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SMR 1331/8

## AUTUMN COLLEGE ON PLASMA PHYSICS

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# Radioastronomy and the Nature of Turbulence in the Interplanetary Medium

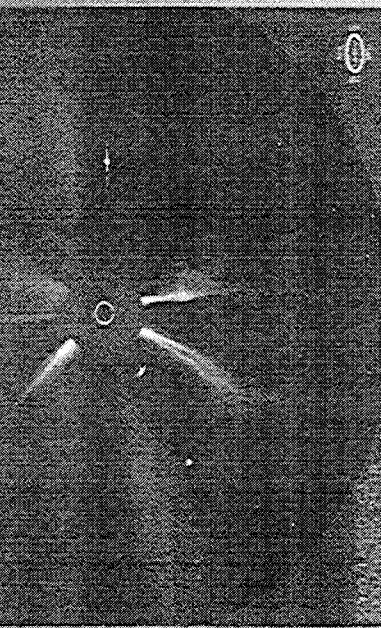
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These are preliminary lecture notes, intended only for distribution to participants.



# Radioastronomy and the Nature of Turbulence in the Interplanetary Medium

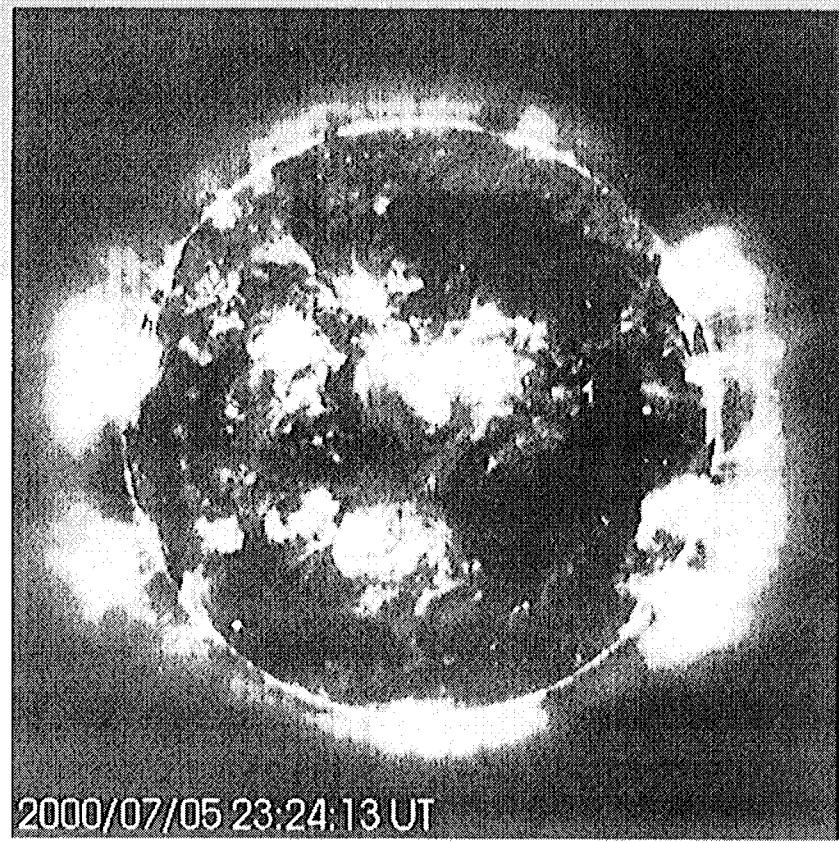


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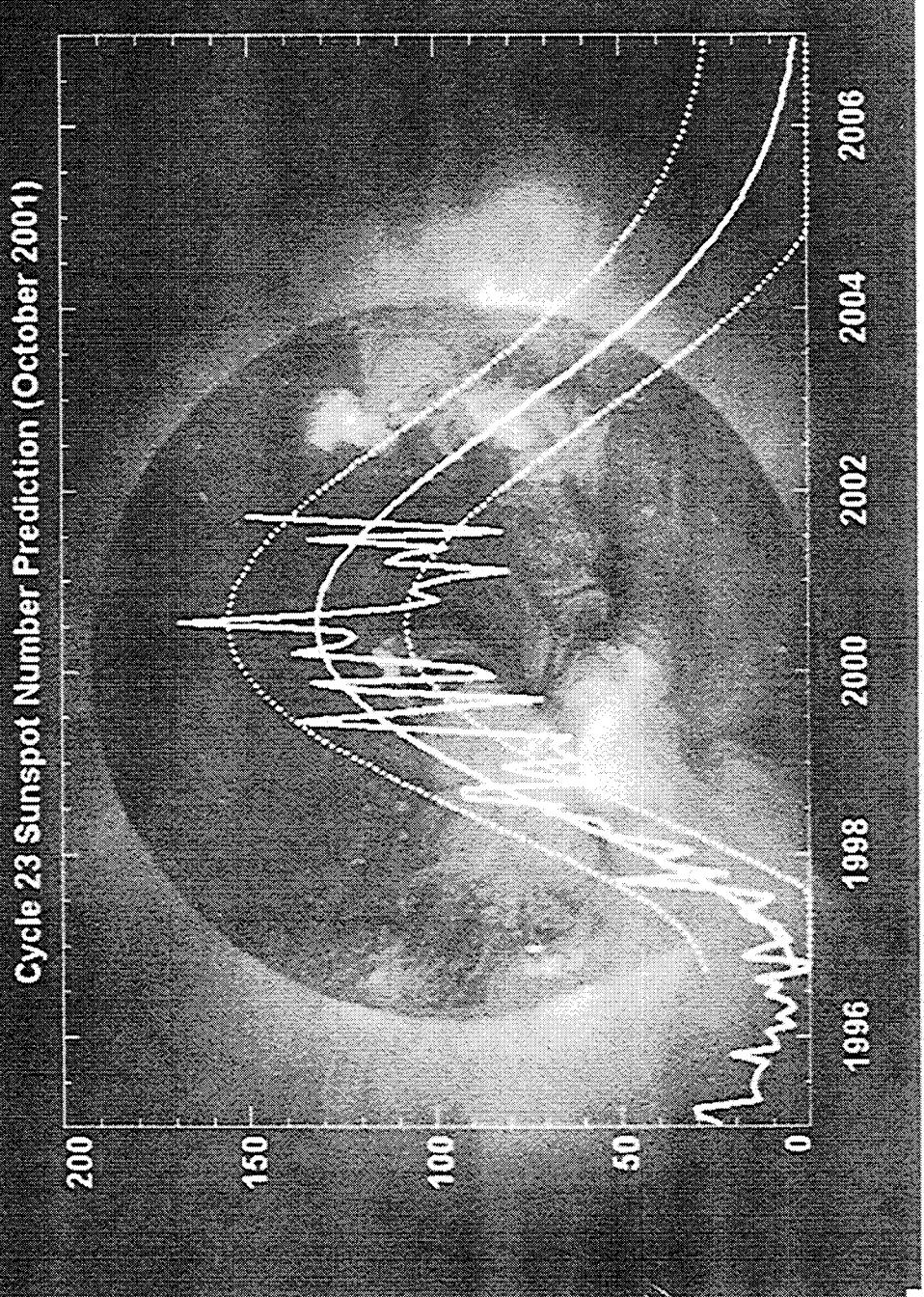
# Magnetism and the Solar Corona

