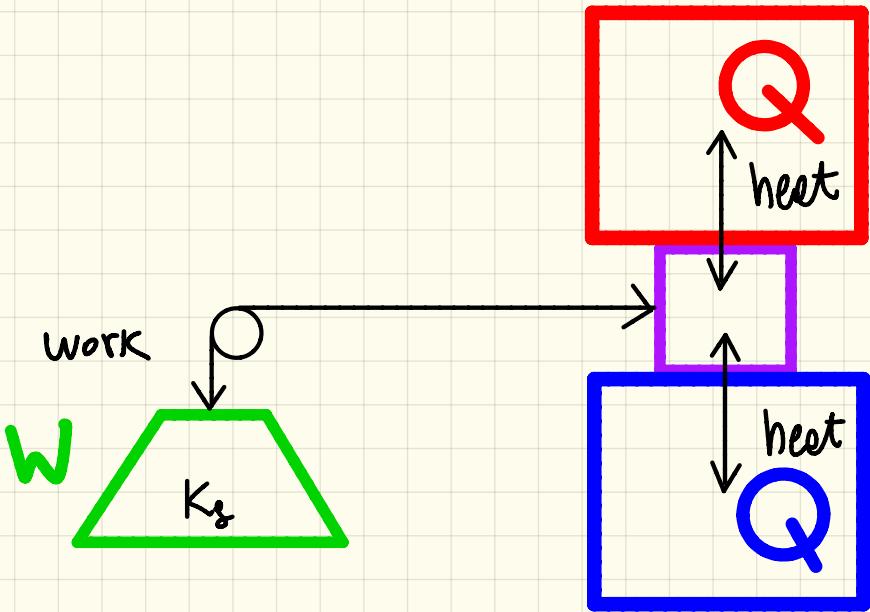
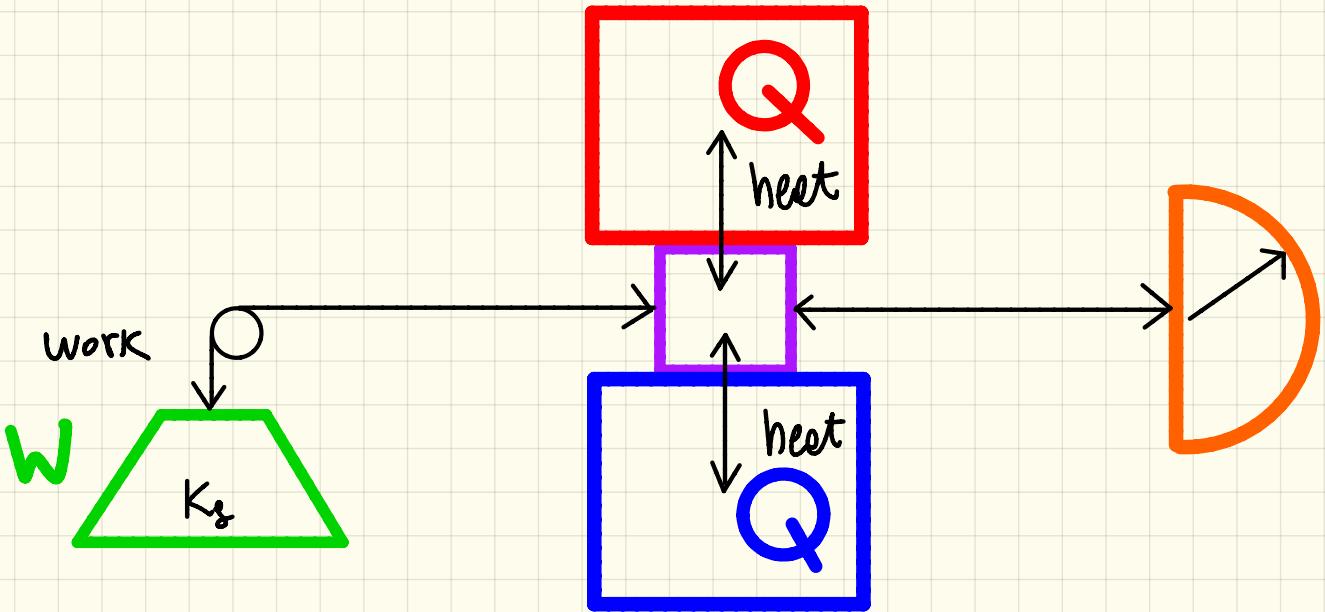


