

Part 2
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Topics



Part 1. Basic concepts. Z-pinch, pulsed power, plasma focus.

Part 2. How to obtain information from a dense transient plasma?

Plasma diagnostics

Basic Research and Applications

Part 3. How to design and to build a small plasma focus? Tricks and Recipes



Part 2: outline



- How to obtain information from a dense transient plasma?
 Plasma diagnostics
- Basic research:
 - Plasma dynamics, singularities structures, filaments, schoks, jets
- Applications:
 - As x-rays and neutron sources
 - To study materials for fusion reactors
 - Film deposition
 - To study the effects of puldsed radiation on life matter
 - Plasma thrusters for nanosatellites



How to obtain information from a dense transient plasma



Diagnostics

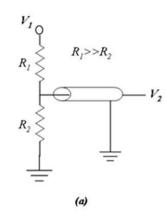
- Electrical signals
- Visible plasma images
- X-ray detections (temporal and spatial resolution)
- Neutron detection (in particular low yield pulses)
- Charged particles
- Optical refractive diagnostics
- Spectroscopy

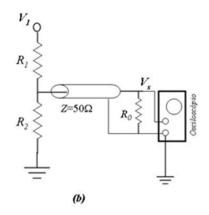


Electrical signals



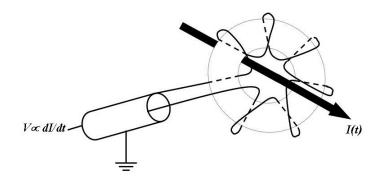
$$V_2 = \frac{R_2}{R_1 + R_2} V_1$$





Voltage monitor: resistive divider

 $V \alpha dI/dt$

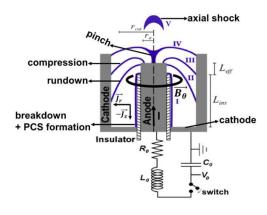


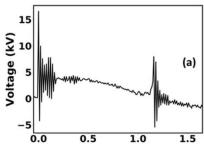
Current monitor: Rogowski coil

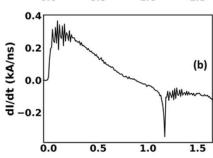


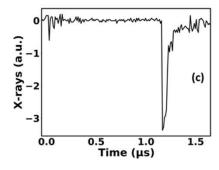
Electrical signals











$$V(t) = \frac{d}{dt} \left[(L_p(t) + L_0)I(t) \right] \tag{1}$$

That gives

$$L_p(t) = \frac{\int_{to}^{t} V(t)dt + (L_0 + L_p(t_0))I(t_0)}{I(t)} - L_0$$

$$L_p(t) + L_0 = \frac{\int_{t_0}^t V(t)dt + (L_0 + L_p(t_0))I(t_0)}{I(t)}$$
 (2)

$$L_P(t) = L_P(t_c) + L'_P(t)$$
(4)

for $t > t_c$:

$$V(t) = [L_0 + L_P(t_c)] \frac{dI}{dt} + \frac{d}{dt} (IL_P')$$

$$V_P = V(t) - (L_0 + L_P(t_c)) \frac{dI}{dt}$$
(6)

$$V_P = V(t) - \left(L_0 + L_P(t_c)\right) \frac{dI}{dt} \tag{6}$$

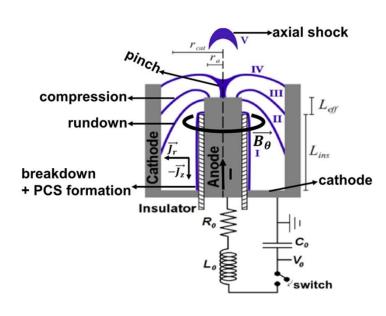
$$V_P = \frac{d(IL_P')}{dt} \tag{7}$$

F. Veloso, C. Pavez, J. Moreno, V. Galaz, M. Zambra and L. Soto, Journal of Fusion Energy 31, 30-37 (2012)

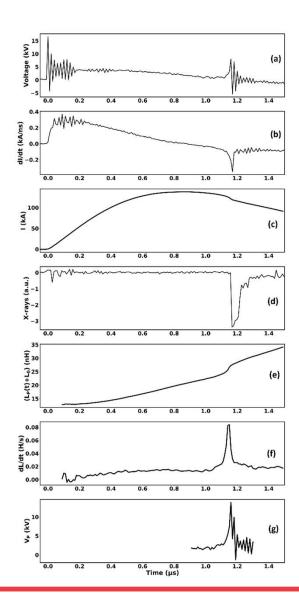


Electrical signals





$$L_p(t) = (\mu_0 / 2\pi) z(t) ln(b/r(t))$$

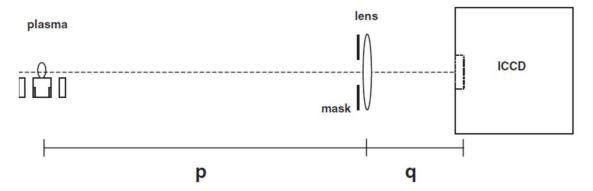




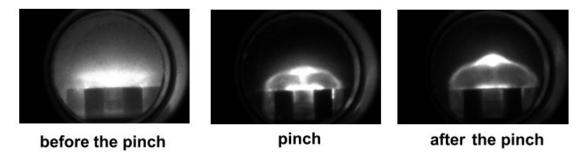
Visible plasma images



Images from plasma light are captured with a ICCD camera, 4ns exposure time



Plasma Dynamics

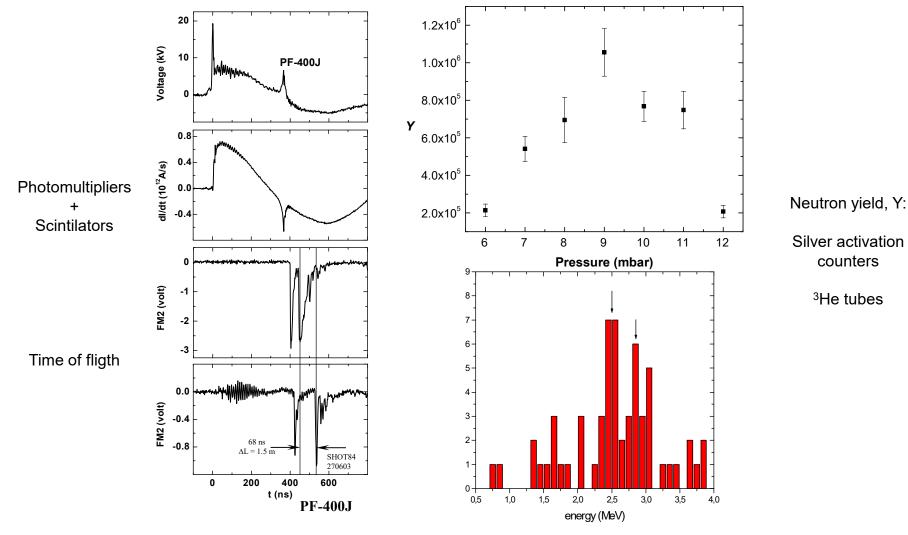


J. Moreno, P. Silva, and L. Soto, Plasma Sources Science and Technology 12, 39 (2003)



X-rays and neutron detection





P. Silva, J. Moreno, L. Soto, L. Birstein, R. Mayer, and W. Kies, App. Phys. Lett. 83, 3269 (2003)



Optical refractive diagnostics

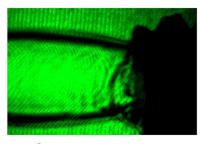


$$\mu_e = 1 - \frac{1}{2} \frac{\omega_p^2}{\omega^2}$$

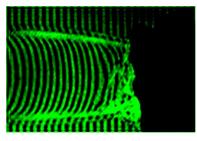
$$\mu_e - 1 = -4,49 \cdot 10^{-16} \lambda^2 n_e$$



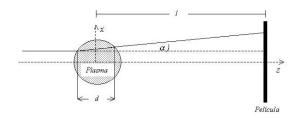
Shadowgraph



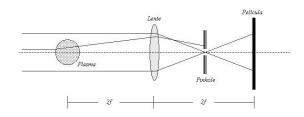
Schlieren



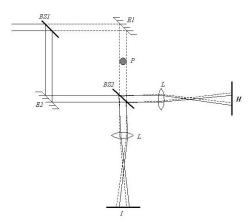
Interferometry



 $\frac{\Delta I}{I} \approx l \int_{z_1}^{z_2} \nabla_{\perp}^2 \mu(x, y, z) dy$



$$\alpha_x = \int_0^d \frac{1}{\mu} (\partial \mu / \partial x) dz$$

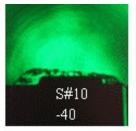


$$\Delta \varphi = \frac{2\pi}{\lambda} \int_{x_1}^{x_2} (\mu(x, y, z) - \mu_0) dx$$



