariant

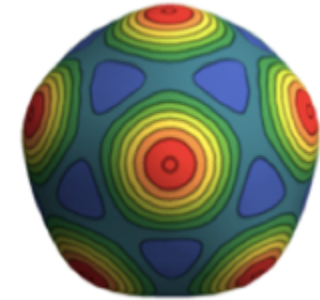
- $r, w, u$ : Phenomenological parameters.
- Simpler: expand free energy in powers of orientational density  $\rho(\Omega)$

$$F(\{\rho(\vec{r})\}) = \int_{Surface} dS \left\{ r \rho(\vec{r})^2 + w \rho(\vec{r})^3 + u \rho(\vec{r})^4 + \dots \right\}$$

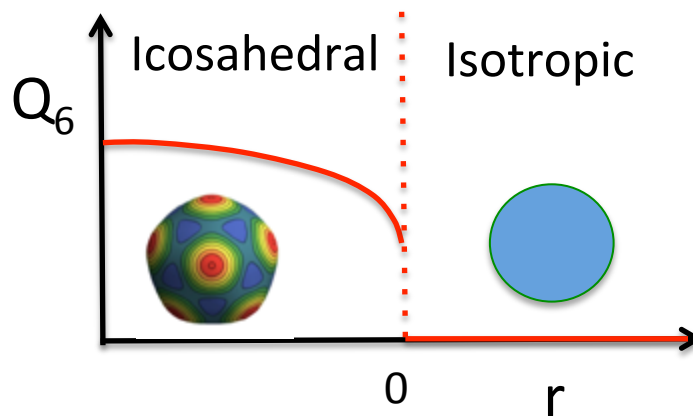
Insert 
$$\rho(\Omega) = \sum_{M=-L}^{+L} Q_{L,M} Y_{L,M}(\Omega)$$

- Free energy minimum for L=6 has icosahedral symmetry

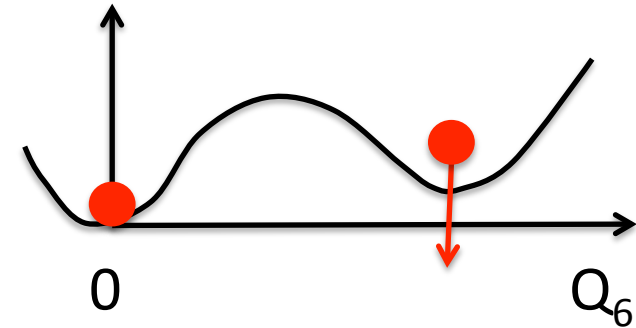
$$\rho(\Omega) = Q_6 \left( Y_{0,0} + \sqrt{\frac{7}{11}} Y_{6,5} - \sqrt{\frac{7}{11}} Y_{6,-5} \right)$$



- L = 6 "icosahedral spherical harmonic"
- $Q_6$  = Order parameter "amplitude"



$$F_6 = rQ_6^2 + wQ_6^3 + uQ_6^4$$



*First-order* transition (cubic term non-zero)

- Icosahedral state is thermodynamically stable against fluctuations.

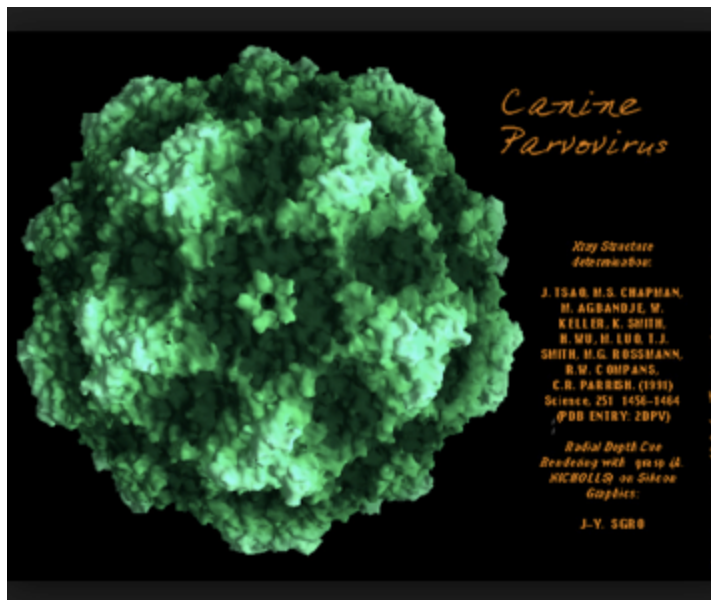
## Landau theory of crystallization and the capsid structures of small icosahedral viruses

V. L. Lorman<sup>1</sup> and S. B. Rochal<sup>1,2</sup>

<sup>1</sup>Laboratoire de Physique Theorique et Astroparticules, CNRS, Universite Montpellier 2, Place Eugene Bataillon, 34095 Montpellier, France

<sup>2</sup>Physical Department, South Federal University, 5 Zorge Street, 344090 Rostov-on-Don, Russia

