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*Empirical study of three-body interactions
in semi-magic nuclei*

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Experiment, Theory and Evaluation*

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Three-body interactions

The Hamiltonian with up to three-body interactions, acting between identical nucleons in a single- j orbital:

$$\hat{H} = \hat{H}_1 + \hat{V}_2 + \hat{V}_3$$

The matrix elements of two and three-body interactions in the case of two and three nucleons in the j -orbital can be given by :

$$V_J = \langle j^2 JM | \hat{V}_2 | j^2 JM \rangle \quad ; \quad W_{J\alpha\alpha'} = \langle j^3 \alpha JM | \hat{V}_3 | j^3 \alpha' JM \rangle$$

Matrix elements

The matrix elements of a scalar k-body interaction V_k between two states belonging to the j^n configuration of n nucleons in the j orbital can be calculated from the recurrence relation:

$$\begin{aligned} & \langle j^n \alpha JM | \hat{V}_k | j^n \alpha' J M \rangle \\ &= \frac{n}{n-k} \sum_{\alpha_1 \alpha'_1 J_1} [j^n \alpha J \{ |j^{n-1}(\alpha_1 J_1) j J \} | j^n \alpha' J] \\ & \quad \times \langle j^{n-1} \alpha_1 J_1 | \hat{V}_k | j^{n-1} \alpha'_1 J_1 \rangle \end{aligned}$$

The coefficient of fractional parentage (cfp) can be determined from the following expression :

$$\begin{aligned} & n [j^{n-1}(\alpha_0 J_0) j J | \{ j^n [\alpha_0 J_0] J \} | j^n [\alpha_0 J_0] J] \\ &= \delta_{\alpha_1 \alpha_0} \delta_{J_1 J_0} + (n-1) \sum_{\alpha_2 J_2} (-1)^{J_0+J_1} \sqrt{(2J_0+1)(2J_1+1)} \\ & \quad \times \left\{ \begin{matrix} J_2 & j & J_1 \\ J & j & J_0 \end{matrix} \right\} [j^{n-2}(\alpha_2 J_2) j J_0 | \{ j^{n-1} \alpha_0 J_0 \} | j^{n-1} \alpha_0 J_0] \\ & \quad [j^{n-2}(\alpha_2 J_2) j J_1 | \{ j^{n-1} \alpha_1 J_1 \} | j^{n-1} \alpha_1 J_1] \end{aligned}$$

Fit

W J Huang, G Audi, M Wang, F G Kondev, S Naimi and X Xu, Chin Phys C 41 (2017) 030002
 +
 ENSDF

$$\sigma = \sqrt{\frac{1}{N} \sum_k [E_{exp}(k) - E_{th}(k)]^2}$$

Root mean square deviation

- $j = 9/2$, $N = 50$, $40 < Z < 50$

44 measured values of binding energies.

