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Trieste, 12 June 2020



Hydrodynamics of inhomogeneous locally integrable models

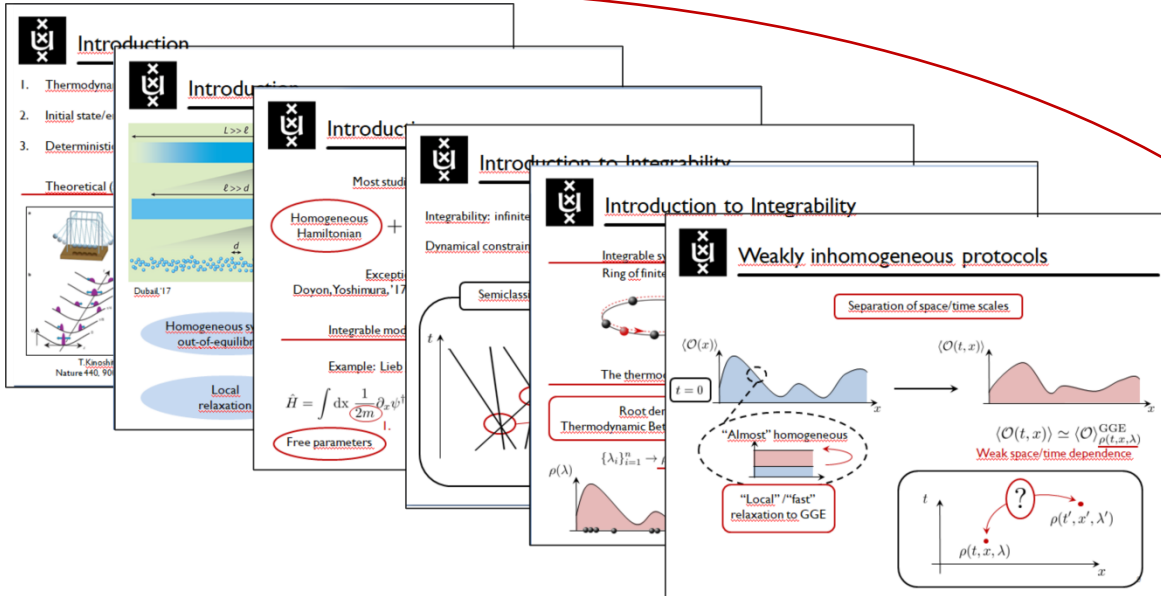
Based on:

AB, A. De Luca, PRL 122 (24), 240606 (2019)

AB, V.Alba, J.-S. Caux, PRL 123 (13) 130602 (2019)



Introduction



Skip \approx 20 min worth talk

Pros of a thematic workshop:
no need of a general introduction

Topics:

- 1 Generalized Hydrodynamics (GHD) can be extended to describe inhomogeneous smooth Hamiltonians
- 2 Inhomogeneities that look smooth sometimes are not: beyond GHD effects and bound-state recombination

From previous talks:
traps (but not only)



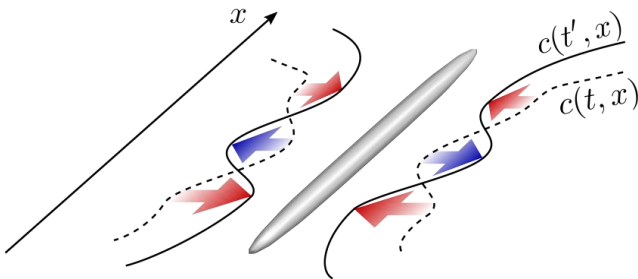
Inhomogeneous interactions

A concrete example: the Lieb-Liniger model

Integrable for arbitrary
(homogeneous) couplings

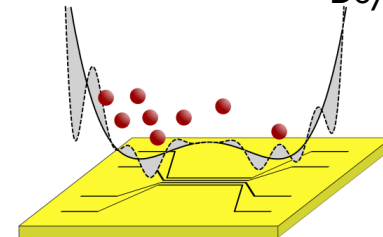
$$\hat{H} = \int dx \partial_x \psi^\dagger \partial_x \psi + c \psi^\dagger \psi^\dagger \psi \psi - \mu \psi^\dagger \psi$$

Inhomogeneous
transverse trap



Inhomogeneous
longitudinal trap

Doyon, Yoshimura '17



Experimentally confirmed!