PF-400J

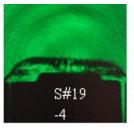




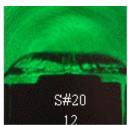


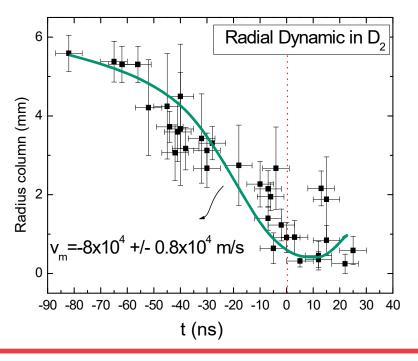


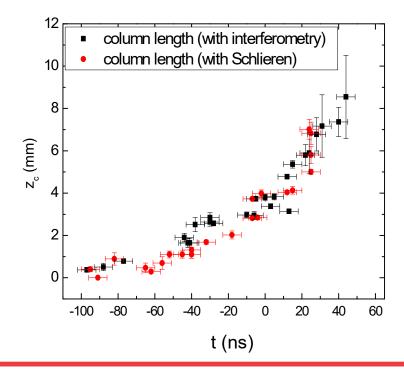








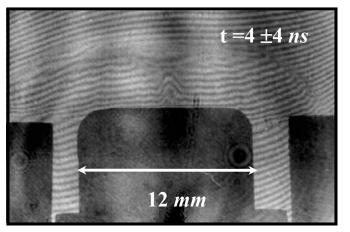


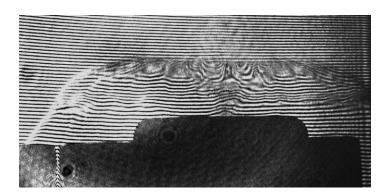


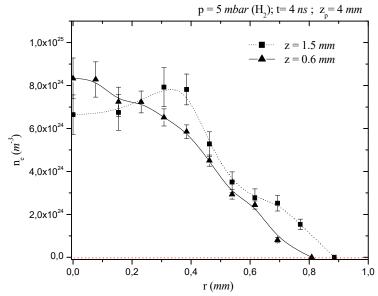


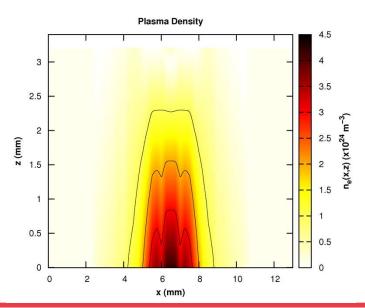
PF-400J









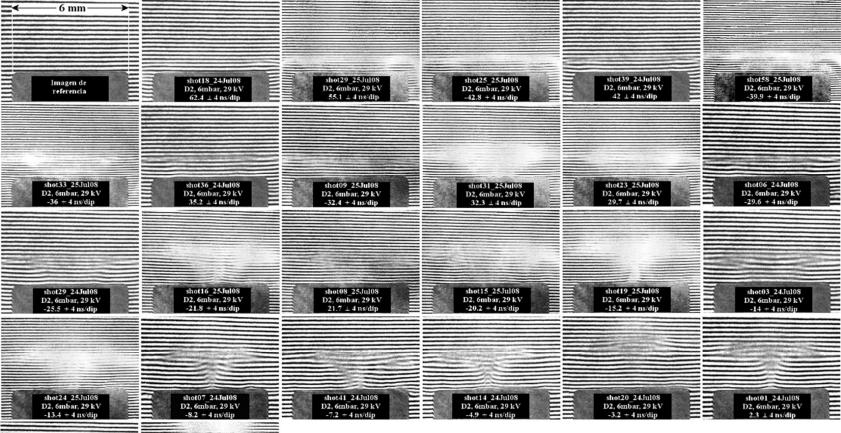


C. Pavez and L. Soto, Physica Scripta **T131**, 014030 (2008)



PF-50J





Radial and pinch phase occurs during the last 55ns before the first quarter of period (~150 ns). After column disruption, remaining plasma propagates in the axial direction as a shock wave, but no structure is observed on the axis at the anode end.

A. Tarifeño, C. Pavez, J. Moreno and L. Soto, IEEE Trans. Plasma Science, 39, 756 (2011)

shot04_25Jul08

D2, 6mbar, 29 kV

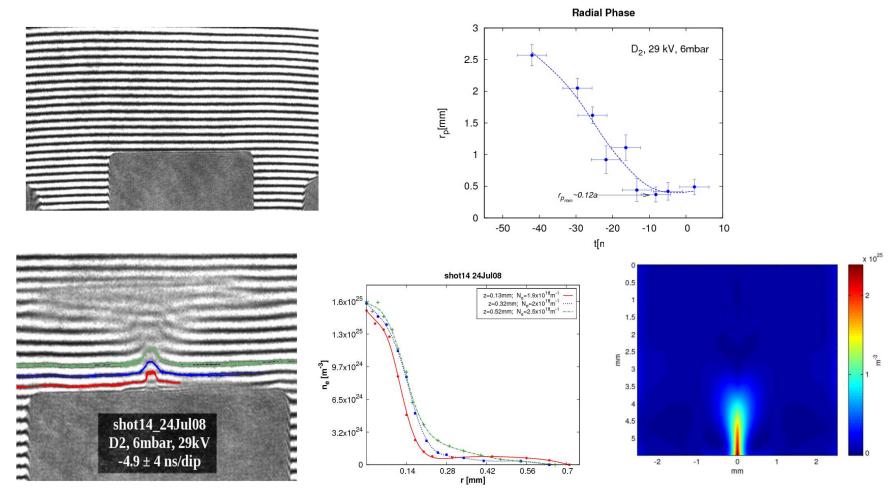
shot23 24Jul08

D2, 6mbar, 29 kV



PF-50J





A. Tarifeño, C. Pavez, J. Moreno and L. Soto, IEEE Trans. Plasma Science, 39, 756 (2011)

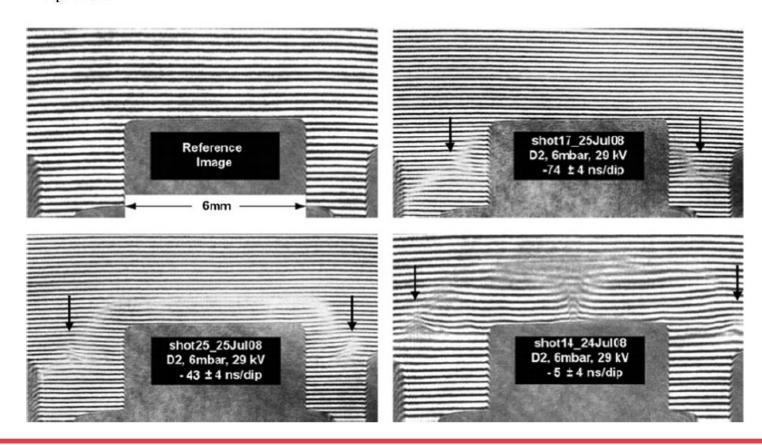
J Fusion Energ (2012) 31:279–283 DOI 10.1007/s10894-011-9469-1

ORIGINAL RESEARCH



Toroidal High-Density Singularity in a Small Plasma Focus

Federico Casanova · Ariel Tarifeño-Saldivia · Felipe Veloso · Cristian Pavez · Alejandro Clausse · Leopoldo Soto



center for research and applications in plasma physics and pulsed power

ORIGINAL RESEARCH

Toroidal High-Density Singularity in a Small Plasma Focus

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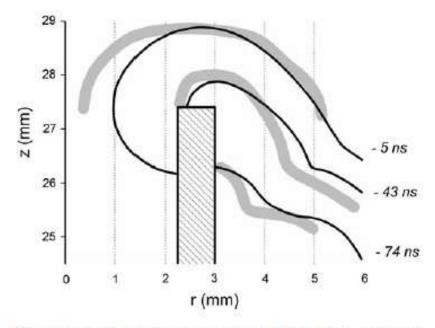


Fig. 5 Shape of the current sheet at different times. Numerical (black), experimental (solid grey). The numbers at the right indicate the corresponding time relative to dip







PHYSICS OF PLASMAS 21, 072702 (2014)

Filamentary structures in dense plasma focus: Current filaments or vortex filaments?

