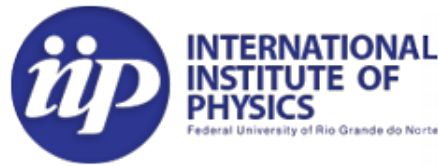


Multi-Qubit Robustness by Local Encoding



Quantum *I*nformation and Quantum *M*atter Group
www.iip.ufrn.br/qiqm

Rafael Chaves

Conference on Quantum Measurement, Trieste 2019



Few words about Natal...



What is this talk about?

Quantum correlations can be fragile...

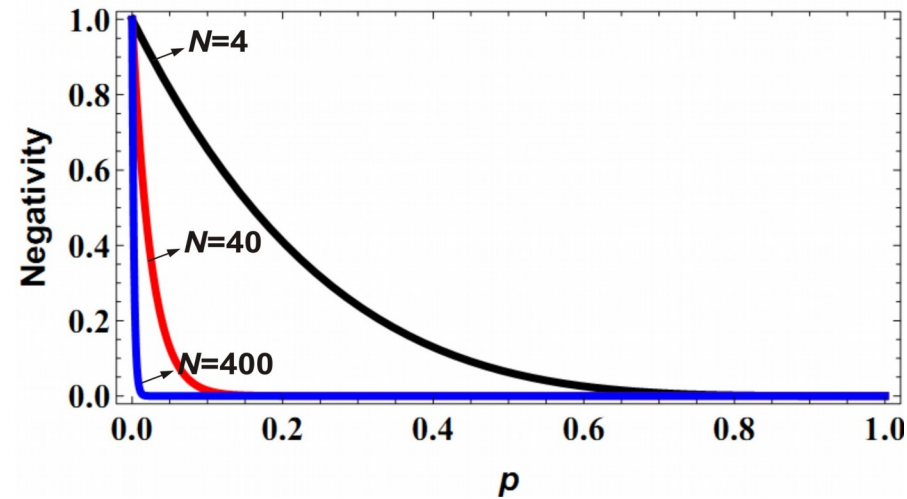
$$|GHZ^N\rangle = \sqrt{\frac{1}{2}} \left(|0\rangle^{\otimes N} + |1\rangle^{\otimes N} \right)$$



$$E(\rho_{GHZ}^N(p)) \leq (1-p)^N \cdot E(\rho_{GHZ}^N(0))$$

***[Aolita, RC, Cavalcanti, Acin, Davidovich PRL 2008]**

***[Aolita, Melo, Davidovich, RPP 2015]**

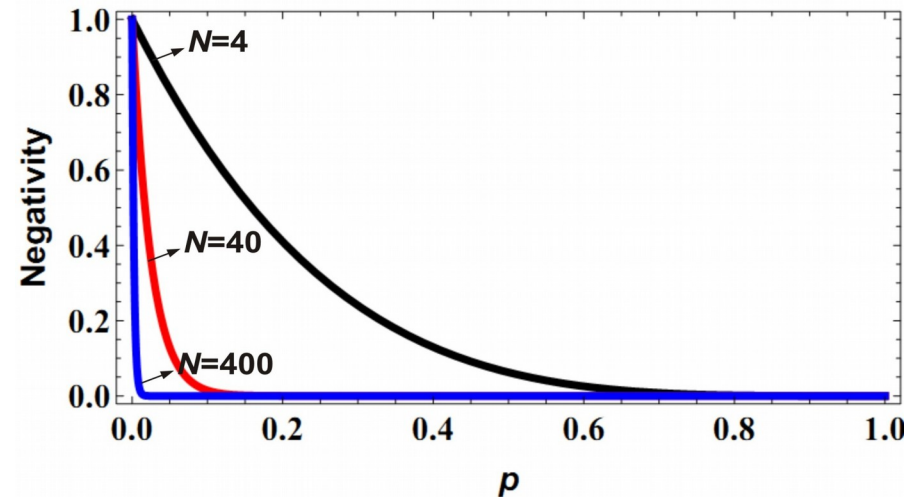


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*[Aolita, RC, Cavalcanti, Acin, Davidovich PRL 2008]

*[Aolita, Melo, Davidovich, RPP 2015]

?

Fragile Entanglement = Fragile resource

Quantum Metrology →

$$\delta\lambda \geq \frac{1}{2\sqrt{\nu\mathcal{F}}}$$

Apparently, any full-rank noise will spoil scaling quantum advantages.

$$\begin{aligned} \mathcal{F}(\rho_{sep}) &\leq N \\ \mathcal{F}(\rho_{GHZ}^N(\rho)) &= (1-p)^{2N} N^2 \end{aligned}$$

*[Huelga et al., PRL 1997]

*[Escher et al. NatPhys 2011]

*[Demkowicz-Dobrzanski et al. NatComm 2012]

Quantum correlations can be made robust...

Solution 1 (general state, any noise)

Quantum error-correction codes

- Error Syndrome
- Feedback correction

Quantum correlations can be made robust...

Solution 1 (general state, any noise)

Quantum error-correction codes

- Error Syndrome
- Feedback correction

Solution 2 (GHZ state, white noise)

Encode each logical qubit into a GHZ state

*[Frowis, Duer, PRL 2011]

- Protection is achieved passively
- Experimental overhead

$$\begin{array}{lcl} |0\rangle & \rightarrow & |GHZ_+^M\rangle = |0\rangle^{\otimes M} + |1\rangle^{\otimes M} \\ |1\rangle & \rightarrow & |GHZ_-^M\rangle = |0\rangle^{\otimes M} - |1\rangle^{\otimes M} \end{array}$$

Decay rate of entanglement is still

$$(1-p)^N$$

but now p exponentially decreases with M .

