

Cavity Optomagnonics

Silvia Viola Kusminskiy



MAX PLANCK INSTITUTE

for the science of light



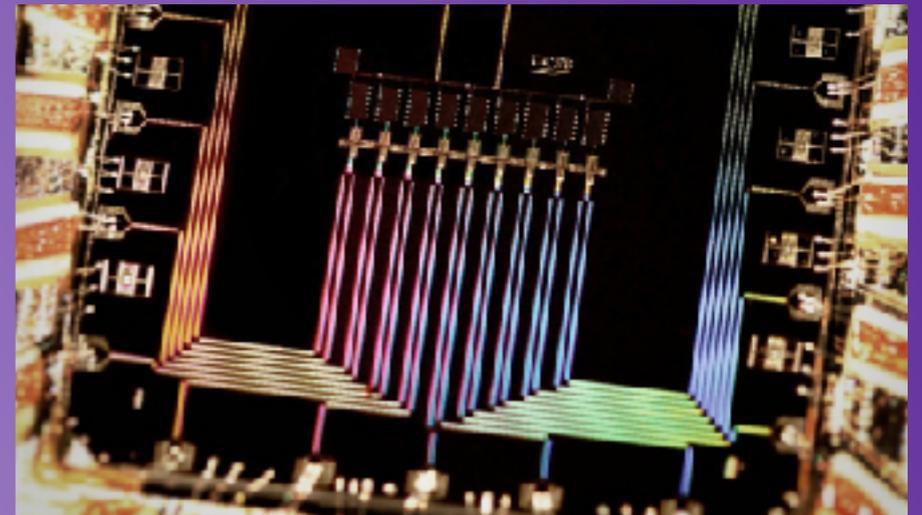
classical technologies

quantum technologies

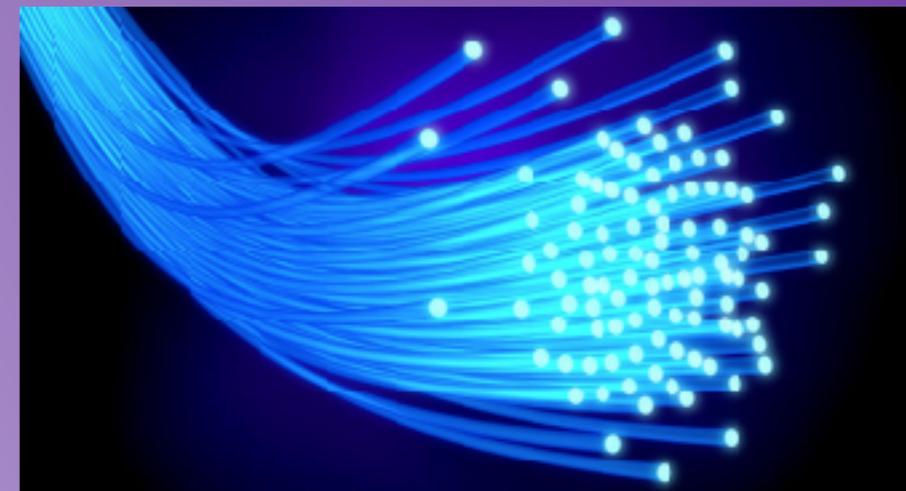
- state preparation
- info processing
- communication



superconducting quantum circuit



Martinis group
UCSB and Google (2015)

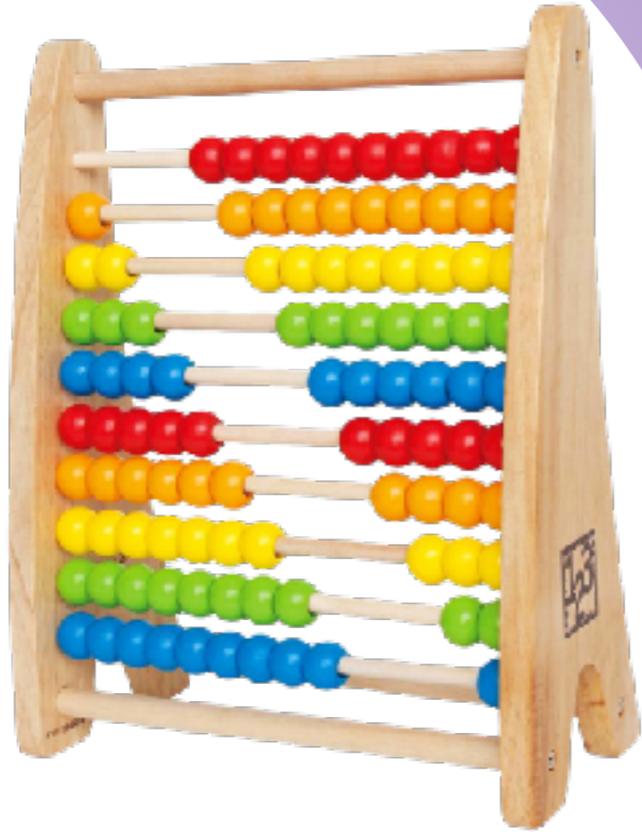


optical fiber

classical technologies

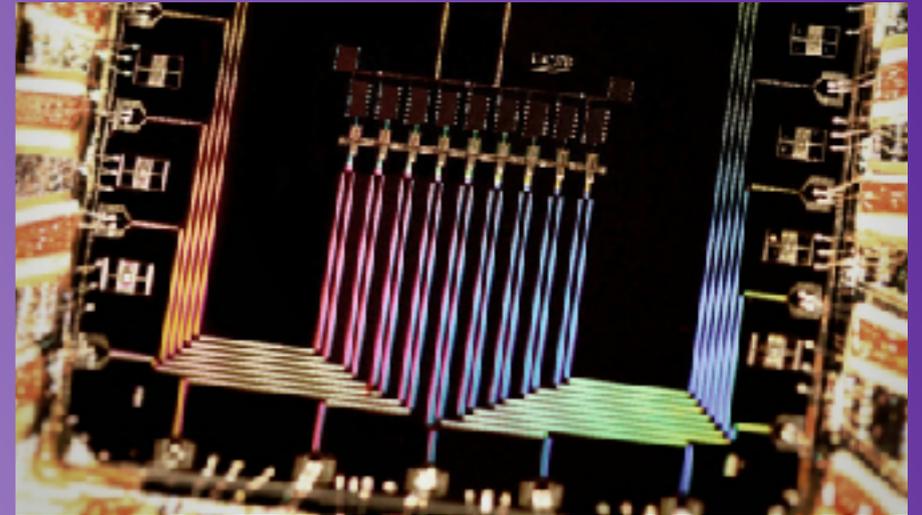
quantum technologies

- state preparation
- info processing
- communication

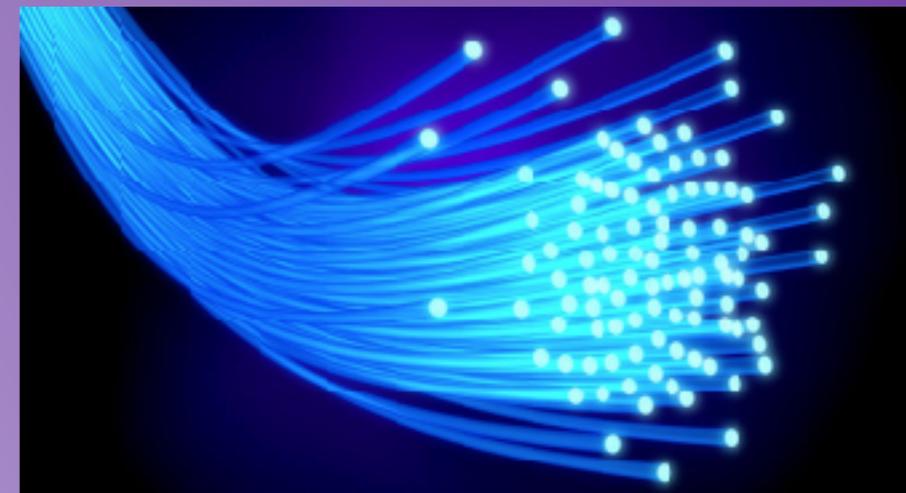


need
hybrid
systems

superconducting quantum circuit



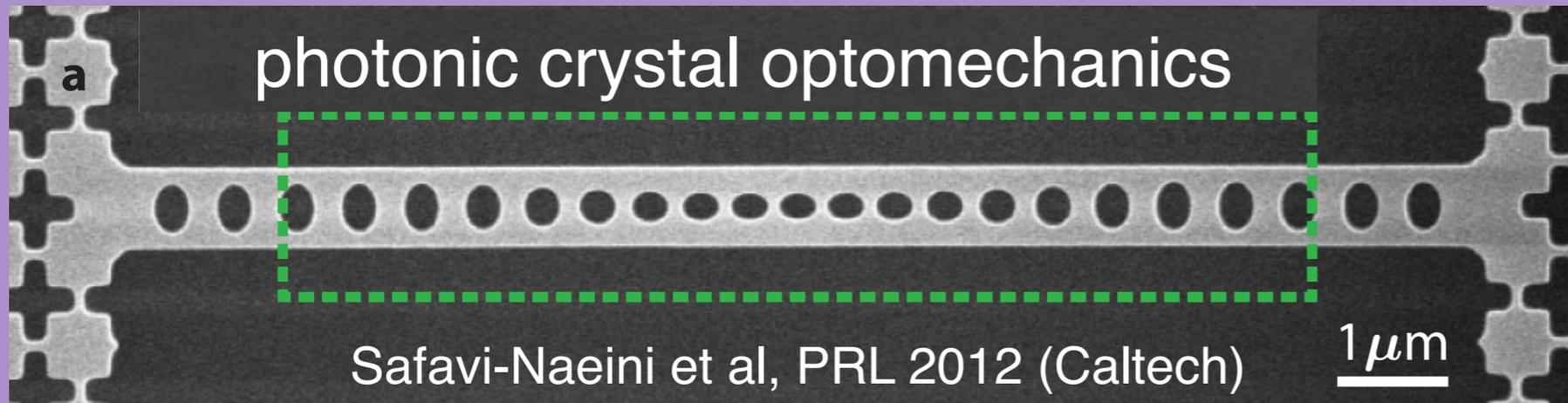
Martinis group
UCSB and Google (2015)



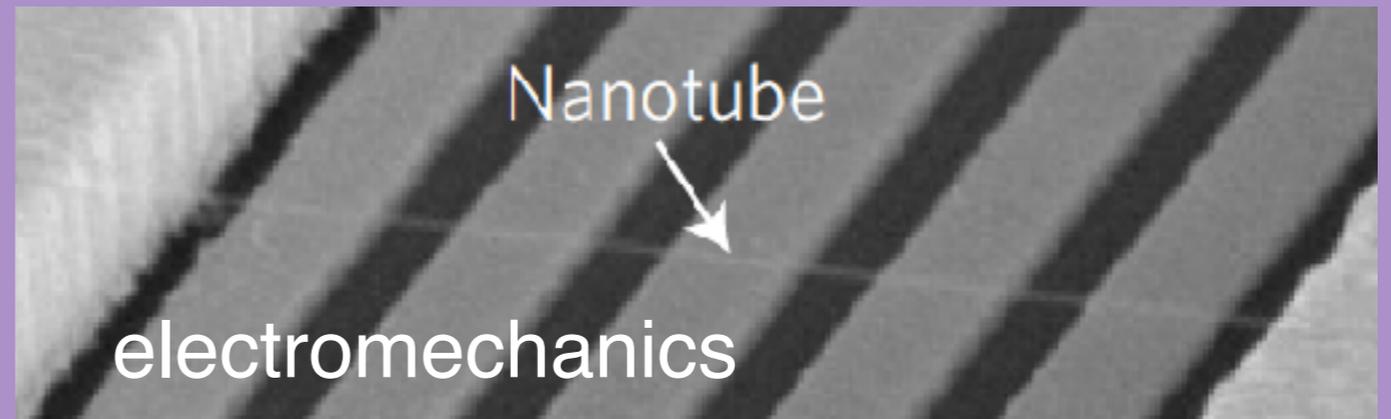
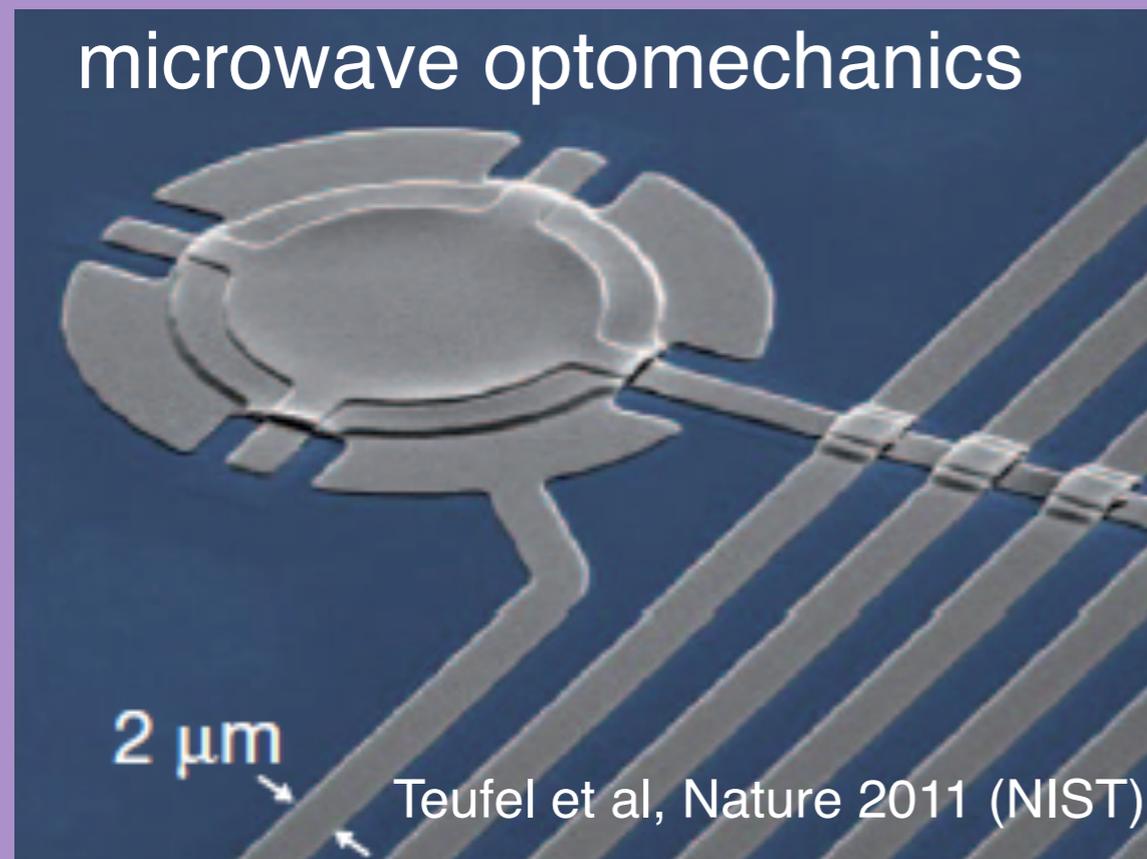
optical fiber

classical technologies

Hybrid Systems for Quantum Technologies



nano/micro scale
systems



Benyamini et al, Nature Physics 10, 151 (2014)



Osada et. al PRL 116, 223601 (2016)

use collective excitations

Optomagnonics



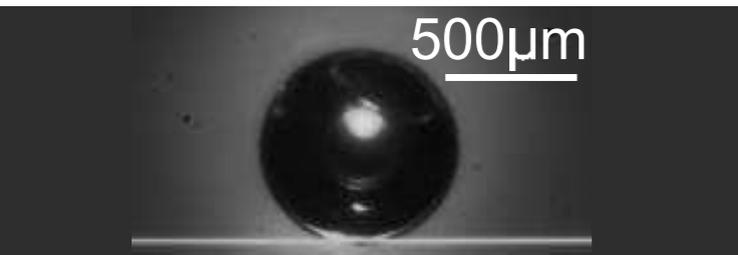
Picture from Tabuchi et al, PRL 113, 083603 (2014)



