Topological transition in weak-measurement-induced geometric phases

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מכוז ויצמו למדע

WEIZMANN INSTITUTE OF SCIENCE

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Project participants

Valentin Gebhart **Thomas Wellens** Andreas Buchleitner



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Alessandro Romito

Lancaster 🌌 University



Gebhart, KS, Wellens, Buchleitner, Romito, Gefen, arXiv:1905.01147



Outline

- 1. Berry phase and
 - Pancharatnam (measurement-induced) phase
- 2. Weak measurements and their back action

- 3. Weak-measurement-induced geometric phase
- 4. The topological transition

Berry phase

 $H(\mathbf{R})|n(\mathbf{R})\rangle = E_n(\mathbf{R})|n(\mathbf{R})\rangle$

R(t) — slowly-changing $R(0) = R(T) = R_0$

$$|\psi(0)\rangle = |n(\mathbf{R}_0)\rangle \rightarrow |\psi(t)\rangle = e^{i\varphi}|n(\mathbf{R}_0)\rangle$$

$$\varphi = \varphi_d + \varphi_B,$$

$$\varphi_d = -\int_0^T E_n dt,$$

$$\varphi_B = i \oint d\mathbf{R} \langle n(\mathbf{R}) | \partial_{\mathbf{R}} | n(\mathbf{R}) \rangle$$

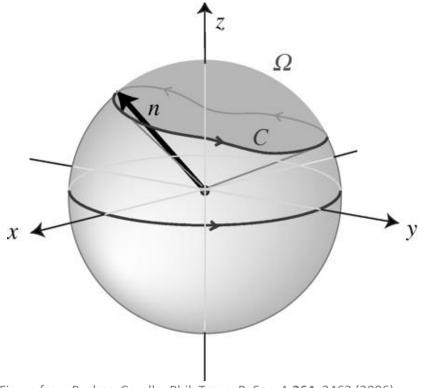


Figure from Pachos, Carollo, Phil. Trans. R. Soc. A 364, 3463 (2006)

Berry, 1984

Berry phase for spin 1/2

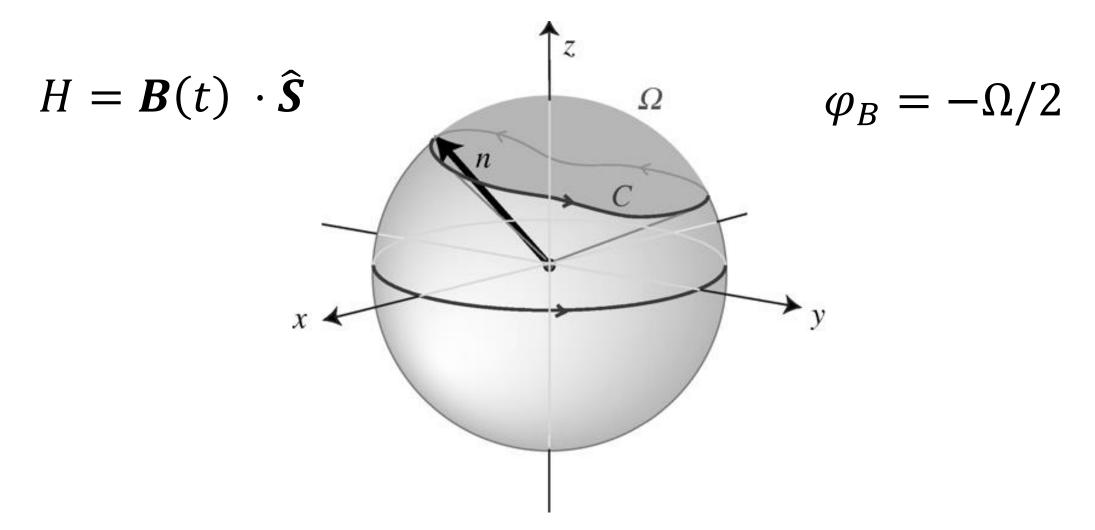


Figure from Pachos, Carollo, Phil. Trans. R. Soc. A 364, 3463 (2006)

Berry, 1984

Pancharatnam phase and measurement

A sequence of states: $|\psi_2\rangle$ $|\psi_0\rangle, |\psi_1\rangle, ..., |\psi_{N-1}\rangle$ $|\psi_0\rangle$ $P_k = |\psi_k\rangle \langle \psi_k|$ $|\psi_1\rangle$...defines a phase φ_P : $e^{i\varphi_P}\sqrt{P} = \langle \psi_0 | P_{N-1} \dots P_2 P_1 | \psi_0 \rangle$ $= \langle \psi_0 | \psi_{N-1} \rangle \langle \psi_{N-1} | \psi_{N-2} \rangle \dots \langle \psi_2 | \psi_1 \rangle \langle \psi_1 | \psi_0 \rangle$

Pancharatnam, 1956

Berry phase = Pancharatnam phase

$$|\psi_{k}\rangle = |n(\mathbf{R}_{k})\rangle$$

$$|\psi_{k+1}\rangle = |\psi_{k}\rangle + |\delta\psi_{k}\rangle$$

$$|\delta\psi_{k}\rangle \approx (\mathbf{R}_{k+1} - \mathbf{R}_{k})\partial_{\mathbf{R}_{k}}|n(\mathbf{R}_{k})\rangle$$

$$\langle\psi_{k+1}|\psi_{k}\rangle \approx$$

$$\approx \exp(-\delta\mathbf{R}\langle n(\mathbf{R}_{k})|\partial_{\mathbf{R}_{k}}|n(\mathbf{R}_{k})\rangle)$$