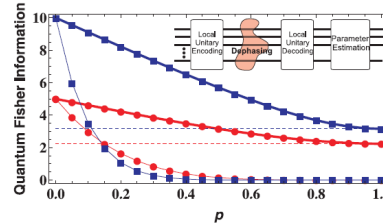
What about noise with a privileged direction?

- Trapped-ion experiments
- Photonic polarization-qubits

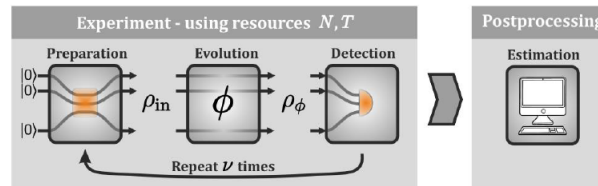
GHZ entanglement can be made exponentially more robust

Outline

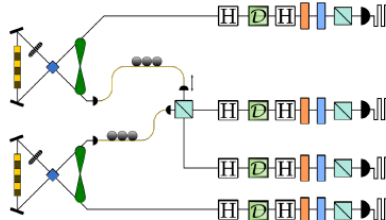
- Experimentally friendly robust GHZ entanglement



- Applications in Metrology



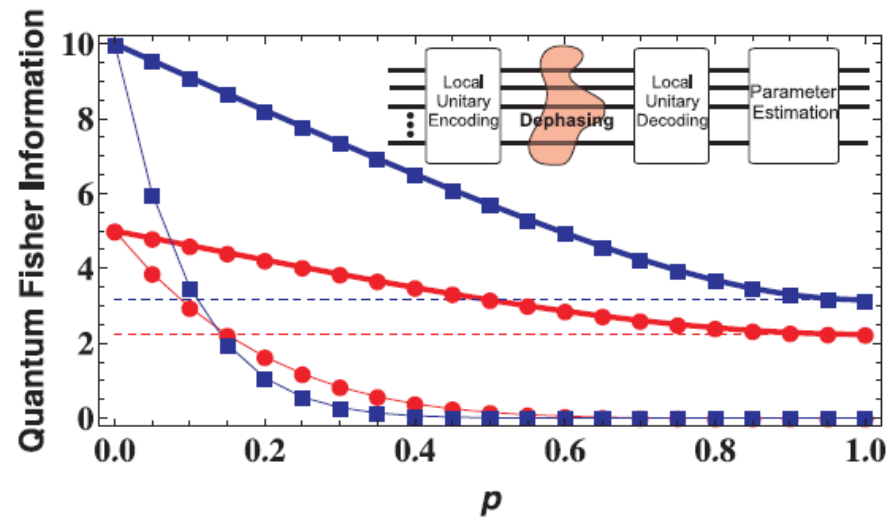
- An experimental investigation



- Some puzzling open questions

Experimentally friendly robust GHZ entanglement

*[RC, Aolita, Acin, PRL 2012]

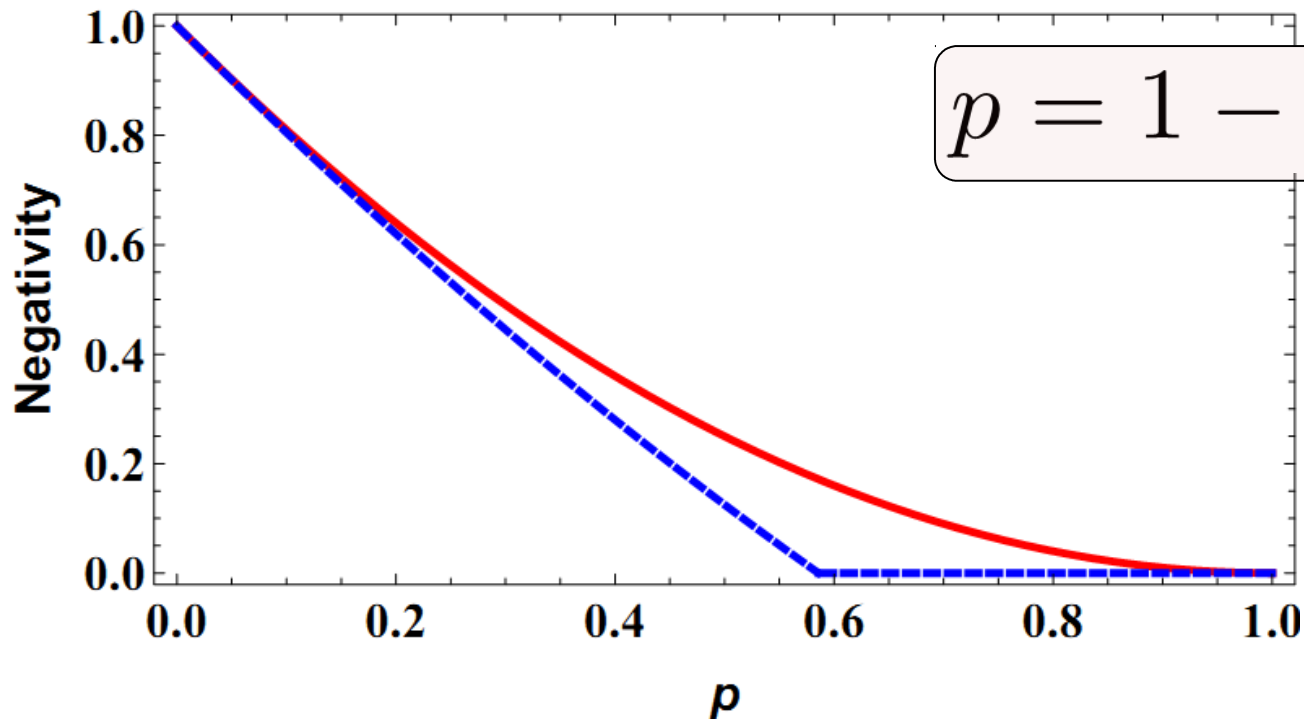


The basic idea

$$\varepsilon(\rho) \rightarrow \left(1 - \frac{p}{2}\right) \rho + \frac{p}{2} Z \rho Z$$