M. Schemmer, I. Bouchoule, B. Doyon, J. Dubail '19

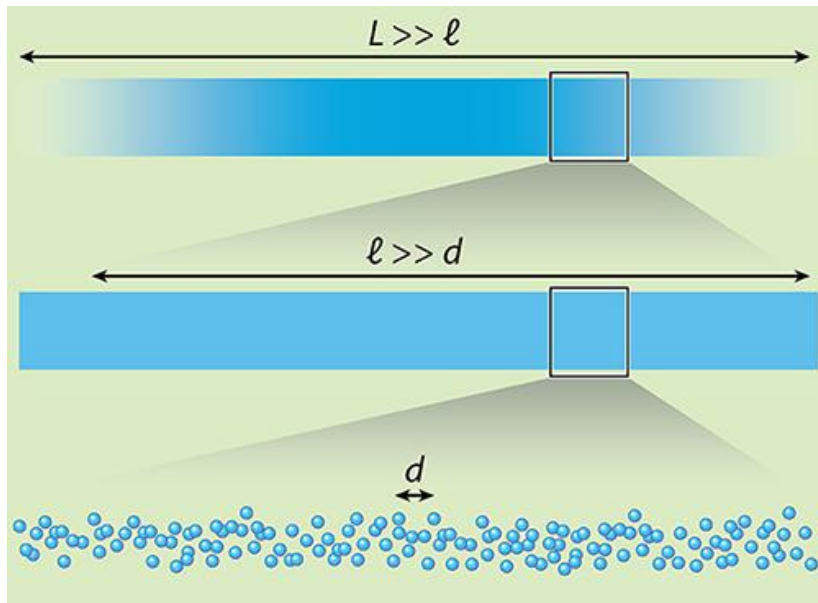
Solution with the Generalized Hydrodynamics (GHD)



Inhomogeneous interactions

GHD in a nutshell

Castro Alvaredo, Doyon, Yoshimura, '16
Bertini, Collura, De Nardis, Fagotti, '16



Dubail, '17

Smooth inhomogeneity



Local relaxation to the
GGE of the local
integrable model

+

Gluing together the
smoothly varying
GGEs



GHD

$$\partial_t \vartheta + v^{\text{eff}} \partial_x \vartheta + F^{\text{eff}} \partial_\lambda \vartheta = 0$$



Inhomogeneous interactions

... + spatial inhomogeneities = ...

$c \rightarrow \alpha$ “generic” coupling

GHD with arbitrary inhomogeneities

AB, V.Alba, J.-S. Caux, PRL 123 (13) 130602 (2019)

$$\partial_t \vartheta + v^{\text{eff}} \partial_x \vartheta + \frac{\partial_t \alpha f^{\text{dr}} + \partial_x \alpha \Lambda^{\text{dr}}}{(\partial_\lambda p)^{\text{dr}}} \partial_\lambda \vartheta = 0$$

“single particle”
effect

“collective”
effect

$$\Theta(\lambda) = -i \log S(\lambda)$$

$$f(\lambda) = -\partial_\alpha p(\lambda) + \int \frac{d\mu}{2\pi} \partial_\alpha \Theta(\lambda - \mu) (\partial_\mu p)^{\text{dr}} \vartheta(\mu)$$

$$\Lambda(\lambda) = -\partial_\alpha \epsilon(\lambda) + \int \frac{d\mu}{2\pi} \partial_\alpha \Theta(\lambda - \mu) (\partial_\mu \epsilon)^{\text{dr}} \vartheta(\mu)$$



