

Brownian dynamics simulations with hard-body interactions: Exact numerical treatment

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PHYSICAL REVIEW E **83**, 065701(R) (2011)

Hard-wall interactions in soft matter systems: Exact numerical treatment

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(Received 29 March 2011; published 20 June 2011)

THE JOURNAL OF CHEMICAL PHYSICS **137**, 164108 (2012)

**Brownian dynamics simulations with hard-body interactions:
Spherical particles**

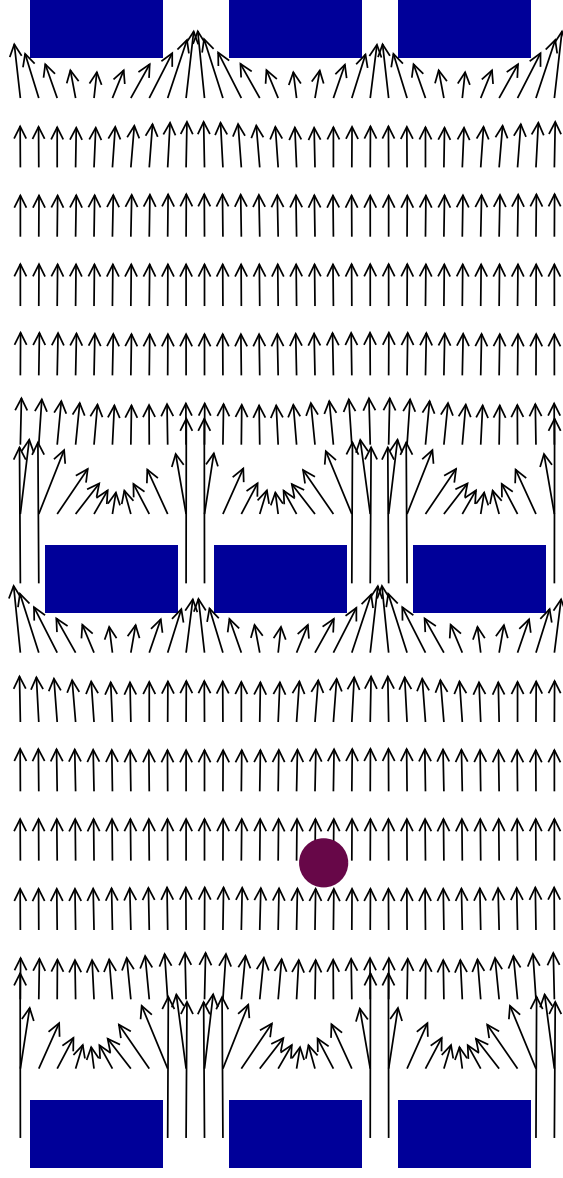
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Brownian motion

- interacting particles in a suspension (e.g. colloids)
- driven by external forces
- in a structured environment
- **solvent**: viscous friction and thermal fluctuations



microfluidics, biomolecules in the cell, self-assembly, polymers, ...

→ **numerical simulation** of particle interaction with **hard walls**

Langevin equation

$$\dot{\vec{r}}(t) = \frac{1}{\eta} \vec{F}(\vec{r}(t), t) + \sqrt{2D} \vec{\xi}(t)$$

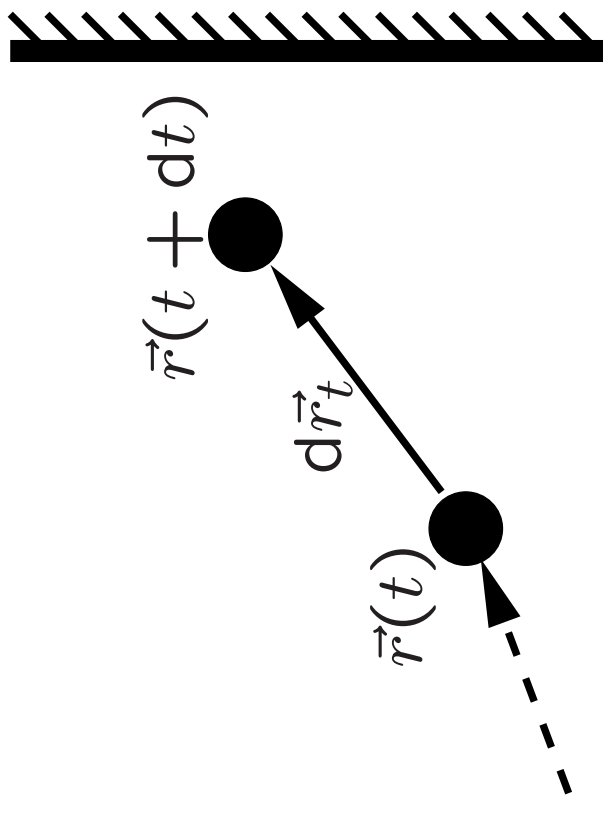
two common idealizations/approximations:

- overdamped limit (“ $m = 0$ ”)
- **hard-body interactions (singular!)**
to represent the extremely short-ranged and strong repulsive contact forces
not included in $\vec{F}(\vec{r}(t), t)$

Euler algorithm

$$d\vec{r}_t = \vec{v}_t dt + \sqrt{2D} dt \vec{G}_t, \quad \vec{r}(t + dt) = \vec{r}(t) + d\vec{r}_t$$

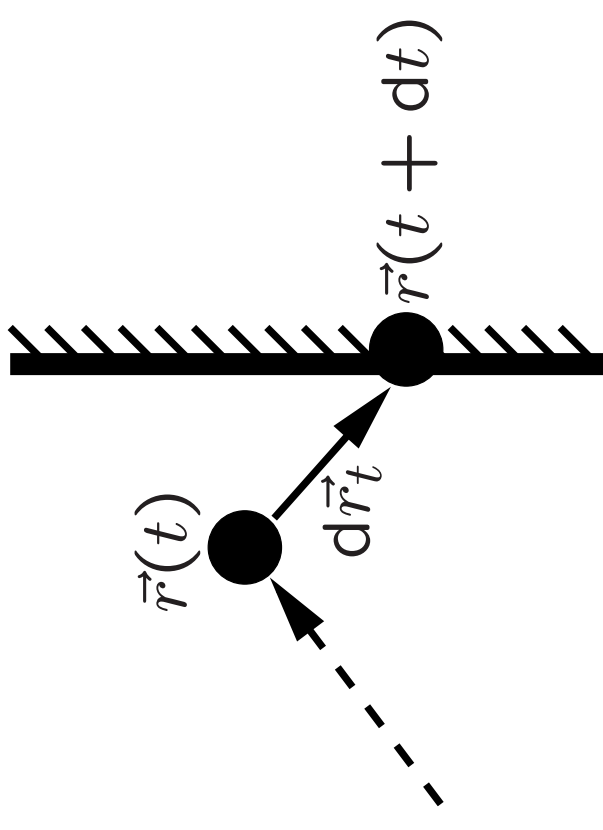
with $\vec{v}_t = \frac{1}{\eta} \vec{F}(\vec{r}(t), t)$, and $\vec{G}_t \in \mathcal{N}(0, 1)^d$



Euler algorithm

$$d\vec{r}_t = \vec{v}_t dt + \sqrt{2D} d\vec{t} \vec{G}_t, \quad \vec{r}(t+dt) = \vec{r}(t) + d\vec{r}_t$$

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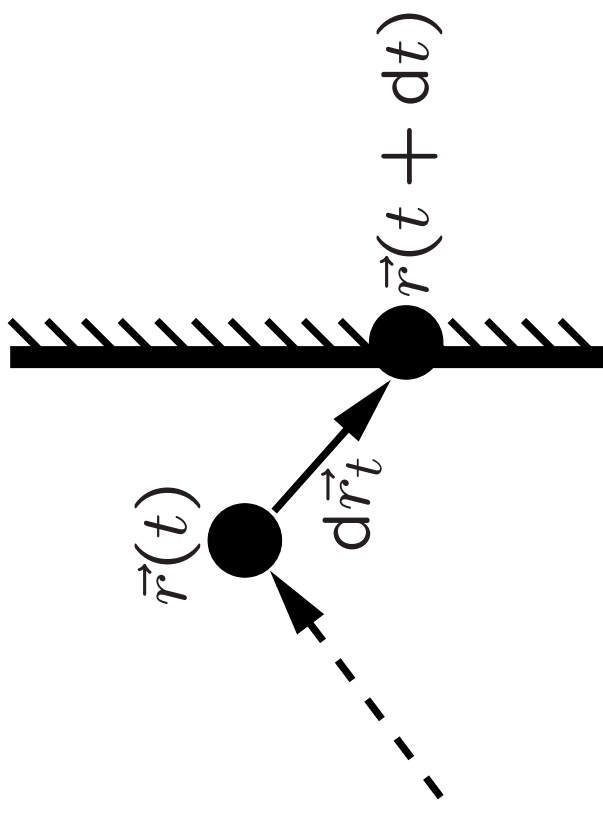
Euler algorithm

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algorithm:

- 1) detect unphysical configurations
(“collisions”)
- 2) **rule to generate physically valid configuration**



Heuristic methods (prominent examples)

- **rejection scheme**
discard unphysical configurations
(advance time or not???)
[B. Cichocki and K. Hinzen, *Physica A* **166**, 473 (1990)]
- **event-driven scheme**
propagate fraction of time step ϵdt until “collision point”
use rejection scheme for remaining time step $(1 - \epsilon)dt$
[Y.-G. Tao et al., *J. Chem. Phys.* **124**, 134906 (2006)]