$$|\Phi^+\rangle = |00\rangle + |11\rangle$$

$$(I \otimes H) |\Phi^+\rangle = |0+\rangle + |1-\rangle$$



What is the most robust local encoding against dephasing?

$$\varepsilon(\rho) \rightarrow \left(1 - \frac{p}{2}\right) \rho + \frac{p}{2} Z \rho Z$$

$$|\Phi_+^N\rangle \doteq \frac{1}{\sqrt{2}}(|0\rangle^{\otimes N} + |1\rangle^{\otimes N})$$

- Considering the negativity and numerical optimizations up to $N=10$ show that the optimal states are the transversal states.

$$|\Phi_{+T}^N\rangle \doteq H^{\otimes N} |\Phi_+^N\rangle = \frac{1}{\sqrt{2}}(|+\rangle^{\otimes N} + |-\rangle^{\otimes N})$$

$$|\pm\rangle \doteq \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$$

- Analytically we can prove that for any entanglement measure

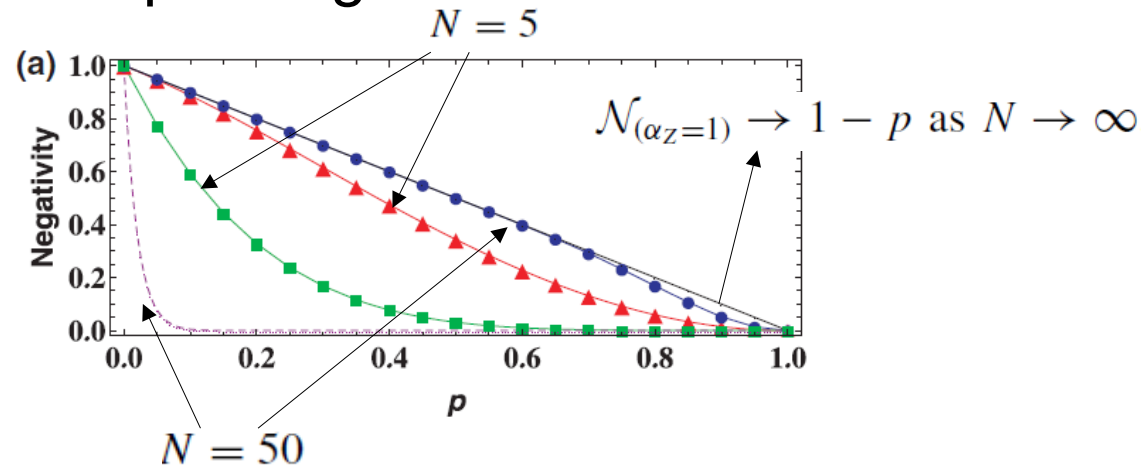
$$\begin{aligned} E(\Lambda^{\text{PD}}(|\Phi_{+T}^N\rangle)) &\geq E(\Lambda^{\text{PD}}(|\Phi_{+T}^{N-1}\rangle)) \\ &\geq \dots \geq E(\Lambda^{\text{PD}}(|\Phi_{+T}^2\rangle)) \end{aligned}$$

- N entangled qubits are at least as robust as 2 entangled ones!**

What is the most robust local encoding against dephasing?

Pure Dephasing

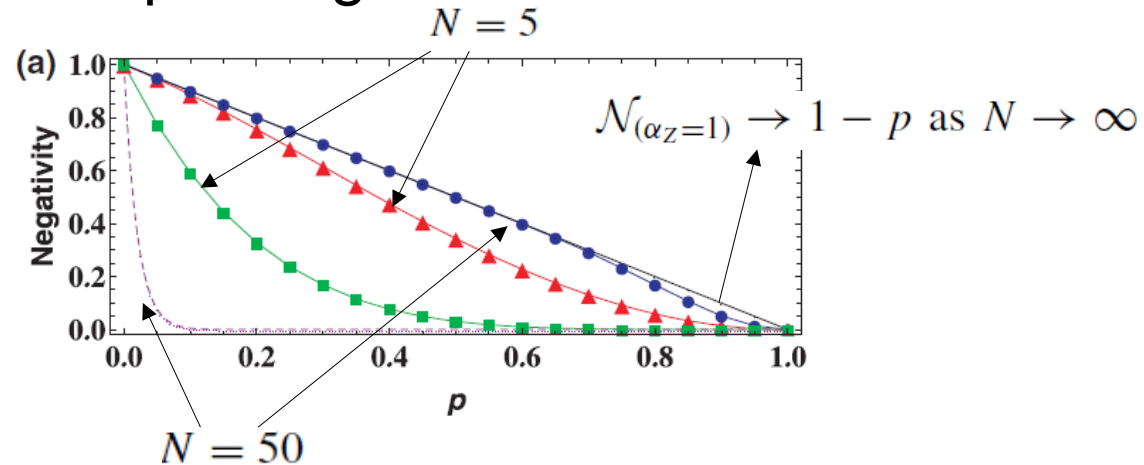
$$\varepsilon(\rho) \rightarrow \left(1 - \frac{p}{2}\right) \rho + \frac{p}{2} Z \rho Z$$



What is the most robust local encoding against dephasing?