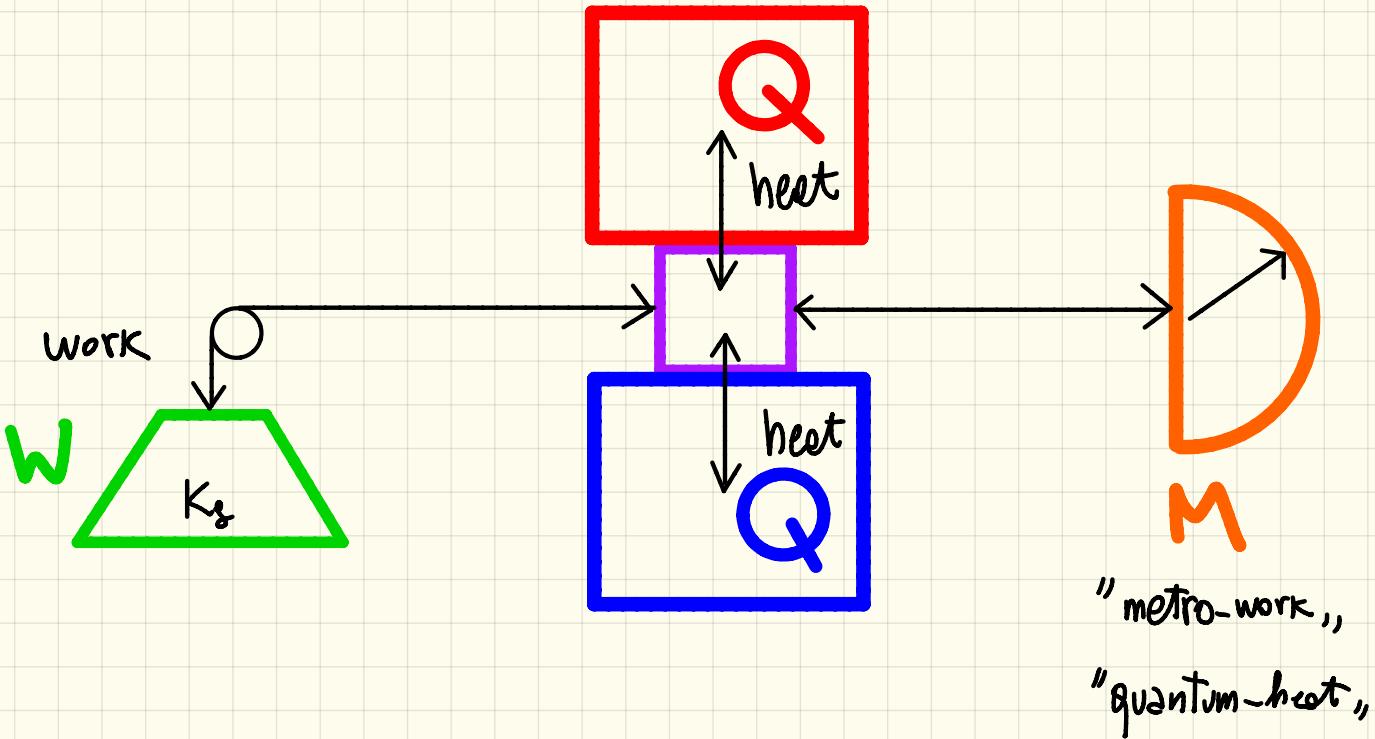
quantum measurement as a thermodynamic Resource

Michele Campisi
University of Florence





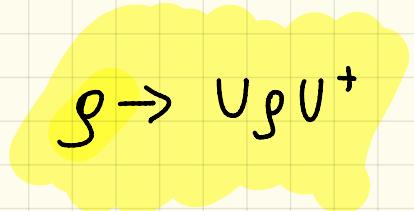
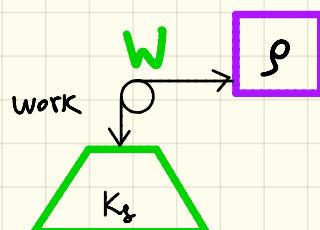




Maximal work extraction from finite quantum systems

A. E. ALLAHVERDYAN^{1,2}, R. BALIAN³ and TH. M. NIEUWENHUIZEN¹

(via unitary evolution)

ERGOTROPY \mathcal{W} 

PAPER: No subject

arXiv 1905.10262

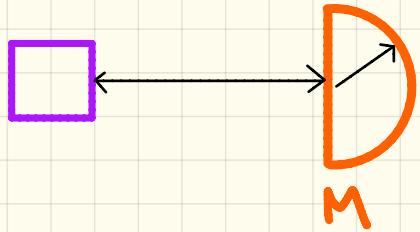
Maximal energy extraction via quantum measurement

