

Sugra localization and AdS black holes (part II)

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Workshop on Susy Localization and Holography:
Black Hole Entropy and Wilson Loops
ICTP, Trieste, 9-13 July 2018

Based on..

- ▶ 1803.05920 with Ivano Lodato and Valentin Reys
- ▶ 1608.07294 with Francesco Benini and Alberto Zaffaroni

Important background literature

- ▶ Susy localization - *[Pestun'07]*
- ▶ (Quantum) entropy function - *[Sen'08]*
- ▶ Localization in supergravity - *[Dabholkar, Gomes, Murthy'10-11]*
- ▶ Topologically twisted index - *[Benini, Zaffaroni'15-16]*

Quantum gravity?

- ▶ Supergravity can be seen as a toy model for quantum gravity at weak coupling
- ▶ Supersymmetric vacua are stable quantum states
- ▶ Supersymmetry allows for extrapolation of results from weak coupling (GR + matter) to strong coupling (novel quantum gravity effects)
- ▶ AdS/CFT gives a dual quantum picture, many exact results accessible via supersymmetric localization
- ▶ Existence of BPS (susy-preserving) black holes - an "integrable" sector of quantum gravity because of AdS_2 near-horizon

Black holes and susy holography

- ▶ "Black holes = statistical ensembles of gravitational degrees of freedom."
- ▶ What are the microscopic states that make up the black hole entropy?

$$S(q, p) = \log d(q, p) , \quad d(q, p) \in \mathcal{Z}_+ \quad (1)$$

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- ▶ Make gradual progress, start with susy case with AdS_2 near-horizon geometry.
- ▶ Look at sugra solutions in various dimensions, assume $SU(1, 1|1)$ near-horizon symmetry algebra ($U(1)_R$ -symmetry, unlike $SU(2)_R$ of [Strominger, Vafa'96]).
- ▶ Holography suggests a dual field theory picture: a susy $\mathcal{N} = 2$ quantum mechanics flowing to an IR conformal point.

The field theory perspective

- ▶ Dual field theory given by a Hamiltonian H_p depending on black hole magnetic charges p^i .
- ▶ Calculate grand-canonical susy partition function

$$Z(\Delta, p) = \text{Tr}_{\mathcal{H}}((-1)^F e^{i\Delta^i J_i} e^{-\beta H_p}) , \quad \langle J_i \rangle = q_i \quad (2)$$

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- ▶ Find the microcanonical partition function via a Legendre transform,

$$Z(\Delta, p) = \sum_q Z(q, p) e^{iq_i \Delta^i} , \quad Z(q, p) = \int_{\Delta} Z(\Delta, p) e^{-iq_i \Delta^i} \quad (3)$$

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- ▶ Assume no cancellation between bosonic and fermionic states in the large charge (large N) limit, find the BH entropy in the microcanonical ensemble by a saddle point approximation,

$$S_{BH}(q, p) \equiv \mathcal{I}(\dot{\Delta}) = \log Z(\dot{\Delta}, p) - iq_i \dot{\Delta}^i , \quad \left. \frac{d\mathcal{I}}{d\Delta} \right|_{\dot{\Delta}} = 0 \quad (4)$$

Weak coupling: easy example

- ▶ Look at gapped $\mathcal{N} = 2$ quantum mechanics with real masses σ^i :
susy ground states $H = \sigma^i J_i$
- ▶ Free chiral multiplet ($y = e^{i(\Delta+i\beta\sigma)}$):

$$H_c = (a^\dagger a + b^\dagger b + 1)|\sigma| - \frac{\sigma}{2} [\bar{\psi}, \psi], \quad J_c = a^\dagger a - b^\dagger b + \frac{1}{2} [\bar{\psi}, \psi] \quad (5)$$

$$Z_c(y) \left(\frac{(a^\dagger)^n}{n!} |0, \uparrow\rangle \right) = \sum_{n=0}^{\infty} y^{n+1/2} = \frac{y^{1/2}}{1-y} \quad (6)$$