Pure Dephasing

$$\mathcal{E}(\rho) \rightarrow \left(1 - \frac{p}{2}\right) \rho + \frac{p}{2} Z \rho Z$$

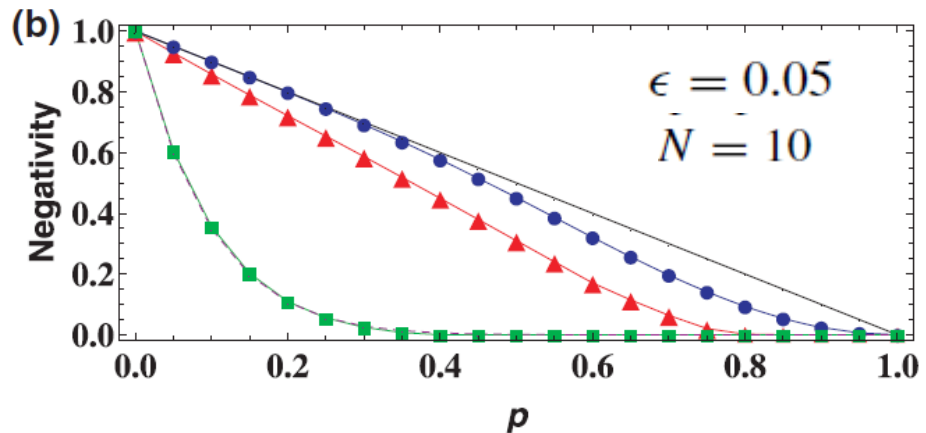


Deviations from pure dephasing

$$\mathcal{E}(\rho) \doteq \left(1 - \frac{p}{2}\right) \rho + \frac{p}{2} (\alpha_X X \rho X + \alpha_Y Y \rho Y + \alpha_Z Z \rho Z),$$

$$\epsilon \doteq \alpha_X + \alpha_Y$$

$$\alpha_Z = 1 - \epsilon$$



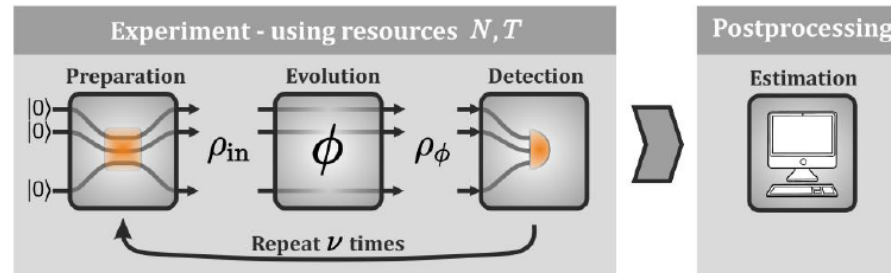
$$\mathcal{N} \approx 1 - p(N\epsilon + 1 - \epsilon) \approx (1 - p)^{N\epsilon + 1 - \epsilon}$$

Applications in Metrology

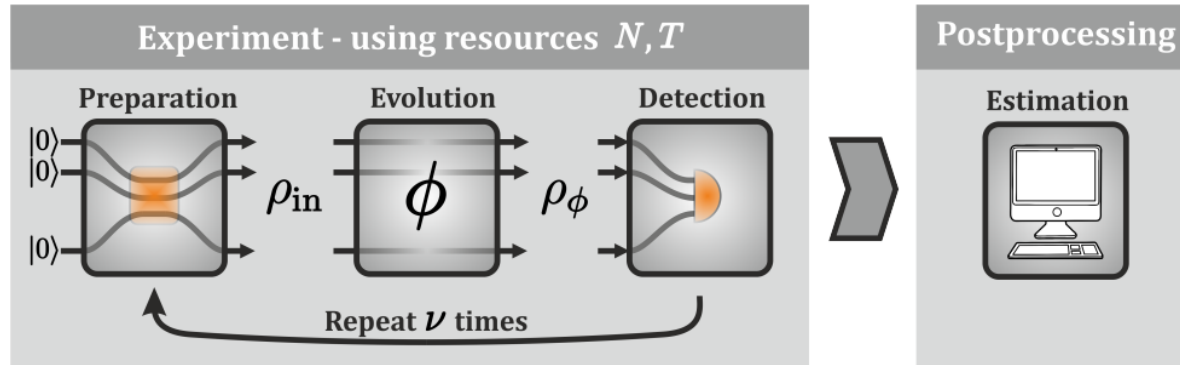
***[RC, Aolita, Acin, PRA 2012]**

***[RC, Brask, Marki Acin, PRL 2013]**

***[Brask, RC, Markiewicz, Kolodynski, PRX 2014]**



Single-Parameter Estimation



Typical experiment: [preparation+interaction+measurement]+postprocessing

As resources we take the **probe size** and the **total time N, T**

- Classical strategies based on separable states, the central limit theorem bounds the precision, the **standard quantum limit (SQL)**

$$\delta\omega \geq c_{cl} \frac{1}{\sqrt{N\nu}}$$

- For entangled states, the precision is limited by the quantum uncertainty relation, the **Heisenberg bound**