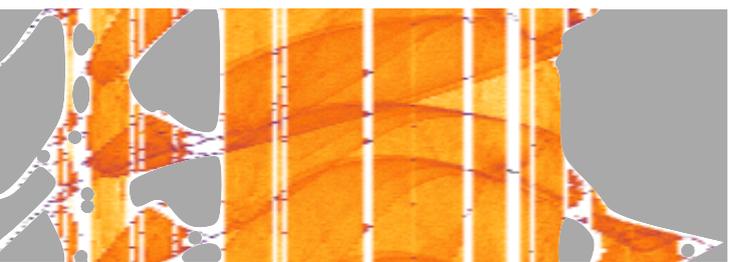
Magnons and the Kittel mode



Microwave regime



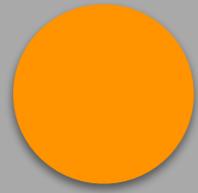
Optomagnonics



Optically induced spin dynamics



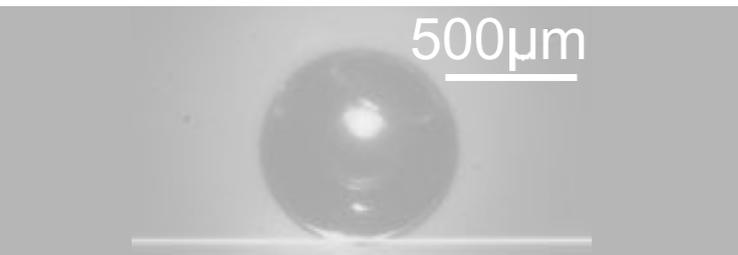
Outlook and Summary



Magnons and the Kittel mode



Microwave regime



Optomagnonics



Optically induced spin dynamics



Outlook and Summary

Magnonics



magnon

elementary magnetic
excitation
(quantum of spin wave)

Magnonics



magnon

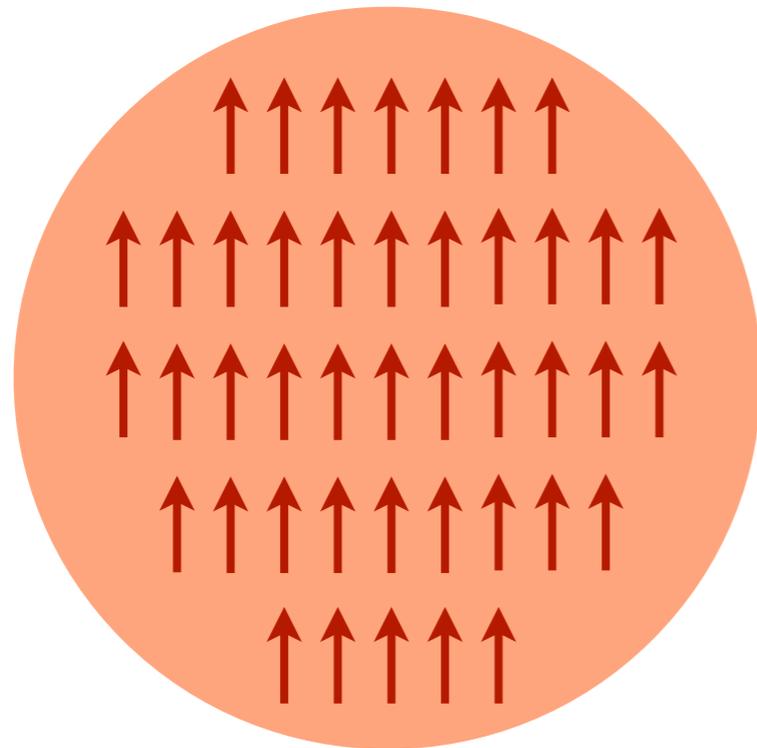
elementary magnetic
excitation
(quantum of spin wave)

Robust

Low Power

Tunable

Kittel mode



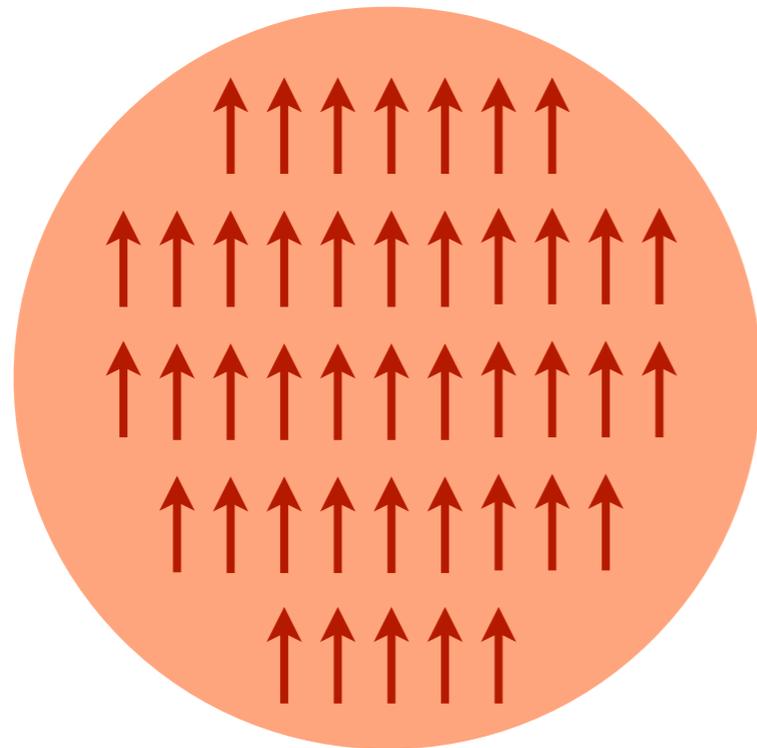
homogeneous
magnetic mode

$$\mathbf{M}(\mathbf{r}) = \mathbf{M}$$

spin wave with $k=0$

Magnonics

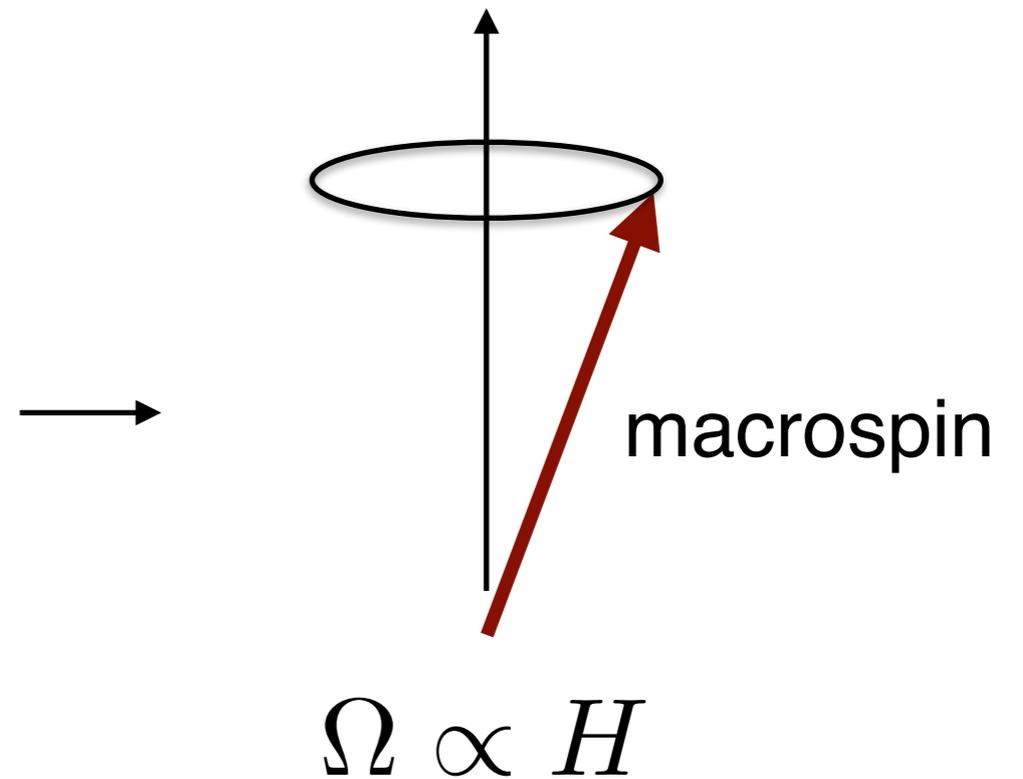
Kittel mode



homogeneous
magnetic mode

$$\mathbf{M}(\mathbf{r}) = \mathbf{M}$$

spin wave with $\mathbf{k} = 0$



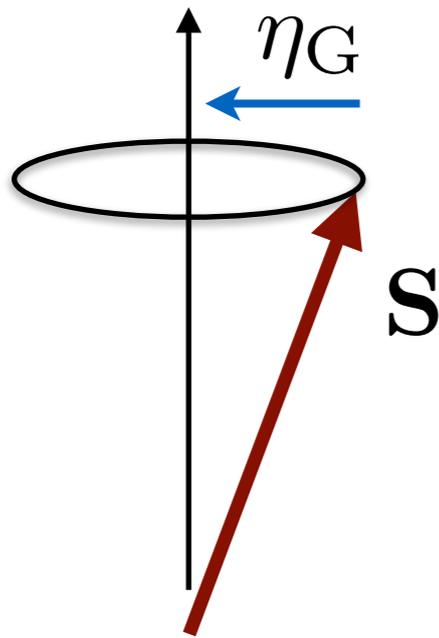
tunable precession frequency

$$\Omega \sim \text{GHz}$$

for 30mT

Landau-Lifschitz-Gilbert Equation

Dynamics of the macrospin



$$\Omega \propto H$$

precession frequency

$$\dot{\mathbf{S}} = -\Omega \mathbf{e}_z \times \mathbf{S} + \frac{\eta_G}{S} \left(\dot{\mathbf{S}} \times \mathbf{S} \right)$$

↓

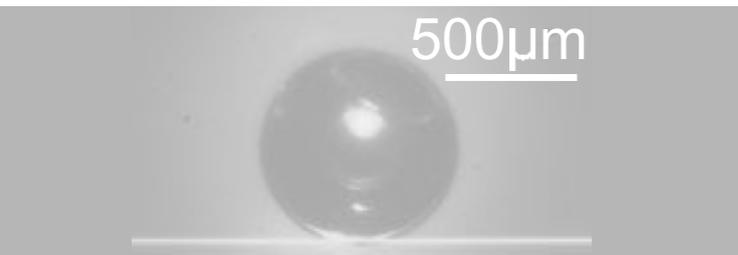
phenomenological damping term
(Gilbert damping)



Magnons and the Kittel mode



Microwave regime



Optomagnonics



Optically induced spin dynamics



Outlook and Summary

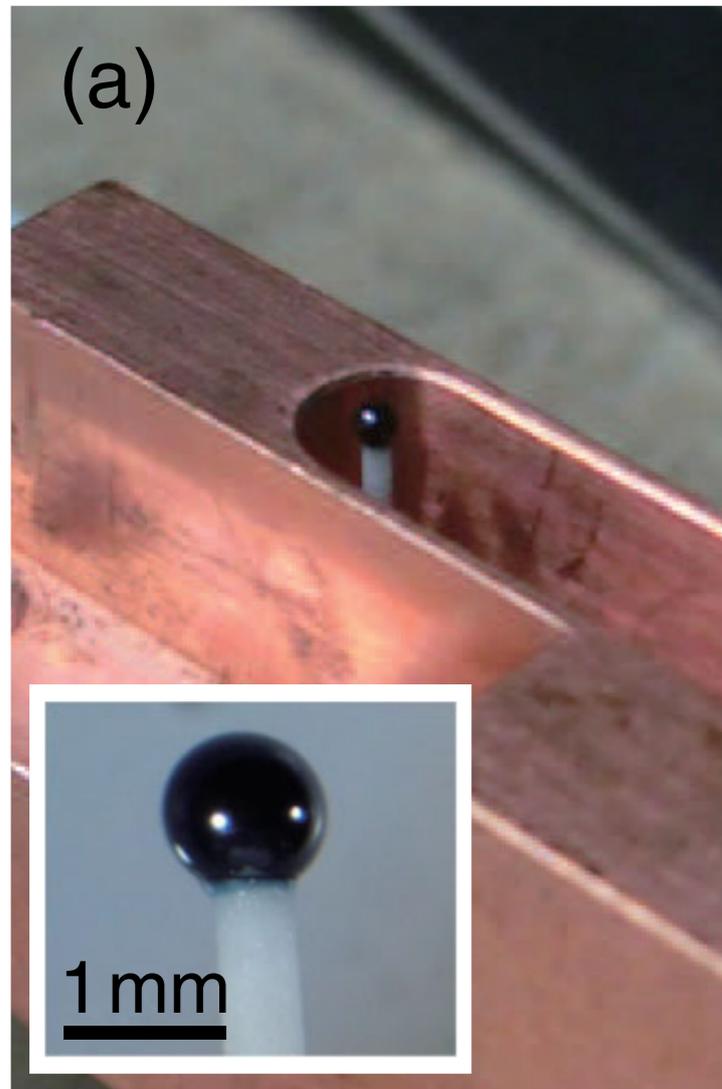
Microwave Regime

Magnons



Microwaves

Strong coupling demonstrated in 2014



- Tabuchi et. al PRL 113, 083603
(Nakamura's group, Tokyo)

- Zhang et. al PRL 113, 156401
(Hong Tang's group, Yale)