- Loops show hot gas confined by Sun's magnetic field

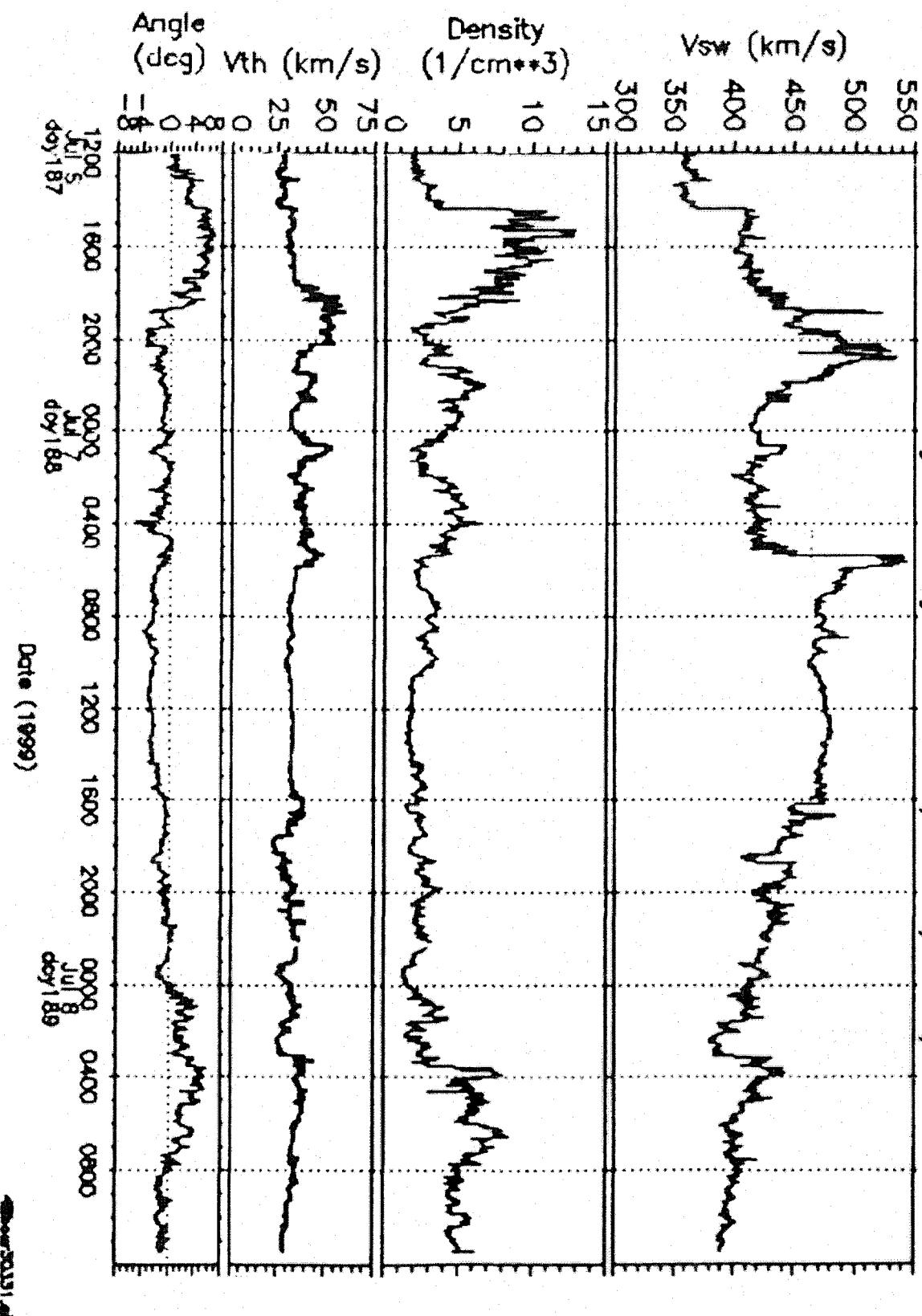


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# The Solar Atmosphere and the Sunspot Cycle



University of Maryland soho/celios/mtof/PM



## A Comparison of the Solar Wind and ISM

Parameter	Diffuse Ionized Phase of ISM	Solar Wind at 1a.u.
$B_0(\mu\text{G})$	5	50
$n_e(\text{cm}^{-3})$	0.080	5
$T_e(\text{K})$	8000	$1.5 \times 10^5$
$T_i(\text{K})$	$\leq 15000$	$5 \times 10^4$
$V_A(\text{km/sec})$	23	50
$\nu_{ION-NEUT}(\text{sec}^{-1})$	$8.3 \times 10^{-10}$	0
"Turbulent Age" <sup>1</sup>	$4 \times 10^6$ years	4 days

### Notes

- (1) Rough estimate of time from generation of turbulence to arrival at point of observation.

# Eine Kleine Nachtphysik ...

## The equations of hydro/magnetohydrodynamics

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \quad (1)$$

With "steady state" spherical symmetry  $\rightarrow$

$$\frac{1}{r^2} \frac{d}{dr} (r^2 \rho v_r) = 0 \quad (2)$$

$$v_r \frac{dv_r}{dr} = -\frac{1}{\rho} \frac{dp}{dr} - \frac{GM_{\odot}}{r^2} - \frac{1}{\rho} \frac{d}{dr} \left( \frac{\delta B^2}{8\pi} \right) \quad (3)$$

