Non-linear processes in X-ray scattering experiments

School on Synchrotron and Free-Electron Laser Methods for Multidisciplinary Applications

Martin Beye, ICTP Trieste, May 18, 2018





Outline

From theoretical field interactions to real experiments

Theory

- From Schrödinger to density matrix

- Field interactions vs. photons: coherence Fundamental dipole transitions and the vacuum Multi-wave mixing schemes

Experiment

- MUSIX
- Heterodyne detection with X-rays Amplified spontaneous emission Concurrent processes
- Stimulated emission





Concept of field interactions and coherences Some theory

How does a laser pulse interact with a quantum system?

Time-dependent Schrödinger equation:

with electric field as perturbation:

Assume Eigenstates:

General time-dependent (perturbation) solution:

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$$i\hbar \frac{\partial}{\partial t} \left|\Psi\right\rangle = \hat{H} \left|\Psi\right\rangle$$

$$\hat{H} = \hat{H}_0 + \hat{W}(t)$$
 $\hat{W}(t) = \hat{\mu}\mathcal{E}\cos(\omega t)$

$$\hat{H}_0 |n\rangle = E_n |n\rangle$$

$$|\Psi(t)\rangle = \sum_{n} c_{n} e^{-iE_{n}t/\hbar} |n\rangle$$



Concept of field interactions and coherences Laser pulse couples ground state and first excited state

After a (short) pulse coupling the ground and first excited states, the system is in a coherent superposition:

 $|\Psi(t)\rangle = c_0 e^{-iE_0t/\hbar} |C|$

 $|0\rangle$ s-orbital

 $|1\rangle$ p-orbital

$$0\rangle + ic_1 e^{-iE_1t/\hbar} |1\rangle$$



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Concept of field interactions and coherences

Polarisation after such a pulse

$$\begin{split} |\Psi(t)\rangle &= c_0 e^{-iE_0t/\hbar} |0\rangle + ic_1 e^{-iE_1t/\hbar} |1\rangle \\ P(t) &= \langle \hat{\mu} \rangle = \langle \Psi(t) | \hat{\mu} | \Psi(t) \rangle \\ &= \left(c_0 e^{iE_0t/\hbar} \langle 0| - ic_1 e^{E_1t/\hbar} \langle 1| \right) \hat{\mu} \left(c_0 e^{-iE_0t/\hbar} |0\rangle + ic_1 e^{-iE_1t/\hbar} |1\rangle \right) \\ &= c_0 c_1 \langle 0 | \hat{\mu} | 1 \rangle \sin \left(\omega_{01} t \right) + c_0^2 \langle 0 | \hat{\mu} | 0 \rangle + c_1^2 \langle 1 | \hat{\mu} | 1\rangle \\ &\equiv c_0 c_1 \mu_{01} \sin \left(\omega_{01} t \right) \\ &\propto + \mu_{01}^2 \sin \left(\omega_{01} t \right) \end{split}$$

Polarisation oscillates with $\omega_{01} = (E_1 - E_0)/\hbar$ radiates field proportional to $-\mu_{01}^2 \cos(\omega_{01}t)$

- Out of phase with exciting field: exciting field is "attenuated" (destructive interference)



Concept of field interactions and coherences The radiation from this coherent state is a second interaction with the field and transfers population $|\Psi(t)\rangle = c_0 e^{-iE_0t/\hbar} |0\rangle + ic_1 e^{-iE_1t/\hbar} |1\rangle$

This state can be represented in a density matrix.

Directly after the laser pulse, without dephasing or population relaxation, this yields:

$$\rho_{00}(t) = c_0^2 = \text{const.} \\
\rho_{11}(t) = c_1^2 = \text{const.} \\
\rho_{01}(t) = -ic_0c_1e^{+i\omega_{01}t} \\
\rho_{10}(t) = ic_0c_1e^{-i\omega_{01}t} \\
\rho\left(t = 0\right) = \begin{pmatrix} 1/2 & -i/2 \\ +i/2 & 1/2 \end{pmatrix}$$

Diagonal elements are called "populations". Off-diagonal elements are "coherences".





Field interactions and level schemes (population spectroscopies) **Described one field interaction. Absorption becomes real with second field interaction**



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What is a "coherence"?

From quantum optics to experiment: Sidebands vs streaking

Coherence is a well-defined phase relation.

Quantum optics of photon / electric field Coherent state:

- -Well-defined (electric field) phase
- -Undefined photon number
- Fock state (number state):
- -Undefined phase
- -Well-defined photon number





What is a "coherence"? From quantum optics to experiment: Sidebands vs streaking

Interaction of an electric field (photons) with electrons:



Sidebands: Well-defined photon number Undefined phase

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Streaking: Undefined photon number Well-defined phase







What is a "coherence"? From quantum optics to experiment: Sidebands vs streaking

Experimental results from FLASH



Radcliffe et al., NJP 14, 043008 (2012)

Sidebands: Well-defined photon number Undefined phase

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Frühling et al., Nat. Phot. 3, 523 (2009)







Field interaction vs. photon interaction

One photon interaction can be regarded as two field interactions ($cc^* = |c|^2$)

In the photon picture the phase is lost.

Population is transferred with two (identical) field interactions.