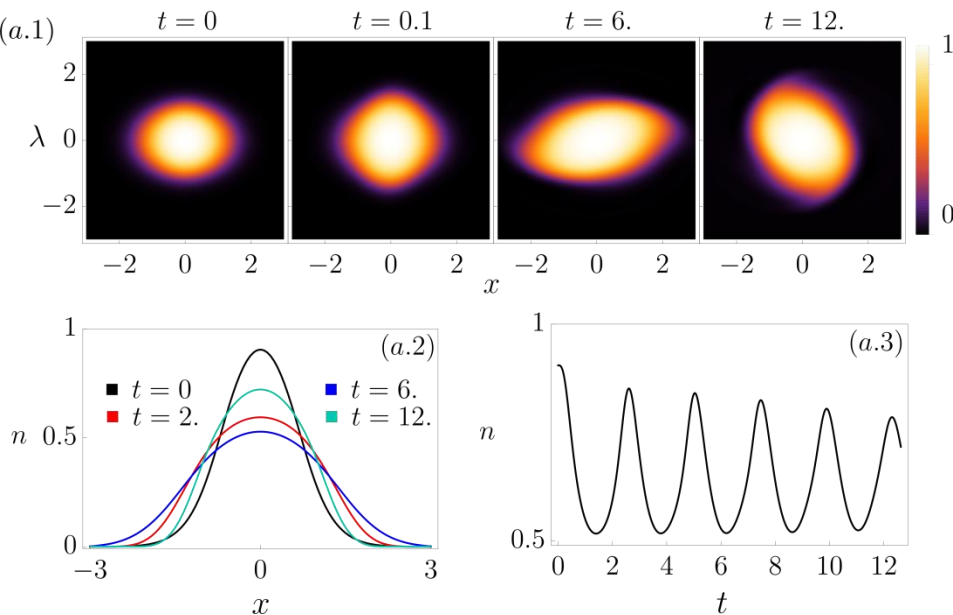
Inhomogeneous interactions

Applications

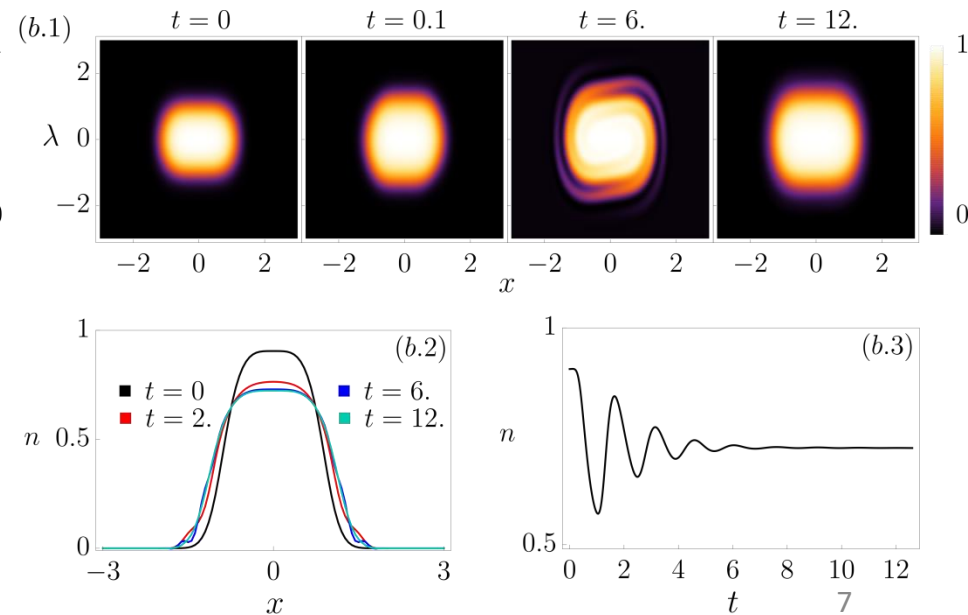
Slow interaction changes in trapped Lieb-Liniger

$$\hat{H} = \int dx \partial_x \psi^\dagger \partial_x \psi + c(t/L) \psi^\dagger \psi^\dagger \psi \psi + V(x/L) \psi^\dagger \psi$$