Andrea Solfanelli^{1,2}, Lorenzo Buffoni^{1,3},
Alessandro Cuccoli^{1,2} and Michele Campisi^{1,2}

J. S.

→ METROTROPy

\mathcal{M}



$$g \rightarrow \sum_k \pi_k g \pi_k$$

ENTROPY \mathcal{W}

$$[H, S] = 0$$

$$E = \sum E_k r_k$$

$$g^l = U g U^\dagger$$

$$E' = \sum_k E_k \langle k | U_p U^\dagger | k \rangle$$

$$= \sum_k E_k P_{k\ell} r_\ell$$

$$= E^T \cdot P \cdot r$$

$$P_{k\ell} = |\langle k | U | \ell \rangle|^2$$

$$\mathcal{W} = \max (E - E')$$

$$= E - \min E'$$

P is uni-stochastic
 \Downarrow

P is doubly stochastic

Birkhoff theorem

$$P = \sum_i \lambda_i \sigma_i$$

$\sigma_i \leftarrow$ permutations

$$\min \rightarrow P = \sigma_w$$

bist.

unist.

given a permutation matrix σ

You can always find a unitary U such that

Permutations

$$\sigma_{ij} = |\langle i | U | j \rangle|^2$$

e.g.

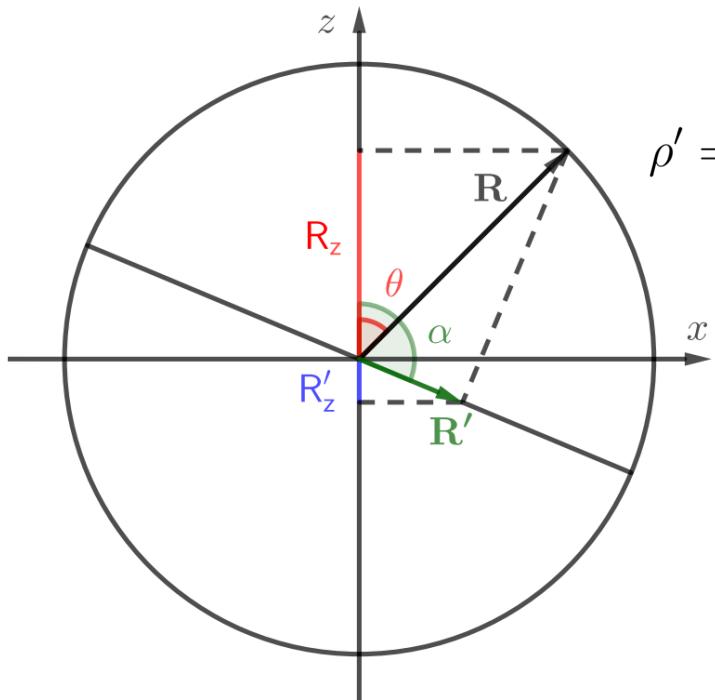
$$\sigma = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$U = \begin{pmatrix} 0 & 0 & e^{i\alpha} \\ e^{i\beta} & 0 & 0 \\ 0 & e^{i\gamma} & 0 \end{pmatrix}$$

ERGOTROPY vs METROTROPY

Two level system

$$H = \sigma_z$$



$$\rho = (\mathbb{1} + \mathbf{R} \cdot \boldsymbol{\sigma})/2$$

$$\rho' = \frac{\mathbb{1} + (\mathbf{R} \cdot \mathbf{n})\mathbf{n} \cdot \boldsymbol{\sigma}}{2} = \frac{\mathbb{1} + \mathbf{R}' \cdot \boldsymbol{\sigma}}{2}$$

$$\mathcal{W} = (|\bar{R}| + R_z) b_z$$

$$\mathcal{M} = \frac{\mathcal{W}}{2}$$

METROTROPY

$$[H, \mathcal{S}] = 0$$

$$\rho' = \sum_k \pi_k \rho \pi_k \quad \pi_k = |\psi_k \times \psi_k|$$

$$\begin{aligned} E' &= \sum_{lk} E_l \langle l | \psi_k \rangle \langle \psi_k | \rho | \psi_k \rangle \langle \psi_k | l \rangle \\ &= \sum_{lkn} E_l \langle l | \psi_k \rangle \langle \psi_k | n \rangle \langle n | \psi_k \rangle \langle \psi_k | l \rangle r_n \\ &= \sum_{lkn} E_l |\langle l | \psi_k \rangle|^2 |\langle n | \psi_k \rangle|^2 r_n \\ &= \mathbf{E}^T \cdot \mathbf{P}^T \cdot \mathbf{P} \cdot \mathbf{r} \end{aligned}$$