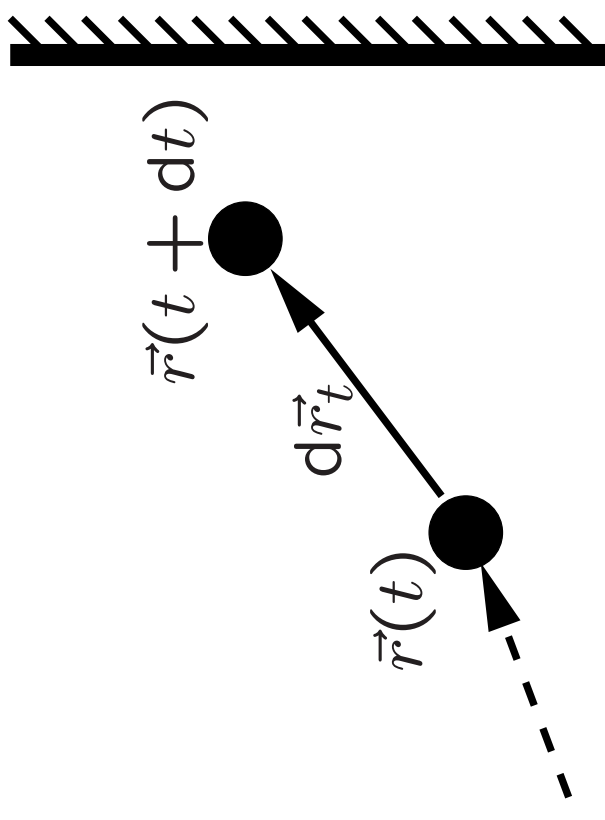
lack thorough justification

Euler algorithm

$$d\vec{r}_t = \vec{v}_t dt + \sqrt{2D} d\vec{t} \vec{G}_t, \quad \vec{r}(t + dt) = \vec{r}(t) + d\vec{r}_t$$

with $\vec{v}_t = \frac{1}{\eta} \vec{F}(\vec{r}(t), t)$, and $\vec{G}_t \in \mathcal{N}(0, 1)^d$

$$p(d\vec{r}_t) = \frac{1}{\sqrt{4\pi D dt}^d} \exp\left(-\frac{[d\vec{r}_t - \vec{v}_t dt]^2}{4D dt}\right)$$



Euler algorithm

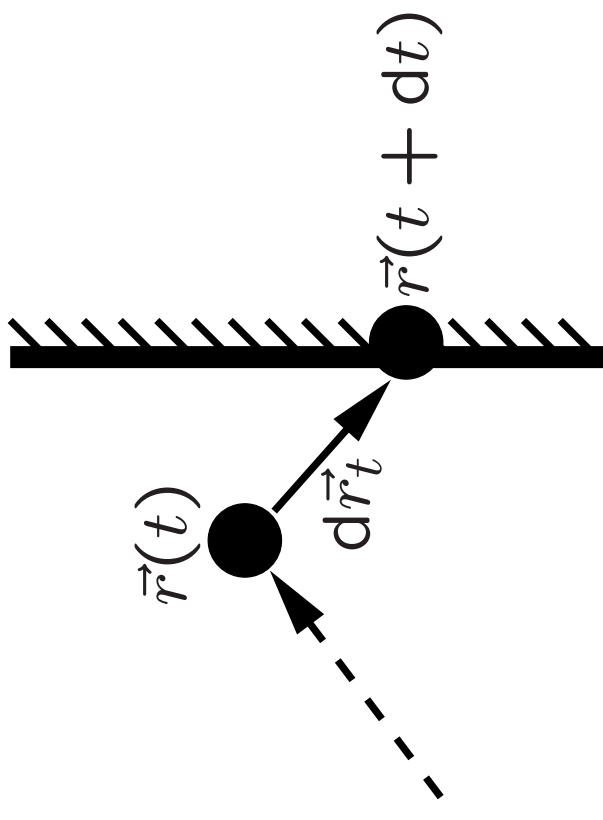
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$$p(d\vec{r}_t) = ???$$



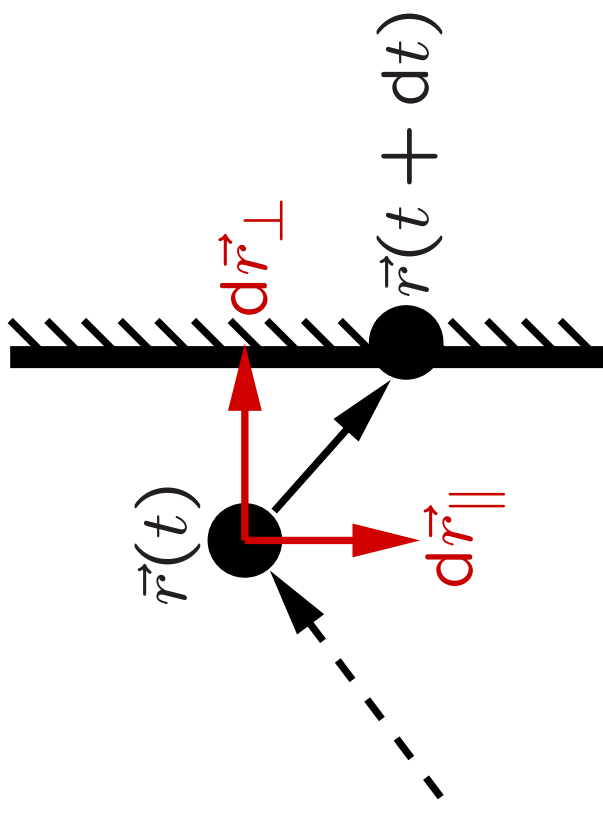
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algorithm:

- 1) detect unphysical configurations
("collisions")
- 2) **rule to generate physically valid configuration**
 $p(d\vec{r}_t) = ???$



Smoluchowski solution

[M. V. Smoluchowski, Phys. Z. **17**, 557 (1916)]

$$\frac{\partial}{\partial t} p = D_q \frac{\partial^2}{\partial q^2} p - v_q \frac{\partial}{\partial q} p$$
$$- \left[D_q \frac{\partial}{\partial q} p - v_q p \right]_{q=0} = 0$$

driven diffusion on a half-line $q \in [0, \infty)$

reflecting boundary at $q = 0$

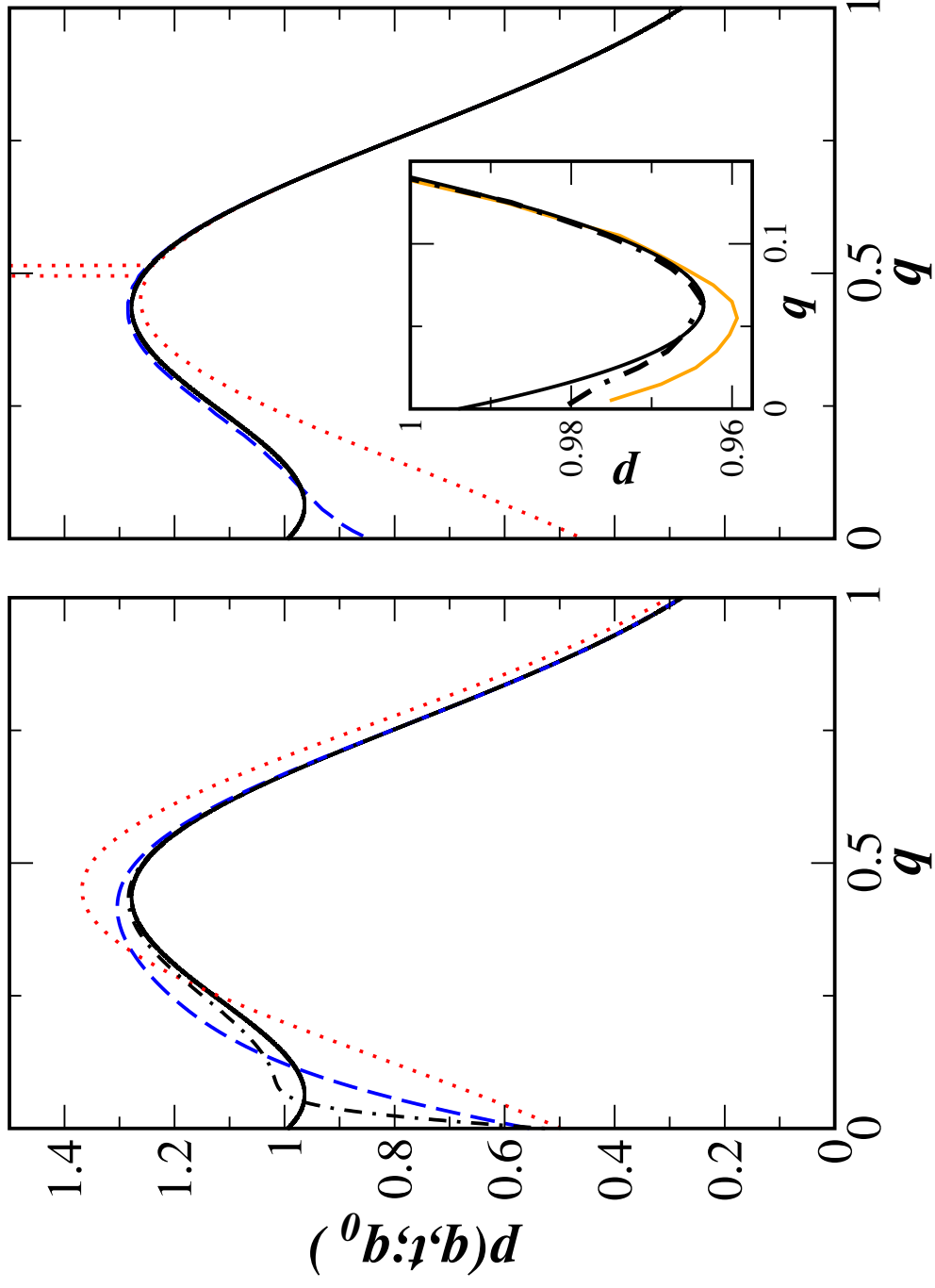
initial position