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- ▶ Free fermion multiplet:

$$H_F = \sigma J_F = \sigma \left(\lambda^\dagger \lambda - \frac{1}{2} \right) \quad (7)$$

$$Z_F(y) (|\uparrow\rangle + |\downarrow\rangle) = y^{-1/2} - y^{1/2} = \frac{1-y}{y^{1/2}} \quad (8)$$

Strong coupling: localization

- ▶ Twisted index on $S^1 \times \Sigma_g$ via Bethe potential / twisted superpotential of the 2d theory on Σ_g [Hosseini, Zaffaroni'16]:

$$Z(\Delta, p) = \int du \frac{Q(u, p)}{\prod_i (1 - e^{i\partial\mathcal{W}(\Delta, u)/\partial u})} \quad (9)$$

- ▶ Large N evaluation, one leading solution \bar{u} of $e^{i\partial\mathcal{W}/\partial u} = 1$,

$$\mathcal{W}(\Delta, \bar{u}) \sim F_{S^3}(\Delta) \quad (10)$$

$$\log Z(\Delta, p) = - \sum_i p^i \frac{\partial \bar{\mathcal{W}}}{\partial \Delta^i} \quad (11)$$

Localization matches with sugra at large N

- ▶ Twisted index of ABJM theory [Benini, KH, Zaffaroni'15] match to asymptotically $\text{AdS}_4 \times S^7$ black hole entropy [Cacciatori, Klemm'09] in 11d sugra (also with mass-deformation [Bobev, Min, Pilch'18])

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$$F_{S^3} \sim N^2 \frac{\Delta^1 \Delta^2 \Delta^3}{\Delta^0}$$

More matches with sugra

- ▶ Universal twist $\bar{\Delta}^i \sim \bar{p}^i$ RG flows match to many 10d and 11d sugra black holes [*Azzurli, Bobev, Cricigno, Min, Zaffaroni'17*]

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- ▶ Evidence of field theory matches to rotating black hole entropy via anomaly coefficients: $\mathcal{N} = 4$ SYM theory [*Hosseini, KH, Zaffaroni'17*] to rotating $\text{AdS}_5 \times S^5$ black holes [*Gutowski, Reall'04*]; 6d (2, 0) theory [*Hosseini, KH, Zaffaroni'18*] to rotating $\text{AdS}_7 \times S^4$ black holes [*Cvetic, Gibbons, Lu, Pope'05; Chow'07*]

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- ▶ Subleading corrections to large N results: $\log N$ corrections from localization computed numerically [*Liu, Pando Zayas, Rathee, Zhao'17*]

Entropy at finite N ?

- ▶ No cancellation between bosons and fermions assumption! (?)
- ▶ Finite N field theory (microscopic) entropy

$$d_{micro}^{SU\mathcal{N}} \equiv Z(q, p) = \int \left(\prod_i d\Delta^i \right) \delta\left(\sum_i \Delta^i - 1\right) Z(\Delta, p) e^{-iq_i \Delta^i} \quad (12)$$

Entropy at finite N ?

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$$d_{micro}^{SUSY} \equiv Z(q, p) = \int \left(\prod_i d\Delta^i \right) \delta\left(\sum_i \Delta^i - 1\right) Z(\Delta, p) e^{-iq_i \Delta^i} \quad (12)$$

- ▶ Infer the exact (macroscopic) black hole entropy via the holographic dictionary

$$e^{S(q, p)} = d_{macro} = d_{micro} \stackrel{?}{=} d_{micro}^{SUSY} \in \mathcal{Z}_+ \quad (13)$$

- ▶ Any putative quantum gravity calculation must lead to d_{macro}
- ▶ However, holographically $d_{macro}^{SUSY} = d_{micro}^{SUSY}$, no assumptions!

