

A machine learning perspective on the many-body problem in classical and quantum physics

Juan Felipe Carrasquilla Alvarez

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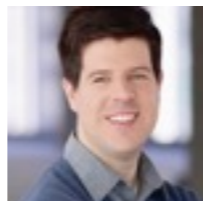
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Alireza Makhzani



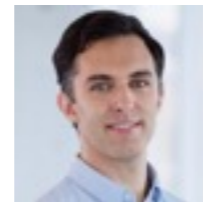
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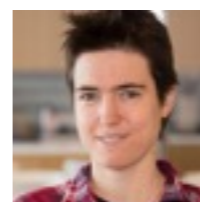
Daniel Roy



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Graham Taylor



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PHYSICS/MACHINE LEARNING FRIENDS



Roger Melko (U. Waterloo and Perimeter)



Giacomo Torlai (U. Waterloo and Perimeter)



Peter Broecker (U. Cologne)



Simon Trebst (U. Cologne)



Ehsan Khatami (San Jose State U)



Kelvin Chng (San Jose State U)



Giuseppe Carleo (ETH)



Matthias Troyer (Microsoft, ETH)



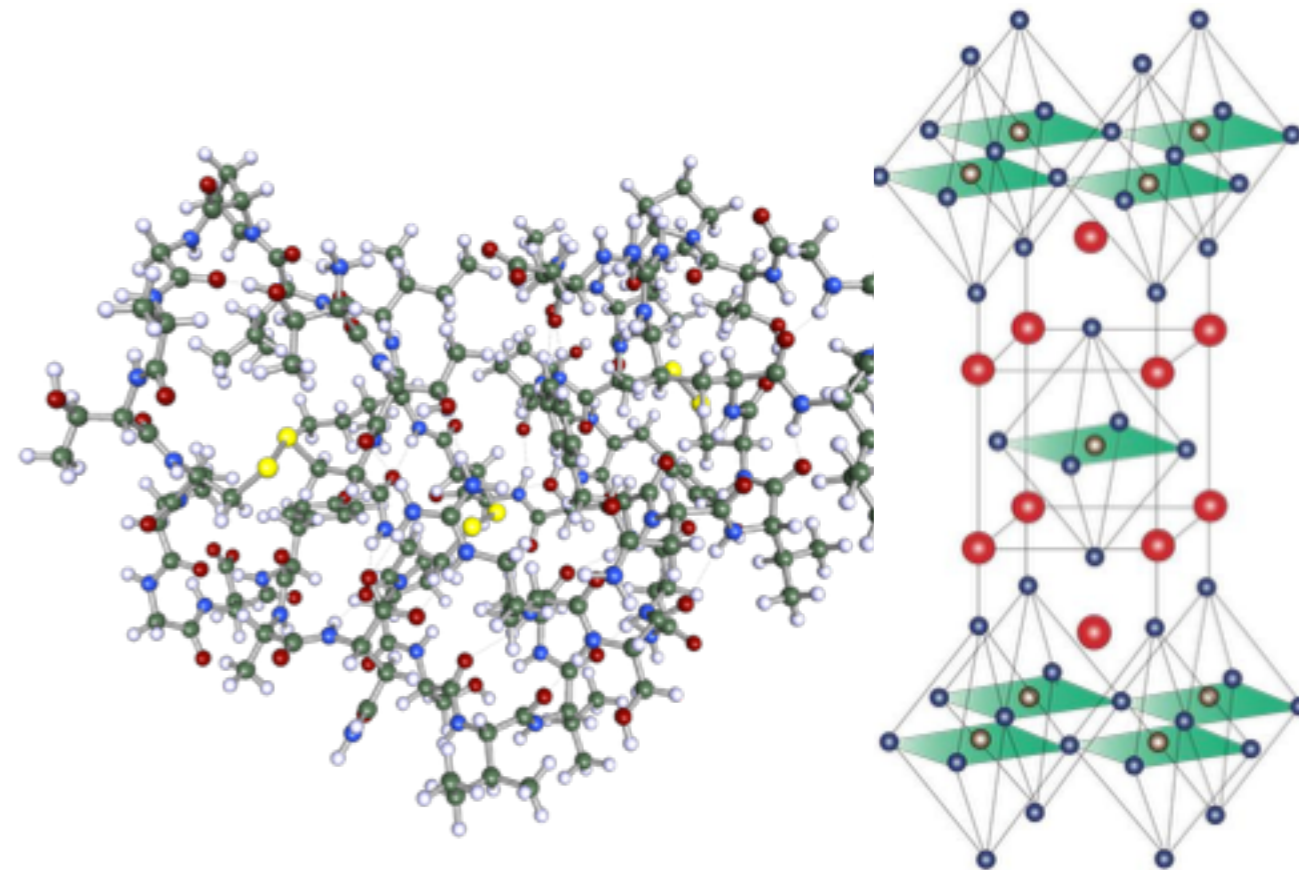
Guglielmo Mazzola (ETH)

**THE COMPLEXITY OF THE
MANY-BODY PROBLEM IN
CLASSICAL AND QUANTUM
MECHANICS**

THE MANY-BODY PROBLEM IN QUANTUM MECHANICS

- Generic specification of a quantum state requires resources exponentially large in the number of degrees of freedom N
- Today's best supercomputers can solve the wave equation **exactly** for systems with a maximum of ~ 45 particles.
- Storing the state of a 273 spin-1/2 system requires a computer with more bits than there are atoms in the universe
- Yet, technologically relevant problems in chemistry, condensed matter physics, and quantum computing are much larger than 273.
- Quantum computing (?)

$|\Psi\rangle$ vector with 2^N

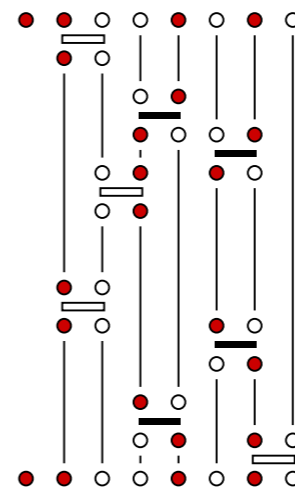


THERE IS STILL HOPE FOR CLASSICAL ALGORITHMS

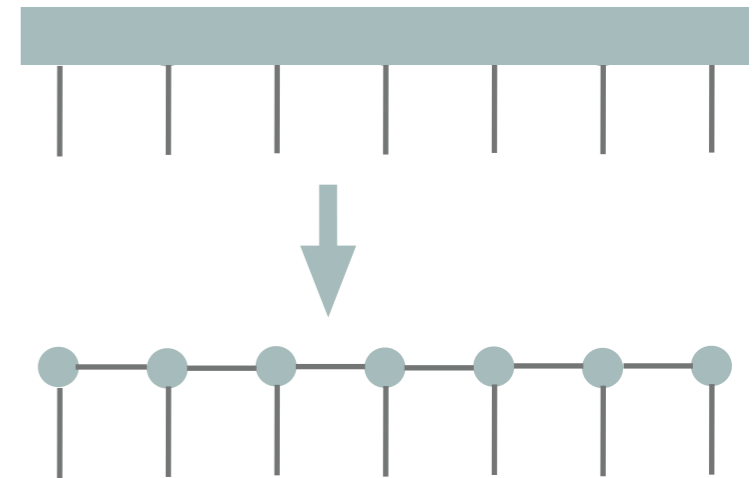
- Nature is sometimes compassionate: many-body systems can be typically characterized by an amount of information smaller than the maximum capacity of the corresponding state space.
- Quantum Monte Carlo and other numerical methods based on Tensor Networks exploit this fact and are able to accurately study large quantum system in practice with limited amount of resources.
- Machine learning community deals with equally high dimensional problems and battle the **curse of dimensionality** successfully with impressive results in a **wide spectrum of scientific and technologically relevant areas of research**.