If far from the Sun, the right hand side  $\rightarrow 0$ .  $\Rightarrow$

$$\frac{1}{2} \frac{dv_r^2}{dr} = 0 \quad (4)$$

Therefore  $v_r^2 = C$

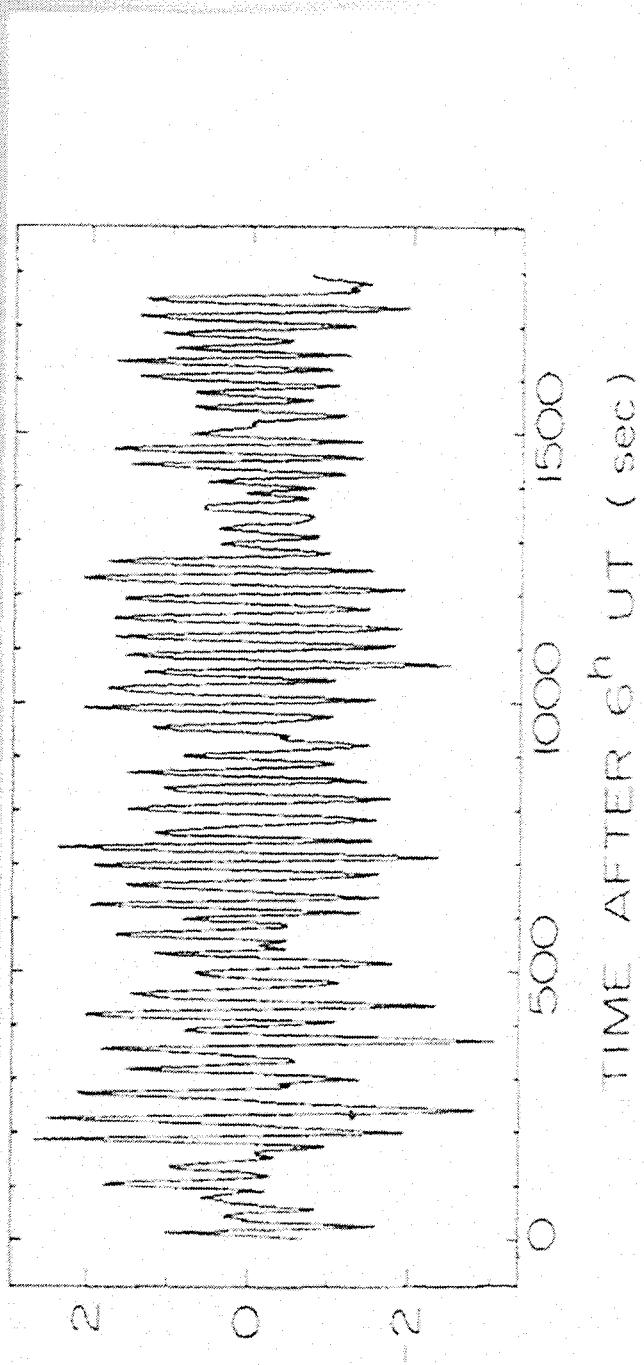
$$\frac{de}{dt} + \vec{v} \cdot \nabla e = e \nabla \cdot \vec{v} - \frac{1}{\tau} \left( \frac{\delta B^2}{4\pi} \right) \quad (5)$$

Therefore "Distant Wind" has  $v = C$ ,  $\rho \propto r^{-2}$ .

# Coronal Magneto-hydrodynamic Waves

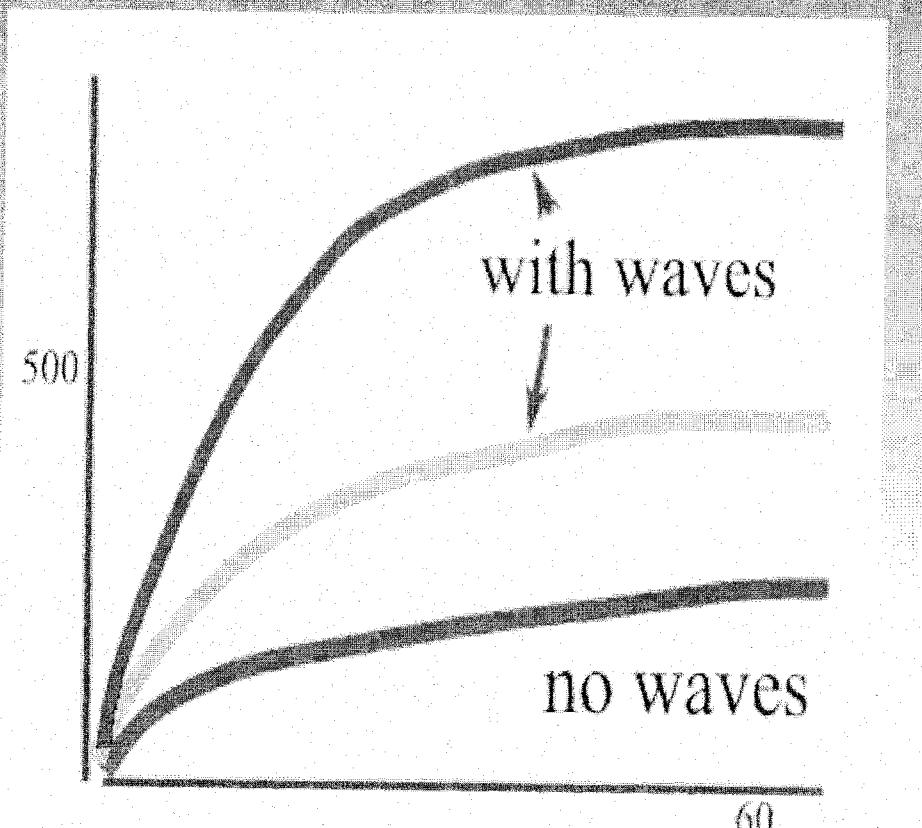
These have been found especially in strong magnetic fields

- Associated phenomenon of Faraday rotation



# Why is Near-Sun Turbulence of Interest?

- Observed wind is faster and hotter than models
- Damping of turbulence provides heat and momentum
- Phenomenon similar to flow in middle atmosphere