Harmonic trap



Anharmonic trap





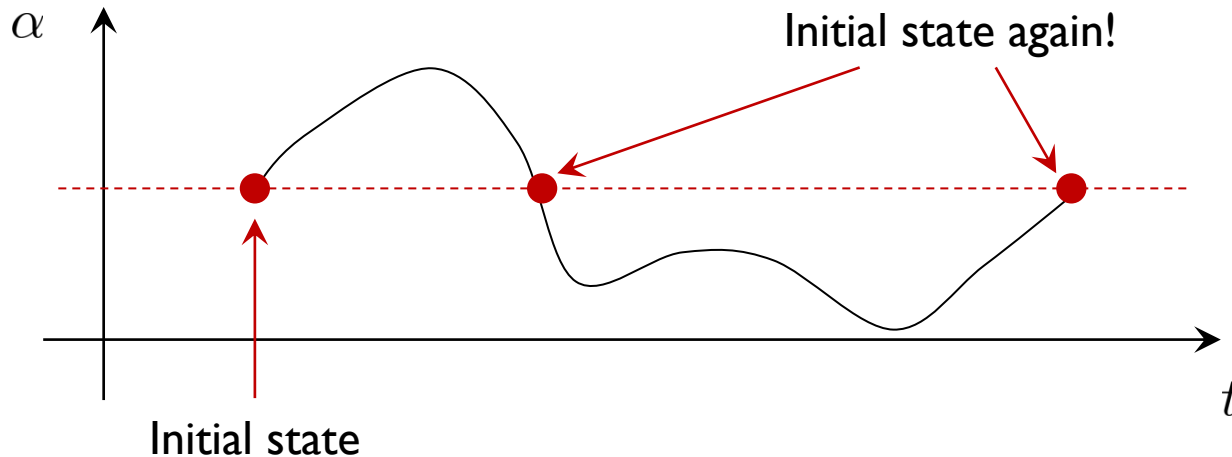
GHD, reversibility and beyond

Homogeneous system, slow coupling changes

No explicit time dependence!

$$\partial_t \vartheta + \frac{\partial_t \alpha f^{\text{dr}}}{(\partial_{\lambda p})^{\text{dr}}} \partial_{\lambda} \vartheta = 0 \quad \xrightarrow[t \rightarrow \alpha]{\text{Change variable}} \quad \partial_{\alpha} \vartheta + \frac{f^{\text{dr}}}{(\partial_{\lambda p})^{\text{dr}}} \partial_{\lambda} \vartheta = 0$$

Reversibility under slow coupling changes



... always?

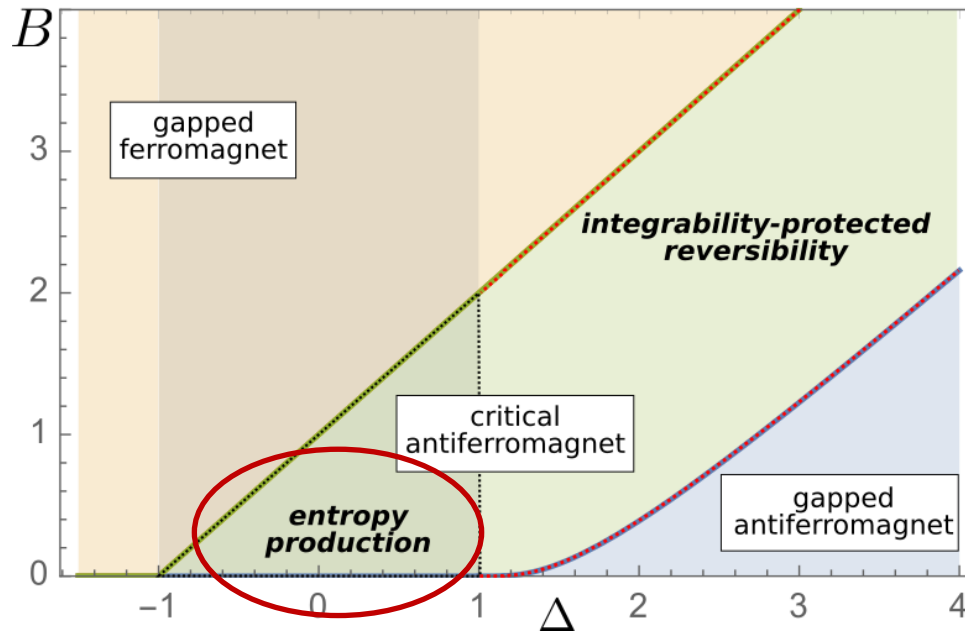
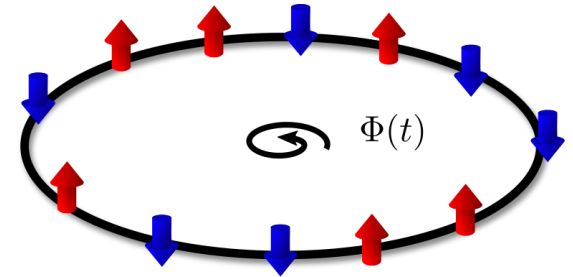


GHD, reversibility and beyond

Homogeneous magnetic flux in XXZ

AB, A. De Luca, PRL 122 (24), 240606 (2019)

$$\hat{H}(\Phi) = \sum_{j=1}^N \frac{1}{2} \left(e^{i\Phi} \hat{s}_j^+ \hat{s}_{j+1}^- + \text{h.c.} \right) + \Delta \hat{s}_j^z \hat{s}_{j+1}^z - B \hat{s}_j^z$$



