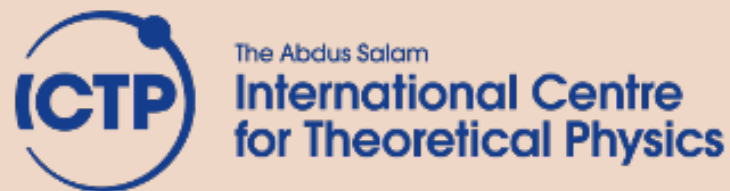


MAX School on Advanced Materials and Molecular Modelling  
with QUANTUM ESPRESSO

QUANTUM ESPRESSO: overview and basic  
functionalities. The self-consistent cycle. PBC:  
supercells and k-point sampling

Ralph Gebauer

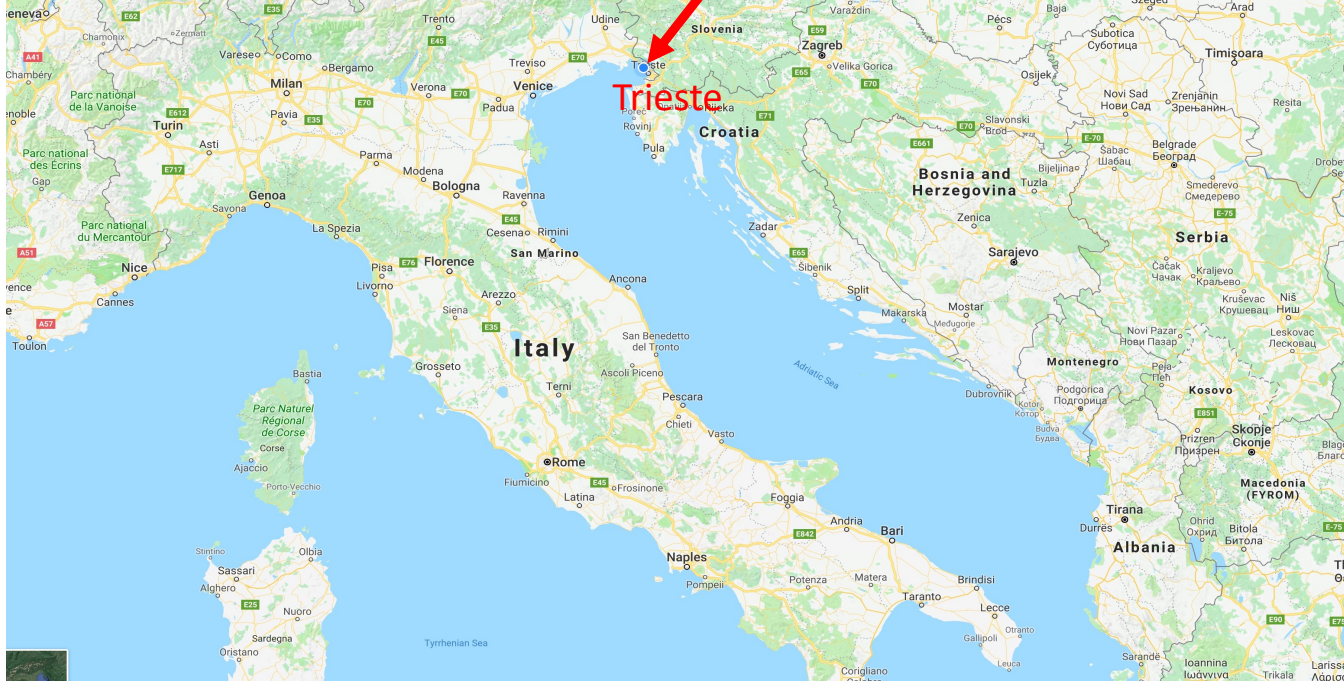
This school is a large-scale collaboration:







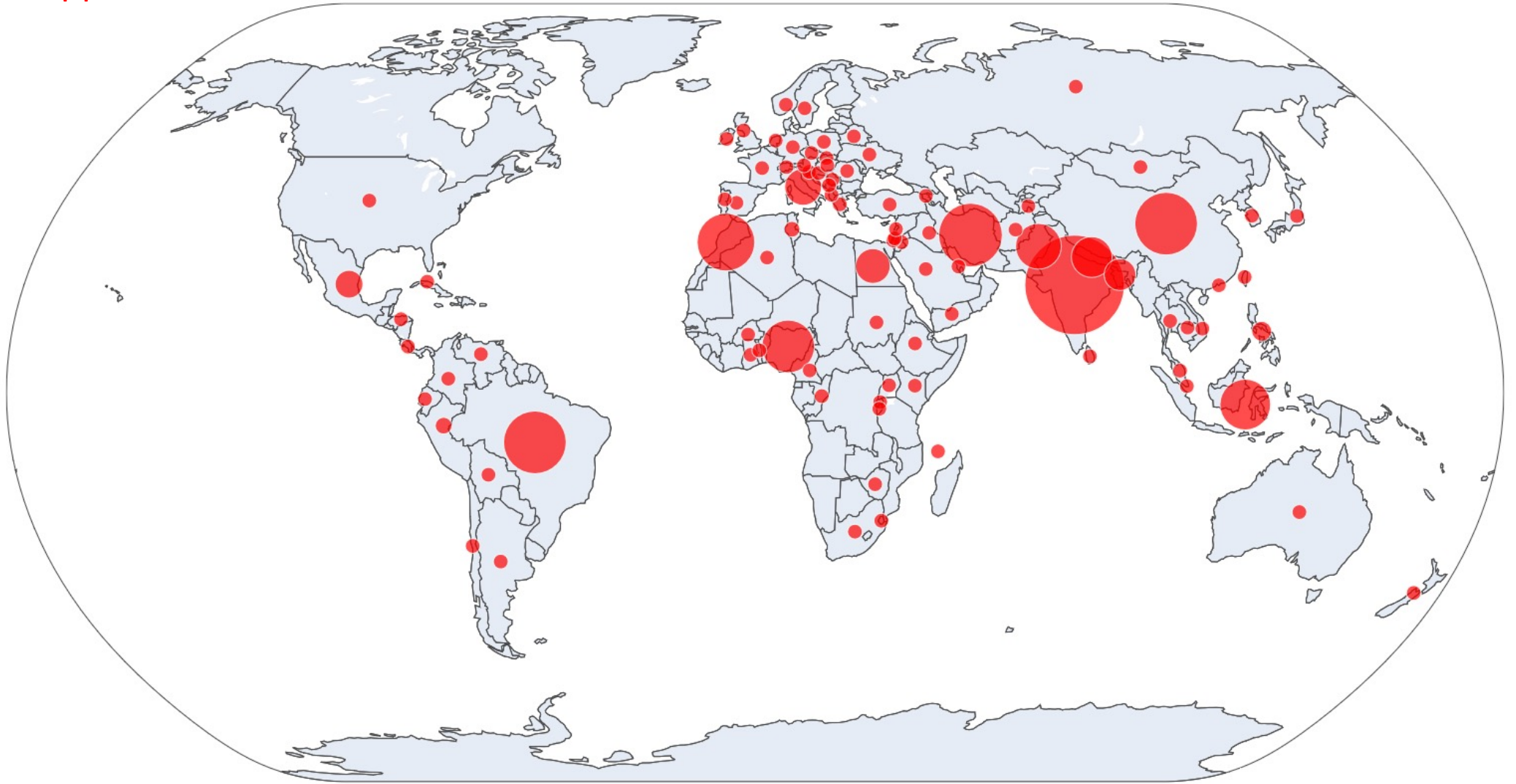




Founded in 1964, the Abdus Salam International Centre for Theoretical Physics (ICTP) operates under the aegis of two United Nations Agencies: UNESCO and IAEA and is regularized by a seat agreement with the Government of Italy.



1292 applicants!



Special thanks to the secretaries: Monica Ancuta, Viktoriya Lvova and Adriana Pinto

The star of this school:

# Quantum-ESPRESSO

Quantum ESPRESSO: Quantum opEn-Source Package for Research in Electronic Structure, Simulation, and Optimization, is a distribution (an integrated suite) of software for first-principle simulations, i.e., atomistic calculations based on electronic structure, using density-functional theory, a plane-wave basis set, pseudopotentials.

QE exists since 2002, resulting from the merge of pre-existing packages; some core components have been under development for ~ 30 years

The main goals of Quantum-ESPRESSO are:

- innovation in theoretical methods and numerical algorithms
- efficiency on modern computer architectures







The star of this school:

# Quantum-ESPRESSO License:

QE is distributed under the GNU (Gnu's Not Unix) General Public License (GPL v.2), probably the most common free-software license.

- The source code is available.
- You can do whatever you want with the sources, but if you distribute any derived work, you have to distribute under the GPL the sources of the derived work.

Advantages:

- Everybody – including commercial entities – can contribute.
- Nobody can “steal” the code and give nothing back to the community.

The most successful example is probably the Linux Kernel.



The star of this school:

# Quantum-ESPRESSO

You can get QE help, answers and general info from various sources:

the documenting papers:

[J. Phys.: Condens. Matter 21, 395502 \(2009\)](#) and

[J. Phys.: Condens. Matter 29, 465901 \(2017\)](#)

the web site: [www.quantum-espresso.org](http://www.quantum-espresso.org)

the documentation in the [Doc subdirectories](#)

the developers' portal: [github.com/QEF/q-e/releases](https://github.com/QEF/q-e/releases)

the mailing lists:

[users@lists.quantum-espresso.org](mailto:users@lists.quantum-espresso.org)

[developers@lists.quantum-espresso.org](mailto:developers@lists.quantum-espresso.org)

(use first the rich [ARCHIVE](#) of these lists)





# Quantum-ESPRESSO as a distribution

QE is not a single, executable code. It is composed of several packages:

**PWscf:** self-consistent electronic structure, (variable-cell) structural optimization, molecular dynamics

**CP:** Car-Parrinello molecular dynamics, also with variable cell

They share a common installation method, input format, pseudopotential format, data output format, large parts of the basic code.

Further codes:

**PHonon:** linear-response calculations (phonons, dielectric properties)

**PostProc:** graphical and postprocessing utilities (density of states, STM, etc.)

**PWneb:** Nudged Elastic Band (NEB) for reaction pathways and barriers

**atomic:** pseudopotential generation code

**PWGui:** a Graphical User Interface for production of input files

**PWcond:** ballistic conductance

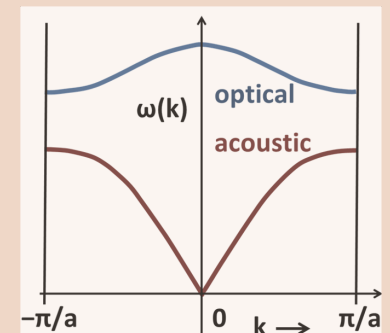
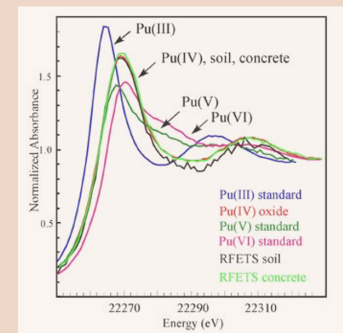
**XSpectra:** Calculation of X-ray near-edge adsorption spectra (XANES)

**GWL:** GW band structure with ultralocalized Wannier functions

**TD-DFPT:** Time-Dependent Density-Functional Perturbation Theory

**EPW:** Electron-phonon coefficients and related properties

**HP:** Hubbard parameters from linear response



# Density Functional Theory

The ground-state energy of a system of  $N$  electrons is a functional of the electronic density  $n(\mathbf{r})$ :

$$E^{DFT}[n] \quad n(\mathbf{r}) \geq 0 \quad \int d\mathbf{r} n(\mathbf{r}) = N$$

The exact functional is unknown. Within the Kohn-Sham (KS) formalism, we write:

$$E^{DFT}[n] = T_s[\{\psi_i\}] + E_{ext}[n] + E_{Har}[n] + E_{xc}[n] + E_{Ions}$$

Where the KS orbitals have been introduced to approximate the kinetic energy:

$$T_s[\{\psi_i\}] = -\frac{\hbar^2}{2m} \sum_{i=1}^N \int d\mathbf{r} \psi_i^*(\mathbf{r}) \nabla^2 \psi_i(\mathbf{r})$$

$$\langle \psi_i | \psi_j \rangle = \delta_{i,j} \quad i, j = 1 \dots N$$

$$n(\mathbf{r}) = \sum_{i=1}^N \psi_i^*(\mathbf{r}) \psi_i(\mathbf{r})$$

$$E_{ext}[n] = \int d\mathbf{r} n(\mathbf{r}) V_{ext}(\mathbf{r})$$

$$E_{Har}[n] = \frac{e^2}{2} \int d\mathbf{r} d\mathbf{r}' \frac{n(\mathbf{r}) n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$E_{Ions} = \sum_{IJ} \frac{Z_I Z_J e^2}{|\mathbf{R}_I - \mathbf{R}_J|}$$

The main difficulty lies in approximating the exchange-correlation functional  $E_{xc}[n]$



