

# A SUSY RG-Flow Using Hamiltonian Methods

Ami Katz

Boston University

w/ L. Fitzpatrick, M. Walters, & Y. Xin

# Why use Light-Cone Quantization?

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

Low-E states are different:  $|\langle \psi_{H_0} | \psi_H \rangle|^2 \sim e^{-cL^{d-1}}$

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

Low-E states are different:  $|\langle \psi_{H_0} | \psi_H \rangle|^2 \sim e^{-cL^{d-1}}$

Problem: You need many particles to properly approximate low-E states of  $H$  for larger Volumes.

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

Low-E states are different:  $|\langle \psi_{H_0} | \psi_H \rangle|^2 \sim e^{-cL^{d-1}}$

Problem: You need many particles to properly approximate low-E states of  $H$  for larger Volumes.

Even more challenging for:  $\Delta > \frac{d+1}{2}$

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

Low-E states are different:  $|\langle \psi_{H_0} | \psi_H \rangle|^2 \sim e^{-cL^{d-1}}$

Problem: You need many particles to properly approximate low-E states of  $H$  for larger Volumes.

Even more challenging for:  $\Delta > \frac{d+1}{2}$

$$|\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}} (\lambda^2 \Lambda^{2\Delta - (d+1)} + \dots)$$

# Why use Light-Cone Quantization?

ET Quantization Orthogonality Catastrophe:

$$H = H_0 + V: \quad |\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}}$$

Low-E states are different:  $|\langle \psi_{H_0} | \psi_H \rangle|^2 \sim e^{-cL^{d-1}}$

Problem: You need many particles to properly approximate low-E states of  $H$  for larger Volumes.

Even more challenging for:  $\Delta > \frac{d+1}{2}$

$$|\langle 0 | \Omega \rangle|^2 \sim e^{-cL^{d-1}} (\lambda^2 \Lambda^{2\Delta - (d+1)} + \dots)$$

True for: Yukawa in 3D or fermion mass in 4D

## Light-cone quantization has a trivial vacuum

$$P_+ \rightarrow P_+ + \delta P_+$$



$$[\delta P_+, P_-] = 0$$

## Light-cone quantization has a trivial vacuum

$$P_+ \rightarrow P_+ + \delta P_+ \quad \longrightarrow \quad [\delta P_+, P_-] = 0$$

**But**  $P_- |0\rangle = 0$  **while**  $P_- > 0$  **for other**  $|\psi\rangle$  **(NEC)**

## Light-cone quantization has a trivial vacuum

$$P_+ \rightarrow P_+ + \delta P_+ \quad \longrightarrow \quad [\delta P_+, P_-] = 0$$

**But**  $P_-|0\rangle = 0$  **while**  $P_- > 0$  **for other**  $|\psi\rangle$  **(NEC)**

**So**  $\langle 0 | \delta P_+ | \psi \rangle = 0!$  (i.e. vacuum is trivial)

## Light-cone quantization has a trivial vacuum

$$P_+ \rightarrow P_+ + \delta P_+ \quad \longrightarrow \quad [\delta P_+, P_-] = 0$$

But  $P_-|0\rangle = 0$  while  $P_- > 0$  for other  $|\psi\rangle$  (NEC)

So  $\langle 0 | \delta P_+ | \psi \rangle = 0!$  (i.e. vacuum is trivial)

There's no dependence of observables on the volume

LC has its own issues

## LC has it own issues

Ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

## LC has it own issues

Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$

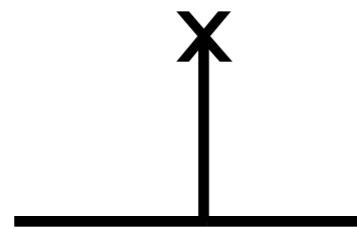
## LC has its own issues

Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$

(all LC modes have  $P_- > 0$ )

## LC has its own issues

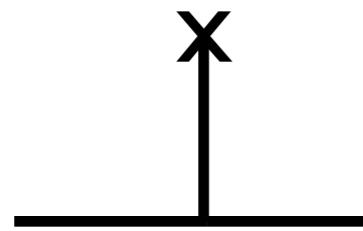
Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$   
**(all LC modes have  $P_- > 0$ )**



$$\delta m^2 = -\frac{\lambda g}{m^2}$$

## LC has its own issues

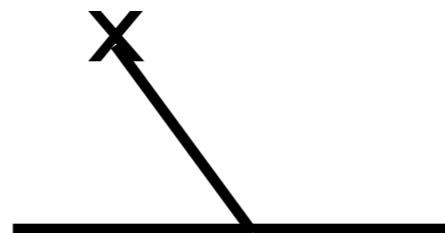
Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$   
(all LC modes have  $P_- > 0$ )



:

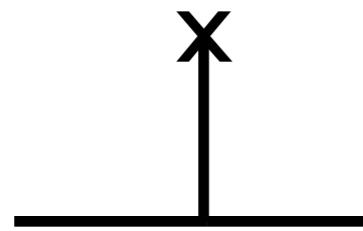
$$\delta m^2 = -\frac{\lambda g}{m^2}$$

ET:



## LC has its own issues

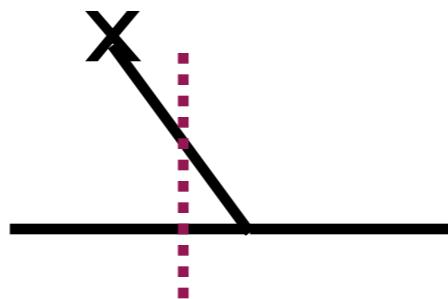
Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$   
**(all LC modes have  $P_- > 0$ )**



:

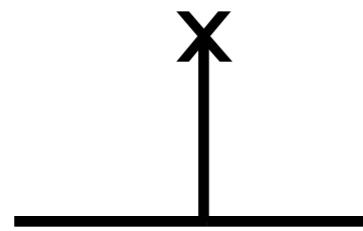
$$\delta m^2 = -\frac{\lambda g}{m^2}$$

ET:



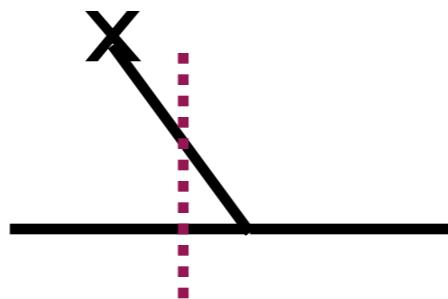
## LC has its own issues

Ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi \rightarrow \int dx^- dx^{d-2} \phi(x^-, \vec{x}^\perp) = 0$   
(all LC modes have  $P_- > 0$ )



:  $\delta m^2 = -\frac{\lambda g}{m^2}$

ET:



LC lacks appropriate intermediate states!

2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_* :$

2D Ex:  $\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_*$ :

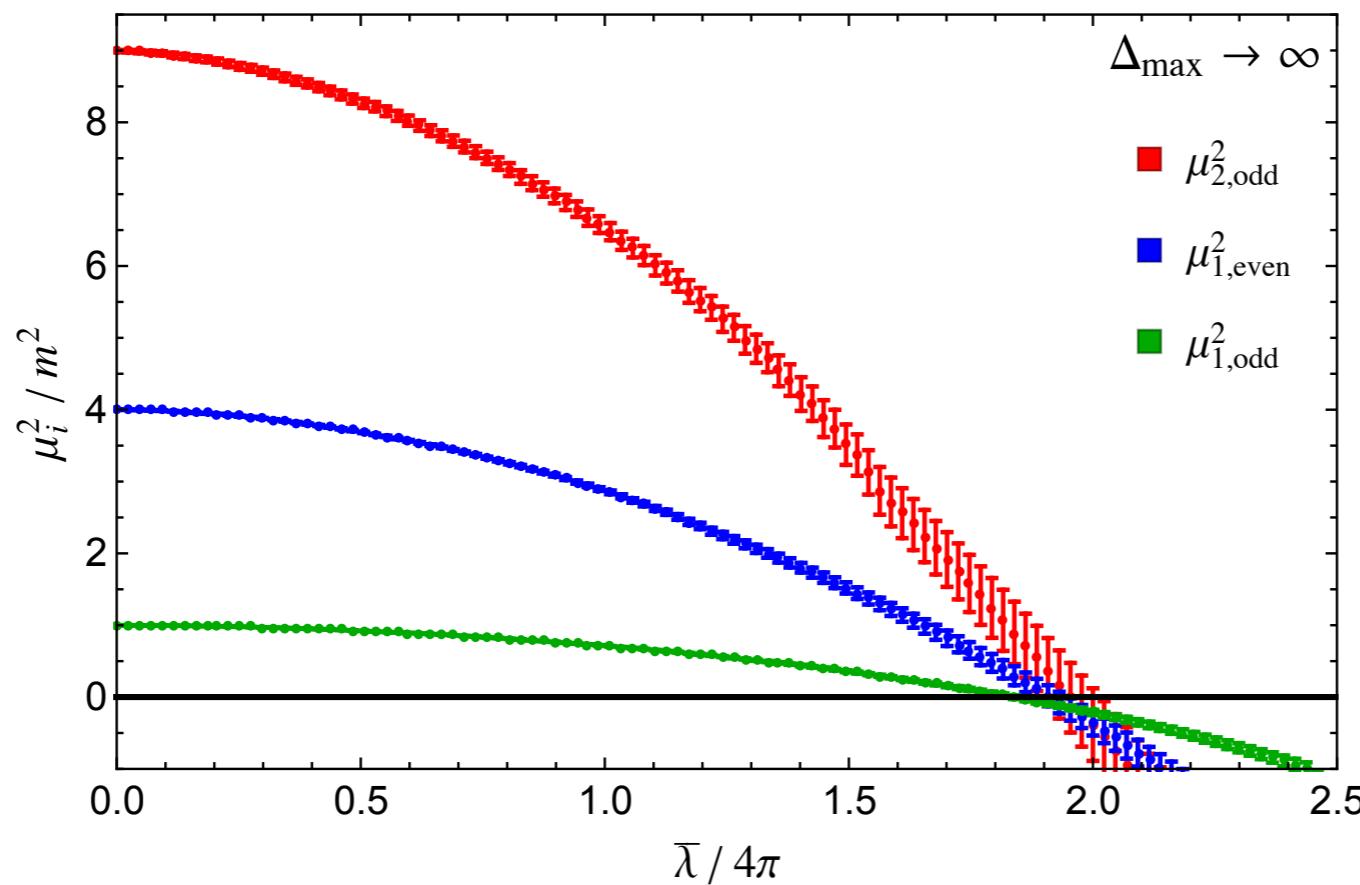
$$\mathcal{L}_{IR} = \frac{1}{2}\psi i\partial_+ \psi + \frac{1}{2}\chi i\partial_- \chi - m_f \chi \psi$$

2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_*$ :

$$\mathcal{L}_{IR} = \frac{1}{2}\psi i\partial_+ \psi + \frac{1}{2}\chi i\partial_- \chi - m_f \chi \psi$$



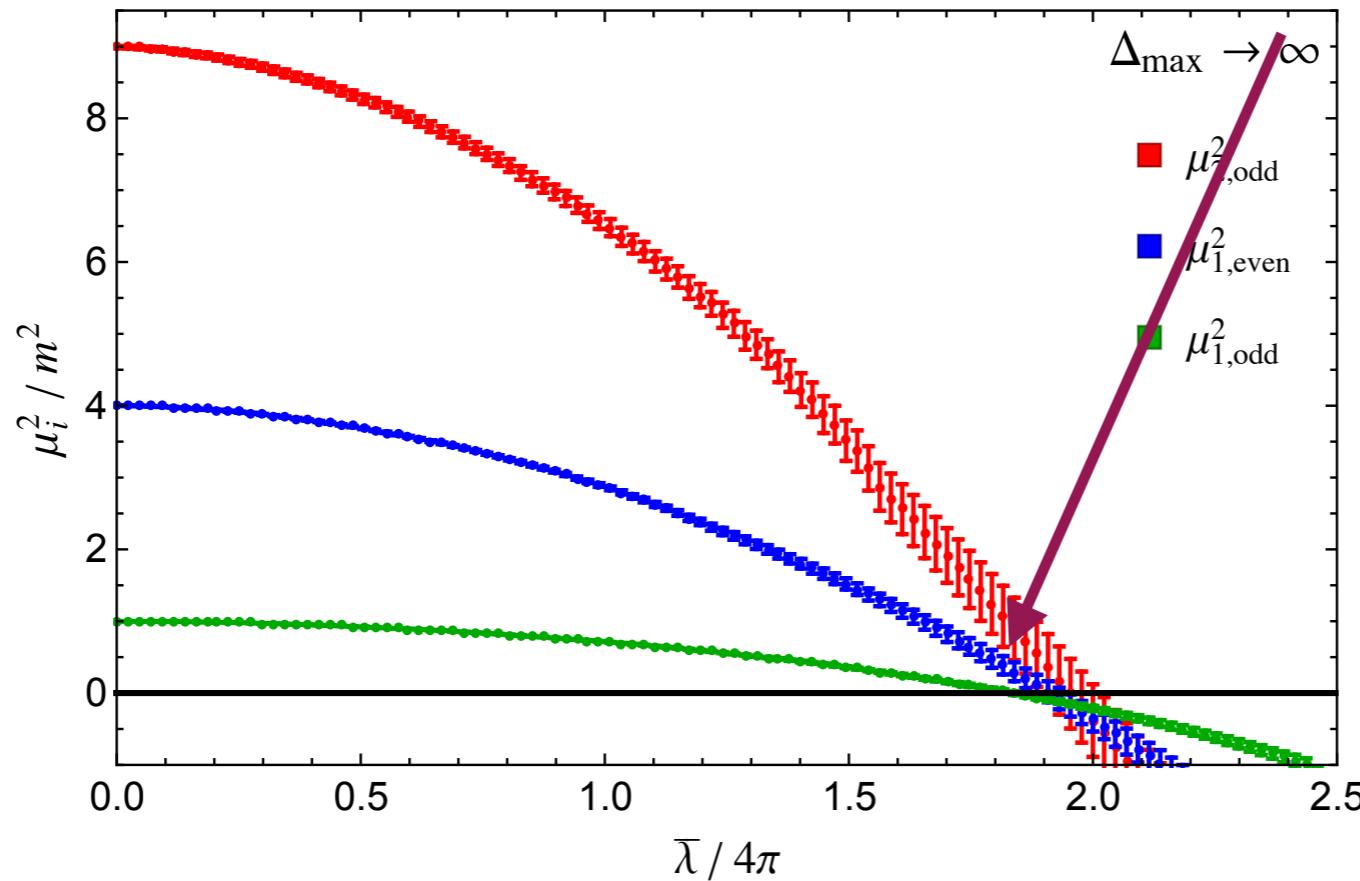
2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_*$ :

$$\mathcal{L}_{IR} = \frac{1}{2}\psi i\partial_+ \psi + \frac{1}{2}\chi i\partial_- \chi - m_f \chi \psi$$

Strange linear behavior!