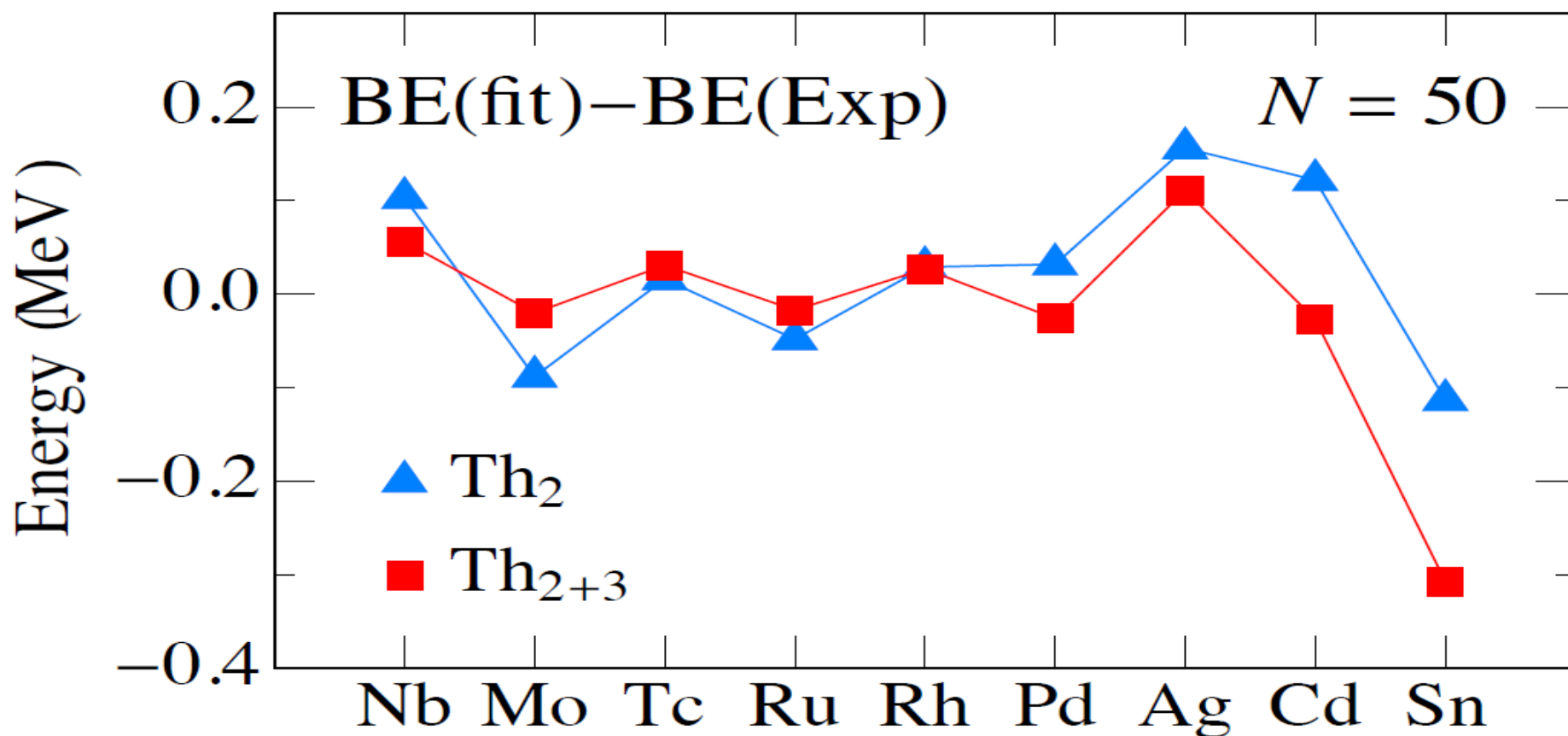
Two-body interaction : $\sigma = 63 \text{ keV}$

+

Three-body interaction $\sigma = 46 \text{ keV}$

The three-body interaction
 meaningfully improves the
 results.

