

JORGE MARIO ESCOBAR, BRAYAN FERNANDO DÍAZ, JESÚS MARÍA CALERO

PHOTONIC CRYSTALS

WITH

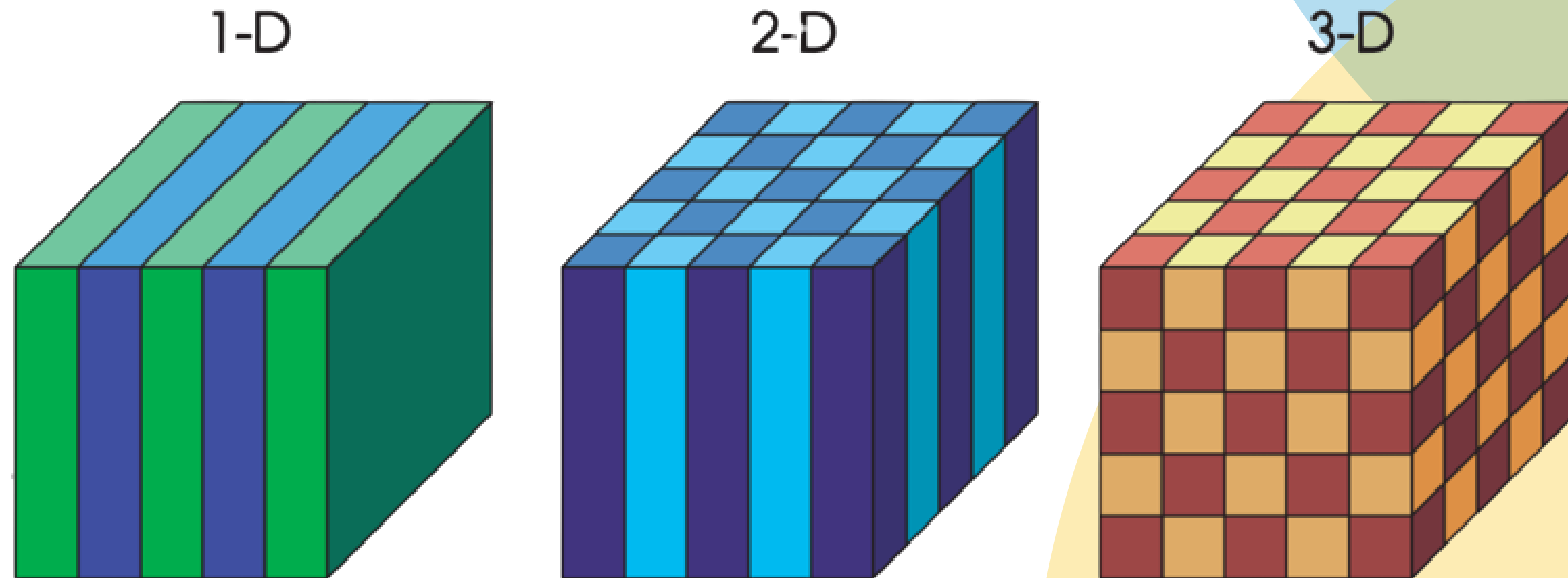
SUPERCONDUCTORS

A STUDY OF THEIR REFLECTANCE IN ANNULAR TERNARY SYSTEMS

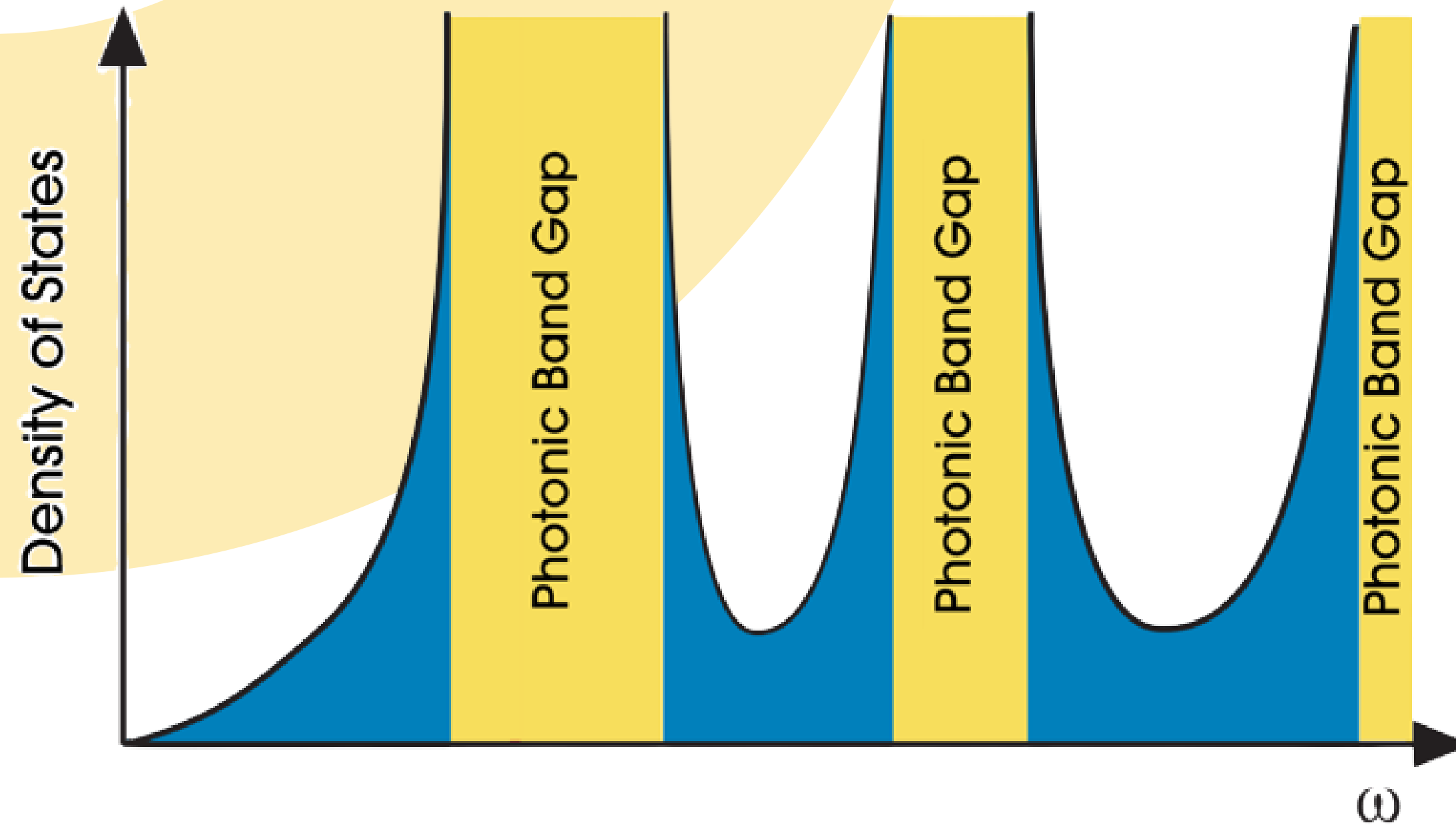


THEORETICAL SOLID STATE PHYSICS GROUP
DEPARTMENT OF PHYSICS
UNIVERSIDAD DEL VALLE
COLOMBIA

PHOTONIC CRYSTALS



PHOTONIC BAND GAPS