Velocity versus Heliocentric Distance

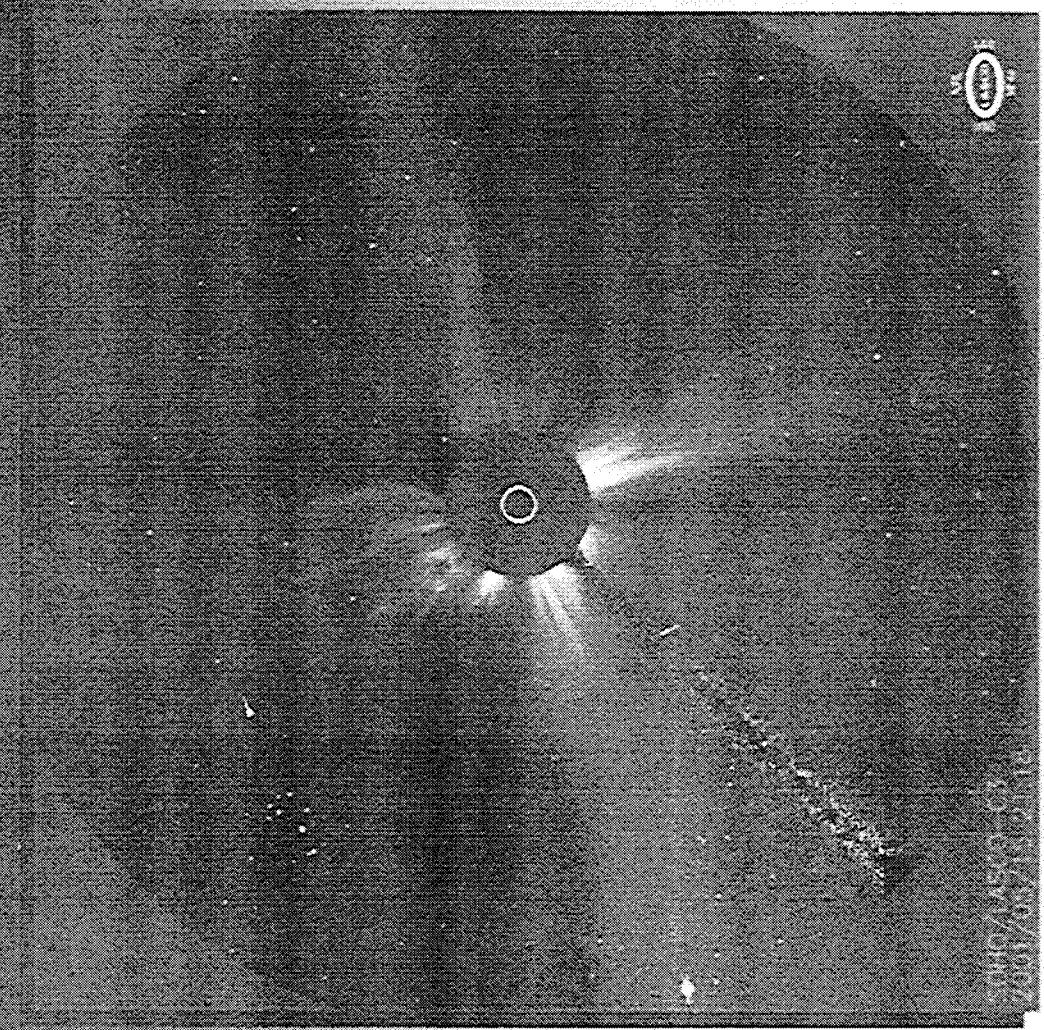
# **The Wave Driven Solar Wind**

For wave-driven models to be viable (1) there must be a sufficient wave flux (turbulent energy density)  $\simeq (2-5) \times 10^5$  ergs/sec/cm<sup>2</sup>, and (2) the waves should damp sufficiently rapidly ( $\leq 10 - 20R_\odot$ ).

## **Two Models for Turbulence**

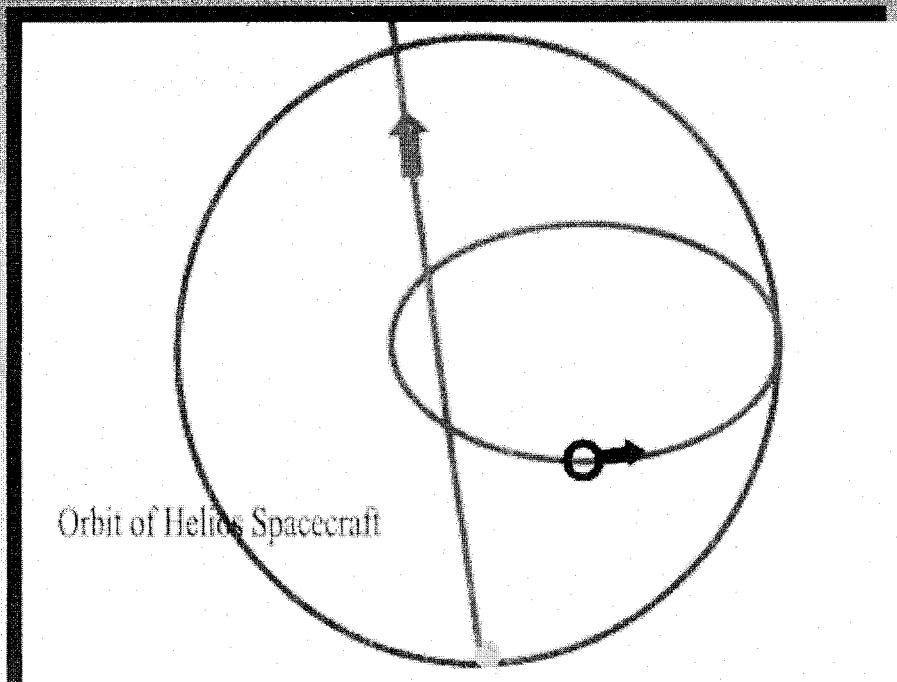
- (1) Low Frequency Waves (periods of tens of minutes to couple of hours)
- (2) High Frequency Waves (plasma frame frequencies of 1 - 1000 Hertz).

# The Sector of Space for Remote Sensing

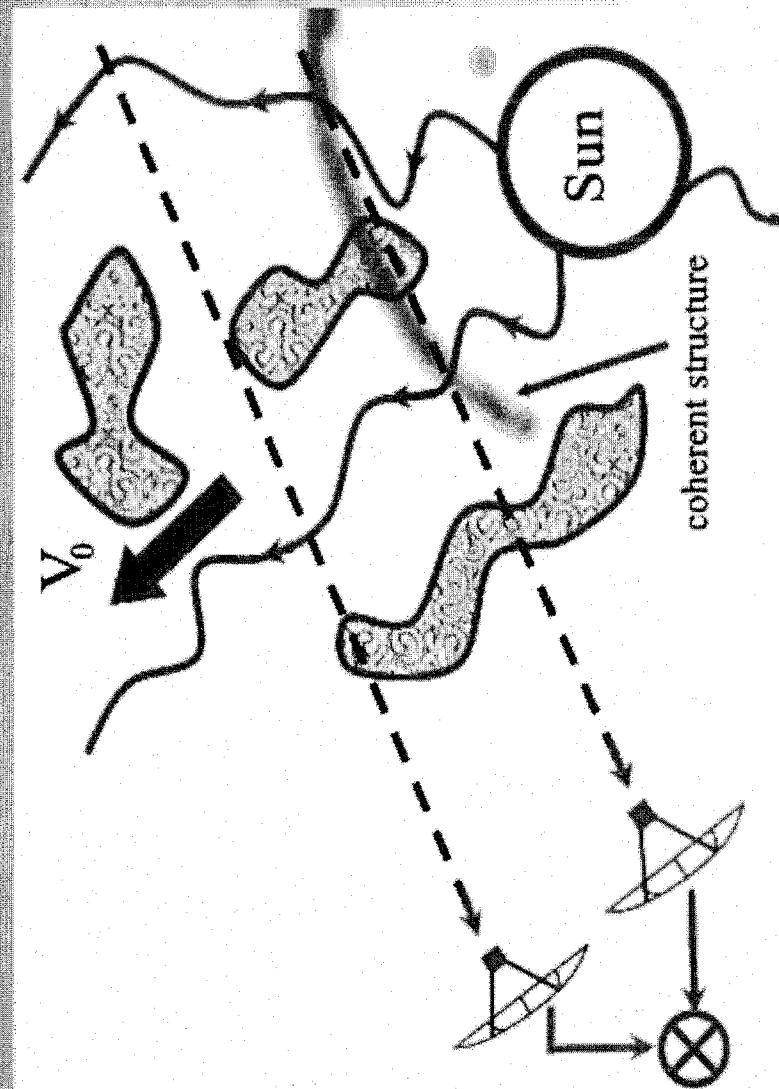


# Remote Sensing of Plasmas and Plasma Turbulence

- Why bother?
- Information important where in-situ data not available
- Helios spacecraft found solar wind fully accelerated at 0.28 au