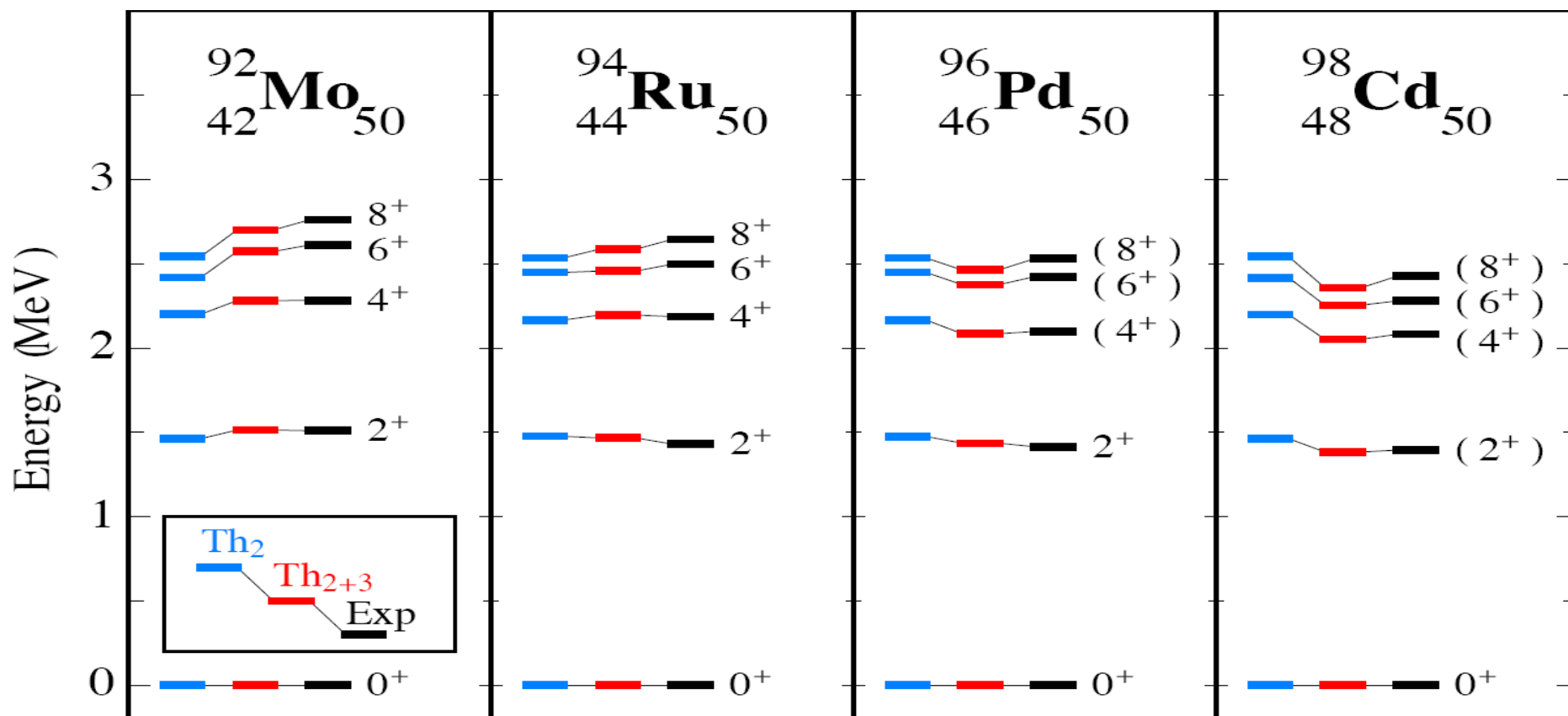
Ground states:



The difference between calculated and experimental ground-state binding energies of $N = 50$ isotones..

- The addition of a three-body interaction does not improve the description of the binding energies of ground states.