**THESE PHOTONIC BAND GAPS
ALLOW US TO DESIGN FILTERS
AND REFLECTORS FOR
ELECTROMAGNETIC WAVES.**

J. D. Joannopoulos, *Photonic Crystals*, Princeton University Press, (2008).

M. S. Chen *et al.*, *Solid State Communications*, 149, 1888–1893 (2009).

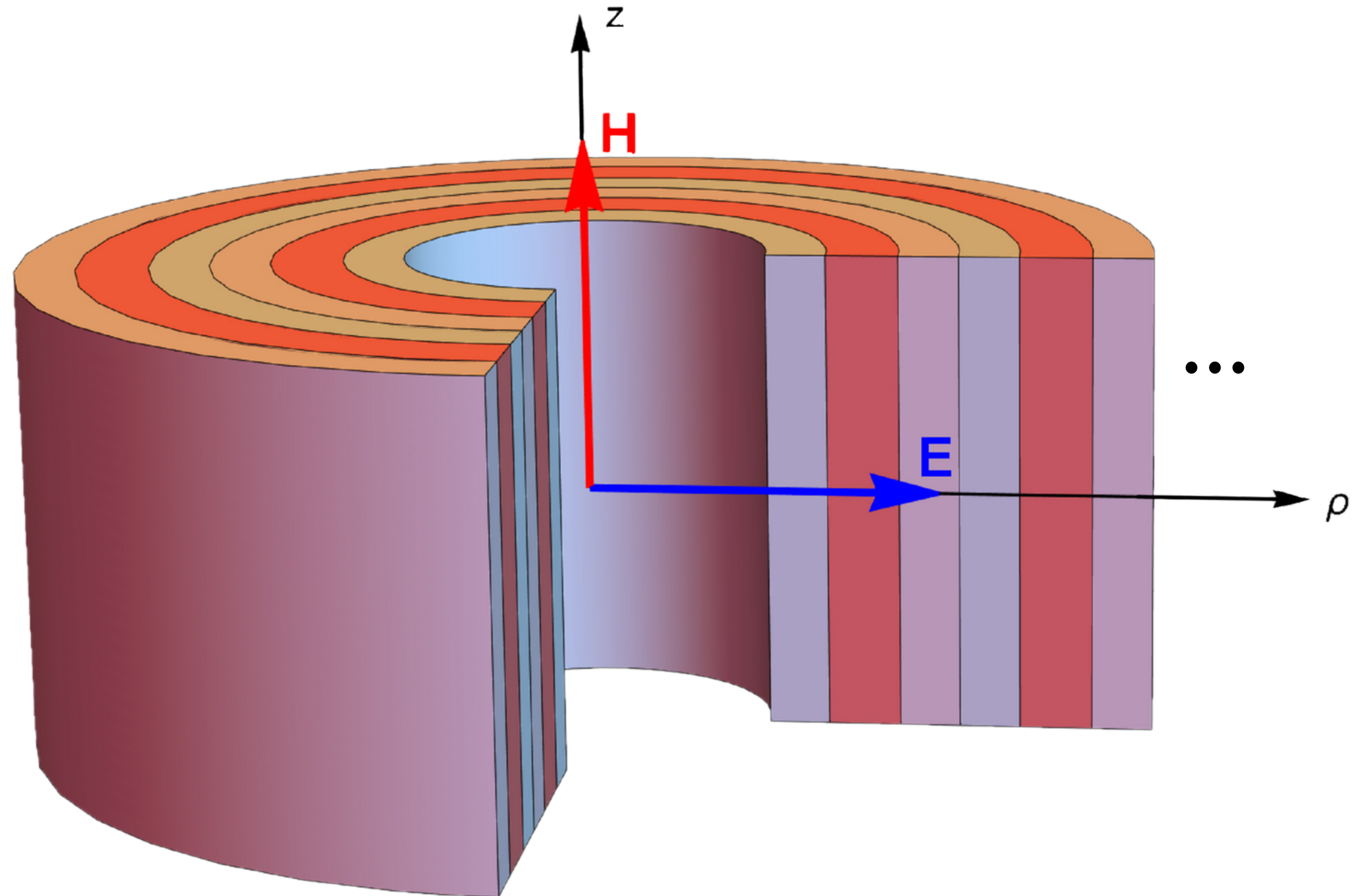
V. V. Nikolaev *et al.*, *Fiz. Tekh. Poluprovodn.*, 33, 174–179 (1999).