Leopoldo Soto, ^{1,2,3,a)} Cristian Pavez, ^{1,2,3} Fermin Castillo, ⁴ Felipe Veloso, ⁵ José Moreno, ^{1,2,3} and S. K. H. Auluck ⁶

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⁵Instituto de Física, Pontificia Universidad Católica de Chile, 7820436 Santiago, Chile

⁶Bhabha Atomic Research Center, Mumbai 400 085, India



PF-400J





Visible images

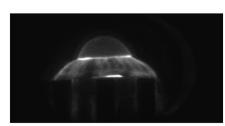
Schlieren



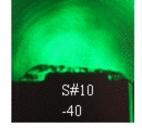
-16ns



- 6ns



49ns



S#21 -32



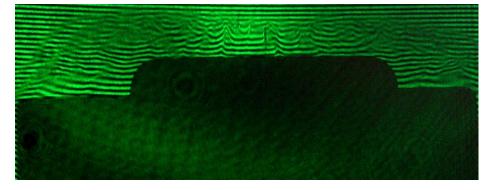








Interferogram



Filaments diameter $\sim 300 \mu m, \quad n_e \sim 10^{25} \ m^{\text{-}3}$

L. Soto, C. Pavez, F. Castillo, F. Veloso, J. Moreno, S. K. Auluck, Physics of Plasmas 21, 072702 (2014)





IOP Publishing

Plasma Physics and Controlled Fusion

Plasma Phys. Control. Fusion 57 (2015) 035008 (6pp)

doi:10.1088/0741-3335/57/3/035008

Neutron energy distribution and temporal correlations with hard x-ray emission from a hundreds of joules plasma focus device

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- Department of Physics, Faculty of Electrical Engineering, Czech Technical University, Technicka 2, 16627 Prague 6, Czech Republic

E-mail: jmoreno@cchen.cl

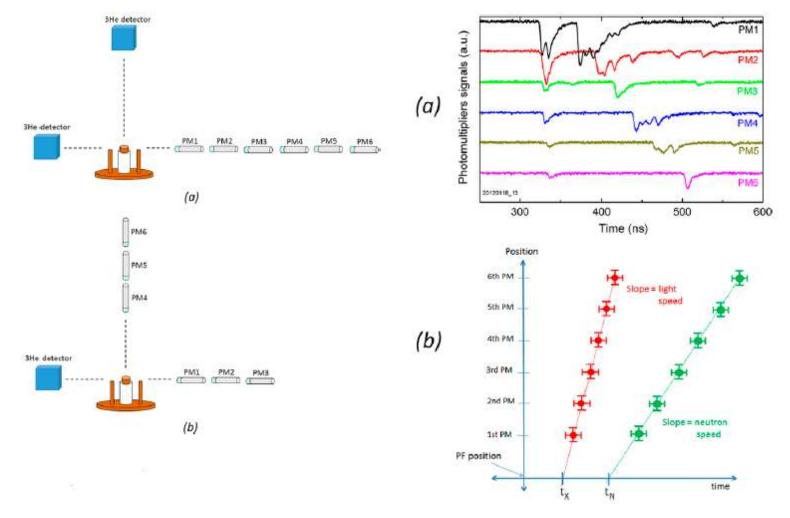
Received 1 October 2014, revised 31 December 2014 Accepted for publication 14 January 2015 Published 18 February 2015





Neutron energy distribution and temporal correlations with hard x-ray emission on 400J



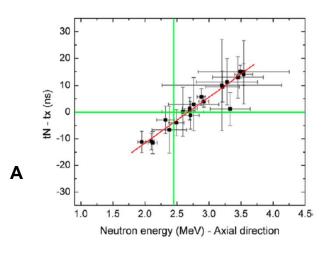


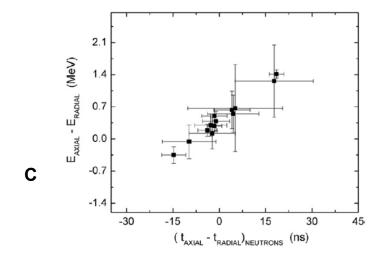
J. Moreno, F. Veloso, C. Pavez, A. Tarifeño-Saldivia, D. Klir, and L. Soto, Plasma Phys. Control. Fusion 57, 035008 (2015)

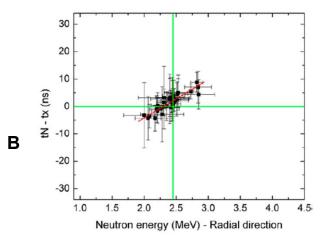


Neutron energy distribution and temporal correlations with hard x-ray emission on PF-400J









These results not only show differences in the production time of hard x-rays and neutrons, but also some correlation on the neutron energy and the $t_N - t_X$ time difference in both directions (i.e., the larger neutron energy corresponds to later times with respect to hard x-rays emission).

The axial-to-radial ratio of both total neutron yield and neutron energies indicates anisotropic emission, which is consistent with a 100 keV kinetic energy of the deuterons in the axial direction. The energy spread among different shots was ~0.5 MeV in the axial direction which is 2.5 times the spread in the radial direction. Furthermore, temporal differences on hard x-rays and neutron production over each direction are found. These differences show correlation with neutron energies. This could be related to the existence of two temporally separated neutron production times corresponding to different moments during the plasma focus discharge.

J. Moreno, F. Veloso, C. Pavez, A. Tarifeño-Saldivia, D. Klir, and L. Soto, Plasma Phys. Control. Fusion 57, 035008 (2015)







PHYSICS OF PLASMAS 21, 122703 (2014)

Characterization of the axial plasma shock in a table top plasma focus after the pinch and its possible application to testing materials for fusion reactors

Leopoldo Soto, ^{1,2,3,a)} Cristian Pavez, ^{1,2,3} José Moreno, ^{1,2,3} María José Inestrosa-Izurieta, ^{1,2} Felipe Veloso, ⁴ Gonzalo Gutiérrez, ⁵ Julio Vergara, ⁶ Alejandro Clausse, ⁷ Horacio Bruzzone, ⁸ Fermín Castillo, ⁹ and Luis F. Delgado-Aparicio ¹⁰

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⁴Instituto de Física, Pontificia Universidad Católica de Chile, Santiago, Chile

⁵Departamento de Física, Facultad de Ciencias, Universidad de Chile, Santiago, Chile

⁶Facultad de Ingeniería, Pontificia Universidad Católica de Chile, Santiago, Chile

⁷CNEA-CONICET and Universidad Nacional del Centro, 7000 Tandil, Argentina

⁸CONICET and Universidad de Mar del Plata, Mar del Plata, Argentina

⁹Instituto de Ciencias Físicas, Universidad Nacional Autónoma de México, Cuernavaca, Morelos, Mexico

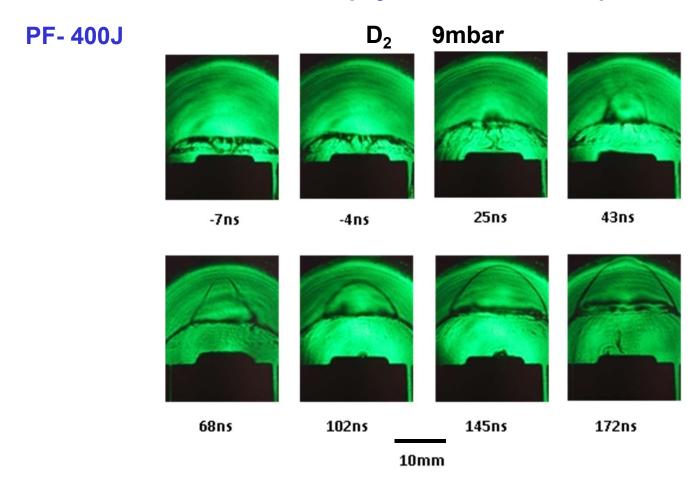
¹⁰Princeton Plasma Physics Laboratory, Princeton University, Princeton, New Jersey 08543, USA



Plasma bursts after the pinch



Previus studies did not pay atention after the pinch disruptions

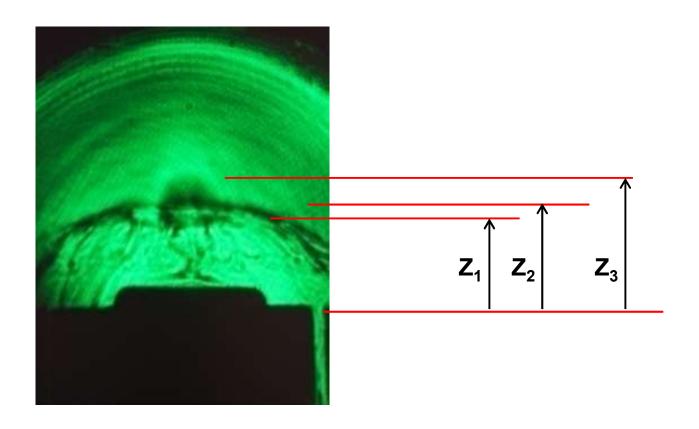


L. Soto, C. Pavez, J. Moreno, M. J. Inestrosa, F. Veloso, G. Gutierrez, J. Vergara, F. Castillo, A. Clausse, H. Bruzzone and L. Delgado-Aparicio, Physics of Plasmas 21, 122703 (2014)