### Capsid Canine Parvovirus



$$\rho_{15}(\Omega) \propto \sum_{M=-15}^{+15} Q_{15,M} Y_{15,M}(\Omega)$$



**L=15** Icosahedral spherical harmonic



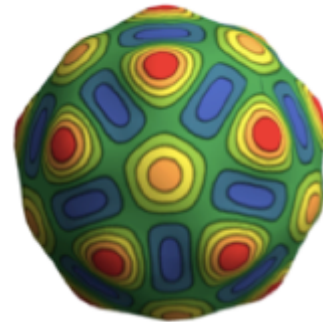
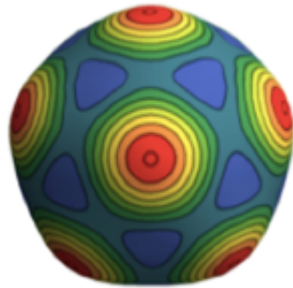
For **certain**  $L$  there is a **unique** linear combination  $Y_h(L)$  of the  $Y_{L,M}$  that transforms as a scalar **or** as a pseudo-scalar under the icosahedral symmetry group.

- **Even  $L$**        $L = 6, 10, 12, 16, \dots$

$L=6$

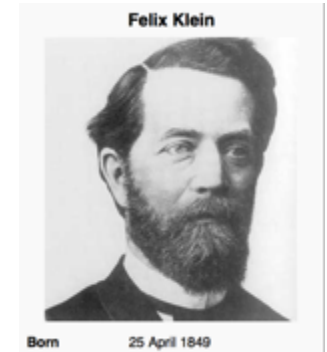
$L=10$

$Y_h(L)$



scalar density

**even** under inversion  $r \rightarrow -r$



- **Odd  $L$**        $L = 15, 21, 25, 27, \dots$

$Y_h(15)$



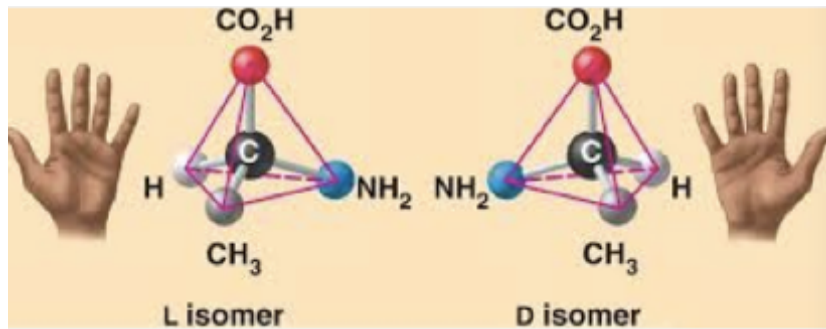
pseudo-scalar density

**odd** under inversion

- *Chiral pairs* (isomers)

Rochal & Lorman: *capsid density cannot be presented by the even L scalars*

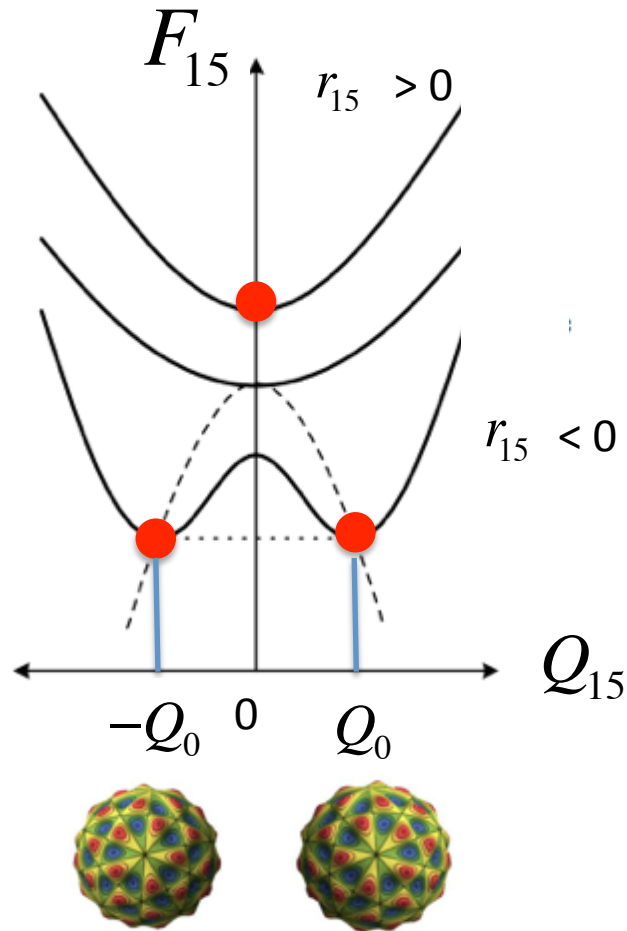
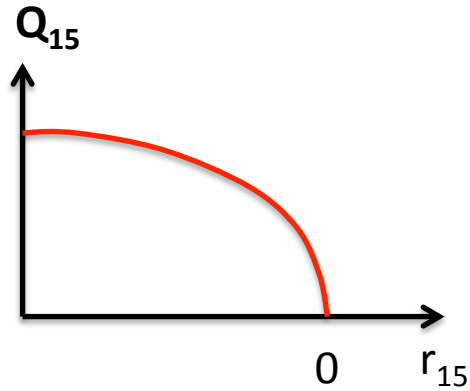
amino-acids



- Capsid densities are not even under inversion: *only odd L spherical harmonics*.
- Smallest viral capsids should correspond to  $Y_h$  (L=15): T=1 viruses.