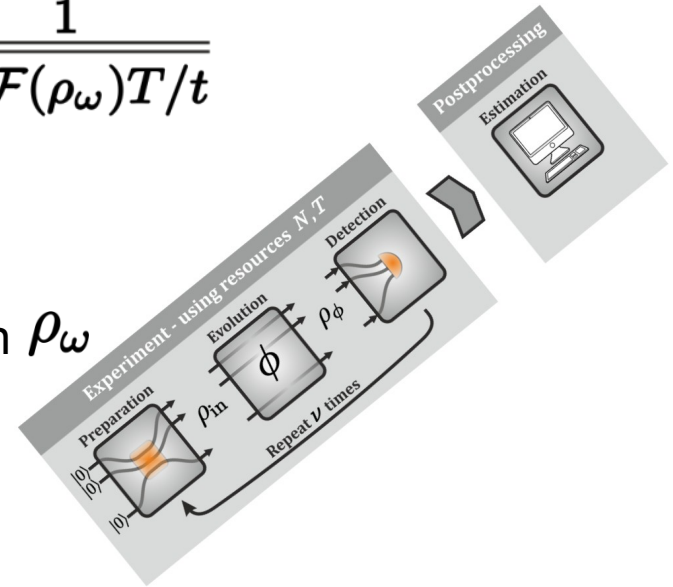
$$\delta\omega \geq c_q \frac{1}{N\sqrt{\nu}} \quad \nu = T/t$$

Computing the precision

The attainable precision can be bounded by the quantum Fisher information

$$\delta\omega \geq \frac{1}{\sqrt{\mathcal{F}(\rho_\omega)\nu}} = \frac{1}{\sqrt{\mathcal{F}(\rho_\omega)T/t}}$$

$\mathcal{F}(\rho_\omega)$ quantifies the information about ω encoded in ρ_ω



The optimal precision requires optimization over inputs

To avoid the optimization we can compute

$$\max_{\rho_{in}} \mathcal{F}(\rho_\omega)$$

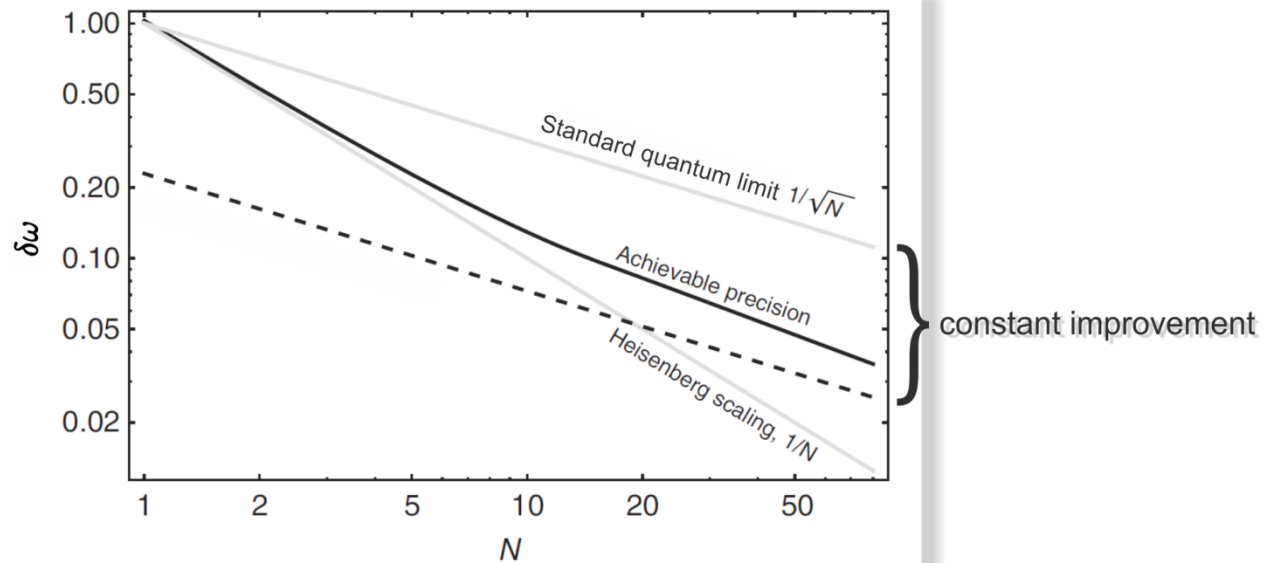
hard!

- Fisher information of a **specific state**
 - gives lower bound on the optimum
- **Bounds** that require no or simpler optimization
 - gives upper bounds on the optimum

Metrology versus noise

- The attainable precision depends strongly on the noise model
 - difficult to obtain general results
- Result indicate that even for arbitrarily small amounts of noise, quantum strategies provide only a **constant improvement over the SQL**

Demkowicz-Dobrzanski *et al.*,
Nat. Comm. 3, 1063 (2012)



- Full rank noise independent of the probe size: SQL like scaling (Fujiwara, Imai, 2008)
- SQL-like scaling for: lossy optical interferometry, atomic spectroscopy under dephasing or spontaneous emission (Knysh *et al.*, 2010, Escher *et al.*, 2011, Demkowicz-Dobrzanski *et al.*, 2012)