$|\Psi\rangle$ vector with 2^N

QMC



Low entanglement
MPS and other TN



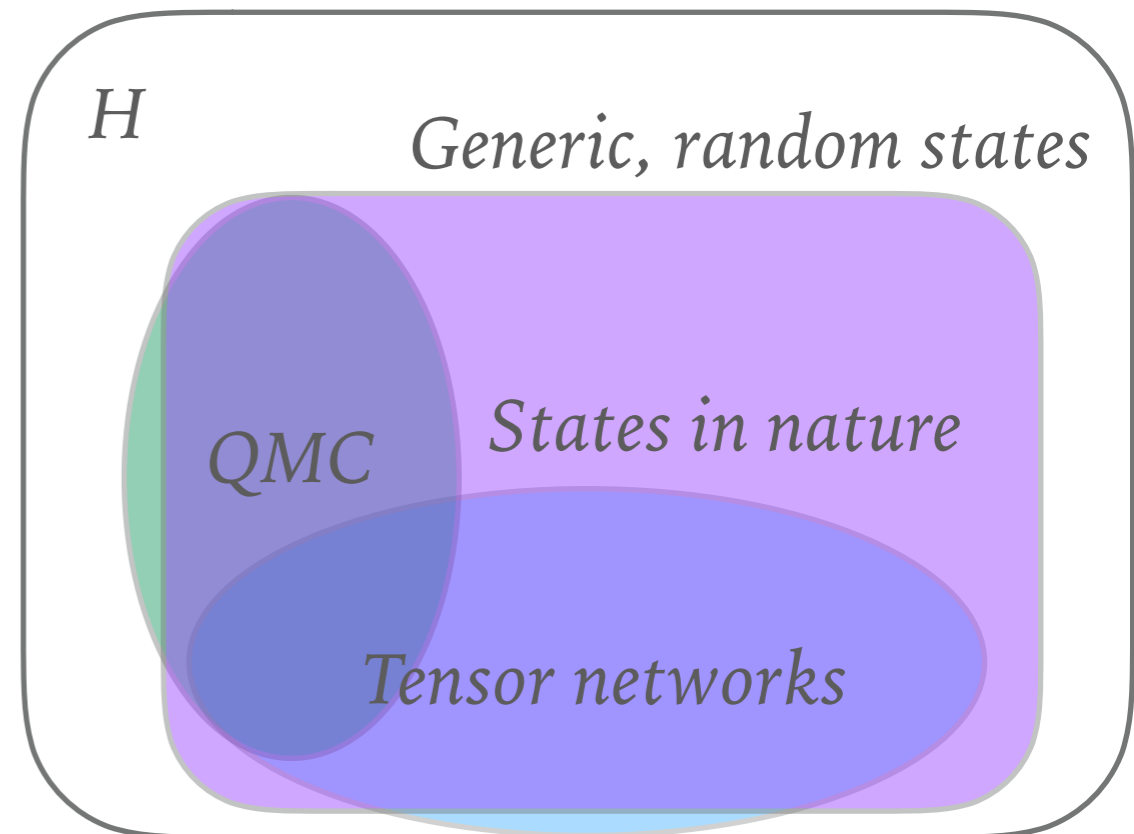
H

Generic, random states

QMC

States in nature

Tensor networks

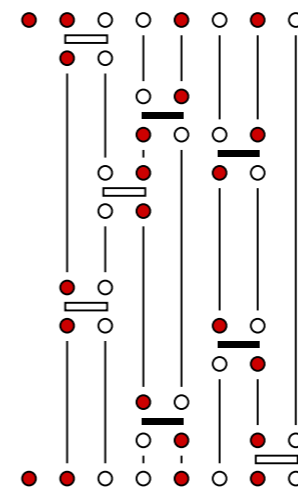


THERE IS STILL HOPE FOR CLASSICAL ALGORITHMS

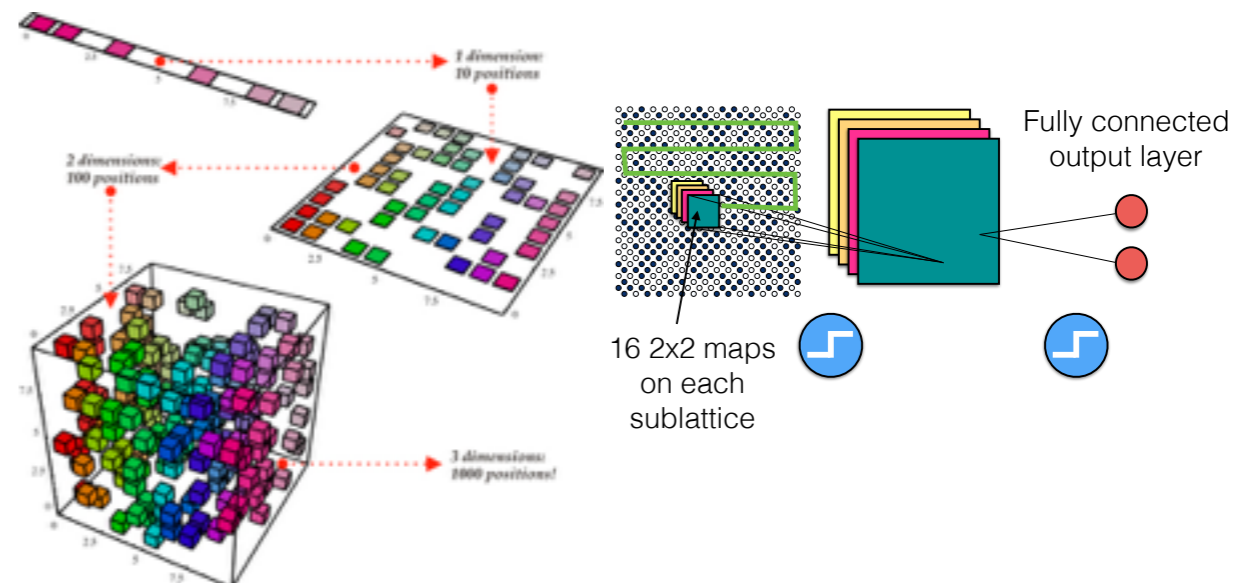
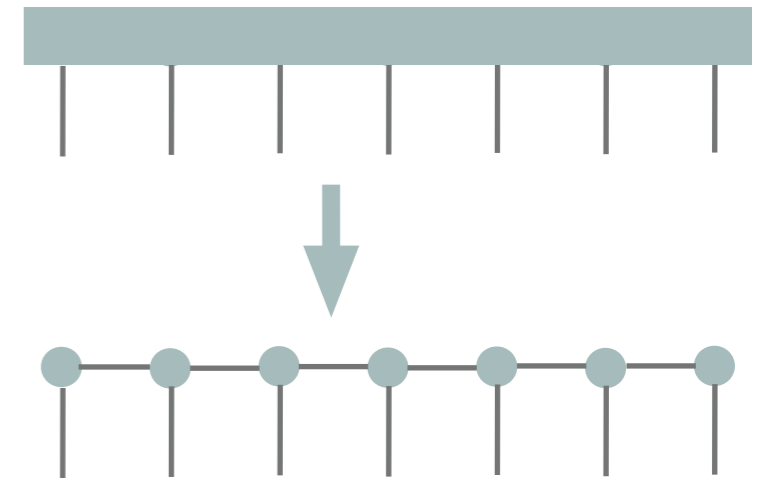
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$|\Psi\rangle$ vector with 2^N

QMC



Low entanglement
MPS and other TN



QUANTUM AND CLASSICAL MANY-BODY PHYSICS HAS NOT BEEN THE EXCEPTION

- [ML phases of matter/phase transitions](#) (Carrasquilla, Melko 1605.01735, Wang 1606.00318, Zhang, Kim, 1611.01518)
- [New ML inspired ansatz for quantum many-body systems](#) (Carleo, Troyer 1606.02318, Deng, Li, Das Sarma, 1701.04844, Deng, Li, Das Sarma 1609.09060, Carrasquilla, Melko 1605.01735)
- [Accelerated Monte Carlo simulations](#) (Huang, Wang 1610.02746)
- [Quantum state preparation guided by ML](#) (Bukov, Day, Sels, Weinberg, Polkovnikov and Mehta 1705.00565)
- [Renormalization group analyses, RBMs, PCA](#) (Bradde, Bialek 1610.09733, Koch-Janusz, Ringel 1704.06279, Mehta, Schwab, 1410.3831)
- [Quantum state tomography based on RBMs](#) (Torlai, Mazzola, Carrasquilla, Troyer, Melko, Carleo, 1703.05334)
- [ML based decoders for topological codes](#) (Torlai, Melko 1610.04238, Varsamopoulos, Criger, Bertels, 1705.00857)
- [Supervised Learning with Quantum-Inspired Tensor Networks](#) (Stoudenmire, Schwab 1605.05775, Novikov, Trofimov, Oseledets, 1605.03795)
- [Quantum Boltzmann machines](#) (Amin, Andriyash, Rolfe, Kulchytsky, Melko, 1601.02036, Kieferova, Wiebe, 1612.05204,)
- [Quantum machine learning algorithms to accelerate learning](#) (Biamonte, Wittek, Pancotti, Rebentrost, Wiebe, Lloyd, 1611.09347)

And many more

IN THIS TALK

- I will discuss several applications of ML ideas to problems in many-body physics.
- Supervised learning approach to classical phase transitions (Ising models)
- Briefly mention two quantum systems:
- Interpretation of the wave function as a neural network and write the ground state of Kitaev's toric code using convolutional neural networks.
- Data intensive problem in quantum mechanics: quantum state tomography with neural networks (RBMs) $|\Psi\rangle$ 2^N

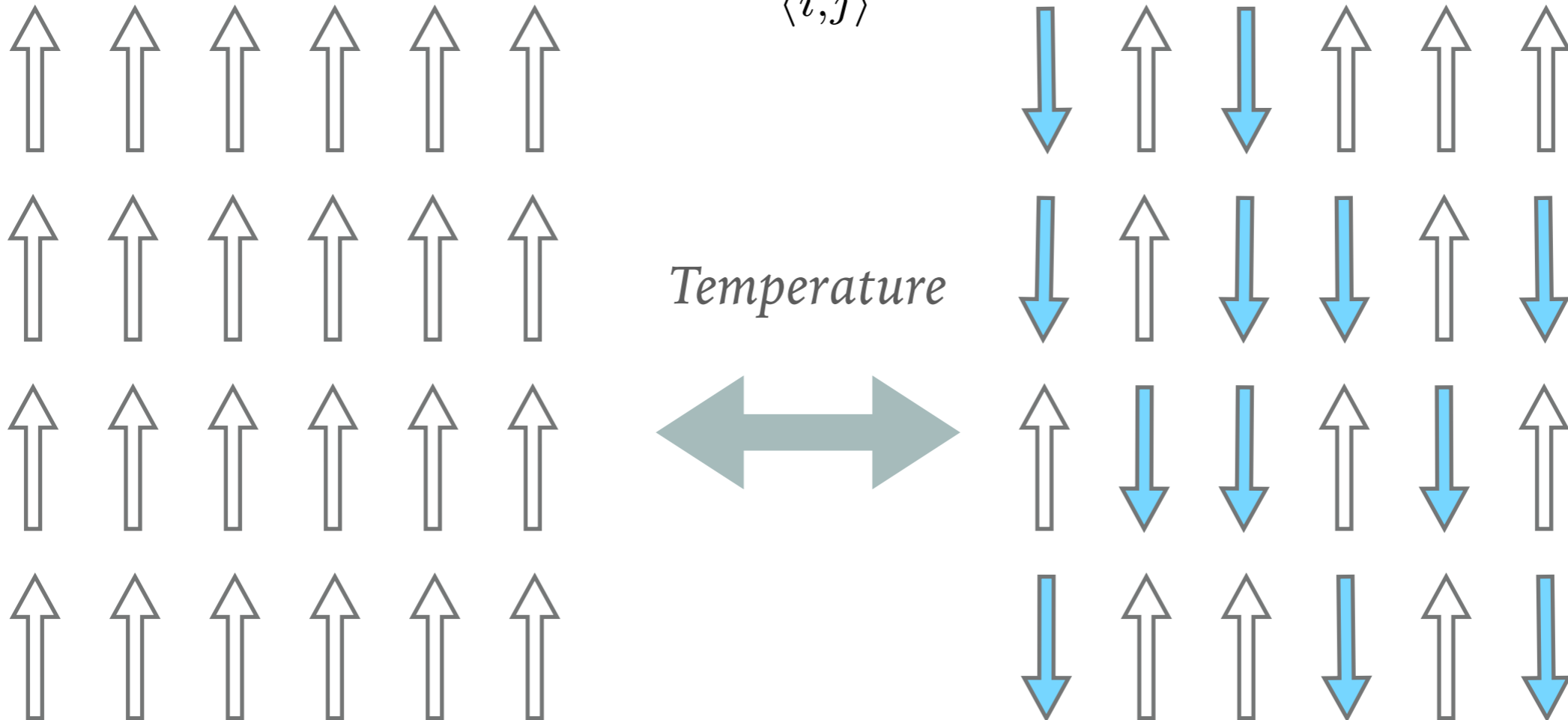
**SUPERVISED LEARNING
PERSPECTIVE OF
PHASES OF MATTER**

TOY PROBLEM: ISING MODEL

PHASES, PHASE TRANSITIONS, AND THE ORDER PARAMETER

Ising ferromagnet in two dimensions

$$E = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j \quad \sigma_i = \pm 1$$