Parvovirus (T=1) :  $\rho(\Omega) = Q_{15} Y_h(15)$

Landau free energy (L=15):  $F_{15} = r_{15}Q_{15}^2 + uQ_{15}^4$



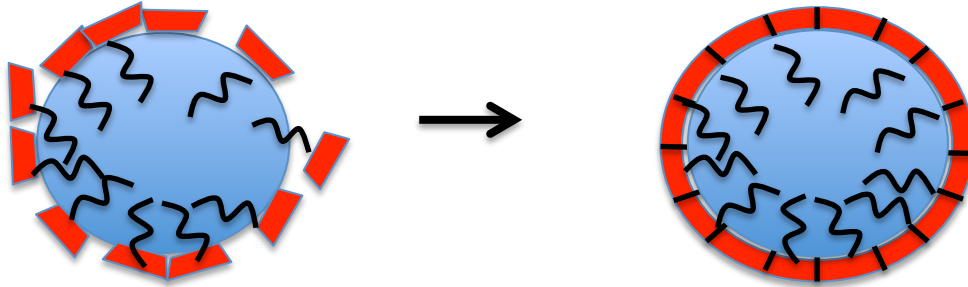
Cubic term zero for any odd L !

“Spontaneous chiral symmetry breaking transition”.

$$Q_0 = \sqrt{\frac{-r_{15}}{2u}}$$

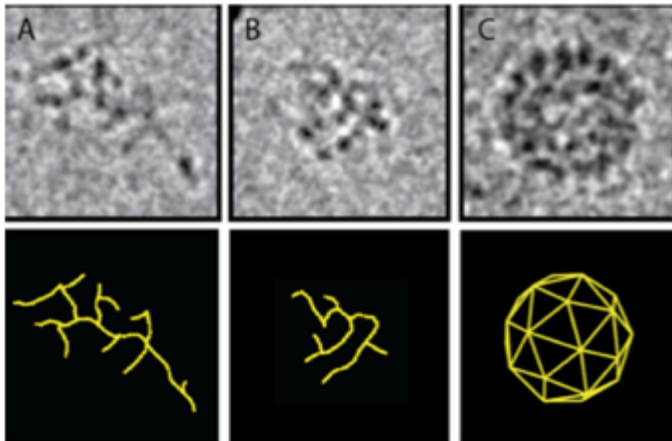
## Questions

1) Second-order transition?



*Collective assembly pathway*

### Cowpea Chlorotic Mottle Virus (CCMV)



R. Cadena-Nava, M. Comas-Garcia,  
R. Garmann, A. Rao C. M. Knobler,  
and W. M. Gelbart  
J. Virol. 86, 3318 (2012).

R. Garmann, M. Comas-Garcia, A. Gopal,  
C. M. Knobler, and W. M. Gelbart  
J. Mol. Biol. 426, 1050 (2013)

- Fluorescence thermal shift assay CCMV: **first-order transition**

Guillaume Tresset, Jingzhi Chen, Maelenn Chevreuil, Naïma Nhiri, Eric Jacquet, and Yves Lansac  
Phys. Rev. Applied 7, 014005, 2017

2) Spontaneous chiral symmetry breaking transition at  $r_{15} = 0$ .  
*Uniform state is already chiral.*

3) Thermodynamic stability against fluctuations: expand free energy around  $Y_h(15)$  solution

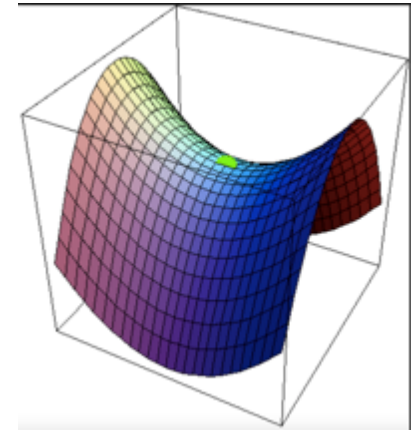
$$\rho(\Omega) = \text{const.} + \rho_{15} Y_h(15) + \sum_{M=-15}^{+15} \delta Q_{15,M} Y_{15,M}(\Omega)$$

fluctuation

Fluctuation free energy: 
$$\delta F_{15} = \frac{1}{2} \sum_{M=-15}^{15} \sum_{M'=-15}^{15} C_{M,M'} \delta Q_{15,M} \delta Q_{15,M'}$$

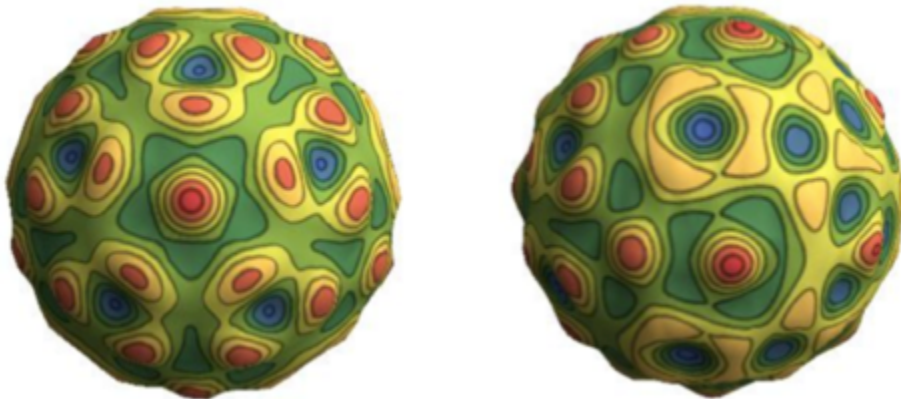
- $C_{M,M'}$  : 31 x 31 stability matrix should have **positive** eigenvalues.  
:

21 positive eigenvalues  
3 zero eigenvalues (rotations)  
7 negative eigenvalues !



saddle-point

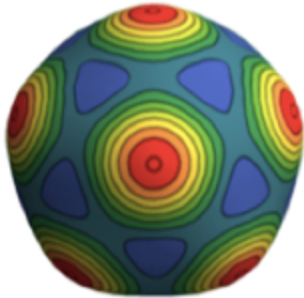
- Lowest free energy state in the  $L = 15$  sector:



- Symmetry group  $D_5$  :  
One five-fold symmetry axis.  
Five “odd” two-fold axes.

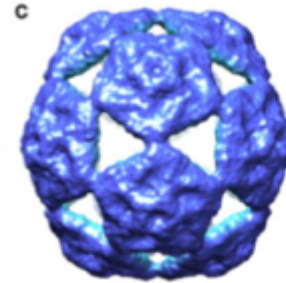
- *All* odd icosahedral spherical harmonic states are unstable.
- *Only*  $L = 6, 10, 12,$  and  $18$  icosahedral states are stable. ( $L=16$  is unstable, Matthews).

Should we perhaps take another look at the even L states?

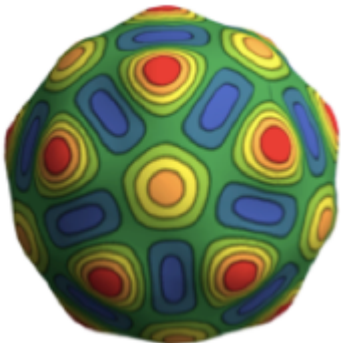


L=6

Picornavirus:  
assembles from 12 identical  
chiral pentagons composed of 15  
proteins: **capsomers**

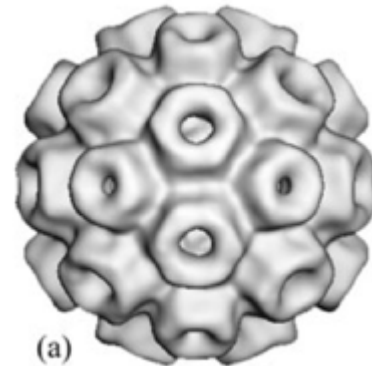


- Interpret  $\rho(\Omega)$  as the density of **capsomer centers**. Add chiral capsomers.



L=10

CCMV:  
assembles from 32 chiral capsomers  
12 pentamers + 20 hexamers.



- But .... no signature of chirality in the density of capsomer centers??

## Extensions of Lorman-Rochal

- **Step 1:** Landau-*Brazovskii* theory
  - smectic-nematic transition
  - chiral liquid crystals (cholesterics)
  - weak solidification (block co-polymers).

$$F_{LB} = \int dS \left\{ \left| (\nabla^2 + k_0^2) \rho(\vec{r}) \right|^2 + r \rho(\vec{r})^2 + w \rho(\vec{r})^3 + u \rho(\vec{r})^4 \right\}$$

- Wavenumber  $k_0 = 2\pi/a$ . Dominant molecular length-scale:  $a =$  protein size.

- First term minimized by density waves  $\rho(\vec{r}) \propto \exp i(k_0 \hat{n} \cdot \vec{r})$

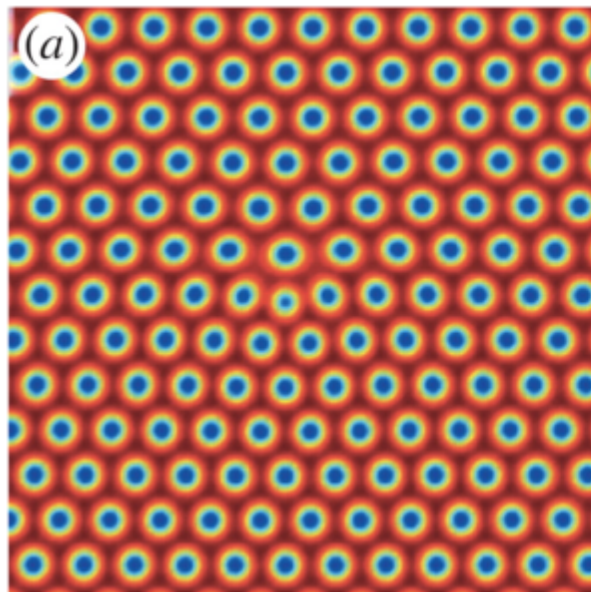
$$(\nabla^2 + k_0^2) \exp i(k_0 \hat{n} \cdot \vec{r}) = 0$$