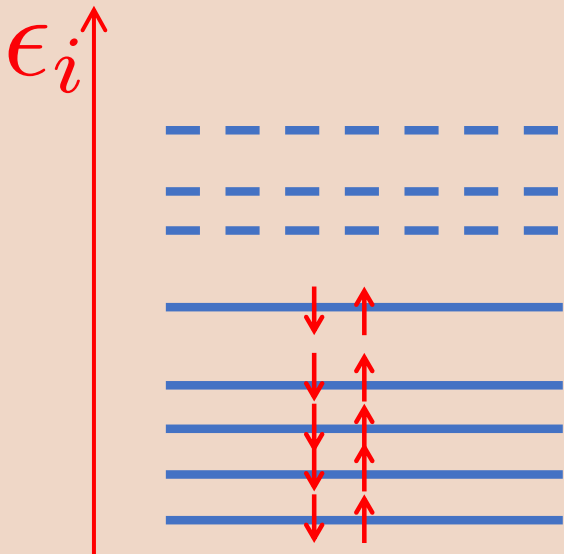
# Density Functional Theory

Minimization of  $E^{\text{DFT}}[n]$  with respect to  $n$  (actually, with respect to the KS orbitals) leads to the KS equations:

$$H^{KS} \psi_i(\mathbf{r}) = \left( -\frac{\hbar^2}{2m} \nabla^2 + V_{ext}(\mathbf{r}) + V_{Har}(\mathbf{r}) + V_{xc}(\mathbf{r}) \right) \psi_i(\mathbf{r}) = \epsilon_i \psi_i(\mathbf{r})$$

$$V_{xc}(\mathbf{r}) = \frac{\delta E_{xc}}{\delta n(\mathbf{r})}$$

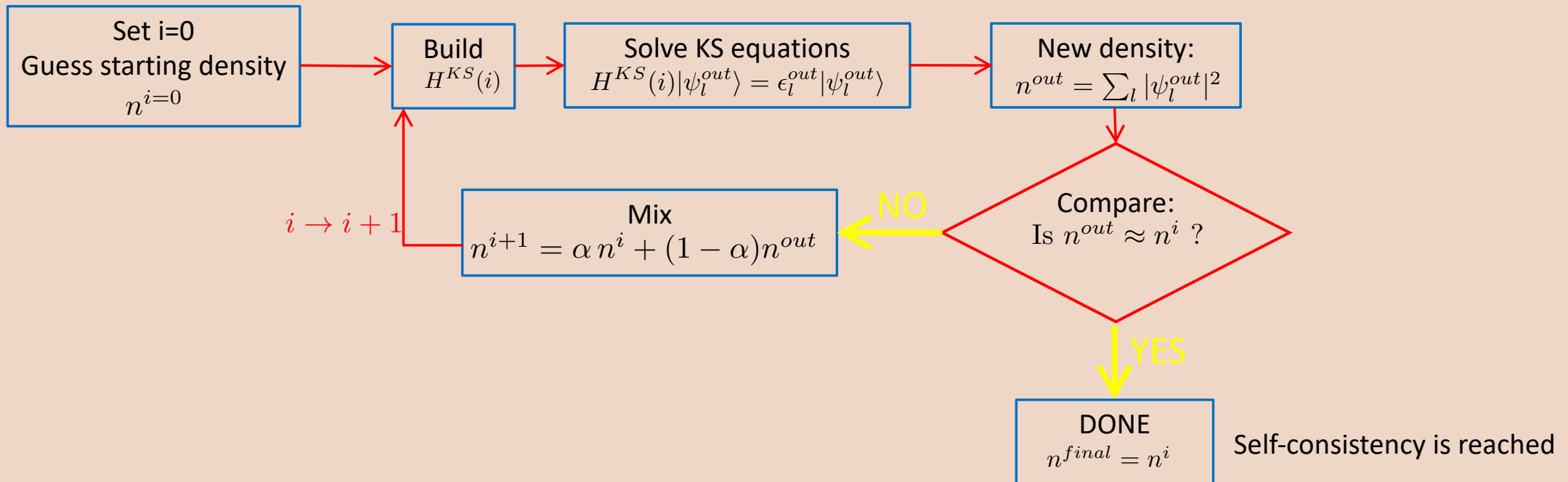
$$V_{Har}(\mathbf{r}) = \frac{\delta E_{Har}}{\delta n(\mathbf{r})} = e^2 \int d\mathbf{r}' \frac{n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$



Problem:  $H^{KS}$  depends on the charge density ....

# Density Functional Theory

## The self-consistent cycle



## The self-consistent cycle

Initial potential from superposition of free atoms

starting charge 29.42246, renormalised to 30.00000

negative rho (up, down): 7.551E-05 0.000E+00

Starting wfcs are 30 randomized atomic wfcs

total cpu time spent up to now is 1.5 secs

Self-consistent Calculation

iteration # 1 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 1.00E-02, avg # of iterations = 3.0

negative rho (up, down): 2.325E-05 0.000E+00

total cpu time spent up to now is 2.4 secs

total energy = -75.13217408 Ry

estimated scf accuracy < 1.92237710 Ry

iteration # 2 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 6.41E-03, avg # of iterations = 2.0

negative rho (up, down): 1.858E-05 0.000E+00

total cpu time spent up to now is 3.1 secs

total energy = -75.22539882 Ry

estimated scf accuracy < 0.08211827 Ry

iteration # 3 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 2.74E-04, avg # of iterations = 6.0

negative rho (up, down): 5.570E-05 0.000E+00

(...)

iteration # 10 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 6.88E-09, avg # of iterations = 7.0

negative rho (up, down): 3.572E-04 0.000E+00

total cpu time spent up to now is 9.8 secs

total energy = -75.28039726 Ry

estimated scf accuracy < 0.00000156 Ry

iteration # 11 ecut= 20.00 Ry beta= 0.70

Davidson diagonalization with overlap

ethr = 5.20E-09, avg # of iterations = 3.0

negative rho (up, down): 3.554E-04 0.000E+00

total cpu time spent up to now is 10.5 secs

End of self-consistent calculation

k = 0.0000 0.0000 0.0000 ( 4318 PWs) bands (ev):

21.1685 -18.3817 -18.3806 -14.7643 -14.7544 -12.8025 -11.1007 -10.8244  
10.1315 -10.1240 -8.9981 -8.1363 -8.1213 -6.2696 -6.2696 -1.1442

highest occupied, lowest unoccupied level (ev): -6.2696 -1.1442

total energy = -75.28039766 Ry

estimated scf accuracy < 0.00000008 Ry

The total energy is the sum of the following terms:

one-electron contribution = -518.40332837 Ry

hartree contribution = 262.27253441 Ry

xc contribution = -25.13397296 Ry

ewald contribution = 205.98436926 Ry

convergence has been achieved in 11 iterations

## Representing charge densities, KS orbitals in a computer:

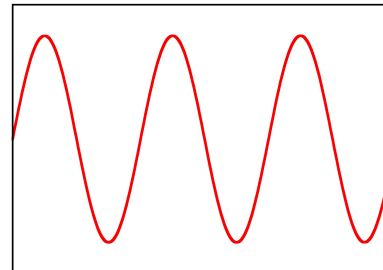
### Basis sets

$$f(\mathbf{r}) = \sum_{\alpha=1}^M c_{\alpha} b_{\alpha}(\mathbf{r})$$

Diagram illustrating the expansion of a function  $f(\mathbf{r})$  in terms of basis functions  $b_{\alpha}(\mathbf{r})$  with expansion coefficients  $c_{\alpha}$ . The summation index  $\alpha$  ranges from 1 to  $M$ , where  $M$  is the size of the basis set. Red arrows point from the terms in the equation to their respective labels:  $f(\mathbf{r})$  to "Function to be represented",  $c_{\alpha}$  to "Expansion coefficient",  $b_{\alpha}(\mathbf{r})$  to "Basis function", and  $M$  to "Size of basis set".

In Quantum-ESPRESSO, plane waves are used as basis set:

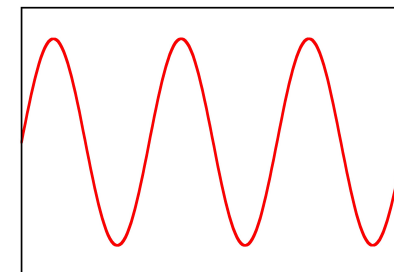
$$b_{\alpha}(\mathbf{r}) = \frac{1}{\sqrt{V}} \exp(i \mathbf{G}_{\alpha} \cdot \mathbf{r})$$





## Advantages of plane waves as basis set

$$b_{\alpha}(\mathbf{r}) = \frac{1}{\sqrt{V}} \exp(i \mathbf{G}_{\alpha} \cdot \mathbf{r})$$



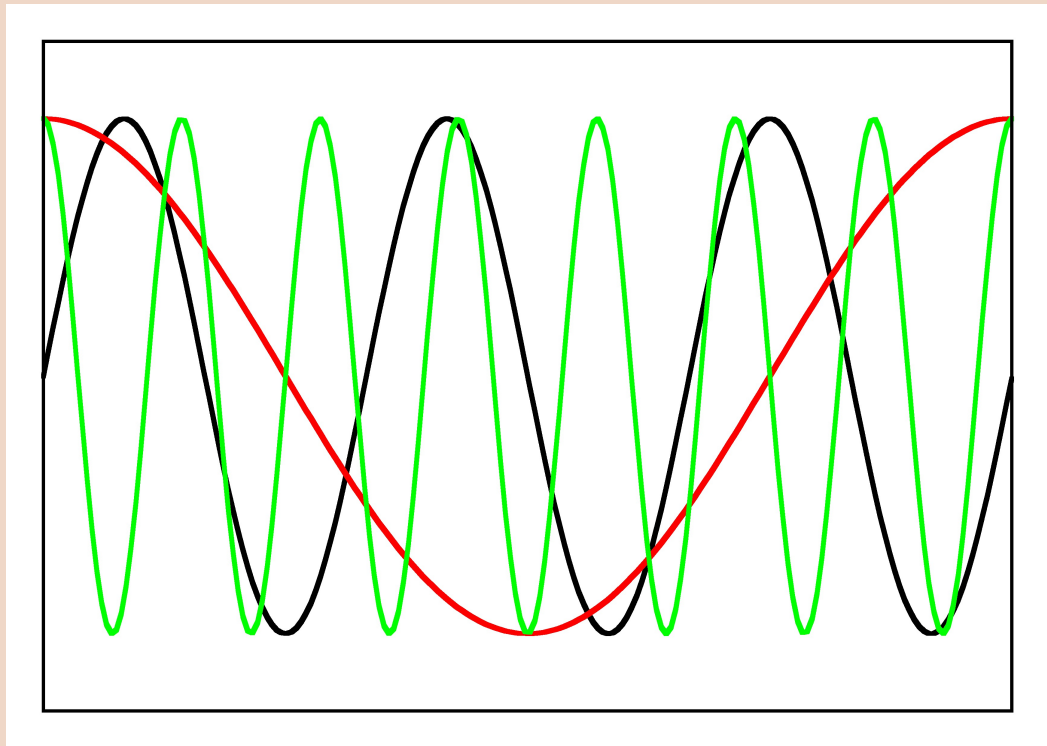
Using plane-waves as a basis set has many advantages:

- Simple analytical form: derivatives and integrals are easy to perform
- PWs are orthonormal
- Unbiased: no assumptions where charges/electrons are localized
- Independent of atomic positions: no “Pulay forces”
- Easy to control convergence of basis set size (see following slides ...)
- Straightforward use of FFTs: easy use of R- and G-space dualities

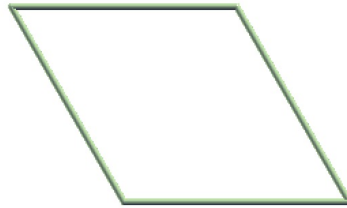
However: PW basis set are typically MUCH bigger than other sets of (localized) basis functions

# Periodicity

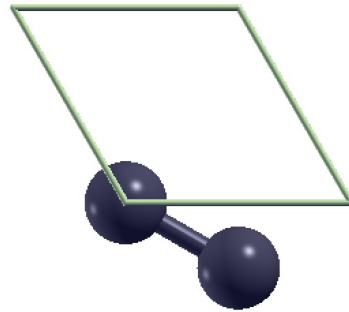
Due to the periodic nature of PWs, the use of this basis set is closely linked with the periodicity of the physical system.



Periodicity

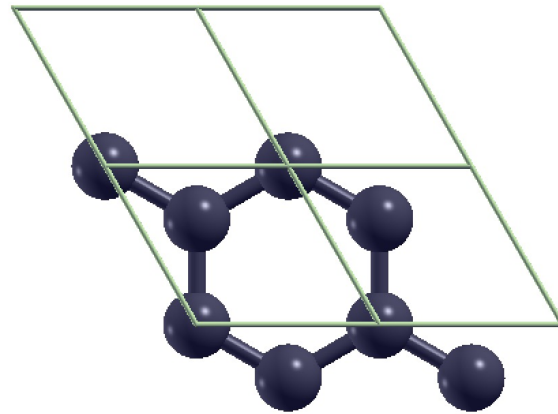


# Periodicity

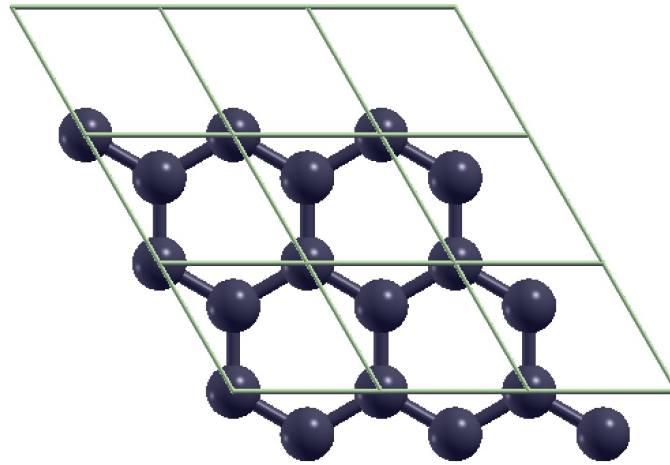




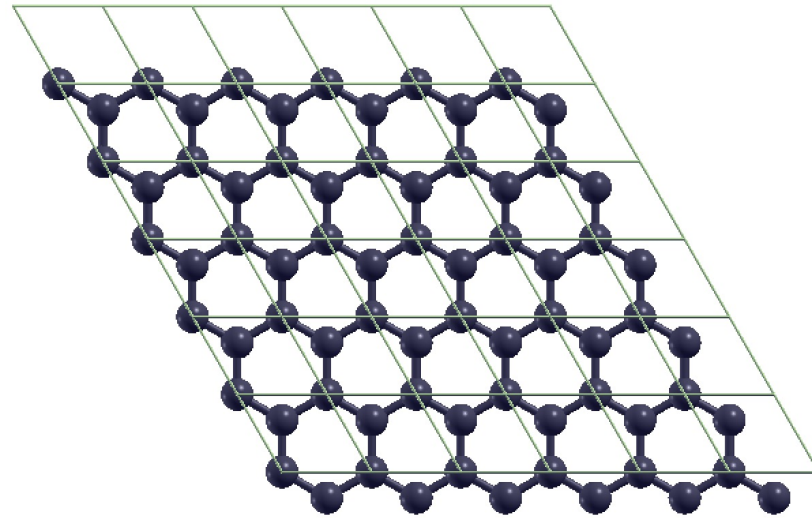
# Periodicity



# Periodicity



# Periodicity



## Periodicity

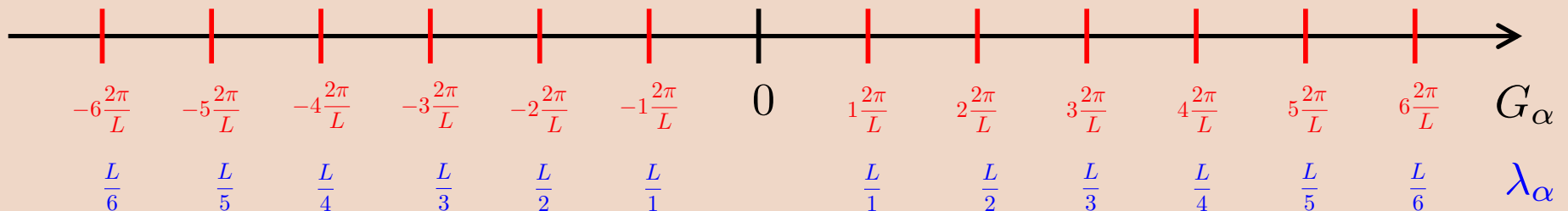
If a function is periodic in real space, its Fourier transform is non-zero only for discrete wavefactors

