

On 2D Materials & Their Defects under electron irradiation

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2D Materials

Xia et al., Nat. Photon. 8, 899-907 (2014)



Nature 499, 419 (2013)





Nan et al., ACS Nano 8, 5738 (2014)

Transmission Electron Microscopy





U. Ludacka, FEI Titan 80-300 @ Vienna

Jinschek, Chem. Commun. 50, 2696-2706 (2014)

Nion UltraSTEM 100 @ Vienna





Krivanek et al., Ultramicroscopy 108, 179 (2008)

Nion UltraSTEM 100 @ Vienna





Krivanek et al., Nature 464, 571 (2010)











Ramasse et al, Nano Lett. 13, 4989 (2013)

Electron Beam Effects in the (S)TEM



Characteristic cross sections and time scales for carbon @ 100 kV

| Knock-on dynamics cross section: time scale: | < 1 barn 10 ⁻¹⁵ s | lonization cross section: time scale: | 10 ⁶ barn 10 ⁻¹⁵ s | Plasmons cross section: time scale: | 10 ⁶ barn 10 ⁻¹³ s |
|--|--|---|---|---|---|
| Phonons (elastic coll cross section: time scale: | ision) 10 ⁶ barn 10 ⁻¹² s | Core hole cross section: time scale: | 10 ⁴ barn 10 ⁻¹⁴ s | Beam Current: ca. 1 e- / nm² , | 30 pA / 10 ⁻⁹ s |

Egerton, Ultramicroscopy 127, 100–108 (2013), Brühwiler et al., Phys. Rev. Lett. 74, 614–617 (1995), Yan et al., Nature Photon. 7, 394–399 (2013), Kang et al., Phys. Rev. B 81, 165405 (2010), Banhart, Rep. Prog. Phys. 62, 1181 (1999), Cosslett, J. Microsc. 113, 113-129 (1978)

Knock-On Process



Figure 1. Geometry of scattering of a light particle (electron) by a nucleus.

Maximum energy transfer occurs for the back scattering electron and yields

$$E_{\max}(E_{e}, v) = \frac{\left(2\sqrt{E_{e}(E_{e}+2mc^{2})}+Mvd\right)^{2}}{2Mc^{2}}$$

or $E_{\max}(E_{e}) = \frac{2E_{e}(E_{e}+2m_{e}c^{2})}{Mc^{2}}$ for $v = 0$

Electron energy: eUInitial velocity of the nucleus



Susi, ..., JK, Nat. Commun. 7, 13040 (2016)

Graphene @ 90 kV, TEM



Meyer, JK et al., Phys. Rev. Lett. (2012)

Graphene @ 95 kV, STEM



Susi, ..., JK, Nat. Commun. 7, 13040 (2016)

Displacement Cross Section

Coulomb Interaction:

Rutherford scattering:

$$V(\mathbf{r}) = \frac{-e^2 Q_1 Q_2}{4\pi\varepsilon_0 r} \qquad \begin{array}{l} \text{Nucleus: } Q_2 = Z \\ \text{Electron: } Q_1 = 1 \end{array}$$

 $rac{d\sigma}{d\Omega} = \left(rac{Z_1 Z_2 lpha(\hbar c)}{4 E_{
m K} \sin^2 rac{\Theta}{2}}
ight)^2$

Extended for spin and relativistic electrons by *Mott* & approximated by *McKinley* and *Feshbach*:

$$\sigma(\theta) = \sigma_R \left[1 - \beta^2 \sin^2(\theta) 2 + \pi \frac{Ze^2}{\hbar c} \beta \sin(\theta) 2(1 - \sin(\theta) 2) \right]$$

Electron scattering angle

Using

 $T(\theta) = T_{max} \sin^2(\theta/2)$ $\sigma(T) = \left(\frac{Ze^2}{4\pi\varepsilon_0 2m_0 c^2} \frac{T_{max}}{T}\right)^2 \frac{1-\beta^2}{\beta^4} \left[1-\beta^2 \frac{T}{T_{max}} + \pi \frac{Ze^2}{\hbar c} \beta \left(\sqrt{\frac{T}{T_{max}}} - \frac{T}{T_{max}}\right)\right]$

one gets

Energy of nucleus after scattering

Assuming isotropic displacement threshold E_d and integrating over $T > E_d$ leads to:

$$\sigma_d = 4\pi \left(\frac{Ze^2}{4\pi\varepsilon_0 2m_0c^2}\right)^2 \frac{1-\beta^2}{\beta^4} \left\{\frac{T_m}{E_d} - 1 - \beta^2 \ln\left(\frac{T_m}{E_d}\right) + \pi \frac{Ze^2}{\hbar c}\beta \left[2\left(\frac{T_m}{E_d}\right)^{1/2} - \ln\left(\frac{T_m}{E_d}\right) - 2\right]\right\}$$

See: Zobelli et al., Phys. Rev. B 75, 245402 (2007) McKinley and H. Feshbach, Phys. Rev. 74, 1759 (1948)



Extended model:
$$\sigma(T, E_e) = \int_{E_{max}(v, E_e) \ge T_d} P(v, T) \sigma(E_{max}(v, E_e)) dv$$

Velocity distribution of nuclei



Susi, ..., JK, Nat. Commun. 7, 13040 (2016)

Sub-threshold Knock-On Effects



Dynamical process due to an impact on one atom



Stone-Wales

"Flower defect"