Ferromagnet

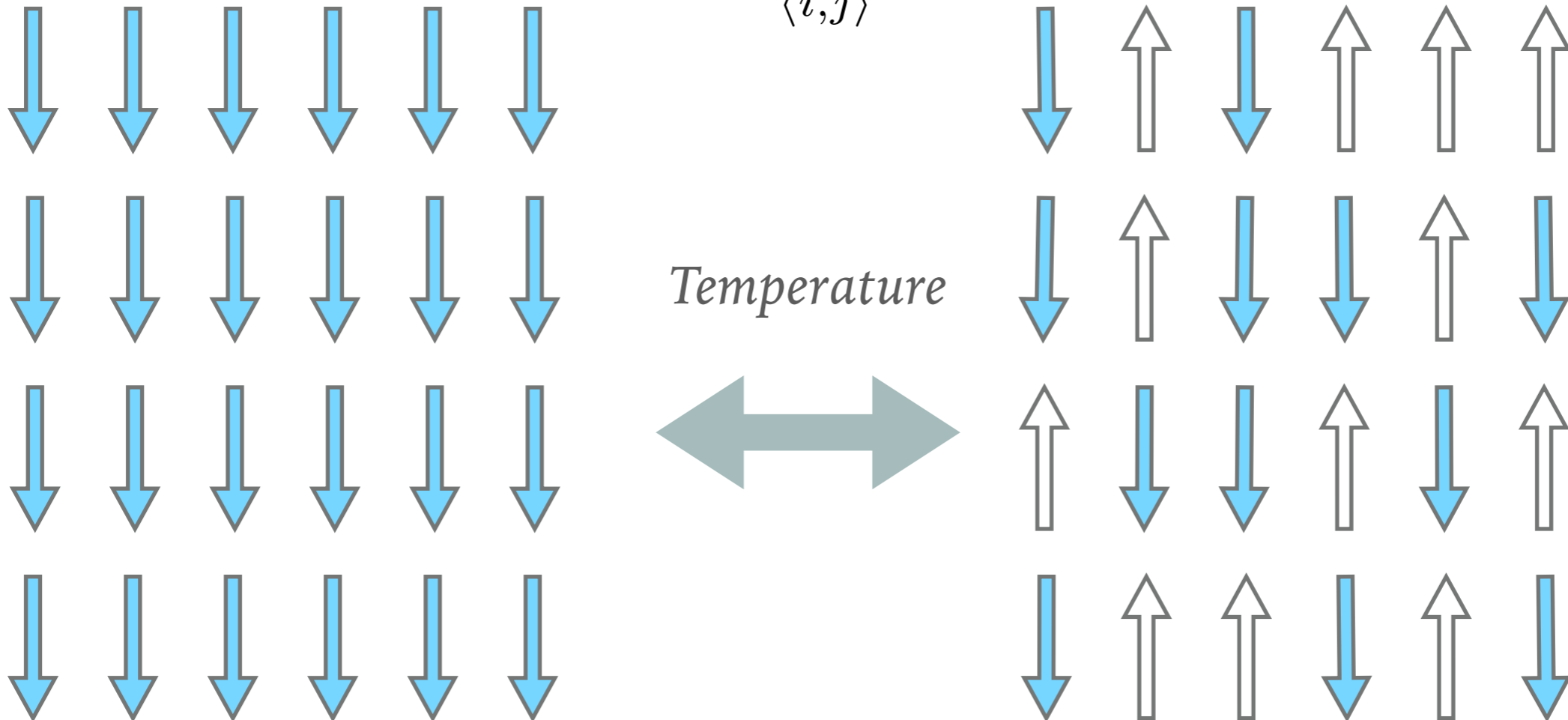
Lars Onsager *Phys. Rev.* 65, 117

Paramagnet

PHASES, PHASE TRANSITIONS, AND THE ORDER PARAMETER

Ising ferromagnet in two dimensions

$$E = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j$$

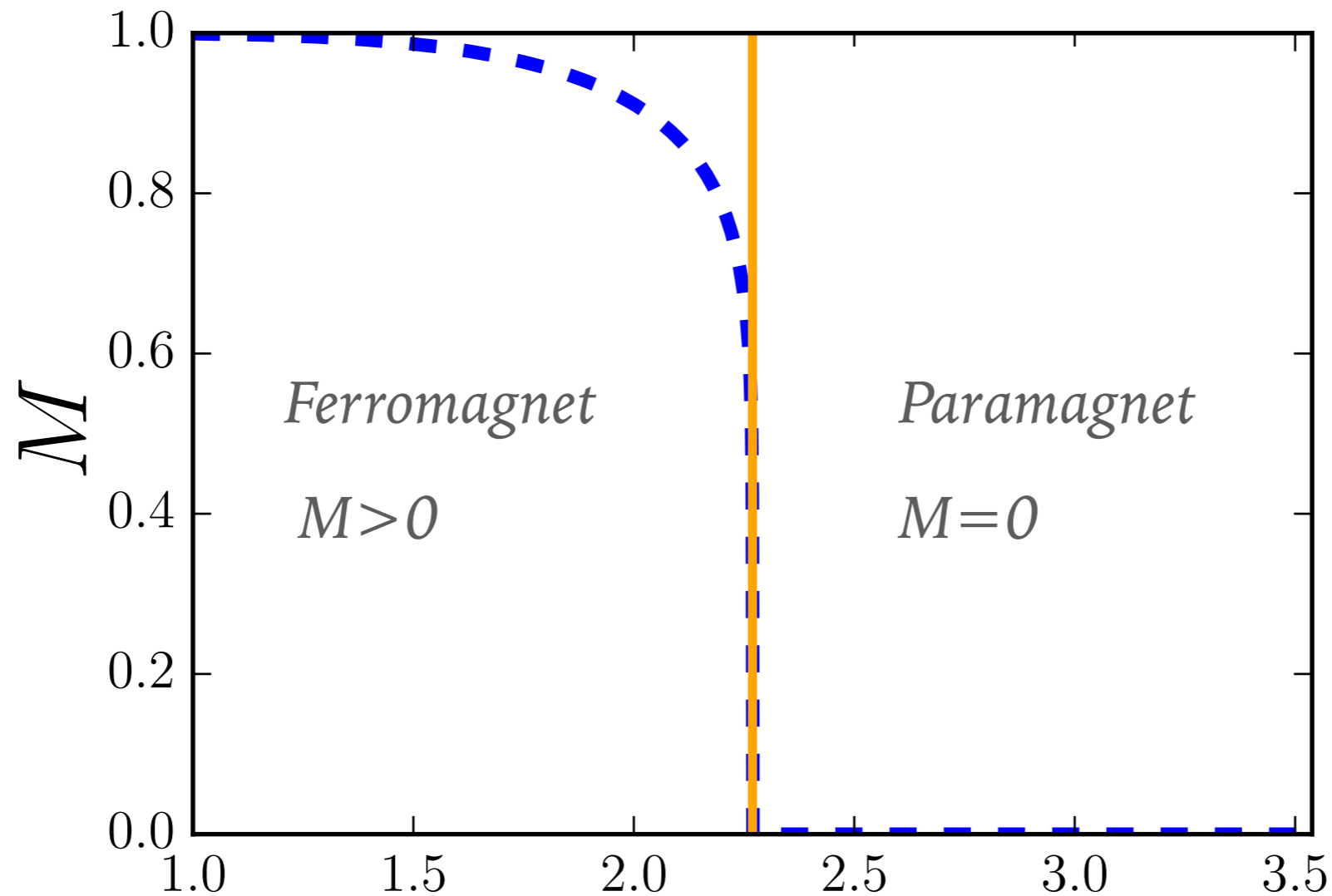


Ferromagnet

Paramagnet

PHASES, PHASE TRANSITIONS, AND THE ORDER PARAMETER

Ferromagnetic transition: order parameter

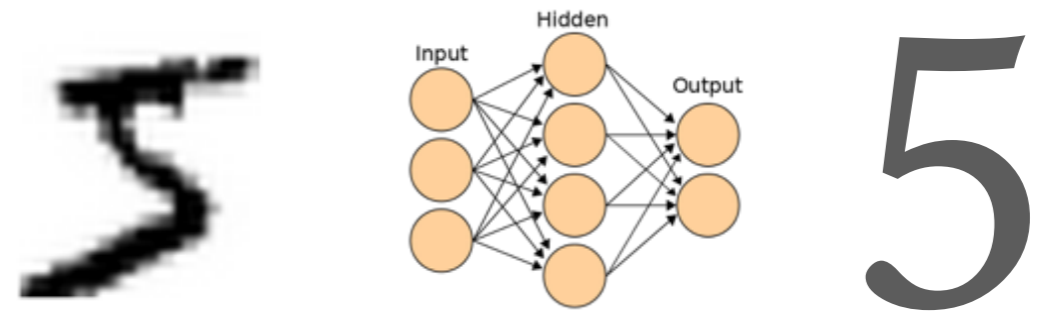


It is a measure of the degree of order in the system

$$M = \frac{1}{N} \sum_i \langle \sigma_i \rangle, \quad T \quad \sigma_i = \pm 1 \quad \text{Lars Onsager Phys. Rev. 65, 117 (1944)}$$

**WHAT DO I MEAN BY MACHINE LEARNING PHASES
OF MATTER?**

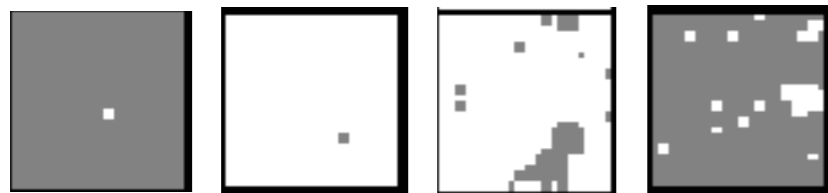
FLUCTUATIONS HANDWRITTEN DIGITS (MNIST)



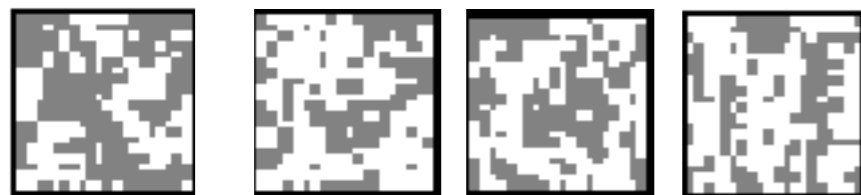
$\mathcal{S} = 5 + \text{fluctuations}$

ML community has developed powerful *supervised* learning algorithms

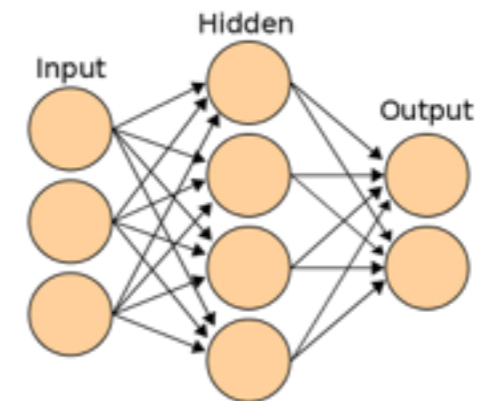
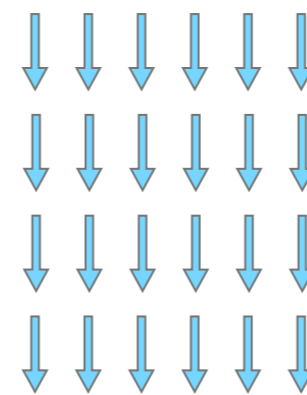
FM phase



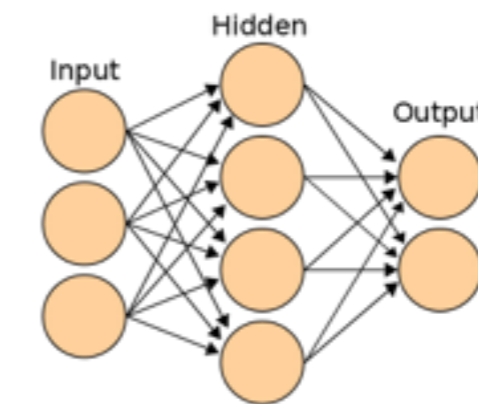
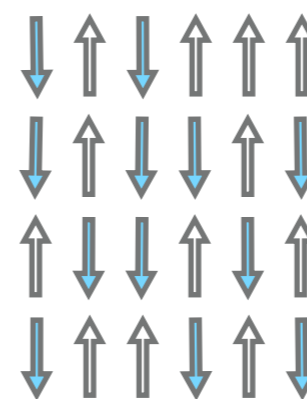
High T phase



gray = spin up
white = spin down



FM (0)



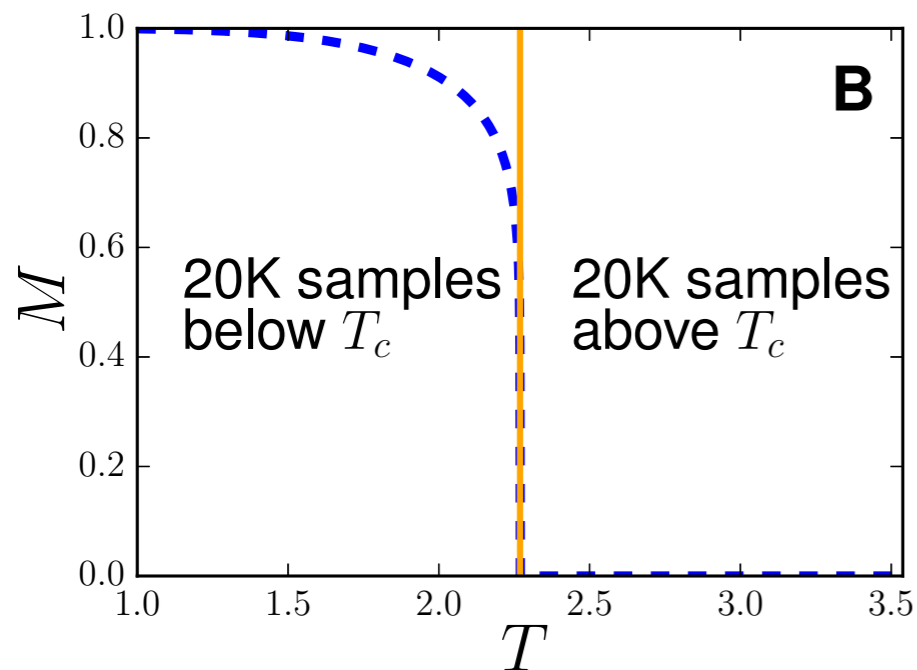
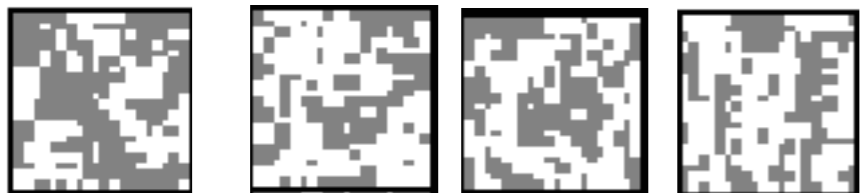
PM (1)