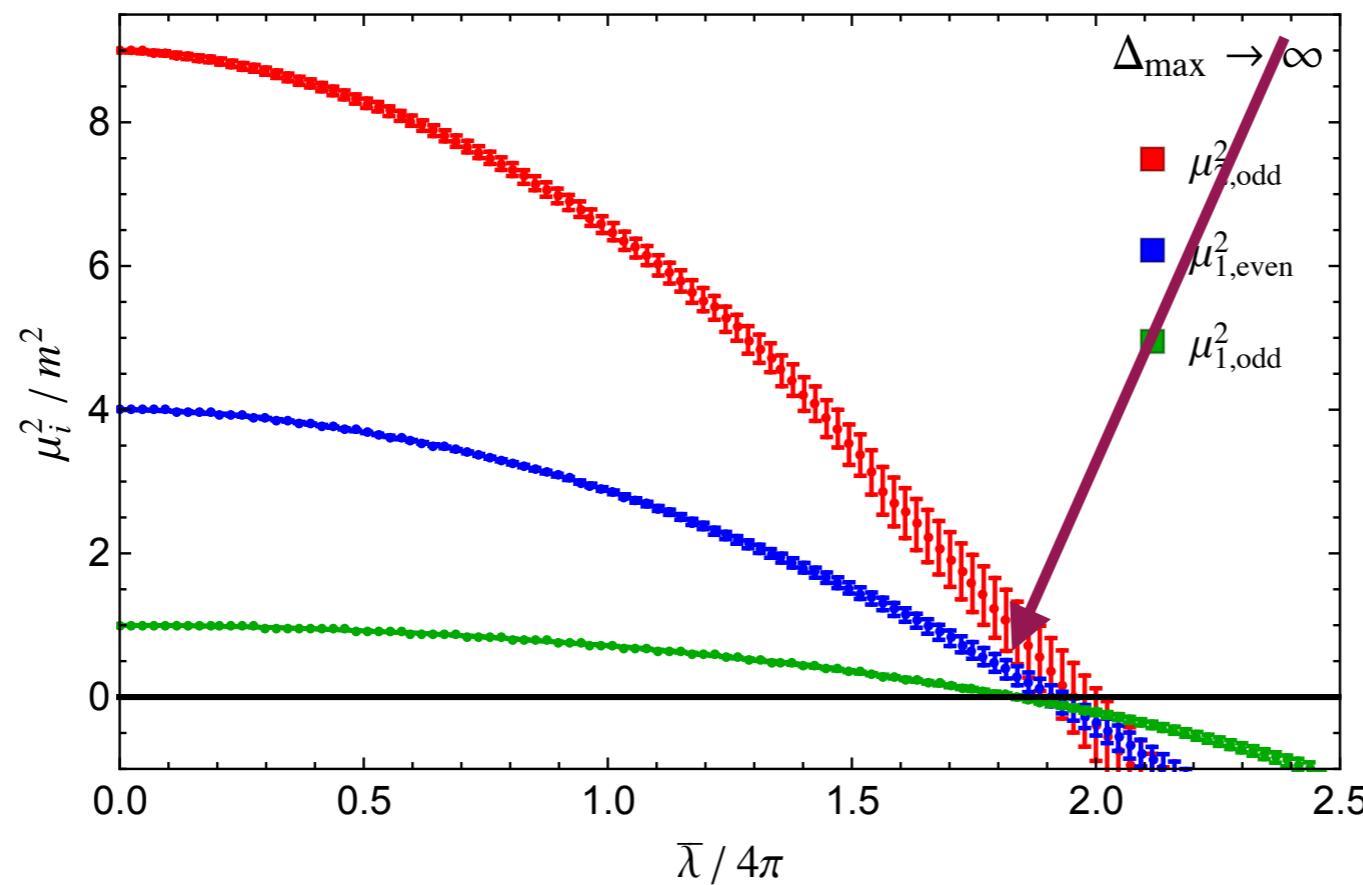
2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_*$ :

$$\mathcal{L}_{IR} = \frac{1}{2}\psi i\partial_+ \psi + \frac{1}{2}\chi i\partial_- \chi - m_f \chi \psi$$

Strange linear behavior!



$$m_{Gap}^2 \neq (\lambda_* - \lambda)^{2\nu}$$

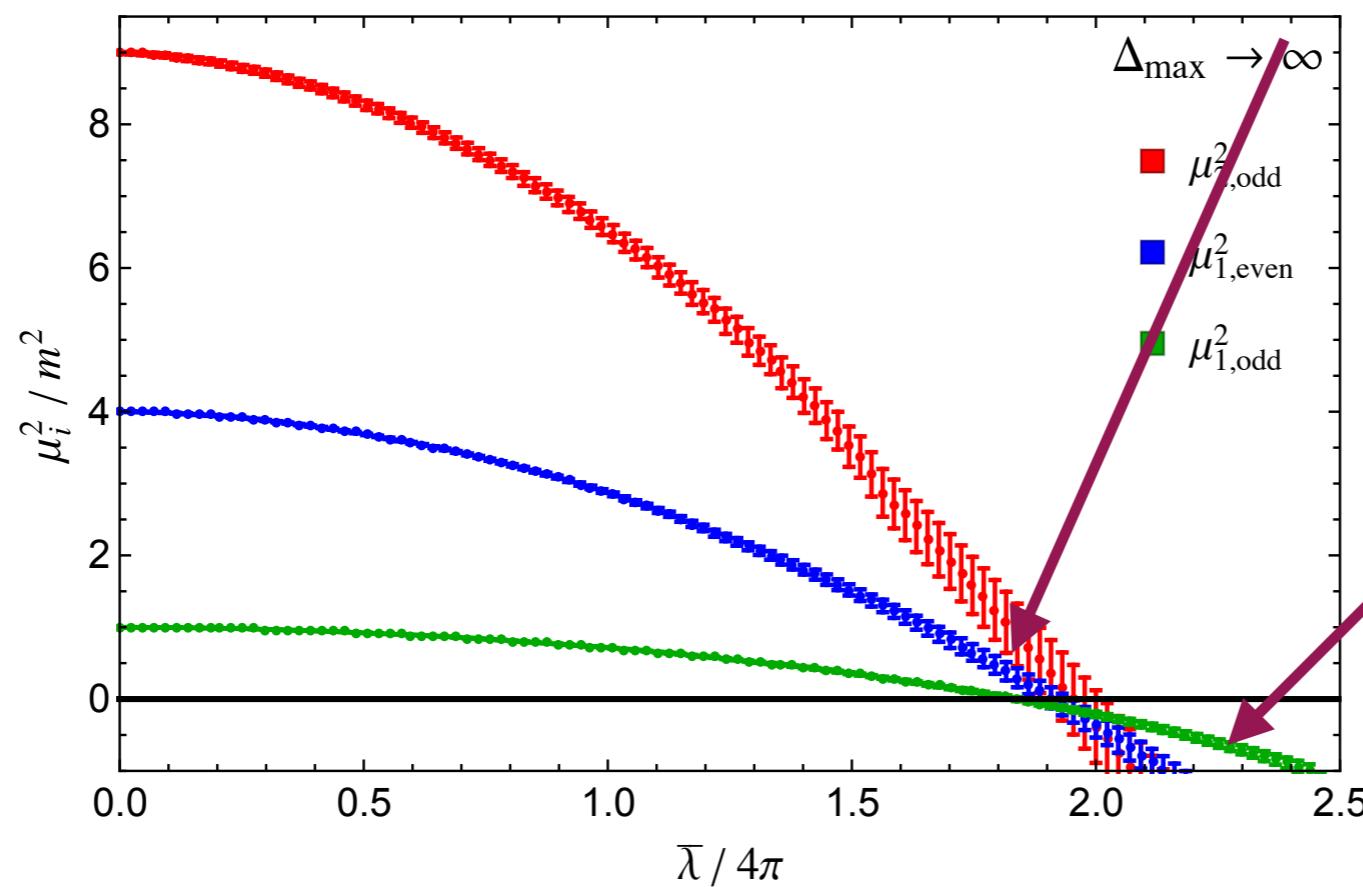
2D Ex:

$$\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - \frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4$$

Flows to the 2D Ising for  $\lambda \rightarrow \lambda_*$ :

$$\mathcal{L}_{IR} = \frac{1}{2}\psi i\partial_+ \psi + \frac{1}{2}\chi i\partial_- \chi - m_f \chi \psi$$

Strange linear behavior!



$$m_{Gap}^2 \neq (\lambda_* - \lambda)^{2\nu}$$

NEC violation!

## Problem: “LC - Zero Modes”

## Problem: “LC - Zero Modes”

$$P_+ = \frac{P_\perp^2 + m^2}{2P_-}$$

## Problem: “LC - Zero Modes”

$$P_+ = \frac{P_\perp^2 + m^2}{2P_-} \rightarrow \infty, P_- \rightarrow 0$$

## Problem: “LC - Zero Modes”

$$P_+ = \frac{P_\perp^2 + m^2}{2P_-} \rightarrow \infty, P_- \rightarrow 0$$

I.e. “Zero-Modes” have large energies  
and need to be integrated out properly!

# A proposal for incorporating LC Zero-Modes

# A proposal for incorporating LC Zero-Modes

$$U(x^+, 0) \equiv \mathcal{T}\{e^{-i \int_0^{x^+} dy^+ V_{LC}(y^+)}\}$$

# A proposal for incorporating LC Zero-Modes

$$U(x^+, 0) \equiv \mathcal{T}\{e^{-i \int_0^{x^+} dy^+ V_{LC}(y^+)}\}$$

$$H_{eff} \equiv \lim_{x^+ \rightarrow 0} i\partial_+ U(x^+, 0)$$

# A proposal for incorporating LC Zero-Modes

$$U(x^+, 0) \equiv \mathcal{T}\{e^{-i\int_0^{x^+} dy^+ V_{LC}(y^+)}\}$$

$$H_{eff} \equiv \lim_{x^+ \rightarrow 0} i\partial_+ U(x^+, 0)$$

$$\langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle = \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

$$-\frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots$$

# A proposal for incorporating LC Zero-Modes

$$U(x^+, 0) \equiv \mathcal{T}\{e^{-i \int_0^{x^+} dy^+ V_{LC}(y^+)}\}$$

$$H_{eff} \equiv \lim_{x^+ \rightarrow 0} i\partial_+ U(x^+, 0)$$

$$\langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle = \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

$$- \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+) V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots$$

ZM associated with  $\delta(y_{ij}^+)$ :

# A proposal for incorporating LC Zero-Modes

$$U(x^+, 0) \equiv \mathcal{T}\{e^{-i \int_0^{x^+} dy^+ V_{LC}(y^+)}\}$$

$$H_{eff} \equiv \lim_{x^+ \rightarrow 0} i\partial_+ U(x^+, 0)$$

$$\langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle = \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

$$- \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+) V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots$$

**ZM associated with  $\delta(y_{ij}^+)$ :**

$$\int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+) V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle \sim \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | \delta H_{eff}(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

Back to our simple ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

Back to our simple ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

$$\langle p|U(x^+)|p\rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p|\mathcal{T}\{\lambda\phi(y_1)\frac{g}{3!}\phi^3(y_2)\}|p\rangle$$

Back to our simple ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

$$\langle p|U(x^+)|p\rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p|\mathcal{T}\{\lambda\phi(y_1)\frac{g}{3!}\phi^3(y_2)\}|p\rangle$$

$$\sim \int_0^t dy_1^+ \ dy_2^+ \int dq_+ \frac{i\lambda g}{-m^2} e^{-i(q_+ + p_+)y_{12}^+}$$

Back to our simple ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

$$\langle p|U(x^+)|p\rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p|\mathcal{T}\{\lambda\phi(y_1)\frac{g}{3!}\phi^3(y_2)\}|p\rangle$$

$$\sim \int_0^t dy_1^+ \ dy_2^+ \int dq_+ \frac{i\lambda g}{-m^2} e^{-i(q_+ + p_+)y_{12}^+}$$

Analytic in  $q^+$

Back to our simple ex:  $\mathcal{V} = \frac{g}{3!}\phi^3 + \lambda\phi$

$$\langle p|U(x^+)|p\rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p|\mathcal{T}\{\lambda\phi(y_1)\frac{g}{3!}\phi^3(y_2)\}|p\rangle$$

$$\sim \int_0^t dy_1^+ \ dy_2^+ \int dq_+ \frac{i\lambda g}{-m^2} e^{-i(q_+ + p_+)y_{12}^+}$$

**Analytic in  $q^+$**

In contrast with ET:

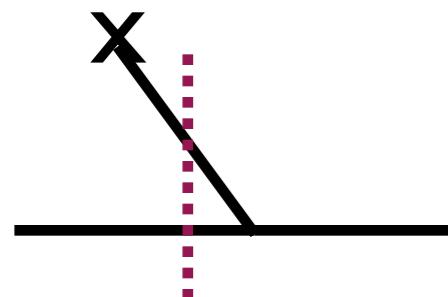
Back to our simple ex:  $\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi$

$$\langle p | U(x^+) | p \rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p | \mathcal{T} \{ \lambda \phi(y_1) \frac{g}{3!} \phi^3(y_2) \} | p \rangle$$

$$\sim \int_0^t dy_1^+ \ dy_2^+ \int dq_+ \frac{i\lambda g}{-m^2} e^{-i(q_+ + p_+)y_{12}^+}$$

**Analytic in  $q^+$**

In contrast with ET:



$$\sim \int_0^t dt_1 \ dt_2 \int dq_0 \frac{i\lambda g}{q_0^2 - m^2} e^{-iq_0 t_{12}}$$

Back to our simple ex:

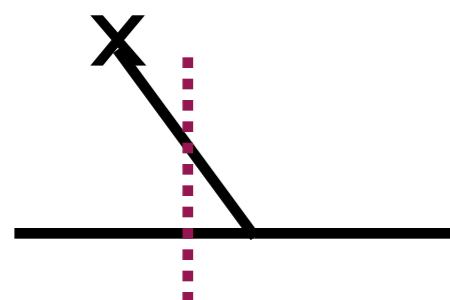
$$\mathcal{V} = \frac{g}{3!} \phi^3 + \lambda \phi$$

$$\langle p | U(x^+) | p \rangle \supset - \int_0^{x^+} d^d y_1 \ d^d y_2 \ \langle p | \mathcal{T} \{ \lambda \phi(y_1) \frac{g}{3!} \phi^3(y_2) \} | p \rangle$$

$$\sim \int_0^t dy_1^+ \ dy_2^+ \int dq_+ \frac{i \lambda g}{-m^2} e^{-i(q_+ + p_+)y_{12}^+}$$

Analytic in  $q^+$

In contrast with ET:



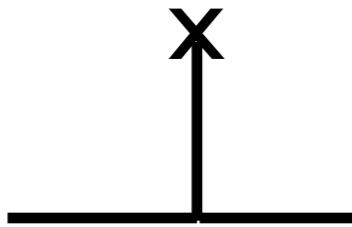
$$\sim \int_0^t dt_1 \ dt_2 \int dq_0 \frac{i \lambda g}{q_0^2 - m^2} e^{-iq_0 t_{12}}$$