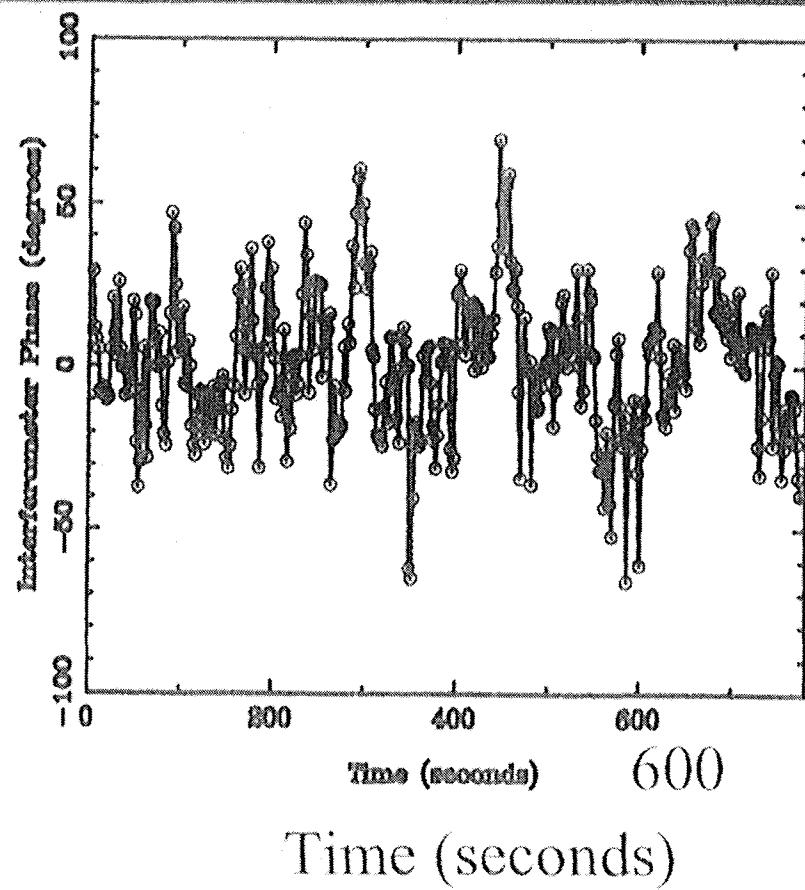


# The Physical Basis of Phase Scintillations



# Measured VLBI Phase Scintillations

- Frequency 8.4 GHz
- Source elongation 22.6 solar radii
- Baseline Medicina-Noto (500 – 800 km)



# Expression for Phase Power Spectrum

R. L. Mutel, 1975

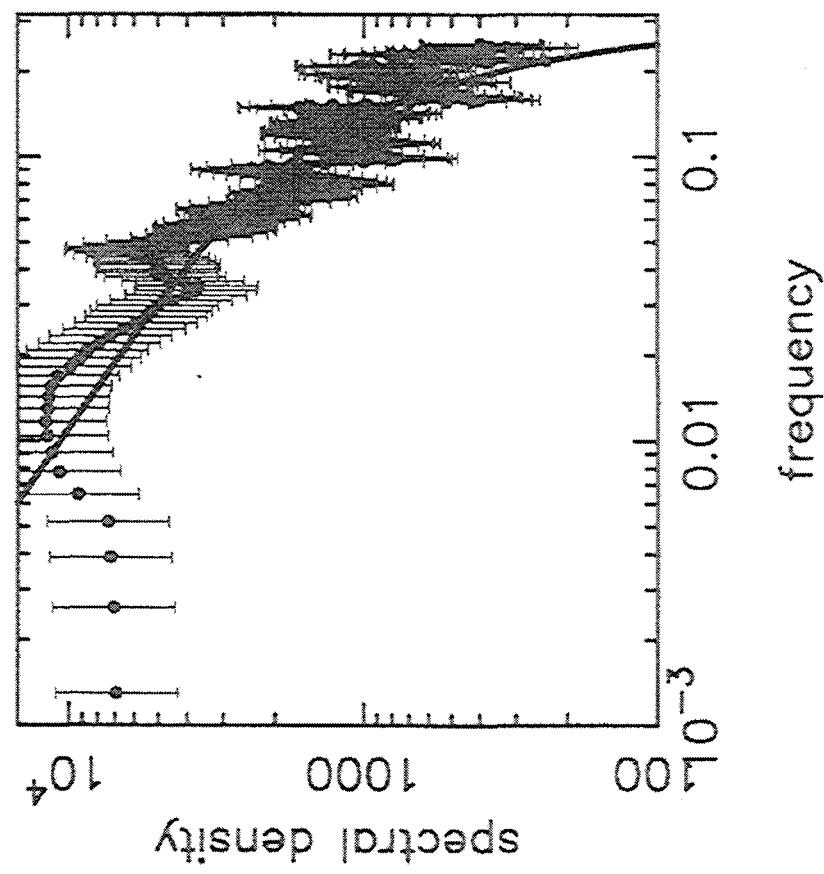
## Physical Variables in Phase Power Spectrum

- Density power spectrum,  $P_{\delta n}(q) = C_N^2 q^{-\alpha}$ .  $C_N^2(r)$  expression given by Spangler *et al* 1996, *Solar Wind 8*.
- Solar Wind flow speed  $V_0$ .
- Effective thickness of solar wind plasma screen,  $L$ .
- Outer scale  $l_0$

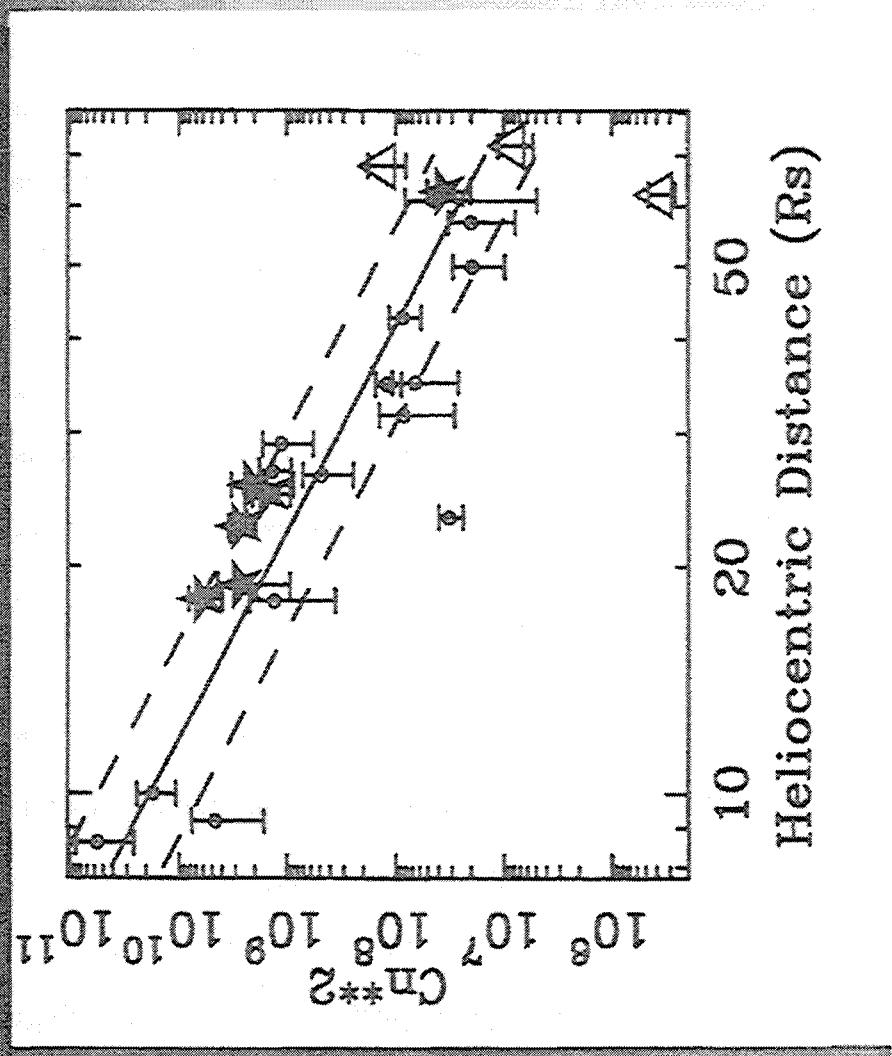
Then ...