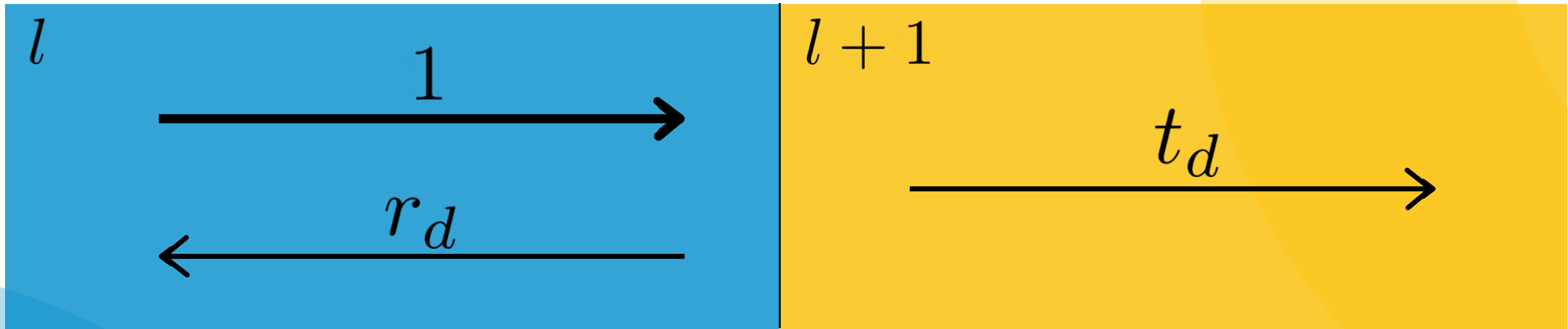
ANNULAR TERNARY SYSTEM

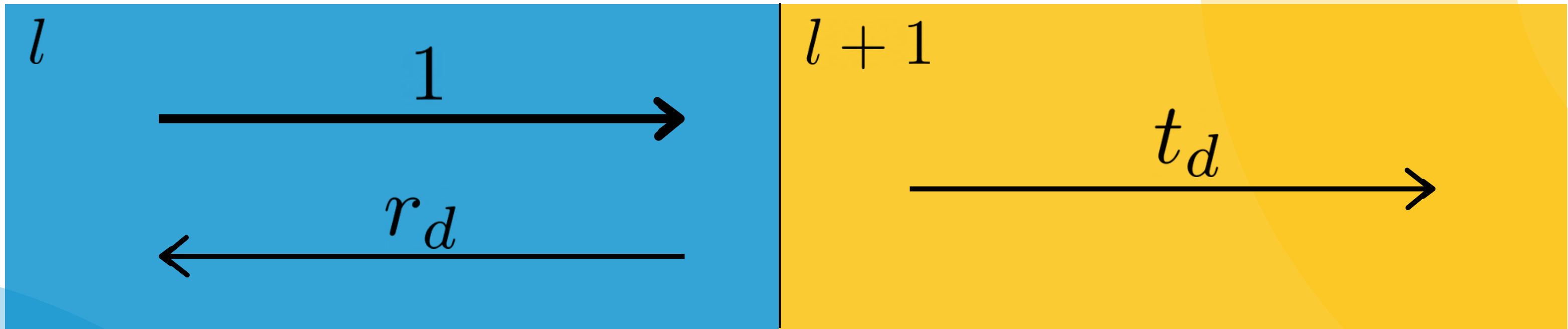
WE WILL ANALYZE THE
OPTICAL RESPONSE OF THE
SYSTEM IN TM
POLARIZATION.

$$\mathbf{H} = H_z \hat{\mathbf{e}}_z$$

$$\mathbf{E} = E_\rho \hat{\mathbf{e}}_\rho + E_\phi \hat{\mathbf{e}}_\phi$$