Non-Analytic in  $q_0$

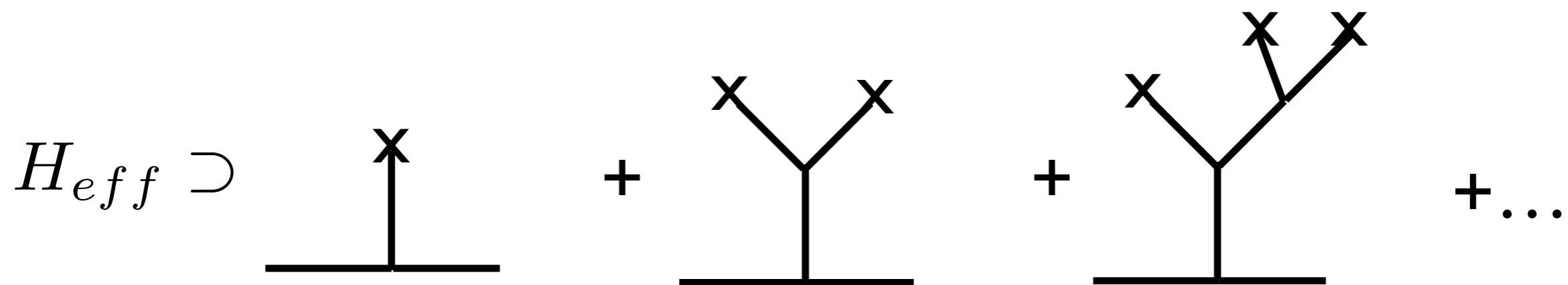
Focusing on tree-level (but to all orders):

Focusing on tree-level (but to all orders):

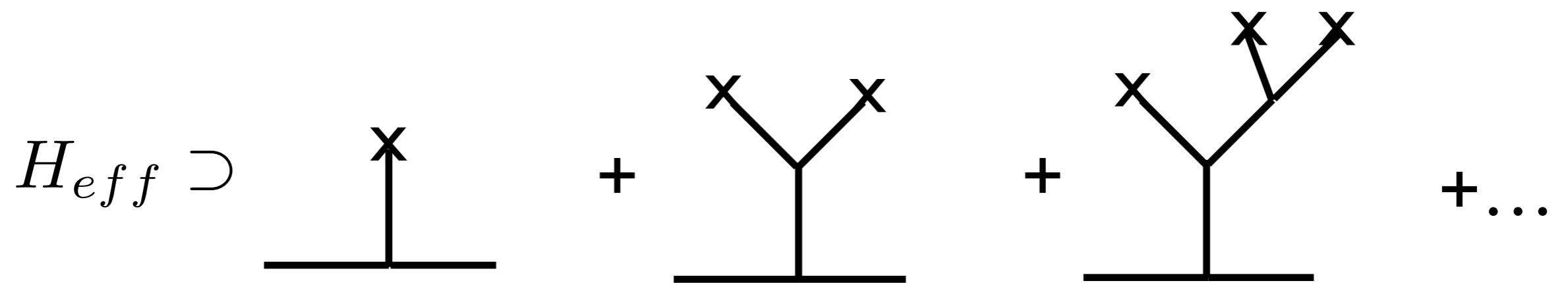
$$H_{eff} \supset$$



Focusing on tree-level (but to all orders):



Focusing on tree-level (but to all orders):



$$= \int d^{d-1}x \ v(\lambda, g) \phi^2$$

Focusing on tree-level (but to all orders):

$$H_{eff} \supset \begin{array}{c} \text{---} \\ | \\ X \end{array} + \begin{array}{c} \text{---} \\ | \\ X \\ | \\ X \end{array} + \begin{array}{c} \text{---} \\ | \\ X \\ | \\ X \\ | \\ X \end{array} + \dots$$
$$= \int d^{d-1}x \ v(\lambda, g) \phi^2$$

Prescription reproduces missing vevs!

## In general:

$$\begin{aligned} \langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle &= \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle \\ &\quad - \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots \end{aligned}$$

In general:

$$\begin{aligned}\langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle &= \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle \\ &\quad - \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots\end{aligned}$$

Focusing on 2nd Ord:

In general:

$$\begin{aligned} \langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle &= \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle \\ &\quad - \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots \end{aligned}$$

Focusing on 2nd Ord:

$$\int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle \sim \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | \delta H_{eff}(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

## In general:

$$\begin{aligned}\langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle &= \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle \\ &\quad - \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots\end{aligned}$$

Focusing on 2nd Ord:

$$\int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle \sim \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | \delta H_{eff}(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

Happens when 4-pt fcn loses  
spectral decomposition on the LC

In general:

$$\begin{aligned} \langle \mathcal{O}, P_-, \mu | U(x^+, 0) | \mathcal{O}, P_-, \mu' \rangle &= \langle \mathcal{O}, P_-, \mu | \mathcal{O}, P_-, \mu' \rangle - i \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | V(y_1^+) | \mathcal{O}, P_-, \mu' \rangle \\ &\quad - \frac{1}{2} \int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle + \dots \end{aligned}$$

Focusing on 2nd Ord:

$$\int_0^{x^+} dy_1^+ dy_2^+ \langle \mathcal{O}, P_-, \mu | \mathcal{T}\{V(y_1^+)V(y_2^+)\} | \mathcal{O}, P_-, \mu' \rangle \sim \int_0^{x^+} dy_1^+ \langle \mathcal{O}, P_-, \mu | \delta H_{eff}(y_1^+) | \mathcal{O}, P_-, \mu' \rangle$$

Happens when 4-pt fcn loses  
spectral decomposition on the LC

Only possible if some 3pt-functions vanish on the LC:

$$\Delta' = \Delta + \Delta_R + 2n$$

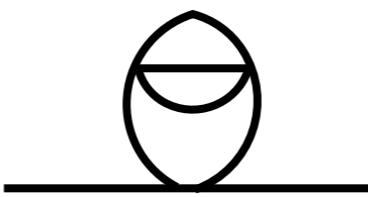
# Prescription maps ET parameters to LC parameters

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

$$H_{eff} \supset \text{Diagram} + \dots$$
A Feynman diagram representing a loop correction to the effective Hamiltonian. It consists of a horizontal line with a vertical loop attached to it. The loop has a small horizontal line segment at its top center.

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

$$H_{eff} \supset \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4 + \dots$$

Serone, Spada, and Villadoro:

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

$$H_{eff} \supset \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4 + \dots$$

Serone, Spada, and Villadoro:

$$m_{LC}^2 = m_{ET}^2 \left[ 1 + \sum_n a_n \left( \frac{\lambda}{m_{ET}^2} \right)^n \right]$$

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

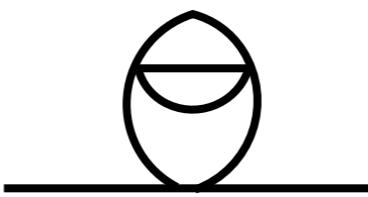
$$H_{eff} \supset \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4 + \dots$$

Serone, Spada, and Villadoro:

$$m_{LC}^2 = m_{ET}^2 \left[ 1 + \sum_n a_n \left( \frac{\lambda}{m_{ET}^2} \right)^n \right] = m_{ET}^2 + 12\lambda \langle \phi^2 \rangle_{PT}$$

## Prescription maps ET parameters to LC parameters

2D Ex:  $\mathcal{V} = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4$

$$H_{eff} \supset \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} + \dots$$


Serone, Spada, and Villadoro:

$$m_{LC}^2 = m_{ET}^2 \left[ 1 + \sum_n a_n \left( \frac{\lambda}{m_{ET}^2} \right)^n \right] = m_{ET}^2 + 12\lambda \langle \phi^2 \rangle_{PT}$$

$$m_{Gap}^2 = m_{ET}^2 \left[ 1 + \sum_n c_n \left( \frac{\lambda}{m_{ET}^2} \right)^n \right]$$

Using entirely ET data you find the LC series:

$$m_{Gap}^2 = m_{LC}^2 \left[ 1 + \sum_n \tilde{c}_n \left( \frac{\lambda}{m_{LC}^2} \right)^n \right]$$

Using entirely ET data you find the LC series:

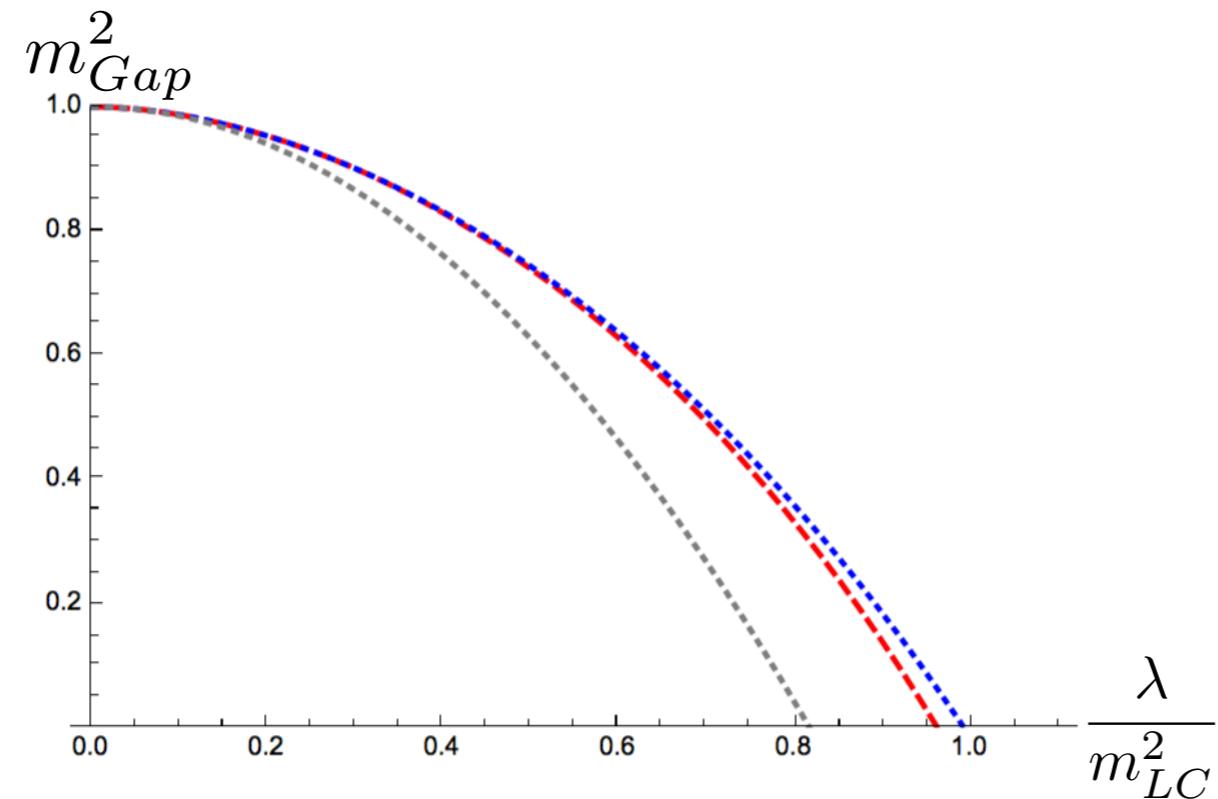
$$m_{Gap}^2 = m_{LC}^2 \left[ 1 + \sum_n \tilde{c}_n \left( \frac{\lambda}{m_{LC}^2} \right)^n \right]$$

Borel resum this series, minimizing over  
Borel “voodoo” parameters (as in Serone et al.):

Using entirely ET data you find the LC series:

$$m_{Gap}^2 = m_{LC}^2 \left[ 1 + \sum_n \tilde{c}_n \left( \frac{\lambda}{m_{LC}^2} \right)^n \right]$$

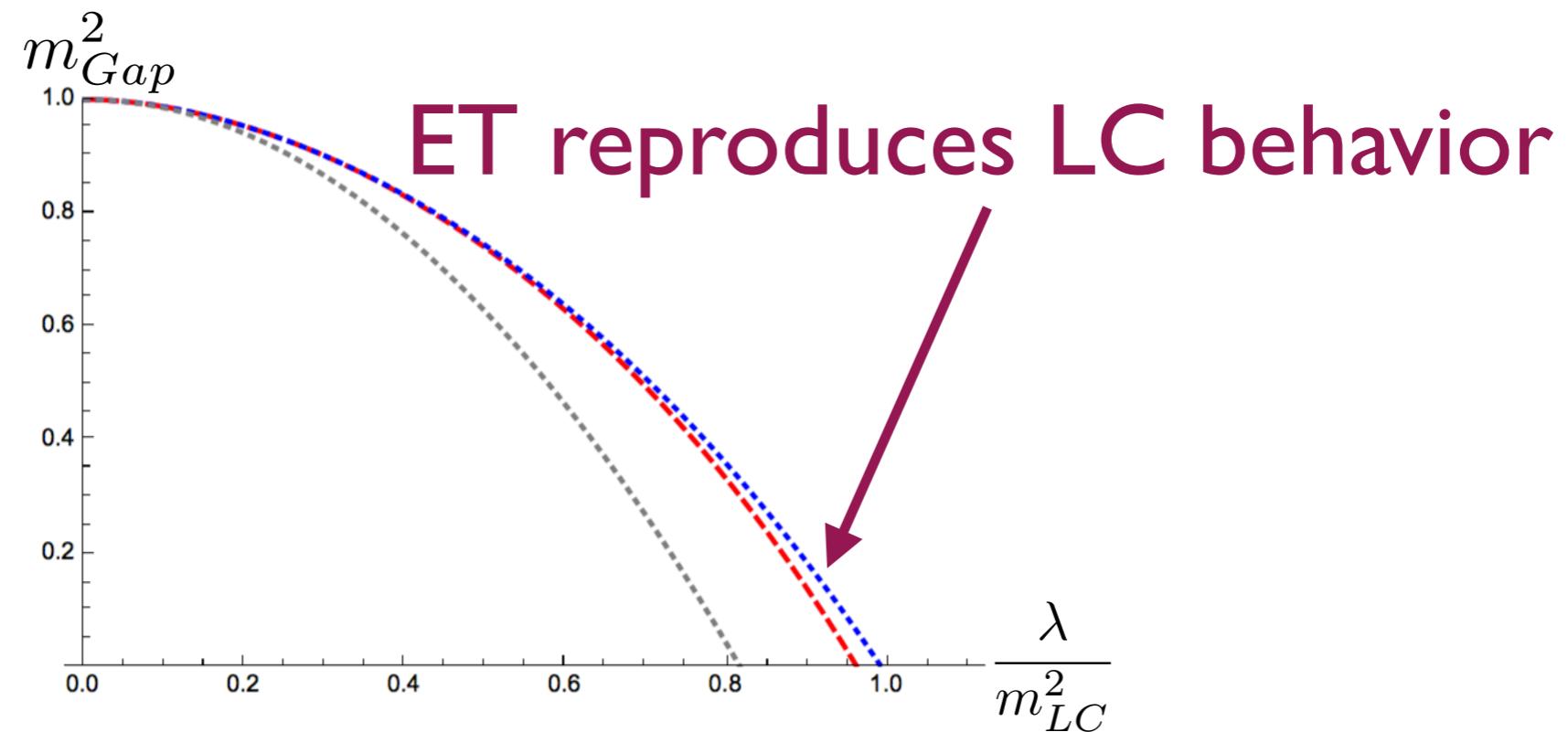
Borel resum this series, minimizing over  
Borel “voodoo” parameters (as in Serone et al.):



Using entirely ET data you find the LC series:

$$m_{Gap}^2 = m_{LC}^2 \left[ 1 + \sum_n \tilde{c}_n \left( \frac{\lambda}{m_{LC}^2} \right)^n \right]$$

Borel resum this series, minimizing over  
Borel “voodoo” parameters (as in Serone et al.):



And now with SUSY

## And now with SUSY

The 2D theory: Real  $\phi, \psi$ :

$$W = h\Phi + \frac{g}{3!}\Phi^3$$

## And now with SUSY

The 2D theory: Real  $\phi, \psi$ :

$$W = h\Phi + \frac{g}{3!}\Phi^3$$

SUSY variation:  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2} \langle \phi^2 \rangle$

## And now with SUSY

The 2D theory: Real  $\phi, \psi$ :

$$W = h\Phi + \frac{g}{3!}\Phi^3$$

SUSY variation:  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2}\langle \phi^2 \rangle$

$h < 0$ : **SUSY preserved but  $Z_2$  broken:**

$$m_\phi = m_\psi \neq 0$$

## And now with SUSY

The 2D theory: Real  $\phi, \psi$ :

$$W = h\Phi + \frac{g}{3!}\Phi^3$$

SUSY variation:  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2}\langle \phi^2 \rangle$

$h < 0$ : **SUSY preserved but  $Z_2$  broken:**

$$m_\phi = m_\psi \neq 0$$

$h > 0$ : **SUSY broken but  $Z_2$  preserved:**

$$m_\phi \neq 0, m_\psi = 0$$

The phase transition is captured by the  
Tri-Critical Ising Model

The phase transition is captured by the  
Tri-Critical Ising Model

$$\Phi_{TIM} = \epsilon + \theta\bar{\psi} + \bar{\theta}\psi + \theta\bar{\theta}\epsilon'$$

# The phase transition is captured by the Tri-Critical Ising Model

$$\Phi_{TIM} = \epsilon + \theta\bar{\psi} + \bar{\theta}\psi + \theta\bar{\theta}\epsilon'$$

Integrable SUSY deformation:  $\delta P_+ = \int dx^- \lambda \epsilon'(x^-)$

The phase transition is captured by the  
Tri-Critical Ising Model

$$\Phi_{TIM} = \epsilon + \theta\bar{\psi} + \bar{\theta}\psi + \theta\bar{\theta}\epsilon'$$

Integrable SUSY deformation:  $\delta P_+ = \int dx^- \lambda \epsilon'(x^-)$

$$\nu = 1.25$$

Can we maintain SUSY (in the SUSY phase)?