Noisy Frequency Estimation Model

We use a master equation description

$$\frac{\partial \rho(t)}{\partial t} = \mathcal{H}(\rho) + \mathcal{L}(\rho)$$

- Unitary Evolution**

$$H = \frac{\omega}{2} \sum_{k=1}^N \sigma_z^k$$

frequency to
be estimated

$$\mathcal{H}(\rho) = -i [H, \rho]$$

- Noise**

$$\mathcal{L}(\rho) = -\frac{\gamma}{2} \left[\rho - \sum_{k=1}^N \alpha_x \sigma_x^k \rho \sigma_x^k - \alpha_y \sigma_y^k \rho \sigma_y^k - \alpha_z \sigma_z^k \rho \sigma_z^k \right]$$

overall noise strength

transversal noise

parallel noise

$\sum_i \alpha_i = 1$
 $\alpha_i \geq 0$

$\alpha_z = 1$ ➡ Parallel noise, commutes, studied before

$\alpha_{x,y} > 0$ ➡ Does not commute, noise is full rank

Beating the shot noise limit

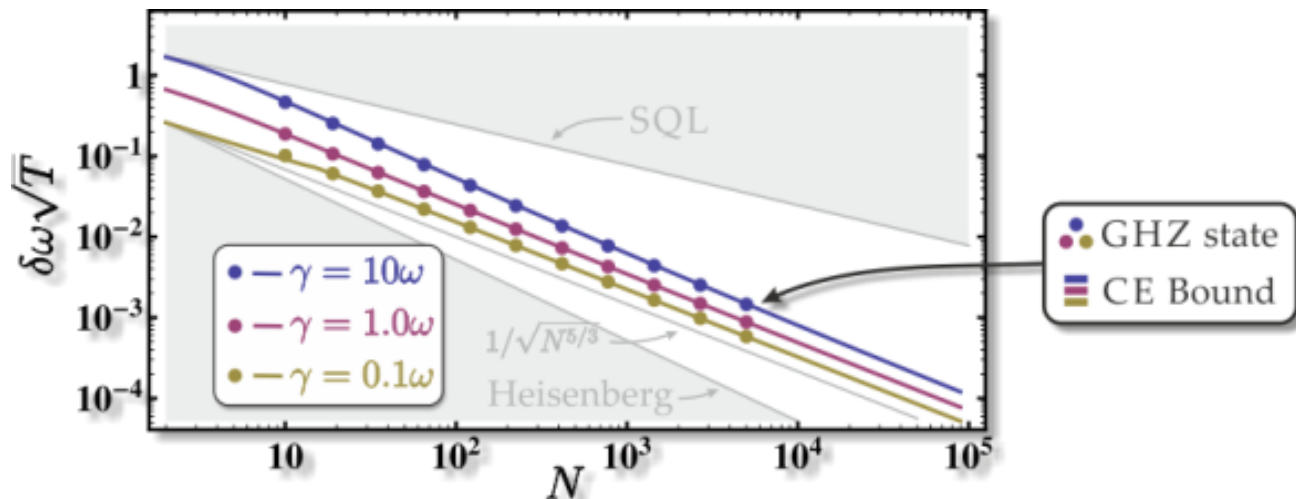
We cannot derive analytical bounds – but we can compute the finite-N Channel extension numerically

$$\delta\omega \sqrt{T} \geq \sqrt{\frac{c_x(\gamma)}{N^{5/3}}} \quad (\text{based on numerics})$$

$$c_x(\gamma) = \frac{3^{2/3}}{2} \gamma^{1/3}$$

$$t_x^{\text{opt}} \xrightarrow{N \rightarrow \infty} 0$$

We can compute the Fisher Information for the GHZ state analytically and optimize the evolution time numerically



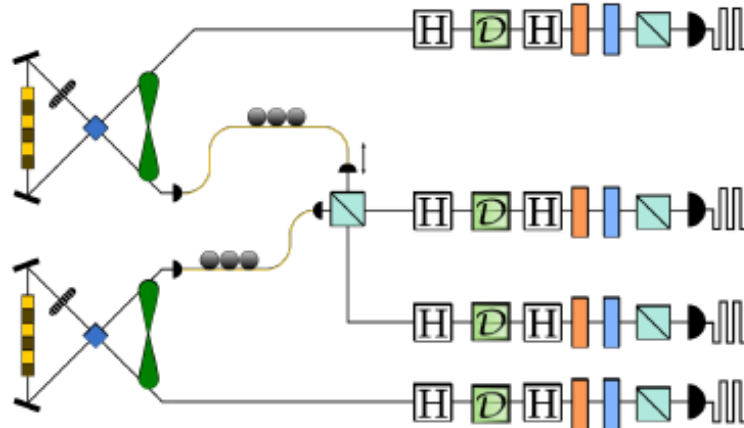
• *The bound is tight*

The GHZ state saturates the bound! 