q_0 at time $t = 0$

solution: $p(q, t; q_0) = p_1(q, t; q_0) + p_2(q, t; q_0) + p_3(q, t; q_0)$

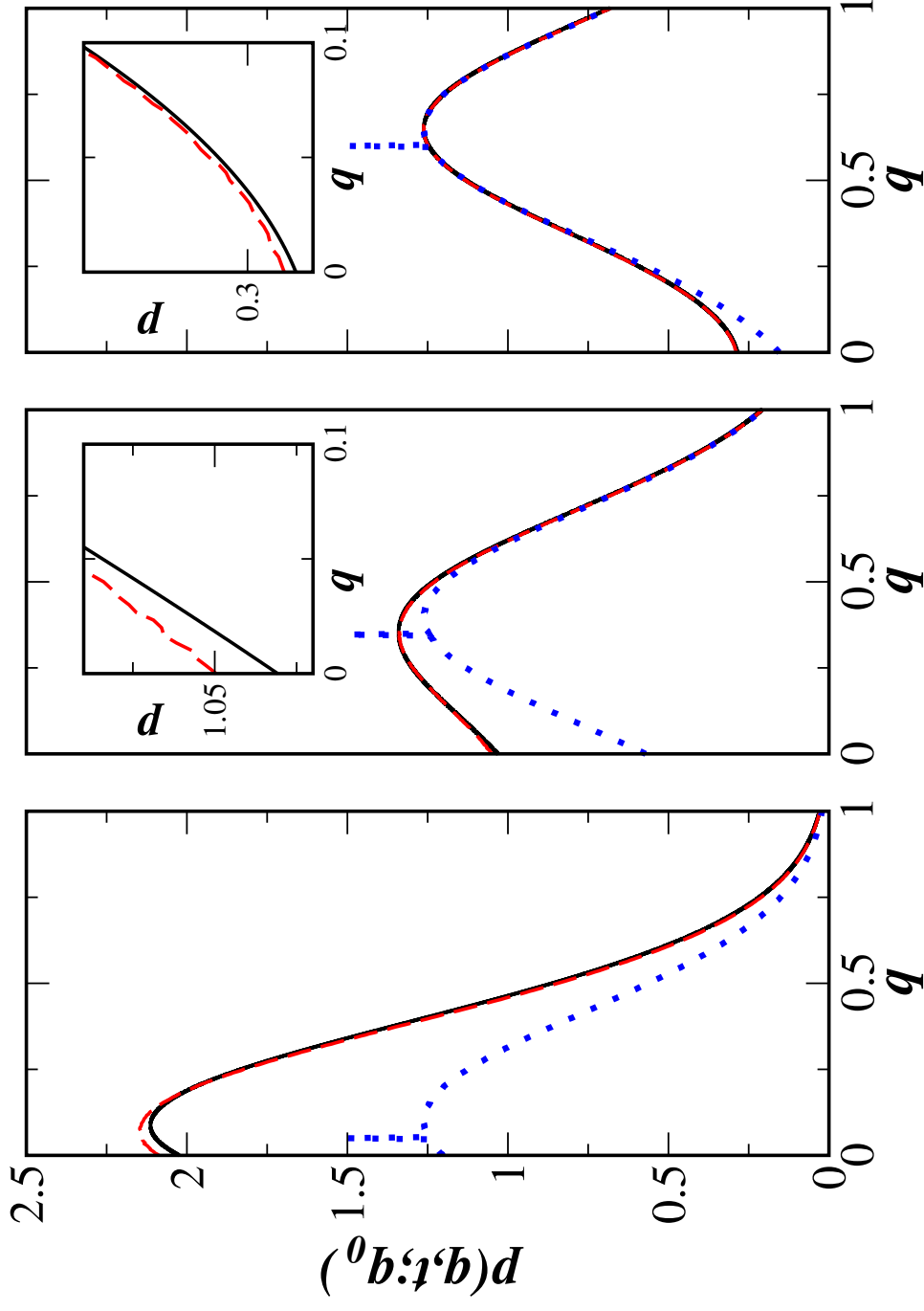
$$\text{with } p_1(q, t; q_0) = \frac{1}{\sqrt{4\pi D_q t}} \exp\left(-\frac{(q - q_0 - v_q t)^2}{4D_q t}\right)$$
$$p_2(q, t; q_0) = \frac{\exp\left(-\frac{v_q q_0}{D_q}\right)}{\sqrt{4\pi D_q t}} \exp\left(-\frac{(q + q_0 - v_q t)^2}{4D_q t}\right)$$
$$p_3(q, t; q_0) = -\frac{v_q}{2D_q} \exp\left(\frac{v_q q}{D_q}\right) \operatorname{erfc}\left(\frac{q + q_0 + v_q t}{\sqrt{4D_q t}}\right)$$