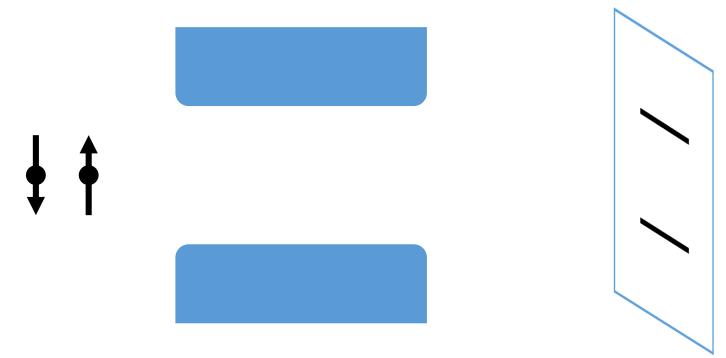
$$\psi_{0}|\psi_{N-1}\rangle\langle\psi_{N-1}|\psi_{N-2}\rangle \dots \langle\psi_{2}|\psi_{1}\rangle\langle\psi_{1}|\psi_{0}\rangle =$$

$$= \exp(-\oint d\mathbf{R}\langle n(\mathbf{R})|\partial_{\mathbf{R}}|n(\mathbf{R})\rangle) = \exp(i\varphi_{B})$$

cf. Chruscinski, Jamiolkowski "Geometric phases in classical and quantum mechanics"

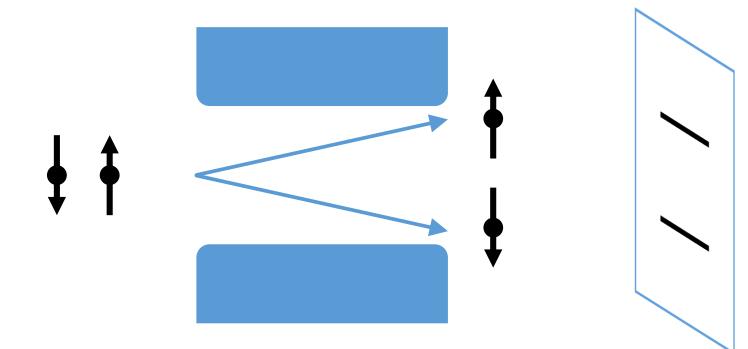
(projective) Quantum measurements

Stern–Gerlach device



(projective) Quantum measurements

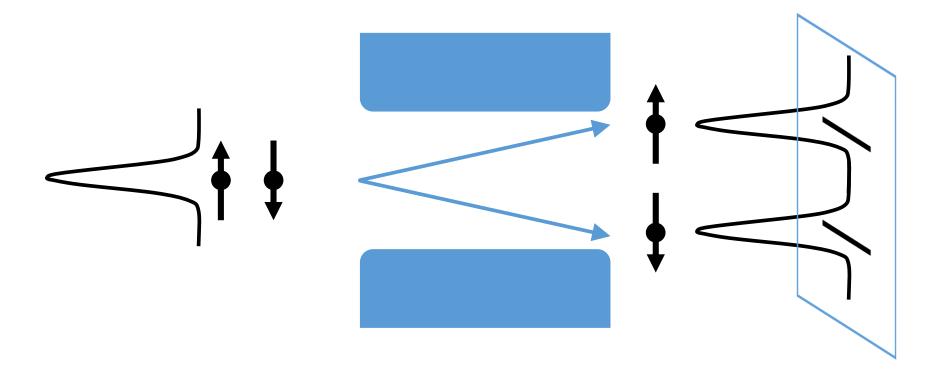
Stern–Gerlach device



$$\begin{aligned} (a|\uparrow\rangle + b|\downarrow\rangle)|Q &= 0\rangle \longrightarrow \\ \longrightarrow a|\uparrow\rangle|Q &= 1\rangle + b|\downarrow\rangle|Q &= -1\rangle \end{aligned}$$

 $\begin{aligned} a|\uparrow\rangle &= P_{\uparrow}(a|\uparrow\rangle + b|\downarrow\rangle) \\ b|\downarrow\rangle &= P_{\downarrow}(a|\uparrow\rangle + b|\downarrow\rangle) \end{aligned}$

Weak measurements



$(a|\uparrow\rangle + b|\downarrow\rangle)\int dQ\psi(Q)|Q\rangle$ $\rightarrow \int dQ (a|\uparrow) \psi(Q-1) + b|\downarrow\rangle \psi(Q+1) |Q\rangle$ $= \int dQ |Q\rangle M_Q(a|\uparrow\rangle + b|\downarrow\rangle)$

Weak measurements

Formalism of generalized measurements

Initial state $|\Psi\rangle$ (e.g., $a|\uparrow\rangle + b|\downarrow\rangle$)

Measurement readout r (e.g., Q)

System state after measurement $M_r |\Psi\rangle$, e.g., $M_Q(a|\uparrow\rangle + b|\downarrow\rangle) = (a|\uparrow\rangle\psi(Q-1) + b|\downarrow\rangle\psi(Q+1))$

Probability of readout *r*:

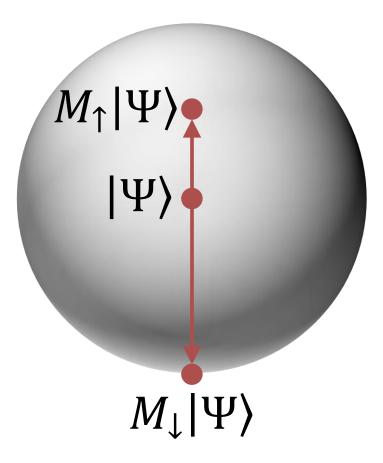
$$P = \langle \Psi | M_r^{\dagger} M_r | \Psi \rangle = \| M_r | \Psi \rangle \|^2$$

Two-outcome detector: the measurement back-action

Two-outcome detector: $r = \uparrow/\downarrow$

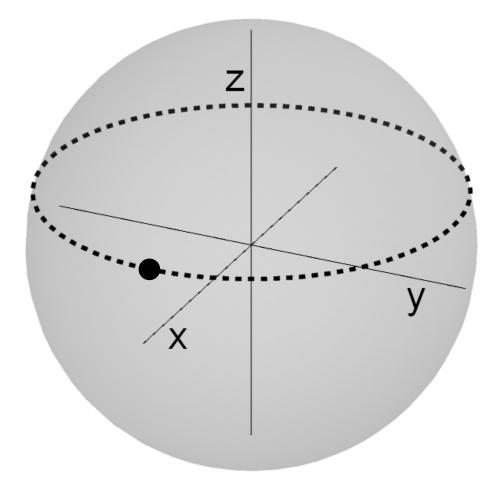
$$M_{\uparrow} = \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix}$$

$$M_{\downarrow} = \begin{pmatrix} 0 & 0 \\ 0 & \sqrt{1 - t^2} \end{pmatrix}$$



Weak-measurement-induced phase

$$\langle \Psi_0 | P_{N-1}^{(\uparrow)} \dots P_1^{(\uparrow)} | \Psi_0 \rangle = e^{i\varphi_P}$$



