Non-equilibrium thermodynamics of quantum processes: two

or an invitation to study stochastic thermodynamics of quantum processes

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Content & structure

 $ho_0^{
m eq}$

 ρ_{τ}^{B}

n

 Bath

 \mathcal{H}^F_t

Non-equilibrium definition of thermodynamic work: fluctuation theorems



Irreversibility & entropy production in closed q-systems

Quantum correlations, coherences and thermodynamics



Plan of the discussion



Landauer principle & quantum (open-system) dynamics

0) Non-equilibrium interpretation of non-unitality



Finite-time

 Non-equilibrium version of Landauer principle: A tighter bound than the original one?



Talkner, Lutz, and Haenggí, Phys. Rev. E 75, 050102 (2007) Goold, Poschinger, and Modí, Phys. Rev. E 90, 020101 (2014)



Link to fluctuation theorems

Statements linking the stochastic properties of thermodynamically relevant quantities to equilibrium features

Jarzínsky equalíty $\langle e^{-\beta W} \rangle = e^{-\beta \Delta F}$ Jarzynski, PRL 78 2690 (1997) free-energy change



Finite-time

Standard fluctuation relations hold unchanged for unital open channels

A. E. Rastegin, J. Stat. Mech. P06016



Mr. Landauer's principle

The principle

"Any logically irreversible manipulation of information... must be accompanied by a corresponding entropy increase in non-information bearing degrees of freedom of the information processing apparatus or its environment"



 $\beta \langle Q \rangle \geq \Delta S$

Rolf Landauer

Can we let Landauer's principle emerge from microscopic equations?



(Im) proving Landauer: Reeb & Wolf



1st result $\beta \Delta Q = \Delta S + I(S':R') + D(\rho'_R || \rho_R) \ge \Delta S$ 2nd result $\beta \Delta Q \ge \Delta S + \frac{2(\Delta S)^2}{\log^2(d-1)+4}$

D. Reeb, and M. Wolf, NJP 16, 103011 (2014)



 $\hat{H} = \hat{H}_{\mathcal{S}} + \hat{H}_{\mathcal{E}} + \hat{H}_{\mathcal{SE}}$

 $\hat{\rho}_{\mathcal{SE}} = \hat{\rho}_{\mathcal{S}} \otimes e^{-\beta \hat{H}_{\mathcal{E}}} / Z_{\mathcal{E}}$

 $\hat{\rho}_{\mathcal{E}}' = \operatorname{tr}_{\mathcal{S}}[\hat{U}(\hat{\rho}_{\mathcal{S}} \otimes \hat{\rho}_{\mathcal{E}})\hat{U}^{\dagger}] = \sum \hat{A}_l \hat{\rho}_{\mathcal{E}} \hat{A}_l^{\dagger}$

 $\sum_{l} \hat{A}_{l}^{\dagger} \hat{A}_{l} = 1_{\mathcal{E}}$ $\hat{A}_{l=jk} = \sqrt{\lambda_j} \langle s_k | \hat{U} | s_j \rangle$

Goold, Paternostro, and Modí, PRL 114, 060602 (2015)



How about heat?



 $P(\mathbf{Q}) = \sum p_m p_{n|m} \delta(\mathbf{Q} - (E_n - E_m))$ mn

Heat probability distribution!

Talkner, Lutz, and Haenggi, Phys. Rev. E 75, 050102 (2007)



 $P(Q) = \sum \langle r_n | \hat{A}_l | r_m \rangle (\hat{\rho}_{\mathcal{E}})_{mm} \langle r_m | \hat{A}_l^{\dagger} | r_n \rangle \delta(Q - E_{nm})$ l.m.n

$$\langle e^{-\beta Q} \rangle = \int e^{-\beta Q} dQ P(Q) = \sum_{l} \operatorname{tr}[\hat{A}_{l}^{\dagger} \hat{\rho}_{\mathcal{E}} \hat{A}_{l}] = \operatorname{tr}[\hat{A} \hat{\rho}_{\mathcal{E}}]$$
$$\hat{\mathbf{A}} = \sum_{l} \hat{A}_{l} \hat{A}_{l}^{\dagger}$$

Thermodynamic interpretation of the 'degree' of non-unitality of a process

Goold, Paternostro, and Modí, PRL 114, 060602 (2015)



 $P(Q) = \sum \langle r_n | \hat{A}_l | r_m \rangle (\hat{\rho}_{\mathcal{E}})_{mm} \langle r_m | \hat{A}_l^{\dagger} | r_n \rangle \delta(Q - E_{nm})$ l.m.n

$$\begin{split} \langle e^{-\beta Q} \rangle &= \int e^{-\beta Q} dQ P(Q) = \sum_{l} \operatorname{tr}[\hat{A}_{l}^{\dagger} \hat{\rho}_{\mathcal{E}} \hat{A}_{l}] = \operatorname{tr}[\hat{A} \hat{\rho}_{\mathcal{E}}] \\ \hat{\mathbf{A}} &= \sum_{l} \hat{A}_{l} \hat{A}_{l}^{\dagger} \\ \text{Jarzinsky equality } \langle e^{-\beta W} \rangle = e^{-\beta \Delta F} \end{split}$$

Not a proper fluctuation theorem Talkner, Campisi, Haenggi, J Stat Mech P02025 (2009) Goold, Paternostro, and Modi, PRL 114, 060602 (2015)



$$P(Q) = \sum_{l,m,n} \langle r_n | \hat{A}_l | r_m \rangle (\hat{\rho}_{\mathcal{E}})_{mm} \langle r_m | \hat{A}_l^{\dagger} | r_n \rangle \delta(Q - E_{nm})$$

$$\langle e^{-\beta Q} \rangle = \int e^{-\beta Q} dQ P(Q) = \sum_{l} \operatorname{tr}[\hat{A}_{l}^{\dagger} \hat{\rho}_{\mathcal{E}} \hat{A}_{l}] = \operatorname{tr}[\hat{\mathbf{A}} \hat{\rho}_{\mathcal{E}}]$$

$$= \operatorname{tr}[\hat{\rho}_{\mathcal{S}} \otimes \mathbb{1}_{\mathcal{E}} \hat{U}^{\dagger} \mathbb{1}_{\mathcal{S}} \otimes \hat{\rho}_{\mathcal{E}} \hat{U}] = \operatorname{tr}[\hat{\mathbf{M}} \hat{\rho}_{\mathcal{S}}]$$

$$\text{Assume unitality} \quad \langle Q \rangle \geq 0$$

$$\beta \langle Q \rangle \geq \mathcal{B}_{Q} = -\ln(\operatorname{tr}[\hat{\mathbf{A}} \hat{\rho}_{\mathcal{E}}]) = -\ln(\operatorname{tr}[\hat{\mathbf{M}} \hat{\rho}_{\mathcal{S}}])$$

$$\text{A tighter bound than Reeb and Wolf's ? Yes!}$$

$$\text{Goold, Paternostro, and Modi, PRL 114, 060602 (2015) }$$