The XXZ chain is not “smooth” under flux changes for $|\Delta| < 1$



GHD, reversibility and beyond

we can still write the GHD...

$$\partial_{\Phi} \vartheta_j + \frac{m_j^{\text{dr}}(\lambda)}{(\partial_{\lambda} p_j(\lambda))^{\text{dr}}} \partial_{\lambda} \vartheta_j = 0$$

$$|\Delta| \geq 1$$

Infinitely many strings
(bound states)

Brillouin zone

$$\lambda \in \left[-\frac{\pi}{2}, \frac{\pi}{2} \right]$$

$$|\Delta| < 1$$

Number of strings
 Δ -dependent

No Brillouin zone

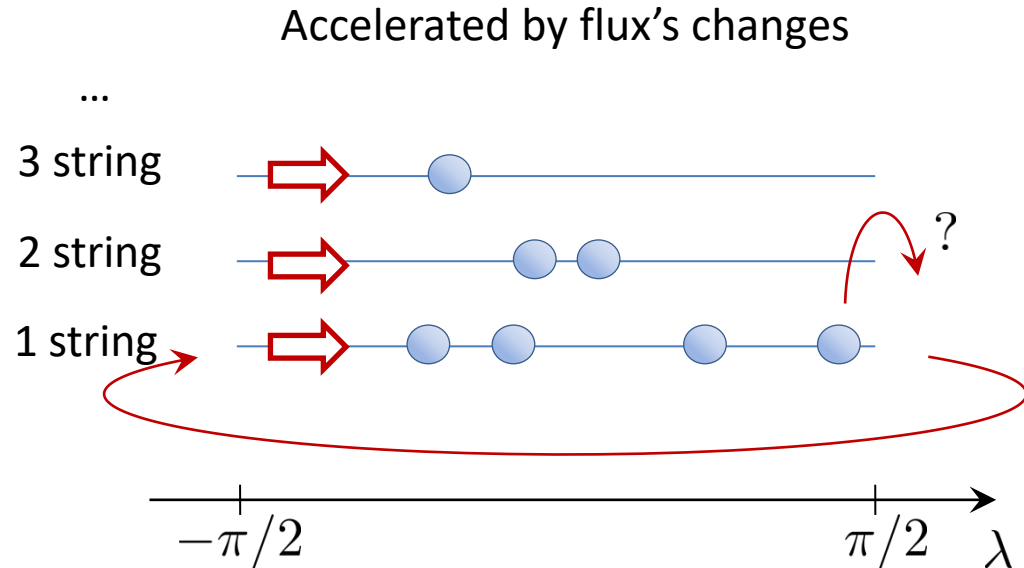
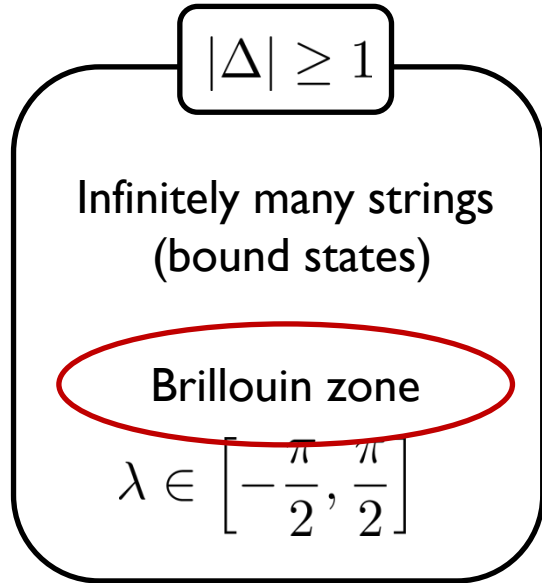
$$\lambda \in \mathbb{R}$$

Where does the entropy production come from?