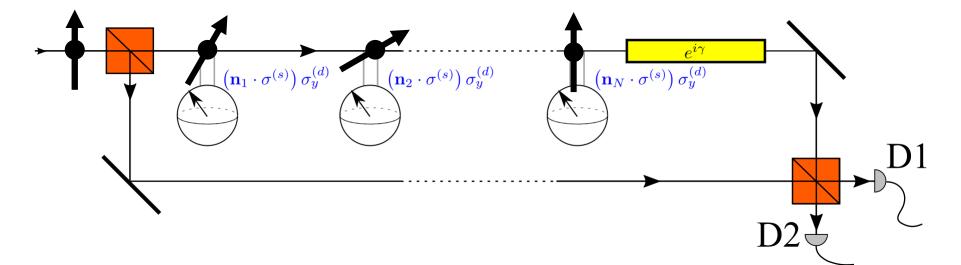
Weak-measurement-induced phase

$$\langle \Psi_{0} | P_{N-1}^{(\uparrow)} \dots P_{1}^{(\uparrow)} | \Psi_{0} \rangle = e^{i\varphi_{P}}$$

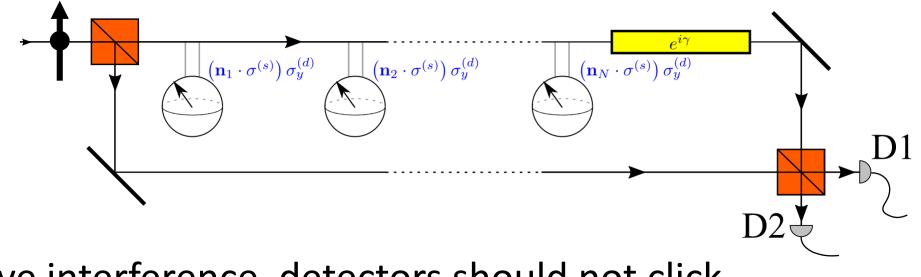
$$\langle \Psi_{0} | M_{N-1}^{(\uparrow)} \dots M_{1}^{(\uparrow)} | \Psi_{0} \rangle = e^{i\chi_{\text{geom}}} \sqrt{P}$$

$$\langle \Psi_{0} | M_{N-1}^{(r_{N-1})} \dots M_{1}^{(r_{1})} | \Psi_{0} \rangle = e^{i\chi_{\text{geom}}} \sqrt{P}$$

Observing measurement-induced geometric phase



Observing measurement-induced geometric phase

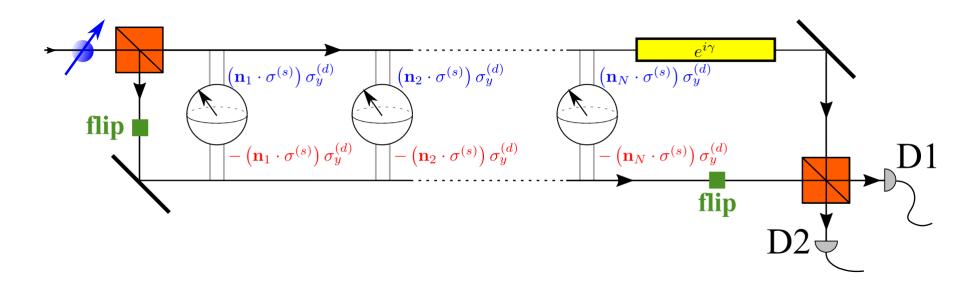


To observe interference, detectors should not click \implies Null-(outcome) measurement

 $\uparrow \text{-postselected phase} \\ \langle \Psi_0 | M_{N-1}^{(\uparrow)} \dots M_1^{(\uparrow)} | \Psi_0 \rangle = \sqrt{P} e^{i \chi_{\text{geom}}}$