$\hat{n} =$  unit vector

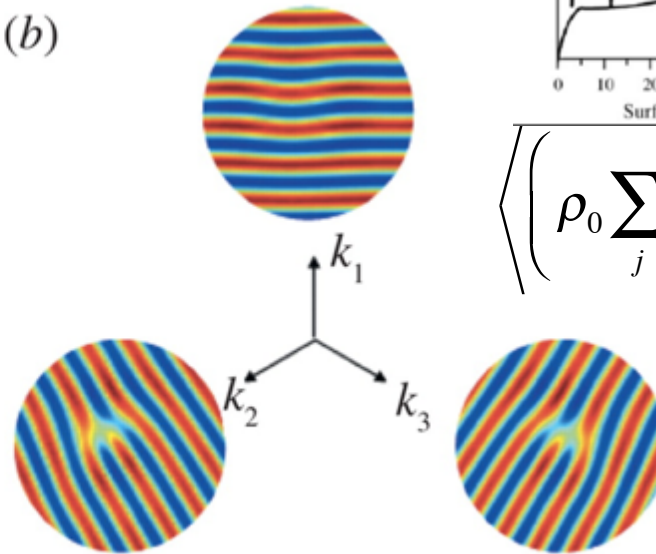


- Look for *combinations* of density waves

$$\rho(\vec{r}) = \rho_0 \sum_{j=1}^3 \cos(k_0 \hat{n}_j \cdot \vec{r})$$



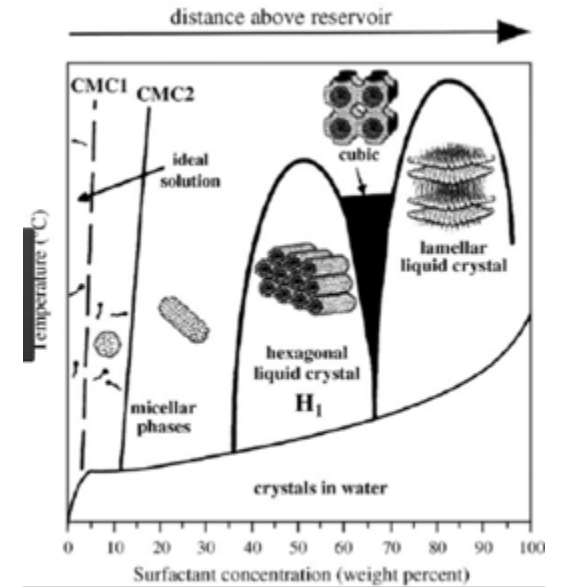
(b)



$$\vec{k}_j = k_0 \hat{n}_j$$

$$\left\langle \left( \rho_0 \sum_j \cos(\vec{k}_j \cdot \vec{r}) \right)^3 \right\rangle \neq 0$$

Pezzutti et al.



- First-order transition: cubic term in Landau energy non-zero
- Block co-polymers (L. Leibler): competition between hexagonal phases, lamellar phases and other symmetries.

## Landau-Brazovskii on a spherical surface.

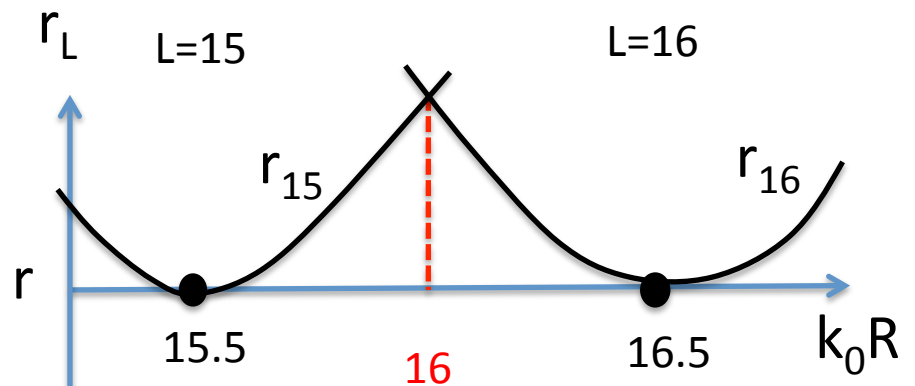
- Insert  $\rho(\Omega) = \sum_{L=1}^{\infty} \sum_{M=-L}^{+L} Q_{L,M} Y_{L,M}(\Omega)$  in Landau-Brazovskii free energy.

$$F_{LB} = \sum_L r_L \sum_{M=-L}^{+L} |Q_{L,M}|^2 + w \sum_{\left\{ \begin{smallmatrix} L_1 L_2 L_3 \\ M_1 M_2 M_3 \\ M_1 + M_2 + M_3 = 0 \end{smallmatrix} \right\}} \begin{pmatrix} L_1 & L_2 & L_3 \\ M_1 & M_2 & M_3 \end{pmatrix} Q_{L_1, M_1} Q_{L_2, M_2} Q_{L_3, M_3} + u\{Q^4\}$$

- Sum over all L.

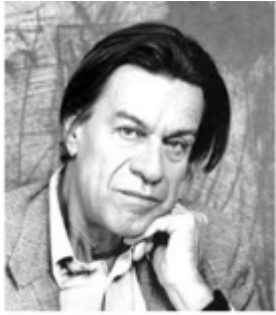
- $r_L = r + \left[ L(L+1) - (k_0 R)^2 \right]^2$

Predict the L value of a capsid !