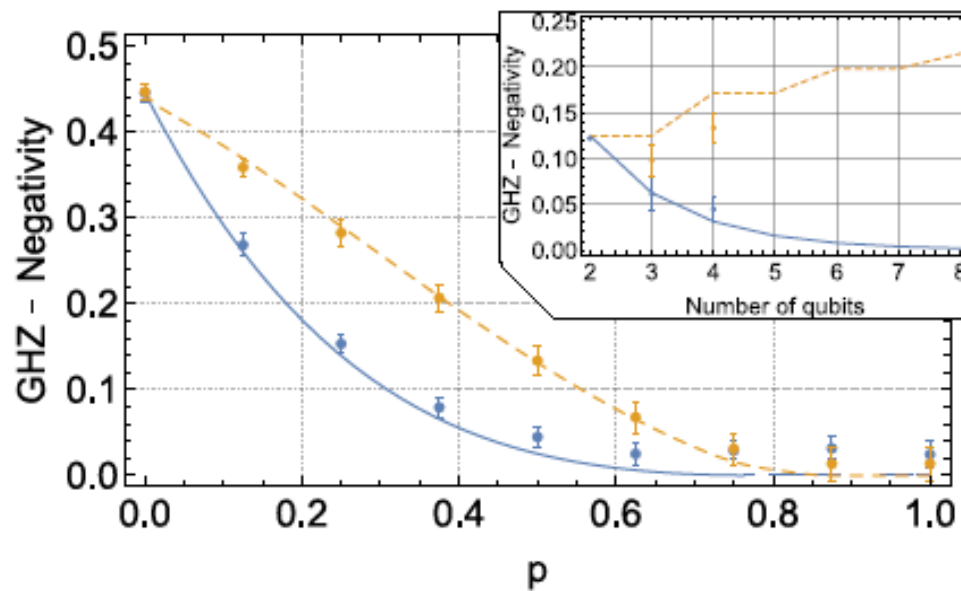
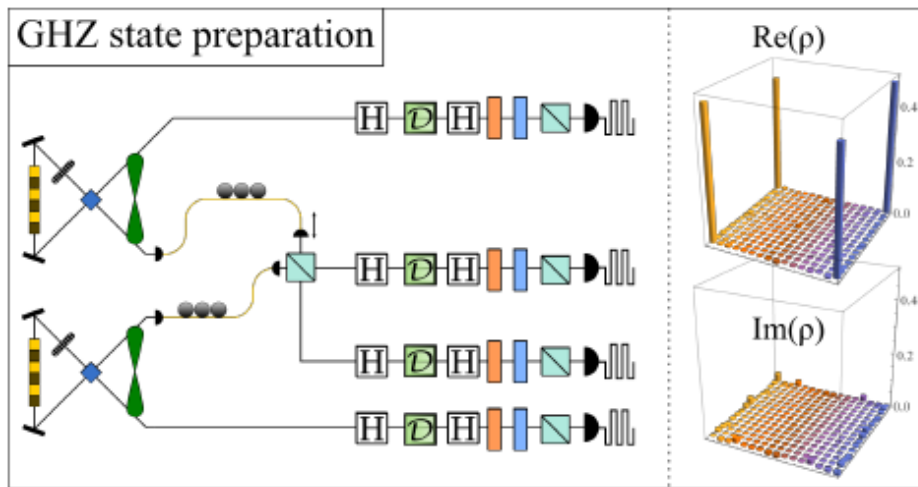
• *Super-classical precision scaling for full-rank noise is indeed possible*

Experimental Investigation

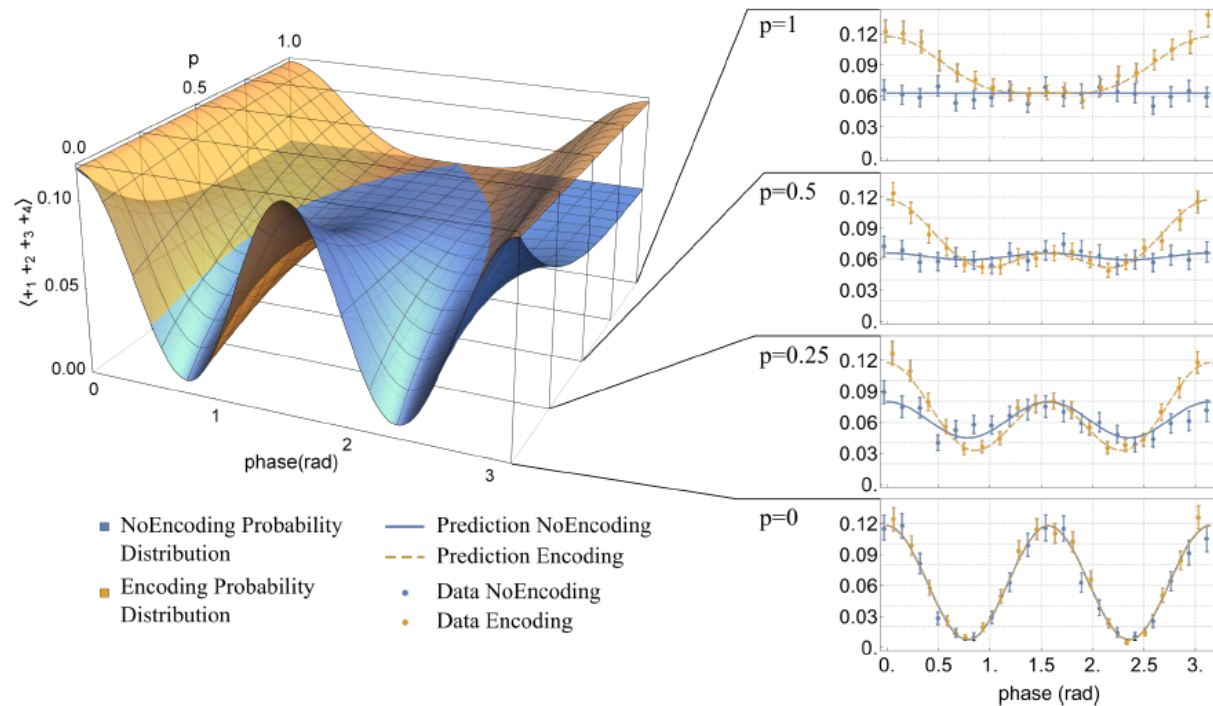
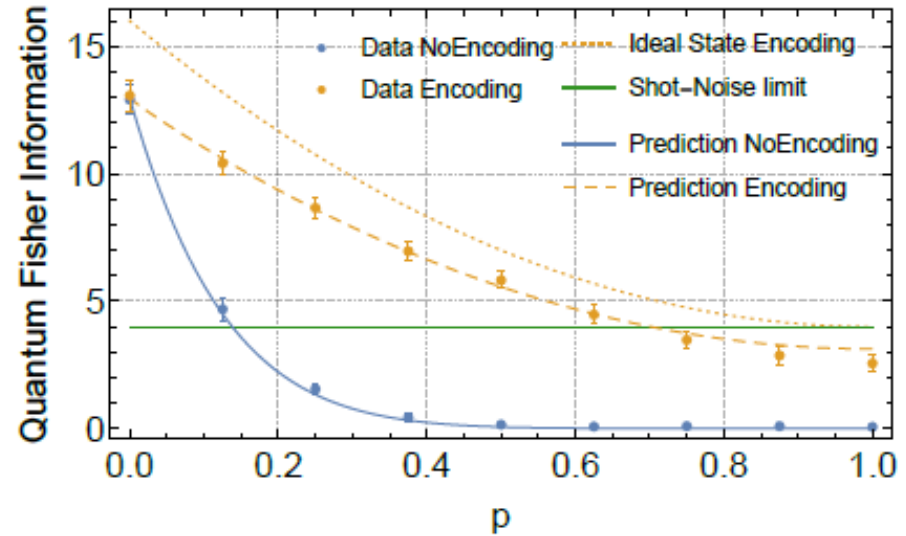
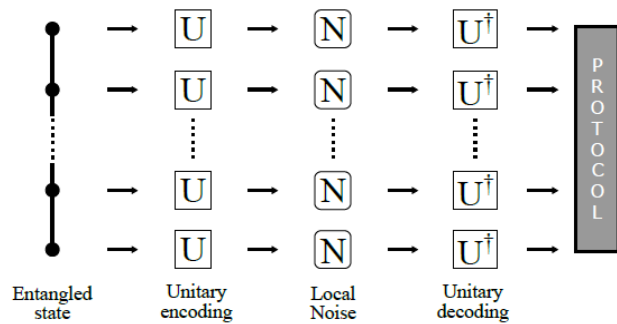
*[M. Proietti, M. Ringbauer, F. Graffitti, P. Barrow, A. Pickston, D. Kundys, D. Cavalcanti, L. Aolita, R. Chaves, and A. Fedrizzi, arXiv: 1903.08667]