$$M_{\uparrow} = \begin{pmatrix} 1 & 0 \\ 0 & 1 - \frac{2c}{N} \end{pmatrix}$$

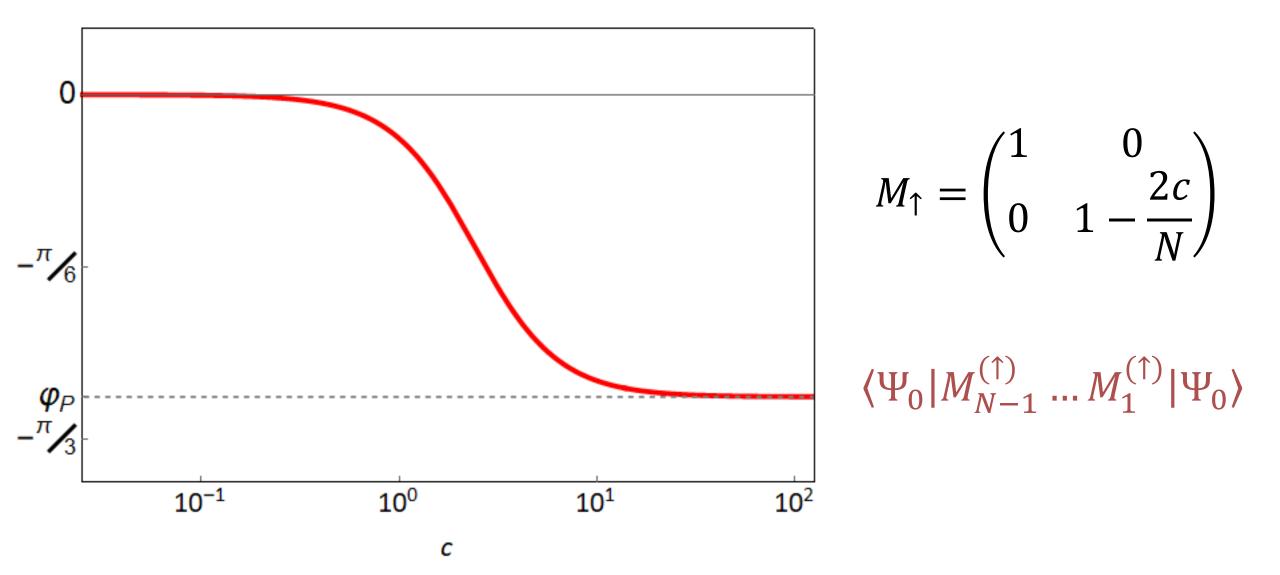
Observing averaged phase



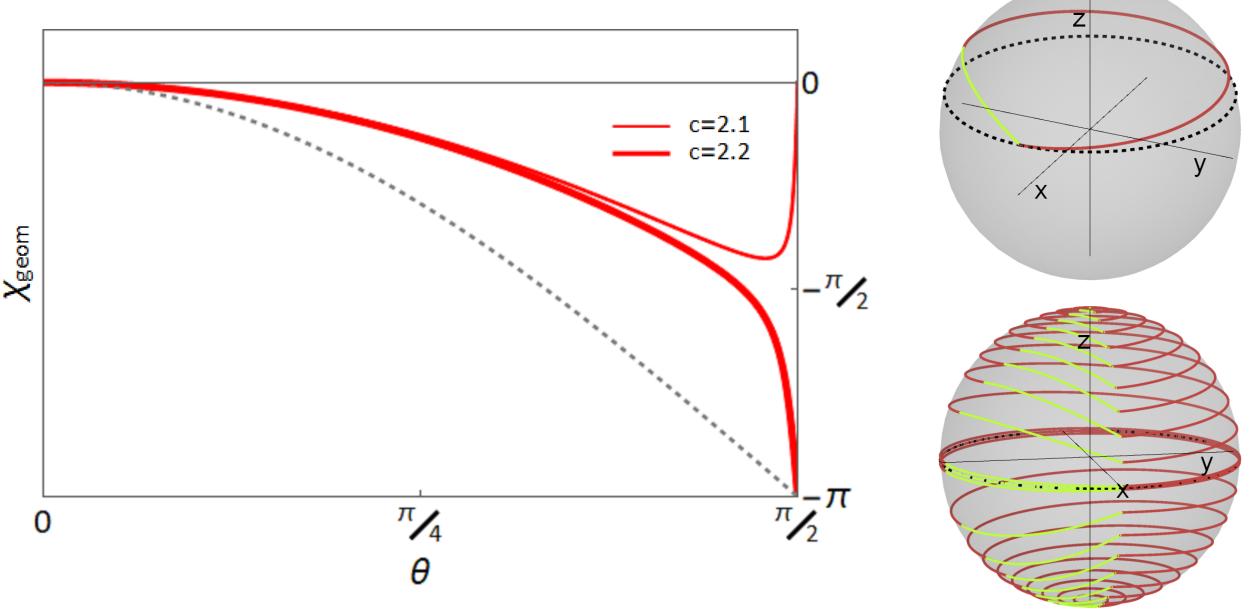
Can measure

$$\sum_{r_1,\ldots,r_{N-1}} \left(\langle \Psi_0 | M_{N-1}^{(r_{N-1})} \dots M_1^{(r_1)} | \Psi_0 \rangle \right)^2 = e^{2i\overline{\chi}_{\text{geom}} - \alpha}$$

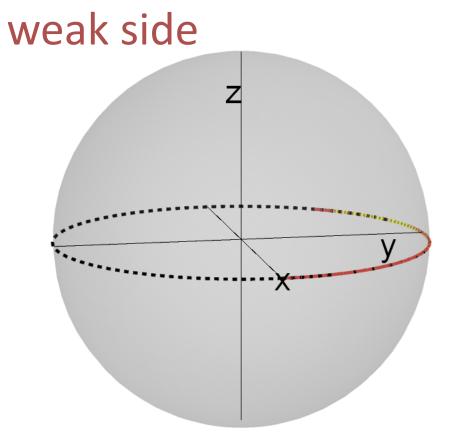
Phase dependence on the measurement strength