$$P_{ek} = |\langle e | \psi_k \rangle|^2 = |\langle e | v_{ik} \rangle|^2$$

$$M = \max(E - E')$$

$$= E - \min E'$$

1) Use Birkhoff theorem

to find

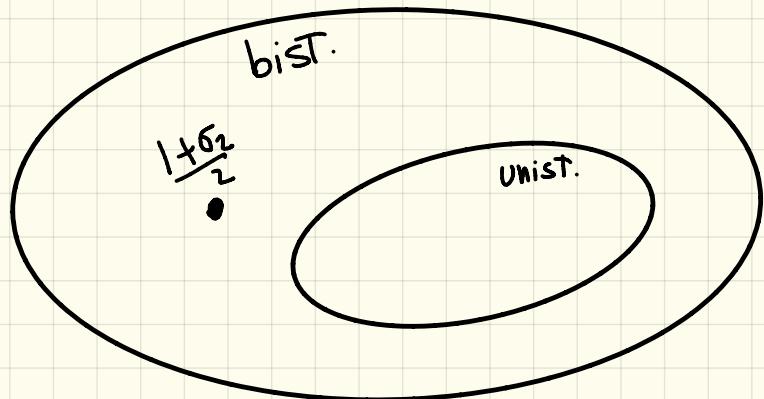
$$\min_{B_n} E^T \cdot P^T \cdot P \cdot r$$

B_n = set of Bistochastic matrices

$$P = \frac{1 + \sigma_2}{2}$$

$$\sigma_2 = \operatorname{argmin} E^T \cdot \frac{\sigma + \sigma^T}{2} \cdot r$$

2) Not all bistochastic matrices are unistochastic!



Example

$$\frac{1}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \\ 1 & 0 & 1 \end{pmatrix} \leftarrow \text{not unistochastic}$$

↓

$$\begin{pmatrix} |1\rangle & |2\rangle & |3\rangle \\ |2\rangle & |3\rangle & |1\rangle \end{pmatrix}$$

Theorem

(Au-Yeung, Linear algebra its appl. 150 (1991))

Any convex combination of permutation matrices which is unistochastic, is such that it involves only permutations that are pairwise complementary.

Two $N \times N$ matrices A and B are said to be complementary if, for any $1 \leq i, j, h, k \leq N$, $A_{ij} = A_{hk} = B_{ik} = 1$ implies $B_{hj} = 1$.

Any permutation that is complementary to the identity is symmetric

Is any convex combination involving only complementary permutations

unistochastic?

OPEN QUESTION

YES: up to $N \leq 15$!

Av-Yaglom, Linear algebra its appl. 150 (1991)

STATEMENT

the minimum of $\mathbf{E}^T \cdot \mathbf{P}^T \cdot \mathbf{P} \cdot \mathbf{r}$ over all
complex combinations containing only pairwise permutations is achieved by

$$\mathbf{P} = (\mathbb{1} + \boldsymbol{\sigma}_{\mathcal{M}})/2$$

$$\boldsymbol{\sigma}_{\mathcal{M}} \doteq \arg \min_{\boldsymbol{\sigma} \in \mathcal{I}_N} \mathbf{E}^T \cdot \boldsymbol{\sigma} \cdot \mathbf{r}.$$

\mathcal{I}_n = Set of symmetric permutations

Symmetric
permutation

$$\sigma_M = \begin{pmatrix} |a\rangle & |b\rangle & |c\rangle & |d\rangle & \dots & |x\rangle \\ |b\rangle & |a\rangle & |d\rangle & |c\rangle & \dots & |x\rangle \end{pmatrix}$$

$$U|a\rangle = \frac{|a\rangle + |b\rangle}{\sqrt{2}}$$

$$U|b\rangle = \frac{|a\rangle - |b\rangle}{\sqrt{2}}$$

$$U|c\rangle = \frac{|c\rangle + |d\rangle}{\sqrt{2}}$$

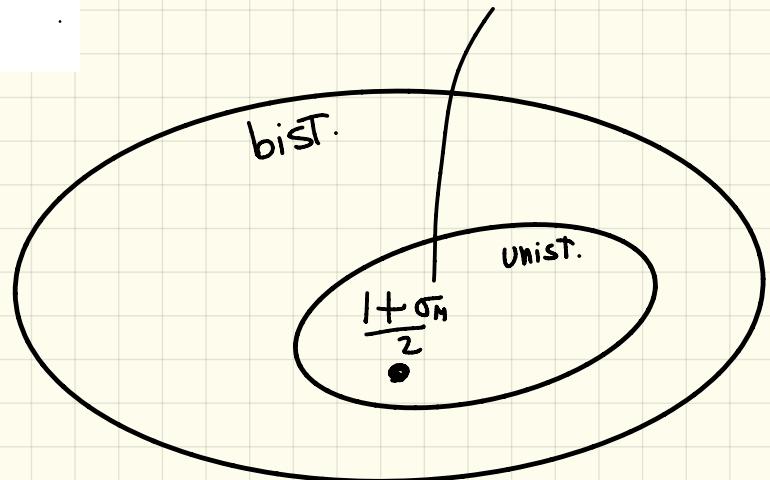
$$U|d\rangle = \frac{|c\rangle - |d\rangle}{\sqrt{2}}$$