Excitation spectra of the even-mass



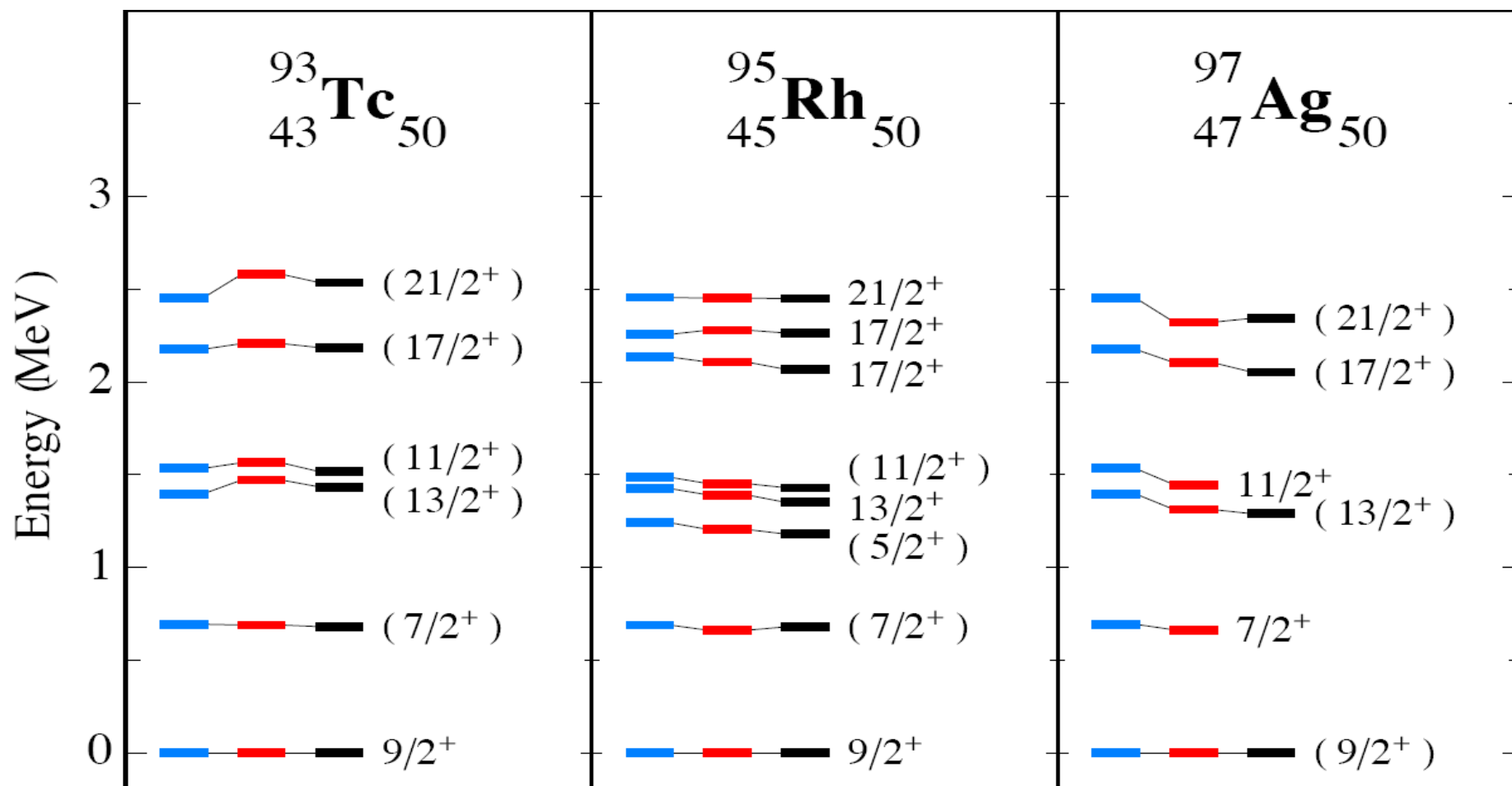
Excitation spectra of the even-mass $N=50$ isotones.

The three-body interaction meaningfully improves the results.

Conservation of the particle–hole symmetry using the two-body interactions.

This symmetry was broken by adding the effect of the three-body interactions.