$$P_{\delta\phi}(f; \alpha, C_N^2, L, V_0) = \frac{32\pi^2 r_e^2 \lambda^2}{V_0} C_N^2 L l_0^\alpha \int_0^\infty dq_\perp \frac{1}{(1 + l_0^2[q_\perp^2 + q_\parallel^2])^{\alpha/2}} \sin^2 \left[ \frac{\pi f \rho_\parallel}{V_0} + \frac{q_\perp \rho_\perp}{2} \right] \quad (1)$$

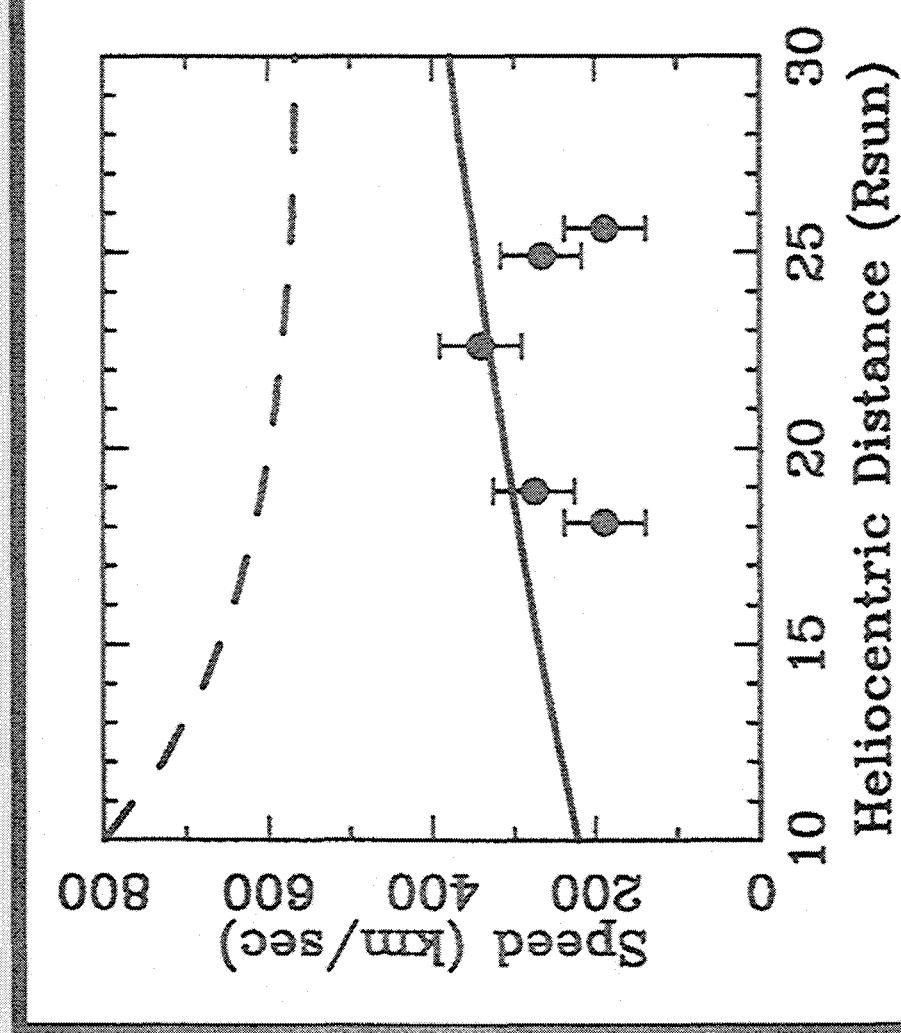
# Measuring Solar Wind Parameters: Fits to Phase Power Spectra



# Evolution of $C_N^2$ with Heliocentric Distance



# Remote Measurement of the Speed of the Solar Wind

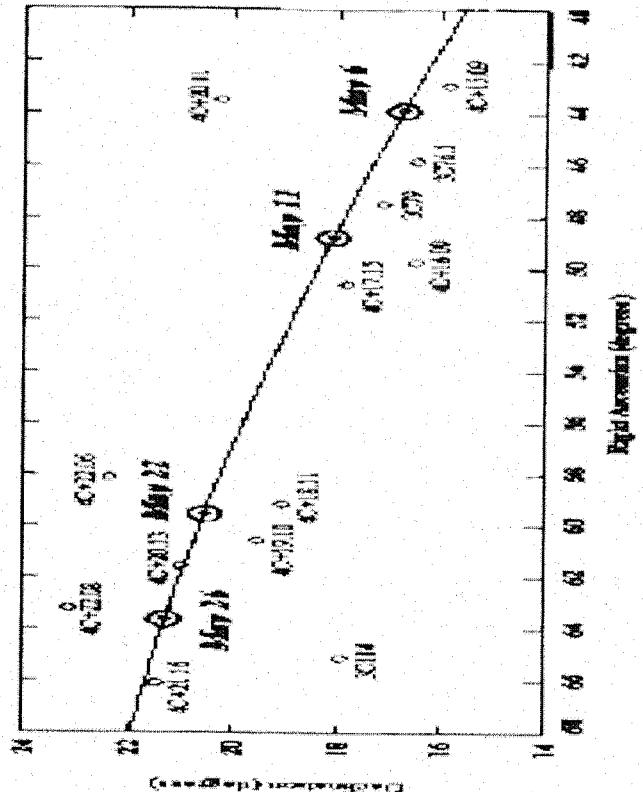


## What can one do with coronal Faraday rotation observations?

- $RM$  fluctuations  $\rightarrow$  coronal Alfvén waves
- $\bar{RM}$   $\rightarrow$  magnetohydrodynamic structure of corona
- *Depolarization*  $\rightarrow$  small scale Alfvén waves

# Tomography of the Corona

- Mancuso & Spangler,  
ApJ 525, 1995, 1999;  
also M&S submitted
- RM along 13 lines of  
sight
- Goal: determine B and  
Ne models