# Can we maintain SUSY (in the SUSY phase)?

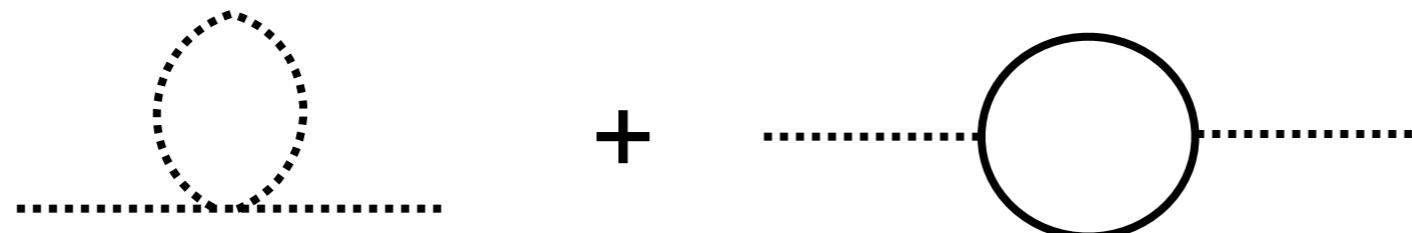
Worry:

Normal ordering and cutoff effects spoil SUSY:

# Can we maintain SUSY (in the SUSY phase)?

Worry:

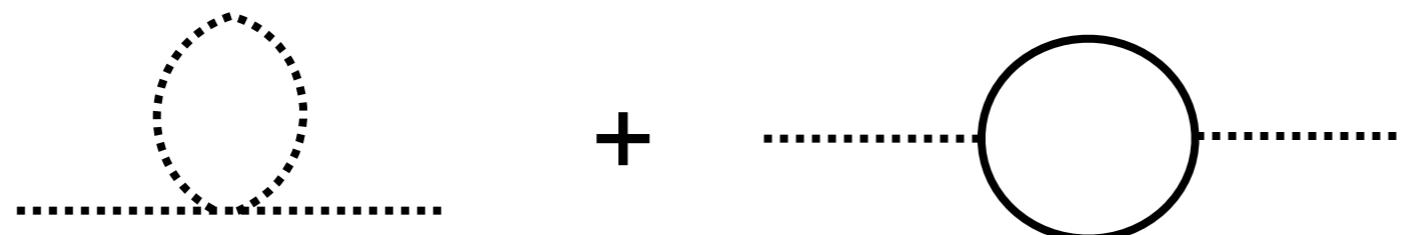
Normal ordering and cutoff effects spoil SUSY:



# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:

$$\text{---} \circ \text{---} + \text{---} \circ \text{---} = 0 \times g^2 \text{Log}(\Lambda)$$


# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:

$$\text{Diagram with crossed lines} + \text{Diagram with circle} = 0 \times g^2 \text{Log}(\Lambda)$$

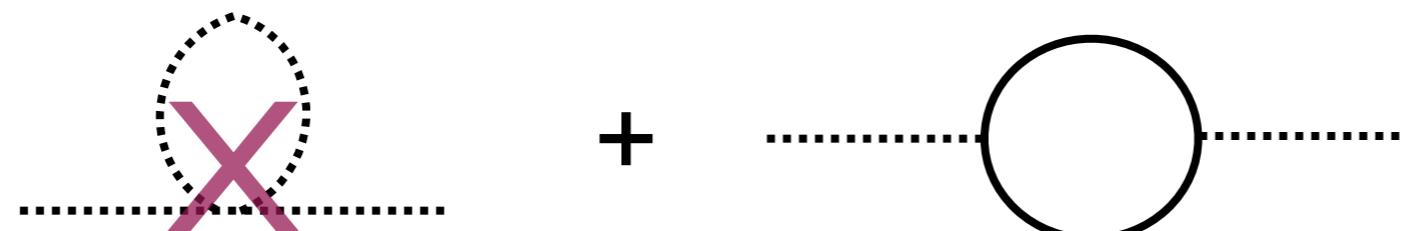
$\therefore \phi^4 :$



# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:



A Feynman diagram for a four-point vertex. It consists of four external lines meeting at a central point. A horizontal dotted line from the left and a vertical dotted line from the bottom meet at the top-left vertex. From the top-left vertex, a horizontal dotted line goes right and a vertical dotted line goes down. From the bottom-right vertex, a horizontal dotted line goes left and a vertical dotted line goes up. A red 'X' is drawn through the top-left and bottom-right vertices, indicating that this diagram is not allowed in normal ordering.

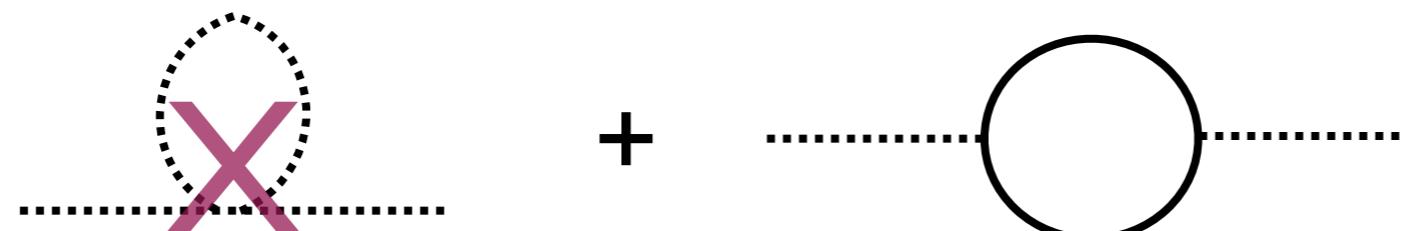
$$: \phi^4 : + \text{ (loop diagram)} = 0 \times g^2 \text{Log}(\Lambda)$$

**Instead:**  $(Q_+)_\text{trunc} = \int dx^- \left[ (m\phi + \frac{g}{2}\phi^2)\psi(x^-) \right]_\text{trunc}$

# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:



A Feynman diagram consisting of a horizontal dotted line with a vertical dashed line crossing it from top to bottom. A red 'X' is drawn over the crossing point. To the right of the diagram is a plus sign, followed by another horizontal dotted line with a circular loop attached to its right end.

$$: \phi^4 : + \text{ (loop diagram)} = 0 \times g^2 \text{Log}(\Lambda)$$

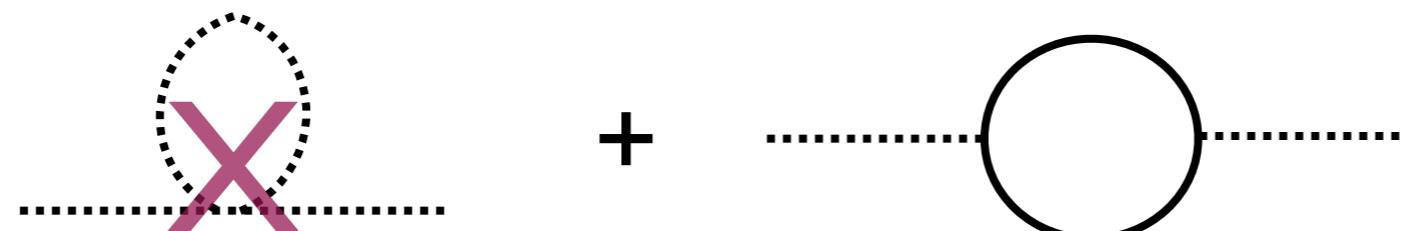
**Instead:**  $(Q_+)_\text{trunc} = \int dx^- \left[ (m\phi + \frac{g}{2}\phi^2)\psi(x^-) \right]_\text{trunc}$

$$P_+ \equiv (Q_+)_\text{trunc}^2$$

# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:



A Feynman diagram consisting of a horizontal dotted line with a vertical dashed line crossing it from top to bottom. A circular loop is attached to the right end of the horizontal line. A large red 'X' is drawn over the crossing point.

$$+ \quad \text{---} \circ \text{---} = 0 \times g^2 \text{Log}(\Lambda)$$

$\therefore \phi^4 :$

**Instead:**  $(Q_+)_{trunc} = \int dx^- \left[ (m\phi + \frac{g}{2}\phi^2)\psi(x^-) \right]_{trunc}$

$$P_+ \equiv (Q_+)_{trunc}^2 \neq (P_+)_{N.O.}$$

# Can we maintain SUSY (in the SUSY phase)?

Worry:

Normal ordering and cutoff effects spoil SUSY:

A Feynman diagram for a four-point vertex. It consists of four external lines meeting at a central point. A horizontal dotted line from the left and a vertical dotted line from the bottom meet at the top-left vertex. From the top-left vertex, a horizontal dotted line goes right and a vertical dotted line goes down. From the bottom-right vertex, a horizontal dotted line goes left and a vertical dotted line goes up. A red 'X' is drawn through the top-left and bottom-right vertices. Below the diagram, the text  $: \phi^4 :$  is written.

+      =  $0 \times g^2 \text{Log}(\Lambda)$

**Instead:**  $(Q_+)_{trunc} = \int dx^- \left[ (m\phi + \frac{g}{2}\phi^2)\psi(x^-) \right]_{trunc}$

$$P_+ \equiv (Q_+)_{trunc}^2 \neq (P_+)_{N.O.}$$

$$Q_- = Q_-^{free}$$

SUSY is healthy for the LC

## SUSY is healthy for the LC

**Recall:**  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2} \langle \phi^2 \rangle$

## SUSY is healthy for the LC

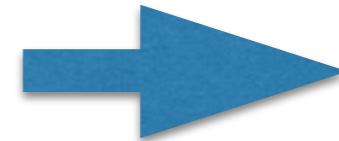
**Recall:**  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2} \langle \phi^2 \rangle$

**SUSY phase:**  $\langle \phi^2 \rangle = \frac{-2h}{g} = \langle \phi^2 \rangle_{cl}$

## SUSY is healthy for the LC

**Recall:**  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2} \langle \phi^2 \rangle$

**SUSY phase:**  $\langle \phi^2 \rangle = \frac{-2h}{g} = \langle \phi^2 \rangle_{cl}$

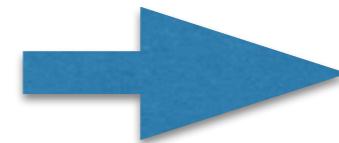


$$H_{eff} = P_{+,naive}$$

## SUSY is healthy for the LC

Recall:  $\langle \delta\psi \rangle = \langle W' \rangle = h + \frac{g}{2} \langle \phi^2 \rangle$