YIG



YIG

Yttrium Iron Garnet



- ferrimagnetic
- insulator
- transparent in the infrared

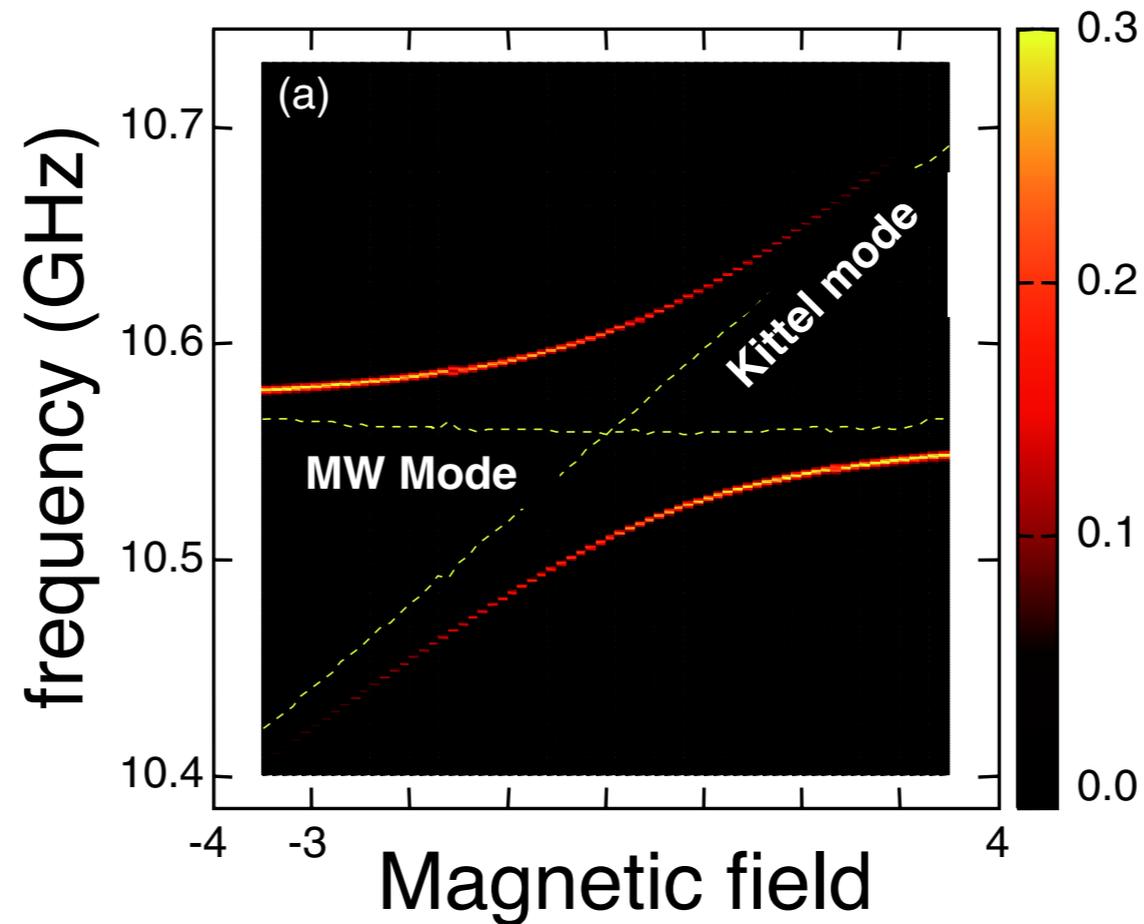
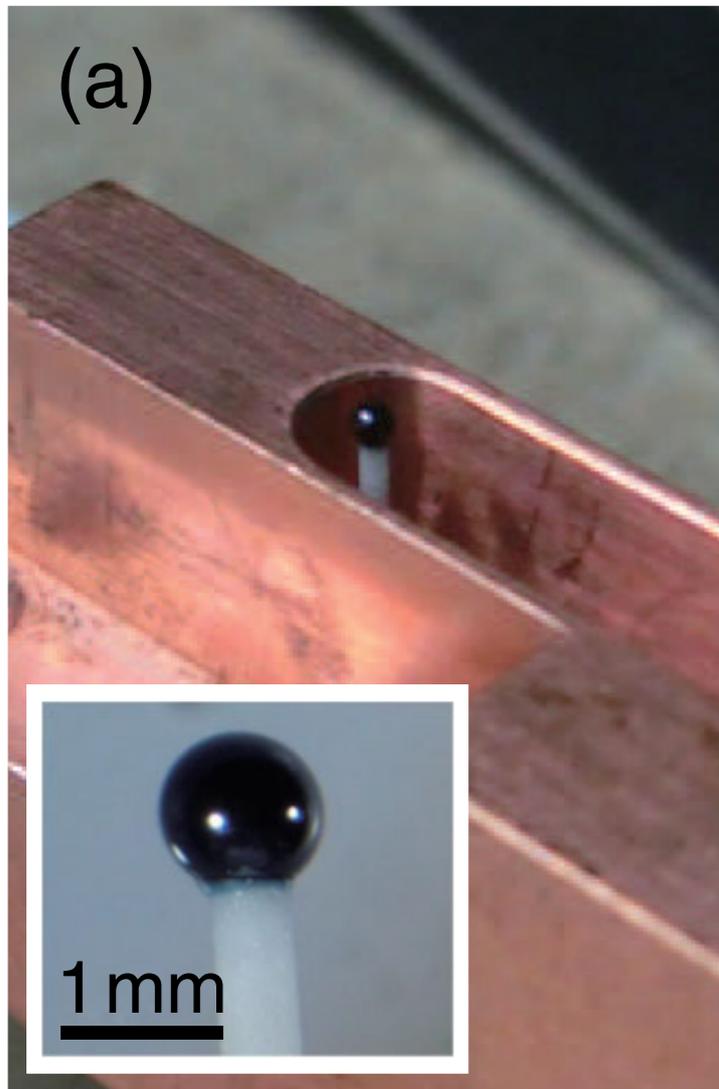
Microwave Regime

Magnons



Microwaves

Strong coupling demonstrated in 2014



- Tabuchi et. al PRL 113, 083603 (Nakamura's group, Tokyo)

- Zhang et. al PRL 113, 156401 (Hong Tang's group, Yale)

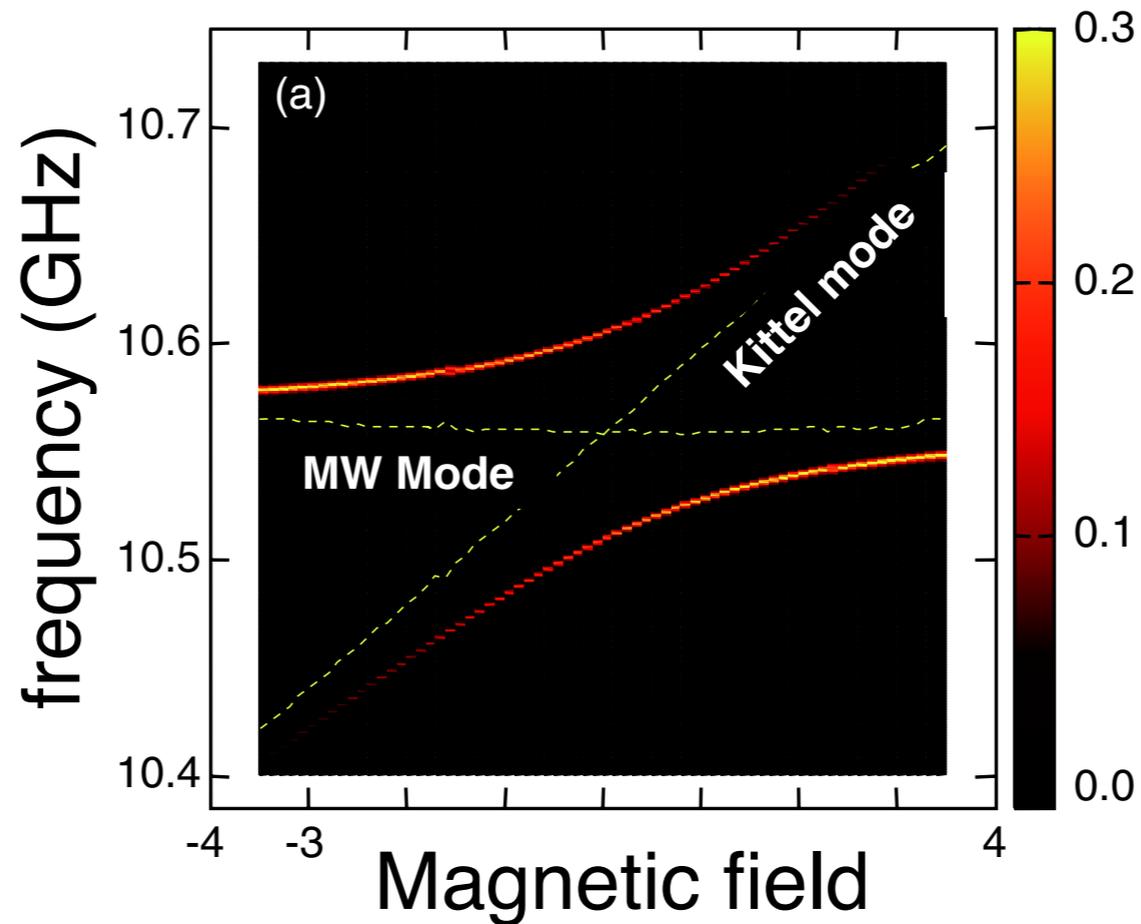
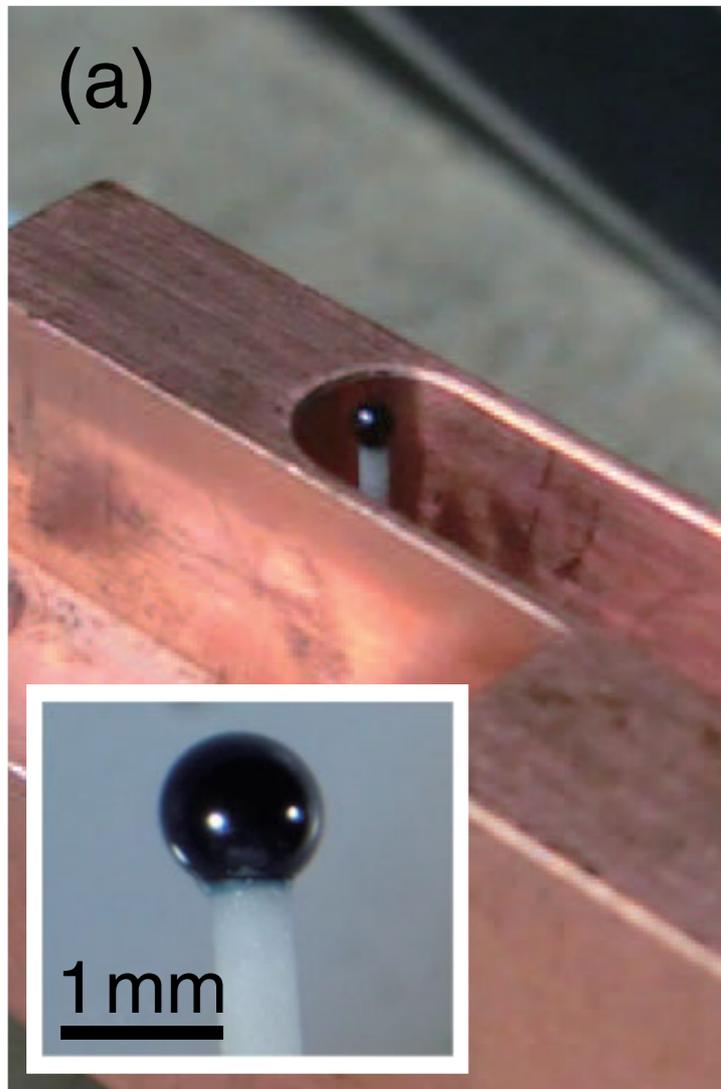
Microwave Regime

Magnons



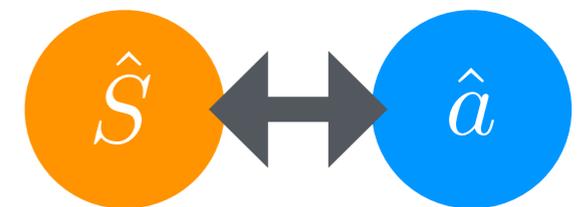
Microwaves

Strong coupling demonstrated in 2014



Resonant coupling

$$\hat{S}^+ \hat{a} + \hat{S}^- \hat{a}^\dagger$$



$\sim 50\text{MHz}$

Cooperativity

$$\mathcal{C} = 3 \times 10^3$$

- Tabuchi et. al PRL 113, 083603 (Nakamura's group, Tokyo)

- Zhang et. al PRL 113, 156401 (Hong Tang's group, Yale)

Microwave Regime

Magnons



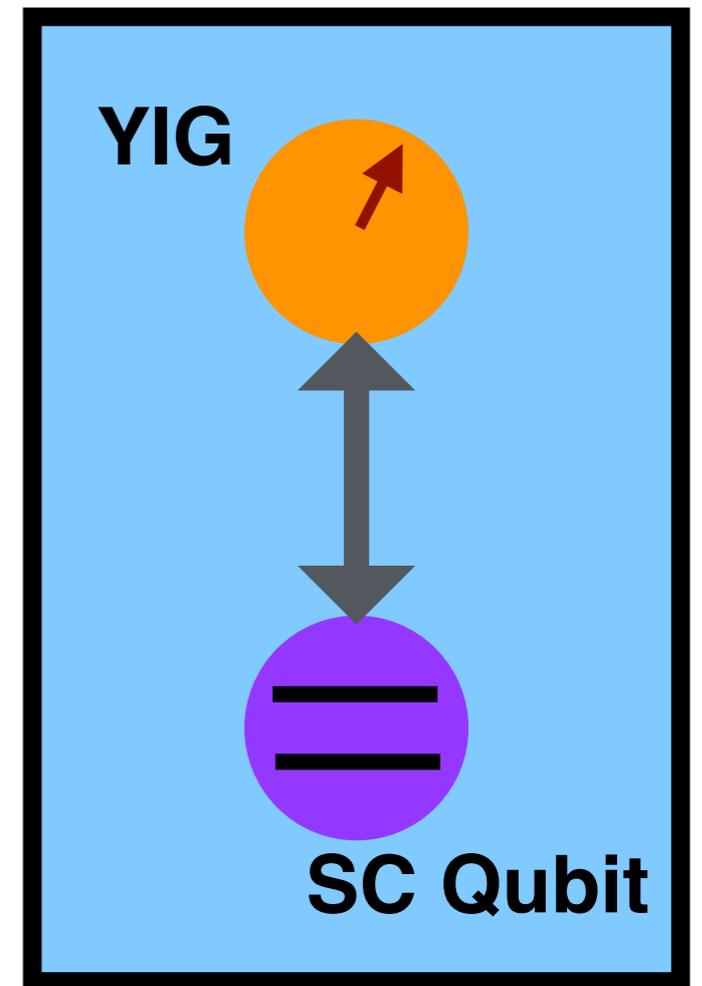
Microwaves

QUANTUM INFORMATION

(Science 2015)

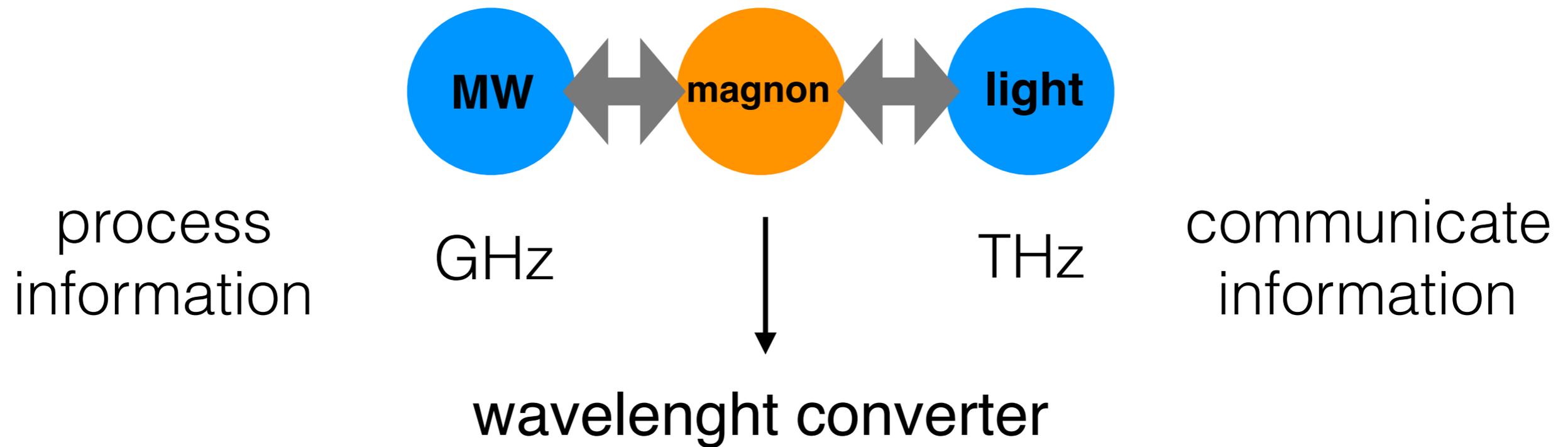
Coherent coupling between a ferromagnetic magnon and a superconducting qubit

Yutaka Tabuchi,^{1*} Seiichiro Ishino,¹ Atsushi Noguchi,¹ Toyofumi Ishikawa,¹
Rekishu Yamazaki,¹ Koji Usami,¹ Yasunobu Nakamura^{1,2}



MW Cavity

Coupling to Optics?