COLLECTING THE TRAINING/TESTING DATA: MC **SAMPLING** ISING MODEL AND **LABELS**

2D Ising model in the **ordered phase**

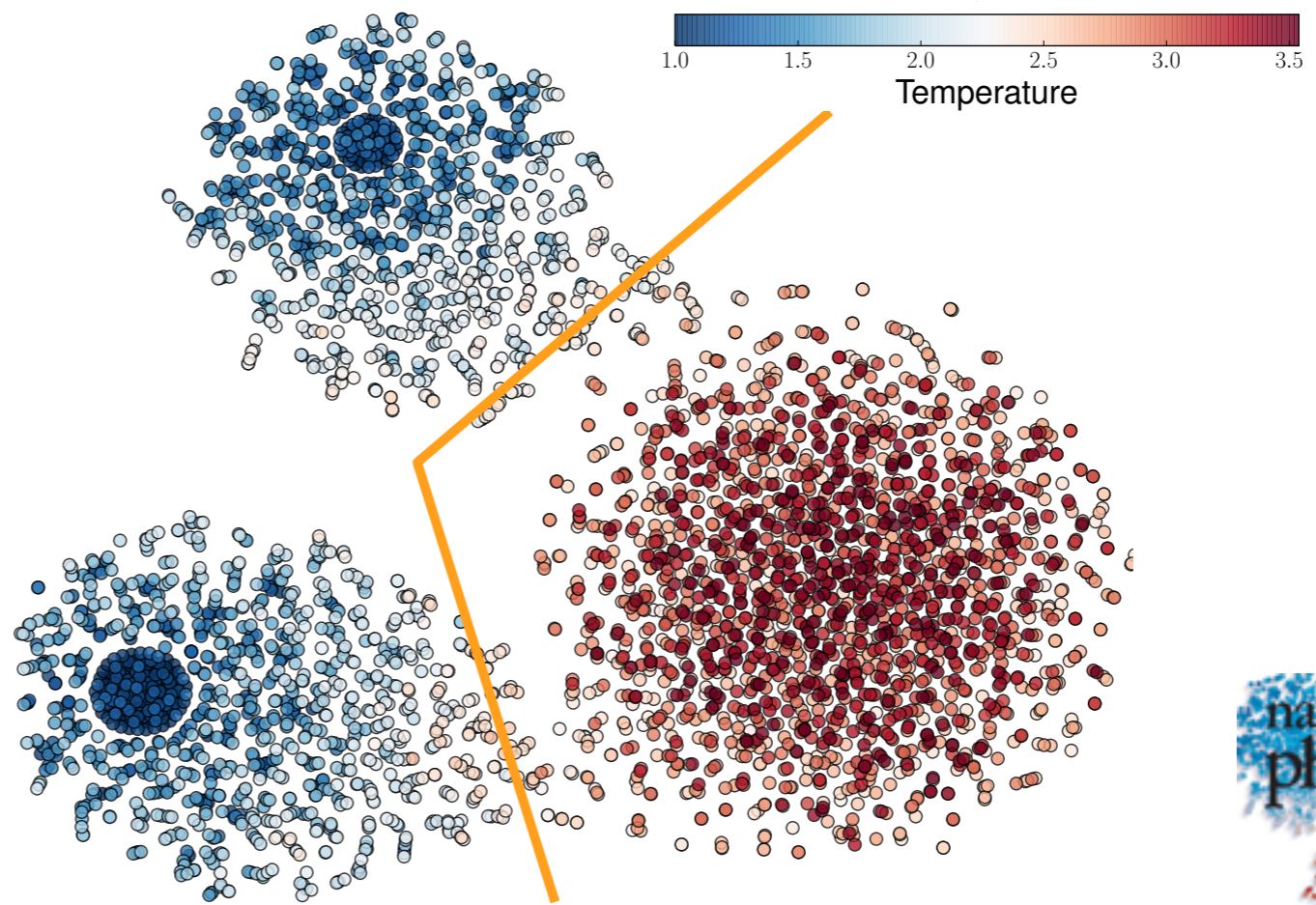


2D Ising model in the **disordered phase**

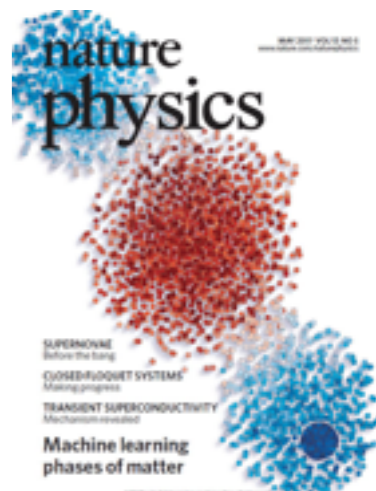


Training/testing data is drawn from the Boltzmann distribution

$$p(\sigma_1, \sigma_2, \dots, \sigma_N) = \frac{e^{-\beta E(\sigma_1, \sigma_2, \dots, \sigma_N)}}{Z(\beta)}$$

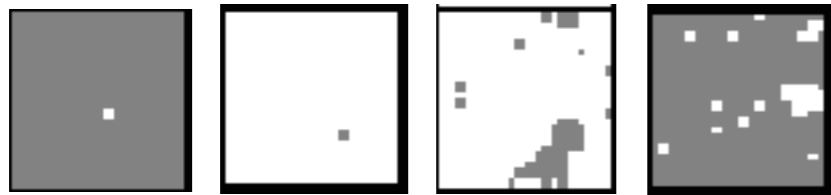


J. Carrasquilla and R. G. Melko. Nature Physics 13, 431–434 (2017)

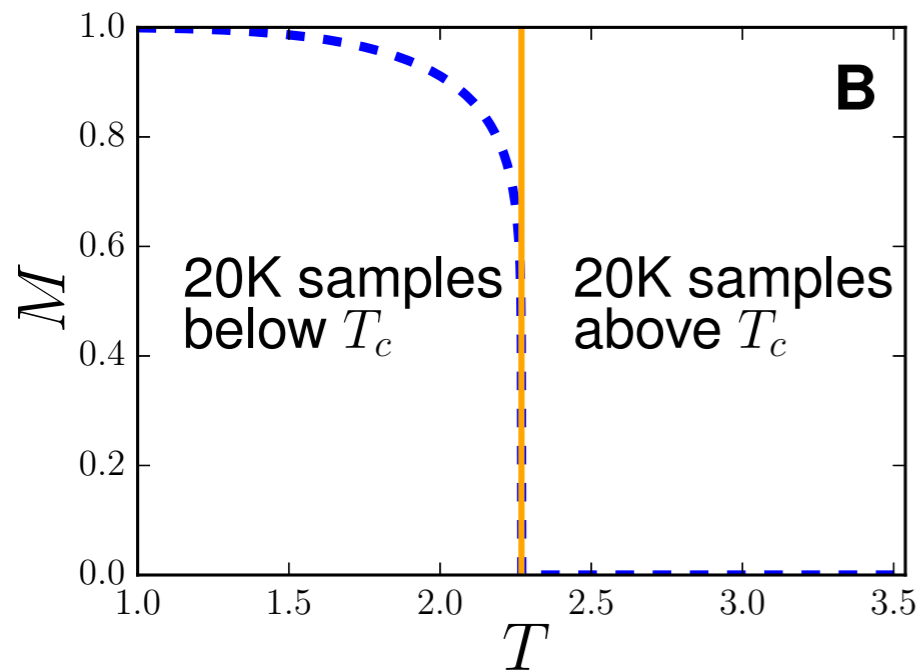
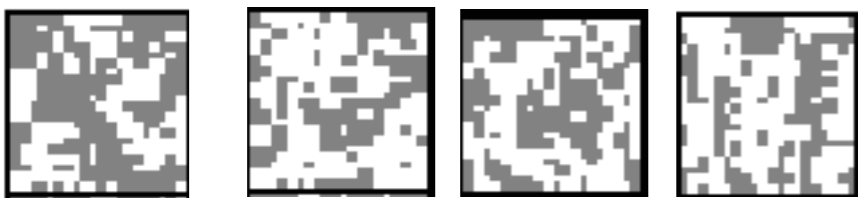


RESULTS: SQUARE LATTICE ISING MODEL (TEST SETS)

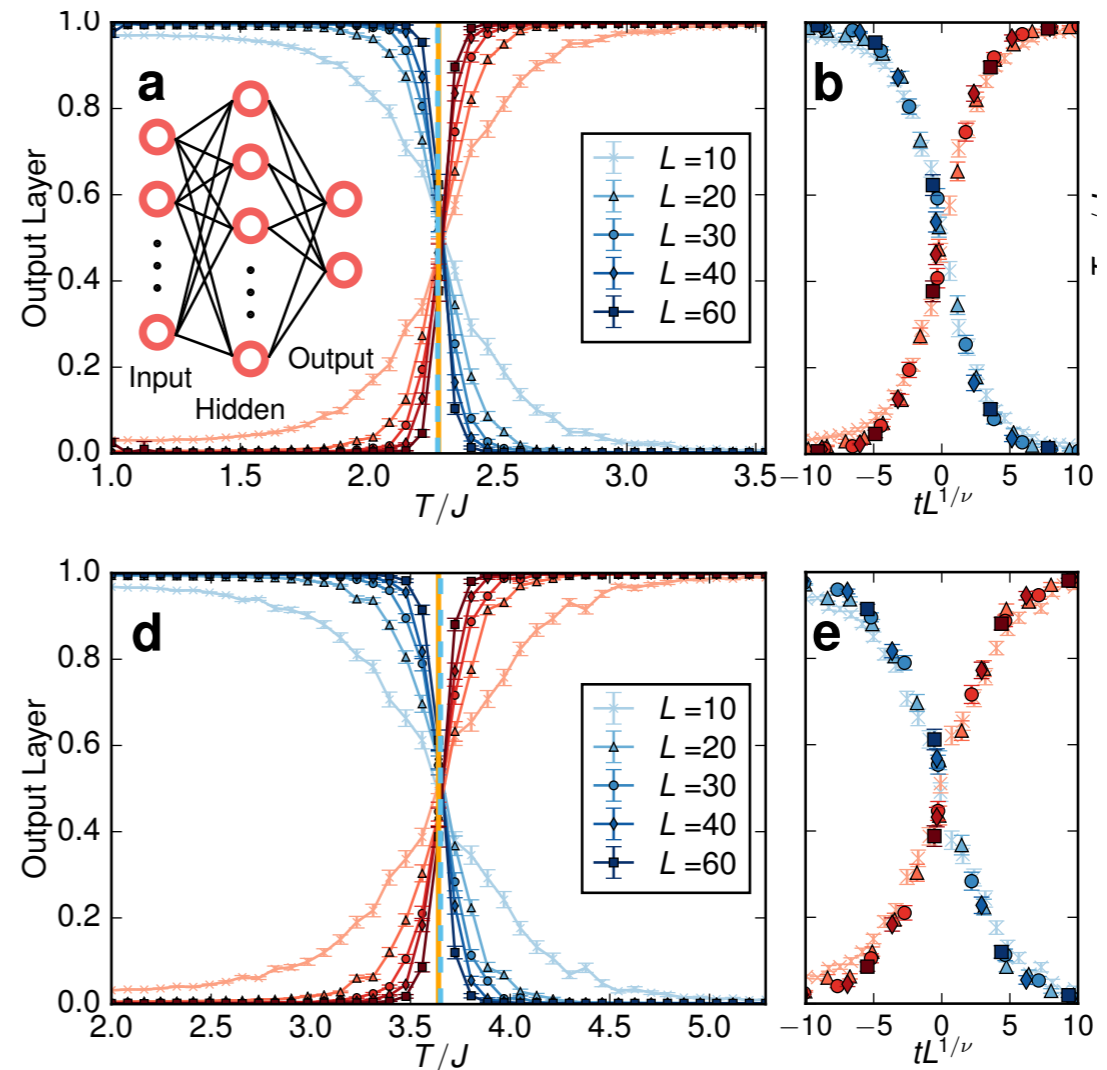
2D Ising model in the ordered phase



2D Ising model in the disordered phase



$$T_c/J = 2/\ln(1 + \sqrt{2})$$



T_c of Triangular within $< 1\%$

NN knows about criticality

$$\nu \approx 1$$

Analytical understanding: toy model with 3 analytically trained neurons. NN relies on the magnetization of the system

J. Carrasquilla and R. G. Melko. Nature Physics 13, 431–434 (2017)

**CAN WE DEAL WITH
DISORDERED AND TOPOLOGICAL
PHASES NOT DESCRIBED BY
ORDER PARAMETERS?**

PHASES OF MATTER **WITHOUT AN ORDER PARAMETER AT T=0**

- **Topological phases of matter**. Examples: Fractional quantum hall effect, quantum spin liquids, Ising gauge theory. Potential applications in topological quantum computing. Interestingly, these phases defy the Landau symmetry breaking classification.
- **Coulomb phases** = Highly correlated “spin liquids” described by electrodynamics. Examples: Common water ice and spin ice materials ($\text{Ho}_2\text{Ti}_2\text{O}_7$ and $\text{Dy}_2\text{Ti}_2\text{O}_7$)