$$\text{In 1-D: } \exp(iG_\alpha x) \rightarrow G_\alpha = n \times \frac{2\pi}{L} \quad n \in \mathbb{Z}$$

$$\text{In 3-D: } \exp(i\mathbf{G}_\alpha \cdot \mathbf{r}) \rightarrow \mathbf{G}_\alpha = m \times \mathbf{B}_1 + n \times \mathbf{B}_2 + p \times \mathbf{B}_3 \quad m, n, p \in \mathbb{Z}$$

(where  $\mathbf{B}_{1,2,3}$  are reciprocal lattice vectors)

1-D illustration:



Wavelength  $\lambda_\alpha$  is inversely proportional to  $|\mathbf{G}_\alpha|$

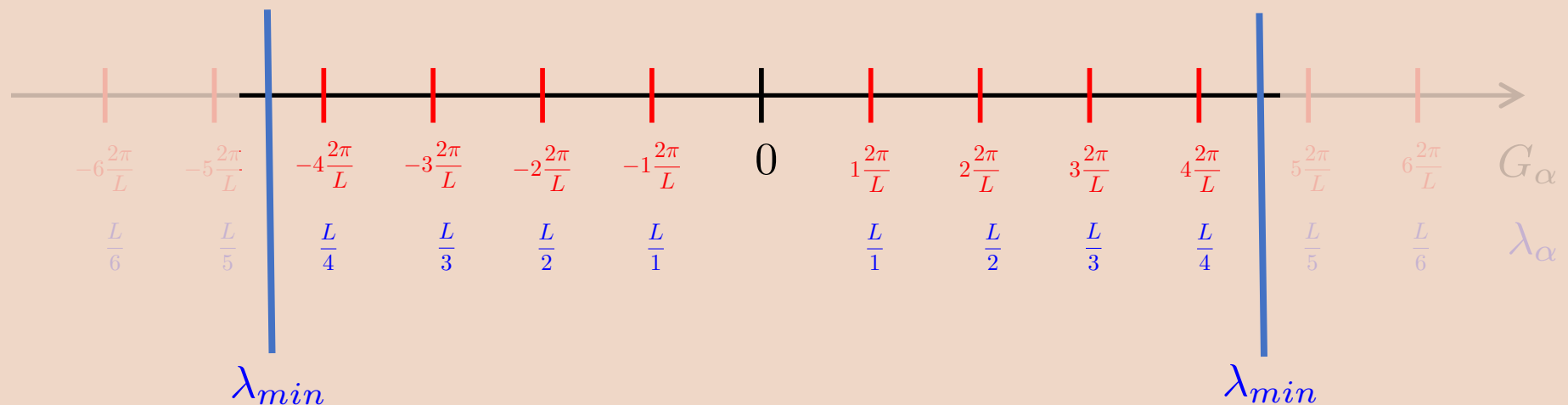


# Periodicity

The number of PWs compatible with a given periodicity is **infinite**.

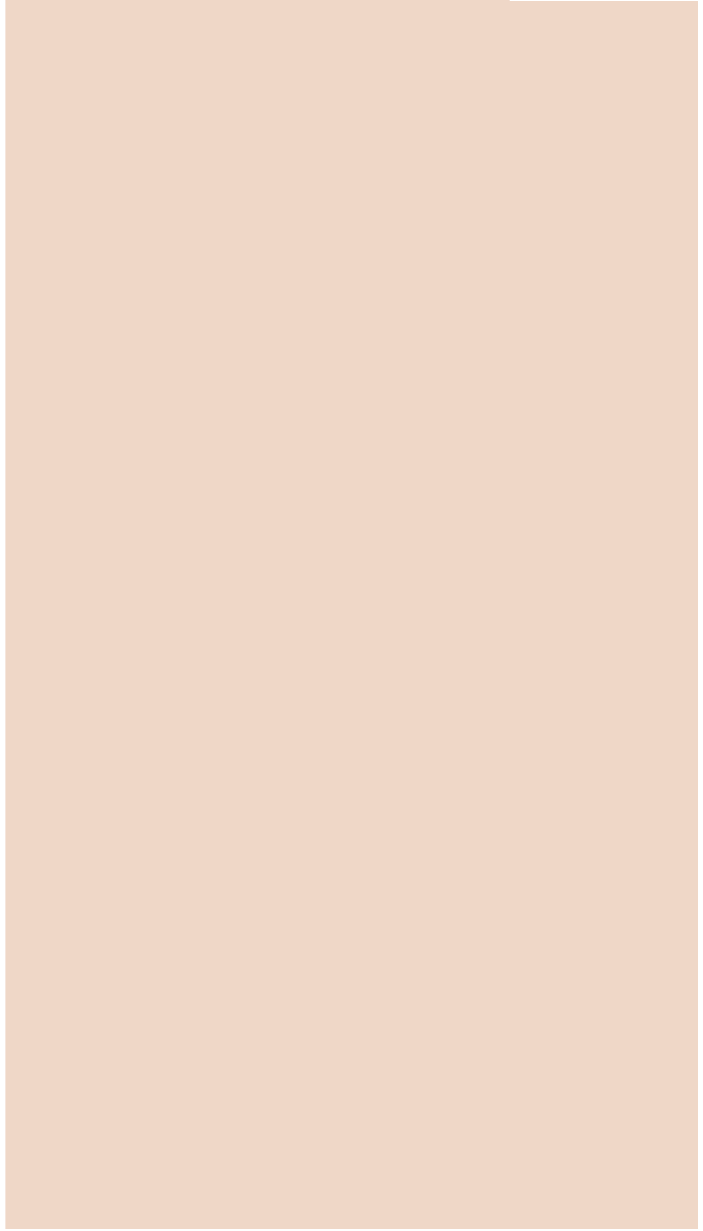
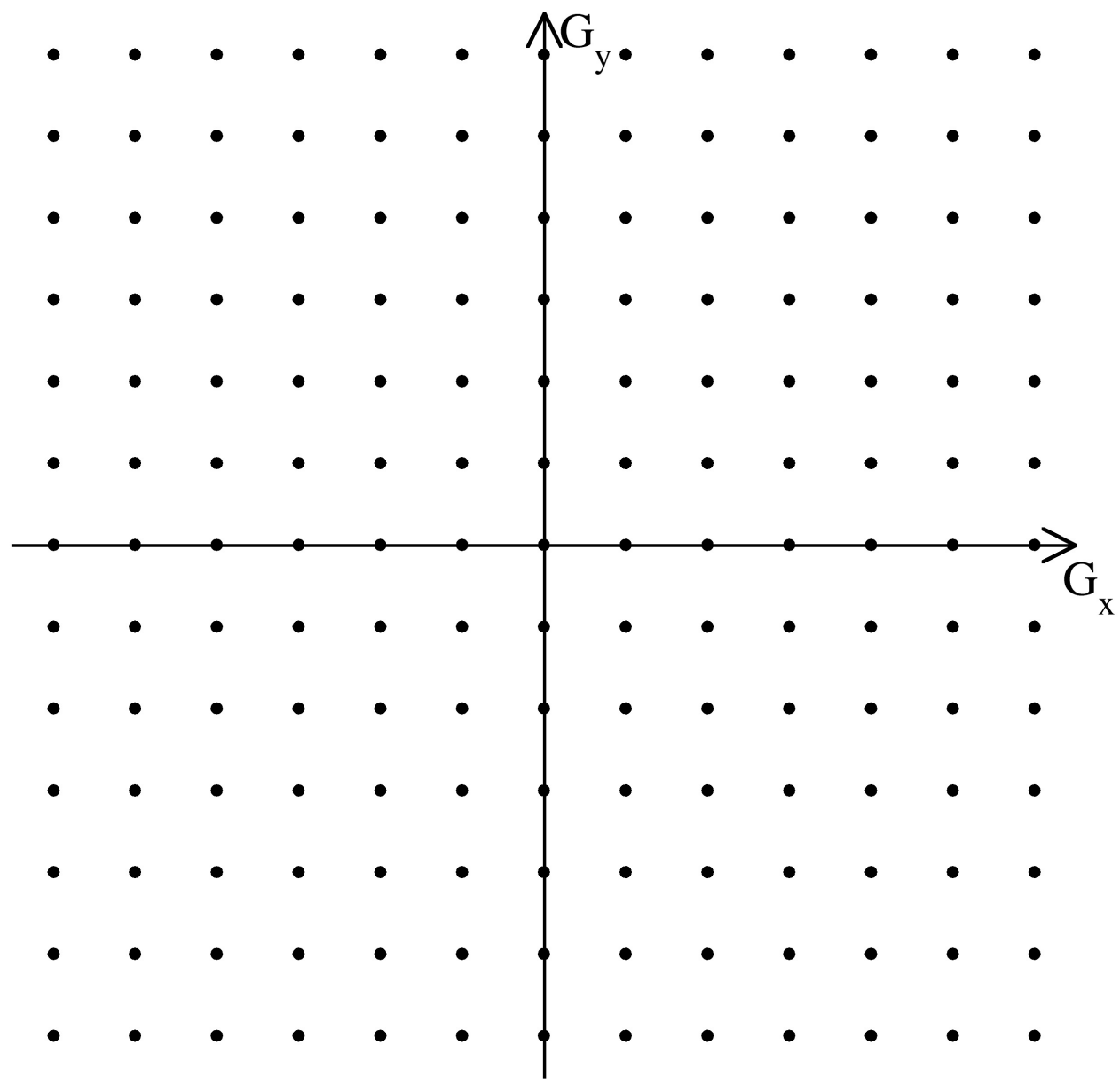
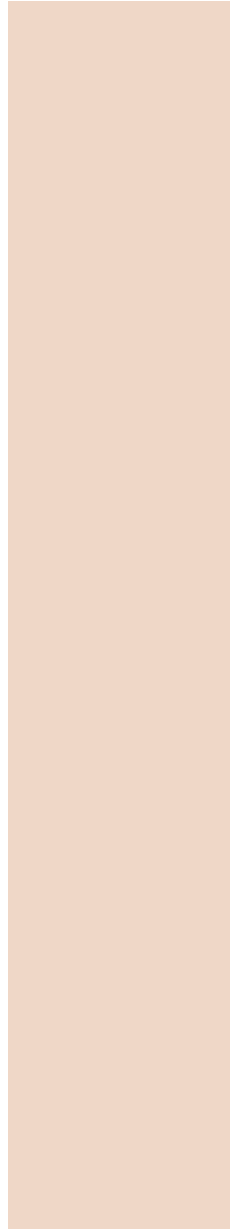
In order to obtain a **finite** number of basis functions, one fixes a “smallest” feature size:  $\lambda_{min}$