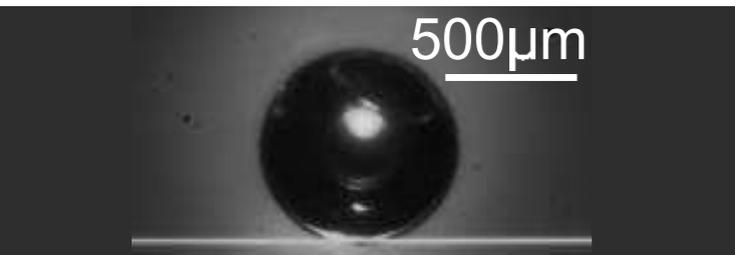
**Motivation:
magnon as a transducer**



Magnons and the Kittel mode



Microwave regime



Optomagnonics



Optically induced spin dynamics

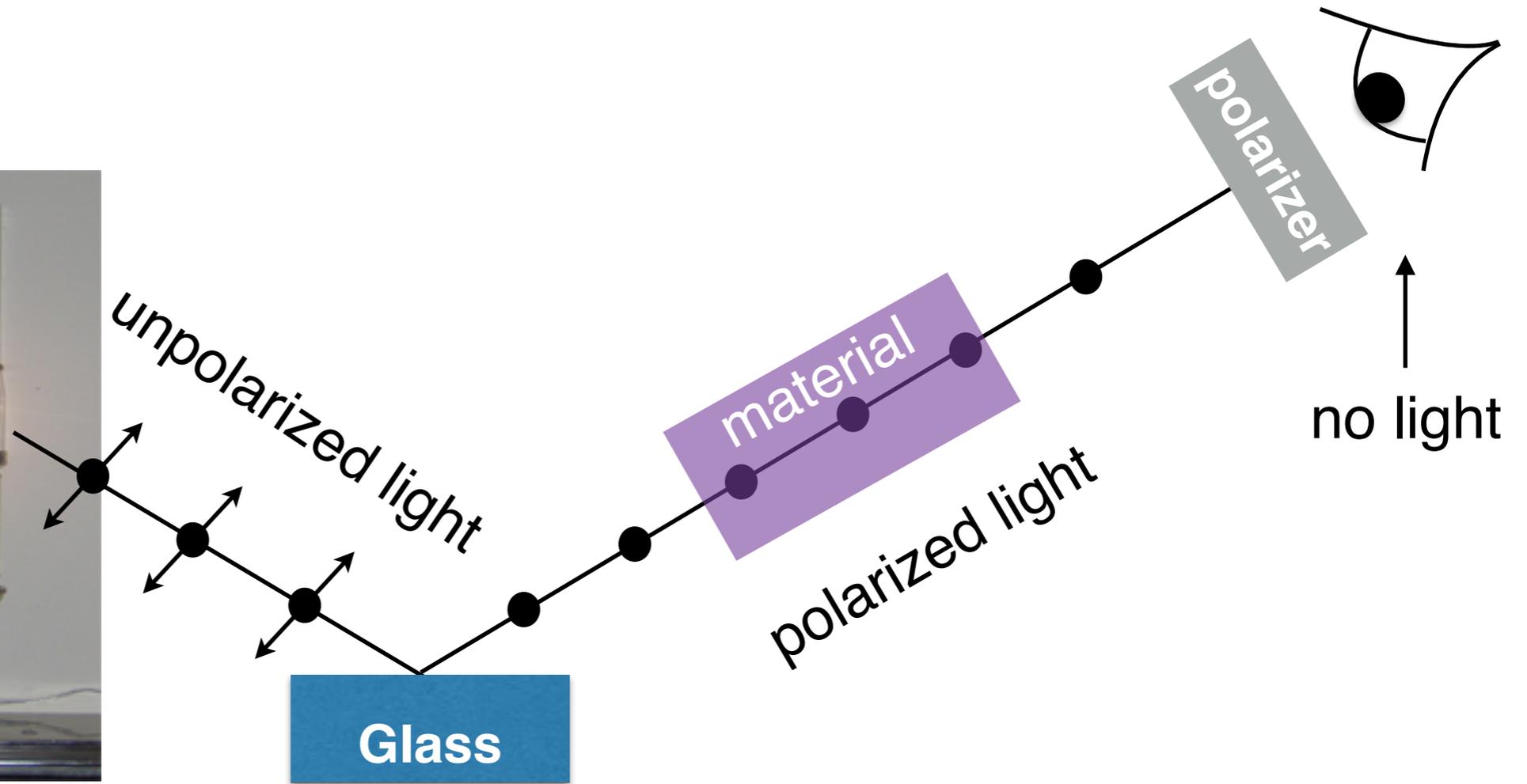


Outlook and Summary

Faraday Effect (1846)



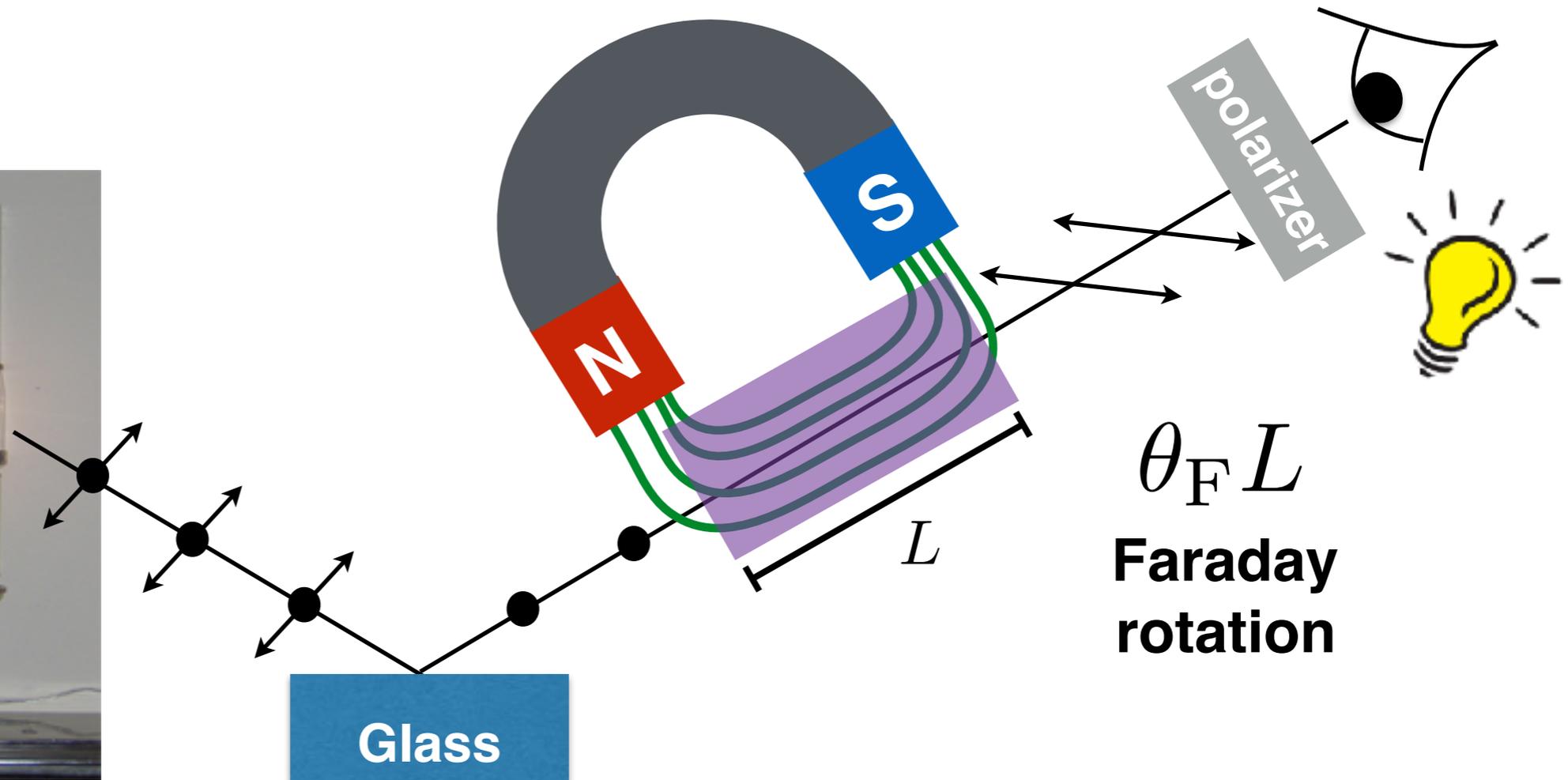
Oil Lamp



Faraday Effect (1846)

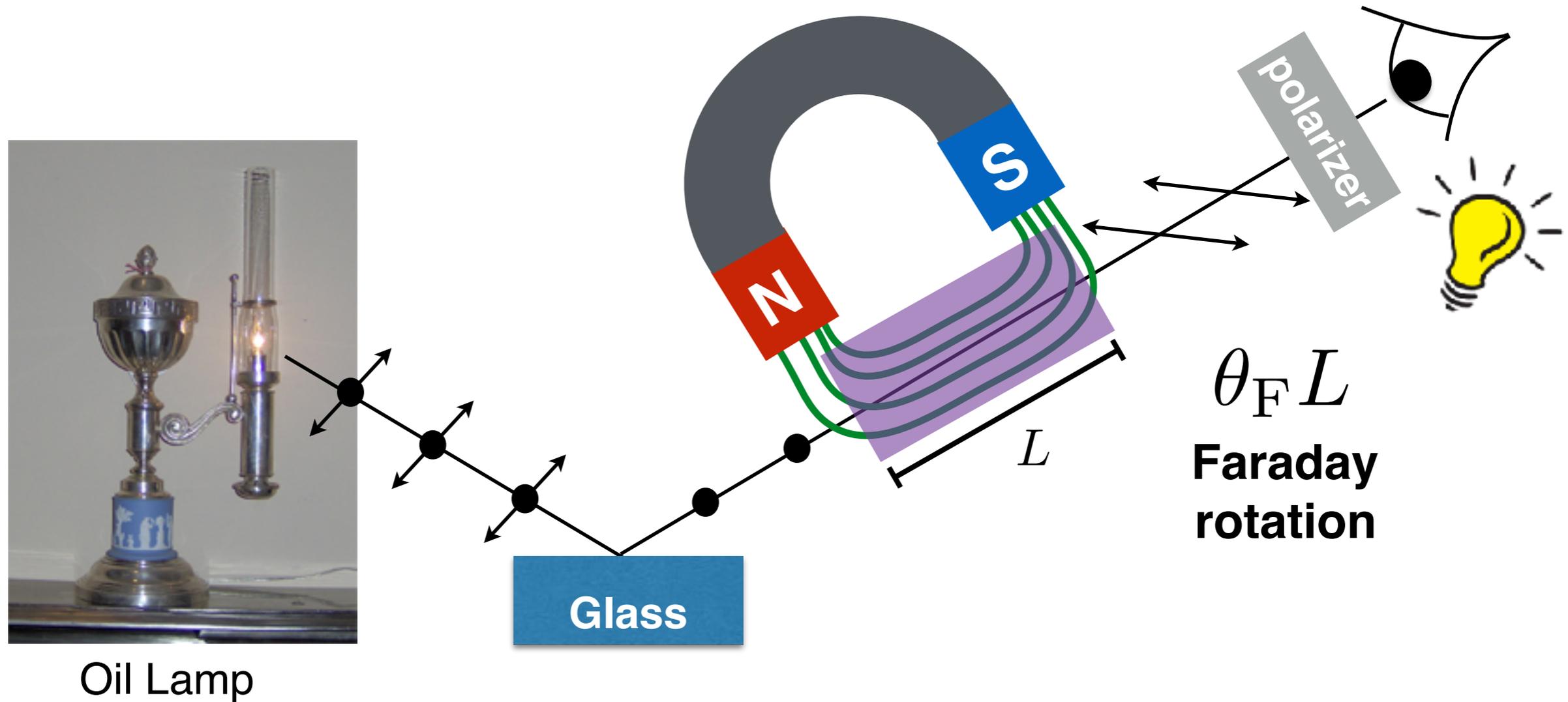


Oil Lamp



$\theta_F L$
Faraday
rotation

Faraday Effect (1846)



RELATION OF LIGHT TO THE MAGNETIC FORCE.

15

¶ iii. *General considerations.*

2221. Thus is established, I think for the first time*, a true, direct relation and dependence between light and the magnetic and electric forces; and thus a great

Optomagnonic Hamiltonian

Faraday rotation

$$\bar{U}_{\text{MO}} = \theta_{\text{F}} \sqrt{\frac{\epsilon}{\epsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_{\text{S}}} \cdot \frac{\epsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$

optical spin density 

magnetization density 

Optomagnonic Hamiltonian

$$\bar{U}_{\text{MO}} = \theta_{\text{F}} \sqrt{\frac{\varepsilon}{\varepsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_{\text{s}}} \cdot \frac{\varepsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$

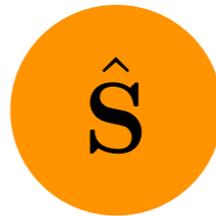
Quantize:



Optomagnonic Hamiltonian

$$\bar{U}_{\text{MO}} = \theta_{\text{F}} \sqrt{\frac{\varepsilon}{\varepsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_{\text{s}}} \cdot \frac{\varepsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$

Quantize:



two-photon process

Optomagnonic Hamiltonian

$$\bar{U}_{\text{MO}} = \theta_{\text{F}} \sqrt{\frac{\epsilon}{\epsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_{\text{s}}} \cdot \frac{\epsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$

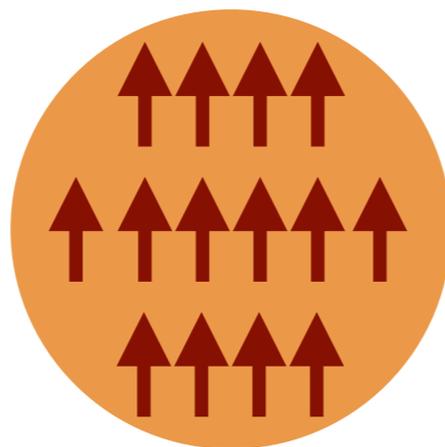
Quantize:



two-photon process

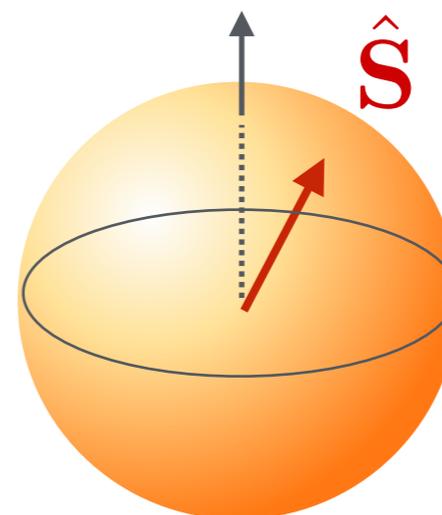
$$\mathbf{M}(\mathbf{r}) = \mathbf{M}$$

Kittel mode



$$\Omega \propto H$$

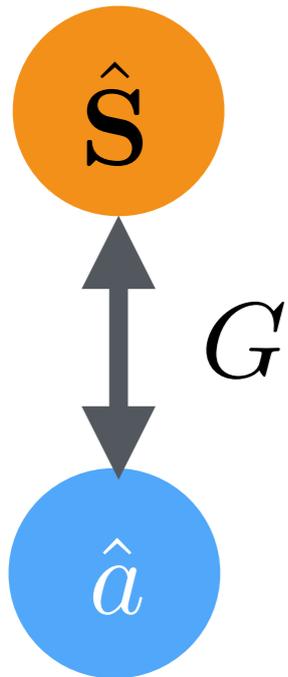
Bloch sphere



Optomagnonic Hamiltonian

Microscopic Hamiltonian

Parametric
coupling



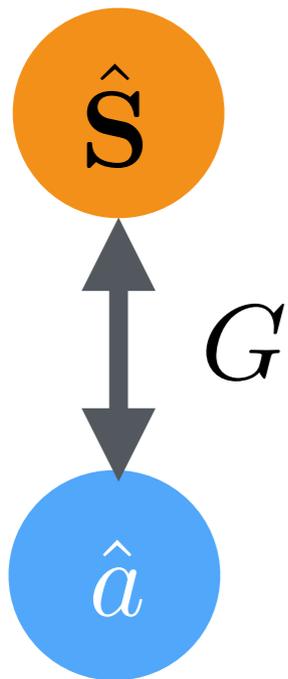
$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