JK et al., Phys. Rev. B 83, 245420 (2011) — Susi, ..., JK, 2D Mater. 4, 042004 (2017) Kurasch, JK, et al., Nano Lett. 12, 3168 (2012) — JK et al., Nat. Commun. 5, 4991 (2014)

From Single- to Multivacancies

@ 100 kV



Kotakoski et al., Phys. Rev. Lett. 106, 105505 (2011) & Kotakoski et al., Phys. Rev. B 89, 201406 (2014)

... and to Amorphous 2D Carbon



Kotakoski et al., Phys. Rev. Lett. 106, 105505 (2011) & Eder, JK et al., Sci. Rep. 4, 4060 (2014)

Impurity Atom Dynamics

Observed dynamics

| Atom | Event | T_d^{\exp} | T_d^{sim} | |
|--------------------|-------|--------------|-------------|--|
| ¹² C | КО | 21.14 | 21.9 | |
| ¹³ C | KO | 21.14 | 21.9 | |
| N _{sub} | КО | >20.4 | 19.1 | |
| C@N _{sub} | KO | 17.5 | 19.2 | |
| C@N _{sub} | Jump | 15.9 | _ | |
| C@N _{pyr} | Jump | 9.7 | | |
| В | КО | 15.9 | 18 | |
| C@B | KO | 15.1 | 19.6 | |
| C@B | Jump | 14.8 | 18.5 | |
| Si | КО | >10.0 | 13.3 | |
| C@Si | KO | 12.8 | 16.9 | |
| C@Si | Jump | 13.0 | 14.8 | |



Susi, JK, et al. Phys. Rev. Lett. 113, 115501 (2014), Yang et al., Angewandte Chemie 126, 9054 (2014), Lee et al., Nature 4, 1650 (2013), Lin et al., Nano Letters 15, 74087413 (2015), Kepaptsoglou et al., ACS Nano 9, 11398 (2015)

Susi, ..., JK, 2D Mater. 4, 042004 (2017)

Hints Towards 2D Silicon Carbide?

Sample: reduced graphene oxide Imaging: Nion UltraSTEM 100 @ 60 kV, ca. 30 pA @ 10⁻⁹ mbar







Susi, JK, et al. ,Sci. Rep. 7, 4399 (2017)



Susi, JK, et al. ,Sci. Rep. 7, 4399 (2017)



| Polytype | 2D-SiC | 6H (α) | 4H | 3C (β) |
|-----------------------------------|-----------|---------------------|---------------------|----------------------|
| Symmetry | hexagonal | hexagonal | hexagonal | cubic |
| In-plane lattice constant (Å) | 3.104 | 3.0810 | 3.0730 | 4.3596 |
| Si-C bond length (Å) | 1.792 | 1.89 | 1.89 | 1.89 |
| Bandgap (eV) | 2.58 | 3.05 | 3.23 | 2.36 |
| Bulk modulus (GPa) | 98.3 | 220 | 220 | 250 |
| Optical phonon energy (meV) | 127 | 102.8 | 104.2 | 104.2 |
| C 1 <i>s</i> – Si 2 <i>p</i> (eV) | 182.19 | 181.9 ⁴⁵ | 182.3 ⁴⁶ | 182.17 ⁴⁷ |



Knock-On vs. Ionization, or Both?

Electron irradiation of hBN

@ 120 kV





Jin et al., Phys. Rev. Lett. 102, 195505 (2009)



Meyer et al., Nano Lett. 9, 2683 (2009)



(a) 80 kV (b) 120 kV (c) 200 kV

Sulfur Vacancies in MoS₂





From Knock-On to Ionization

Heavier TMDs: MoTe₂



Non-knock-on dynamics



Elibol, ..., JK, Chem. Mater. 30, 1230 (2018)



Elibol, ..., JK, Chem. Mater. 30, 1230 (2018)

Intrinsic defects



TMD-protection through sandwiching: Algara-Siller et al., Appl. Phys. Lett. 103, 203107 (2013)

Graphene Edges under the Electron Beam

Typical Results in a HRTEM: Growing Holes





What Should We Expect?

Knock-on damage at various edges



Experimental setup in Vienna



1st FLOOR

01 - Nion UltraSTEM 100

02 - EELS Gatan PEELS 666 spectrometer Andor iXon 897 EMCCD

- 03 Customized objective area
- 04 Leak valve
- 05 Gas distribution line
- 06 Non-evaporating getter (NEG) Pressure @ sample: 1.4e-10 mbar
- 07 6W diode laser (445 nm)
- 08 Custom UHV sample entry

09 - Getec self-sensing AFSEM

(installation to finish in 2018)

TRANSFER

10 Line out from microscope room 11 To AFM and 2nd floor 12 Entry to 2nd floor 13 To loadlock and manipulation

2nd FLOOR 14 2nd floor setup

15 Storage for samples

16 Load lock and bake (150°C)

17 Manipulation system18 Knudsen evaporation cell19 Ion beam entry & 6W diode laser (ion beam to arrive in 2018)20 e-beam evaporator

21 Plasma ion source

SUITCASE 22 Ar+ suitcase (sample transfer)

... in a STEM at 10^{-7} to 10^{-6} mbar (O₂)

STEM imaging at 60 kV



Near-barrierless reconstruction



Absorbing more O atoms on C–C bridges...



...leads to the edge unraveling with no barrier

Leuthner et al., in preparation

How About in UHV?

STEM imaging at 60 kV and ca. 2x10⁻¹⁰ mbar



Etching-Induced Cleaning @ 10⁻⁷ mbar (Air)



Leuthner et al., in preparation



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