**SUSY phase:**  $\langle \phi^2 \rangle = \frac{-2h}{g} = \langle \phi^2 \rangle_{cl}$

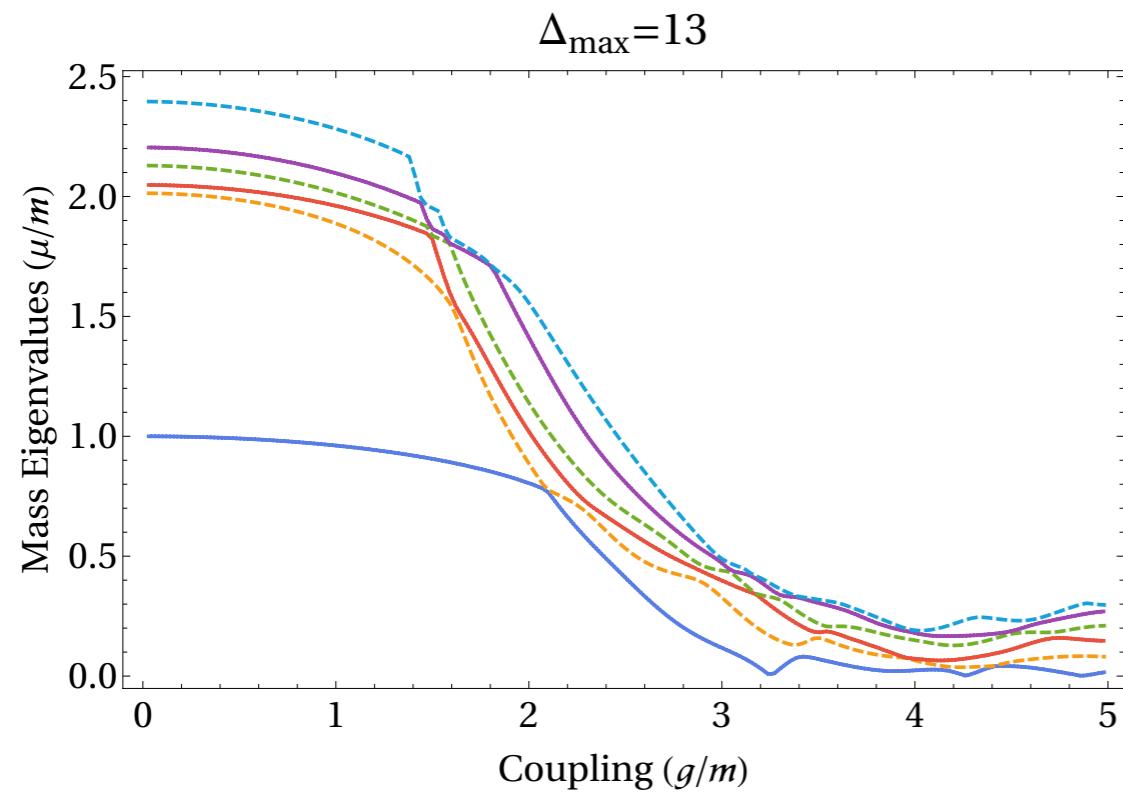


$$H_{eff} = P_{+,naive}$$

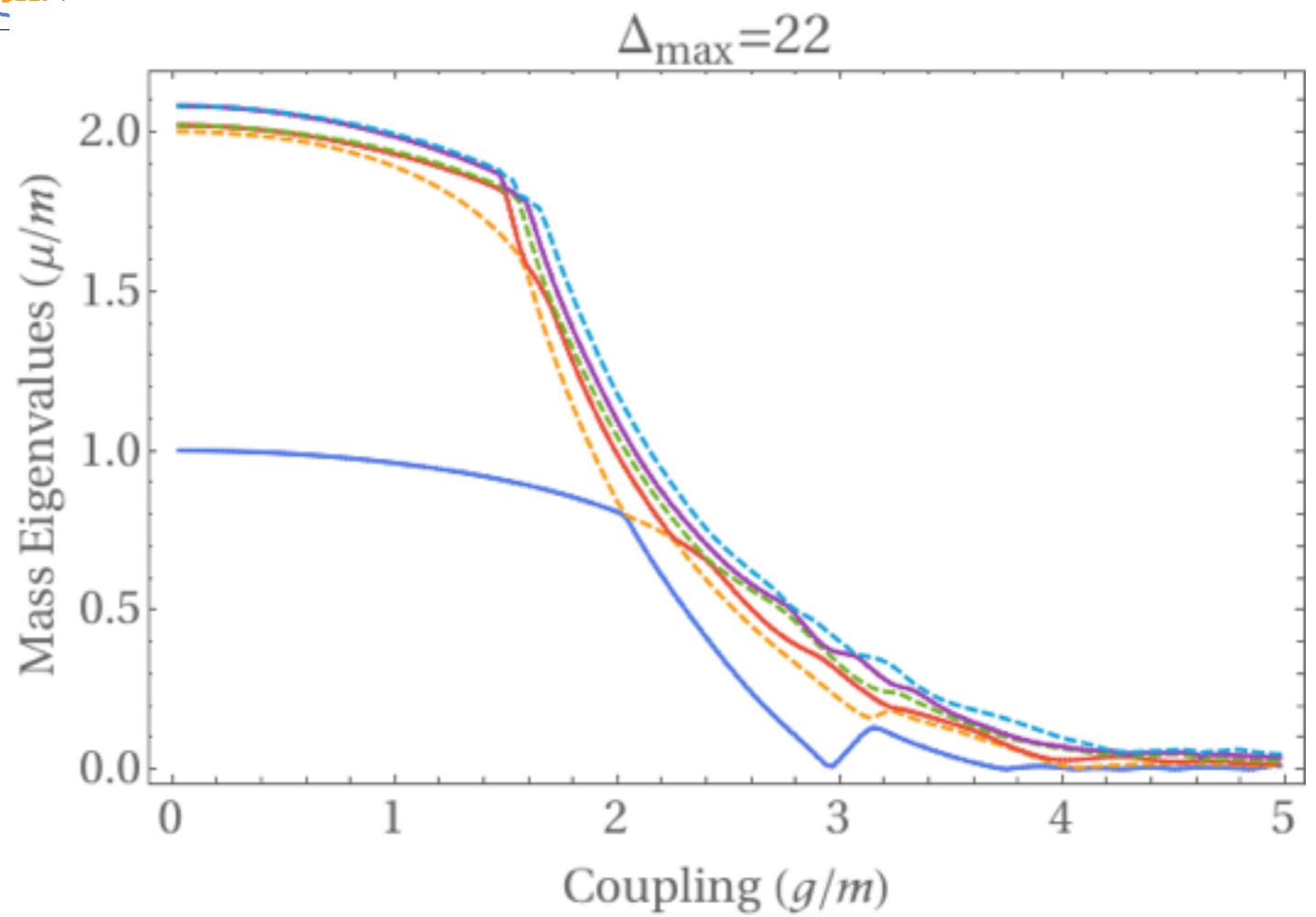
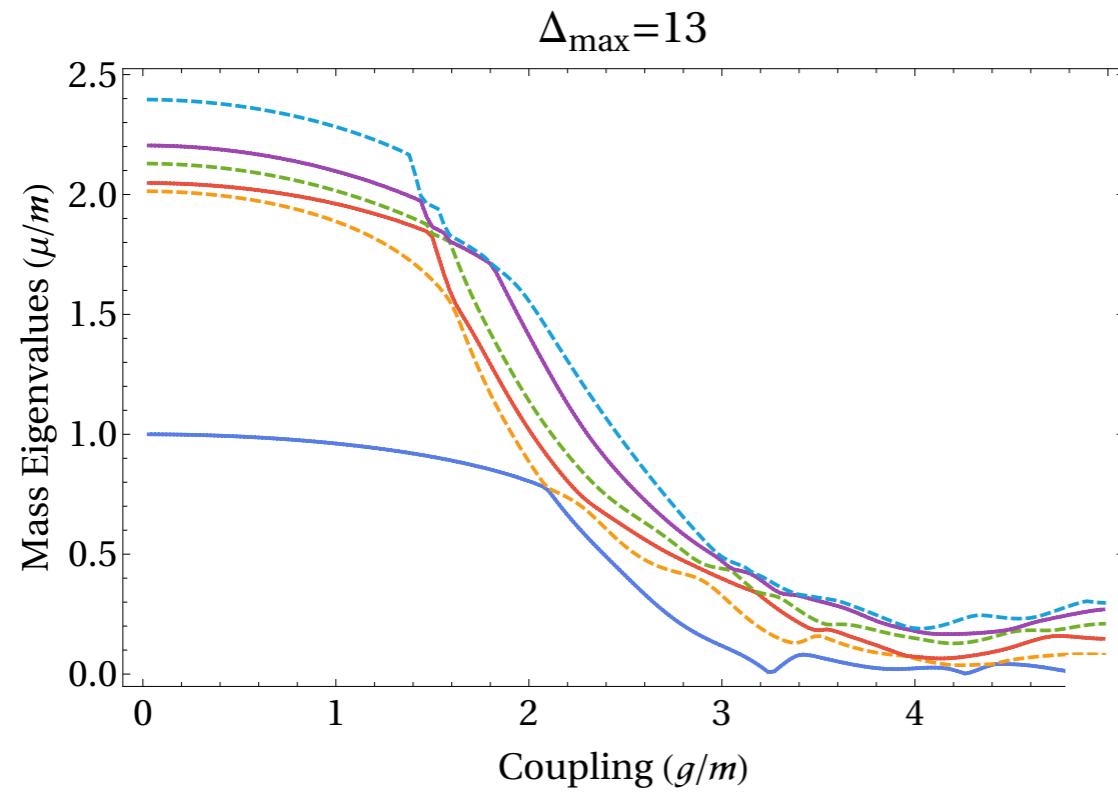
$$P_+ = Q_+^2 > 0 \text{ (NEC)}$$