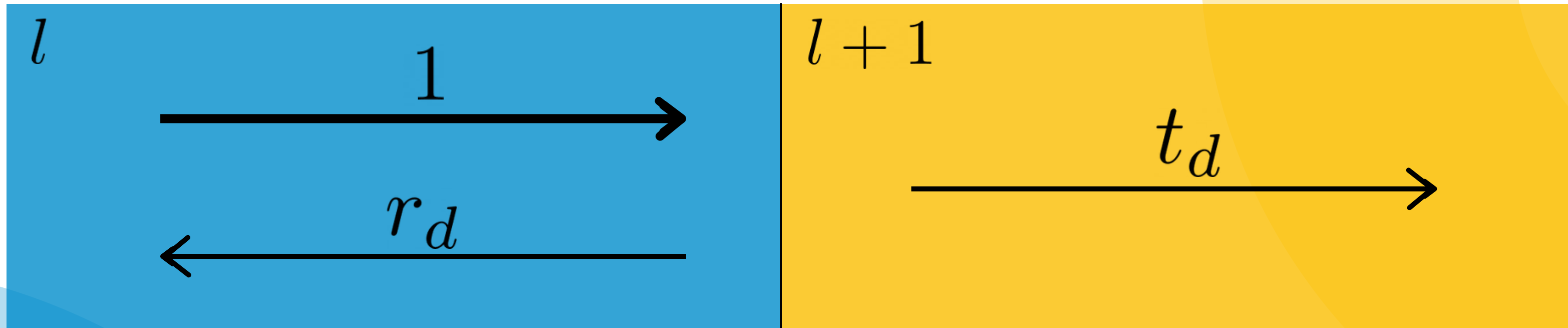






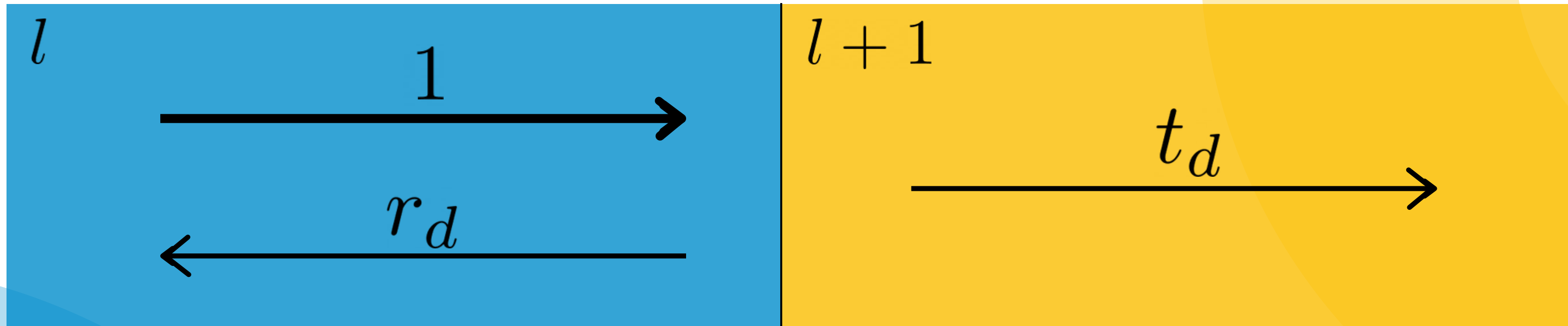
$$\begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_l \longrightarrow \begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_{l+1}$$