$r_L$  has minimum when  $k_0 R \approx L$

## Step 2: Include chirality directly in free energy



P-G De Gennes

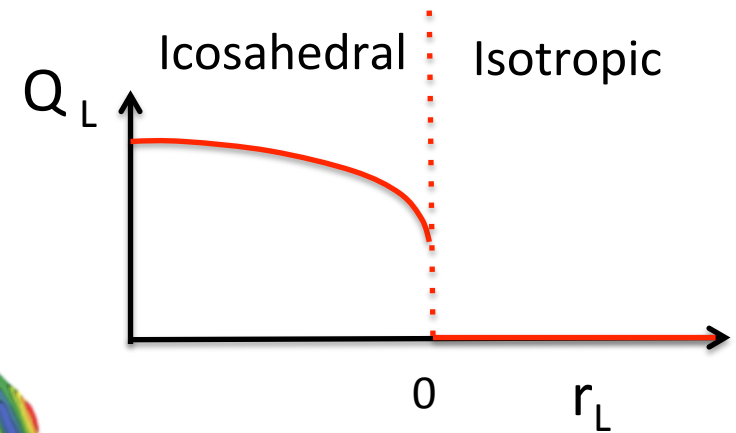
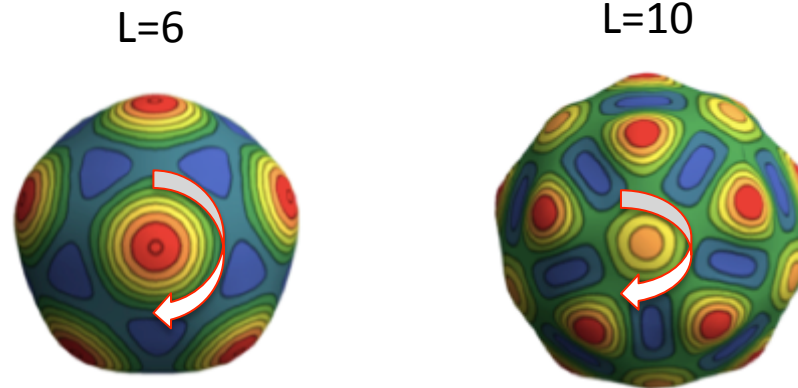
Chiral liquid crystals: add **lowest-order chiral pseudoscalar** to the Landau free energy of the achiral liquid crystal.

$$F = \int_{S^2} \left[ \left( \frac{1}{2} \left( (\Delta + k_o^2) \rho \right)^2 + \frac{r}{2} \rho^2 + \frac{u}{3} \rho^3 + \frac{v}{4} \rho^4 + \dots \right) + \chi \left( \nabla \rho \cdot (\nabla \nabla \rho)^2 \cdot (\mathbf{n} \times \nabla \rho) \right) \right] dS.$$

scalars

- lowest non-zero pseudo-scalar
  - variant of ``Helfrich-Prost''
  - Sanjay Dharmavaram
- 
- Now let's try again

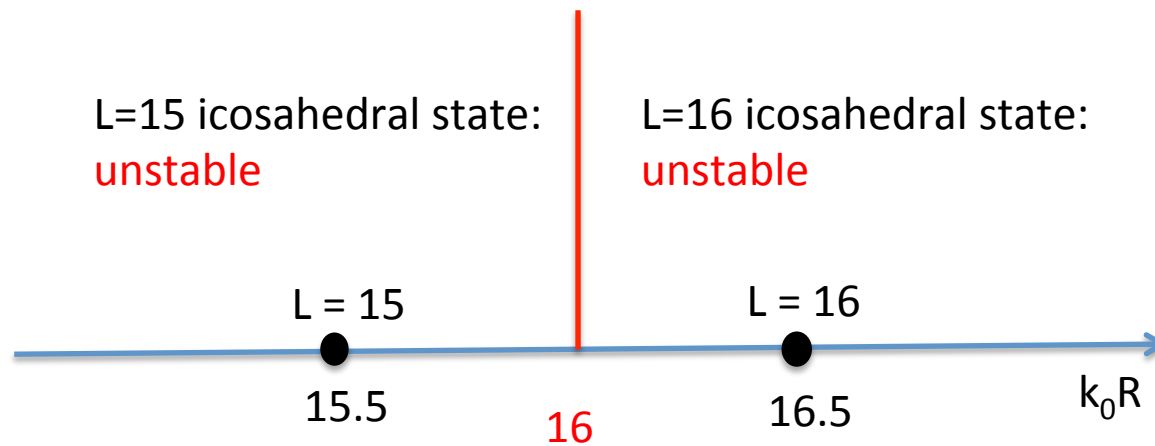
- Stable minima for  $L=6, 10, 12,$  and  $16$ .