## A Model for the Plasma Structure of the Solar Corona

- Density distribution

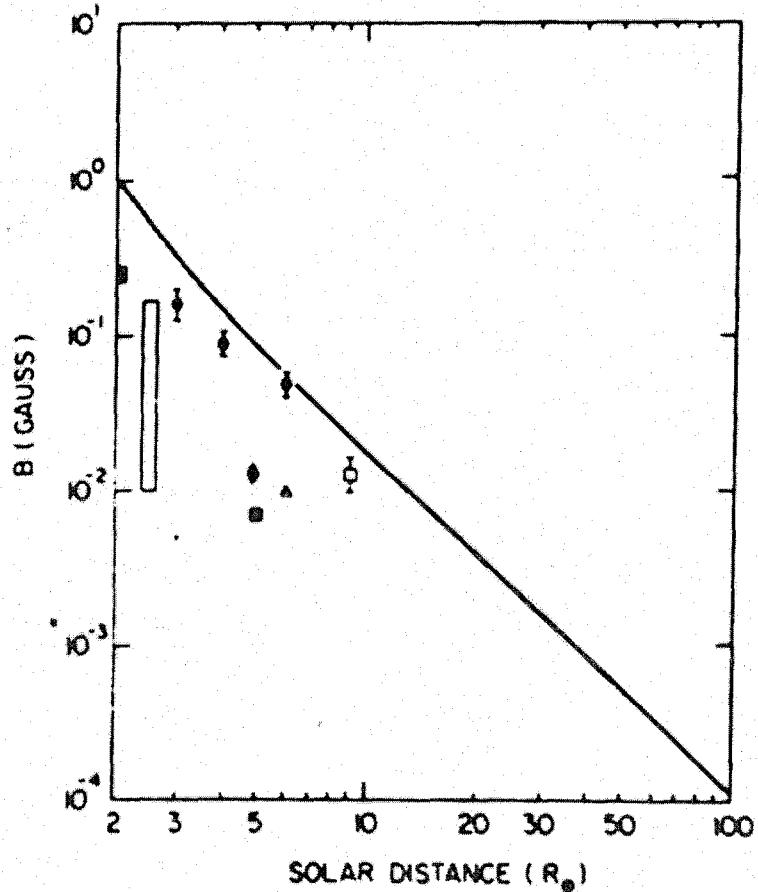
$$n(r, \theta, \phi) = n_{CH}(r) + (n_{CS}(r) - n_{CH}(r)) \exp [-(\theta/w(\phi))^2] \quad (8)$$

$n_{CH}$  and  $n_{CS}$  compound power laws in  $r$ .

- Magnetic Field model

$$B(r) = 0.06(r/R_\oplus)^{-3} + 3.1(r/R_\oplus)^{-2} \text{ nT} \quad (9)$$

with  $R_\oplus = 1$  astronomical unit (Pätzold, 1987).



# Comparison of Model Comparisons with Observations

Figure 10 shows the correlation coefficient between the observed and model RM values for the various models.

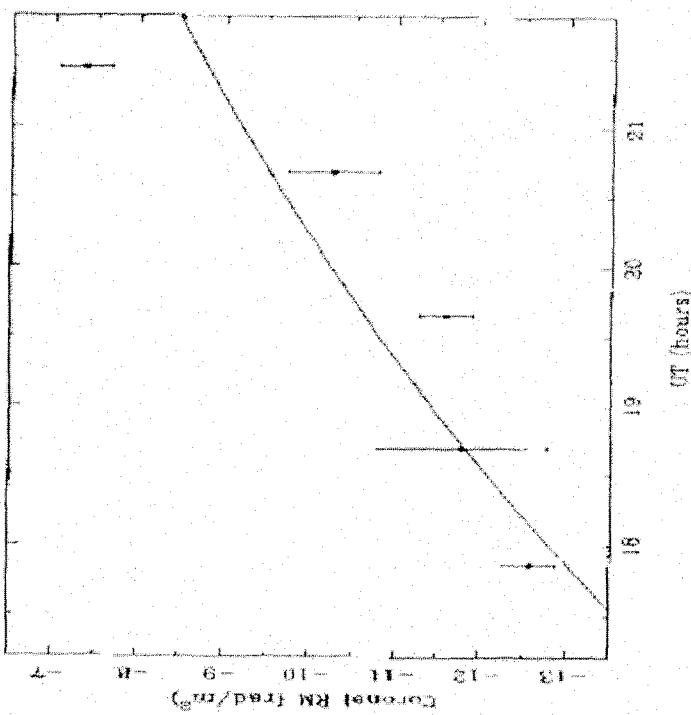
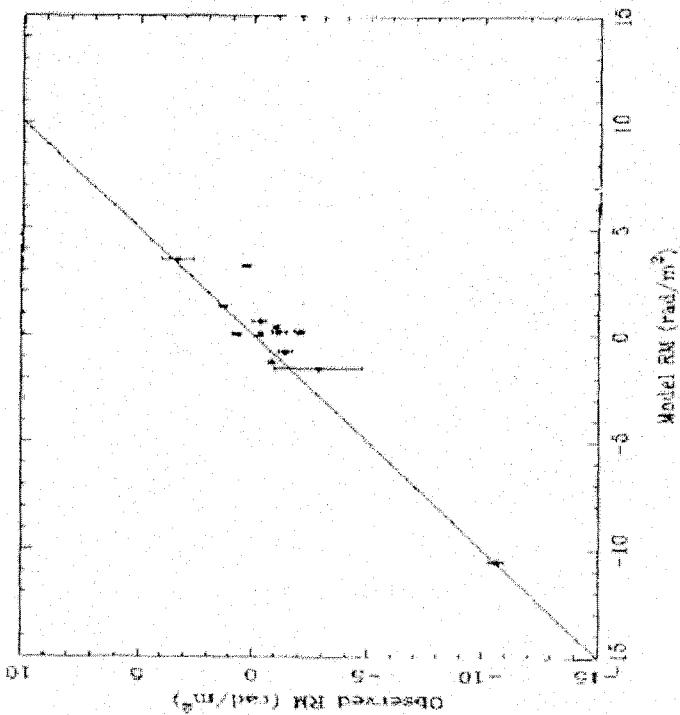


Figure 11 shows the correlation coefficient between the observed and model RM values for the various models.