Surprise: sharp transition



Equator trajectory and the topological transition

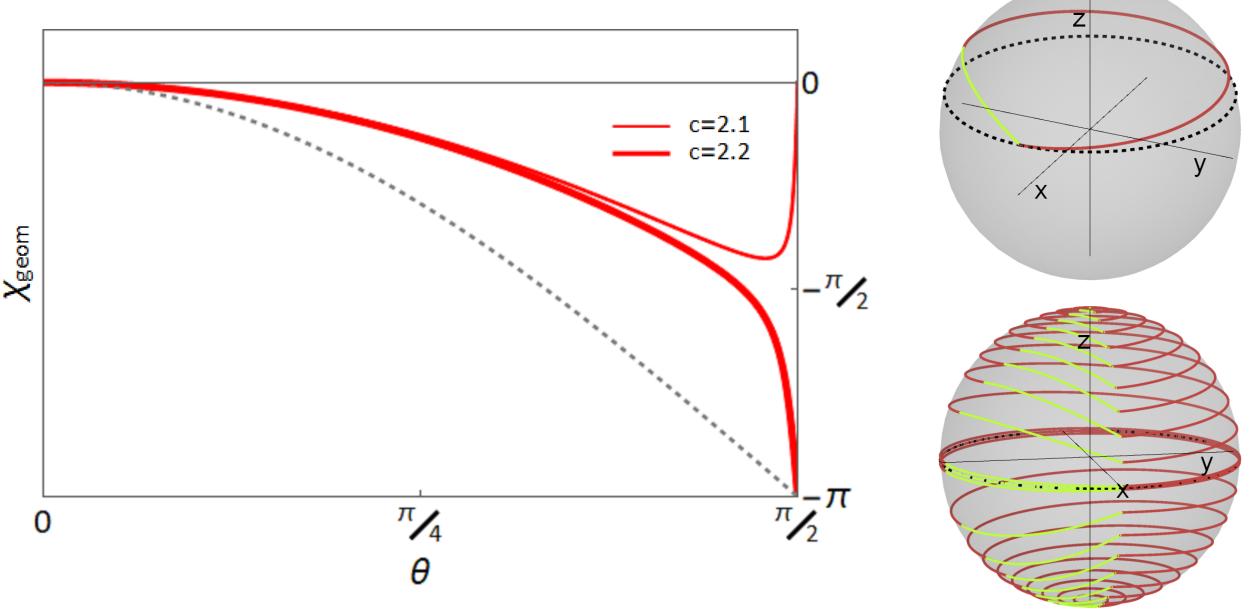


strong side

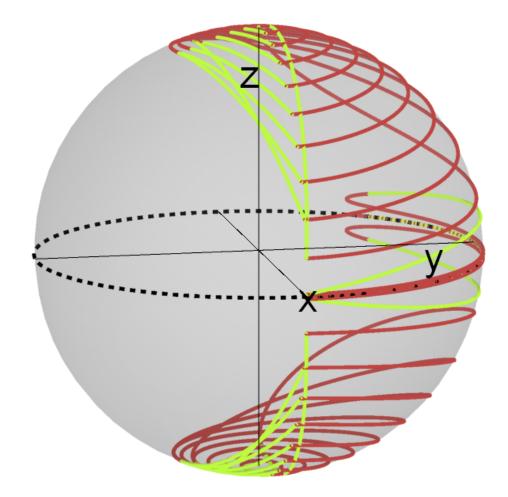
 $geometric \ phase = 0$

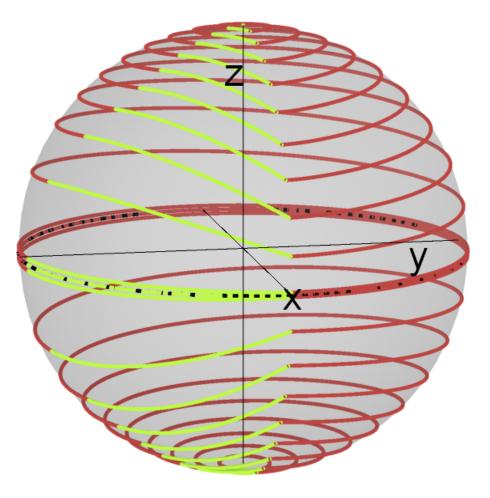
geometric phase = $-\pi$

Surprise: sharp transition



Trajectories and the topological transition (postelected)





Conclusions

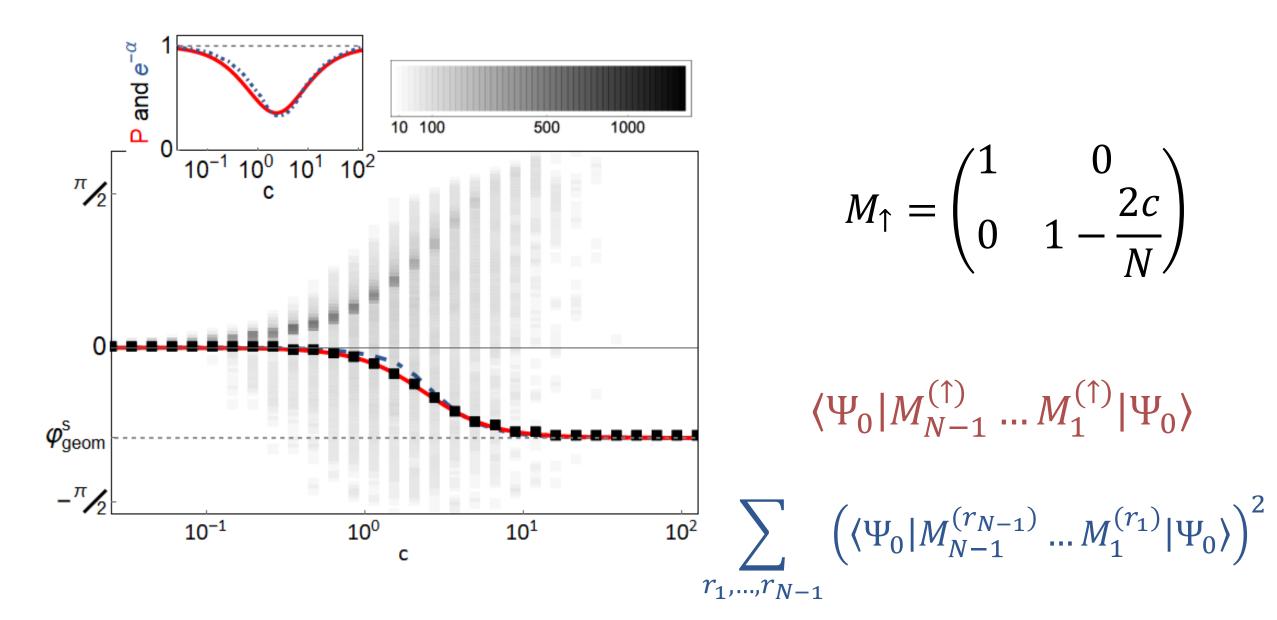
• Use measurement back-action to manipulate states