TRANSFER MATRIX METHOD



$$\hat{\mathcal{M}}_l \begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_l = \begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_{l+1}$$

TRANSFER MATRIX METHOD



$$\hat{\mathfrak{M}}_l \begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_l = \begin{pmatrix} \text{EM} \\ \text{wave} \end{pmatrix}_{l+1}$$

$$t_d = \frac{|\hat{\mathfrak{M}}_l|}{\mathfrak{M}_l^{22}} \quad \text{y} \quad r_d = -\frac{\mathfrak{M}_l^{21}}{\mathfrak{M}_l^{22}} \quad \longrightarrow \quad R_d = \left| \frac{\mathfrak{M}_l^{21}}{\mathfrak{M}_l^{22}} \right|^2, \quad \text{con } R_d = \boxed{R_d(\lambda)}$$

M. Born, E. Wolf, *Principles of Optics*, P. 77, Moscow, (1970).

P. Markos, C. M. Soukoulis, *Wave Propagation*, Princeton University Press, (2008).

THEORETICAL TREATMENT

MASTER EQUATION

$$\nabla \times \left[\frac{1}{\epsilon(\mathbf{r}, \omega)} \nabla \times \mathbf{H}(\mathbf{r}) \right] = \left(\frac{\omega}{c} \right)^2 \mathbf{H}(\mathbf{r})$$

THEORETICAL TREATMENT

MASTER EQUATION

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Description in
cylindrical coordinates

BESSEL PARTIAL DIFFERENTIAL EQUATION

$$\left[\frac{1}{\rho^2} \frac{\partial}{\partial \phi} \left(\frac{1}{\epsilon(\mathbf{r}, \omega)} \frac{\partial}{\partial \phi} \right) + \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{1}{\epsilon(\mathbf{r}, \omega)} \frac{\partial}{\partial \rho} \right) \right] H_z(\mathbf{r}) = - \left(\frac{\omega}{c} \right)^2 H_z(\mathbf{r})$$

J. D. Joannopoulos, *Photonic Crystals*, Princeton University Press, (2008).

C. A. Hu et al., *Optics Communications*, 291, 424–434, (2013).

THEORETICAL TREATMENT

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Solutions in traveling
wave basis

$$\begin{cases} H_z^\pm = A H_m^{(1/2)}(k\rho) e^{im\phi} \\ E_\phi^\pm = ipA \frac{\partial}{\partial(k\rho)} H_m^{(1/2)}(k\rho) e^{im\phi} \end{cases}$$

THEORETICAL TREATMENT

MASTER EQUATION

$$\nabla \times \left[\frac{1}{\epsilon(\mathbf{r}, \omega)} \nabla \times \mathbf{H}(\mathbf{r}) \right] = \left(\frac{\omega}{c} \right)^2 \mathbf{H}(\mathbf{r})$$

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cylindrical coordinates

BESSEL PARTIAL DIFFERENTIAL EQUATION

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TRANSFER MATRIX

$$\hat{\mathcal{M}}_l = \begin{pmatrix} H_m^{(1)}(k_l \rho_l) & H_m^{(2)}(k_l \rho_l) \\ \frac{1}{k_l} H_m'^{(1)}(k_l \rho_l) & \frac{1}{k_l} H_m'^{(2)}(k_l \rho_l) \end{pmatrix}$$