Optomagnonic Hamiltonian

Microscopic Hamiltonian

$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

Parametric
coupling



Optomagnonic coupling

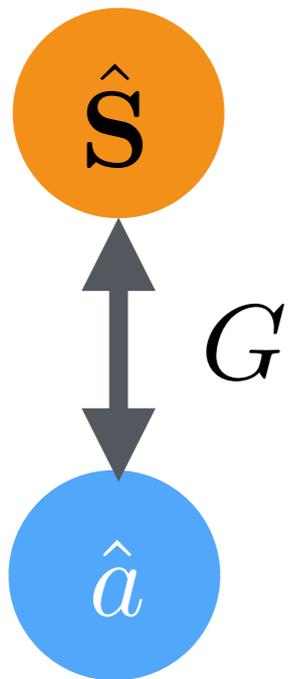
$$G_{\beta\gamma}^j = -i \frac{\theta_F \lambda}{2\pi \hbar S} \frac{\epsilon_0 \epsilon}{2} \epsilon_{jmn} \int d\mathbf{r} E_{\beta m}^*(\mathbf{r}) E_{\gamma n}(\mathbf{r})$$

Optomagnonic Hamiltonian

Microscopic Hamiltonian

$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

Parametric
coupling



Optomagnonic coupling

$$G_{\beta\gamma}^j = -i \frac{\theta_F \lambda}{2\pi \hbar S} \frac{\epsilon_0 \epsilon}{2} \epsilon_{jmn} \int d\mathbf{r} E_{\beta m}^*(\mathbf{r}) E_{\gamma n}(\mathbf{r})$$

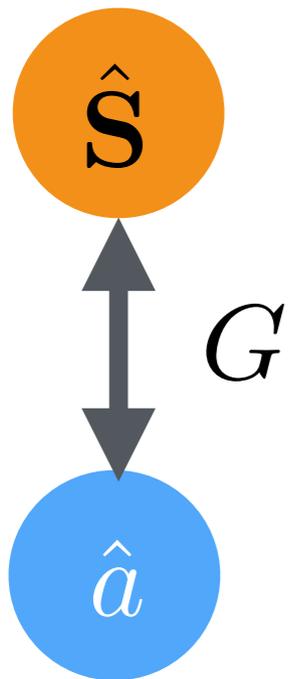
overlap electric field
mode functions

Optomagnonic Hamiltonian

Microscopic Hamiltonian

$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

Parametric
coupling



Optomagnonic coupling

$$G_{\beta\gamma}^j = \left(i \frac{\theta_F \lambda}{2\pi \hbar S} \frac{\epsilon_0 \epsilon}{2} \epsilon_{jmn} \int d\mathbf{r} E_{\beta m}^*(\mathbf{r}) E_{\gamma n}(\mathbf{r}) \right)$$

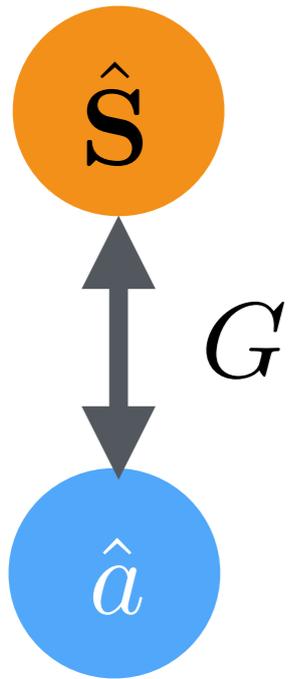
Faraday rotation

Optomagnonic Hamiltonian

Microscopic Hamiltonian

$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

Parametric
coupling

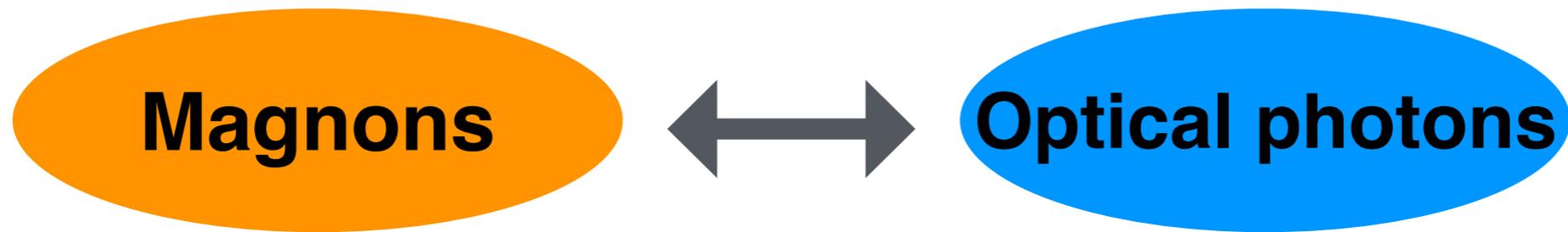


Optomagnonic coupling

$$G_{\beta\gamma}^j = -i \frac{\theta_F \lambda}{2\pi \hbar S} \frac{\epsilon_0 \epsilon}{2} \epsilon_{jmn} \int d\mathbf{r} E_{\beta m}^*(\mathbf{r}) E_{\gamma n}(\mathbf{r})$$

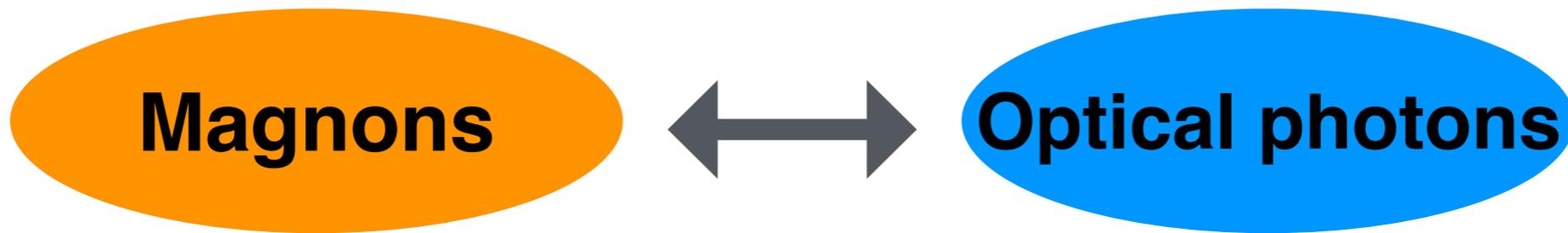
number of spins

Cavity Optomagnonics

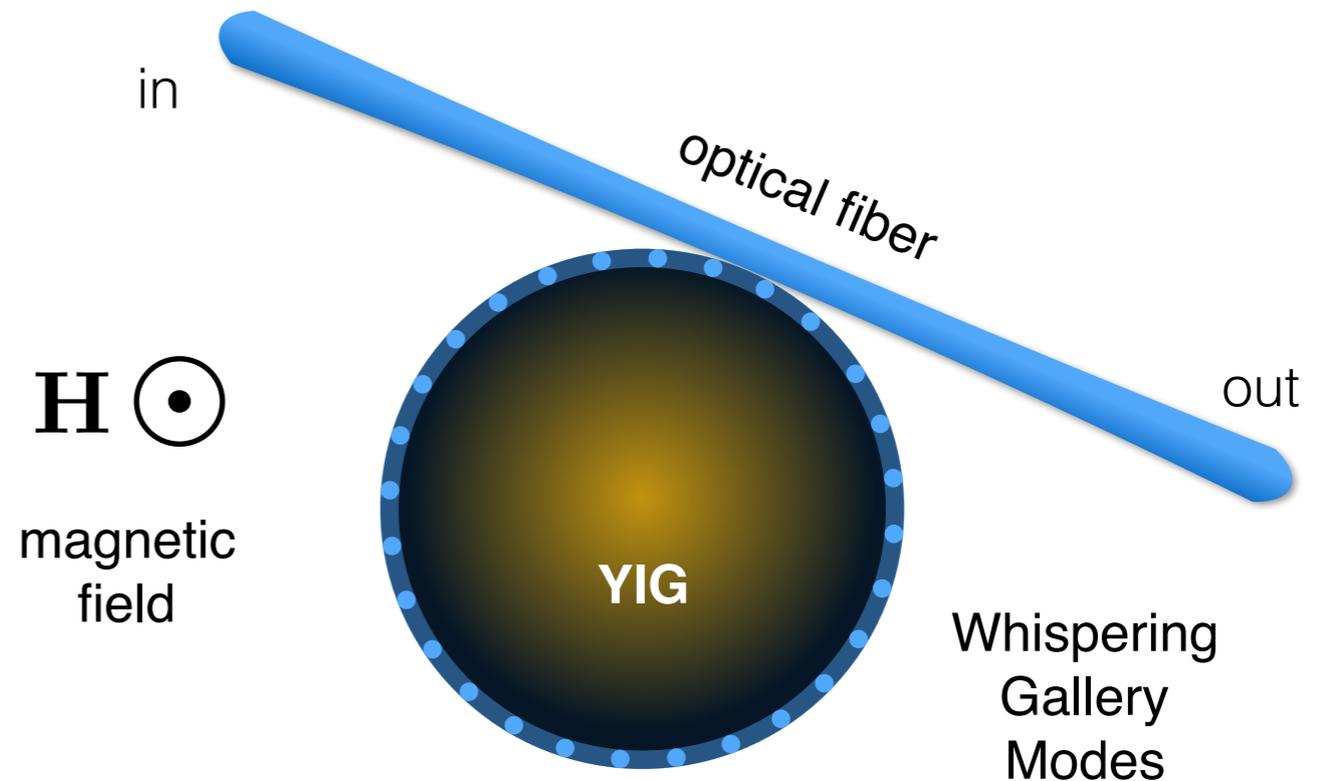
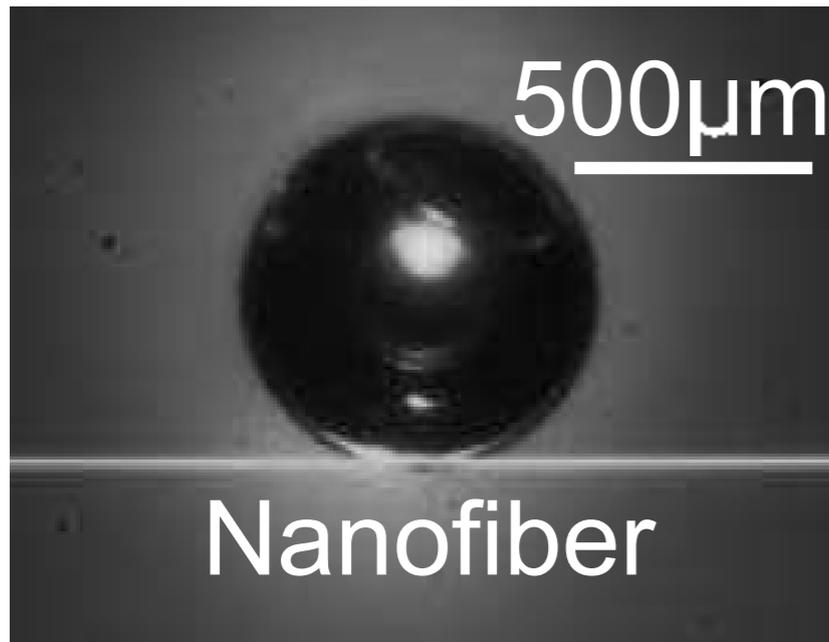


Coupling demonstrated in 2016

Cavity Optomagnonics



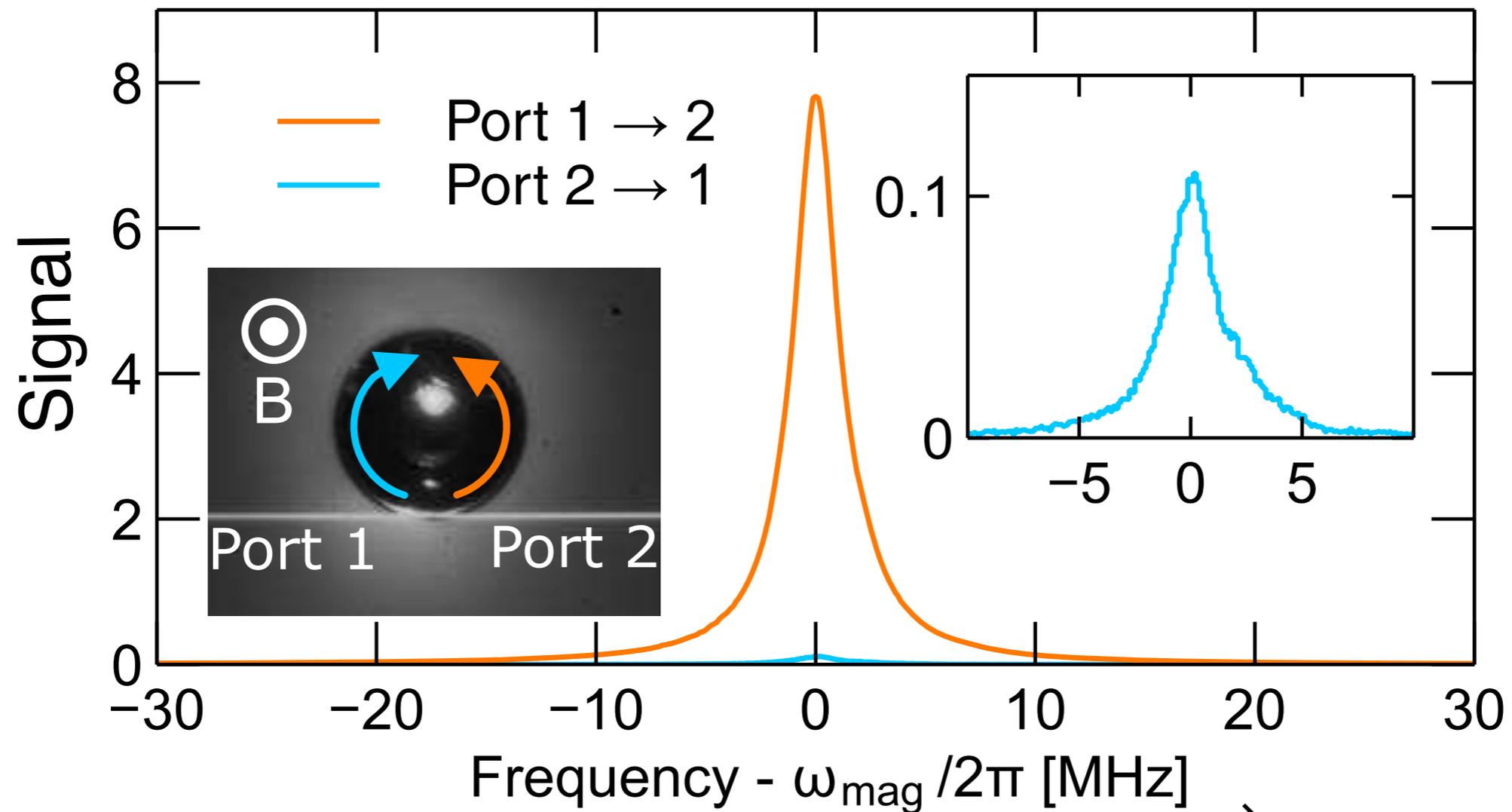
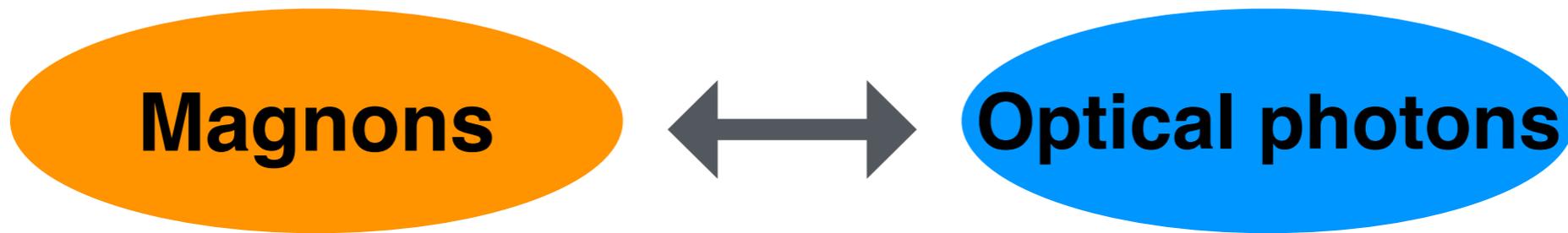
Coupling demonstrated in 2016