Plasma bursts after the pinch



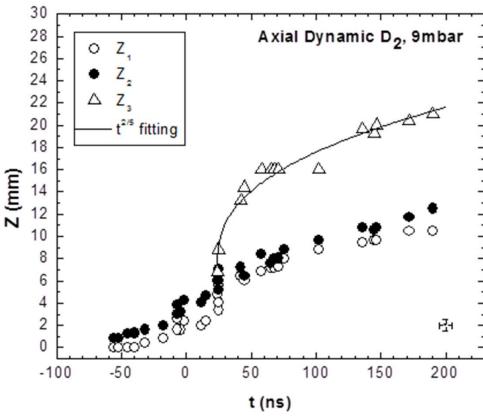


L. Soto, C. Pavez, J. Moreno, M. J. Inestrosa, F. Veloso, G. Gutierrez, J. Vergara, F. Castillo, A. Clausse, H. Bruzzone and L. Delgado-Aparicio, Physics of Plasmas 21, 122703 (2014)



Plasma bursts after the pinch





$$Z_3(t) - Z_3(t_0) = \left[\frac{75}{16\pi} \frac{(\gamma - 1)(1 + \gamma)^2}{(3\gamma - 1)} \frac{E}{\rho_0} \right]^{\frac{1}{5}} (t - t_0)^{2/5}$$

L. Soto, C. Pavez, J. Moreno, M. J. Inestrosa, F. Veloso, G. Gutierrez, J. Vergara, F. Castillo, A. Clausse, H. Bruzzone and L. Delgado-Aparicio, Physics of Plasmas 21, 122703 (2014)







PHYSICS OF PLASMAS 22, 040705 (2015)

Observation of plasma jets in a table top plasma focus discharge

Cristian Pavez, 1,2,3 José Pedreros, 1,4 Ariel Tarifeño-Saldivia, 1,2,a) and Leopoldo Soto 1,2,3,b)

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³Departamento de Ciencias Físicas, Facultad de Ciencias Exactas, Universidad Andrés Bello, República 220, Santiago, Chile

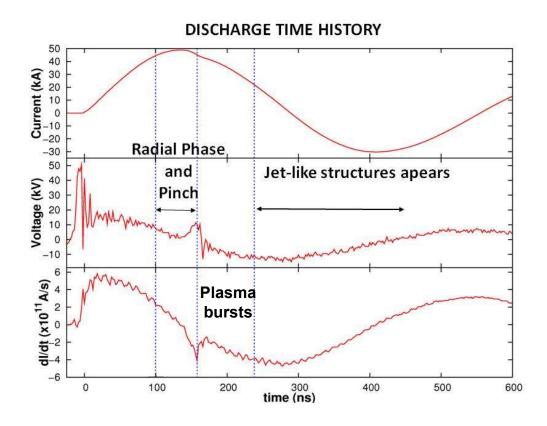
⁴Departamento de Ingeniería Eléctrica, Universidad de Santiago de Chile, Santiago, Chile



After plasma jets are observed



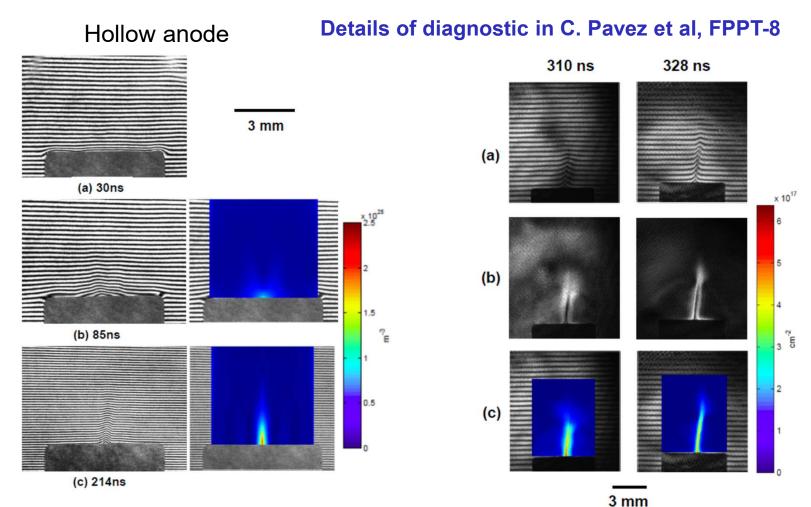
PF-50J





After plasma jets are observed





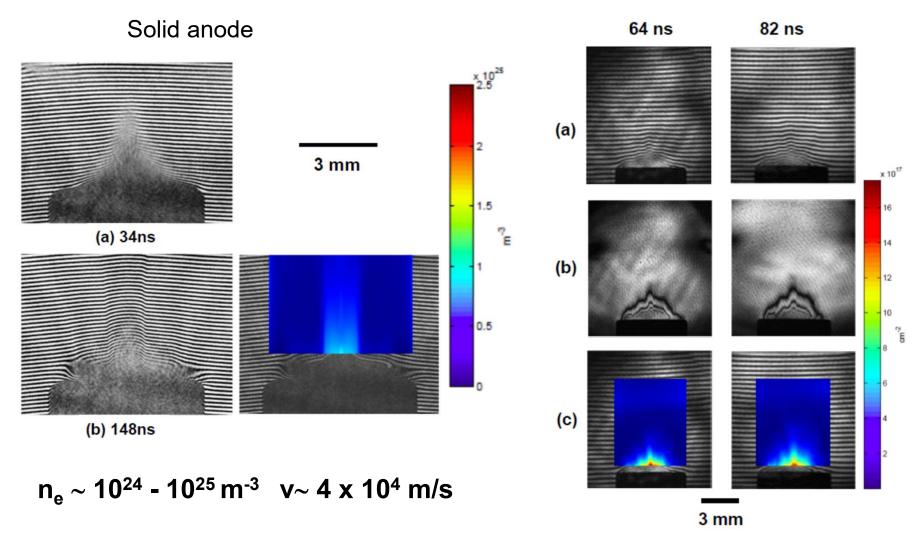
 $n_e \sim 10^{24}$ - $10^{25} \, m^{\text -3} \, v \sim 4 \, x \, 10^4 \, m/s$

C. Pavez, J. Pedreros, A Tarifeño-Saldivia and L. Soto, Physics of Plasmas 22, 040705 (2015)



Plasma jets





C. Pavez, J. Pedreros, A Tarifeño-Saldivia and L. Soto, Physics of Plasmas 22, 040705 (2015)



PF dynamics including times after the pinch disruption









Plasma Focus Aplications



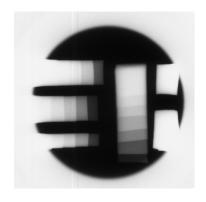


- X-ray pulses, ns
- Neutron pulses, ns
- Filaments
- Plasma shocks
- Jets

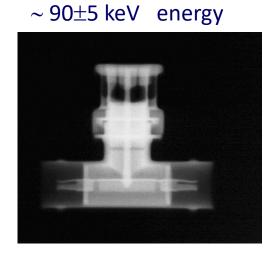


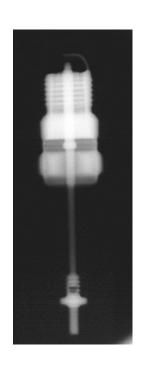
Hard X-ray nanoflash

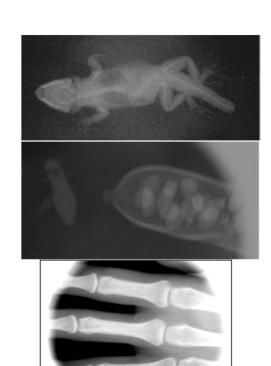




X-ray from PF-400J

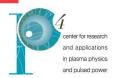






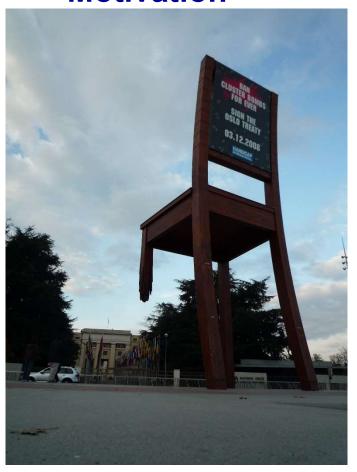
M. Zambra, P. Silva, M. Moreno, C. Pavez and L. Soto, Plasma Physics Controlled Fusion 51, 125003 (2009)
C Pavez, J. Pedreros, M Zambra, F Veloso, J Moreno, A Tarifeño-Saldivia and L. Soto, Plasma Phys. and Control. Fusion. 54 105018 (2012)