⋮

$$U|x\rangle = |x\rangle, \quad x \neq a, b, c, d, \dots$$

$$P_{ij} = |\langle i | U | j \rangle|^2$$

$$\sigma_M = \sigma_M^T$$



Theorem

the minimum of $\mathbf{E}^T \cdot \mathbf{P}^T \cdot \mathbf{P} \cdot \mathbf{r}$ over all

UNISTOCHASTIC MATRICES

is achieved by

$$\mathbf{P} = (\mathbb{1} + \sigma_{\mathcal{M}})/2$$

$$\sigma_{\mathcal{M}} \doteq \arg \min_{\sigma \in \mathcal{I}_N} \mathbf{E}^T \cdot \sigma \cdot \mathbf{r}.$$

\mathcal{I}_n = Set of symmetric permutations

Example 4-level system

$$\sigma_M = \begin{pmatrix} 1 & & & \\ 0 & 1 & & \\ 1 & 0 & & \\ & & & 1 \end{pmatrix}$$

$$P = \begin{pmatrix} 1 & & & \\ \frac{1}{2} & \frac{1}{2} & & \\ \frac{1}{2} & \frac{1}{2} & & \\ \frac{1}{2} & \frac{1}{2} & & 1 \end{pmatrix}$$

$$U = \begin{pmatrix} 1 & & & \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & & \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & & \\ & & & 1 \end{pmatrix}$$

Corollaries

- $M = \frac{\underset{I_n}{\text{E} - \min} \mathbf{E}^T \cdot \boldsymbol{\sigma} \cdot \mathbf{r}}{2}$
- if $\boldsymbol{\sigma}_W$ is symmetric, then $M = \frac{W}{2}$
- $M \leq \frac{W}{2}$ ($[H_{ij}] = 0$)
- $M \leq W$ (in general)

Conjecture

- $M \leq \frac{W}{2}$ (in general)

Three level system

$$E = \varepsilon \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}$$

$$r = \begin{pmatrix} r_0 \\ r_1 \\ r_2 \end{pmatrix}$$

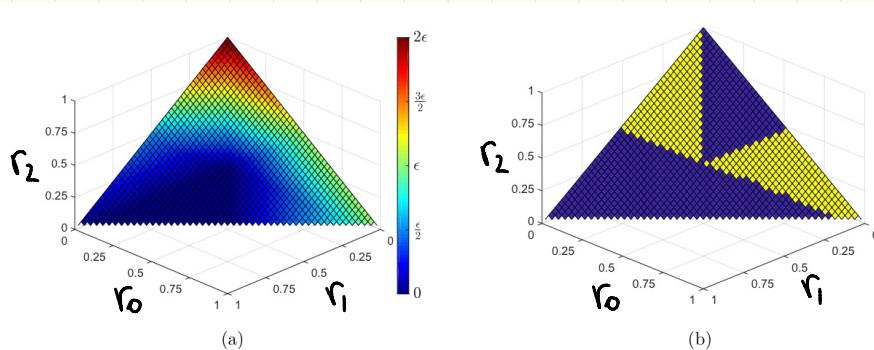


Figure 1: Panel a): Ergotropy, \mathcal{W} . Panel b): Symmetry of the ergotropy permutation $\sigma_{\mathcal{W}}$. Blue denotes symmetric $\sigma_{\mathcal{W}}$, yellow denotes non-symmetric $\sigma_{\mathcal{W}}$.

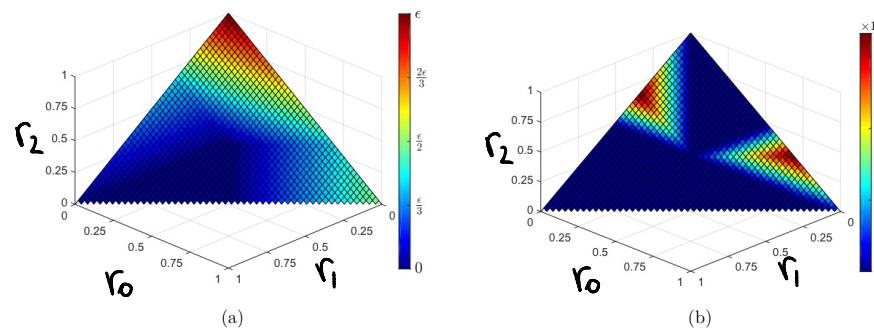


Figure 2: Panel (a): Metroropy \mathcal{M} . Panel b): $\mathcal{W}/2 - \mathcal{M}$.