# The Spectrum

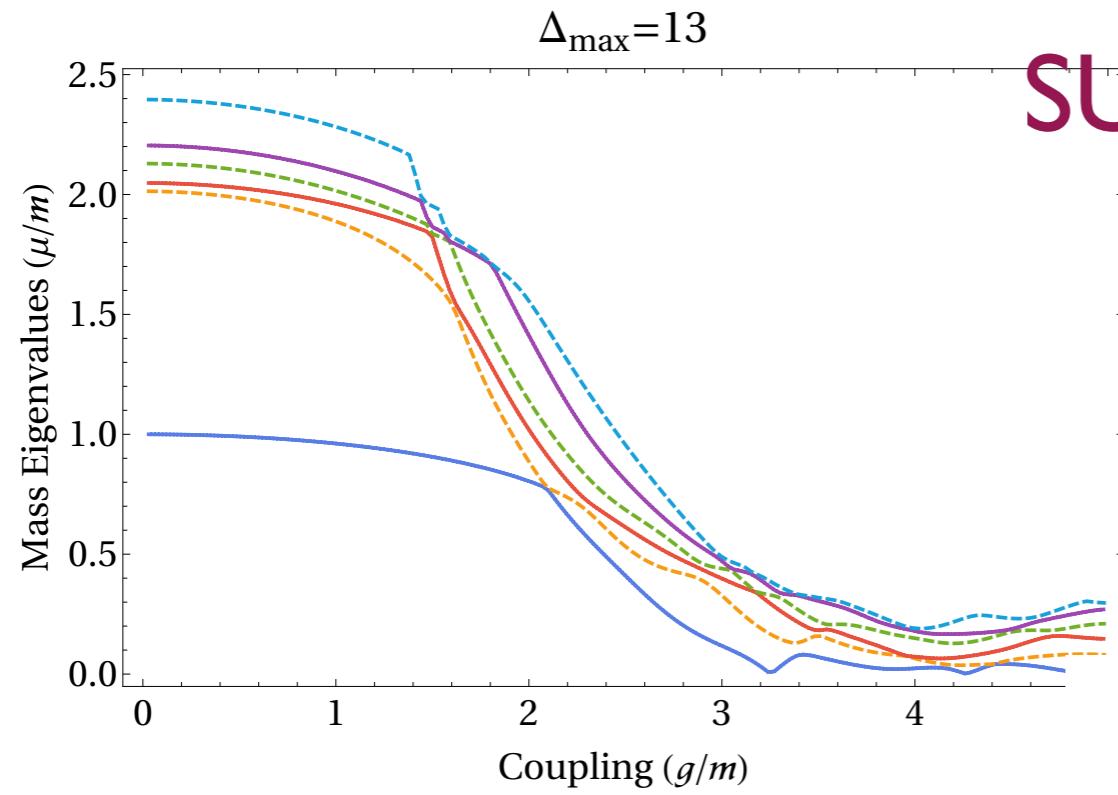
# The Spectrum



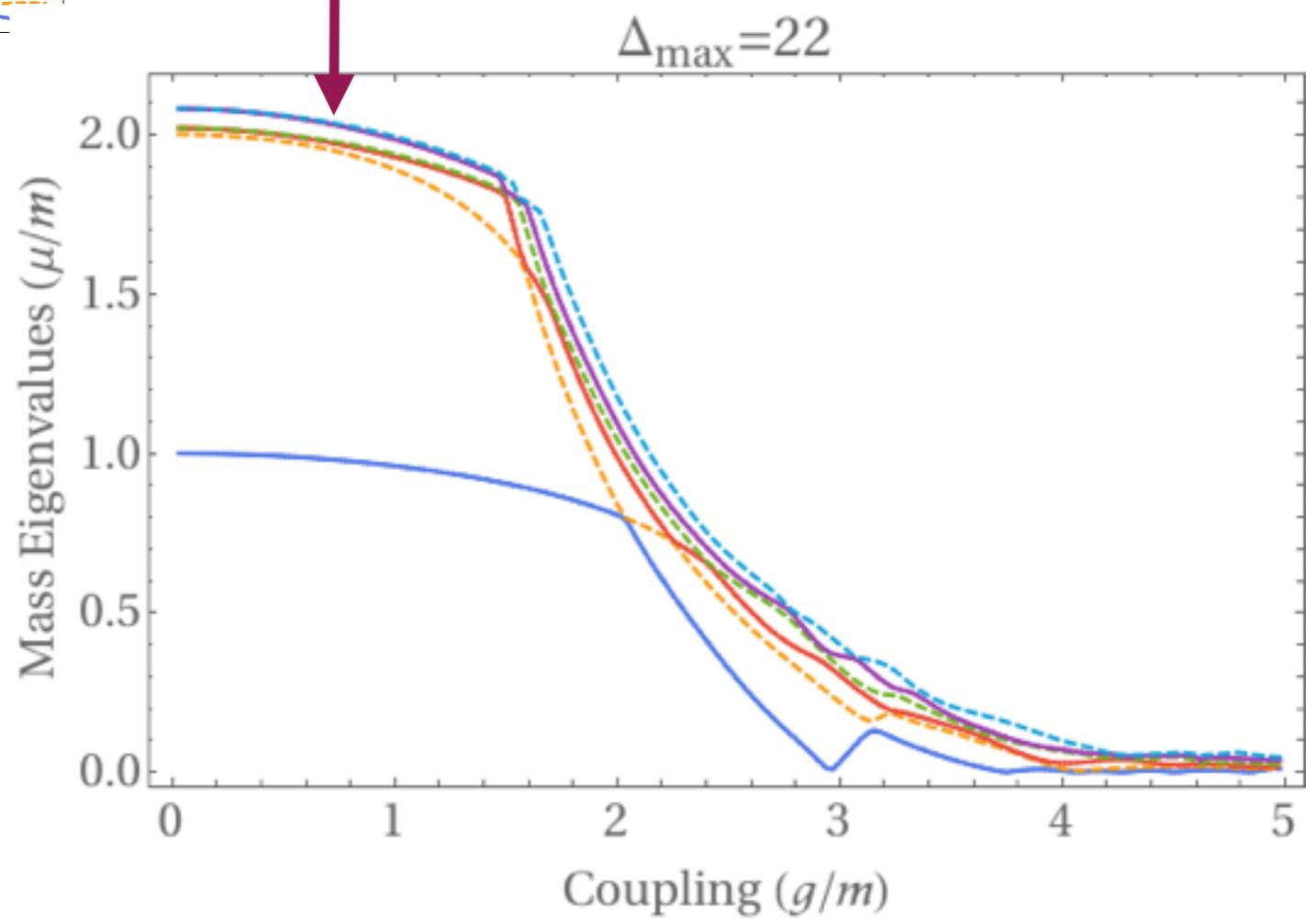
# The Spectrum



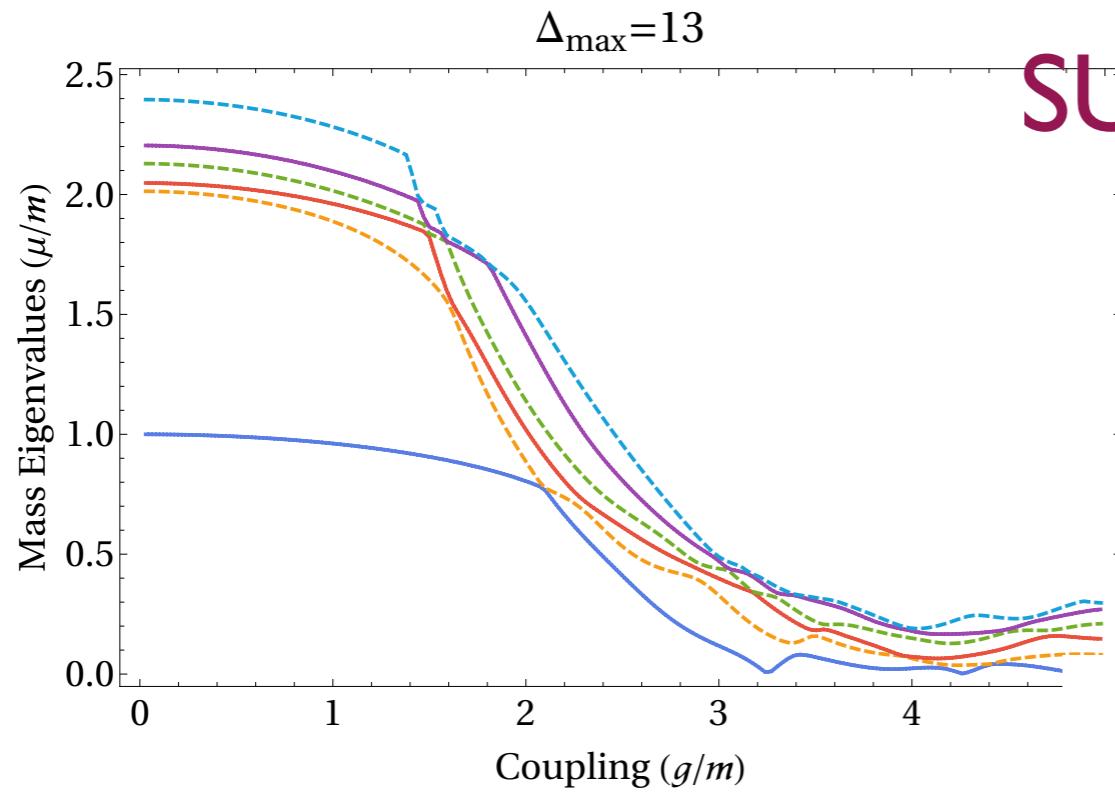
# The Spectrum



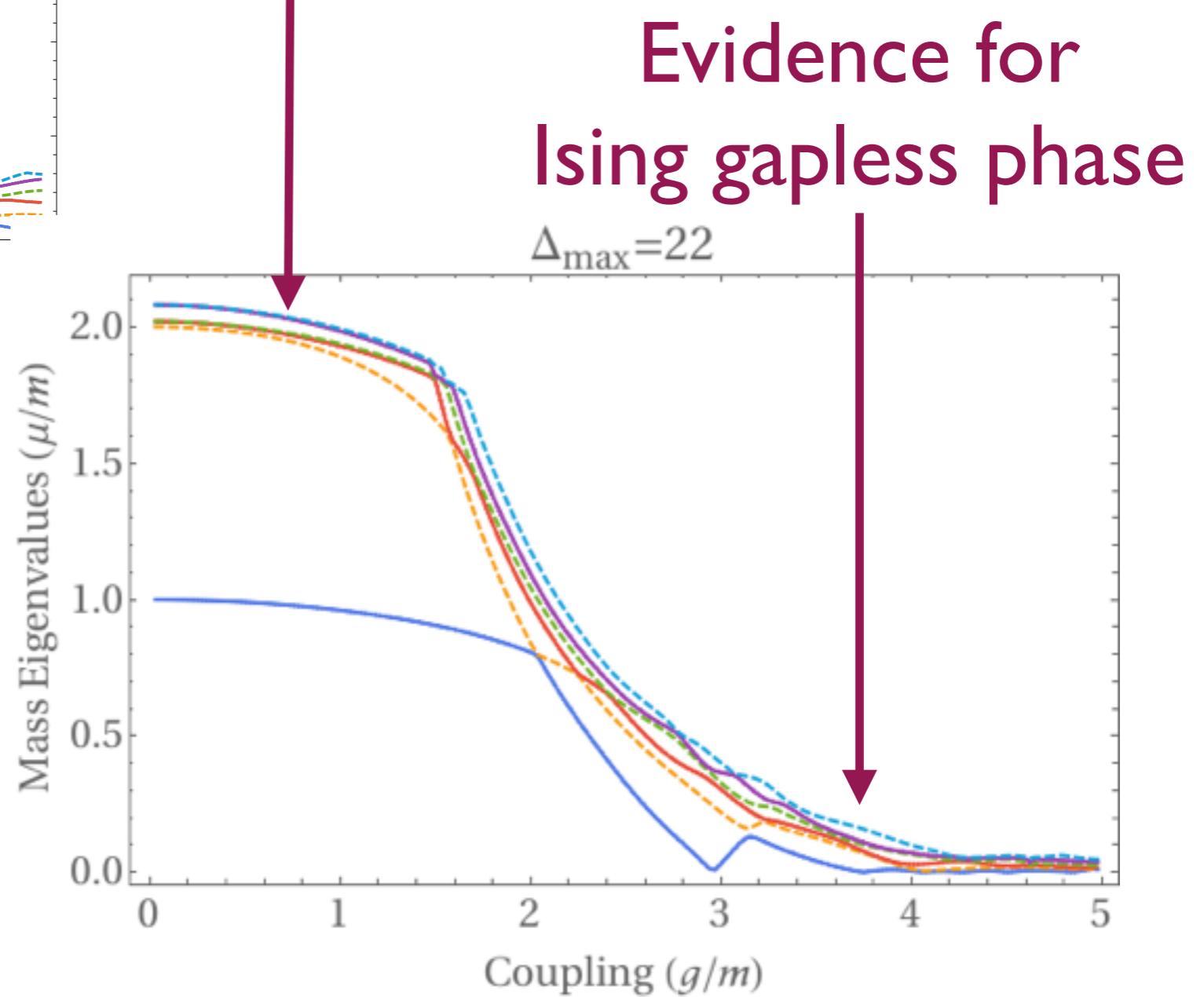
SUSY degeneracy  
recovered



# The Spectrum



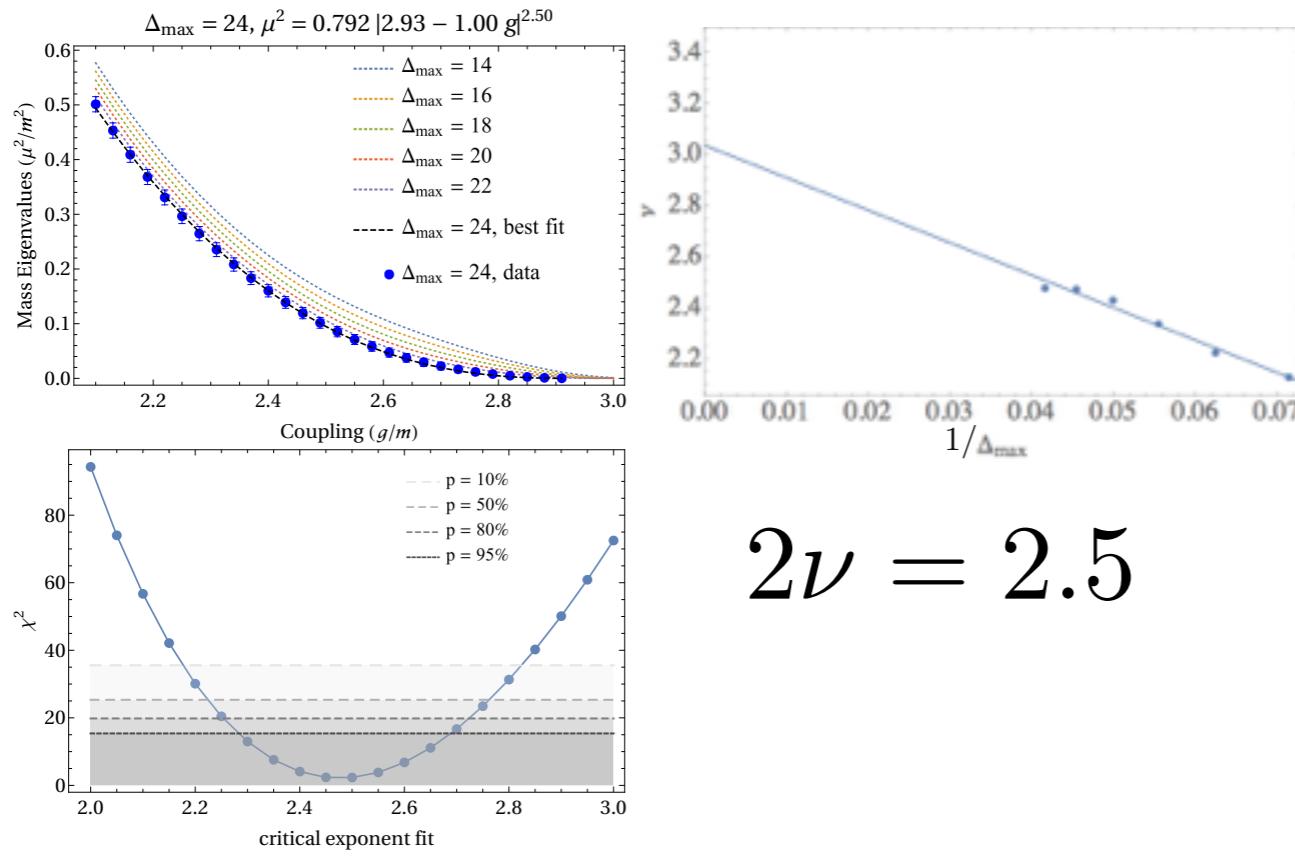
SUSY degeneracy  
recovered



Evidence for  
Ising gapless phase

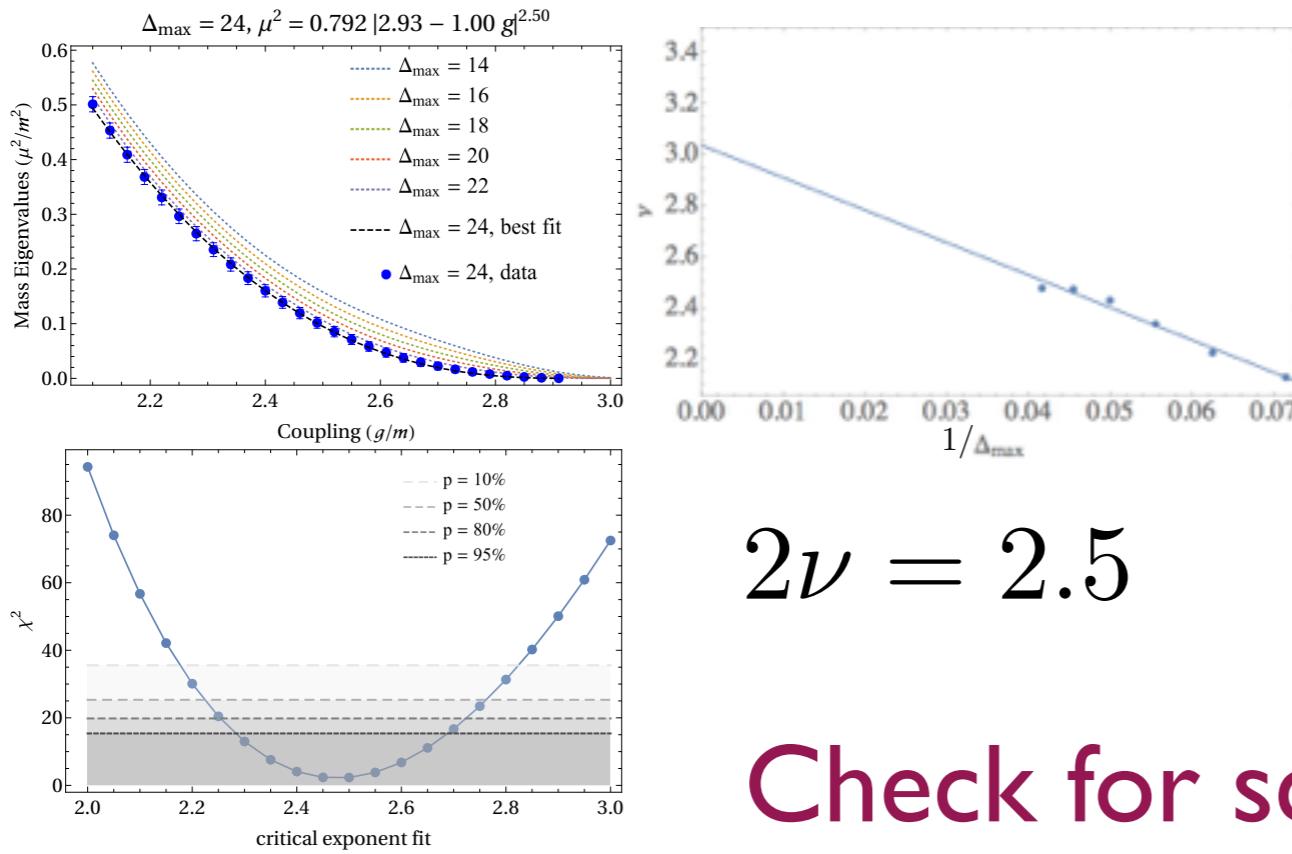
# The Critical Exponent

# The Critical Exponent



$$2\nu = 2.5$$

# The Critical Exponent

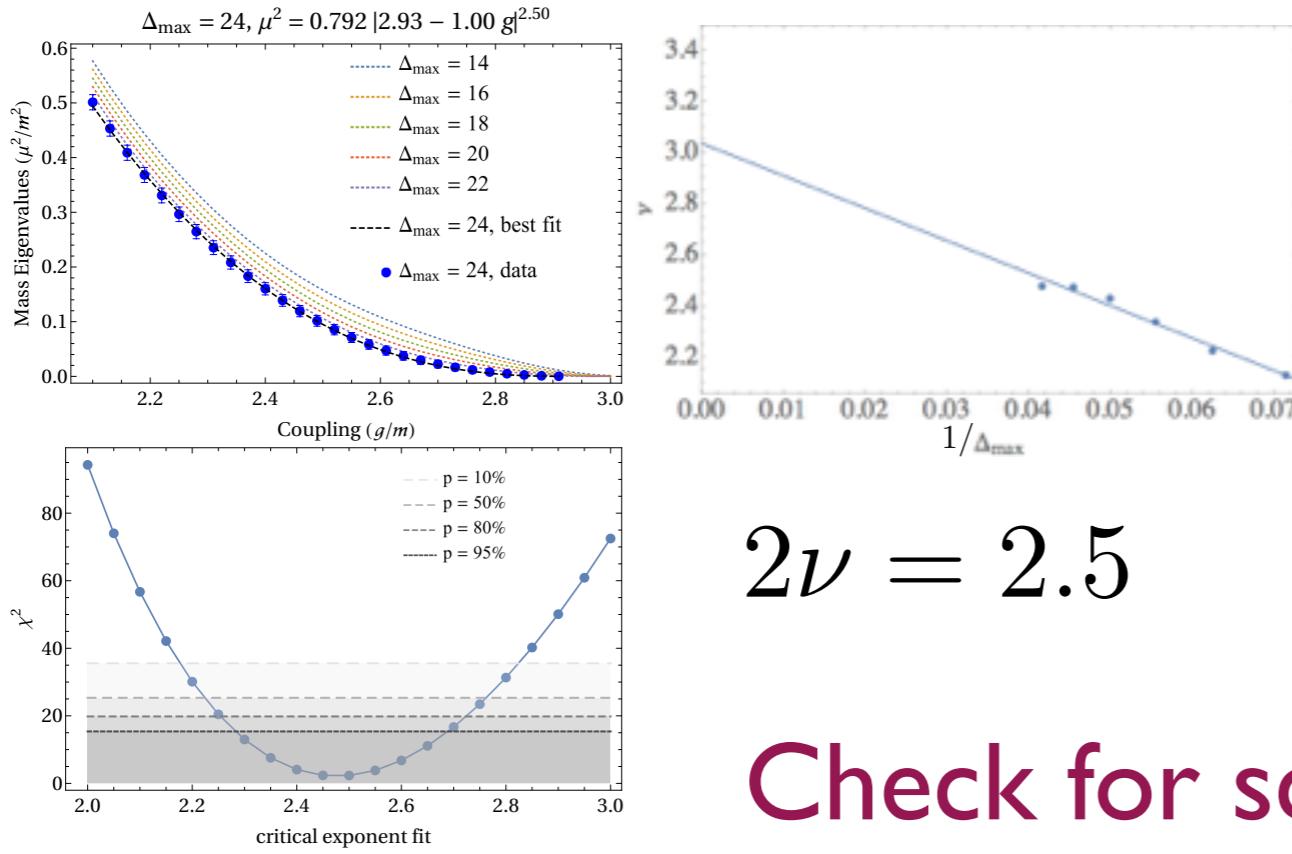


$$2\nu = 2.5$$

**Check for scaling collapse:**

$$m^2 = (g_* - g)^{2\nu} F_{IR} [(g_* - g)^\nu \Delta_{max}^\alpha]$$

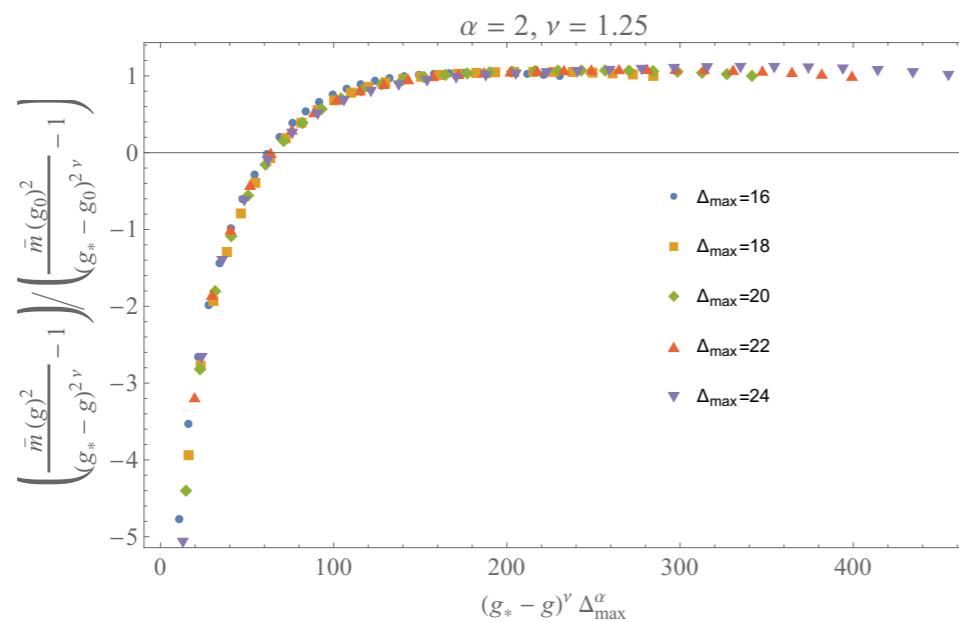