(Synchrotron spectroscopies are nearly always "population spectroscopies": -incoherent (no phase relation) -well-defined number of transitions. Coherence only becomes relevant at FELs: -large number of coherent photons)



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$$\vec{d}|e, N-1
angle$$

Spontaneous Emission $\langle g, 1 | \vec{d} | e, 0 \rangle$

Spontaneous emission Is stimulated by the zero-point field



 $N > N_{vac}$



Everything is four-wave mixing RIXS / XES (population spectroscopies - incoherent, zero-point field)



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3. Create coherence





Everything is four-wave mixing Pump probe—XAS (population spectroscopies - incoherent)



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3. Create coherence 4. Absorb photon



Everything is four-wave mixing Claudio's transient gratings (but now coherent!)



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3. Create coherence



4. Scatter photon





Everything is four-wave mixing

Many more: CARS (coherent anti-Stokes Raman spectroscopy) CSRS (coherent Stokes Raman spectroscopy) **Stimulated Raman Inverse Raman**

Coherent spectroscopies need phase matching! k-vectors need to add up > condition for angles

- > angular separation of signal possible
- > angular focussing of signal

. . .







Everything is four-wave mixing Polarisation expansion

from Claudio:

$$P = \varepsilon_0 \cdot \left[\left(\sum_i \chi^{(1)} \cdot E_i \right) + \left(\sum_{i,j} \chi^{(2)} \cdot E_j \right) \right]$$

two-wave mixing is a $\chi^{(1)}$ process three-wave mixing is a $\chi^{(2)}$ process four-wave mixing is a $\chi^{(3)}$ process

. . .

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$(E_{j} \cdot E_{j}) + \left(\sum_{i \neq k} \chi^{(3)} \cdot E_{i} \cdot E_{j} \cdot E_{k}\right) + \cdots$

one-wave mixing does not exist... (only a pure coherence cannot be observed)



Three-wave mixing SFG (SHG, DFG)



1. Create coherence



2. Transfer coherence

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3. Emit field

E



Six-wave mixing Pump-probe (anti-Stokes) RIXS



coherence

2. Transfer population

3. Create coherence 4. Transfer population

5. Create coherence

6. Emit field





Experimental designs to realise wave-mixing spectroscopies

Things to consider:

-phase matching condition (angles, geometry) -concurrent processes and their relative strengths



SFG vs. anti-Stokes RIXS





Possible SFG geometry at an M-edge Anti-Stokes RIXS has no phase matching condition SFG beam (67eV) 5.8° 10° Incoming X-rays 63.4° 10° (64eV) 57.6° **Reflected X-rays** 33.2° (64eV) Incoming and Reflected Laser Reflected SFG beam (3.0eV) (67eV)





Experiments

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MUSIX

Multi-dimensional spectroscopy and (in-)elastic X-ray scattering instrument



- Designed at DESY
- UHV diffractometer modified from Uni Köln design, in use at BESSY • spectrometer design in house
- in (full) operation within the next two weeks
- flexible for various beam lines / light sources
- fulfill phase matching condition
- spectral analysis
- materials' science environment
- open for collaborations @ FLASH / PETRA (P04) or elsewhere







MUSIX in reality **Completed in next two weeks, ideas for future collaborations welcome**

-Nitrogen K-edge (400eV) XES -Triple RZP analyser -Resonant excitation with third harmonic @ FLASH2







Heterodyne detection: Coherent reference beams

Optimised split and delay units



Standard scheme: Split with edge of mirror







Collaboration with G. Brenner (FLASH) et al. + C. David (PSI) et al.

Amplified spontaneous emission @ FLASH

Silicon at high intensities

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Core Excitation



Core Excitation

Beye et al., Nature 501, 191 (2013)

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Geometry optimisation

Maximise interaction length with core-excited volume



Stimulated Emission

$$P_{stim} = \sigma_{stim} \rho_{ch} d \stackrel{!}{=} 1$$

$$\sigma_{stim} \approx \sigma_{abs}$$

$$\sigma_{abs} = \frac{1}{\lambda_{abs}\rho_{atom}}$$
$$\frac{\rho_{ch}d}{\lambda_{abs}\rho_{atom}} = 1$$

$$\left(\frac{\rho_{ch}}{\rho_{atom}} = \frac{\lambda_{abs}}{d}\right)$$



Stimulation front

Primarily along the surface for soft X-rays on solids



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Angular dependence Si (100), 115eV, 1J/cm², 30 fs









Fluence and spectral dependence



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b

a









Amplified spontaneous emission from water? Reabsorption of emission for long pulses and short core hole lifetimes



Liquid jet speed ~ 20 m/s ~ kHz to MHz rep. rate

PRL 113, 153002 (2014)

PHYSICAL REVIEW LETTERS

Reabsorption of Soft X-Ray Emission at High X-Ray Free-Electron Laser Fluences

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Continuous fluence variation by orders of magnitude LCLS SXR has bendable KB-mirrors







Saturation in spectra and count rate

Simple model taking Auger cascade into account, models dependences quantitatively



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Summary

From theoretical field interactions to real experiments

- Some non-linear processes have been observed
- Mostly population spectroscopy
- Electronic coherences in the X-ray regime still await discovery Potential to circumvent Auger cascade, if no population transfer Coherence lifetime undetermined, potentially very short
- WE NEED SHORTER PULSES!!!!





THANK YOU FOR YOUR ATTENTION.

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