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- ▶ Field theory: at weak coupling states in the grand-canonical ensemble; at strong coupling use localization / anomalies.
- ▶ Quantum gravity: at weak coupling sugra calculation (?); at strong coupling string theory (definition via the field theory dual?)
- ▶ In QFT well-defined microcanonical and grand-canonical ensemble; in quantum gravity so far only microcanonical ensemble?

Quantum entropy - see V. Reys' talk!

- ▶ Formal definition:

$$d_{macro} \equiv \left\langle \exp \left(4\pi q_i \int_0^{2\pi} W_{\tau}^i d\tau \right) \right\rangle_{EAdS_2=\mathcal{H}_2}^{finite} \quad (14)$$

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- ▶ Explicit approach for d_{macro}^{SUSY} : pick a sugra theory and a black hole near-horizon geometry (a BPS *solution*). Assume/prove the gravitational background to be fixed (*freeze gravity multiplet*) and perform susy localization on a curved background of the remaining (vector-, hyper-, tensor-) multiplets in the bulk.

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- ▶ Conceptual issues (see [[de Wit, Murthy, Reys'18](#)]): 1-loop determinant of the gravity multiplet, non-compact space (Euclidean AdS_2), exact integration measure. Let's neglect them at a first approximation!

4d $\mathcal{N} = 2$ off-shell gauged supergravity

- ▶ Conformal sugra formalism developed in *[de Wit, van Holten, van Proeyen'80], ...*
- ▶ use Euclidean version for full consistency *[de Wit, Reys'17]*
- ▶ Weyl multiplet: metric $g_{\mu\nu}$, auxiliary tensor T_{ab}^{\pm} and scalar D , gauge fields b_{μ} (dilatation), A_{μ} ($SO(1,1)_R$), $\mathcal{V}_{\mu j}^i$ ($SU(2)_R$), gravitini and dilatini

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- ▶ (compensating) hypermultiplet: four real scalars A_i^{α} , hyperini
- ▶ after gauge fixing equivalent to Poincare sugra with n_V vector multiplets and FI gauging ξ_I

On-shell solution and superalgebra

- ▶ half-BPS near-horizon geometry in off-shell formalism [*de Wit, van Zalk'12*], [*KH, Katmadas, Lodato'16*]
- ▶ gravity: $ds^2 = v_1 ds_{AdS_2}^2 + v_2 ds_{S^2}^2$, $b_\mu = A_\mu = 0$,
 $D = -(v_1^{-1} + 2v_2^{-2})/6$, $T_{12}^\pm = w^\pm = \pm 2v_1^{-1/2}$

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- ▶ vectors: $\dot{F}_{12}^I = e^I$, $\dot{F}_{34}^I = p^I$, $\dot{Y}^{ij,I}(\dot{X}_\pm^I)$, scalars $X_\pm^I(\xi_I, q_I, p^I)$
subject to attractor mechanism
- ▶ hypers: gauge fix to break $SU(2)_R$ to $U(1)_R$, $\mathcal{V}_{\mu j}^i = -2i\xi_I \dot{W}_\mu^I \sigma_{3j}^i$,
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- ▶ $SU(1, 1|1)$ superalgebra, bosonic subgroups $SU(1, 1) \times U(1)_R$
- ▶ pick localizing supercharge Q , s.t. $Q^2 = \mathcal{L}_\tau + \delta_{U(1)_R} + \delta_{gauge}$
equivariant differential

Localization locus

- ▶ Weyl multiplet frozen
- ▶ vector multiplet: arbitrary functions $C_k^I(\theta, \varphi), D_k^I(\theta, \varphi), Y^{ij,I}(C_k^I, D_k^I)$

$$\delta X_{\pm}^I = \sum_{k=1}^{\infty} (D_k^I \pm C_k^I) r^{-k} \quad (15)$$

$$\delta W_{\tau}^I = \sqrt{v_1} \sum_{k=2}^{\infty} (D_{k-1}^I + C_k^I) (r^{1-k} - 1) \quad (16)$$