Apply boundary conditions
to build the matrix

J. D. Joannopoulos, *Photonic Crystals*, Princeton University Press, (2008).

C. A. Hu et al., *Optics Communications*, 291, 424–434, (2013).

SUPERCONDUCTOR REFRACTIVE INDEX

$$n(\lambda, T) = \sqrt{1 - \left(\frac{\lambda}{2\pi\lambda_L(T)} \right)^2}$$

TWO FLUIDS MODEL

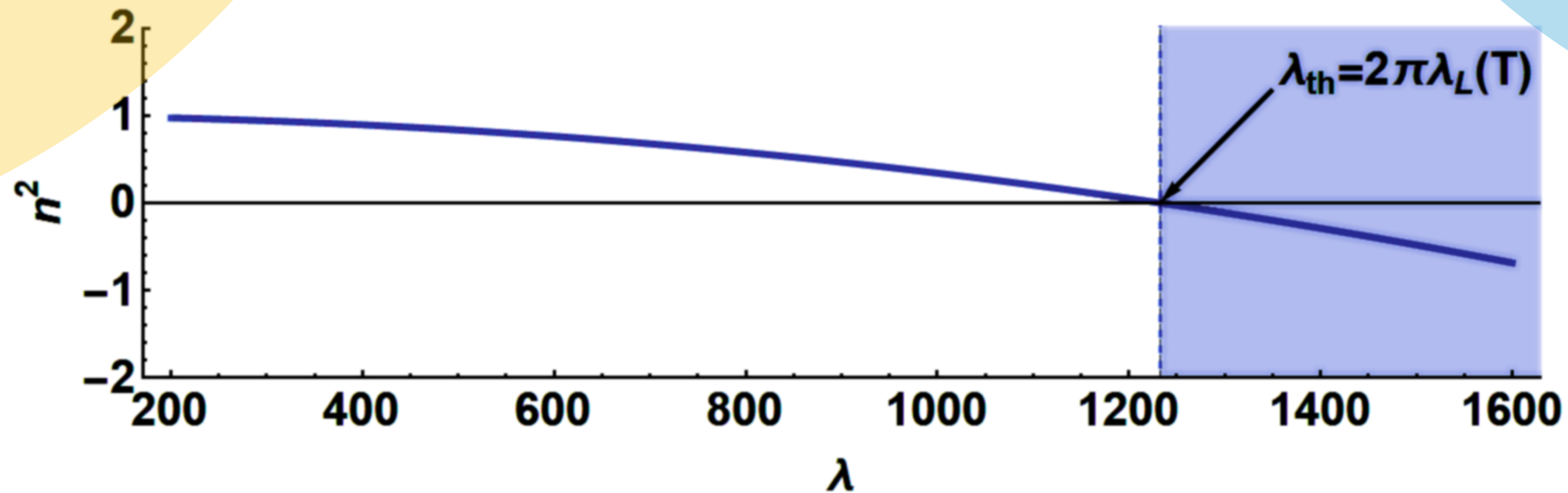
LONDON PENETRATION DEPTH

$$\lambda_L(T) = \frac{\lambda_0}{\sqrt{1 - \left(\frac{T}{T_c} \right)^4}}$$

SUPERCONDUCTOR REFRACTIVE INDEX

$$n(\lambda, T) = \sqrt{1 - \left(\frac{\lambda}{2\pi\lambda_L(T)} \right)^2}$$

AT FIXED TEMPERATURE



K. P. Sreejith, V. Mathew, *Journal of Superconductivity and Novel Magnetis*, 32:2397, (2019).

C. P. Poole Jr. *et al*, *Superconductivity*, Elsevier (2007).

OUR SYSTEM

Superconductor



(YBCO)