A PF for field applications

Motivation



Development of a confirmation method using the neutron backscattering technique for detection of landmines in arid soils

TC IAEA Project



A portable PF device as neutron source for field applications, PF-2J



HYDAD-D at a simulated field with hydrogenated objects



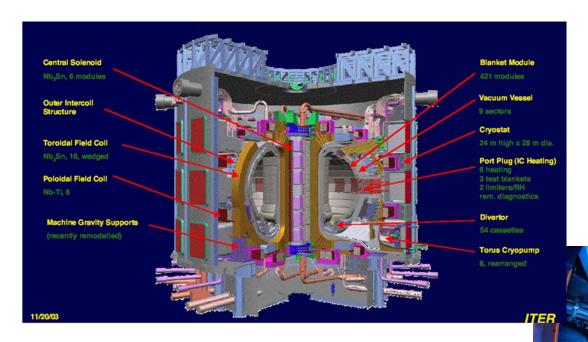
Arica, Atacama desert, North of Chile, September 11, 2009

C. Pavez, F. D. Brooks, F. D Smit, J. Moreno, L. Altamirano, L. Soto "Tests of the HYDAD Landmine Detector on Dry Soil in Northen Chile, VII Latin American Symposium on Nuclear Physics and Applications, Santiago, Chile, Dec. 2009.



Applications to study materials for fusion reactors





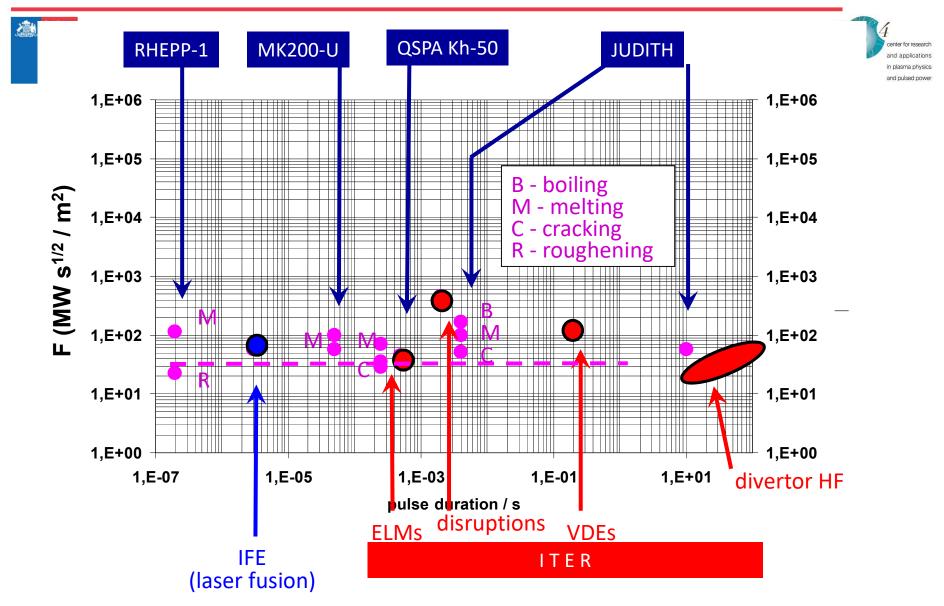




Damage factor

$$F \sim q \cdot \tau^{1/2} = E/S \tau^{1/2}$$

q: power flux, τ : interaction time, S: interaction area



J. Linke et al, J. Nuclear Mat. 367-370, 1422 (2007)



Expected Damage in Fusion Reactor



ITER:

$$F \sim q \cdot \tau^{1/2} \sim 10^8 (W/m^2) s^{1/2} = 10^4 (W/cm^2) s^{1/2}$$

at $0.5 - 1 Hz$, 10^3 pulses

IFE:

$$F \sim q \cdot \tau^{\frac{1}{2}} \sim 10^4 (W/cm^2) s^{\frac{1}{2}}$$

at 10 Hz

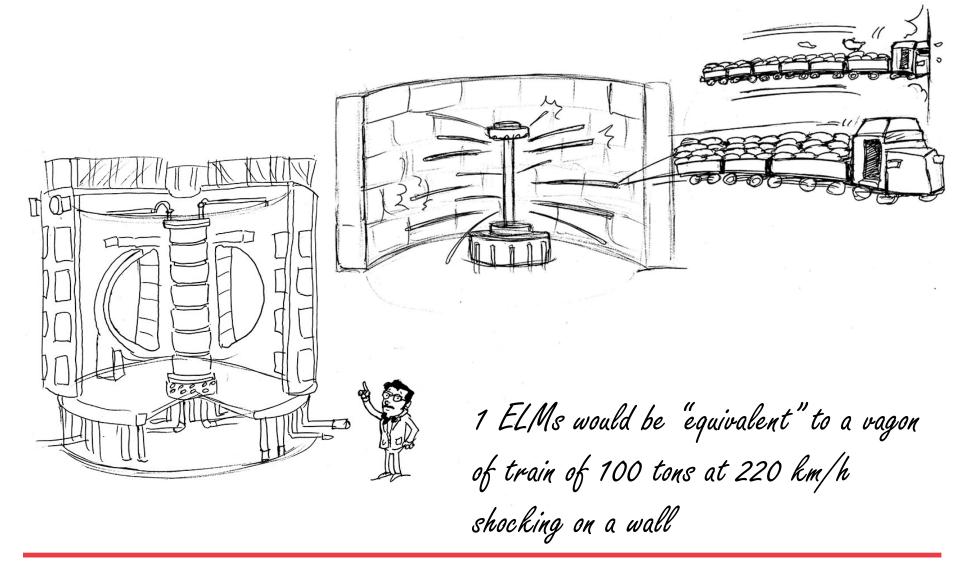
PF-400J:

$$F \sim q \cdot \tau^{1/2} \sim 10^3 - 10^5 (W/cm^2) s^{1/2}$$

at 0.05 Hz



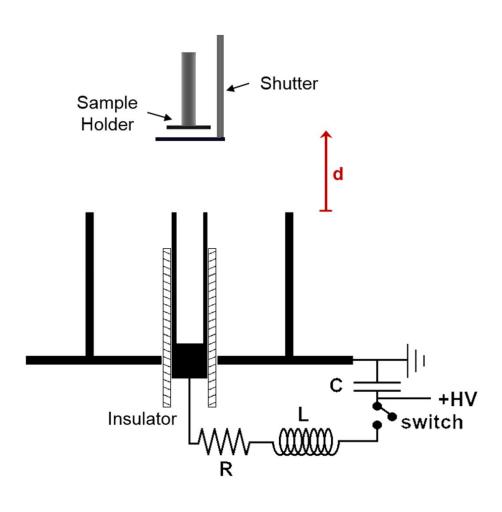






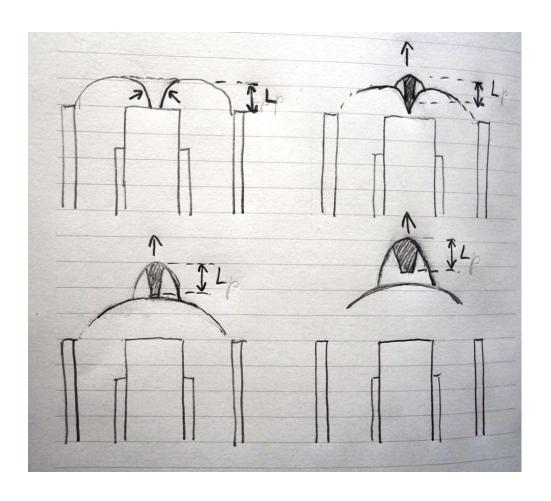
Advanced Materials







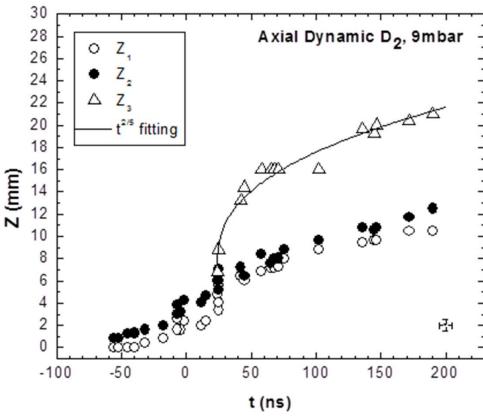






Plasma bursts after the pinch





$$Z_3(t) - Z_3(t_0) = \left[\frac{75}{16\pi} \frac{(\gamma - 1)(1 + \gamma)^2}{(3\gamma - 1)} \frac{E}{\rho_0} \right]^{\frac{1}{5}} (t - t_0)^{2/5}$$