# The Critical Exponent



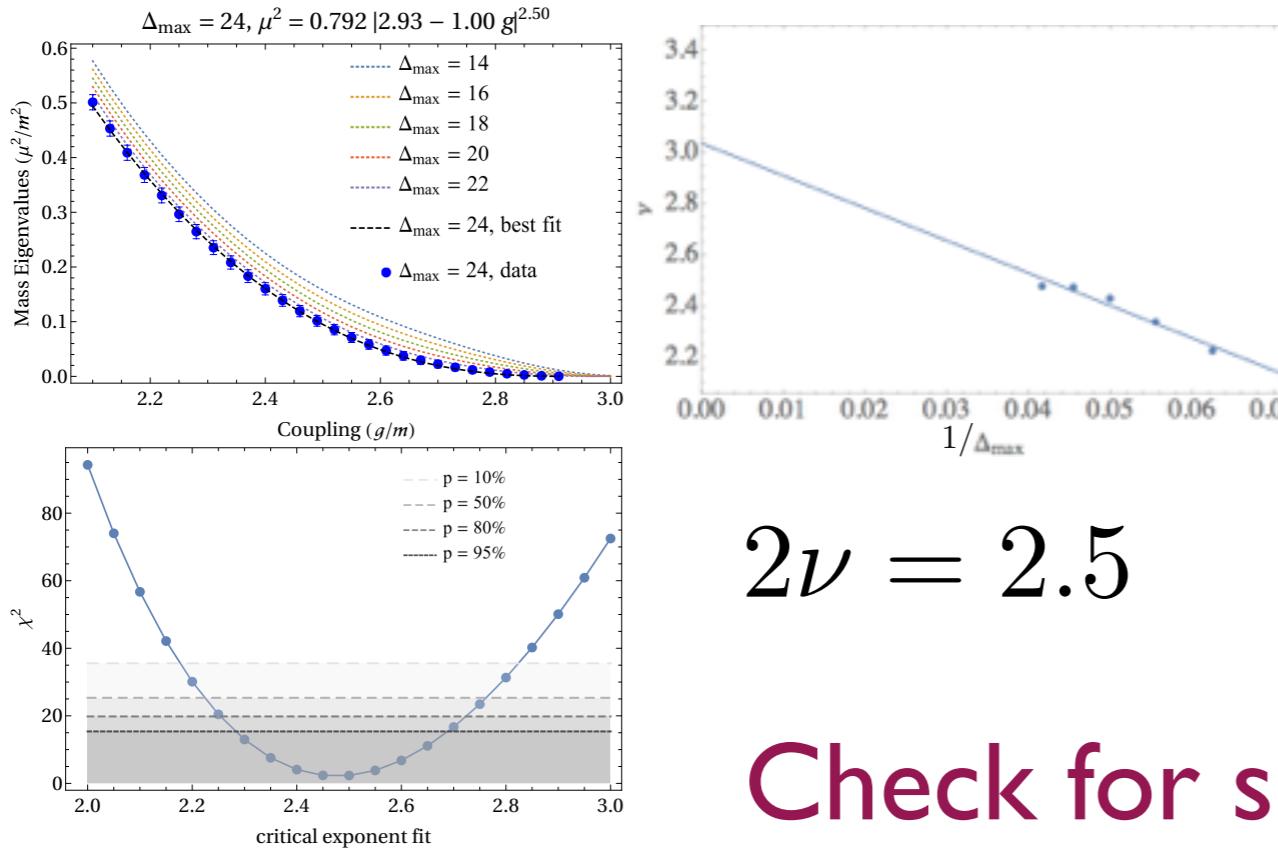
$$2\nu = 2.5$$

**Check for scaling collapse:**

$$m^2 = (g_* - g)^{2\nu} F_{IR} [(g_* - g)^\nu \Delta_{\max}^\alpha]$$



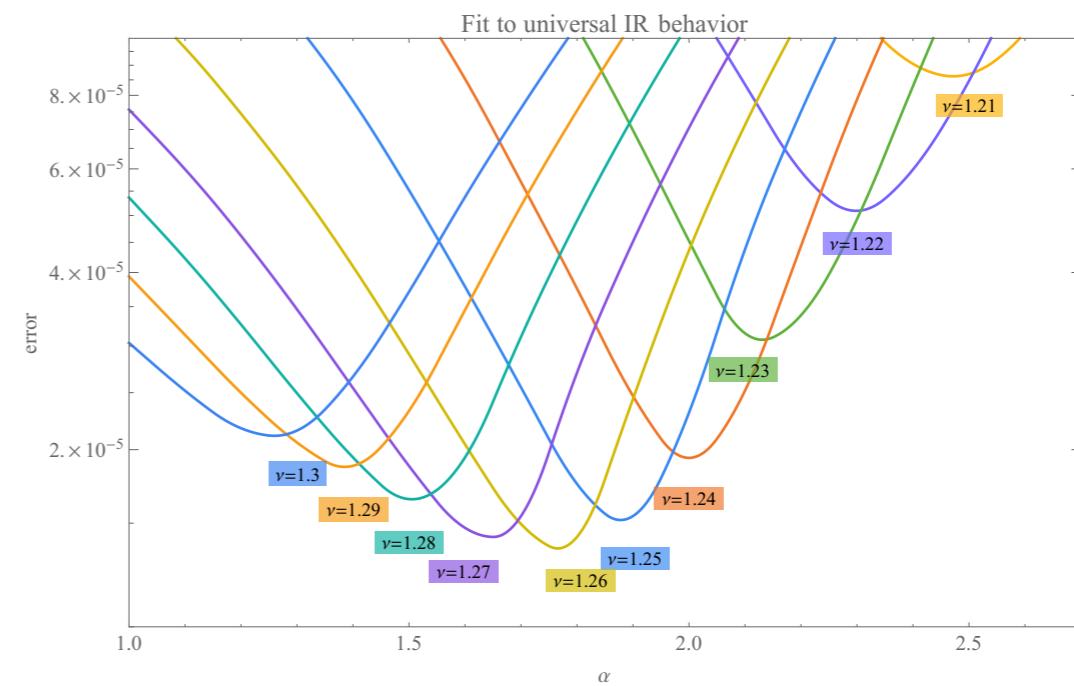
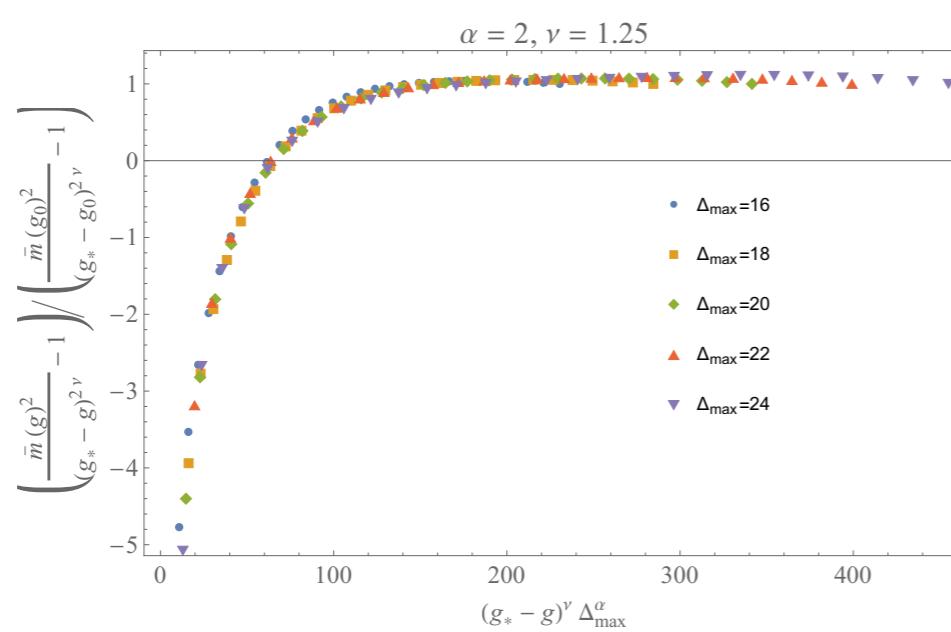
# The Critical Exponent



$$2\nu = 2.5$$

**Check for scaling collapse:**

$$m^2 = (g_* - g)^{2\nu} F_{IR} [(g_* - g)^\nu \Delta_{\max}^\alpha]$$



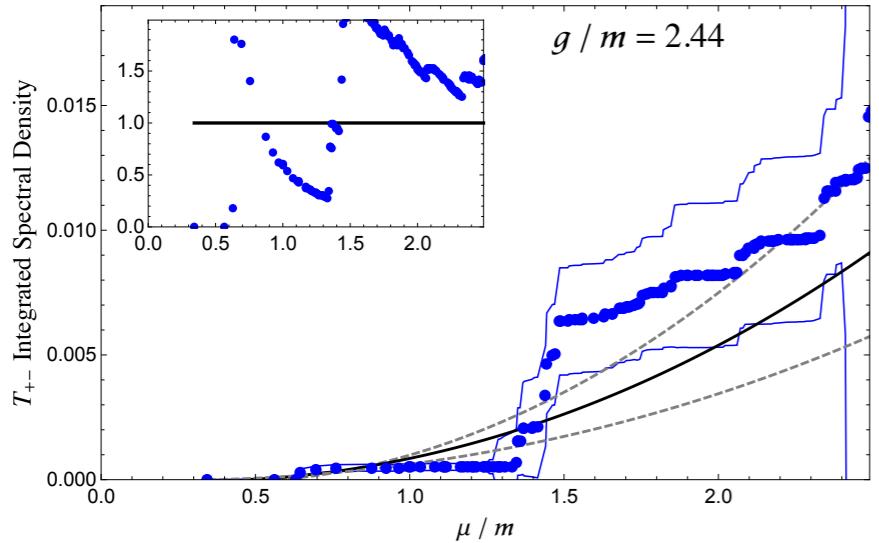
The spectral density of:  $T_\mu^\mu = 2T_{+-}$

The spectral density of:  $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{\frac{4}{5}} \epsilon' + \dots$

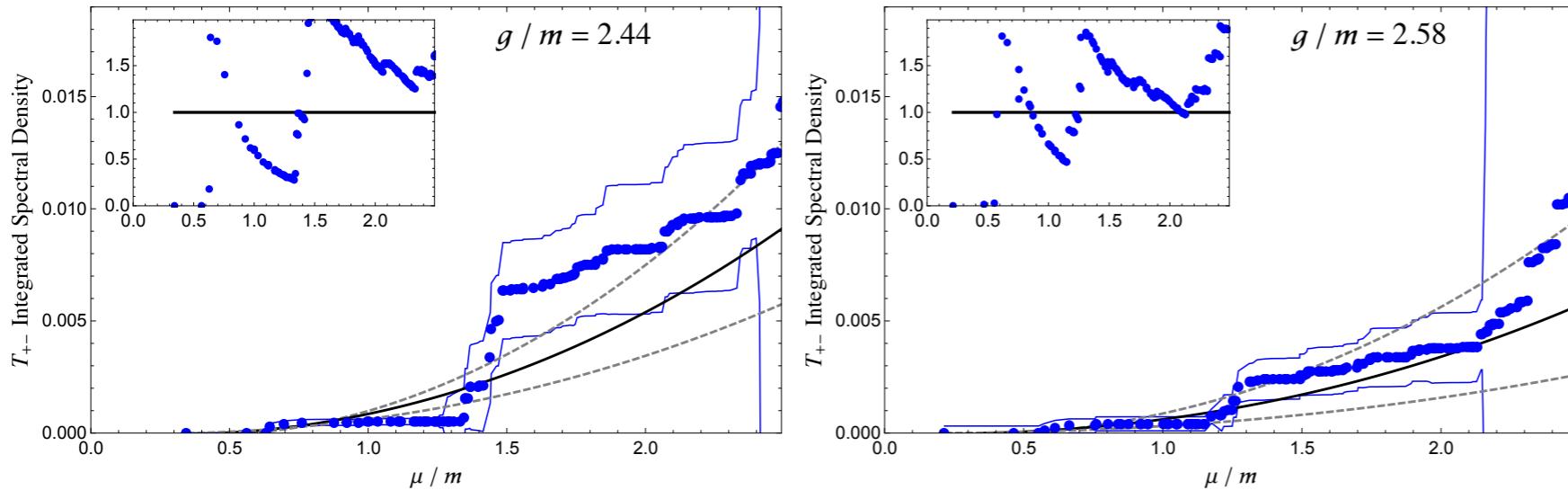
# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$