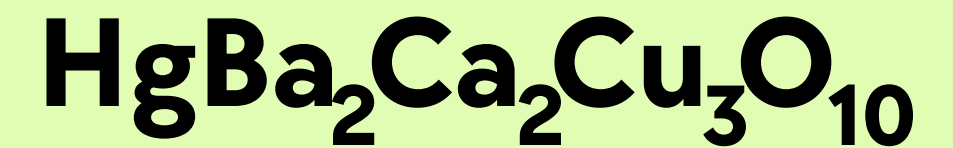
$T_c=92\text{K}$

Dielectric

SrTiO

$n_D=2.437$

Superconductor

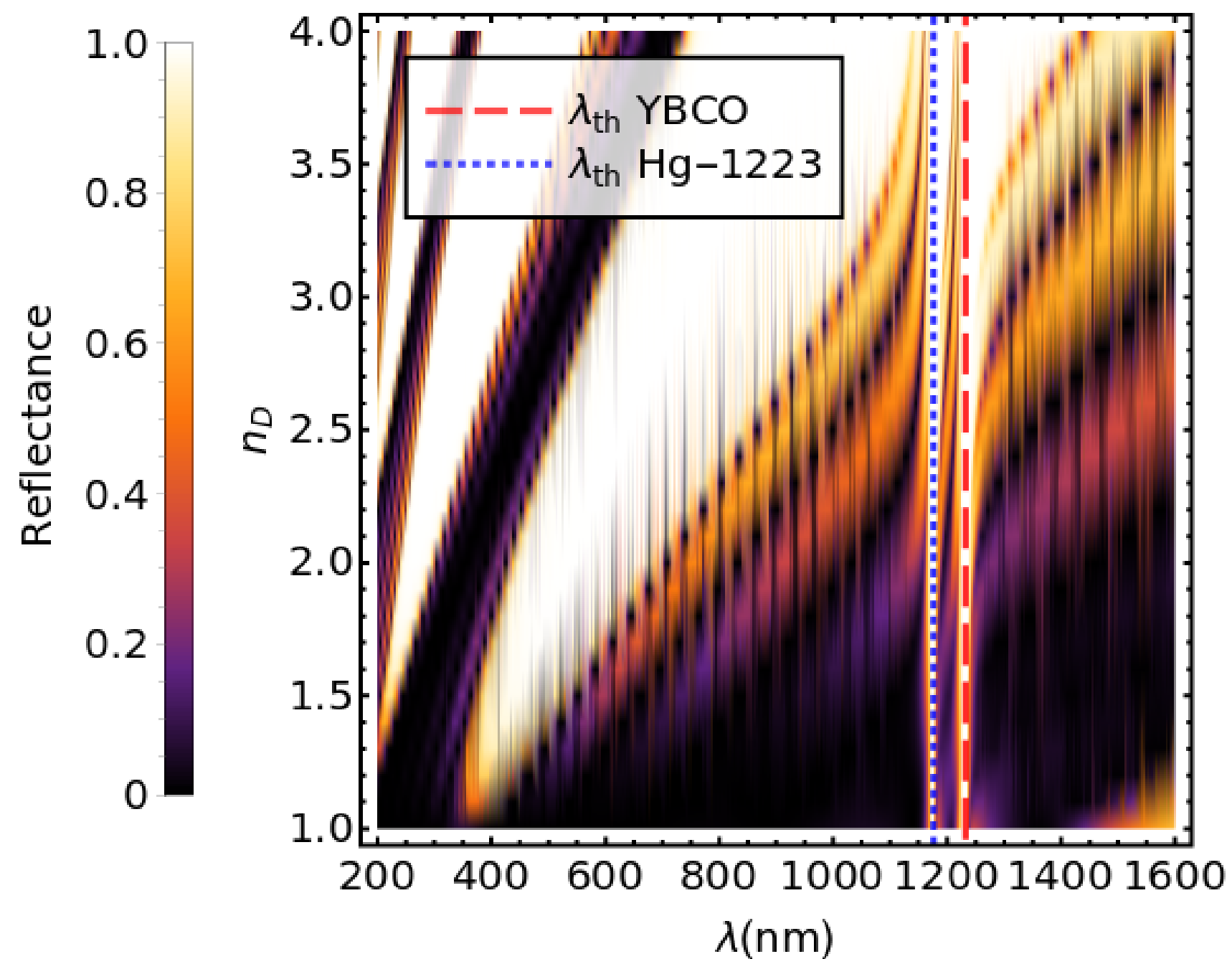


(Hg-1223)

$T_c=135\text{K}$

RESULTS

Dielectric refractive index variation



$$T = 77K$$

Superconductor refractive index

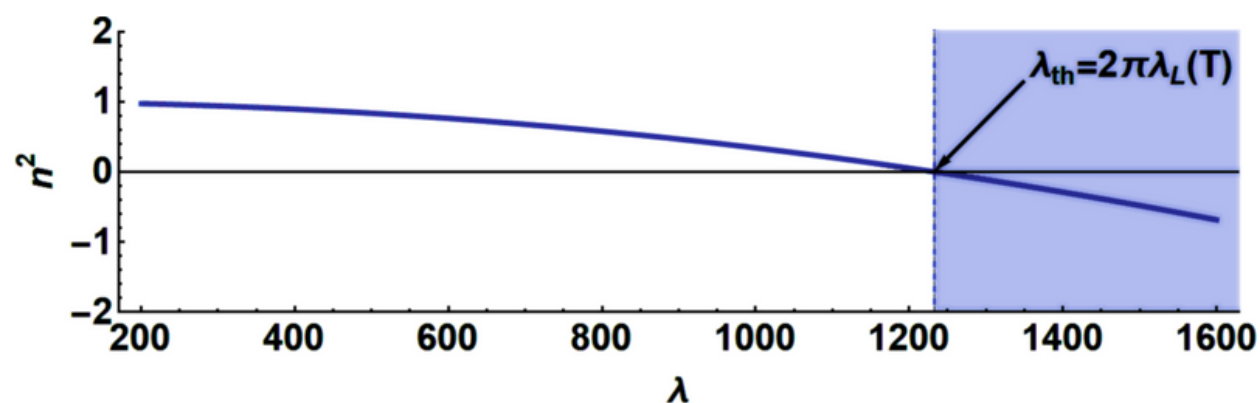
$$n(\lambda, T) = \sqrt{1 - \left(\frac{\lambda}{2\pi\lambda_L(T)} \right)^2}$$