L. Soto, C. Pavez, J. Moreno, M. J. Inestrosa, F. Veloso, G. Gutierrez, J. Vergara, F. Castillo, A. Clausse, H. Bruzzone and L. Delgado-Aparicio, Physics of Plasmas 21, 122703 (2014)



Damage Factor produced by Plasma bursts after the pinch



Total mass inside the bubble, m: ~ total pinch mass

(the pinch is ejected trough Z2, creating so the bubble)

The pinch density was previously measured using pulsed interferometry, thus the total pinch mass is $m \sim 1.5 \times 10^{-10} \text{ kg}$

C. Pavez and L. Soto, Physica Scripta T131, 014030 (2008)

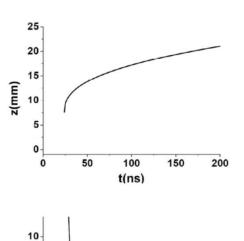
Length of the ejected mass: ~ pinch length, L = 5.6 mm

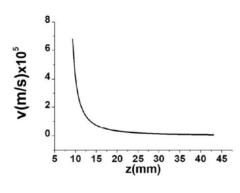
Time of interaction, $\tau \sim L / v$

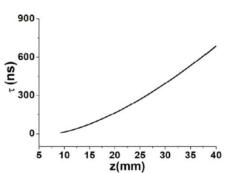


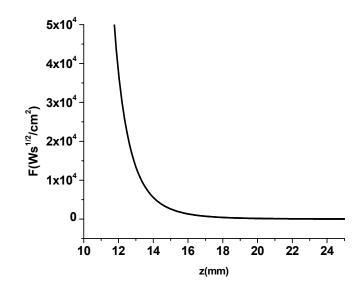
Tunable Damage Factor

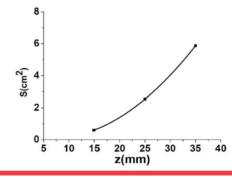












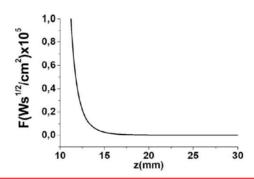
20 25 z(mm)

25

30 35

15

10



L. Soto et al, in preparation



Power flux density does not depend on PF energy



For PF devices:

$$E/a^3 \sim const$$
 a: anode radius $n \sim const$

v ~ const

T ~ const

 $r_p \sim 0.1a$

 $z_p \sim 0.8a$

Plasma ejected from the pinch (burst) on a target at Z, \sim 1.5a < Z < 2.7a

$$m \sim m_p \alpha V_p \alpha a^3$$

Thickness ~ pinch length $L_p \alpha$ a

Time of interaction $\tau = L / v \alpha$ a

Cross section S: a²

 $q \sim KE/ \tau S \alpha m/ \tau S \alpha a^3 / a a^2 \sim const$

L. Soto et al, in preparation



6 order of magnitud in energy translates in only 1 order of magnitude in damage factor



damage factor, F ~ $\mathbf{q} \cdot \tau^{1/2} \alpha \tau^{1/2} \alpha a^{1/2} \alpha (E^{1/3})^{1/2}$

 $F \alpha E^{1/6}$



PF, 1MJ F

PF, 1kJ ~ 1/3 F

PF, 100J ~ 1/5 F

PF, 10J ~ 1/7 F

PF, 1J ~ 1/10 F

Roughly speaking

The damage factor for the PF-1000 (1MJ) at Poland is only 3.65 times greater than the damage factor for the PF- 400J (400J) at Chile.

L. Soto et al, in preparation





IOP Publishing | International Atomic Energy Agency

Nuclear Fusion

Nucl. Fusion 55 (2015) 093011 (8pp)

doi:10.1088/0029-5515/55/9/093011

Morphological and structural effects on tungsten targets produced by fusion plasma pulses from a table top plasma focus

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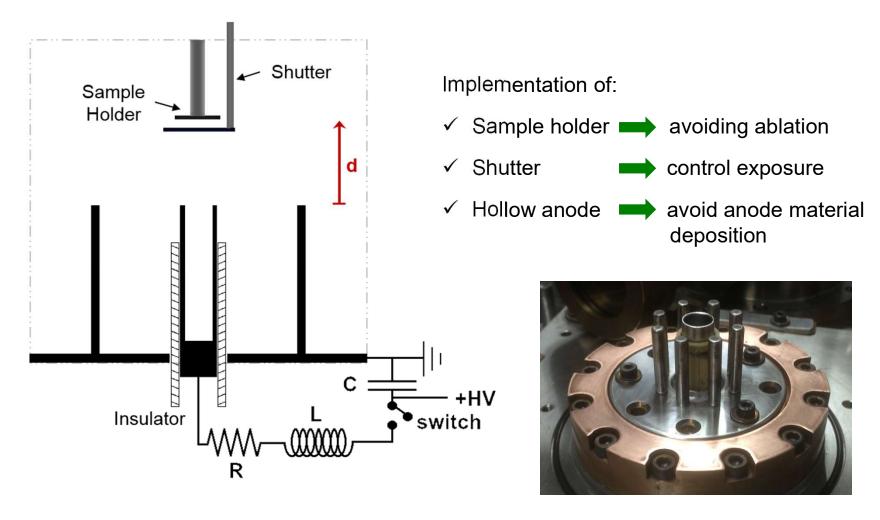
Received 30 January 2015, revised 17 June 2015 Accepted for publication 30 June 2015 Published 5 August 2015





PF400J: Fusion Plasma Pulses Source





M. J. Inestrosa Izurieta, E. Ramos-Moore and L. Soto, Nuclear Fusion 55, 093011 (2015)



PF400J: Fusion Plasma Pulses Source Tungsten target



 $m \sim 1.5 \times 10^{-10} \text{ kg}$

L = 5.6 mm

 $\tau \sim L/v$

| Z (mm) | V (m/s) | S (cm ²) | E/S (J/cm ²) | τ (ns) | q (W/cm²) | F=qτ ^{1/2} |
|--------|----------------------|----------------------|--------------------------|--------|---------------------|--------------------------------------|
| | | | | | | (W/cm ²)s ^{1/2} |
| 15 | 7.5×10 ⁴ | 0.6 | 0.69 | 75 | 9.2×10 ⁶ | 2.5×10 ³ |
| 25 | 2.08×10 ⁴ | 2.54 | 1.3×10 ⁻² | 270 | 4.7×10 ⁴ | 24 |
| 35 | 1.05×10 ⁴ | 5.87 | 1.4×10 ⁻³ | 533 | 2.6×10 ³ | 1.9 |

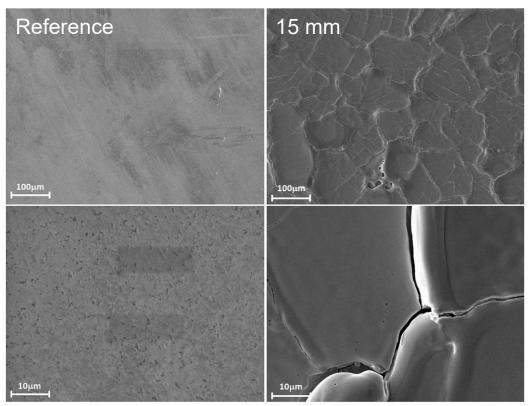
M. J. Inestrosa Izurieta, E. Ramos-Moore and L. Soto, Nuclear Fusion 55, 093011 (2015)



Morphological Effects on W SEM



15 mm

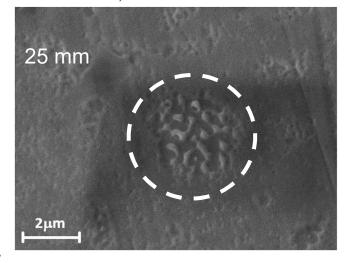


Scanning Electron Microscope images to comparison the extreme irradiation targets

Scanning Electron Microscope image showing targets.