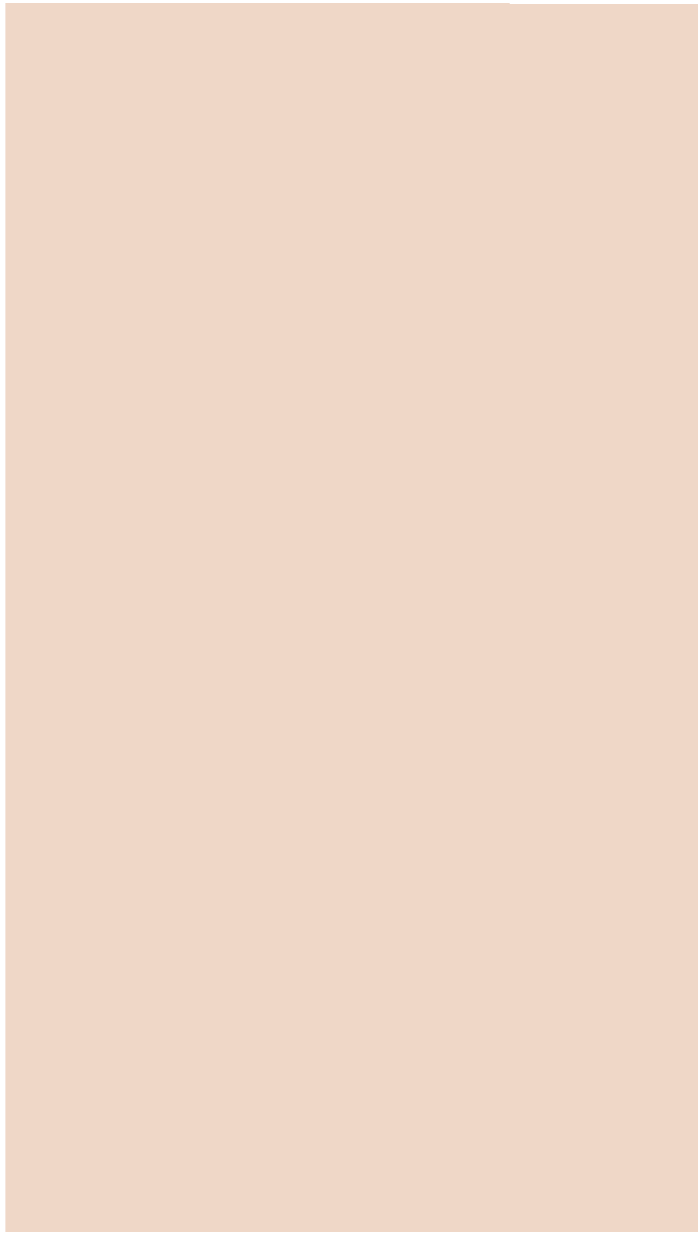
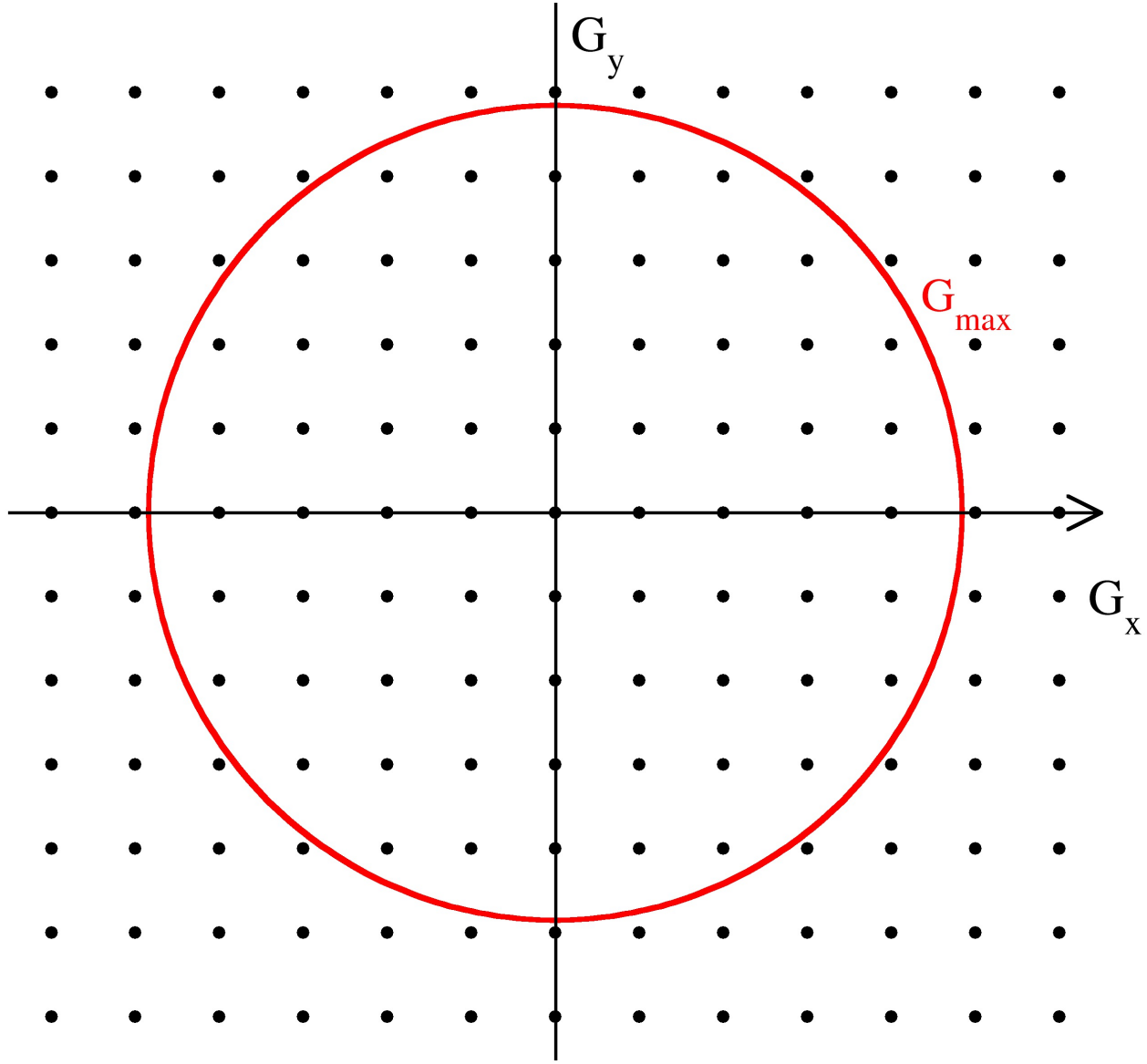
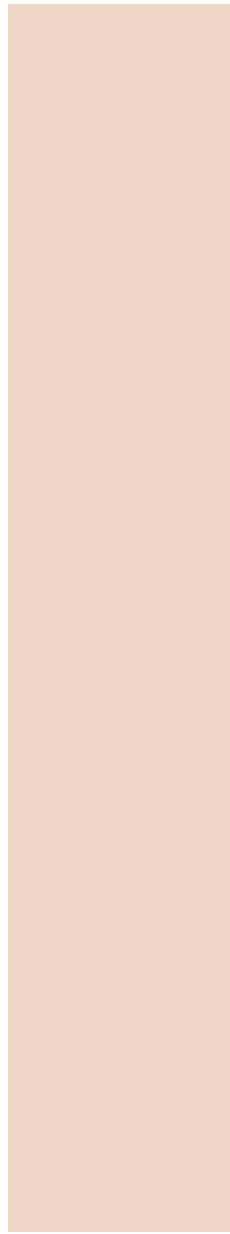
1-D illustration:

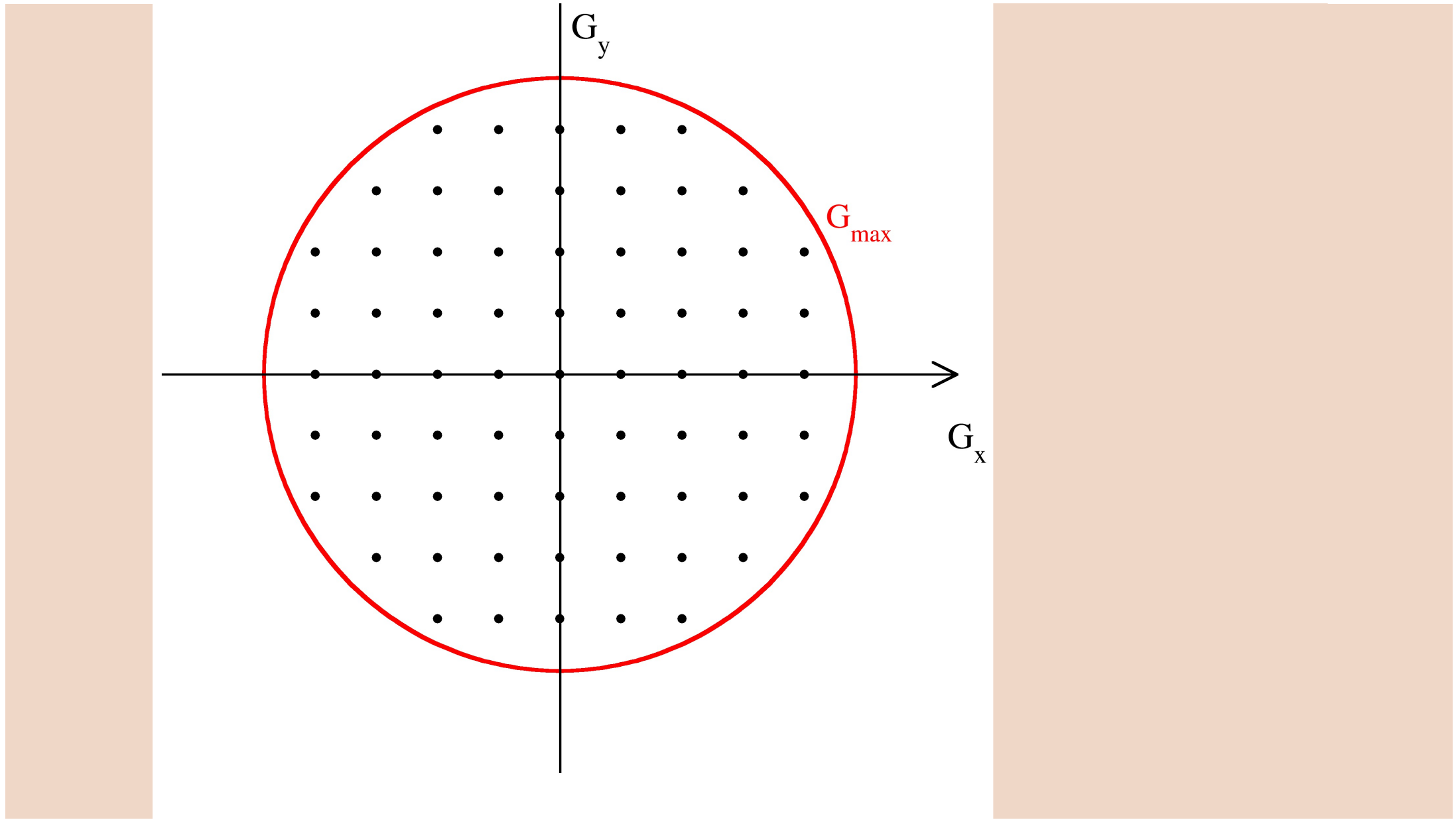


In practice, one describes a **maximal norm** for the wavevector:  $|\mathbf{G}|_{max} = \frac{2\pi}{\lambda_{min}}$

By setting a **cutoff kinetic energy** for the plane-waves:  $E_{cut} = \frac{\hbar^2}{2m} |\mathbf{G}|_{max}^2$









## From KS orbitals to the charge density

Imagine that orbitals are represented in PWs, with a cutoff  $G_{\max}$

$$\psi_l(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{\max}} c(\mathbf{G}) \exp(i\mathbf{G} \cdot \mathbf{r})$$

The charge density is given by the squared modulus of the orbitals:

$$n(\mathbf{r}) = \sum_l \psi_l^*(\mathbf{r}) \psi_l(\mathbf{r})$$
$$\tilde{n}(\mathbf{G}) = \sum_l \sum_{|\mathbf{G}'| < G_{\max}} \tilde{\psi}_l^*(\mathbf{G}') \tilde{\psi}_l(\mathbf{G} - \mathbf{G}')$$

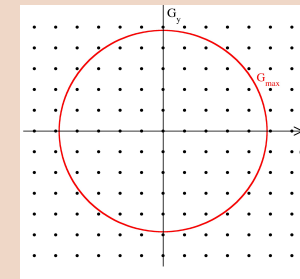
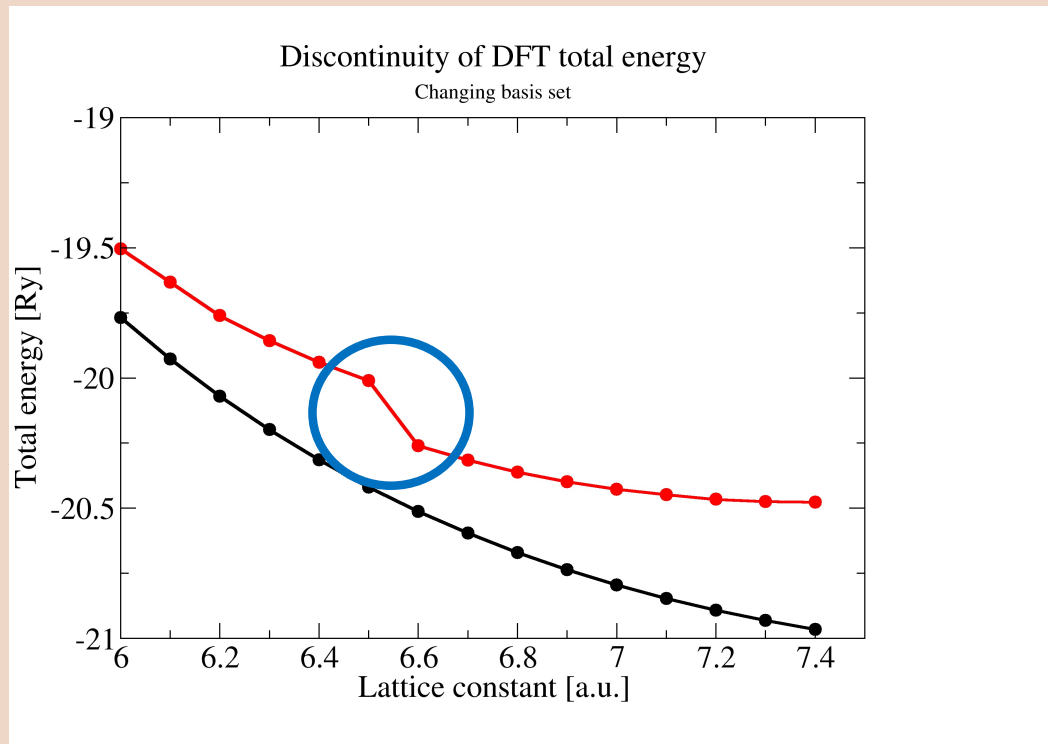
If the orbitals are represented with a cutoff of  $G_{\max}$ , then the charge density has non-zero Fourier components up to  $2 * G_{\max}$  !!! This means:  $E_{\text{cut}}(\text{rho}) = 4 * E_{\text{cut}}(\text{wfc})$

## From KS orbitals to the charge density

```
&control
  calculation='scf'
  prefix='diamond'
  pseudo_dir='./'
/
&system
 ibrav = 2
  celldm(1)=7.4
  nat=2
  ntyp=1
  ecutwfc=60.
/
&electrons
/
```

```
bravais-lattice index      =          2
lattice parameter (alat)  =       7.4000  a.u.
unit-cell volume          =    101.3060 (a.u.)^3
number of atoms/cell      =          2
number of atomic types    =          1
number of electrons       =         8.00
number of Kohn-Sham states =          4
kinetic-energy cutoff      =    60.0000  Ry
charge density cutoff     =    240.0000  Ry
scf convergence threshold =    1.0E-06
mixing beta               =         0.7000
number of iterations used =          8  plain  mixing
Exchange-correlation= SLA  PW  PBX  PBC
                        (  1  4  3  4  0  0  0)
```

Beware: changing the lattice constant means changing the PW basis set



Larger cell size

→ grid of G-vectors more dense

→ at fixed  $G_{\max}$ , "new" plane waves

enter the calculation discontinuously

→ sudden decrease of energy when basis set is enlarged

Curing the discontinuity: Check for input variables `ecfixed`, `qcutz`, `q2sigma`

So have seen how periodic functions are expressed in PWs ....