Induce geometric phase by weak measurements

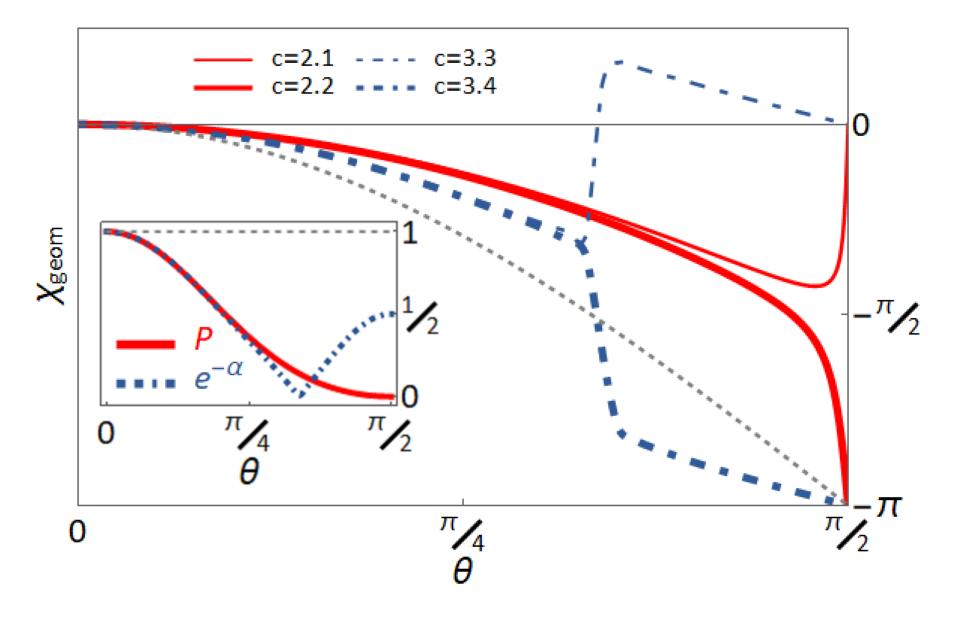
 Topological transition in weak-measurementinduced geometric phase as a function of measurement strength

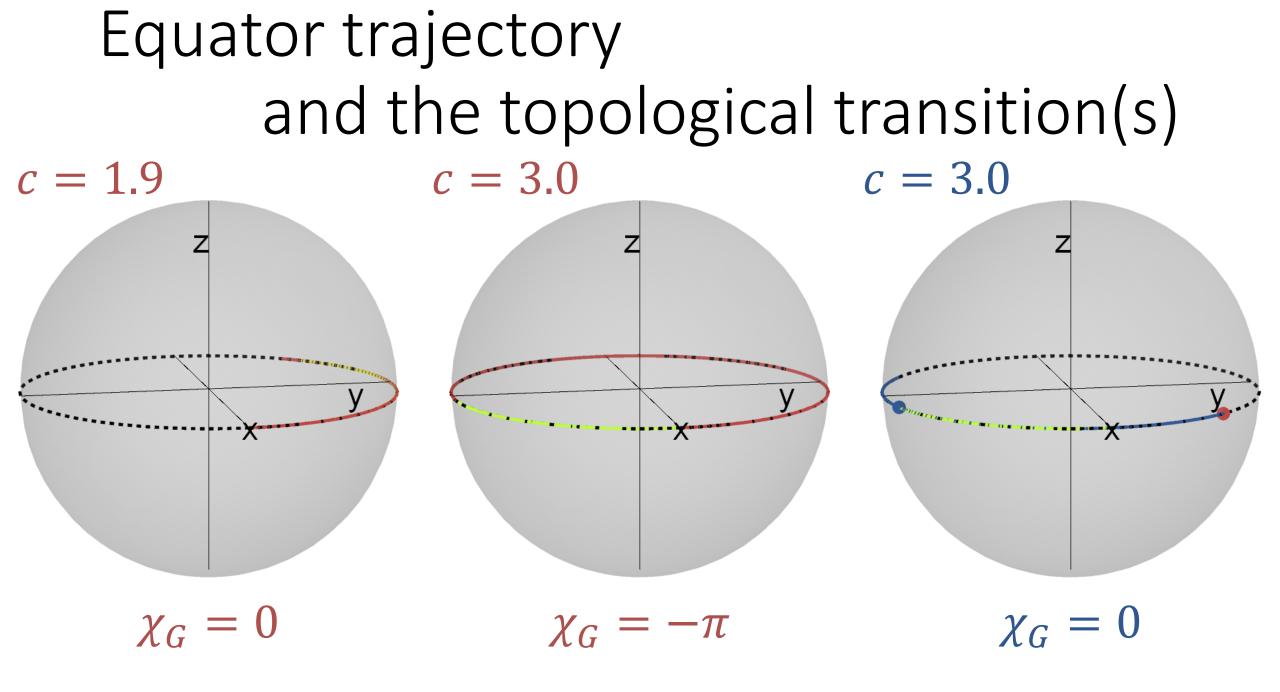
Thank you for attention!

Phase dependence on the measurement strength

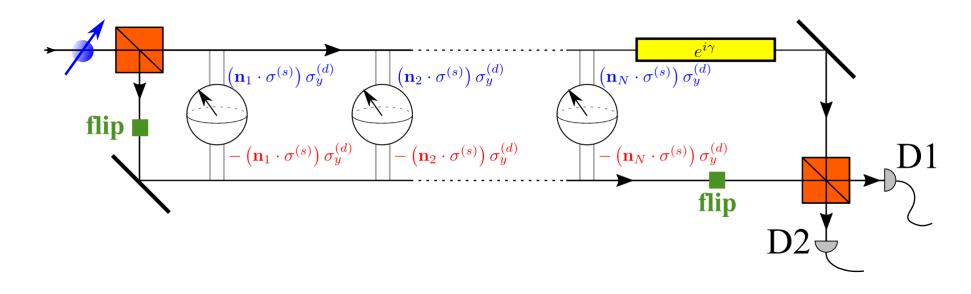


Surprise: sharp transition(s)