Boundary conditions in the rapidity space



GHD, reversibility and beyond



Time-reversible
GHD equation

+

Time-reversible
boundary conditions

=

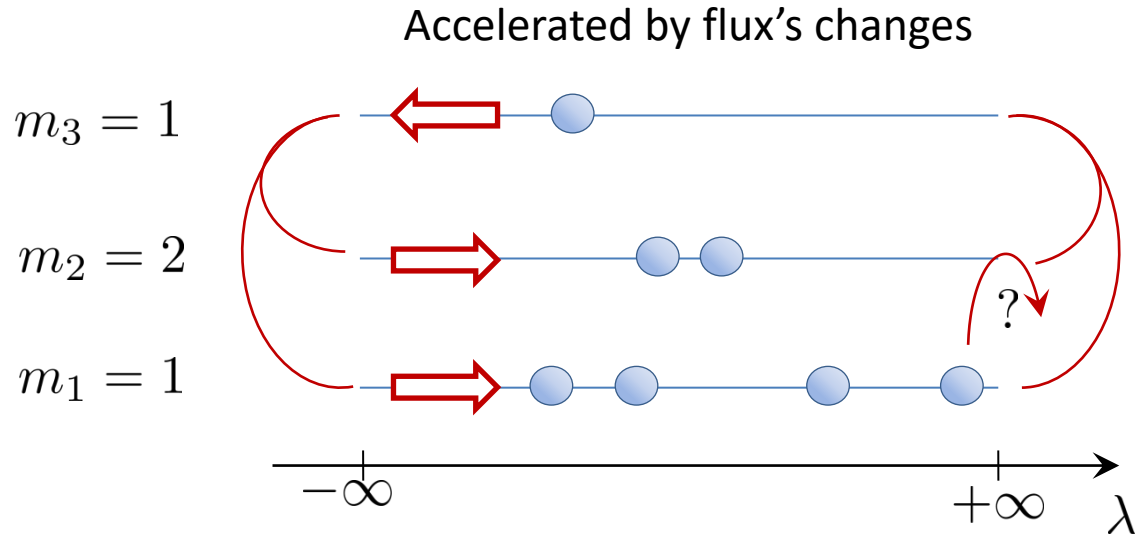
Time-reversible
dynamics



GHD, reversibility and beyond

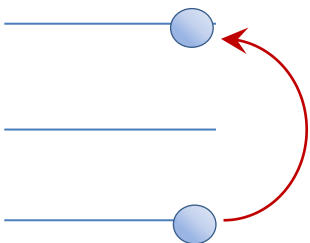
$$\Delta = \frac{1}{2}$$

3 strings
No Brillouin zone
 $\lambda \in \mathbb{R}$

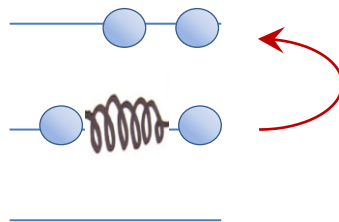


$$\lim_{\lambda \rightarrow \pm\infty} q_j(\lambda) = m_j \lim_{\lambda \rightarrow \pm\infty} q_1(\lambda) \quad \text{Strings indistinguishable at the boundaries}$$