- ✓ Ref. → smooth surface
- √ 15 mm

 microcracks and holes surface melting
- ✓ 25 mm some melting
- ✓ 35 mm → no melting

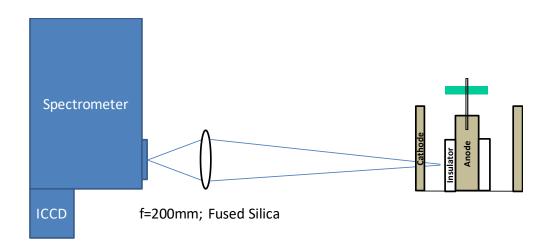


M. J. Inestrosa Izurieta, E. Ramos-Moore and L. Soto, Nuclear Fusion 55, 093011 (2015)



Studies of the plasma interacting with a target material on front of the anode using visible spectroscopy





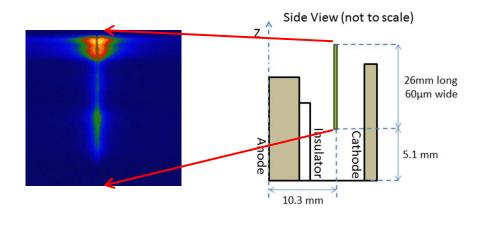
G. Avaria et al, in preparation

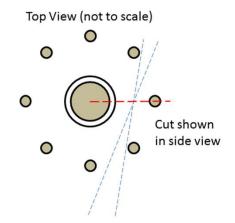
Discharge:

- Charging voltage ~27kV (~310 J)
- Frequency 0.06Hz (~ 16 s)

Diagnostics

- 0.5 m Cerny-Turner *Imaging* Spectrometer
 - 300 l/mm
 - Optical resolution (FWHM):0.4 nm
- ICCD
 - FWHM: 3 ns

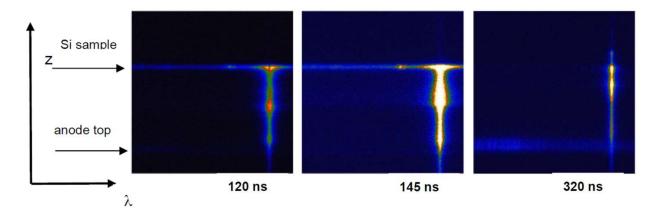


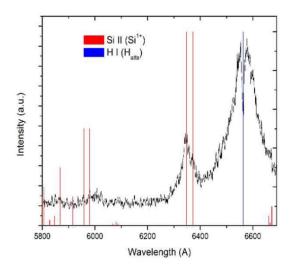




Studies of the plasma interacting with a target material on front of the anode using visible spectroscopy







Visible spectra at different times for the region between the anode and the sample Si wafer, z axis is the axial position between the anode top and the sample, x axis is the wavelength λ . At 120 ns after the pinch Si II (Si¹⁺) is observed in the sample region. Also, the $H\alpha$ emission is observed from the anode towards the sample. The spectrum close to the sample is shown at left. A Si and hydrogen plasma can be identified at the sample and a hydrogen-only plasma at the anode top. At 145 ns there is a similar emission as seen at 120 ns, but with more dense hydrogen plasma at the sample (inferred from the $H\alpha$ line broadening) and hydrogen plasma at middle position between the anode top and sample. At 320 ns on the anode region Fe III (Fe²⁺), Al II (Al^{l+}) , Al III (Al^{2+}) , O II (O^{l+}) can be identified. Si II is not observed on the sample region at that time. In addition, a plasma beyond the sample is observed.

G. Avaria et al, in preparation



Repetitive table top plasma focus to reproduce an equivalent damage on materials than the expected by type I ELMs in ITER



| C (µF) | 12 |
|------------------------------|-----|
| L (nH) | 50 |
| V (kV) | 8.2 |
| E (J) | 403 |
| I (kA) | 127 |
| T/4 (µs) | 1,2 |
| Anode radius (mm) | 6 |
| Anode lenght (mm) | 60 |
| Maximum repetition rate (Hz) | 1 |



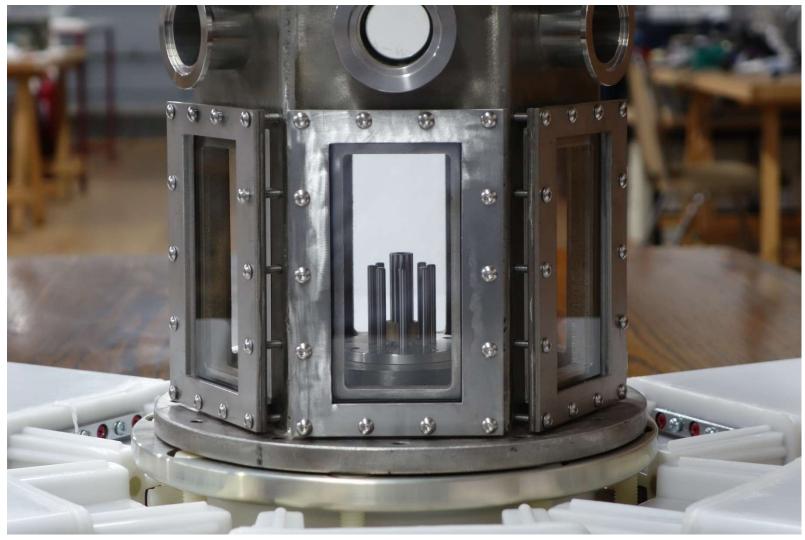


A tabletop PF devices to study the effects of thermonuclear plasmas on materials



Repetitive table top plasma focus to reproduce an equivalent damage on materials than the expected by type I ELMs in ITER







Film deposition



AIP ADVANCES 7, 105026 (2017)

Ti film deposition process of a plasma focus: Study by an experimental design

M. J. Inestrosa-Izurieta, 1,2,3,a J. Moreno, 1,2,3 S. Davis, 1,2,3 and L. Soto 1,2,3

(Received 27 July 2017; accepted 24 October 2017; published online 31 October 2017)

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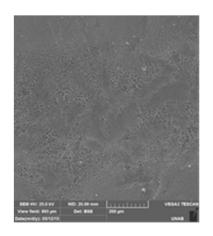
³Universidad Andres Bello, Departamento de Ciencias Fisicas, Facultad de Ciencias Exactas, Republica 220, Santiago, Chile

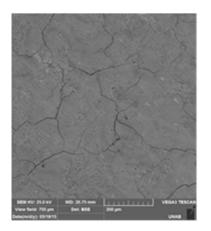


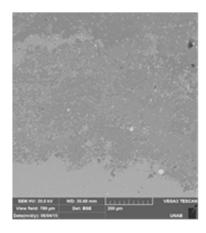
Film deposition



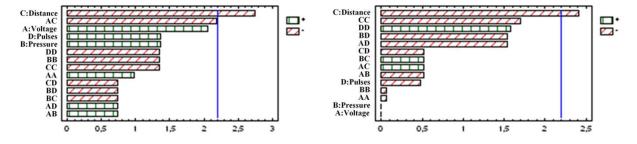
A study varying: Voltage [kV] / Pressure [mbar] / Distance [mm] / # Pulses







SEM images of selected Ti samples with marker length of $200\mu m$ on all the micrographs. From left to right 28/7/21/20, 27/6/9/50 and 27/6/15/80.



Standardized effects of Pareto chart for TiSi₂ (left) and Ti₅Si₃ (rigth) formation.

Inestrosa-Izurieta, J. Moreno, S. Davis and L. Soto, AIP Advances 7, 105026 (2017)



Applications to biology and biomedicine



Effects of pulsed radiation in cell

AIP ADVANCES 7, 085121 (2017)



Hundred joules plasma focus device as a potential pulsed source for *in vitro* cancer cell irradiation

J. Jain,^{1,2} J. Moreno,^{2,3,6} R. Andaur,⁴ R. Armisen,^{5,7} D. Morales,² K. Marcelain,^{4,a} G. Avaria,^{2,3,6} B. Bora,^{2,3,6} S. Davis,^{2,3,6} C. Pavez,^{2,3,6} and L. Soto,^{2,3,6,a}

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⁵Centro de Investigación y Tratamiento del Cáncer, Facultad de Medicina, Universidad de Chile, Independencia 1027, Independencia, Santiago, Chile

⁶Universidad Andres Bello, Departamento de Ciencias Fisicas, Republica 220, Santiago, Chile ⁷Current affiliation: Center for Excellence in Precision Medicine, Pfizer Chile, Santiago 7810305, Santiago, Chile





Mathematics is all that is not understood.

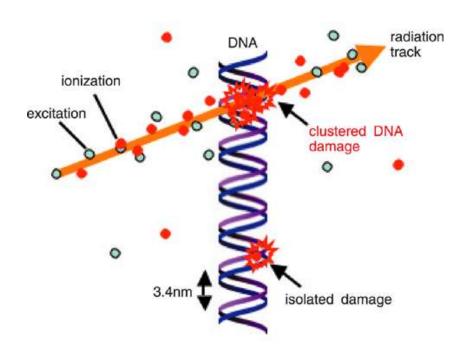
Physics is all that does not work.

Chemistry is everything that smells bad.

Biology is all that is green and that it crawls.