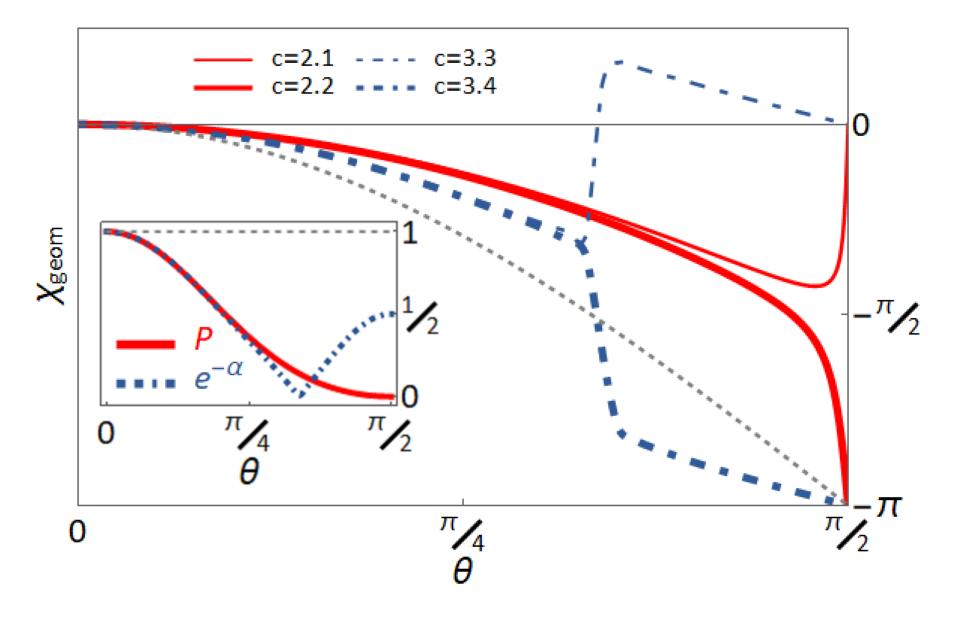
Observing averaged MIGP



Can measure

$$\sum_{r_1,\dots,r_{N-1}} \left(\langle \Psi_0 | \mathcal{M}_{N-1}^{(r_{N-1})} \dots \mathcal{M}_1^{(r_1)} | \Psi_0 \rangle \right)^2 = e^{2i\overline{\chi}_G - \alpha}$$

Surprise: sharp transition(s)

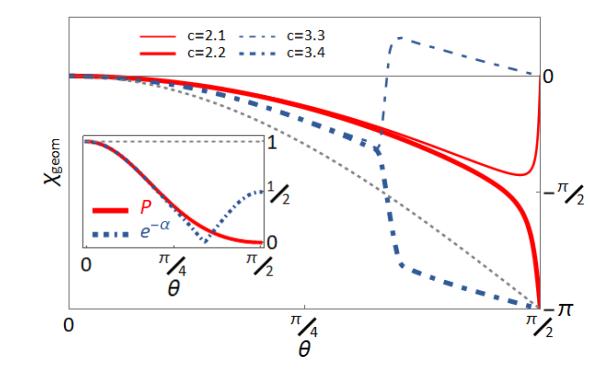


What do the transitions tell about?

• New property of measurements?

Not really, different critical strengths depending on what one looks

• New way of revealing topological features of the measured system? Seems so, topological features of the Bloch sphere in our example.



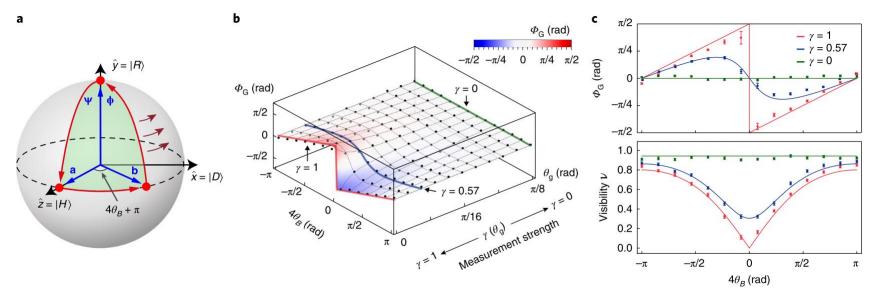
Recent experiments observe MIGP

nature physics

Article | Published: 15 April 2019

Emergence of the geometric phase from quantum measurement back-action

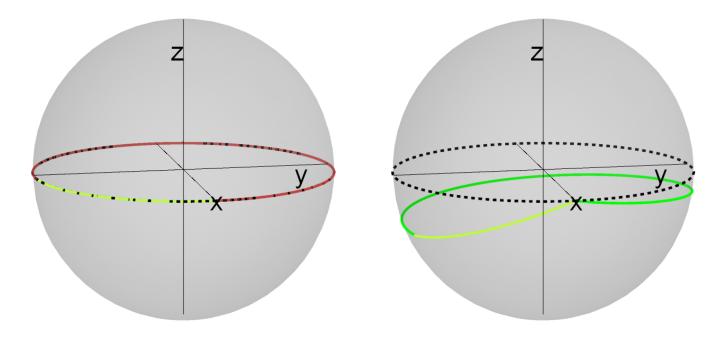
Young-Wook Cho ⊠, Yosep Kim, Yeon-Ho Choi, Yong-Su Kim, Sang-Wook Han, Sang-Yun Lee, Sung Moon & Yoon-Ho Kim ⊠

