Analogue of many-body Berry-Esseen theorem for critical systems

Karen Hovhannisyan



Adriatic Conference on Strongly Correlated Systems, ICTP, 22 Mar 2021

More details at

arXiv:2011.10513 [quant-ph]

In collaboration with

Mathias Jørgensen (Technical University of Denmark)

Gabriel Landi (Universidade de São Paulo)

Álvaro Alhambra (Max-Planck-Institut für Quantenoptik)

Jonatan Brask (Technical University of Denmark)

Martí Perarnau-Llobet (Université de Genève)



$$H = \sum_{\langle i,k \rangle} [atom_i] \otimes [atom_k] + \sum_i V(atom_i)$$





Energy distribution





Energy distribution





$$J(E) = \sum_{E_i \le E} \langle E_i | \rho | E_i \rangle$$



Brandão & Cramer, arXiv:1502.03263 [quant-ph]







 $\|H_v\| \leq h$









$$\frac{\langle A \otimes B \rangle_{\rho} - \langle A \rangle_{\rho} \langle B \rangle_{\rho}}{\|A\| \cdot \|B\|} \le e^{-D(A,B)/\xi}$$



$$\frac{\langle A \otimes B \rangle_{\rho} - \langle A \rangle_{\rho} \langle B \rangle_{\rho}}{\|A\| \cdot \|B\|} \le e^{-D(A,B)/\xi}$$

$$\mu^2 \coloneqq \operatorname{var}(H) \coloneqq \langle H^2 \rangle_\rho - \langle H \rangle_\rho^2 = s^2 N$$

Brandão & Cramer, arXiv:1502.03263 [quant-ph]



$$\sup_{E} |J(E) - G(E)| \le O\left(\frac{\ln^{2d} N}{\sqrt{N}}\right)$$

$$G(E) = \frac{1}{\sqrt{2\pi\mu^2}} \int_{-\infty}^{E} dE \ e^{-(E - \langle H \rangle)^2 / (2\mu^2)}$$

Brandão & Cramer, arXiv:1502.03263 [quant-ph]



$$\sup_{E} |J(E) - G(E)| \le O\left(\frac{\ln^{2d} N}{\sqrt{N}}\right) \xrightarrow{N \to \infty} 0$$

$$G(E) = \frac{1}{\sqrt{2\pi\mu^2}} \int_{-\infty}^{E} dE \ e^{-(E - \langle H \rangle)^2 / (2\mu^2)}$$

What about critical systems?

An analogue of Berry-Esseen theorem holds

... but is valid only for

- Thermal states
- Finite-temperature phase transitions
- Translation-invariant lattices



$$\Gamma(E) = \#\{E_i \le E\}$$











$$q(E) \coloneqq \frac{dJ(E)}{dE} = \frac{e^{-\beta E}}{Z} \Omega(E)$$



$$\lim_{N\to\infty}\frac{\ln\Gamma(uN)}{N}=s(u)$$







$$\lim_{N \to \infty} \frac{\ln \Gamma(uN)}{N} = \underbrace{s(u)}_{N}$$
density of canonical
entropy corresponding to
average energy density *u*





$$\lim_{N \to \infty} \frac{\ln \Gamma(uN)}{N} = s(u)$$

$$\approx \text{microcanonical entropy}$$
density at u



$$\lim_{N \to \infty} \frac{\ln \Gamma(uN)}{N} = s(u)$$

expresses the equivalence
between canonical and
microcanonical ensembles



$$q(uN) = e^{\beta N f_N(\beta) - \beta N [u - Ts_N(u)] + O(1)}$$



$$q(uN) = e^{\beta N f_N(\beta) - \beta N [u - Ts_N(u)] + O(1)}$$

$$\min_{\mathbf{u}}[u - Ts_N(u)] = f_N(\beta)$$



$$q(uN) = e^{\beta N f_N(\beta) - \beta N [u - Ts_N(u)] + O(1)}$$

$$\min_{u} [u - Ts_{N}(u)] = f_{N}(\beta)$$
$$u - Ts_{N}(u) = f_{N}(\beta) - \frac{\beta}{2c_{N}(\beta)} [u - u_{N}(\beta)]^{2} + \cdots$$
specific heat



$$q(uN) = e^{\beta N f_N(\beta) - \beta N [u - Ts_N(u)] + O(1)}$$

$$\min_{\mathbf{u}} [u - Ts_N(u)] = f_N(\beta)$$
$$u - Ts_N(u) = f_N(\beta) - \frac{\beta}{2c_N(\beta)} [u - u_N(\beta)]^2 + \cdots$$

canonic energy density

$$c(\beta) \propto |\beta - \beta_c|^{-\alpha}$$

The scaling of $c(\beta_c)$ with *N*:

When $\alpha = 0$, $c(\beta) \propto \ln N$ When $\alpha > 0$, $c(\beta) \propto N^{\alpha/(2-\alpha)}$ Whenever

$$\frac{|E - \langle H \rangle|}{\sqrt{\operatorname{var}(H)}} \le \sqrt{\ln N}$$

$$q(E) \propto e^{-\frac{(E - \langle H \rangle)^2}{2 \operatorname{var}(H)} + O\left(\frac{1}{\ln N}\right)}$$

Whenever

$$\frac{|E - \langle H \rangle|}{\sqrt{\operatorname{var}(H)}} \le \sqrt{\ln N}$$

$$q(E) \propto e^{-\frac{(E-\langle H \rangle)^2}{2\operatorname{var}(H)} + O\left(\frac{1}{\ln N}\right)}$$

Whenever

$$\frac{|E - \langle H \rangle|}{\sqrt{\operatorname{var}(H)}} > \sqrt{\ln N}$$

$$\int_{|E-\langle H\rangle|>\sqrt{N}\operatorname{var}(H)} dE \ q(E) \propto \frac{1}{N}$$

2D Ising model is a paradigmatic example of $\alpha = 0$.

Its free energy can be computed exactly, and therefore one has access to all moments of the energy.

$$\frac{|\kappa_n|^{1/n}}{\sqrt{\operatorname{var}(H)}} \propto \frac{1}{\sqrt{\ln N}}$$

2D Ising model is a paradigmatic example of $\alpha = 0$.

Its free energy can be computed exactly, and therefore one has access to all moments of the energy.



Not much can be said except that q(E) is a unimodal distribution peaked around $\langle H \rangle$ and decaying exponentially in the tails.

Thank you for your attention!