Experimentally Robust GHZ entanglement



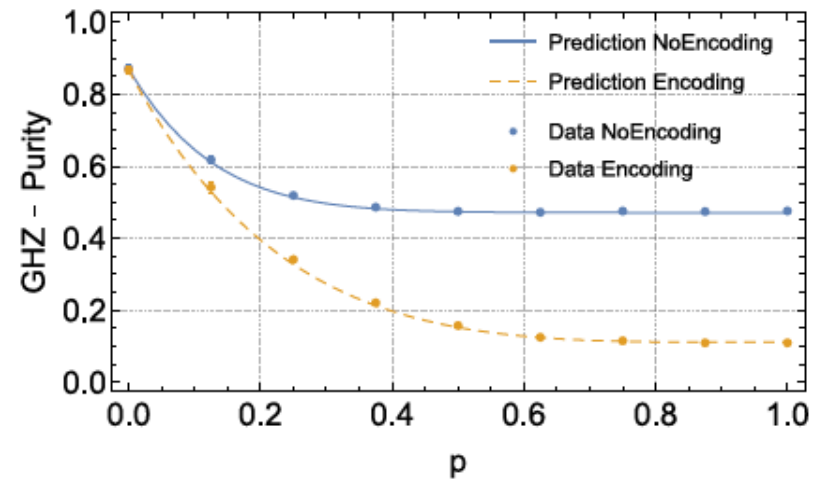
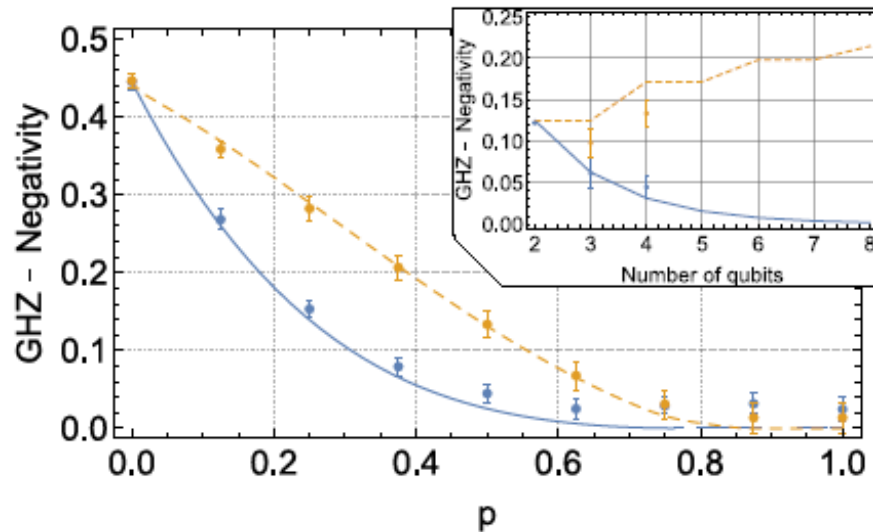
Experimentally Robust metrology



A puzzling question



Puzzle



Most robust state is the one generating far more entanglement with the environment.

What to remember (if anything)

- **i) If you have noise with privileged direction local encoding might improve a lot the robustness of correlations**
- **ii) This robustness is also reflected in the use of these correlations as a resource: metrology, Bell inequalities violation, communication complexity problems, etc**
- **iii) Is there a principle underlying what is the most robust local basis?**