Smoluchowski solution



$$v_q = -1.0, D_q = 1.0, q_0 = 0.5, t = 0.05$$

Smoluchowski solution



$q_0 = 0.05$ $q_0 = 0.35$ $q_0 = 0.6$

$v_q = 1.0, D_q = 1.0, t = 0.05$

Smoluchowski solution

[M. V. Smoluchowski, Phys. Z. **17**, 557 (1916)]

$$\frac{\partial}{\partial t} p = D_q \frac{\partial^2}{\partial q^2} p - v_q \frac{\partial}{\partial q} p$$
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driven diffusion on a half-line $q \in [0, \infty)$

reflecting boundary at $q = 0$

initial position

q_0 at time $t = 0$

solution: $p(q, t; q_0) = p_1(q, t; q_0) + p_2(q, t; q_0) + p_3(q, t; q_0)$

$$\text{with } p_1(q, t; q_0) = \frac{1}{\sqrt{4\pi D_q t}} \exp\left(-\frac{(q - q_0 - v_q t)^2}{4D_q t}\right)$$
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$$p_3(q, t; q_0) = -\frac{v_q}{2D_q} \exp\left(\frac{v_q q}{D_q}\right) \operatorname{erfc}\left(\frac{q + q_0 + v_q t}{\sqrt{4D_q t}}\right)$$

The algorithm

1. Perform standard integration as long as the displacements $d\vec{r}_t$ do not lead to unphysical configurations (“collisions”)
2. If a suggested displacement $d\vec{r}_t$ results in a “collision”, replace its component along the “collision axis” by a new $dq = q - q_0$ drawn from
$$\frac{p_2(q, dt; q_0) + p_3(q, dt; q_0)}{w}$$

with

$$\begin{aligned} w &= \int_0^\infty dq [p_2(q, dt; q_0) + p_3(q, dt; q_0)] \\ &= 1 - \int_0^\infty dq p_1(q, dt; q_0) \\ &= \int_{-\infty}^0 dq p_1(q, dt; q_0) \end{aligned}$$

The algorithm

“collision axis” $\hat{=}$ **half-line** (with “collision point” at the origin)

random number created on the “collision axis” :

$$Q = Q_G \Theta(Q_G) + Q_C \Theta(-Q_G)$$

$$\text{with } p_G(q) = p_1(q) \quad (q \in \mathbb{R}) \quad \text{and} \quad p_C(q) = \frac{p_2(q) + p_3(q)}{w} \quad (q \in \mathbb{R}_{>0})$$

$$\begin{aligned} \Rightarrow p_{\text{alg}}(q) &= \int_{-\infty}^{\infty} dq_1 \int_0^{\infty} dq_2 p_G(q_1) p_C(q_2) \delta(q - [q_1 \Theta(q_1) + q_2 \Theta(-q_1)]) \\ &= \int_{-\infty}^{\infty} dq_1 \int_0^{\infty} dq_2 p_G(q_1) p_C(q_2) [\Theta(q_1) \delta(q - q_1) \\ &\quad + \Theta(-q_1) \delta(q - q_2)] \\ &= \Theta(q) p_G(q) + p_C(q) \int_{-\infty}^0 dq_1 p_G(q_1) \\ &= p_1(q) + p_2(q) + p_3(q) \end{aligned}$$

Euler algorithm for hard wall

$$d\vec{r}_t = \vec{v}_t dt + \sqrt{2D} dt \vec{G}_t, \quad \vec{r}(t + dt) = \vec{r}(t) + d\vec{r}_t^*$$

with $\vec{v}_t = \frac{1}{\eta} \vec{F}(\vec{r}(t), t)$, and $\vec{G}_t \in \mathcal{N}(0, 1)^d$

algorithm:

1) “suggest” $d\vec{r}_t$ using standard Euler scheme

2a) no “collision”:

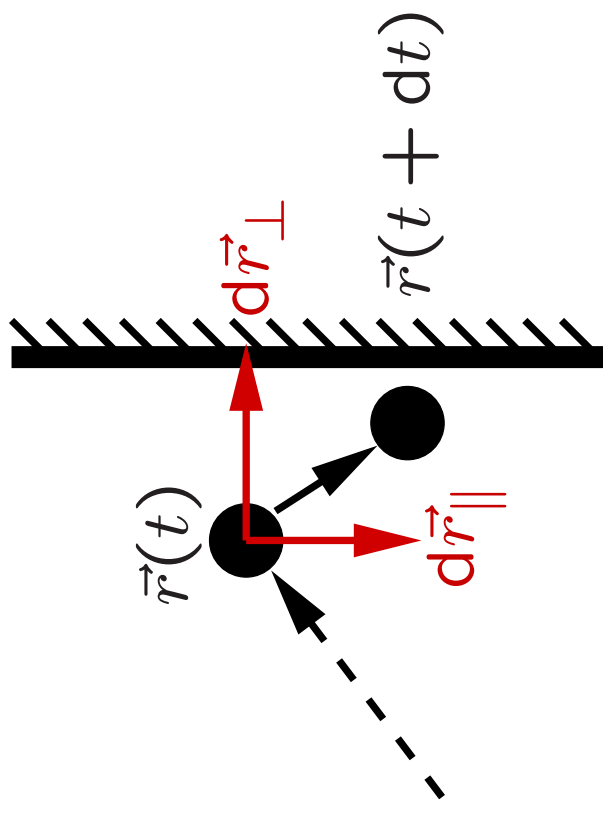
$$d\vec{r}_t^* = d\vec{r}_t$$

2b) “collision”:

$$d\vec{r}_t^* = d\vec{r}_t + (q - q_0 - \vec{n} \cdot d\vec{r}_t) \vec{n}$$

$$\text{with } \vec{n} = -d\vec{r}_\perp / |d\vec{r}_\perp|$$

note: $d\vec{r}_\perp$ and $d\vec{r}_\parallel$ are uncorrelated



Euler algorithm for spherical particles

$$\boxed{d\vec{r}_1 = \vec{v}_1 dt + \sqrt{2D_1} dt \vec{G}_1, \quad d\vec{r}_2 = \vec{v}_2 dt + \sqrt{2D_2} dt \vec{G}_2}$$

$$\vec{r}_1(t + dt) = \vec{r}_1(t) + d\vec{r}_1^*, \quad \vec{r}_2(t + dt) = \vec{r}_2(t) + d\vec{r}_2^*$$

algorithm:

1) “suggest” $d\vec{r}_1$, $d\vec{r}_2$ using standard Euler scheme

2a) no “collision”:

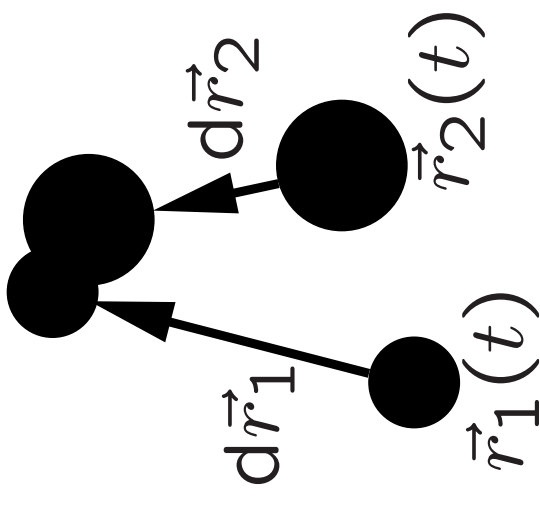
$$d\vec{r}_1^* = d\vec{r}_1 \quad \text{and} \quad d\vec{r}_2^* = d\vec{r}_2$$

2b) “collision”:

$$d\vec{r}_1^* = d\vec{r}_1 + \frac{\eta_2}{\eta_1 + \eta_2} [(d\vec{r}_2 - d\vec{r}_1) \cdot \vec{e} - (q - q_0)] \vec{e}$$