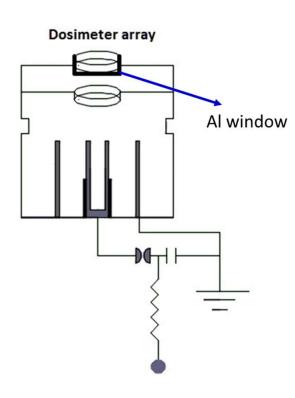




Effects of pulsed radiation in cell









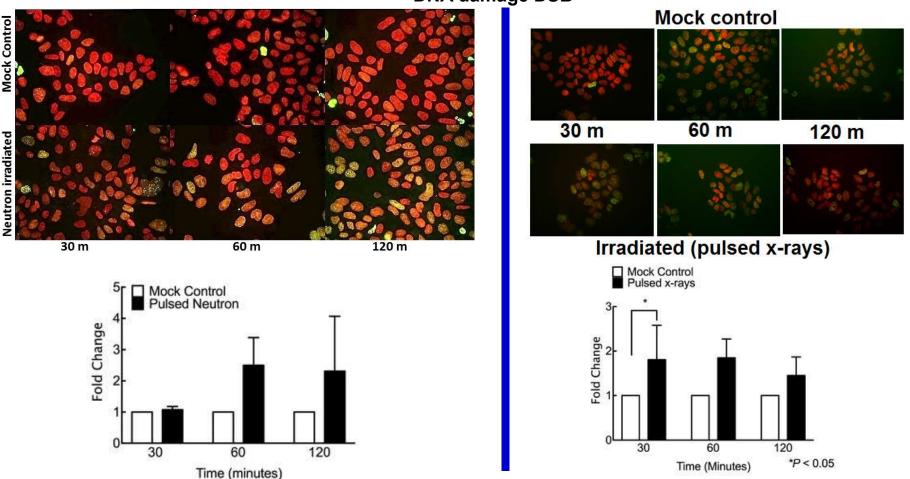
"J. Jain, J. Moreno, R. Andaur, R. Armisen, D. Morales, K. Marcelain, G. Avaria, B. Bora, S. Davis, C. Pavez, and L. Soto, AIP Advances **7**, 085121 (2017)



Effects of neutron and x-ray pulses on cancer cell



DNA damage DSB



J. Jain, J. Moreno, R. Andaur, R. Armisen, D. Morales, K. Marcelain, G. Avaria, B. Bora, S. Davis, C. Pavez, and L. Soto, AIP Advances 7, 085121 (2017)





- Cell death was absent in case pulsed x-rays irradiation.
- Neutron irradiation provides cell death at ultralow doses but DNA damage with higher statistical insignificance.
- The effect depends on the repair time of the cells. Therefore, the time between pulses is a relevant parameter



Joint ICTP-IAEA Workshop on Dense Magnetized Plasma and Plasma Diagnostics, Triste, Italy, November 15-17, 2010

Jalaj Jain was a participant and after he made his PhD in Chile in our group in combination with biologist of University of Chile.

J. Jain, J. Moreno, R. Andaur, R. Armisen, D. Morales, K. Marcelain, G. Avaria, B. Bora, S. Davis, C. Pavez, and L. Soto, **AIP Advances 7, 085121 (2017)**



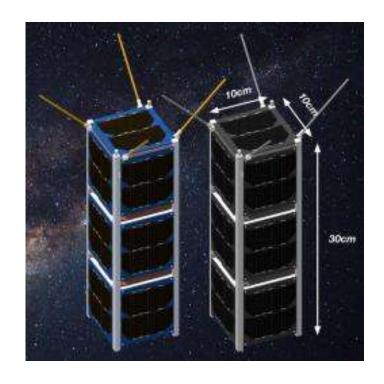
Pulsed Plasma Thruster for nanosatellites based on ultra-miniaturized Plasma Focus



Nanosatellite SUCHAI 2017 University of Chile



SUCHAI-1 10 cm x 10 cm x 10 cm At present in orbit and in operation



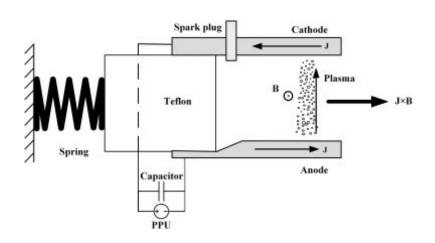
SUCHAI-2 and 3 At present under construction

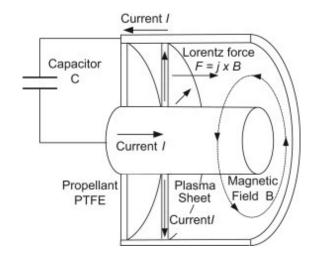
M. Diaz et al., Advances in Space Research, 58, 2134-2147 (2016)



Typical PPT diagram







parallel plate electrodes

Coaxial electrodes



Thrust estimations from our previous world in PF devices

On the one hand, to obtain an estimation for a miniature plasma thruster operating with an energy of the order of 1J, based on plasma focus technology and its scaling laws, we can assume an ejected mass m_e ~ 4 x 10⁻¹³ kg with a velocity v_e ~ 5 x 10⁵ m/s. Thus, an impulse bit

 $I_{bit} = \Delta p = m_e v_e \sim 2 \times 10^{-7} \text{ Ns}$ is estimated.

Considering that the mean propulsion force during a second as <F> = I_{bit} f, with f the operation frequency, the mean thrust could be **0.2**, **2 and 20 µN** for an operation frequency of **1, 10 and 100 Hz** respectively.

 On the other hand, from electromagnetic estimation for a coaxial plasma gun (axial phase of a plasma focus), imposing the condition that the plasma reaches the end of the electrodes coincident with maximum current, i.e. at a time of quarter of period of the discharge,

 $\Delta p = F_{mag}(\tau/4) = (\mu_0/4) I^2 \ln(b/a) (LC)^{1/2}$, with I the peak current, C de capacitance of the capacitor, L the total inductance, a anode radius and b cathode radius. Using I = V (C/L)^{1/2},

$$\Delta p = (\mu_0/4) V^2 \ln(b/a) C^{3/2} L^{-1/2}$$

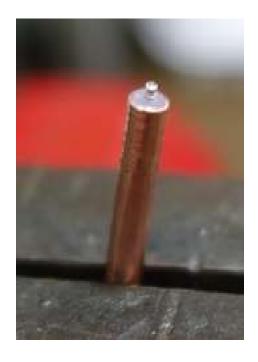
For a device with a capacitor of 225nF charging at 3kV an energy stored of 1J is achieved. Assuming a=0.5mm and b=1.25mm, and 5nH of inductance (that is possible achieved in compact devices, like Nanofocus designed and built at CCHEN a value for $\Delta p \sim 3.8 \times 10^{-6} \, \text{Ns}$ is obtained. Thus, with 1, 10 and 100 Hz, the mean thrust could be 3.8, 38 and 380 μ N respectively.

Both estimations are consistent with the literature for orientation systems for CubeSats, and are enough to support our hypothesis and encourage to pursue a research project.



Experiements on PPT in modified Nanofocus







Electrodes

 R_{ci} : 1.1 mm, R_{ce} : 0.85 mm

R_a: 0.35 mm

Modified Nanofocus with PPT electrodes $p \sim 10^{-4} \text{ mbar}$

"Pulsed Plasma Thruster Based On Ultra-miniaturized Plasma Focus" L. Soto, J. Pedreros, R. Silva, P. Maldonado, G. Avaria, C. Pavez, J. Moreno, and M. Diaz, 19th International Congress on Plasma Physics, ICPP 2018, Vancouver, Canada, June 2018.



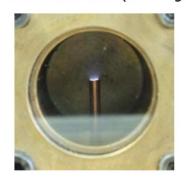
Experiements on PPT in modified Nanofocus

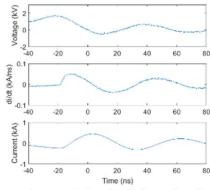


A) Coaxial plasma gun with cathode and anode extended (as in figure 1b)







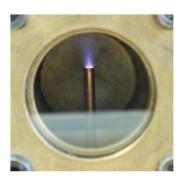


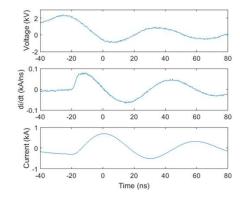
Left: Photograph of plasma gun. Center: plasma discharge. Right: voltage, current derivative and current signals. Voltage breakdown $1.75\pm0.2 \text{ kV}$.

B) Coaxial plasma gun with cathode extended









Left: Photograph of plasma gun. Center: plasma discharge. Right: voltage, current derivative and current signals. Voltage breakdown 2.2±0.1 kV.

"Pulsed Plasma Thruster Based On Ultra-miniaturized Plasma Focus" L. Soto, J. Pedreros, R. Silva, P. Maldonado, G. Avaria, C. Pavez, J. Moreno, and M. Diaz, 19th International Congress on Plasma Physics, ICPP 2018, Vancouver, Canada, June 2018.



Next Lecture



How to build a small Plasma Focus Recipes and tricks