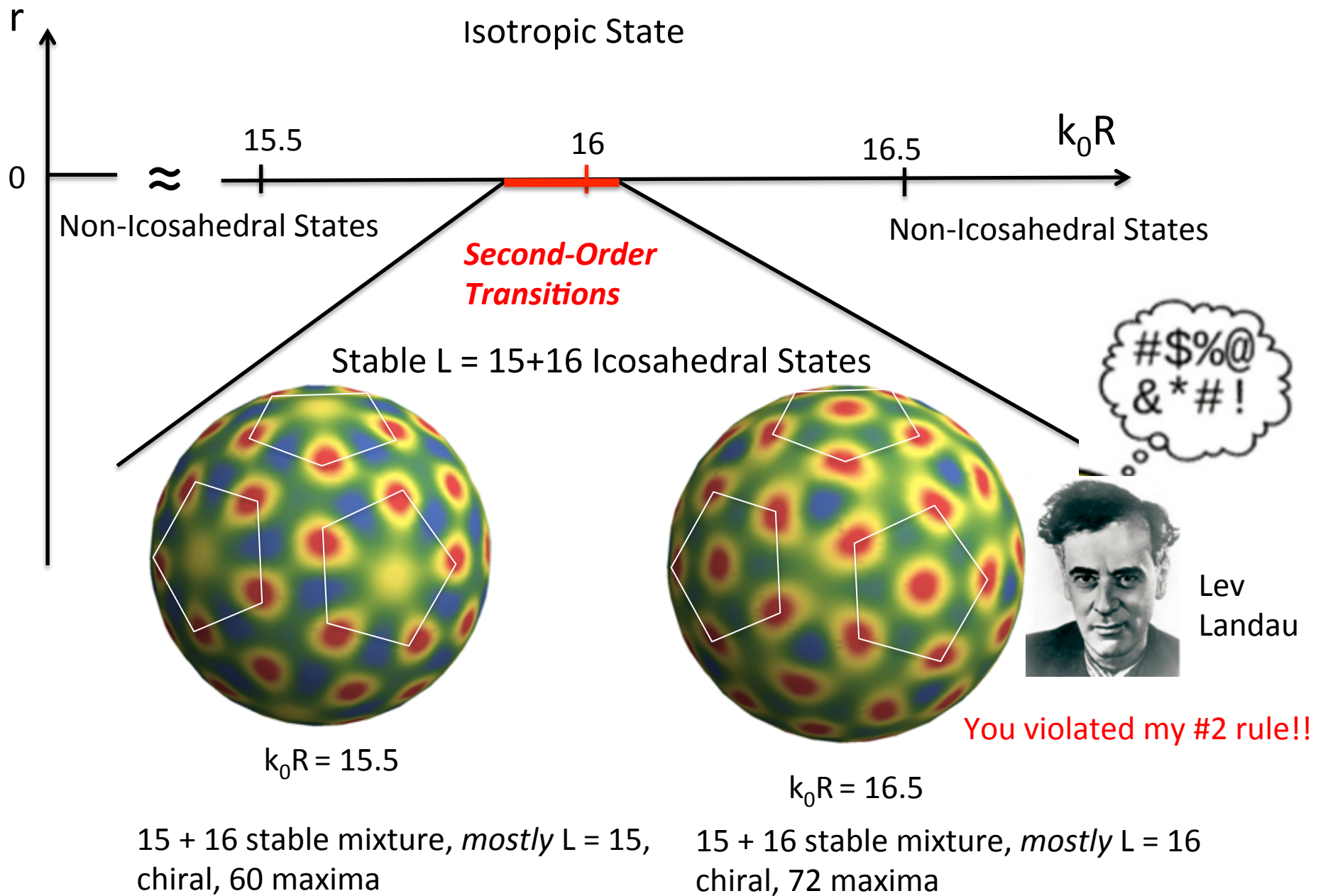
- *Thermal fluctuations* around the isotropic and  $Y_h(L)$  states **are** chiral.
- *No spontaneous chiral symmetry breaking* at  $r_L = 0$ .

➔ Conclusion:  $L = 6, 10, 12,$  and  $16$  icosahedral spherical harmonics are possible representations for collective capsid assembly. ←

- But .... odd  $L$  icosahedral spherical harmonics states remain unstable!

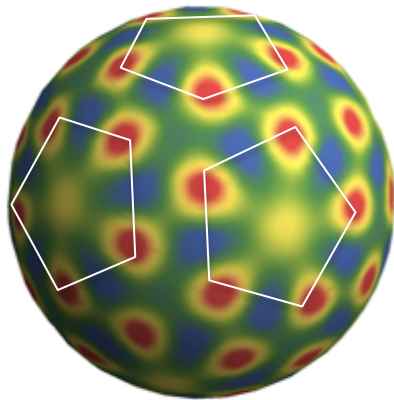


- $r_{15} = r_{16}$  when  $k_0R = 16$
- Could there be a stable icosahedral order parameter composed of **two** one-dimensional “irreducible representations” of  $SO(3)$  near  $k_0R = 16$  ?
- $L = 15 + 16$  space has  $31+33 =$  **64 dimensions**

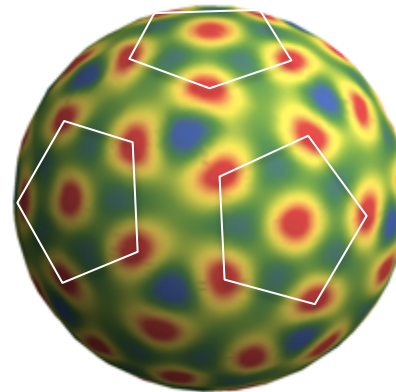


Dharmavaram, S., Xie, F., Klug, W., Rudnick, J., & Bruinsma, R. (2017). *Orientalional phase transitions and the assembly of viral capsids*. Physical Review E, 95(6), 062402.