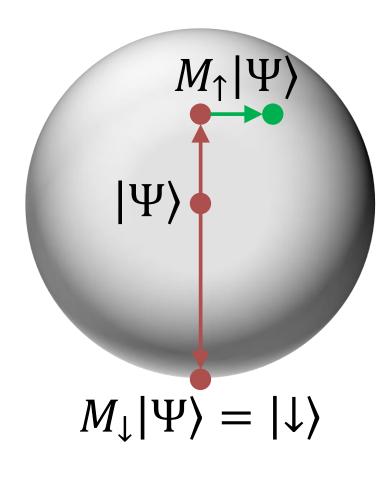


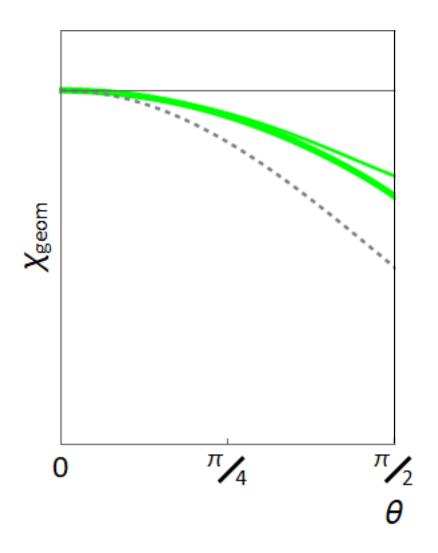


Two-outcome detector: $r = \uparrow/\downarrow$

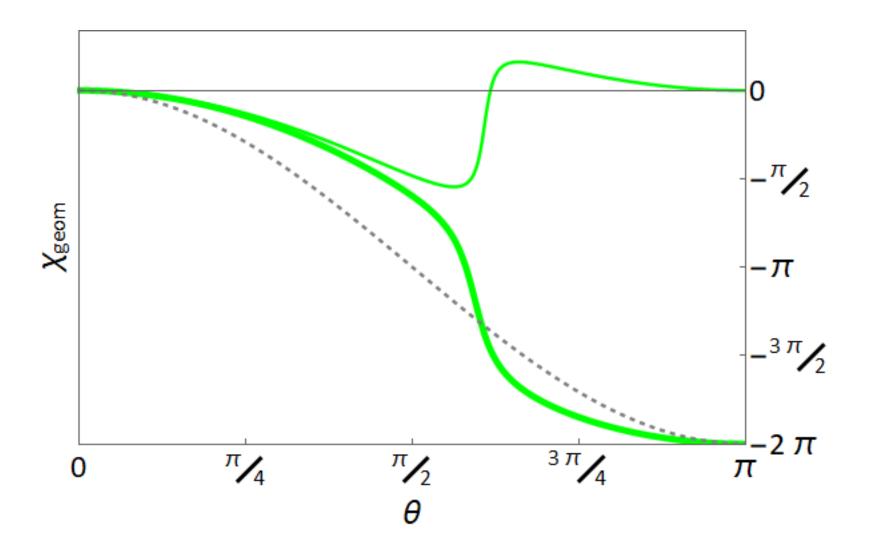
 $M_{\uparrow} = \begin{pmatrix} 1 & 0 \\ 0 & te^{i\phi} \end{pmatrix}$ $M_{\downarrow} = \begin{pmatrix} 0 & 0 \\ 0 & \sqrt{1 - t^2} \end{pmatrix}$



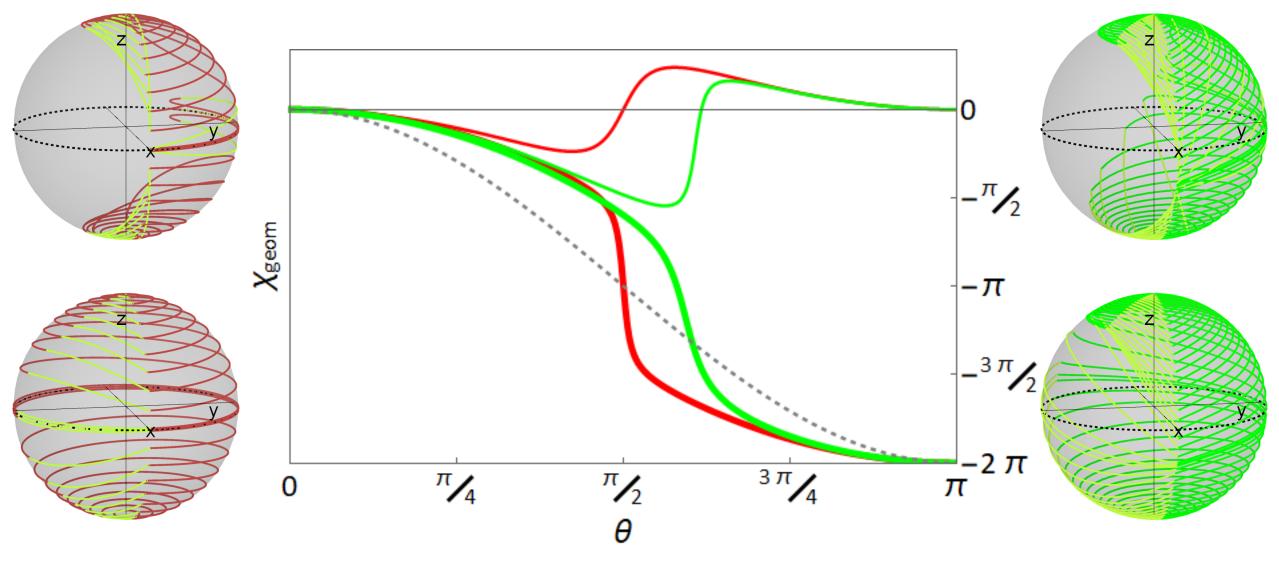




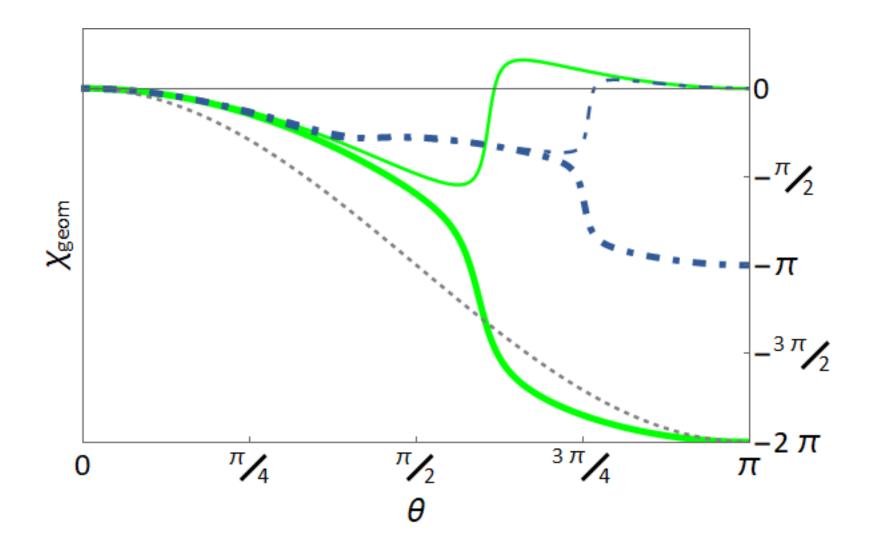








Also for the averaged phase



Phase diagram