BUT: orbitals in periodic systems are NOT lattice periodic!

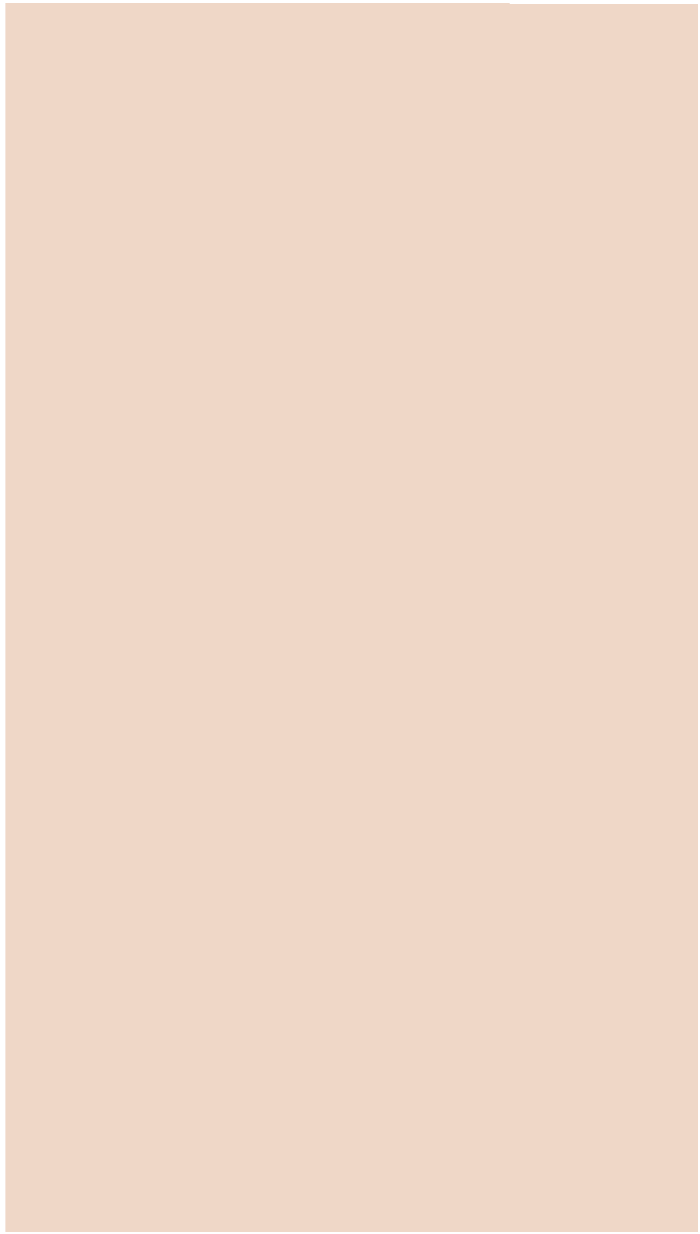
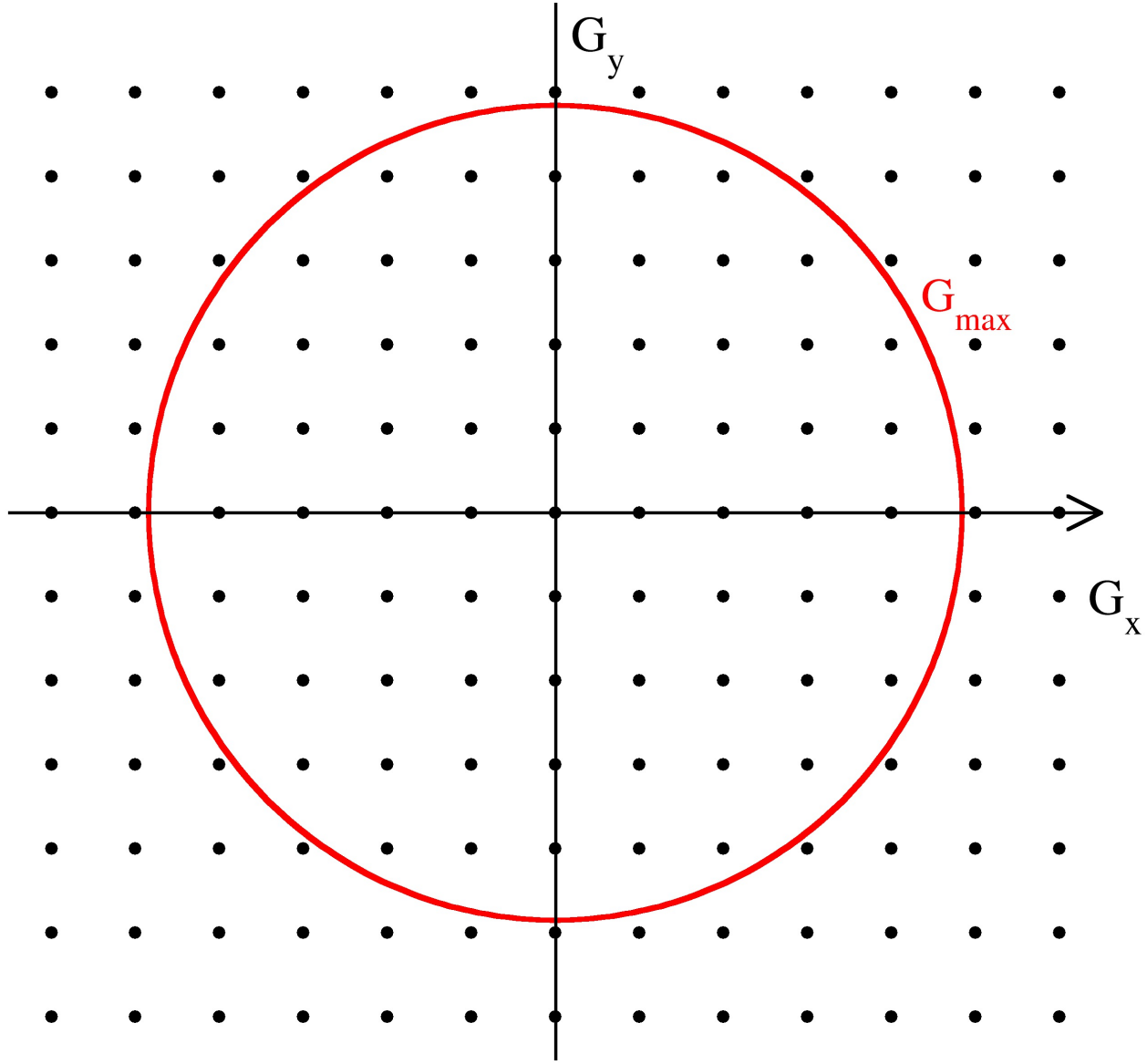
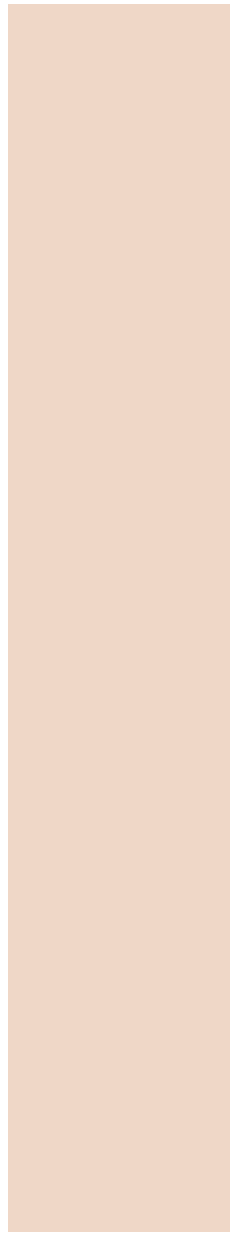
Bloch's theorem:  $\psi_{\mathbf{k},l}(\mathbf{r}) = \exp(i\mathbf{k} \cdot \mathbf{r}) \underbrace{u_{\mathbf{k},l}(\mathbf{r})}_{\text{periodic}}$

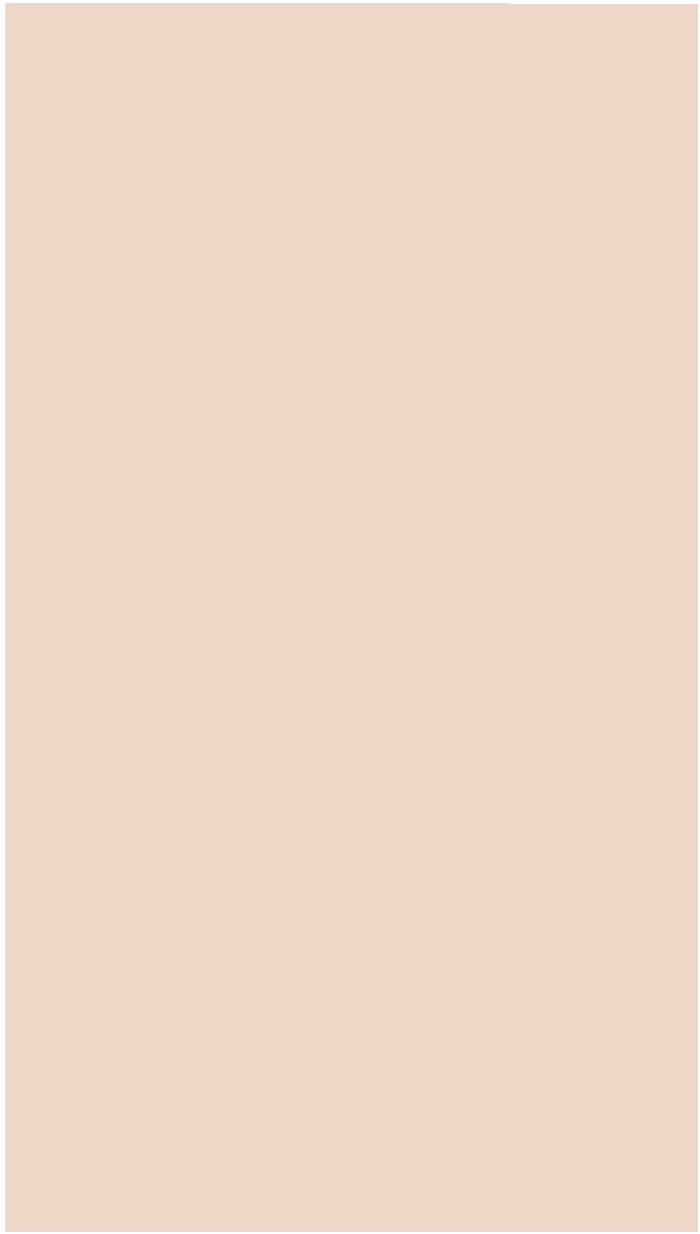
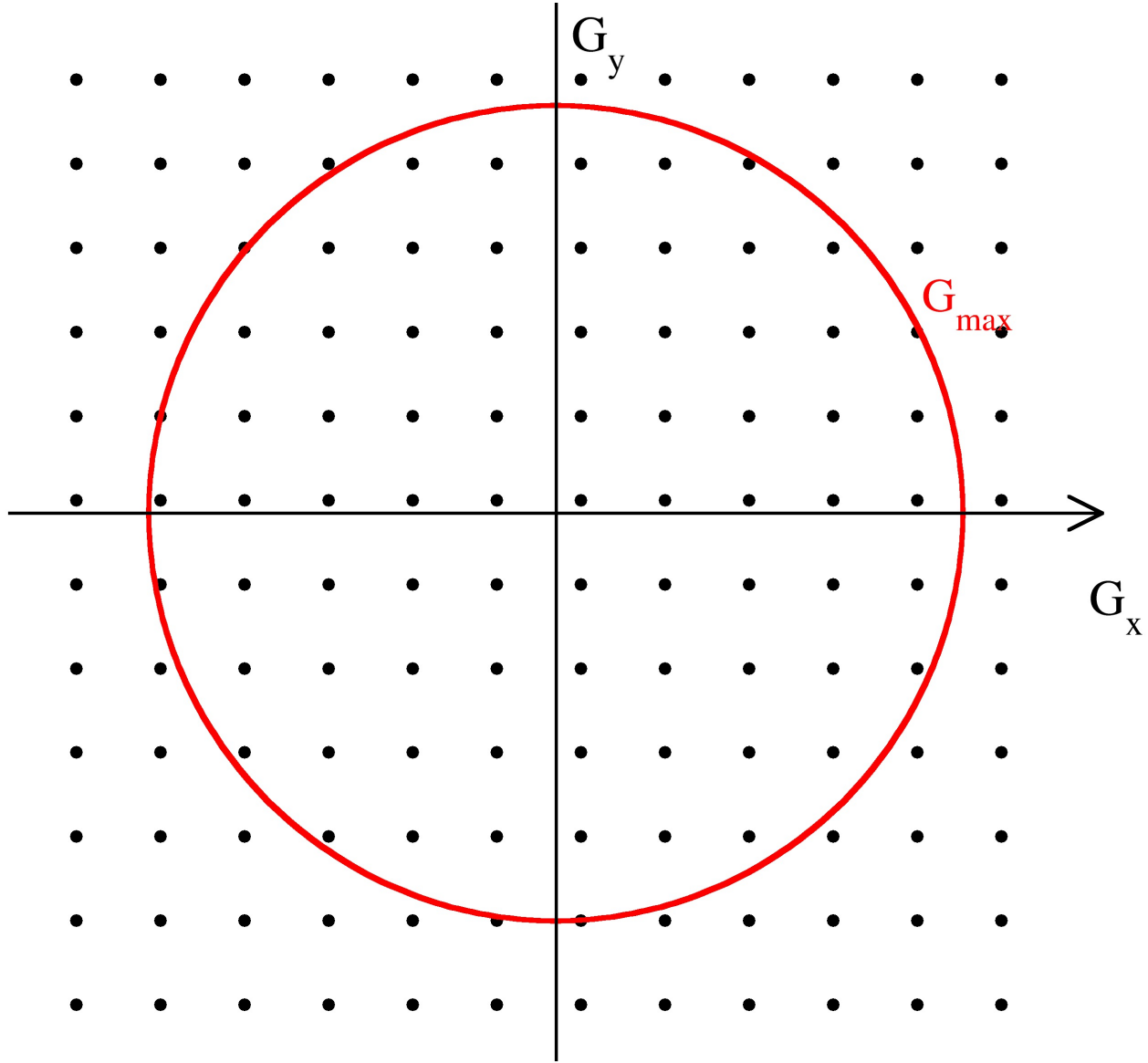
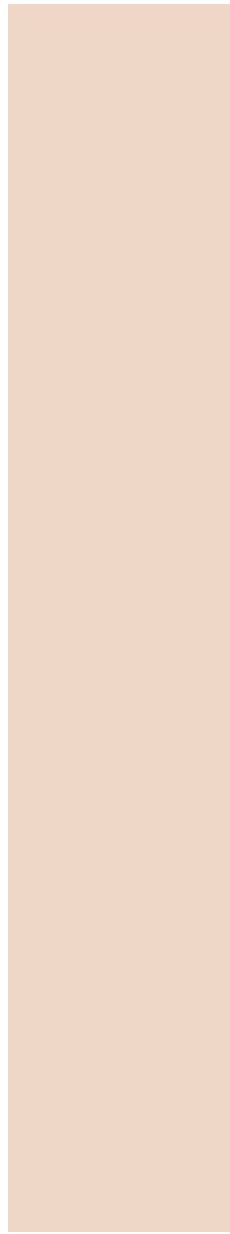
$$u_{\mathbf{k},l}(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{max}} c_{\mathbf{k},l} \exp(i\mathbf{G} \cdot \mathbf{r})$$