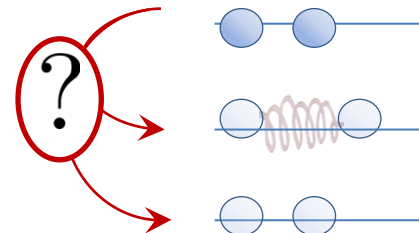
“Trivial” process



“Breaking” of a bound state

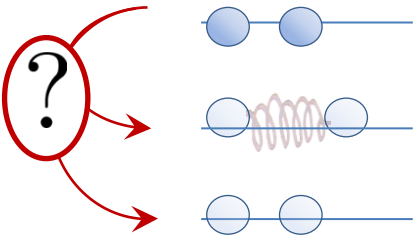


Possible “formation” of a bound state





GHD, reversibility and beyond



$$\text{GGEs} = \text{Charges conservation} + \text{Entropy maximization}$$

✓

Entropy rate maximization fixes recombination rate

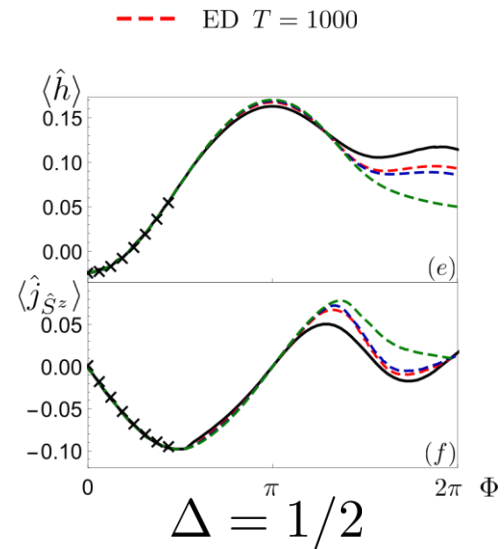
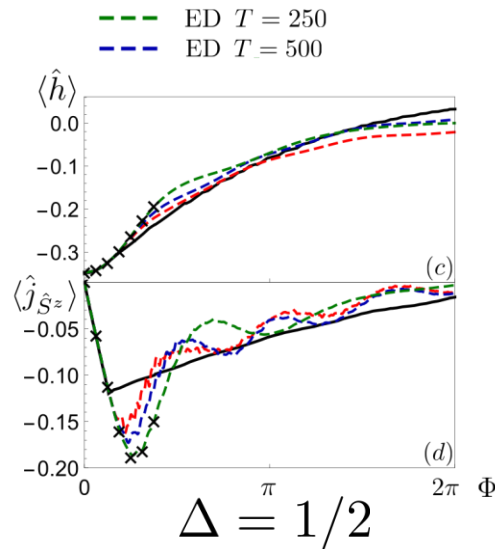
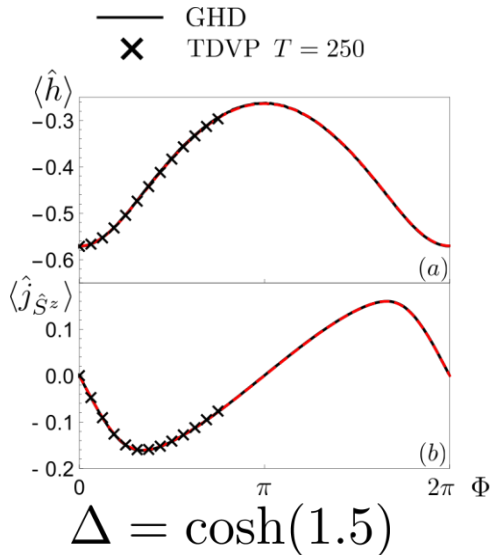
$$\mathcal{S} = \sum_j \int \frac{d\lambda}{2\pi} (\partial_\lambda p_j)^{\text{dr}} \eta(\vartheta_j(\lambda))$$

$$\eta(x) = -x \log x - (1-x) \log(1-x)$$

$$\lambda = -\infty \quad \lambda = +\infty$$

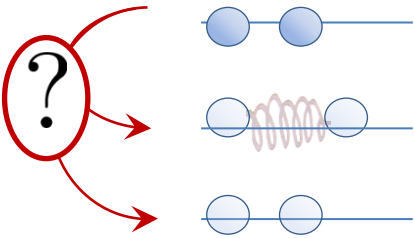
$$\partial_\Phi \mathcal{S} = \partial_\Phi \mathcal{S}^- + \partial_\Phi \mathcal{S}^+$$