A cavity enhances the effect

- Osada et. al PRL 116, 223601 (Nakamura's group, Tokyo)
- Haigh et. al PRL 117, 133602 (Ferguson's group, Cambridge)
- Zhang et. al PRL 117, 123605 (Hong Tang's group, Yale)

Cavity Optomagnonics



- Osada et. al PRL 116, 223601
(Nakamura's group, Tokyo)

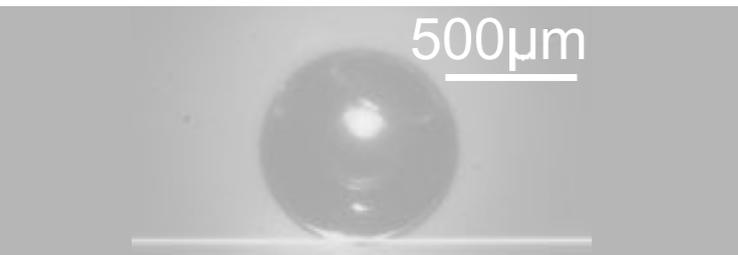
Sidebands at the magnon frequency



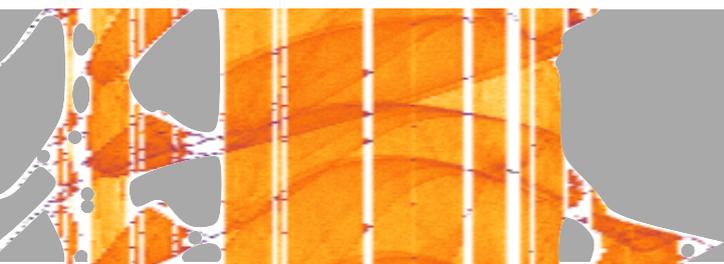
Magnons and the Kittel mode



Microwave regime



Optomagnonics

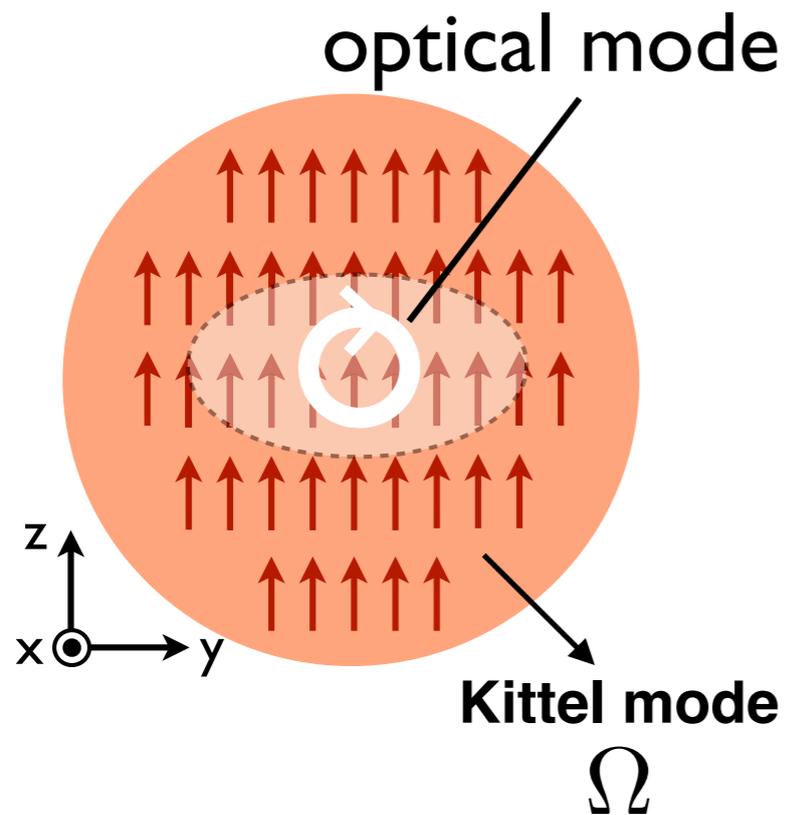


Optically induced spin dynamics



Outlook and Summary

Cavity Optomagnonics: 1 optical mode

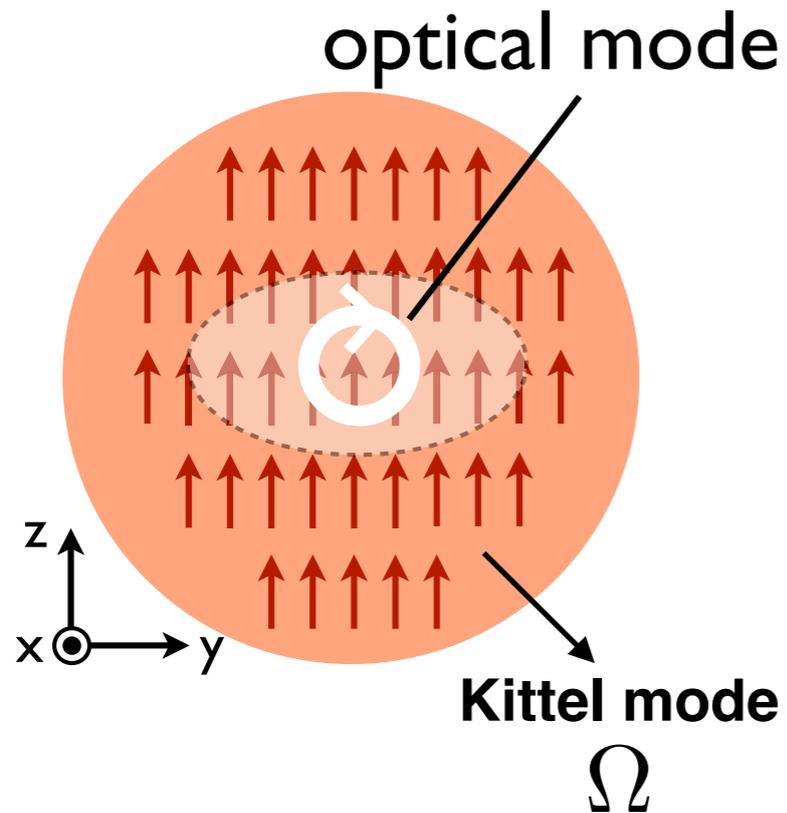


Coupling

$$\hat{H}_{MO} = \hbar \sum_{j\beta\gamma} \hat{S}_j G_{\beta\gamma}^j \hat{a}_\beta^\dagger \hat{a}_\gamma$$

acquires a simple form

Cavity Optomagnonics: 1 optical mode

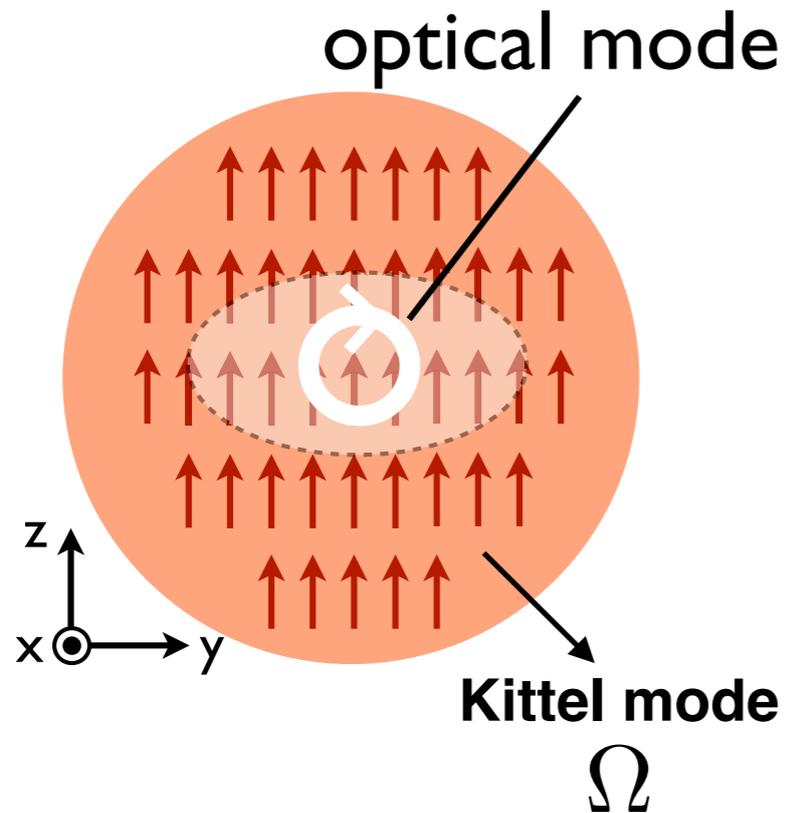


Total Hamiltonian for one optical mode

$$H = -\hbar\Delta\hat{a}^\dagger\hat{a} - \hbar\Omega\hat{S}_z + \hbar G\hat{S}_x\hat{a}^\dagger\hat{a}$$

driving laser detuning $\Delta = \omega_{las} - \omega_{cav}$

Cavity Optomagnonics



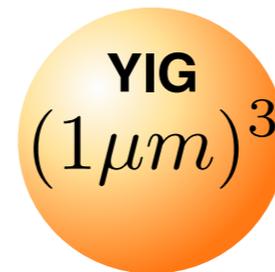
Total Hamiltonian for one optical mode

$$H = -\hbar\Delta\hat{a}^\dagger\hat{a} - \hbar\Omega\hat{S}_z + \hbar G\hat{S}_x\hat{a}^\dagger\hat{a}$$

driving laser detuning $\Delta = \omega_{las} - \omega_{cav}$

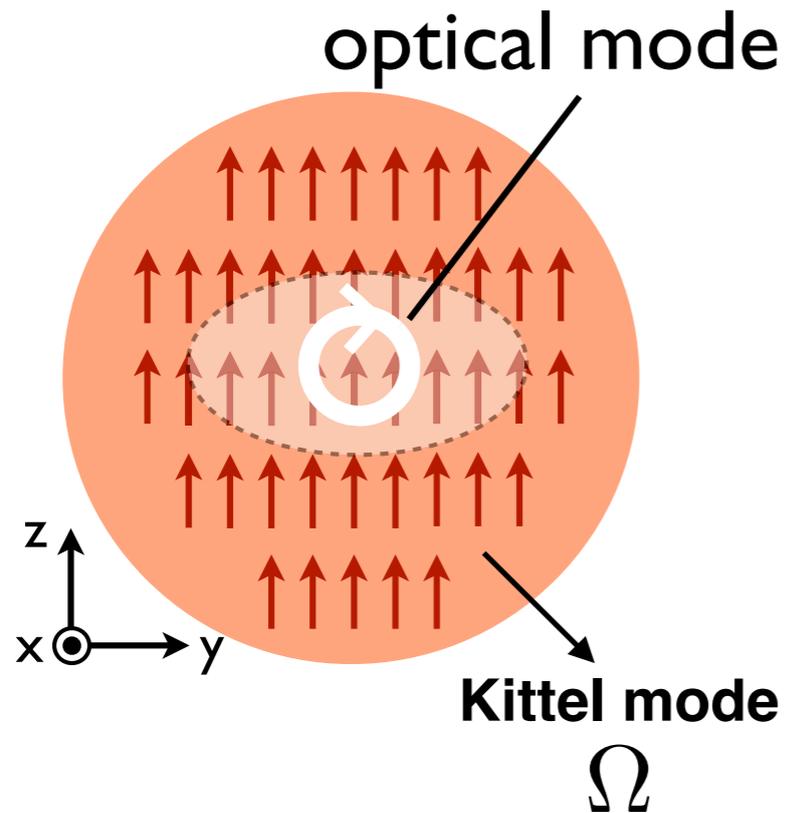
$$G = \frac{1}{S} \frac{c\theta_F}{4\sqrt{\epsilon}} \xi$$

mode overlap factor



$$G \approx 1 \text{ Hz}$$

Cavity Optomagnonics



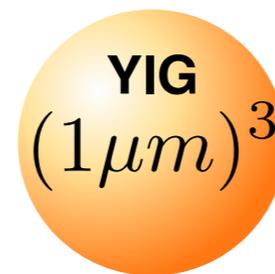
Total Hamiltonian for one optical mode

$$H = -\hbar\Delta\hat{a}^\dagger\hat{a} - \hbar\Omega\hat{S}_z + \hbar G\hat{S}_x\hat{a}^\dagger\hat{a}$$

driving laser detuning $\Delta = \omega_{las} - \omega_{cav}$

$$G = \frac{1}{S} \frac{c\theta_F}{4\sqrt{\epsilon}} \xi$$

mode overlap factor



$$G \approx 1\text{Hz}$$

Optical magnetic field density

$$b_{\text{opt}} \sim \frac{10^{-11}\text{T}}{\text{photon}/(\mu\text{m})^3}$$

Cavity Optomagnonics

Classical Equation of Motion

Cavity decay rate

initial light amplitude

$$\dot{a} = -i(GS_x - \Delta)a - \frac{\kappa}{2}(a - \alpha_{\max})$$
$$\dot{\mathbf{S}} = (Ga^*a \mathbf{e}_x - \Omega \mathbf{e}_z) \times \mathbf{S} + \frac{\eta_G}{S}(\dot{\mathbf{S}} \times \mathbf{S})$$

Fast Cavity Limit

$\kappa \gg \Omega$ integrate out the light field

Effective equation of motion for \mathbf{S} :

$$\dot{\mathbf{S}} = \mathbf{B}_{\text{eff}} \times \mathbf{S} + \frac{\eta_{\text{opt}}}{S} \left(\dot{S}_x \mathbf{e}_x \times \mathbf{S} \right)$$

Fast Cavity Limit

$\kappa \gg \Omega$ integrate out the light field

Effective equation of motion for \mathbf{S} :

$$\dot{\mathbf{S}} = \mathbf{B}_{\text{eff}} \times \mathbf{S} + \frac{\eta_{\text{opt}}}{S} \left(\dot{S}_x \mathbf{e}_x \times \mathbf{S} \right)$$

optically induced

effective field

$$\mathbf{B}_{\text{eff}} = -\Omega \mathbf{e}_z + \mathbf{B}_{\text{opt}} \quad \mathbf{B}_{\text{opt}} = \frac{G}{\left[\left(\frac{\kappa}{2} \right)^2 + (\Delta - GS_x)^2 \right]} \left(\frac{\kappa}{2} \alpha_{\text{max}} \right)^2 \mathbf{e}_x$$

tunable by the external laser drive

Fast Cavity Limit

$\kappa \gg \Omega$ integrate out the light field

Effective equation of motion for \mathbf{S} :

$$\dot{\mathbf{S}} = \mathbf{B}_{\text{eff}} \times \mathbf{S} + \frac{\eta_{\text{opt}}}{S} \left(\dot{S}_x \mathbf{e}_x \times \mathbf{S} \right)$$

optically induced

effective field

$$\mathbf{B}_{\text{eff}} = -\Omega \mathbf{e}_z + \mathbf{B}_{\text{opt}}$$

$$\mathbf{B}_{\text{opt}} = \frac{G}{\left[\left(\frac{\kappa}{2} \right)^2 + (\Delta - GS_x)^2 \right]} \left(\frac{\kappa}{2} \alpha_{\text{max}} \right)^2 \mathbf{e}_x$$