RESULTS

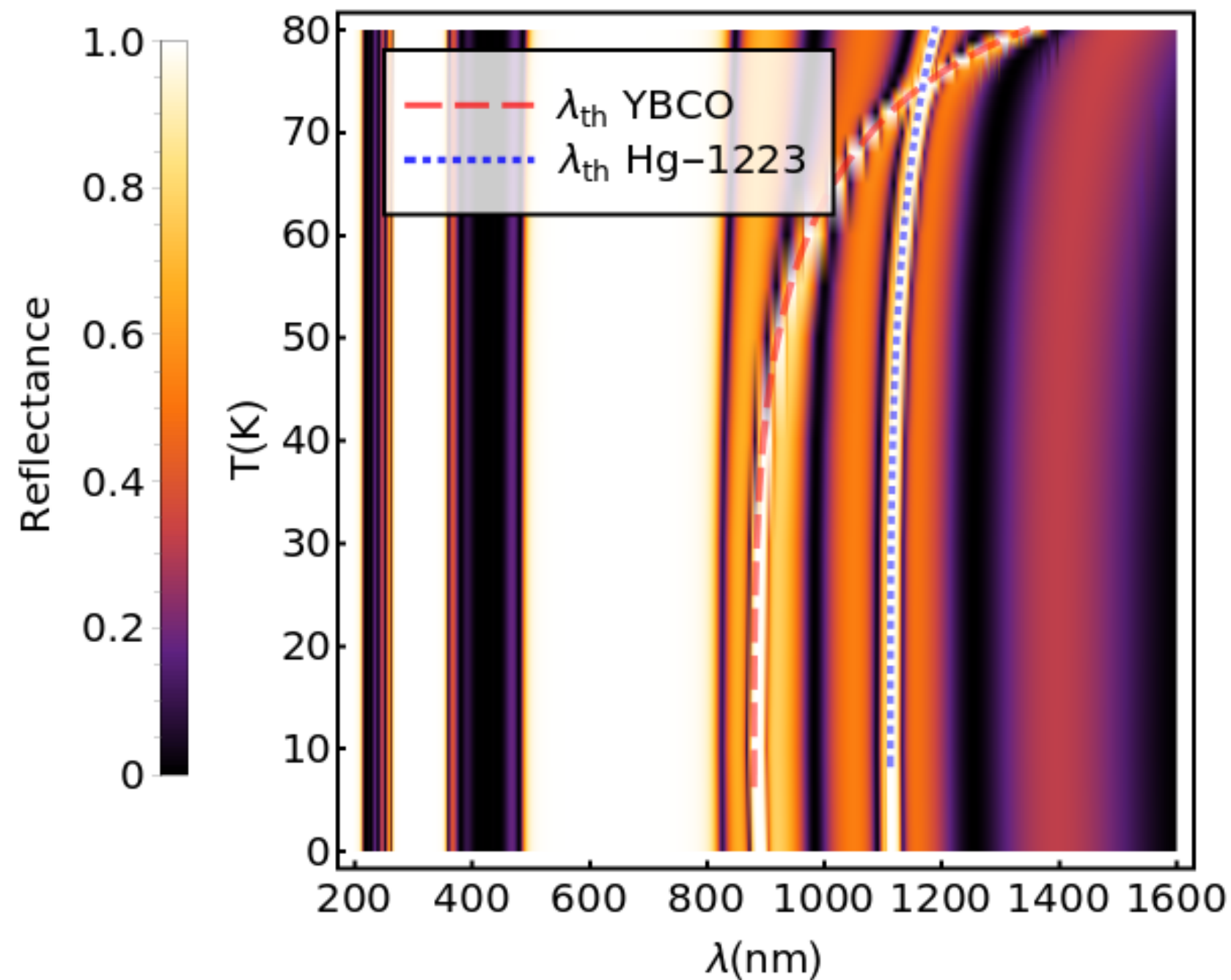
Temperature variation

Superconducting threshold

$$\lambda_{th} = 2\pi\lambda_L(T) = 2\pi \frac{\lambda_0}{\sqrt{1 - \left(\frac{T}{T_c}\right)^4}}$$



$$n_D = 2.437$$



CONCLUSIONS

- **Characterizing the optical response allows us to control the propagation of electromagnetic waves in the crystal.**
- **Introducing superconductors into the system produces new gaps, presenting new technological applications such as thermal sensors.**

PERSPECTIVES

- **It is possible to study the optical response of the system to variations of other parameters such as crystal radius or layer thickness.**
- **Other combinations of materials can be proposed and even alternate the order in which the superconductors are configured.**

THANK YOU

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