Starting from the GS





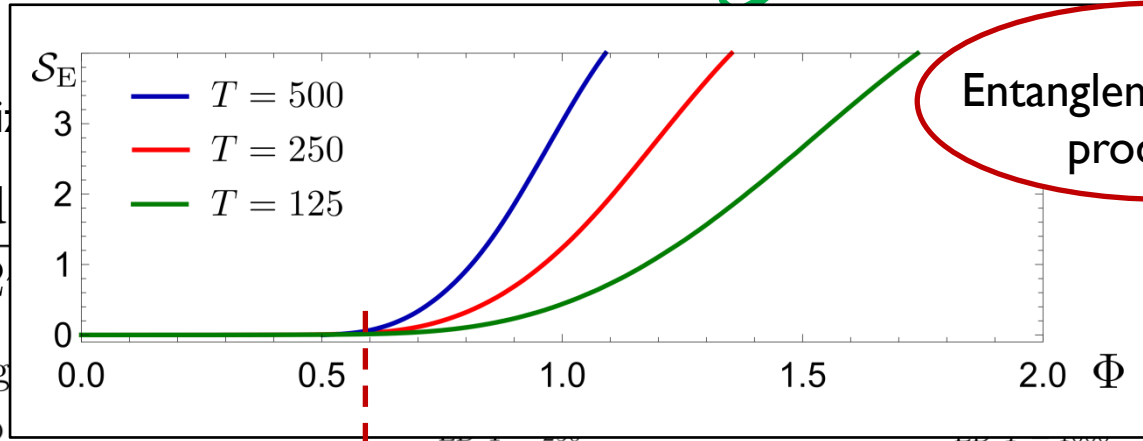
GHD, reversibility and beyond



$$\text{GGEs} = \text{Charges conservation} + \text{Entropy maximization}$$

Entropy rate maximization

$$S = \sum_j \int \frac{d}{2} \eta(x) = -x \log$$

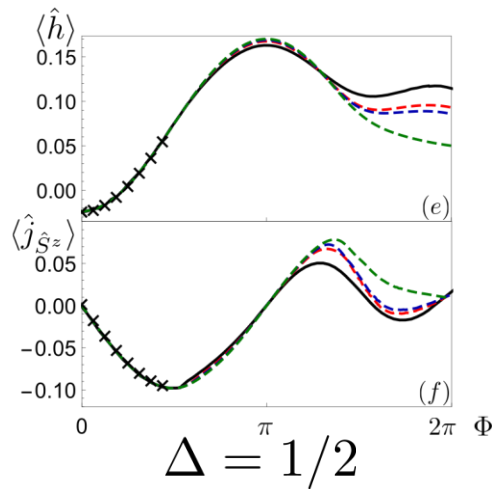
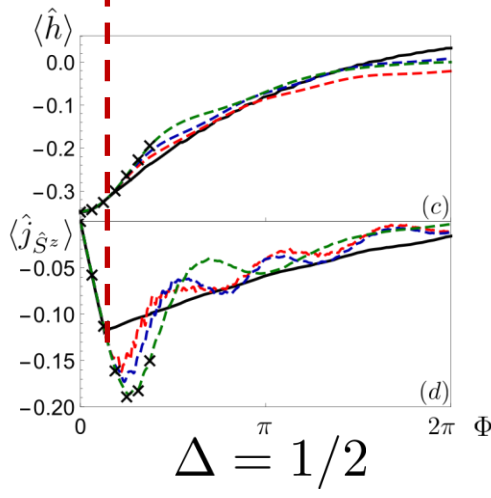
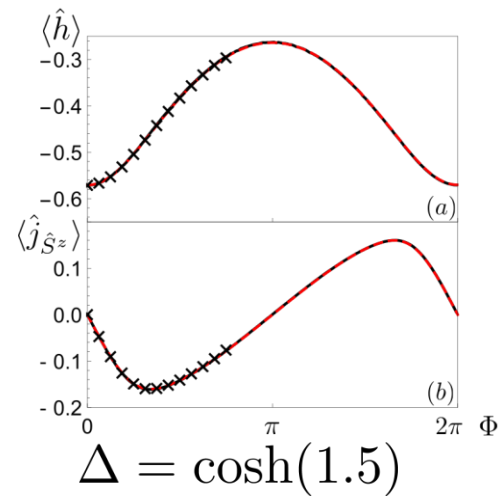


Entanglement entropy production

$$= +\infty$$

$$\frac{\partial}{\partial \Phi} S^+$$

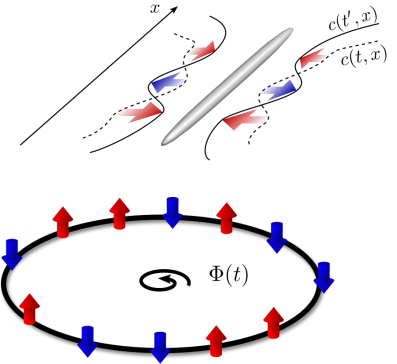
Starting from the GS





Conclusions and outlook

- ★ GHD can describe (locally) integrable systems with (smooth) inhomogeneous couplings
- ★ Sometimes smooth inhomogeneities are not smooth

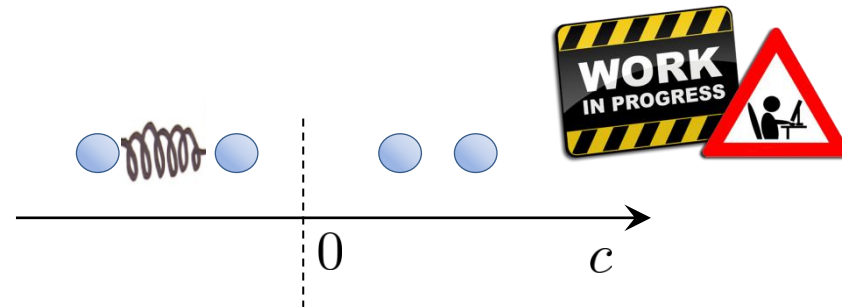


What's next?

General framework to handle smooth inhomogeneities of “non-smooth” integrable models ?

First step

Interaction changes from repulsive to attractive phase in Lieb-Liniger





THANK YOU!



A. De Luca



V. Alba



J.-S. Caux