Quantum Measurement Cooling

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⁴*INFN Sezione di Firenze, via G.Sansone 1, I-50019 Sesto Fiorentino (FI), Italy*

⁵*Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California 93106, USA*

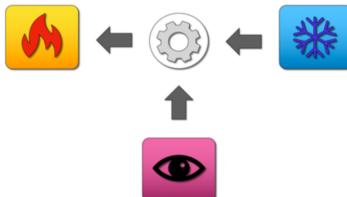


ABOUT BROWSE PRESS COLLECTIONS

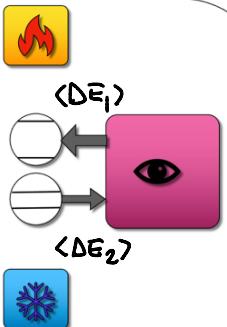
Synopsis: Refrigeration by Quantum Measurements

February 21, 2019

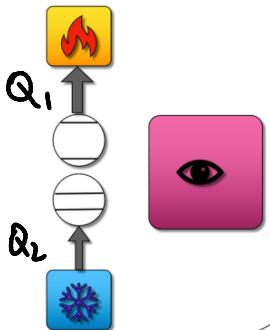
A proposed noise-tolerant approach to quantum refrigeration eliminates the need for feedback control by exploiting the invasiveness of quantum measurements.



First Stroke



Second Stroke



$$H = H_1 + H_2 + V(t)$$

$\frac{w_1}{z} \sigma_1^z$ $\frac{w_2}{z} \sigma_2^z$

$$g = \frac{e^{-\beta_1 H_1}}{Z_1} \otimes \frac{e^{-\beta_2 H_2}}{Z_2}$$

$$g \rightarrow g' = \sum_k \pi_k g \pi_k$$

$$T_{F_i} H_i (g'_i - g_i) = \langle \Delta E_i \rangle = Q_i$$

$$\begin{aligned} \langle M \rangle &= \langle \Delta E \rangle = \langle \Delta E_1 \rangle + \langle \Delta E_2 \rangle \\ &= Q_1 + Q_2 \end{aligned}$$

$$\beta_1 Q_1 + \beta_2 Q_2 = D[\rho' || \rho_1] + D[\rho' || \rho_2] + I_{1/2}[\rho'] + \Delta S \geq 0$$

$$\begin{array}{c}
 \text{IV} \\
 | \\
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 \quad
 \left\{ \begin{array}{c} \text{IV} \\ | \\ \text{O} \end{array} \right\}$$

$I_{1/2}[\rho']$
 \parallel
 $D[\rho' || \rho'_1 \otimes \rho'_2]$
 IV
 $|$
 O

$\rho' = \sum_k \Pi_k \rho \Pi_k$
 is unital!
 \Downarrow
 $\Delta S \geq 0$

$$D[\rho || \sigma] = \text{Tr}(\rho \ln \rho - \rho \ln \sigma)$$

$$I_{1/2}[\rho'] = \sum_i S[\rho'_i] - S[\rho_i]$$

$$S[\rho] = -\text{Tr} \rho \ln \rho$$

$$\Delta S = S[\rho_t] - S[\rho_0]$$

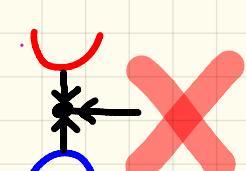
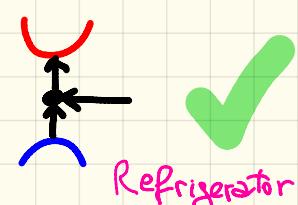
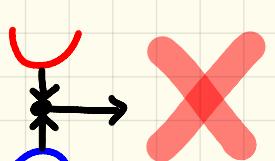
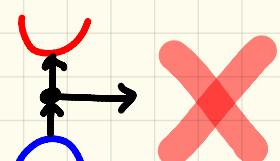
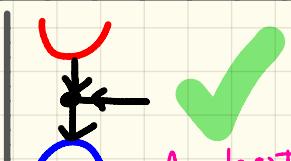
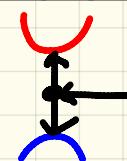
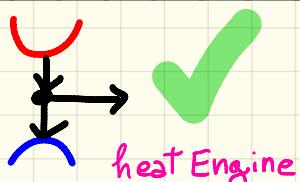
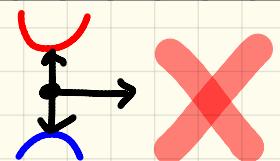
$$\beta_1 Q_1 + \beta_2 Q_2 \geq 0$$