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- ▶ hypermultiplet extra constraint $\xi_I \delta X_{\pm}^I = 0$
- ▶ use integration variables $\phi_{+}^I \equiv 2\dot{X}_{+}^I \sum_{k=1}^{\infty} (C_k^I + D_k^I)$ and $\phi_{\perp}^I(C_k, D_k)(\theta, \varphi)$:

$$\partial_{\theta, \varphi} \phi_{+}^I = 0 = \partial_{\theta, \varphi} \delta W_{\tau}^I, \quad \xi_I \phi_{+}^I = 1 \quad (17)$$

Classical action

- ▶ two-derivative + Wilson line action:

$$S^{2d} + S^W = 2\pi r_0 \left(p^I (\dot{\mathcal{F}}_I^+ + \dot{\mathcal{F}}_I^-) + q_I e^I \right) - 2\pi \left(p^I \mathcal{F}_I^+(\phi_+) + q_I \phi_+^I \right) \quad (18)$$

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- ▶ reinstate Newton's constant,

$$d_{macro}^{SUSY} = \int_{\phi_+} \exp \left[-\frac{\pi}{2\mathcal{G}_N} \left(p^I \mathcal{F}_I^+(\phi_+) + q_I \phi_+^I \right) \right] Z_{\text{ind}}^{\text{reg}}(\phi_+) \quad (19)$$

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- ▶ 1-loop, measure, gravity localization - hidden inside $Z_{\text{ind}}^{\text{reg}}(\phi_+)$
- ▶ higher derivative F-terms: additional $256F_A$ to the classical action, string theory origin?

Saddle point = attractor mechanism

- ▶ Saddle point evaluation

$$\frac{\partial}{\partial \phi_+^I} (p^I \mathcal{F}_I^+(\phi_+) + q_I \phi_+^I) |_{\dot{\phi}_+^I} = 0, \quad S_{BH} = -\frac{\pi}{2\mathcal{G}_N} (p^I \mathcal{F}_I^+(\dot{\phi}_+) + q_I \dot{\phi}_+^I)$$

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- ▶ Precise match with attractor mechanism (Lorentzian) of *[Cacciatori, Klemm'09]*, *[Dall'Agata, Gecchi'10]* after Wick rotation.
- ▶ Saddle point agreement between d_{macro}^{SUSY} and $d_{micro}^{SUSY} = Z(q, p)$ in all known examples:

$$\mathcal{F}(11d/S^7) \sim \sqrt{\phi^0 \phi^1 \phi^2 \phi^3}$$

$$\mathcal{F}(mIIA/S^6) \sim (\phi^1 \phi^2 \phi^3)^{2/3}$$

$$\mathcal{F}(IIB/S^5 \times S^1) \sim \frac{\phi^1 \phi^2 \phi^3}{\phi^0}$$

Grand-canonical ensemble?

- ▶ Field theory microcanonical partition function

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- ▶ Supergravity localization result

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- ▶ Supergravity localization result

$$d_{macro}^{SUSY} = \int \left(\prod_i d\phi_+^I \right) \delta\left(\sum_i \phi_+^I - 1\right) e^{-\frac{\pi}{2\mathcal{G}_N} p^I \mathcal{F}_I^+} Z_{\text{ind}}^{\text{reg}} e^{-\frac{\pi}{2\mathcal{G}_N} q_I \phi_+^I} \quad (21)$$

- ▶ Define grand-canonical ensemble in sugra,

$$Z^{\text{sugra}}(\phi_+, p) = \exp\left[-\frac{\pi}{2\mathcal{G}_N} (p^I \mathcal{F}_I^+(\phi_+))\right] Z_{\text{ind}}^{\text{reg}}(\phi_+) \quad (22)$$

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- ▶ Look for the answer in Euclidean theory, [*Freedman, Pufu'13*], [*Bobev, Elvang, Freedman, Pufu'13*] and [*Cassani, Martelli'14*]?

Future work

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- ▶ Go beyond susy and extremality - near AdS_2 geometries, SYK, ...