← grid of G-vectors like before

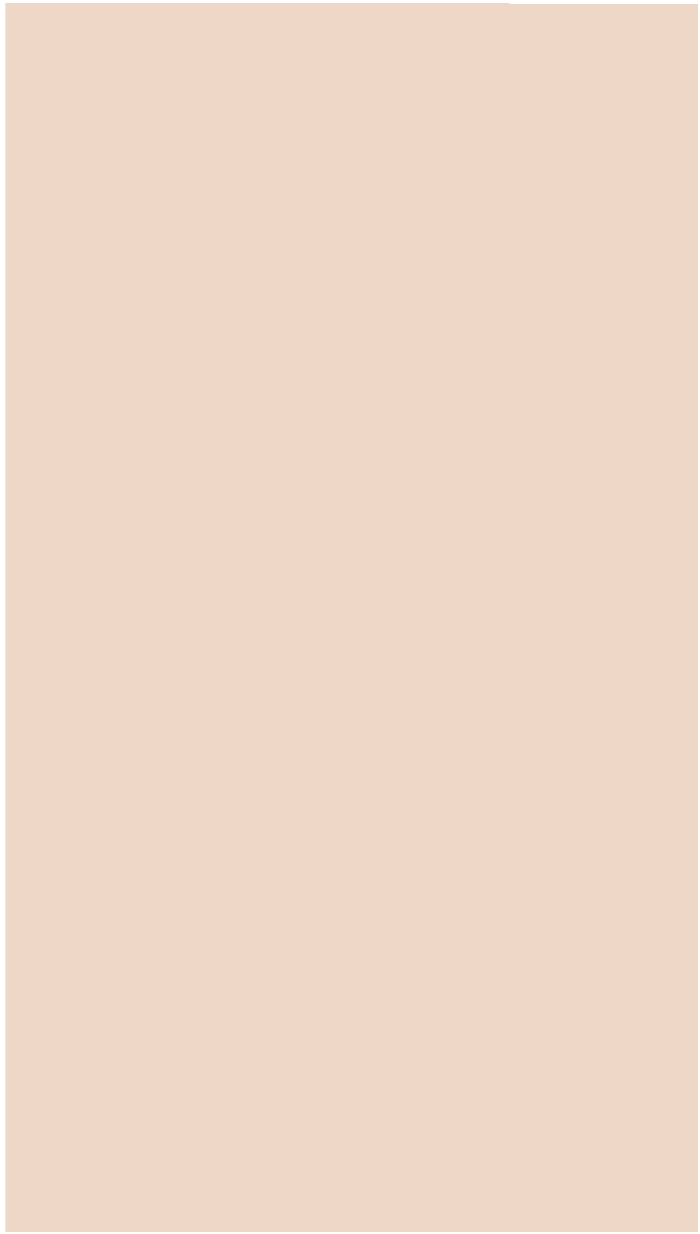
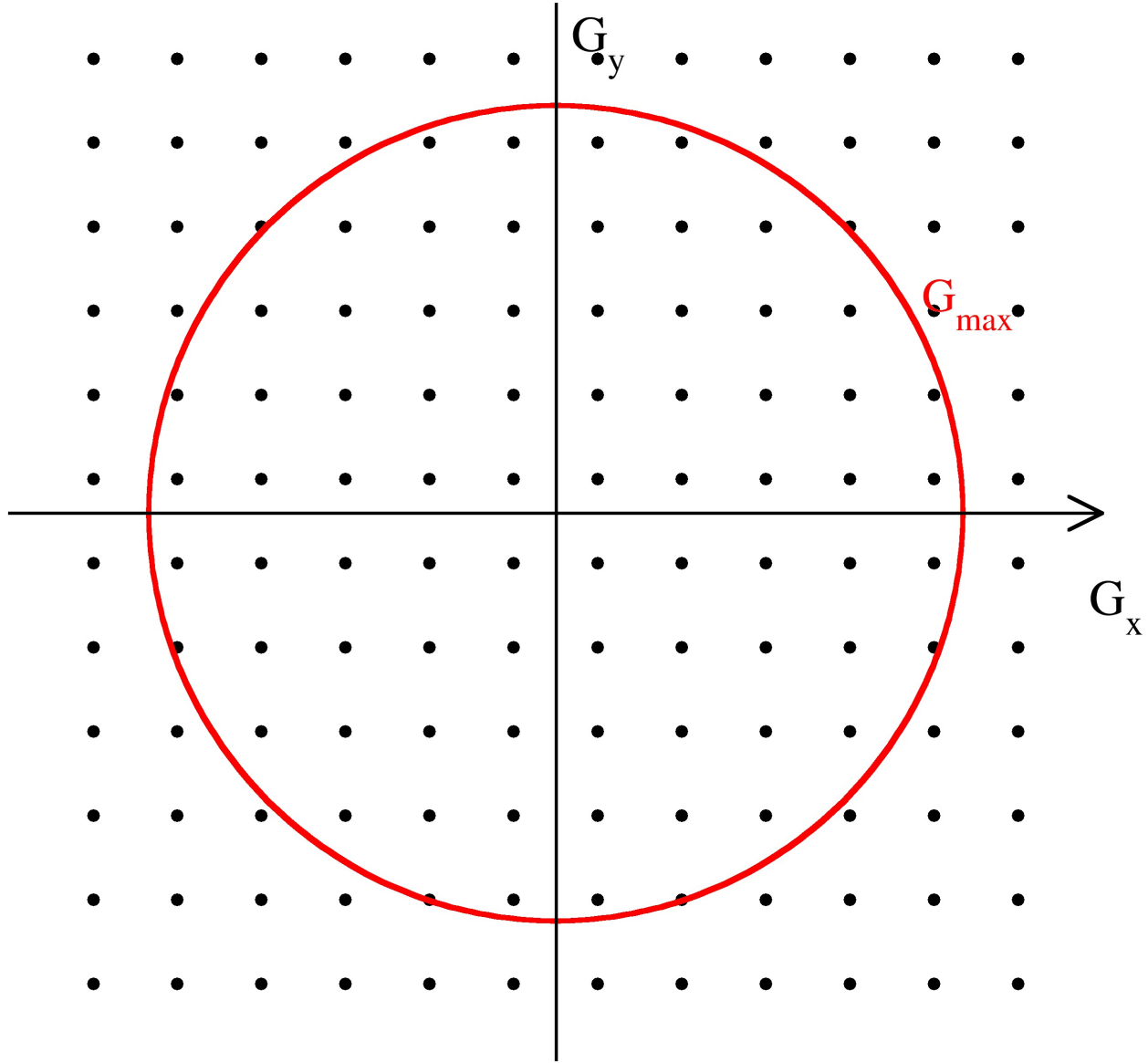
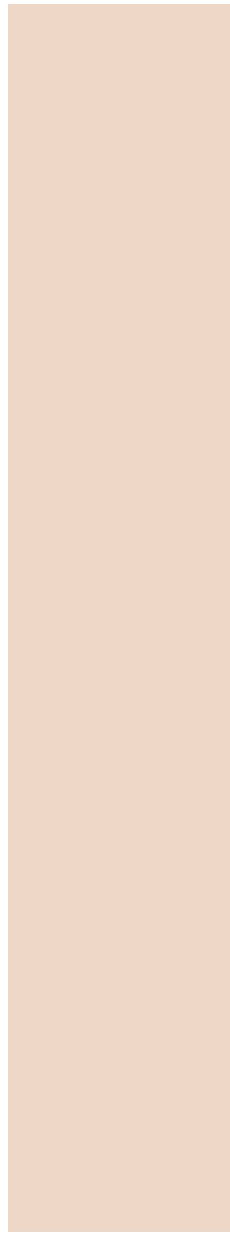
$$\psi_{\mathbf{k},l}(\mathbf{r}) = \sum_{|\mathbf{G}| < G_{max}} c_{\mathbf{k},l} \exp(i(\mathbf{k} + \mathbf{G}) \cdot \mathbf{r})$$