$$\langle \Delta E \rangle = Q_1 + Q_2$$

$\langle M \rangle$

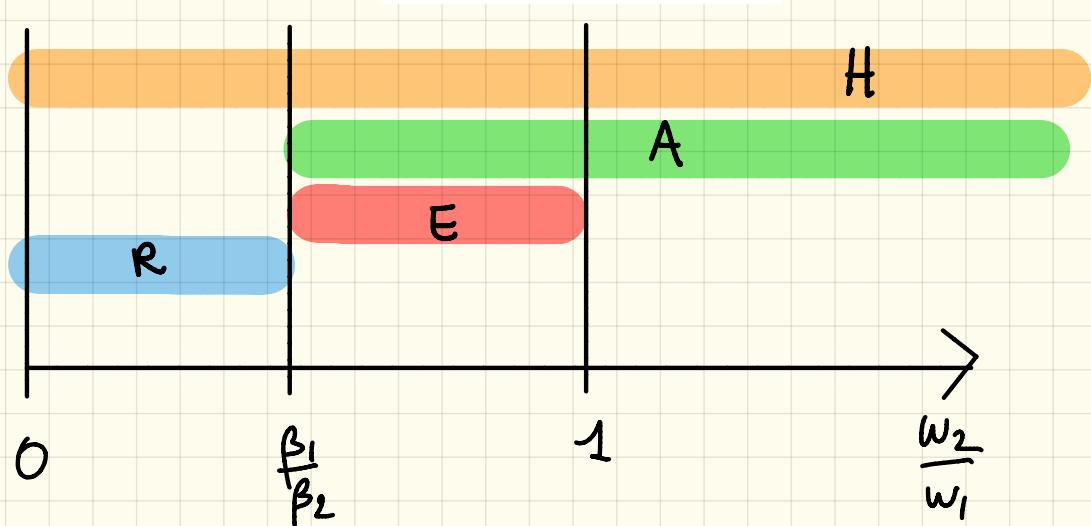
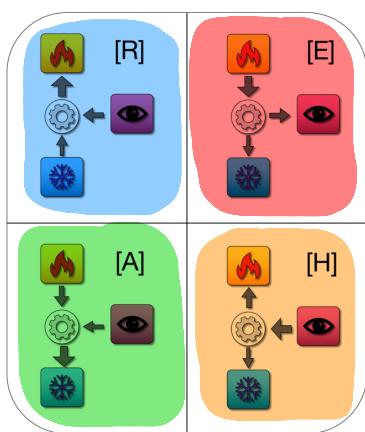
A. Solfanelli

B. Sc. thesis UNIFI



Results

①



Results

②

$$\left\{ \begin{array}{l} |\psi_1^*\rangle = |\uparrow\uparrow\rangle \\ |\psi_2^*\rangle = \frac{|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle}{\sqrt{2}} \\ |\psi_3^*\rangle = \frac{|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle}{\sqrt{2}} \\ |\psi_4^*\rangle = |\downarrow\downarrow\rangle \end{array} \right.$$

maximises

$$\eta^{[R]}, -Q_2 \quad \text{in } [R]\text{-range}$$

$$\eta^{[E]}, \langle \Delta E \rangle \quad \text{in } [E]\text{-range}$$

$$\langle \Delta E_{1,2} \rangle = \frac{\pm \omega_i}{2} \left(\frac{1}{1 + e^{\beta_1 \omega_1}} - \frac{1}{1 + e^{\beta_2 \omega_2}} \right)$$

$$\eta^{[R]} = \frac{1}{\frac{w_1}{N_2} - 1}$$

$$\eta^{[E]} = 1 - \frac{w_2}{w_1}$$

Results

③

Maximal efficiency and $\begin{cases} \text{net-work output [E]} \\ \text{heat extraction [R]} \end{cases}$

can be achieved with higher rank projectors

$$\text{e.g.: } q_1 = |\psi_1^* \times \psi_1^*| + |\psi_2^* \times \psi_2^*|$$

$$q_2 = |\psi_3^* \times \psi_3^*| + |\psi_4^* \times \psi_k^*|$$

Results

(4)

Generally, for two qudits

In order to realize anything other than $[H]$, some of the measurement projectors
must be entangled
(cannot be written in factorised form)