Excitation spectra of the odd-mass

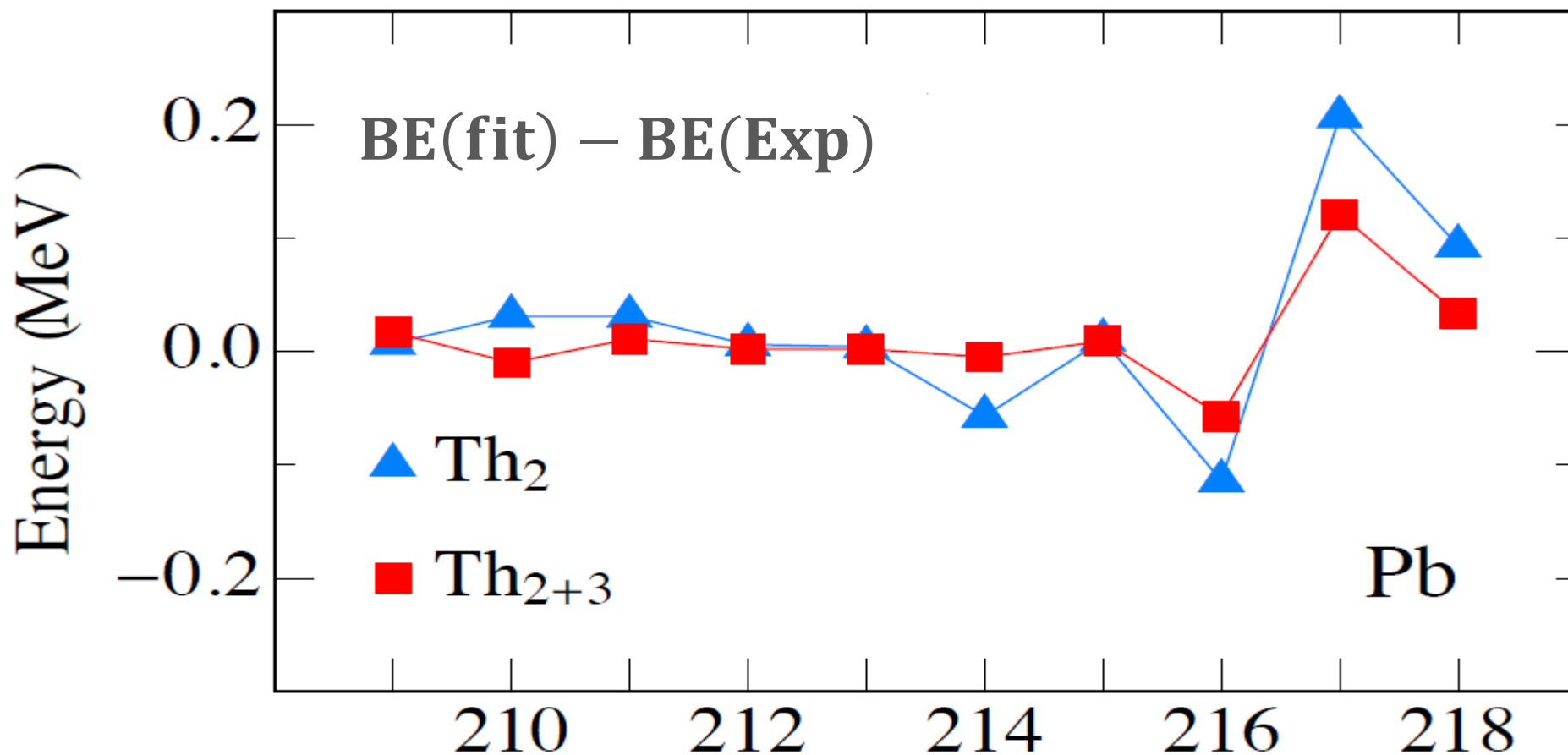


Excitation spectra of the odd-mass $N=50$ isotones.

- A good agreement between theory and experience is obtained by adding the effect of the three-body interaction.

Lead isotopes:

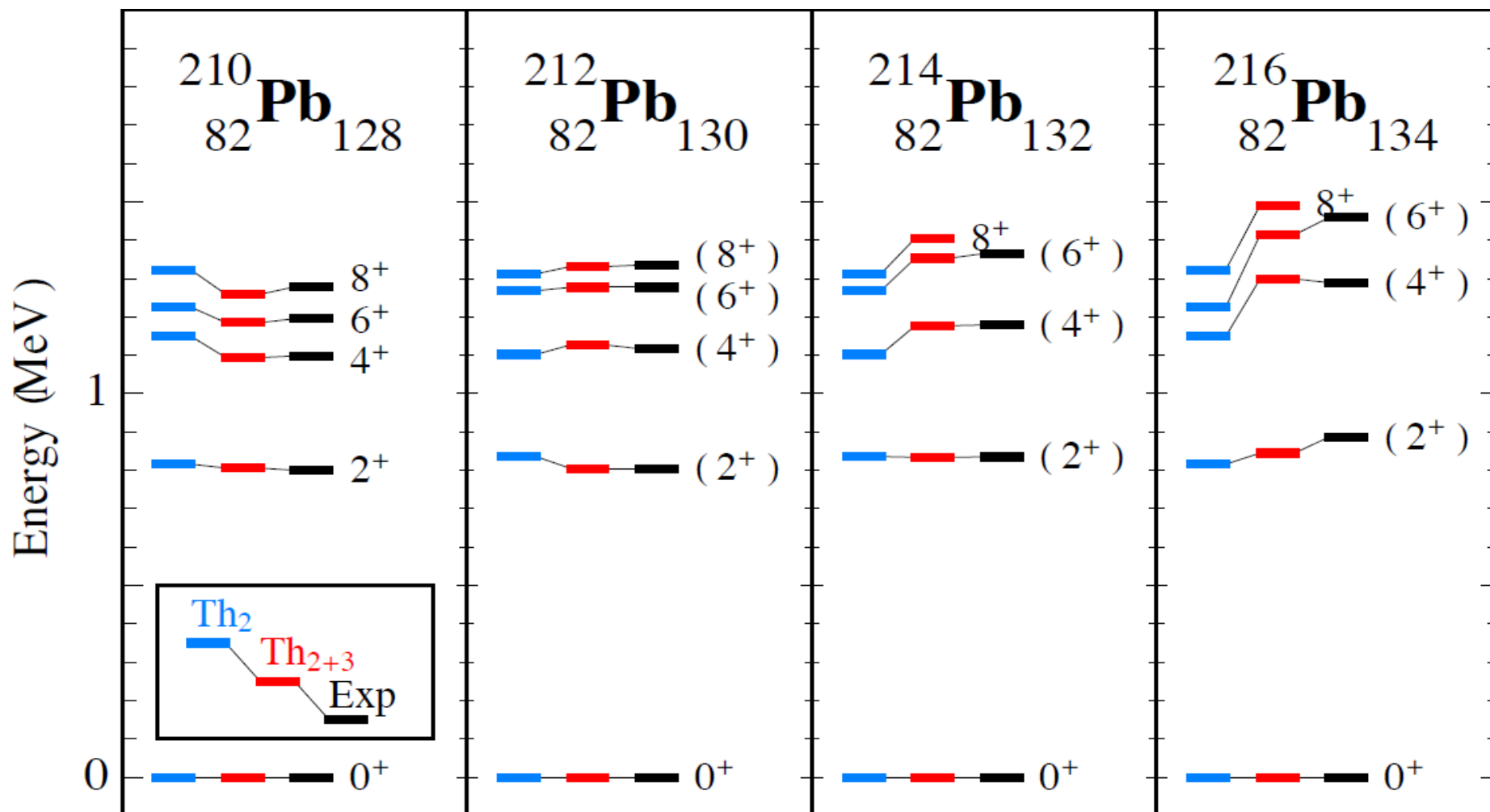
Binding energies of the ground states



The difference between calculated and experimental ground-state binding energies of Pb isotopes.

The addition of a three-body interaction does not notably improve the quality of the fit.

Excitation spectra of the even-mass

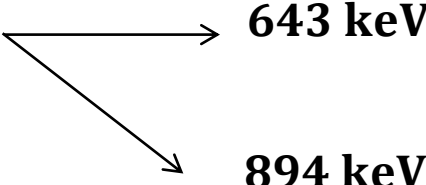


Excitation spectra of the even-mass Pb isotones.

Breaking of the particle–hole symmetry by adding the three-body component of the interactions.

Improved results following the inclusion of the effect of the 3-body interaction.

$N = 50$ isotones: 6^+ level ^{94}Ru

For the ^{211}Pb : two states $11/2^+$ 

$\left\{ \begin{array}{l} \sigma - \text{minimum} \\ \sigma_2 - \text{minimum} \end{array} \right.$

- By adding a new level to the input data, one observes stability of the Th_2 and large fluctuations of the Th_{2+3} matrix elements.

The separate components of the three-body interaction could not be determined reliably from an empirical fit : A small variation in the input data leads to large fluctuations in the three-body matrix elements.

Thank you

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