**damping
can change sign**

$$\eta_{\text{opt}} = -2G\kappa S |\mathbf{B}_{\text{opt}}| \frac{(\Delta - GS_x)}{\left[\left(\frac{\kappa}{2} \right)^2 + (\Delta - GS_x)^2 \right]^2}$$

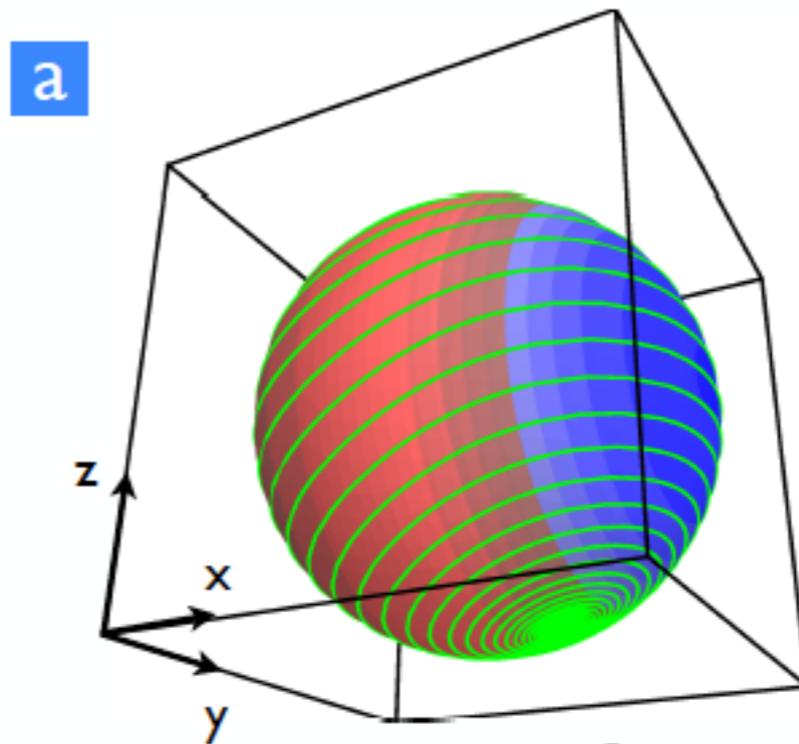
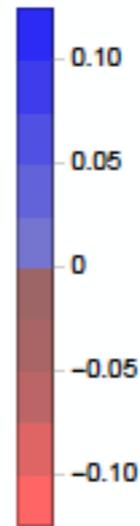
tunable by the external laser drive

Fast Cavity Limit: Spin Dynamics

magnetic
switching

$$G\alpha_{\max}^2/\Omega = 0.6$$

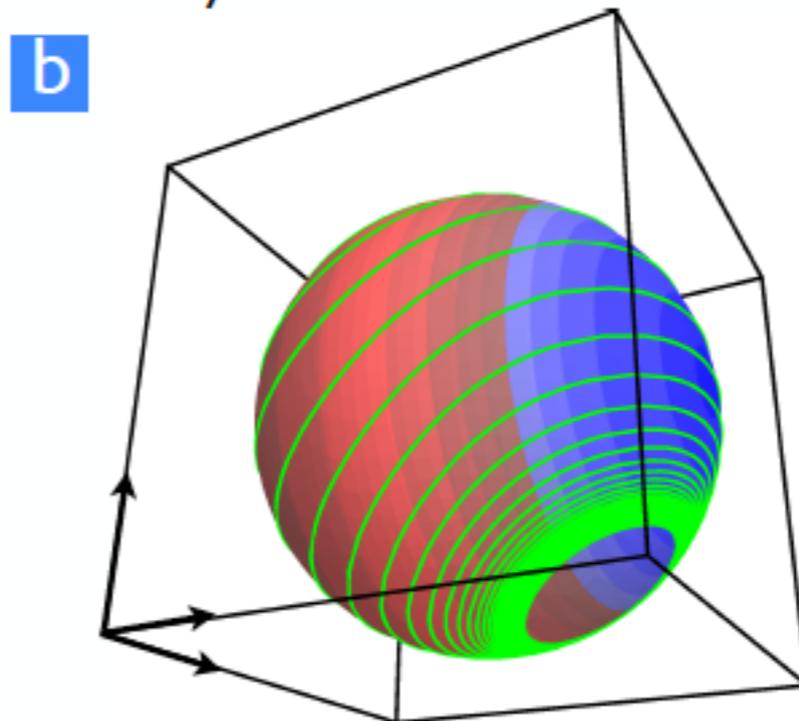
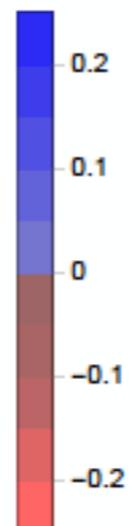
η_{opt}