PHASES OF MATTER WITHOUT AN ORDER PARAMETER AT $T=0, \infty$

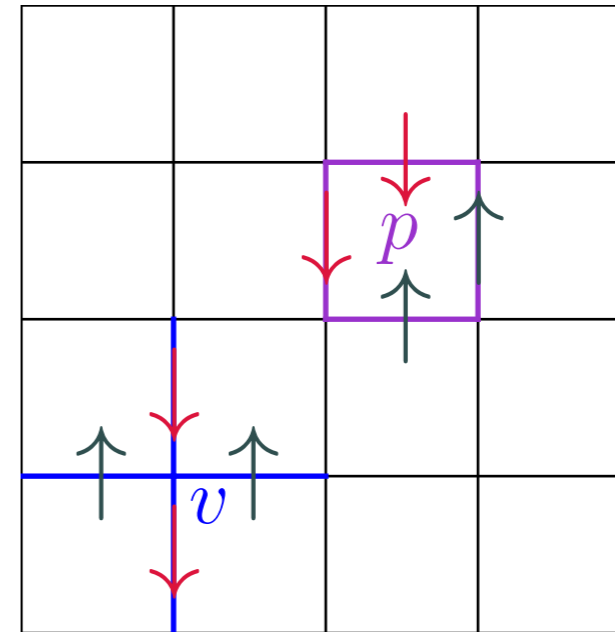
Wegner's Ising gauge theory:

$$H = -J \sum_p \prod_{i \in p} \sigma_i^z$$

F.J. Wegner, *J. Math. Phys.* 12 (1971) 2259

(Kogut *Rev. Mod. Phys.* 51, 659 (1979))

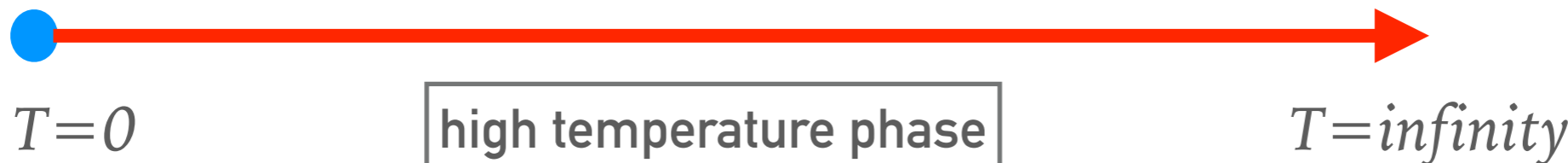
The ground state is a highly degenerate manifold with exponentially decaying spin-spin correlations.



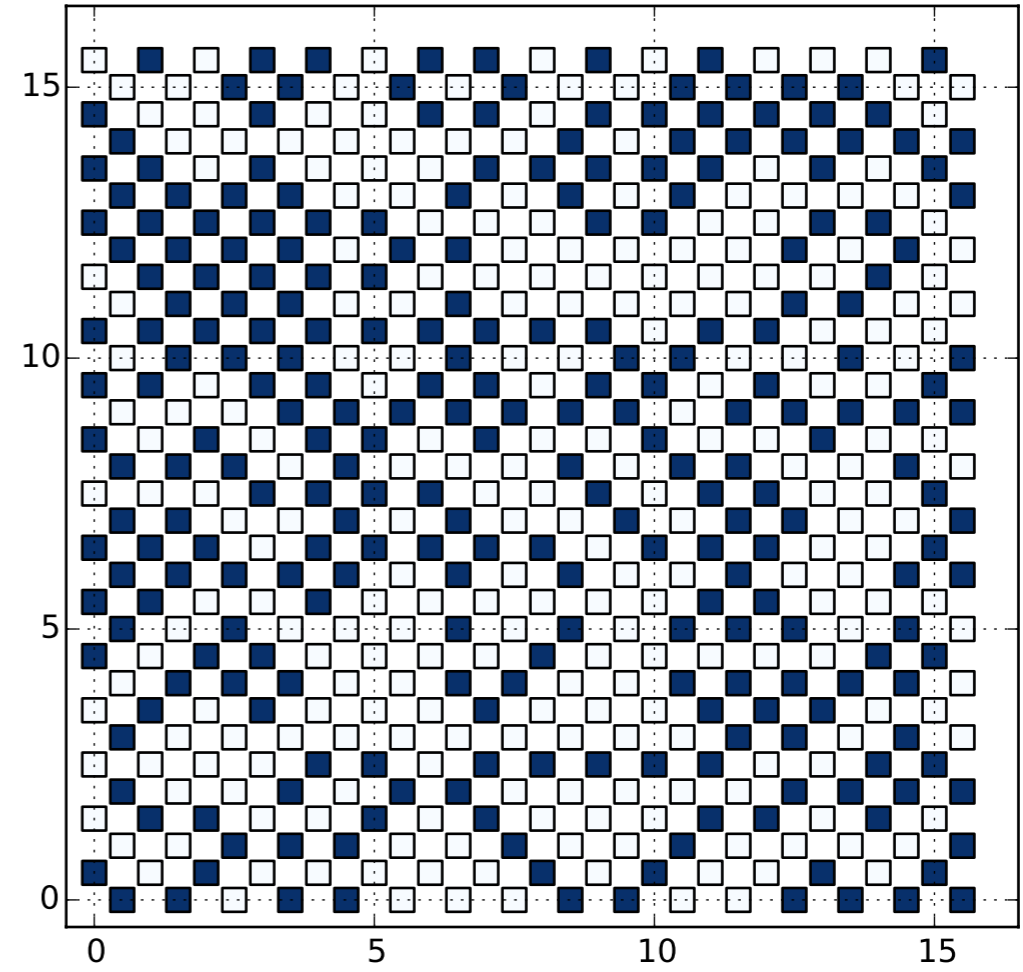
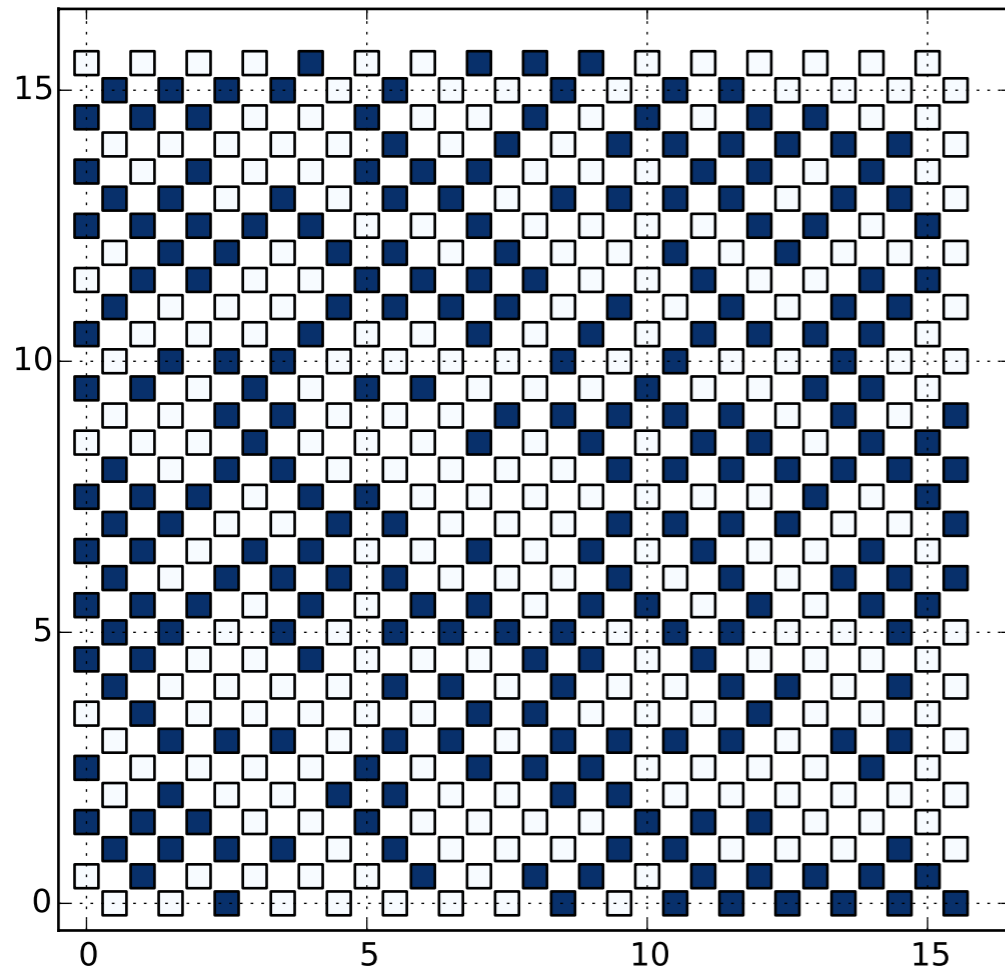
The grandmother of all lattice models for topological quantum computation

Ground state is a classical disordered topologically ordered phase

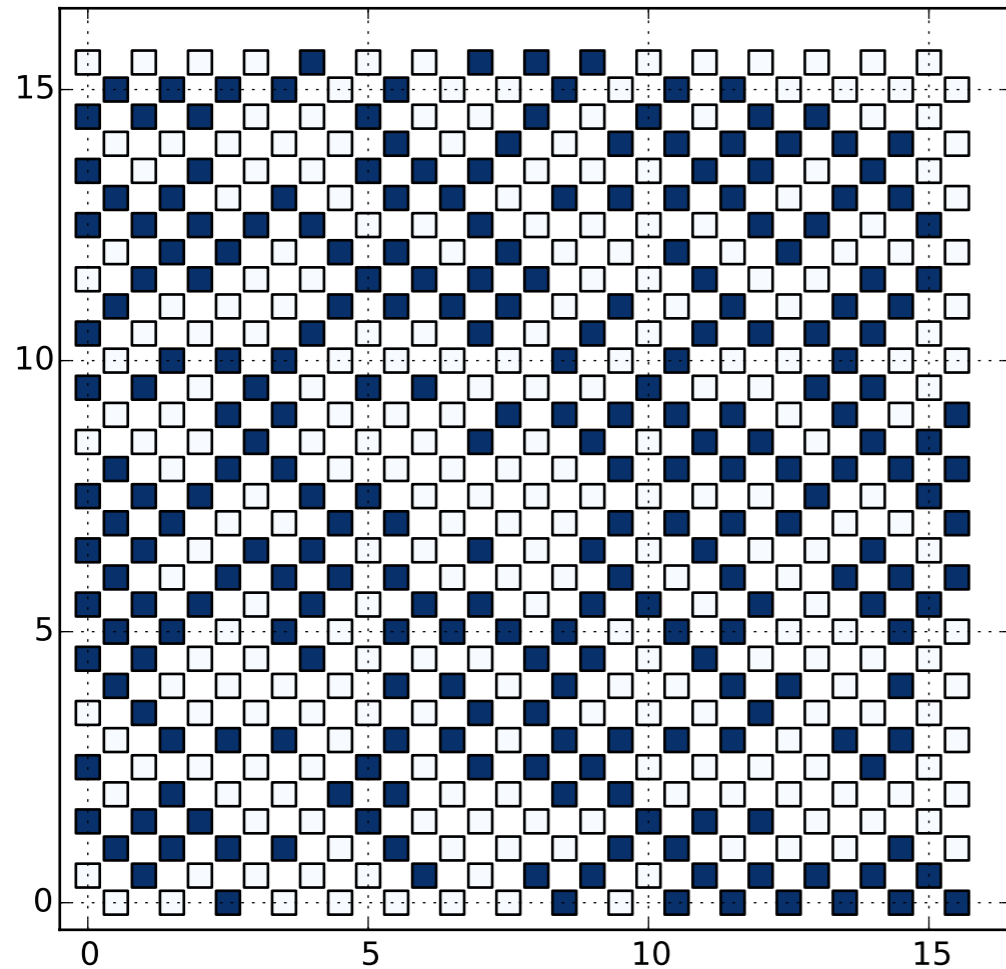
Castelnovo and Chamon *Phys. Rev. B* 76, 174416 (2007)