(4.1)

However, extremal efficiency occurs when

$$S^I = \sum p_n^I \ln x_n I$$

$$|h\rangle = |n_1\rangle |n_2\rangle$$

Results

(5)

$$\text{Let } |\Psi_K\rangle = U|K\rangle$$

Pick U randomly from the invariant $SU(2N)$ measure
then

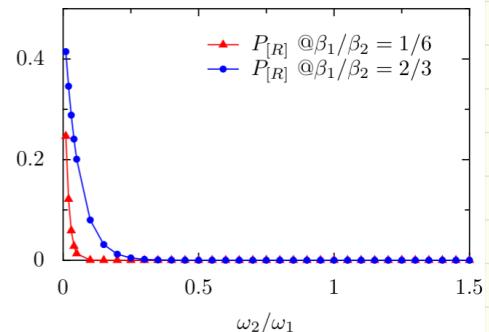
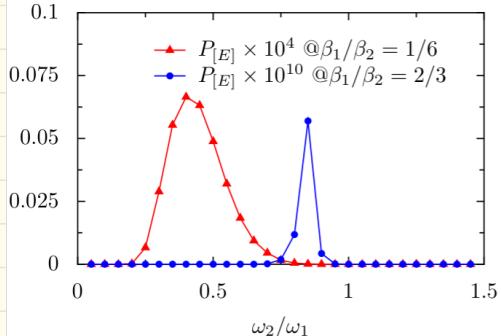
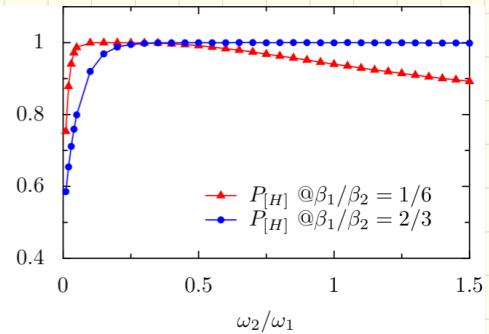
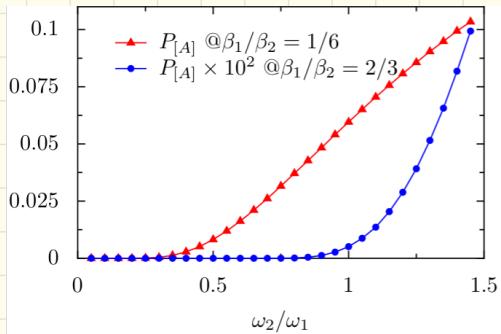
$$\overline{\langle \Delta E_i \rangle} \geq 0 \Rightarrow [H]$$

$$\left(\bar{f} = \int_{SU(2N)} d^m f \right)$$

Results

6

Monte Carlo Sampling of SU(4)



Experiment

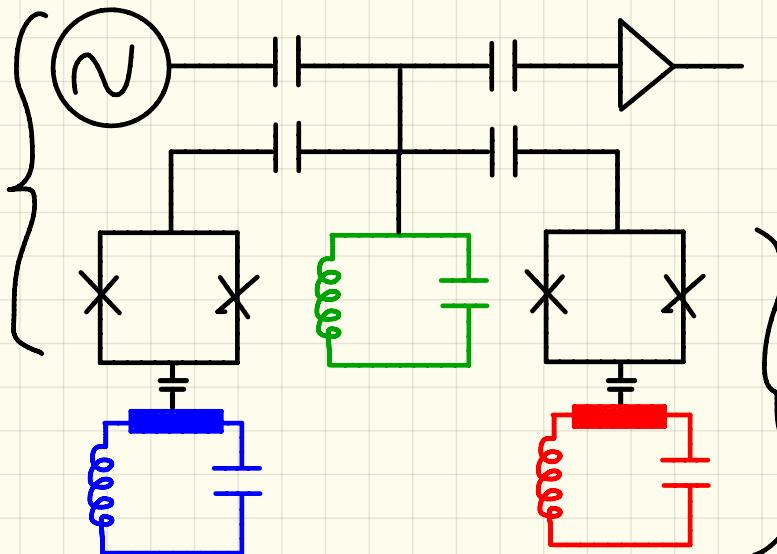
circuit QED + circuit QTD

$$\begin{aligned} g^I &= \sum_k \Pi_k g \Pi_k \\ &= \sum_k U P_k U^+ g U P_k U^+ \end{aligned}$$

circuit Quantum
Thermo
Dynamics

→ Pekola, Giacobbo....

Filipp et al.,
PRL 102
200402 (2009)



Ronzani et al
Nat phys 14 991
(2018)

thank you