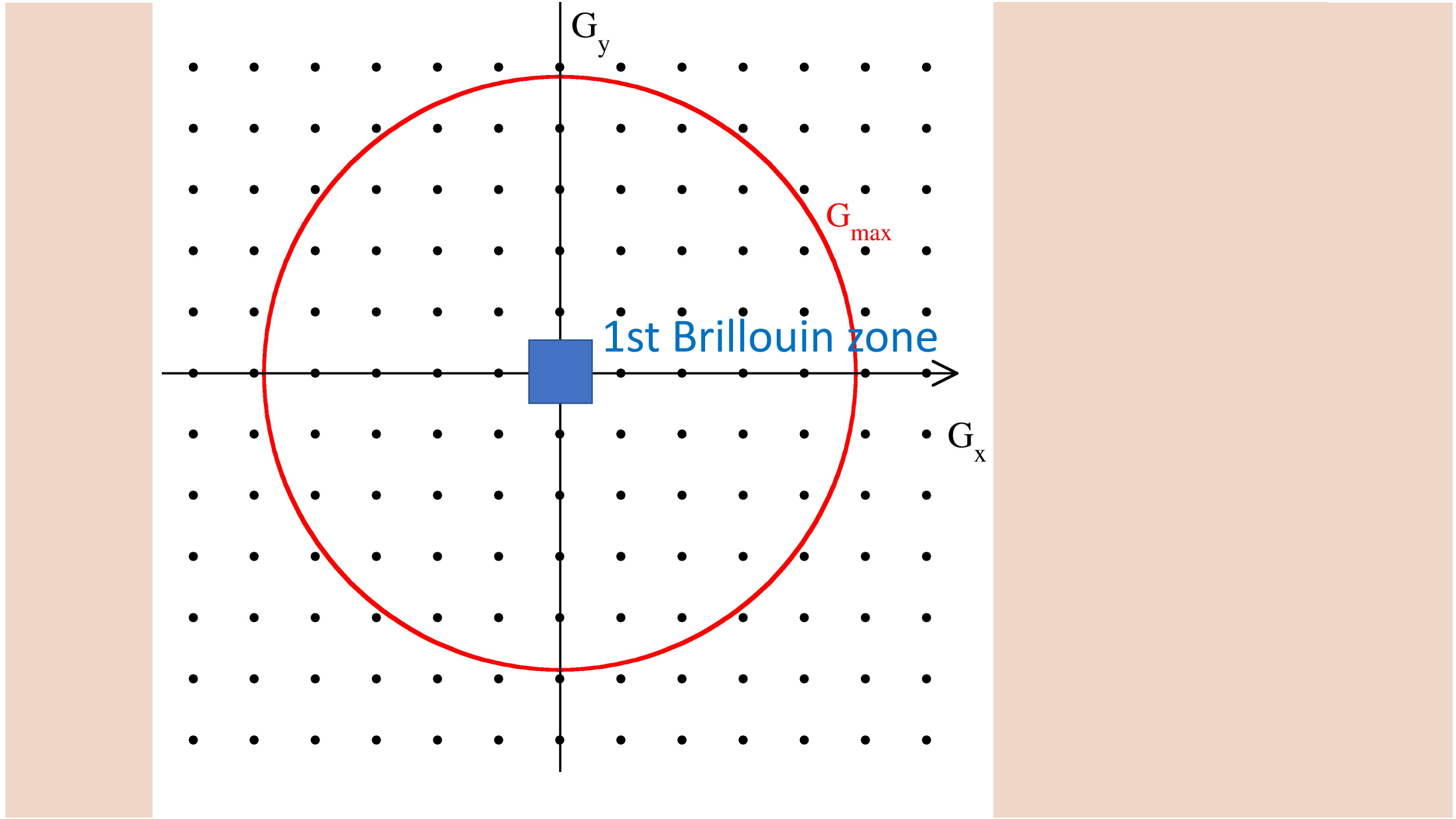
← shifted grid of (k+G)-vectors





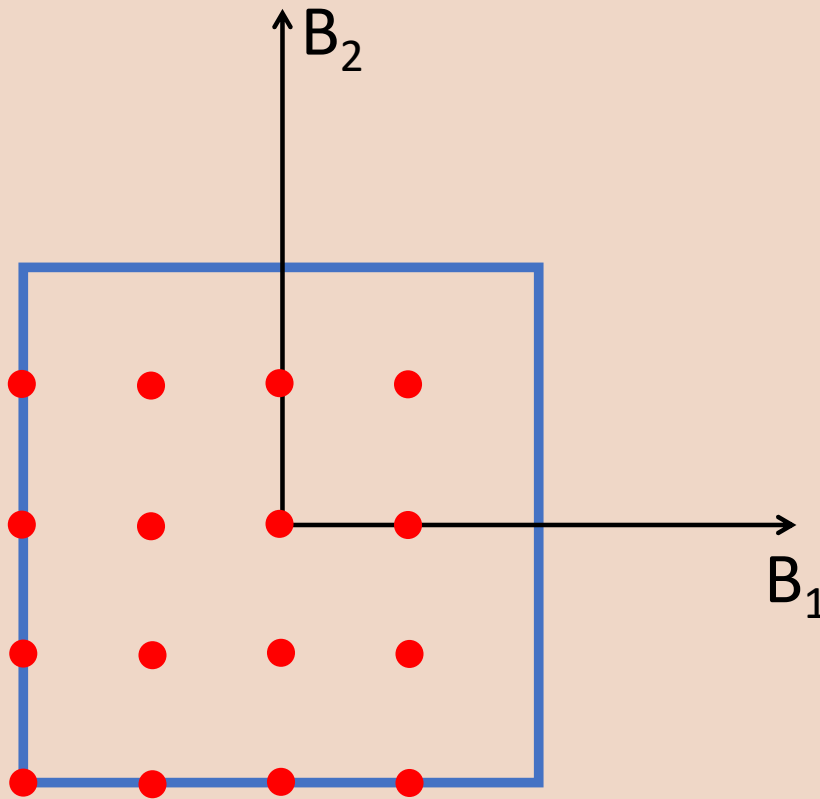
$G_{max}$





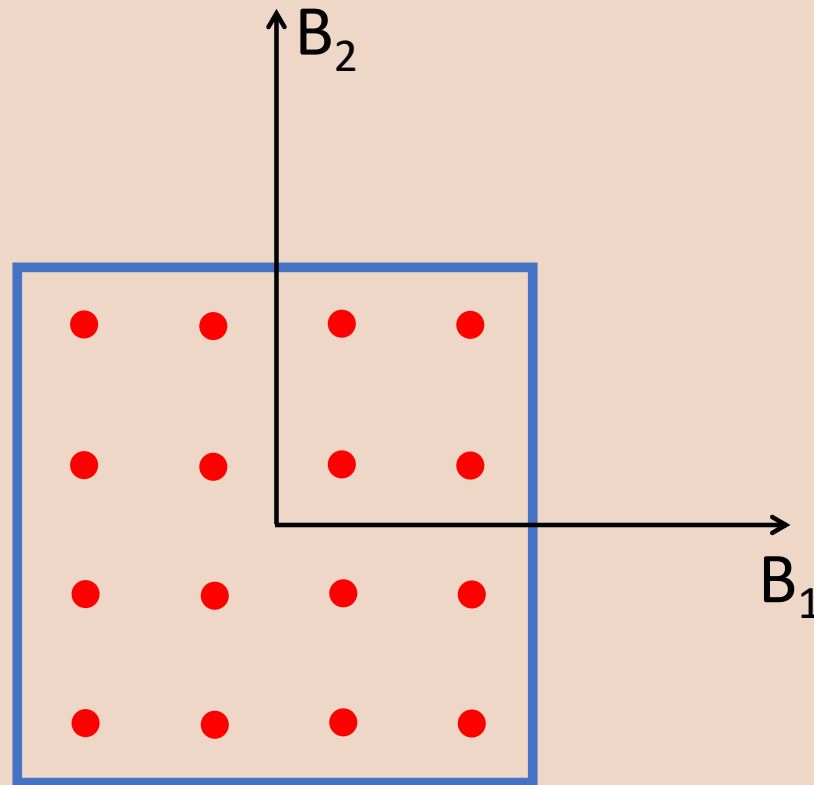


## Sampling the Brillouin zone



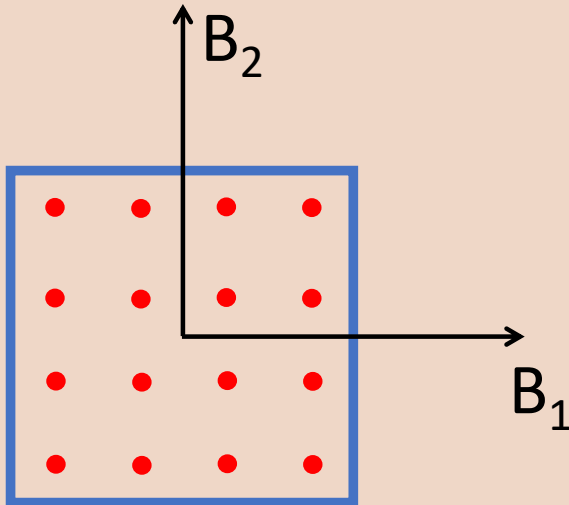
BZ sampled with a 4x4 k-point grid

## Sampling the Brillouin zone



BZ sampled with a **shifted** 4x4 k-point grid  
Shifted k-point grids usually converge faster (symmetry!)

## Sampling the Brillouin zone



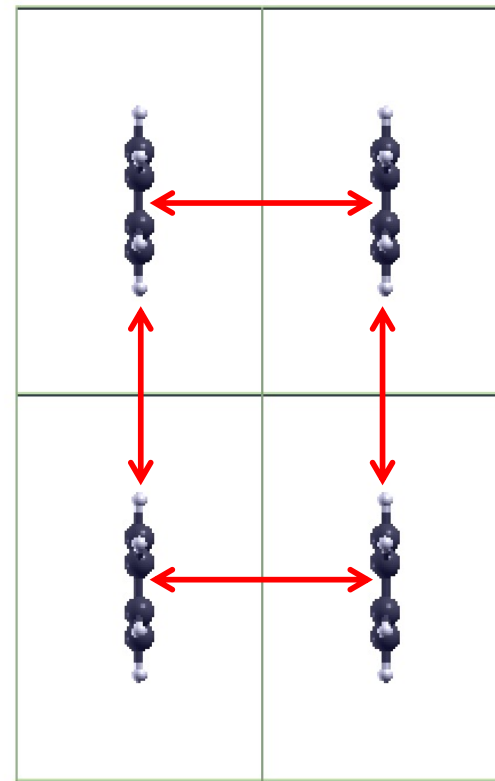
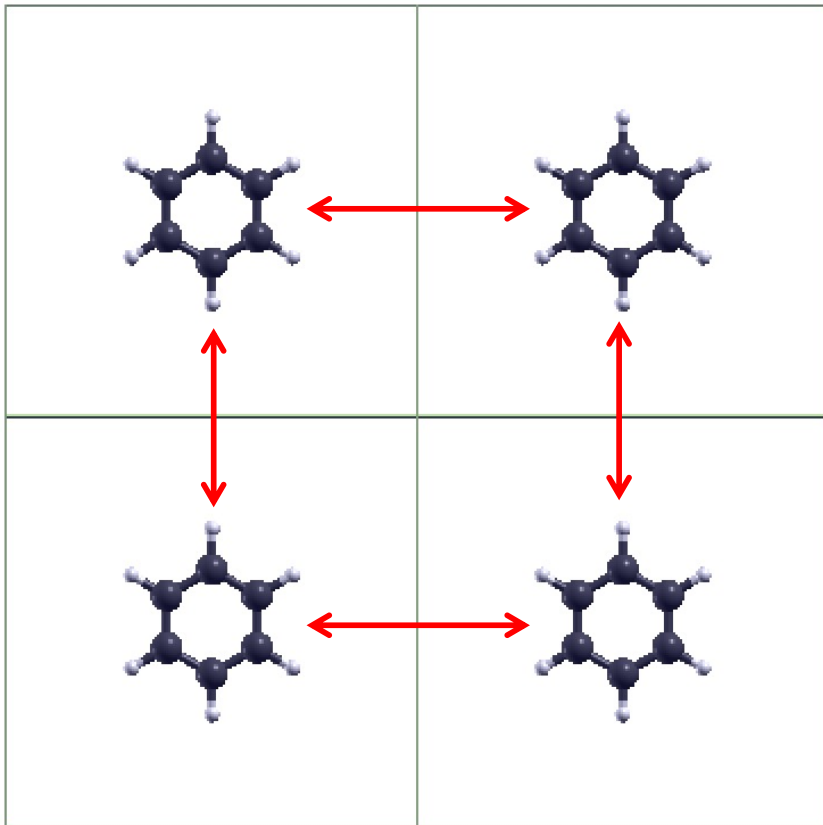
```
K_POINTS automatic
9 9 1 1 1 0
```

```
number of k points= 25
cart. coord. in units 2pi/alat
k( 1) = ( 0.0555556 0.0962250 0.0000000), wk = 0.0493827
k( 2) = ( 0.0555556 0.2245251 0.0000000), wk = 0.0987654
k( 3) = ( 0.0555556 0.3528252 0.0000000), wk = 0.0987654
k( 4) = ( 0.0555556 0.4811252 0.0000000), wk = 0.0987654
k( 5) = ( 0.0555556 -0.5452753 0.0000000), wk = 0.0987654
k( 6) = ( 0.0555556 -0.4169752 0.0000000), wk = 0.0987654
k( 7) = ( 0.0555556 -0.2886751 0.0000000), wk = 0.0987654
k( 8) = ( 0.0555556 -0.1603751 0.0000000), wk = 0.0987654
k( 9) = ( 0.0555556 -0.0320750 0.0000000), wk = 0.0493827
k(10) = ( 0.1666667 0.2886751 0.0000000), wk = 0.0493827
k(11) = ( 0.1666667 0.4169752 0.0000000), wk = 0.0987654
k(12) = ( 0.1666667 0.5452753 0.0000000), wk = 0.0987654
k(13) = ( 0.1666667 -0.4811252 0.0000000), wk = 0.0987654
k(14) = ( 0.1666667 -0.3528252 0.0000000), wk = 0.0987654
k(15) = ( 0.1666667 -0.2245251 0.0000000), wk = 0.0987654
k(16) = ( 0.1666667 -0.0962250 0.0000000), wk = 0.0493827
k(17) = ( 0.2777778 0.4811252 0.0000000), wk = 0.0493827
k(18) = ( 0.2777778 0.6094253 0.0000000), wk = 0.0987654
k(19) = ( 0.2777778 -0.4169752 0.0000000), wk = 0.0987654
k(20) = ( 0.2777778 -0.2886751 0.0000000), wk = 0.0987654
k(21) = ( 0.2777778 -0.1603751 0.0000000), wk = 0.0493827
k(22) = ( 0.3888889 0.6735753 0.0000000), wk = 0.0493827
k(23) = ( 0.3888889 -0.3528252 0.0000000), wk = 0.0987654
k(24) = ( 0.3888889 -0.2245251 0.0000000), wk = 0.0493827
k(25) = ( -0.5000000 -0.8660254 0.0000000), wk = 0.0246914
```

```
Dense grid: 3243 G-vectors FFT dimensions: ( 15, 15, 40)
```

## What if a system is NOT 3D-periodic ?

Use a **supercell**: introduce enough space between periodic replicas that the interaction is negligible



## What if a system is NOT 3D-periodic ?

Use a **supercell**: introduce enough space between periodic replicas that the interaction is negligible

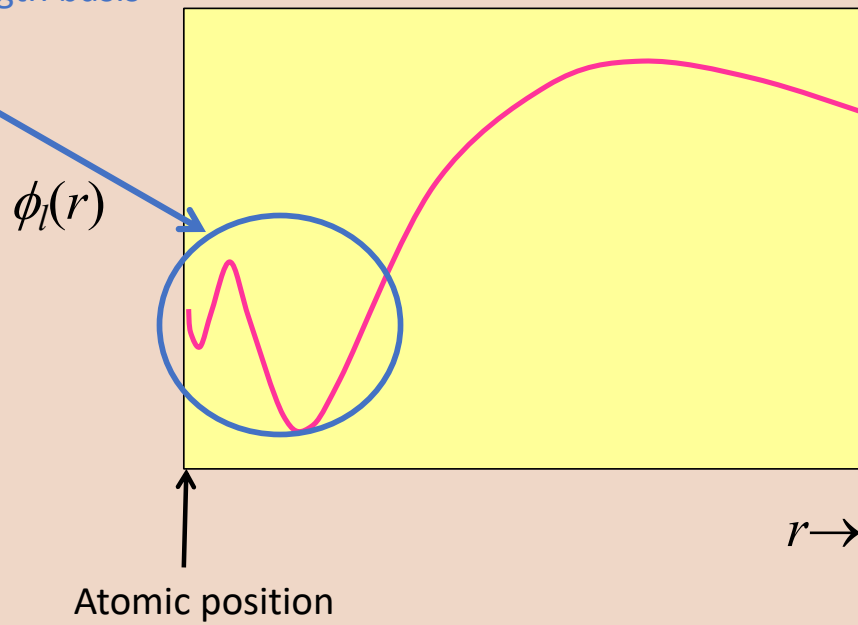
Can be problematic if the object is charged or has a dipole!