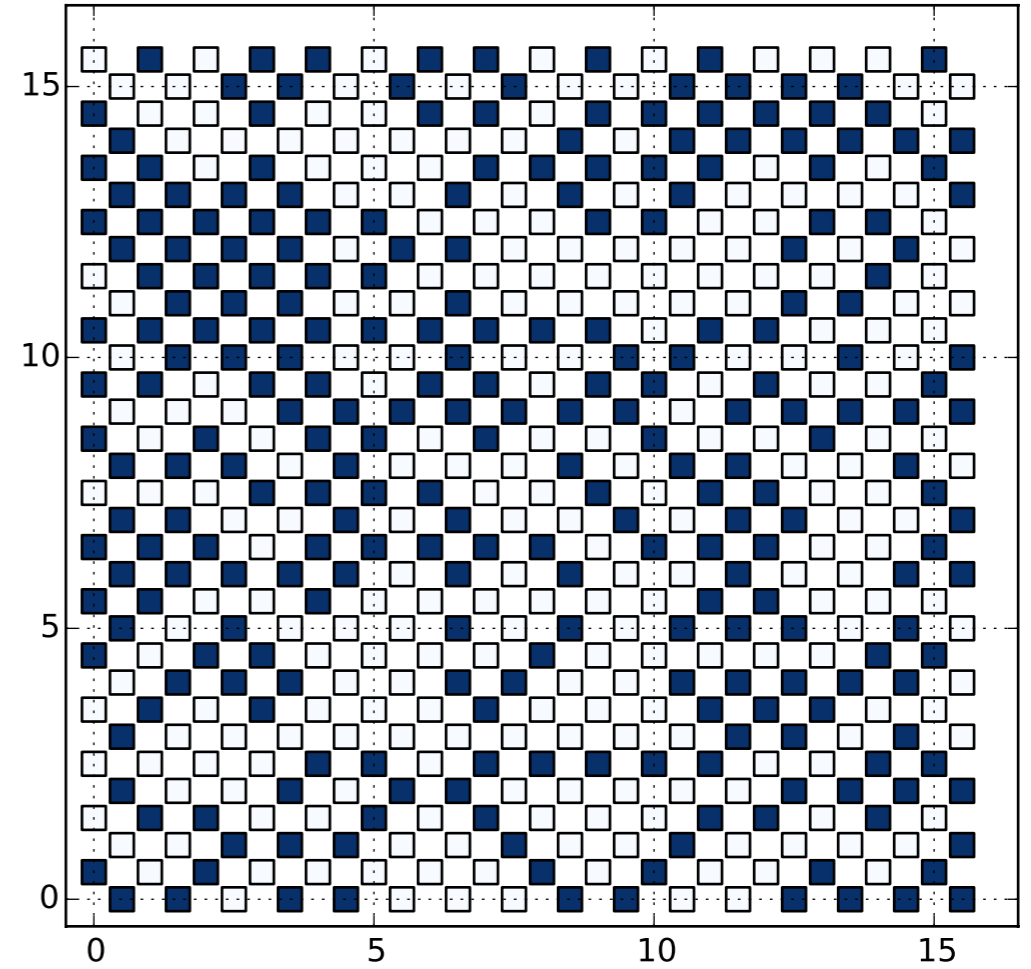
For two configurations



For two configurations



high-temperature state



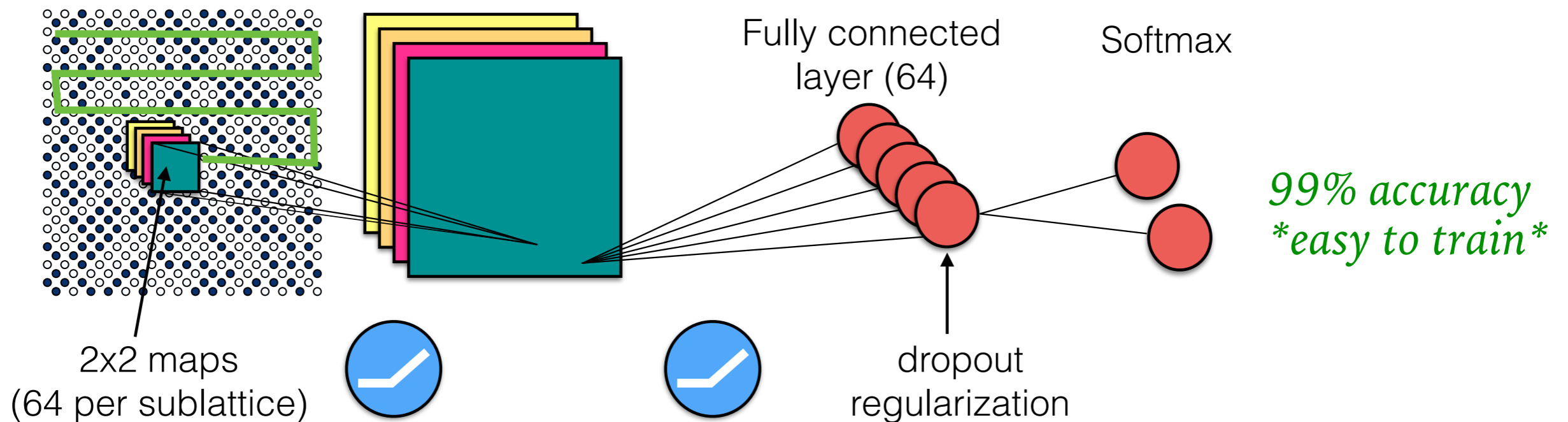
Ground state

Feedforward NN are difficult to apply to this problem and lead to 50% accuracy

ISING GAUGE THEORY

F.J. Wegner, J. Math. Phys. 12 (1971) 2259

$$H = -J \sum_p \prod_{i \in p} \sigma_i^z$$

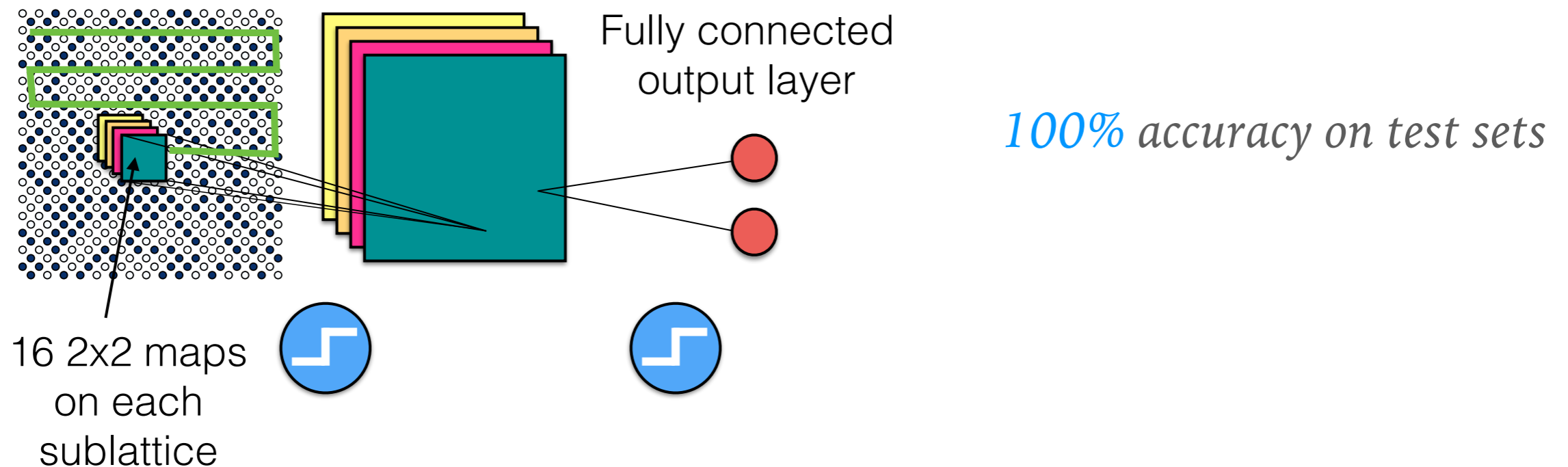


The picture we draw for what the CNN is using to distinguish the phases is that of the detection of satisfied local constraints. In few words, the neural network figures out the energy and uses it to classify states

ANALYTICAL UNDERSTANDING: WHAT DOES THE CNN USE TO MAKE PREDICTIONS?

-
- Based on this observation we derived the weights of a streamlined convolutional network *analytically* designed to work well for this problem:

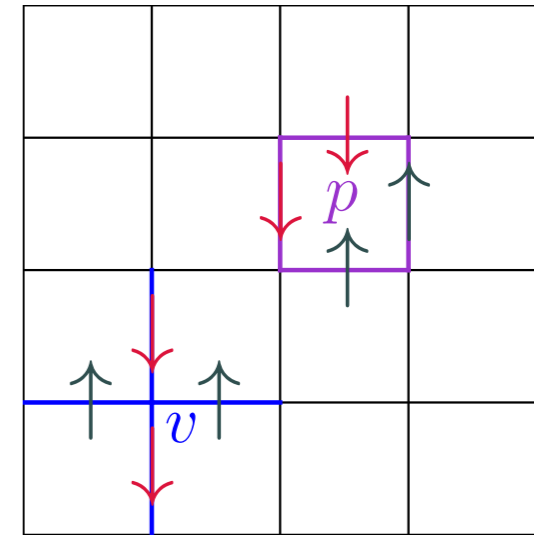
$$O_{\text{cold}}(\sigma_1, \dots, \sigma_N) \propto \lim_{\beta \rightarrow \infty} \exp \beta J \sum_p \prod_{i \in p} \sigma_i^z$$



**KITAEV'S QUANTUM ERROR
CORRECTING CODE WITH
CONVOLUTIONAL NEURAL
NETWORKS**

KITAEV'S TORIC CODE GROUND STATE

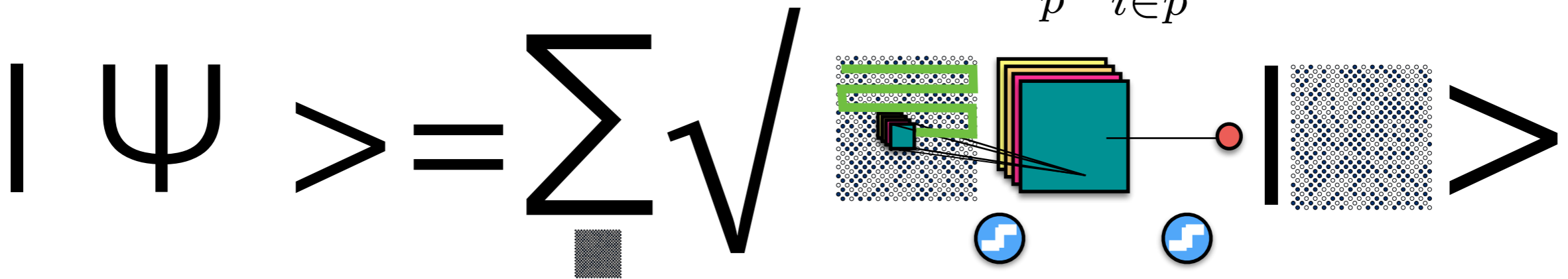
$$H = -J_p \sum_p \prod_{i \in p} \sigma_i^z - J_v \sum_v \prod_{i \in v} \sigma_i^x$$



$$|\Psi_{\text{TC}}\rangle \propto \lim_{\beta \rightarrow \infty} \sum_{\sigma_1, \dots, \sigma_N} e^{\frac{\beta}{2} J \sum_p \prod_{i \in p} \sigma_i^z} |\sigma_1, \dots, \sigma_N\rangle$$

PEPS : F. Verstraete, M. M. Wolf, D. Perez-Garcia, J. I. Cirac *Phys. Rev. Lett.* 96, 220601 (2006).

$$O_{\text{cold}}(\sigma_1, \dots, \sigma_N) \propto \lim_{\beta \rightarrow \infty} \exp \beta J \sum_p \prod_{i \in p} \sigma_i^z$$



J. Carrasquilla and R. G. Melko. *Nature Physics* 13, 431–434 (2017)

Carleo & Troyer, *Science* (2017)

Dong-Ling Deng et al *Phys. Rev. X* 7, 021021 (2017)

Jing Chen, Song Cheng, Haidong Xie, Lei Wang, Tao Xiang *arXiv:1701.04831 RBMs*

MESSAGES

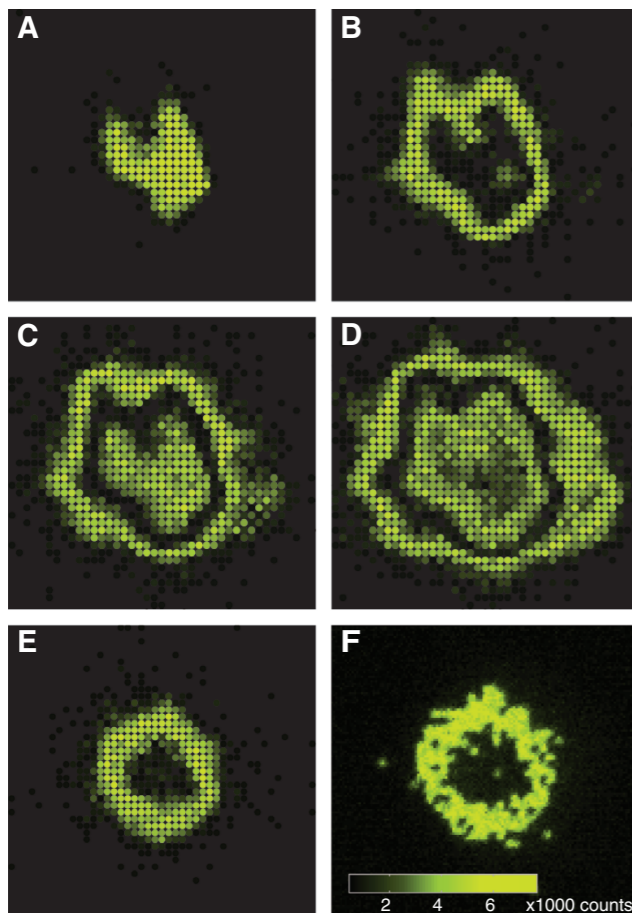
- With a neural network with a small number of parameters we are able to write down analytically the ground state of a system.
- Neural networks seem to enable very good compression quantum many-body states. (analogous to tensor networks)
- No limitations in the dimensionality of the systems
- More importantly, numerical procedures can be constructed to study other systems for where analytical results are elusive (many of which have been developed here in Trieste).
- Potential applications such as solving for ground states in condensed matter and quantum chemistry Hamiltonians, quantum state tomography, etc.