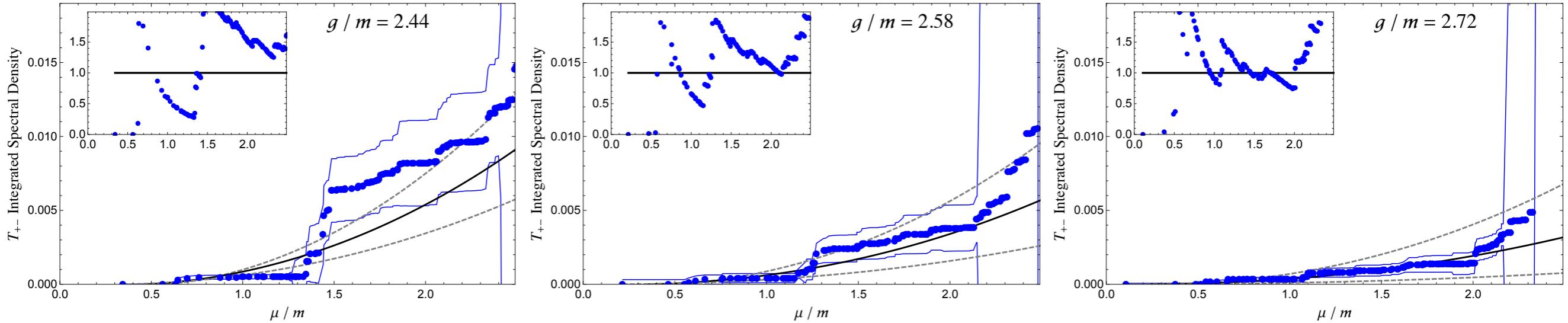
# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$



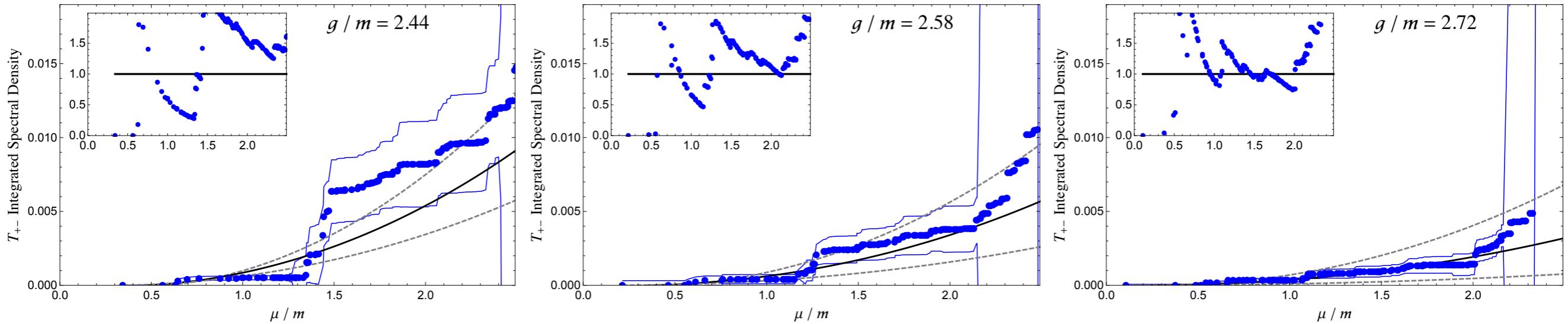
# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$



# The spectral density of: $T_\mu^\mu = 2T_{+-}$

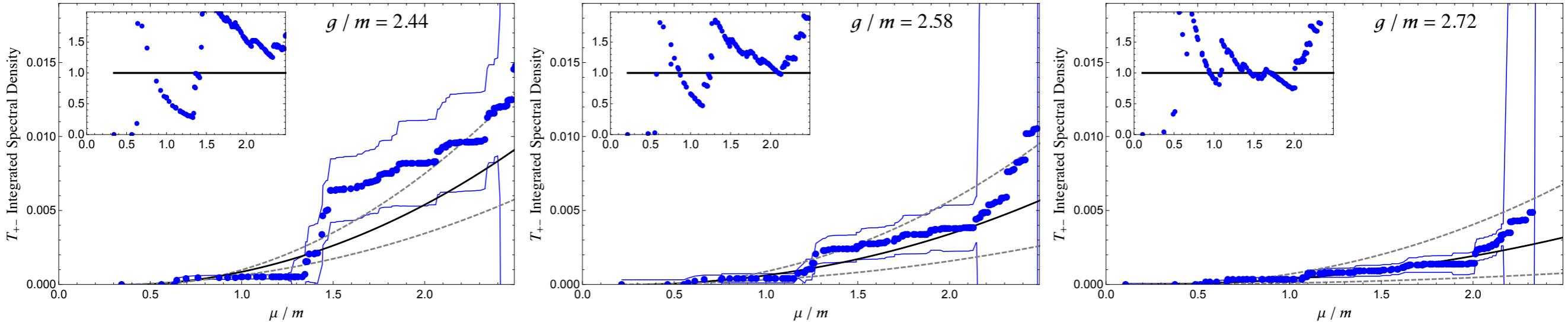
Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$



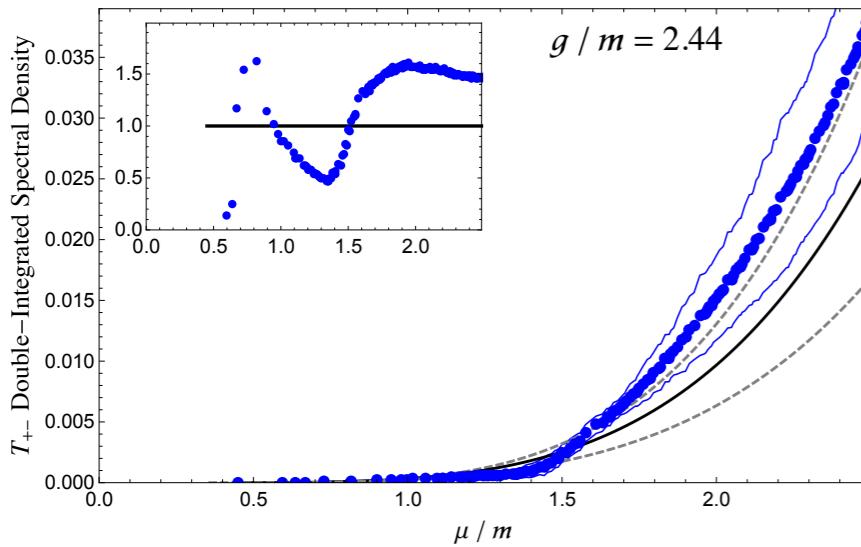
For a smoother taste, let's integrate again:

# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$

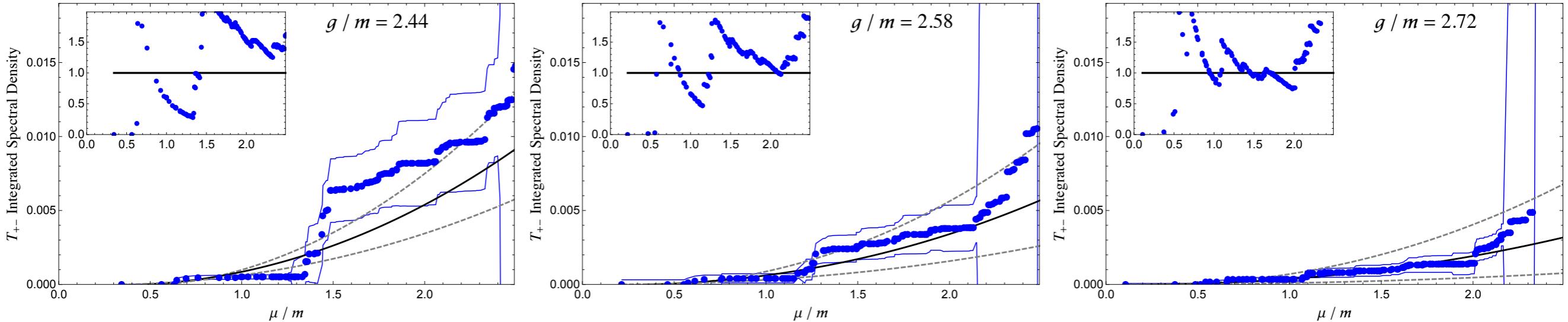


For a smoother taste, let's integrate again:

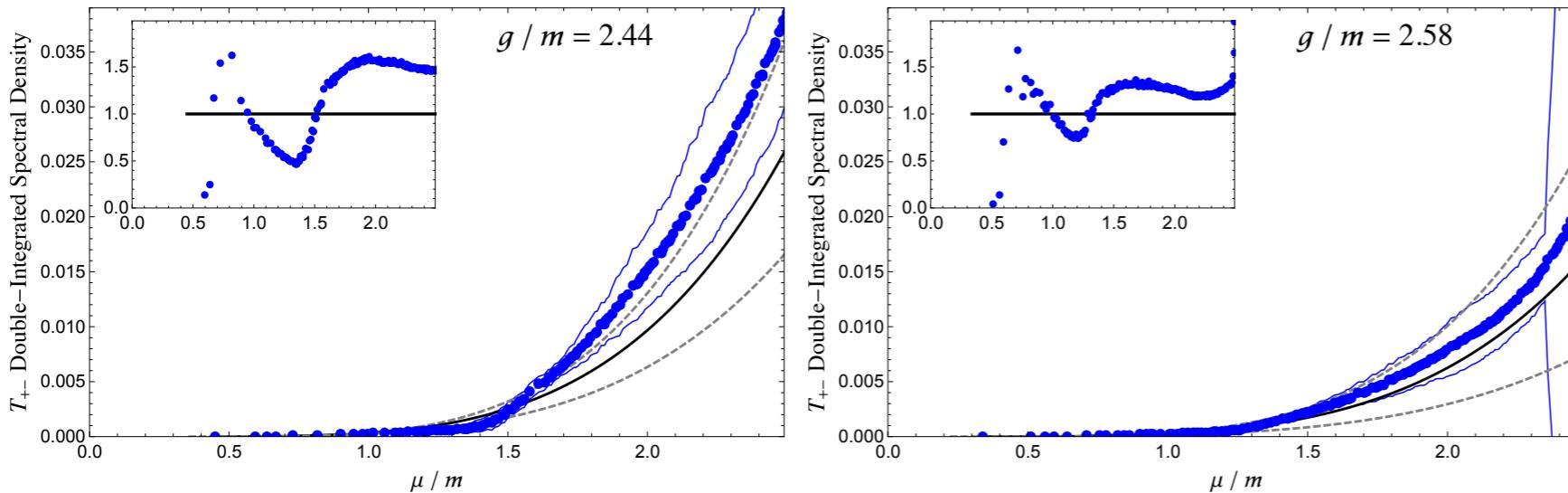


# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$

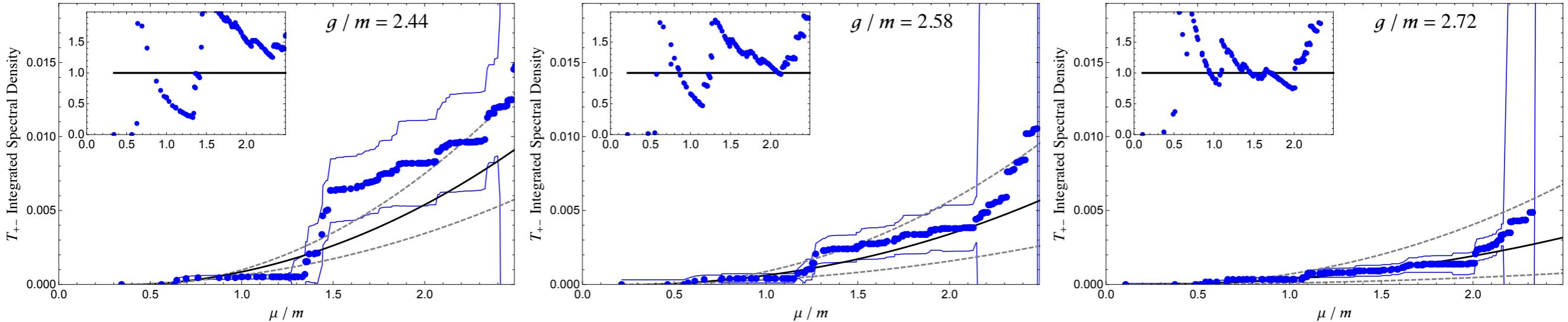


For a smoother taste, let's integrate again:

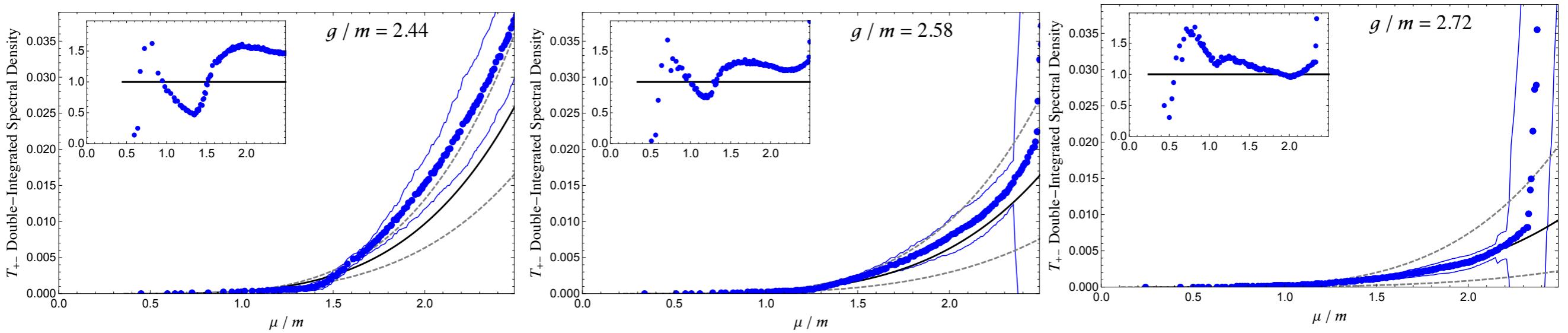


# The spectral density of: $T_\mu^\mu = 2T_{+-}$

Known TIM behavior in the IR:  $T_{+-} = m_{kink}^{5/4} \epsilon' + \dots$



For a smoother taste, let's integrate again:



# A Dream

# A Dream

Large-N RG-flow:

$$\mathcal{N} = 4 \text{ } SYM \rightarrow \mathcal{N} = 2 \text{ } SYM$$

# A Dream

Large-N RG-flow:

$$\mathcal{N} = 4 \text{ } SYM \rightarrow \mathcal{N} = 2 \text{ } SYM$$

$$\delta Q_+ = \int d^{d-1}x \ m \ tr(\phi\psi)$$

# A Dream

## Large-N RG-flow:

$$\mathcal{N} = 4 \text{ } SYM \rightarrow \mathcal{N} = 2 \text{ } SYM$$

$$\delta Q_+ = \int d^{d-1}x \ m \ tr(\phi\psi)$$

SUSY Non-Renorm Thms protect  
naive LC from “zero-modes”!