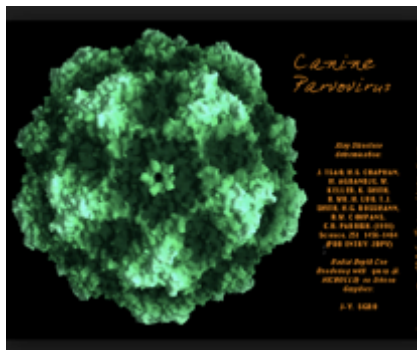
# Examples



$$k_0R = 15.5$$

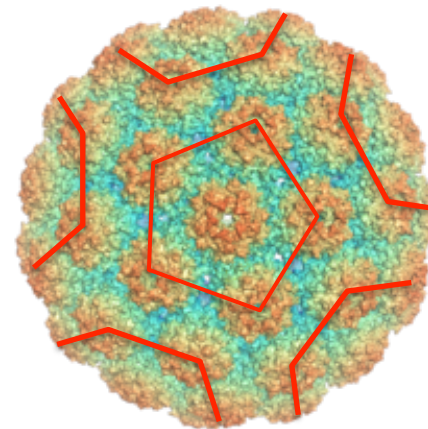


$$k_0R = 16.5$$



Parvovirus

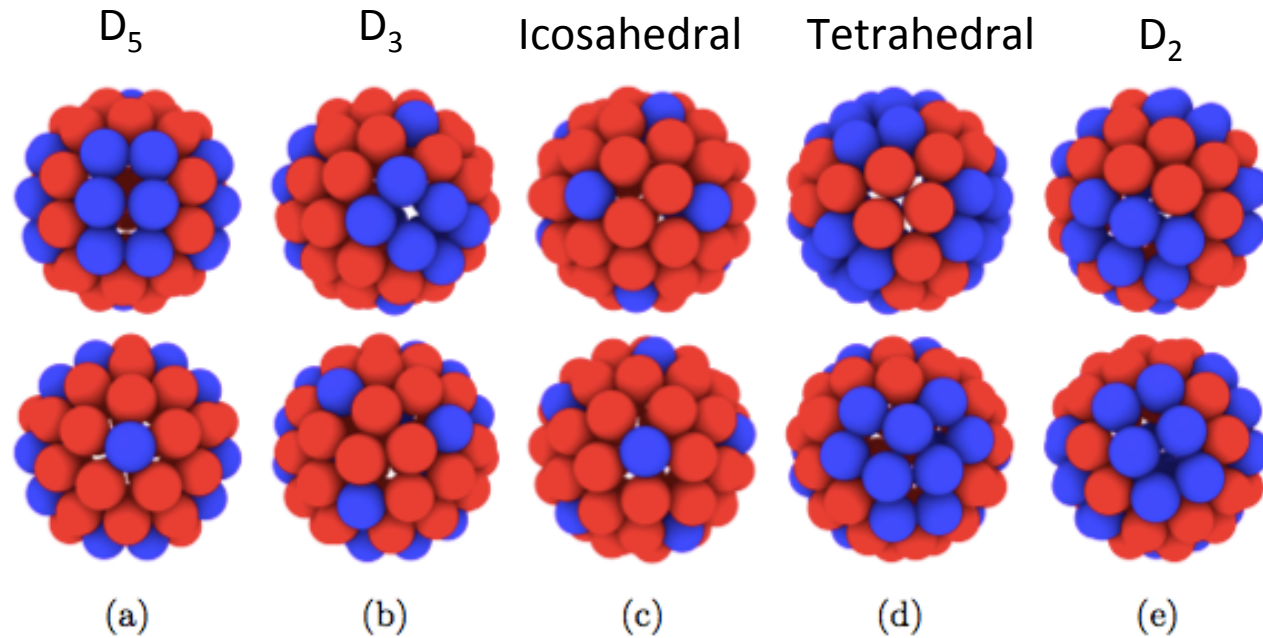
Maxima: *single capsid proteins*



Polyomavirus

72 maxima: *capsid protein pentamers*

## Numerical simulations: 72 particles on a spherical surface



- Icosahedral state: *very fragile*. Competes with other states that have  $D_3$ ,  $D_5$ , T symmetry.

S. Paquay, H. Kusumaatmaja, D. Wales, R. Zandi, and P. van der Schoot, Energetically favoured defects in dense packings of particles on spherical surfaces <http://arxiv.org/abs/1602.07945>

**WHY WAS ICOSAHEDRAL SYMMETRY SELECTED FOR VIRUSES?**



## Conclusion

Lorman-Rochal theory (+ some extensions) is a powerful mathematical tool for the study of capsid assembly.

- **Two modes** of viral assembly :

Assembly of mixed L, *intrinsically chiral* shells via (quasi) continuous transitions.

Assembly of pure L, *effectively achiral* shells via strongly first-order transitions.

- Viral assembly: laboratory for testing and (maybe) violating Landau theory!
- Substantial experimental problems!



Kevin Zhang

Luigi Perotti



Joseph Rudnick

Emmy

Amit Singh

Sanjay Dharmavaram

Josh Kelly

William Klug



with thanks to:

Bill Gelbart

Rees Garmann

Chuck Knobler