self-sustained
oscillations

$$G\alpha_{\max}^2/\Omega = 0.8$$

η_{opt}



Increasing light intensity

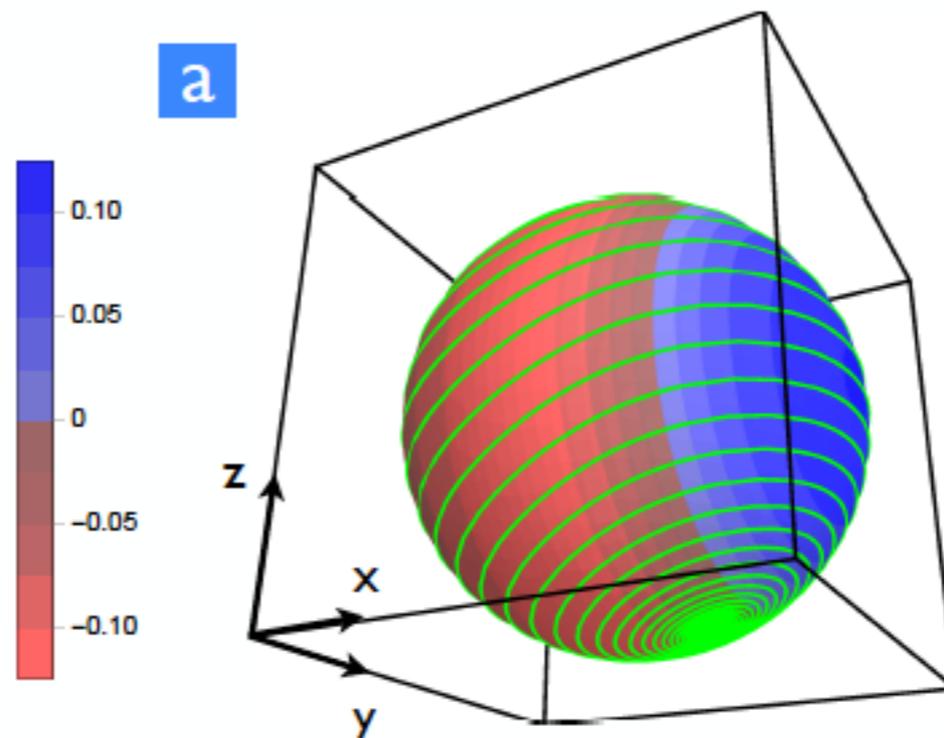


$$\Delta = \Omega, GS/\Omega = 2, \kappa/\Omega = 5$$

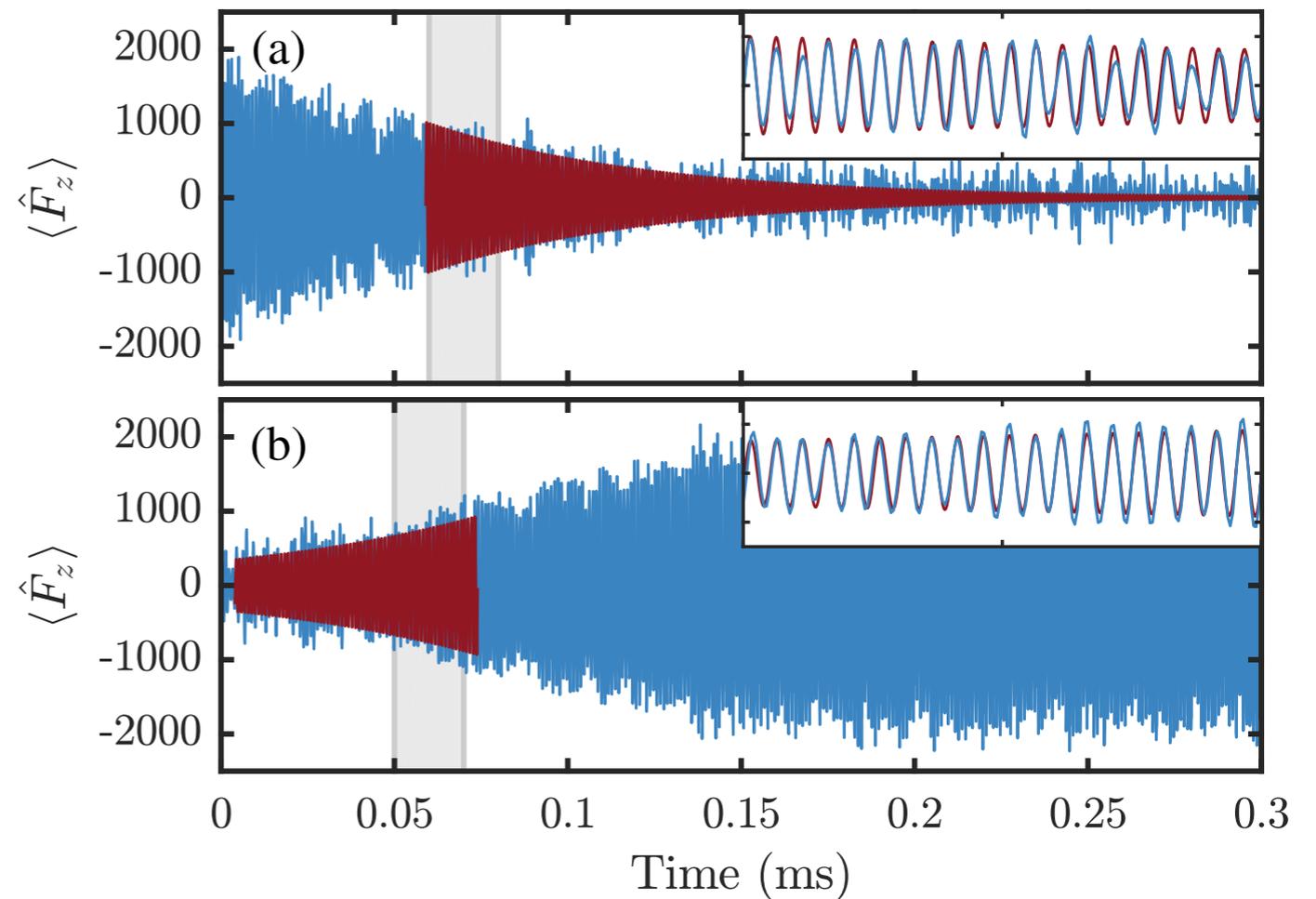
Fast Cavity Limit: Spin Dynamics

magnetic
switching

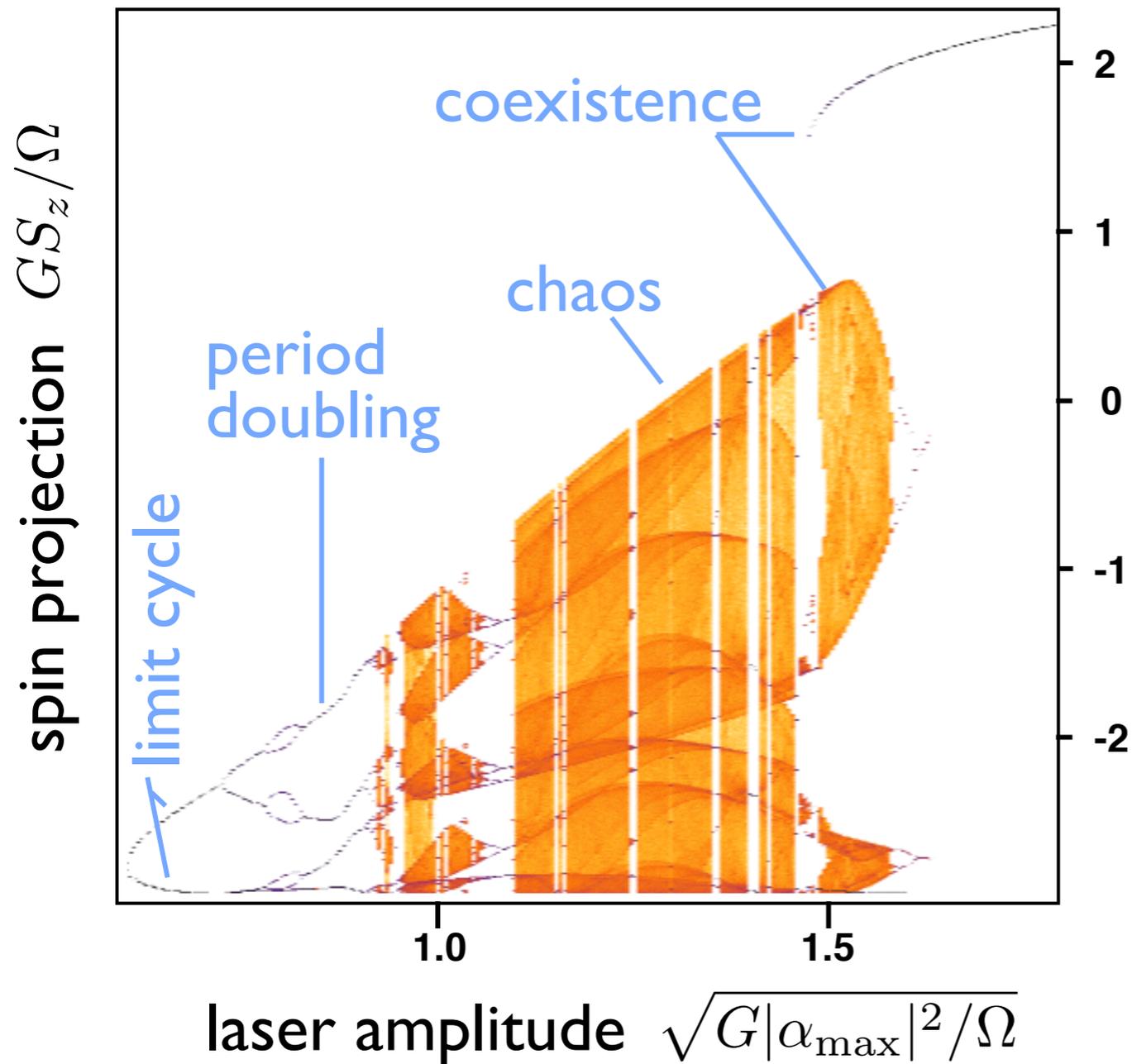
η_{opt}



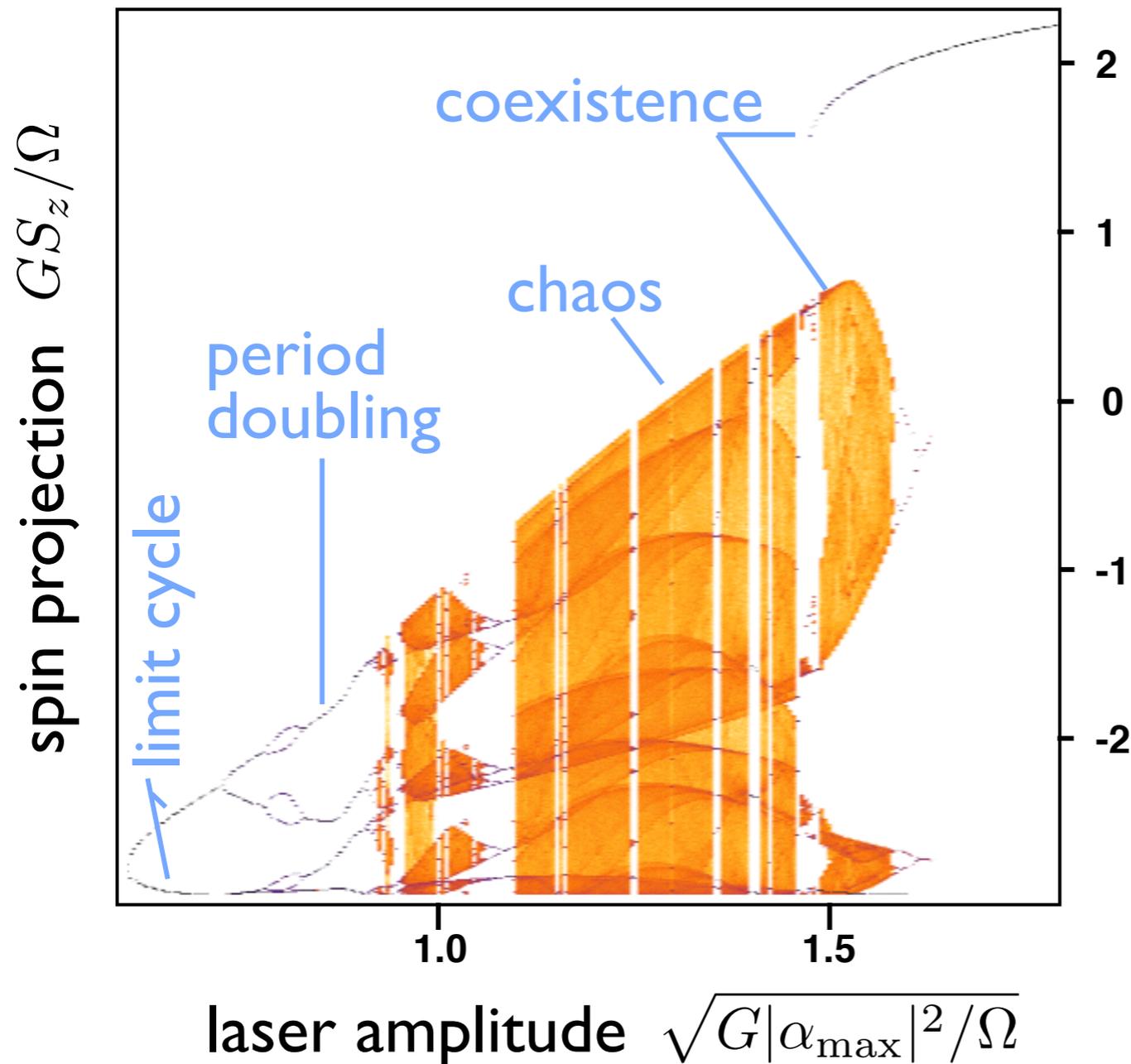
See experimental realization
with cold atoms,
Dan M. Stamper-Kurn Group
Phys. Rev. Lett. **118**, 063604
(2017)



Full Nonlinear Dynamics



Full Nonlinear Dynamics



» **Coherent optical control**

» **Magnetic switching**

» **Self-sustained oscillations**

» **Optically induced route to chaos**

Collaborators

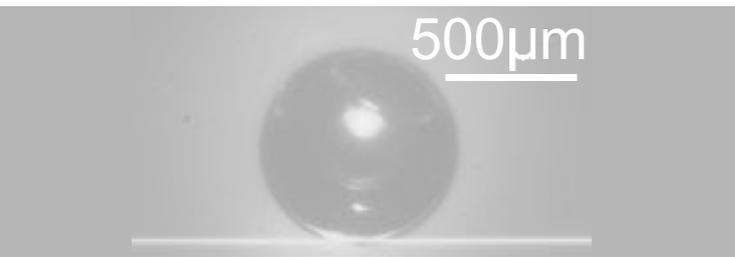
- Florian Marquardt (Erlangen)
- Hong Tang (Yale)



Magnons and the Kittel mode



Microwave regime



Optomagnonics



Optically induced spin dynamics



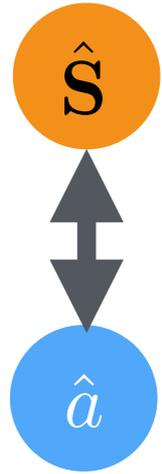
Outlook and Summary

Outlook

Problem

the state of the art optomagnonic coupling is too small

Coupling per photon $g \approx 60 \text{ Hz}$ Cooperativity $\mathcal{C} \approx 10^{-7}$



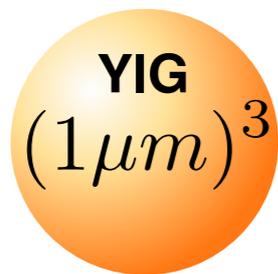
Cavity Optomagnonics

for small oscillations: spin \longrightarrow harmonic oscillator

$$\longrightarrow \hbar G \hat{S}_x \hat{a}^\dagger \hat{a} \approx \hbar G \sqrt{S/2} \hat{a}^\dagger \hat{a} (\hat{b} + \hat{b}^\dagger)$$

same form as the optomechanical Hamiltonian

coupling per magnon



$$g_0 = G \sqrt{S/2} \approx 0.1 \text{ MHz}$$

Outlook

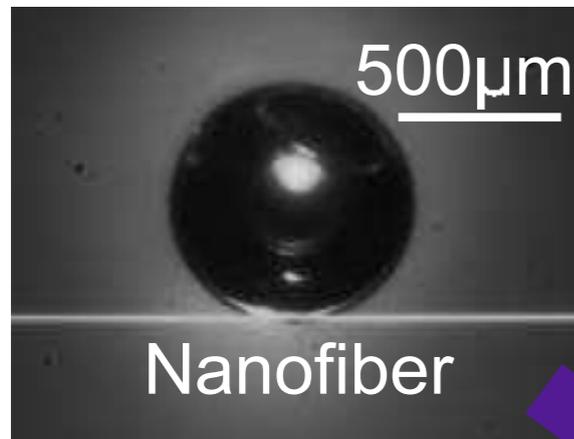
Problem

the state of the art optomagnonic coupling is too small

Coupling per photon $g \approx 60 \text{ Hz}$ Cooperativity $\mathcal{C} \approx 10^{-7}$

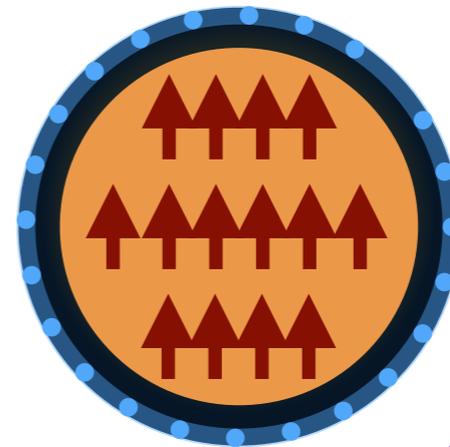
Some solutions

smaller systems



YIG
 $(1 \mu\text{m})^3$

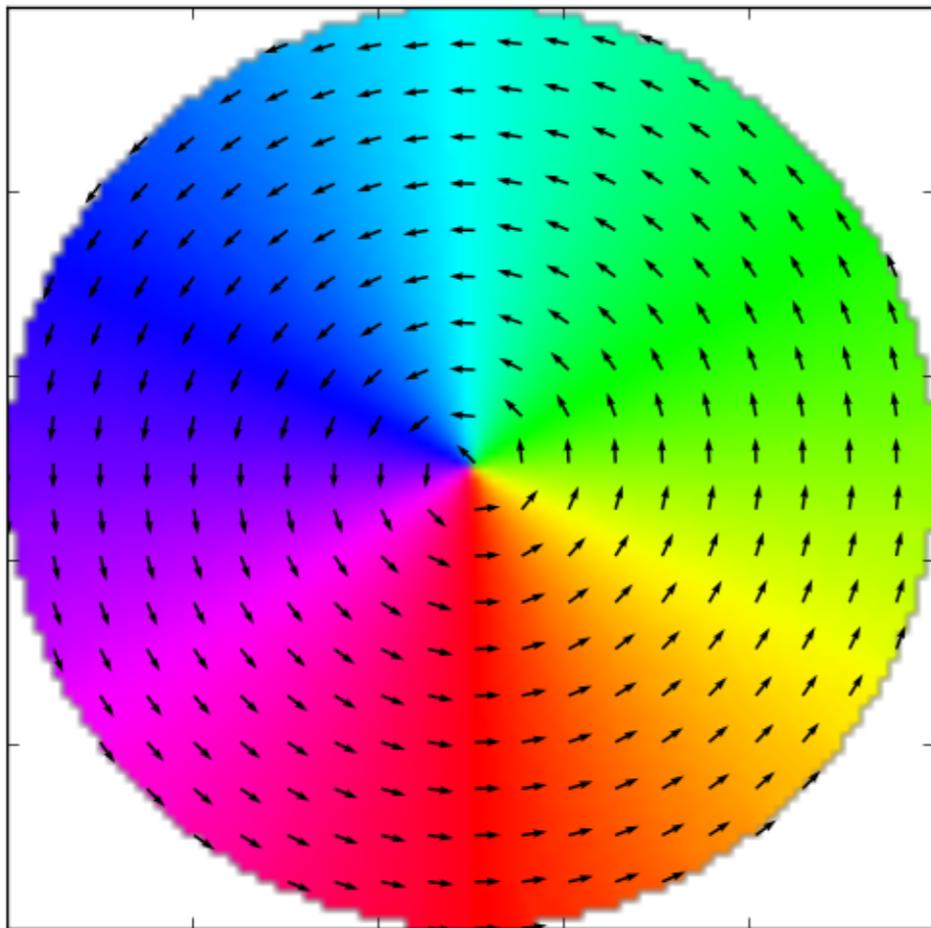
better overlap of modes



?

smaller systems

Magnetic textures

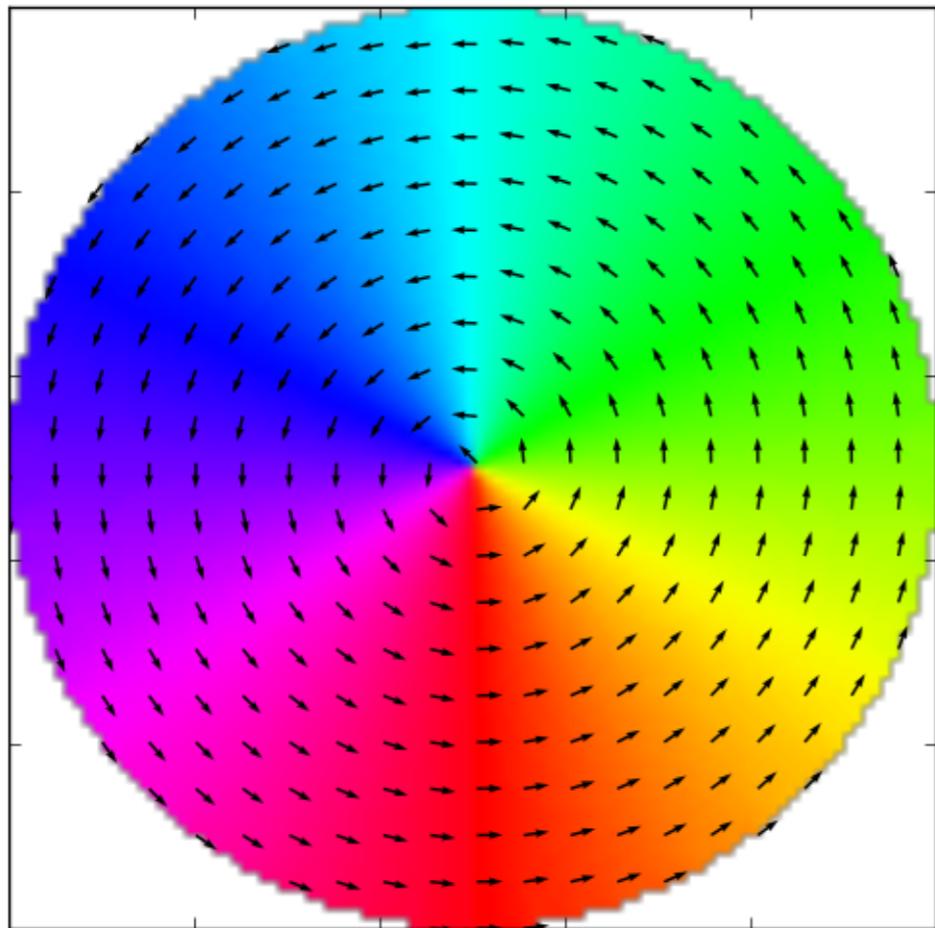


Vortex in a micro disk

Outlook

smaller systems

Magnetic textures

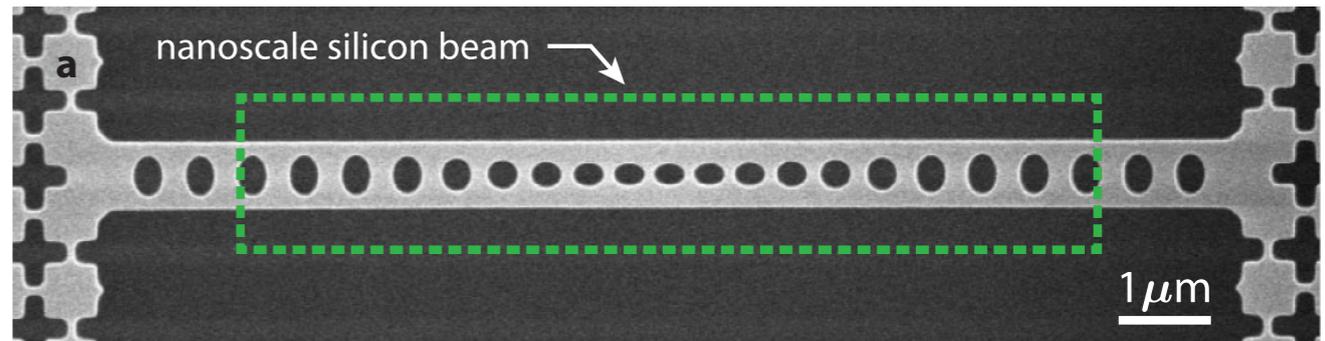


Vortex in a micro disk

better overlap of modes

optomechanical crystals

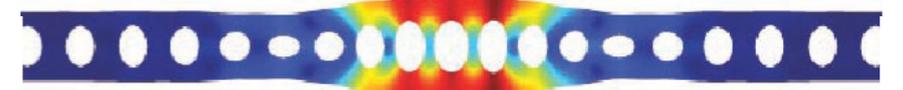
Safavi-Naeini et al, PRL 2012 (Caltech)



optical mode



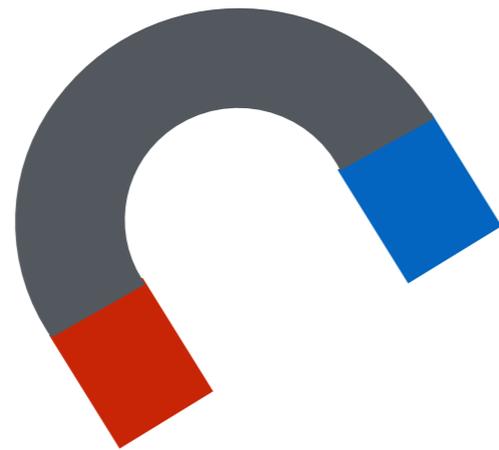
mechanical mode



Optomagnonic crystals?

Summary

- Hybrid systems for quantum technologies
- Magnetic excitations: robust, designable, quantum
- Cavity optomagnonics: promising new field



Open positions starting January 2018!



MAX PLANCK INSTITUTE
for the science of light

New Max Planck Research Group
“Theory of hybrid systems
for quantum technologies”

Erlangen, Germany