Solution:

- input variable **assume\_isolated**: useful in 0-D (molecules, clusters, etc)
- introduce a dipole layer in the vacuum between surfaces (**dipfield**)

## The need for pseudopotentials

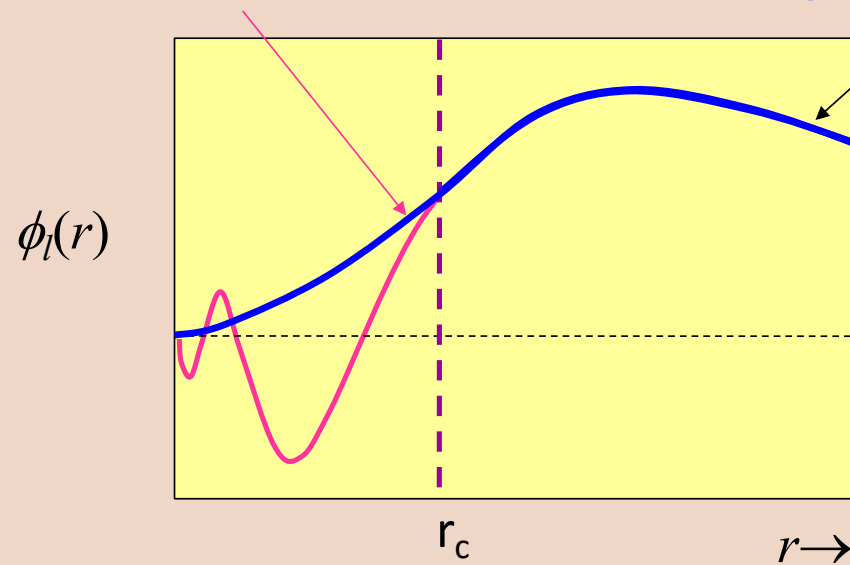
rapid oscillations would need  
a very small wavelength basis  
functions



# Pseudopotentials

all-electron wavefunction

pseudowavefunction

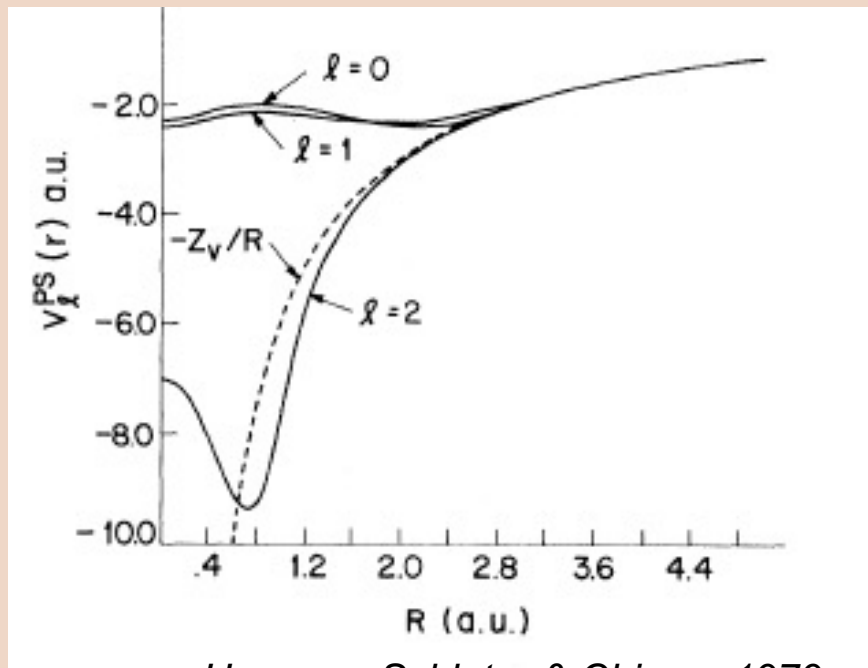


Necessary condition for the pseudopotential to be transferable: **NORM-CONSERVATION**

$$\int_0^{r_c} \phi^{*AE}(r) \phi^{AE}(r) dr = \int_0^{r_c} \phi^{*PS}(r) \phi^{PS}(r) dr$$

## What does a pseudopotential look like?

Example for Mo:



*Hamann, Schluter & Chiang, 1979.*

- Weaker than full Coulomb potential
- No singularity at  $r=0$
- Different pseudopotential for each  $l$  (example of semilocal pseudopotential)
- Will be  $V_{ion}$  (replacing nuclear potential)



## Pseudopotentials in quantum-ESPRESSO:

PPs in QE are provided by the user in the form of files, which contain the PP on a radial grid:  
in the PW input file, the PP filenames are specified for every element in the system

```
ATOMIC_SPECIES
Mo 95.96 Mo.pbe-spn-rrkjus_psl.1.0.0.UPF
W 183.8 W.pbe-spn-rrkjus_psl.1.0.0.UPF
S 32.07 S.pbe-n-rrkjus_psl.1.0.0.UPF
O 16.00 O.pbe-n-rrkjus_psl.1.0.0.UPF
H 1.00 H.pbe-rrkjus_psl.1.0.0.UPF
```

PPs are DFT-functional specific. The code will by default apply the functional to your system with which the PP has been calculated

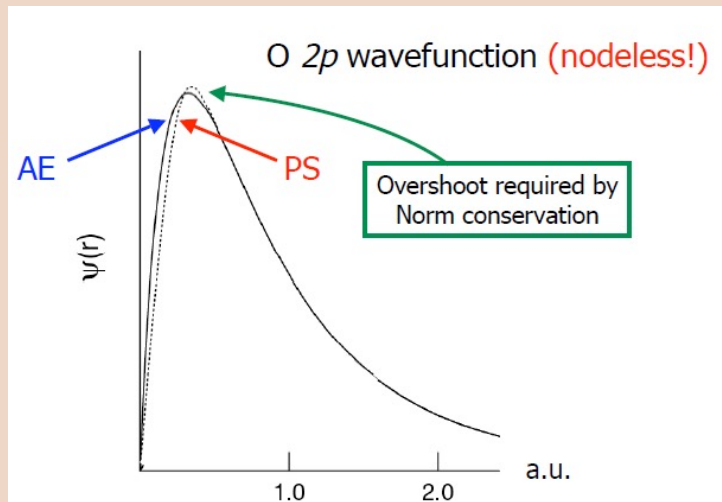
It is generally a bad idea to mix pseudopotentials generated with different functionals!

## Beyond norm-conservation:

Cut-offs still higher than we would like, especially for

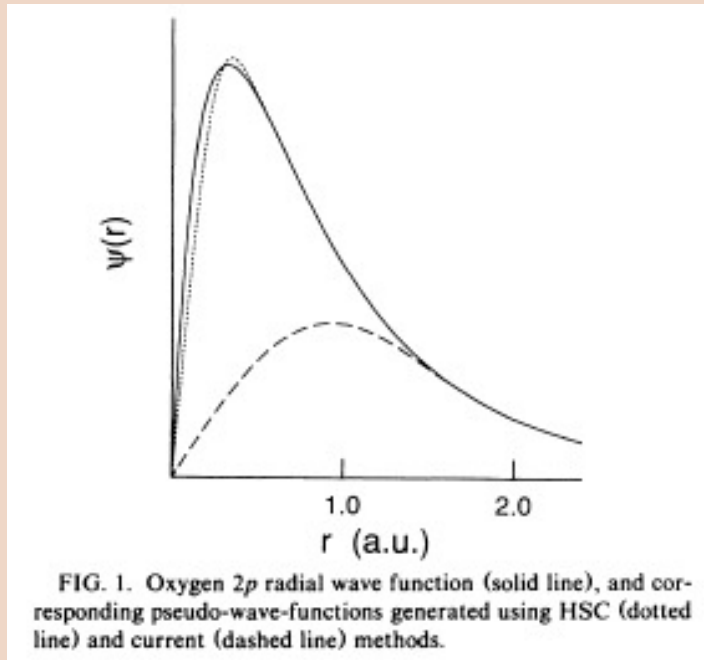
- > first row elements (1s, 2p nodeless)
- > transition metals (3d nodeless)
- > rare-earths (4f nodeless)

This is because of the constraint of norm conservation...



## Beyond norm-conservation:

Vanderbilt's ultrasoft pseudopotentials:



Remove the constraint of norm-conservation

The price to be paid: a more complicated expression for the electronic charge density:

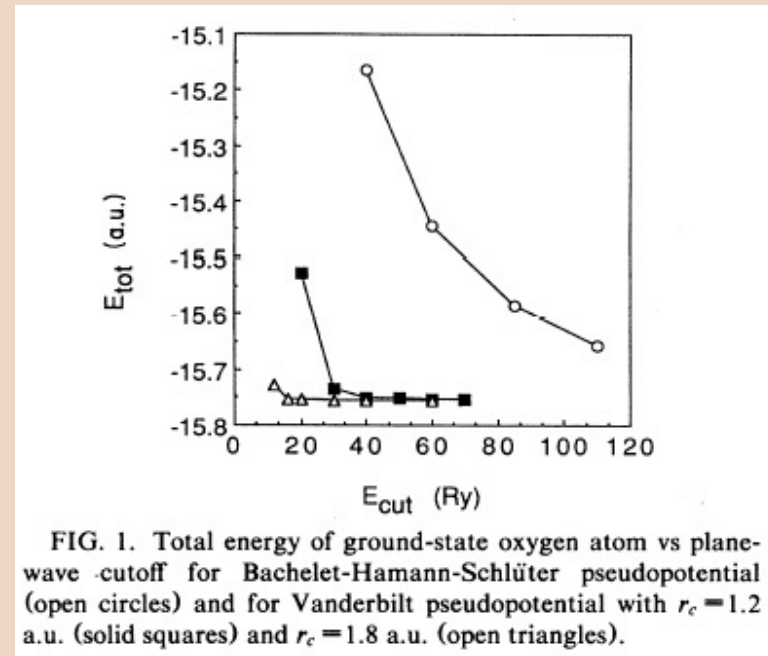
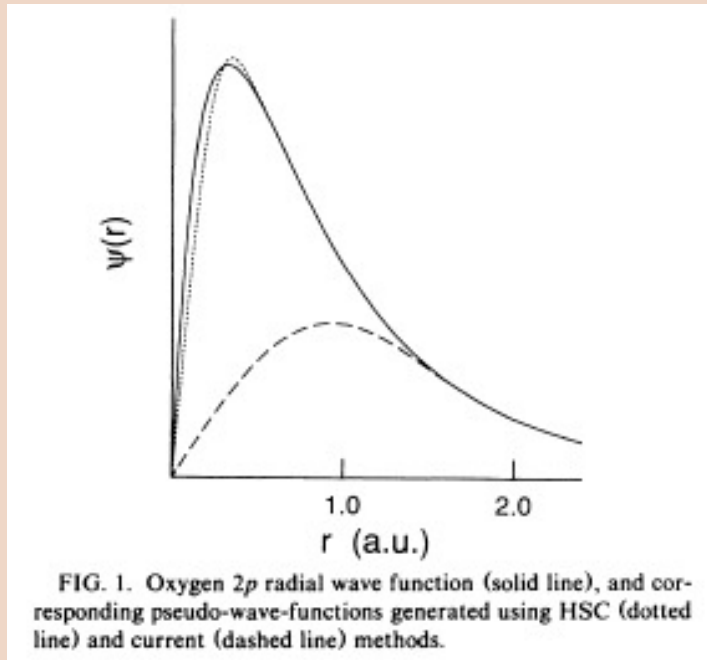
$$n(\mathbf{r}) = \sum_{i=1}^N \psi_i^*(\mathbf{r})\psi_i(\mathbf{r}) + n^{aug}(\mathbf{r})$$

Augmentation charges, localized close to the nucleus account for “missing norm” expressed using projection functions

NB: Ultrasoft (US) and PAW pseudopotentials are (from the user's perspective) very similar

## Beyond norm-conservation:

Vanderbilt's ultrasoft pseudopotentials:



## As a user: BEWARE!

Using US or PAW pseudopotentials, you need a cutoff energy  $E_{\text{cut}}$  for the ccharge density which is  
**MORE THAN  $4 \cdot E_{\text{cut}}$  for the orbitals**

### ATOMIC\_SPECIES

```
Mo 95.96 Mo.pbe-spn-rrkjus_psl.1.0.0.UPF
W 183.8 W.pbe-spn-rrkjus_psl.1.0.0.UPF
S 32.07 S.pbe-n-rrkjus_psl.1.0.0.UPF
O 16.00 O.pbe-n-rrkjus_psl.1.0.0.UPF
H 1.00 H.pbe-rrkjus_psl.1.0.0.UPF
```

Employing US or PAW pseudopotentials and NOT specifying `ecutrho` is the most common beginner's mistake in QE!

```
&system
 ibrav=0,
 nat=78,
 ntyp=5,
 ecutwfc = 40. ,
 ecutrho = 400. ,
 nspin=2,
 starting_magnetization(1) = 0.
 starting_magnetization(2) = 0.1
 starting_magnetization(3) = 0.
 starting_magnetization(4) = 0.2
 starting_magnetization(5) = 0.2
 vdw_corr='grimme-d3'
 occupations='smearing'
 smearing='cold'
 degauss=0.001
/
```

HAPPY QE-COMPUTING!