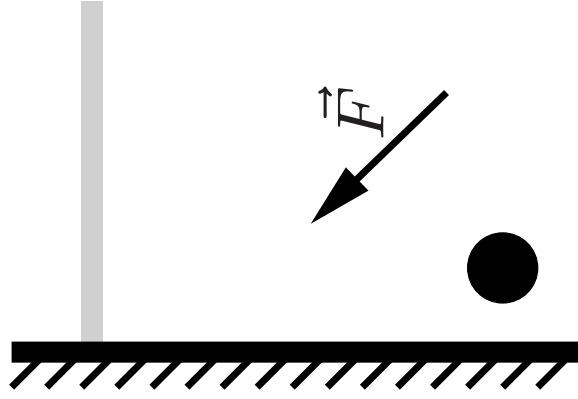
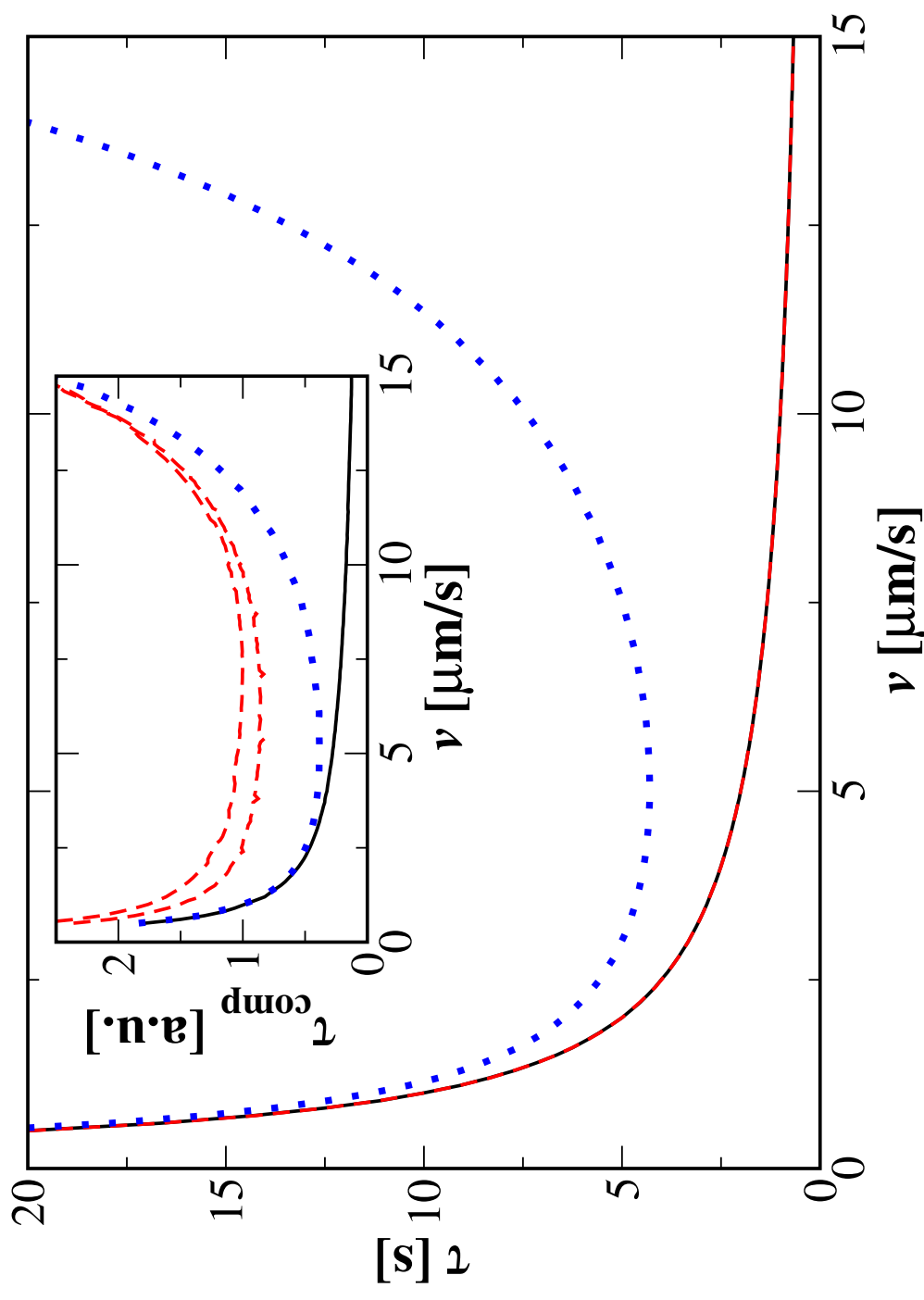
$$d\vec{r}_2^* = d\vec{r}_2 - \frac{\eta_1}{\eta_1 + \eta_2} [(d\vec{r}_2 - d\vec{r}_1) \cdot \vec{e} - (q - q_0)] \vec{e}$$

$$\text{with } \vec{e} = (\vec{r}_2(t) - \vec{r}_1(t)) / |\vec{r}_2(t) - \vec{r}_1(t)|$$



note: center of friction and relative motion are uncorrelated

Example: Mean first passage time



$$\vec{F} = (-f, f), \quad v = f/\eta, \quad \text{particle radius } 1 \mu\text{m}, \quad dt = 0.01 \text{ s}$$

In practice

- detection of “collisions”
(any integration scheme)
- integration time step dt is determined by
variations of \vec{F} and **curvature of structures/particles**
- generation of random number q according to distribution
 $p_C(q) = [p_2(q) + p_3(q)]/w$

$$F(q) = \int_0^q dq' p_C(q') = \frac{\operatorname{erfc}\left(\frac{q_0 + v_q dt}{\sqrt{4D_q dt}}\right) - \exp\left(\frac{v_q q}{D_q}\right) \operatorname{erfc}\left(\frac{q + q_0 + v_q dt}{\sqrt{4D_q dt}}\right)}{\operatorname{erfc}\left(\frac{q_0 + v_q dt}{\sqrt{4D_q dt}}\right)}.$$

Then: $q = F^{-1}(x)$ with x uniformly distributed on $[0, 1]$
is distributed according to $p_C(q)$

numerical solution (Brent’s scheme from the GNU scientific library)

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extensions:

non-spherical particles
corners & wedges
many particles
(crowding)

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