**NEURAL-NETWORK QUANTUM
STATE TOMOGRAPHY FOR
MANY-BODY SYSTEMS**

QUANTUM STATE TOMOGRAPHY

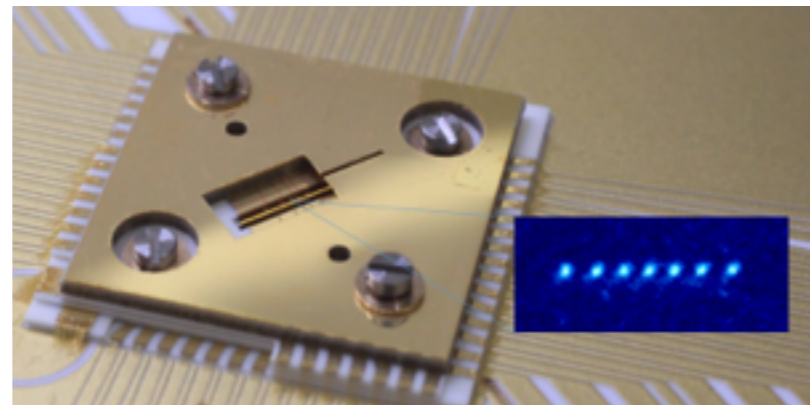
- Problem: Can we reconstruct the quantum state of a physical system from a limited set of experimentally accessible set of measurements?

Cold atomic gases



W. Bakr et al, Science (2010)

Trapped ions



ETH Trapped Ion Quantum Information Group

Quantum devices



QST is used as a diagnostic tool in experiments and implementation of technologically relevant quantum algorithms. Measurements required for QST are routinely available in these devices and other systems

THE PROBLEM AND THE REQUIREMENTS OF QST

- Problem: Can we reconstruct the quantum state of a physical system from a limited set of experimentally accessible set of measurements?

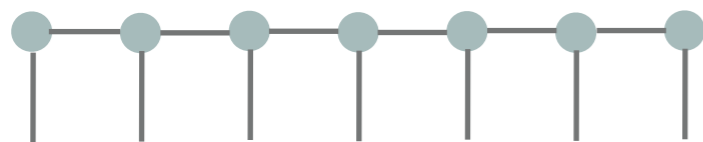
Requirements for QST of **large** systems (for **small** systems QST traditionally requires exponential resources)

-Compact representation of the state: **Neural Networks, MPS.**

-Set of projective measurements in different bases $|\psi_{\lambda,\mu}(\sigma^{[b]})|^2 \simeq P_b(\sigma^{[b]})$

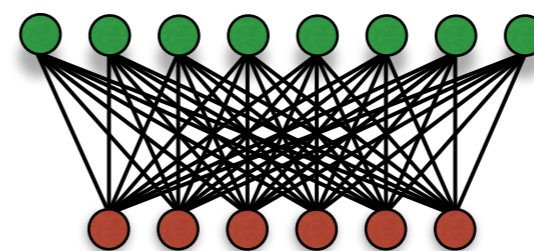
-A learning procedure that makes use of the data to learn the state. It is inherently a big-data problem: **Unsupervised learning (maximum likelihood estimation MLE)**

MPS QST



Cramer et al, Nat. Comm. (2010)

RBM QST



Torlai, Mazzola, Carrasquilla, Troyer, Melko and Carleo 1703:05334 to appear in Nature Physics

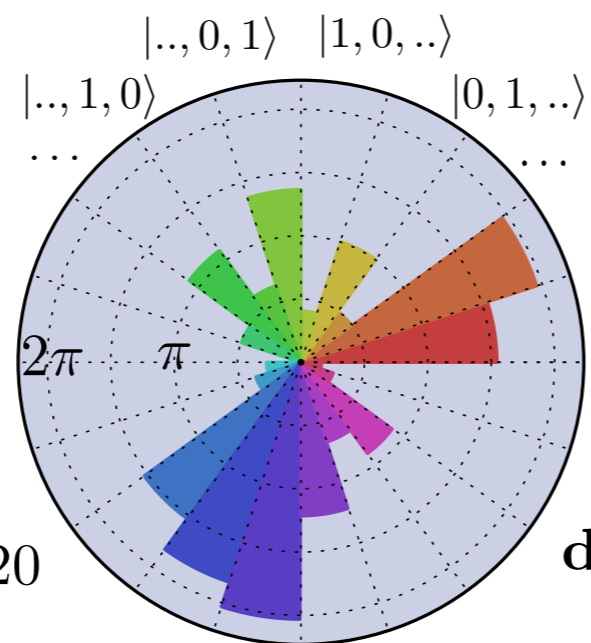
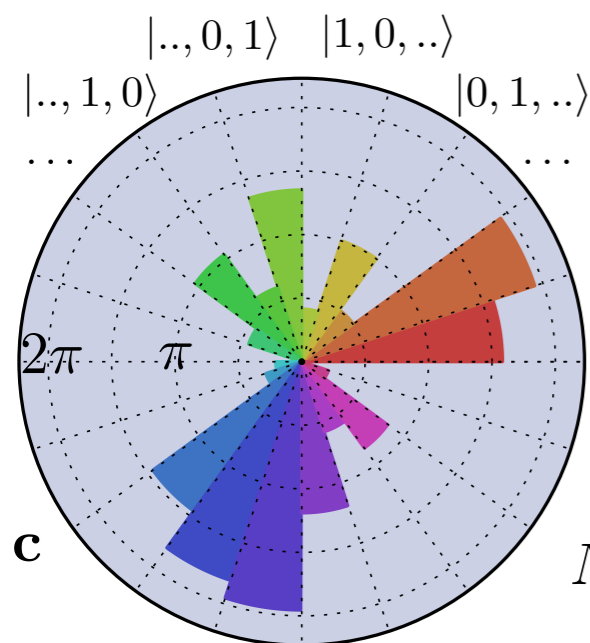
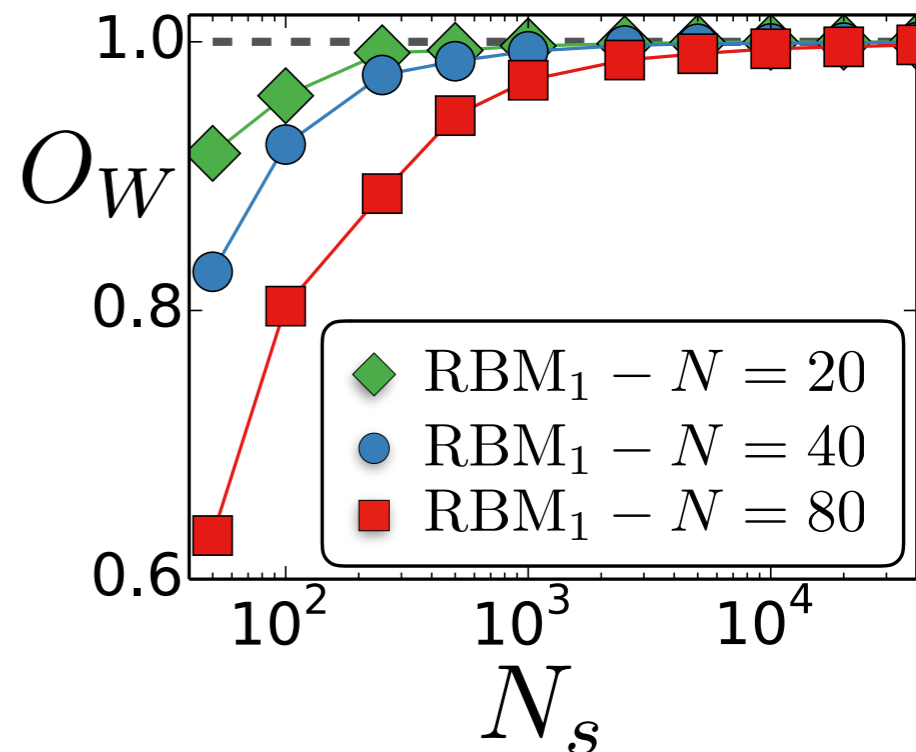
W state

Target wavefunction:

$$|\Psi_W\rangle = \frac{1}{\sqrt{N}} (|100\dots\rangle + |010\dots\rangle + \dots |0\dots 01\rangle)$$

$$O^2 = |\langle \psi_\lambda | \Psi_W \rangle|^2 = \left\langle \frac{\Psi_W(\boldsymbol{\sigma})}{\sqrt{p_\lambda(\boldsymbol{\sigma})}} \right\rangle_{p_\lambda} \times \left\langle \frac{\sqrt{p_\lambda(\boldsymbol{\sigma})}}{\Psi_W(\boldsymbol{\sigma})} \right\rangle_{|\Psi_W|^2}$$

$$|\tilde{\Psi}_W\rangle = \frac{1}{\sqrt{N}} \left(e^{i\theta_1} |100\dots\rangle + e^{i\theta_2} |010\dots\rangle + \dots + e^{i\theta_N} |0\dots 01\rangle \right)$$



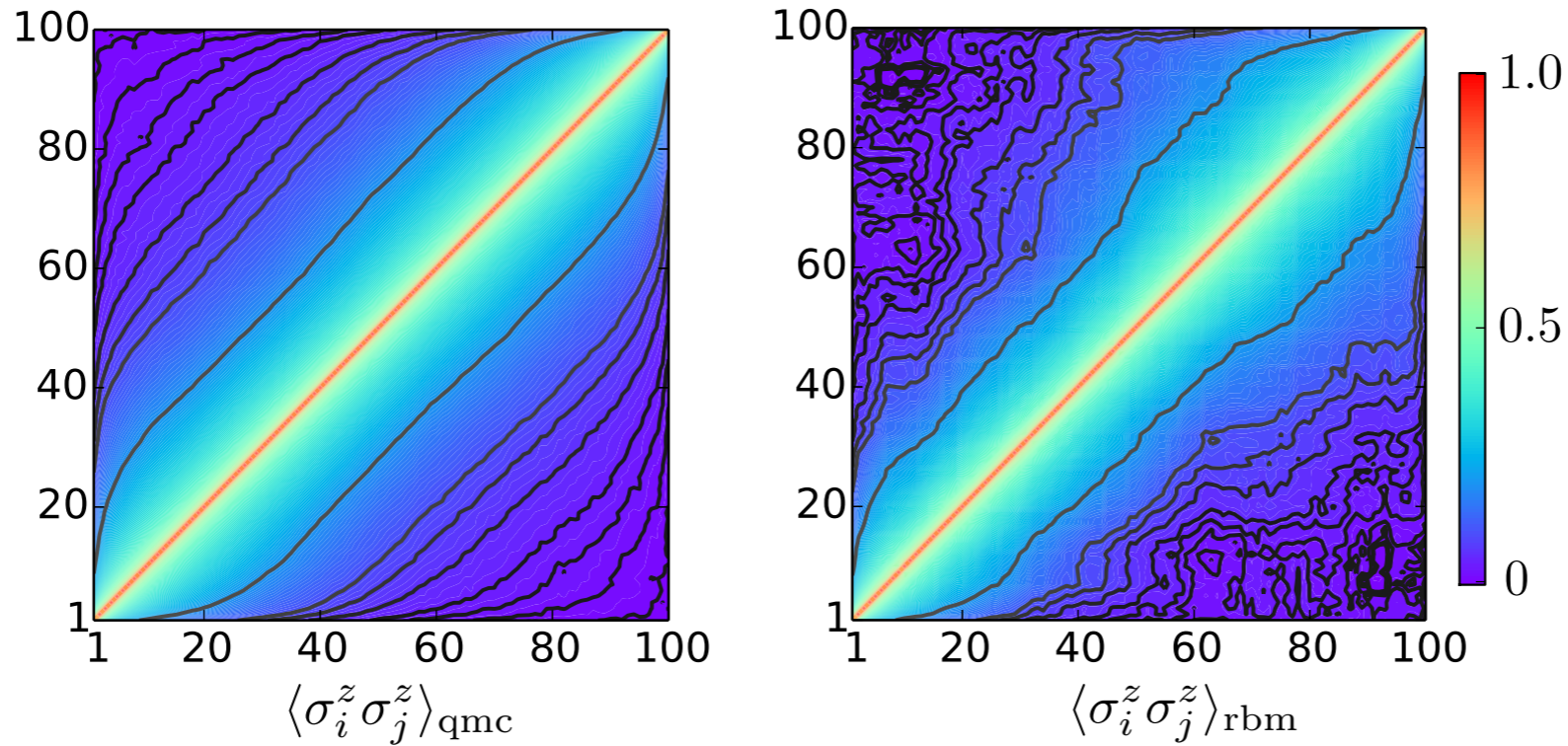
Bases: $\{X, X, Z, Z, \dots\}, \{Z, X, X, Z, \dots\}, \dots$
 $\{X, Y, Z, Z, \dots\}, \{Z, X, Y, Z, \dots\}, \dots$

$\theta(\boldsymbol{\sigma}_k)$ - Exact

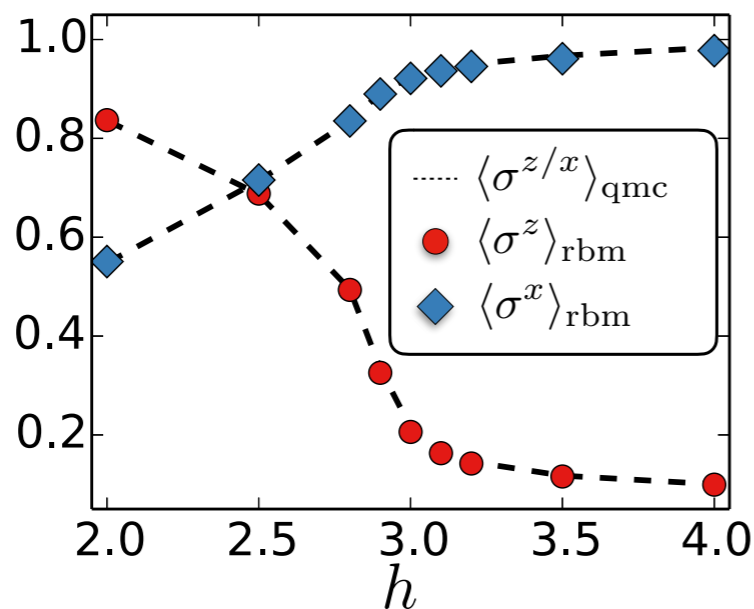
$\phi_\mu(\boldsymbol{\sigma}_k)$ - RBM

Many-body Hamiltonians: ground state Basis: $\{Z, Z, Z, Z, \dots\}$

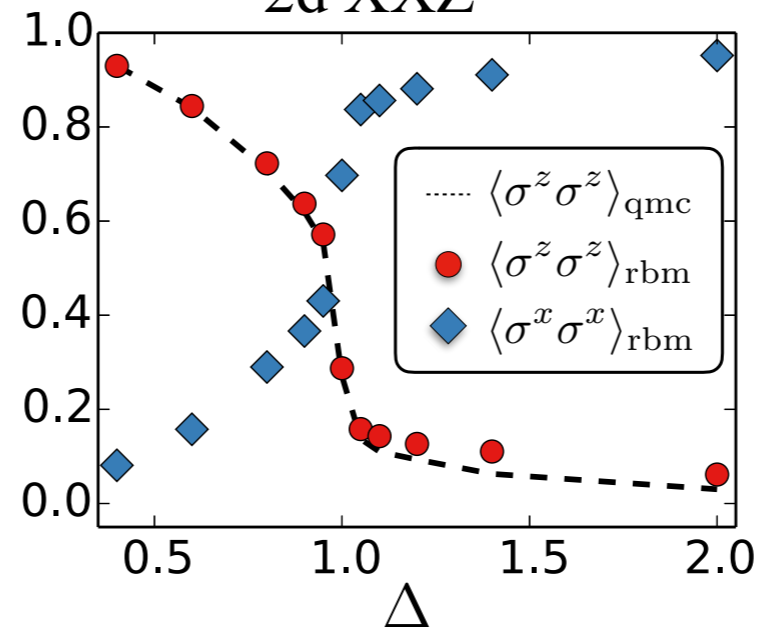
1d TFIM



2d TFIM



2d XXZ



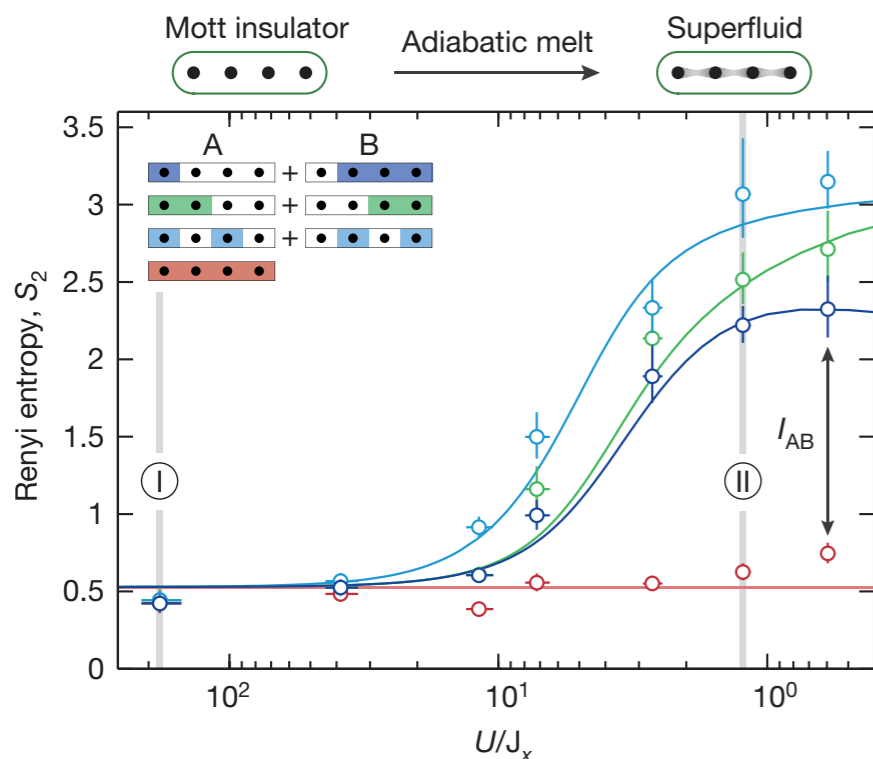
Many-body Hamiltonians: entanglement Basis: $\{Z, Z, Z, Z, \dots\}$

Renyi entropies:

$$S_\alpha(\rho_A) = \frac{1}{1-\alpha} \log(\text{Tr}(\rho_A^\alpha))$$

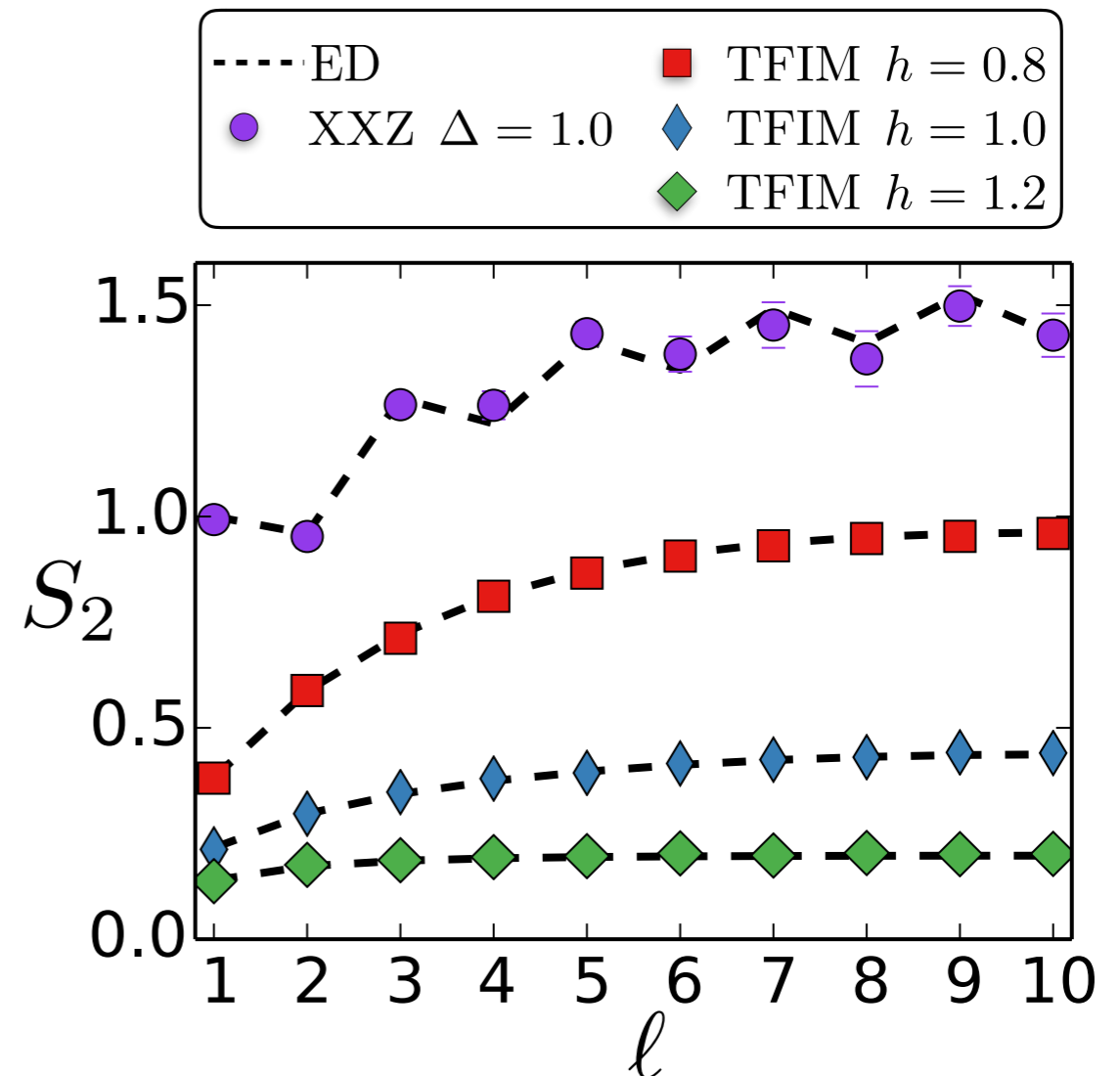
Hastings *et al*, PRL (2010)

$$S_2(\rho_A) = -\log(\langle \text{Swap} \rangle)$$



R. Islam *et al*, Nature (2015) (4 particles)

$N = 20$



For experiments: unequivocal evidence of the quantum behaviour and some of the most profound and striking manifestations of quantum behaviour.

For RBMs: Shows that the learned RBMs generalize very well

Hard to get experimentally. Our approach suggests an experimentally viable way to do it.

CONCLUSION

- We encode and discriminate phases and phase transitions, both conventional and topological, using neural network technology.
- We have a solid understanding of what the neural nets do in those cases through controlled analytical models.
- We have performed QST based on neural networks with some results that are better than the state-of-the-art.
- Potential for discovery enabled by ML. Let's explore this some more!