# How Observables Relate to the Theory

(1) RMS Rotation Measure Fluctuation

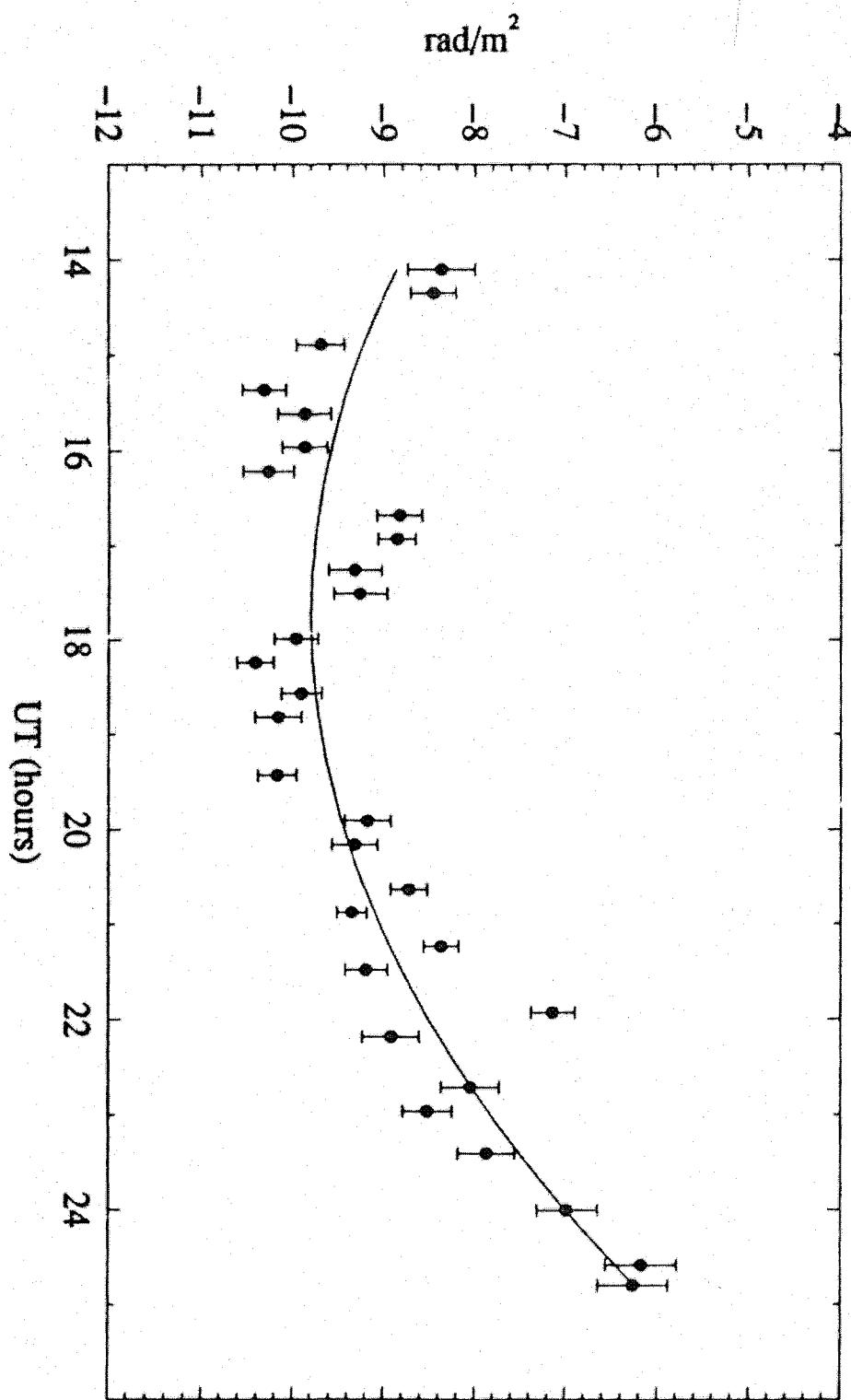
$$\langle (\delta RM)^2 \rangle = \left( \frac{e^3}{2\pi m_e^2 c^4} \right)^2 \int_0^\infty n_e^2 \langle (\delta B_z)^2 \rangle L ds \quad (1)$$

(2) Relating  $\delta B_z$  to wave flux

$$F_A = \frac{\langle \delta B^2 \rangle}{8\pi} (3U + 2V_A) \quad (2)$$

(3) Relating flux to coronal base

$$F_{A0} = \frac{\langle \delta B^2 \rangle}{4\pi V_A} (U + V_A)^2 \left[ \frac{A}{A_0} \right] \quad (3)$$



HOLLWEG ET AL.: CORONAL WAVES

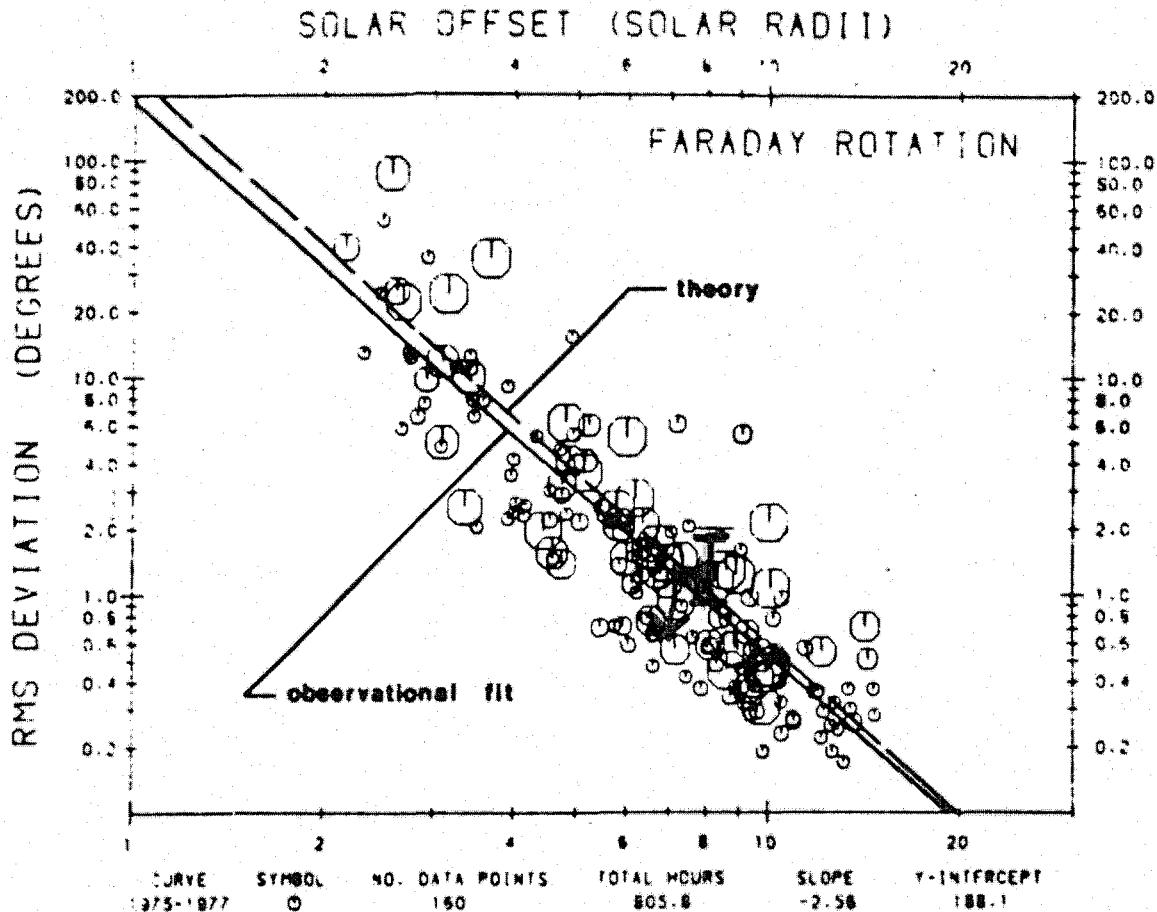


Fig. 6. Standard deviation of the Faraday rotation from the mean linear trend over each observation pass versus solar offset; all data from 1975 through 1977. The solid line is a power law fit to the observations. The dashed line is the theoretical expression (25), which is based on the assumption that most of the Faraday rotation variance is caused by fluctuations of the coronal magnetic field and not the electron concentration.

# Summary: Radio Remote Sensing of the Corona and Solar Wind

Magnetic field fluctuations are consistent with small amplitude, passive, dynamically unimportant fluctuations

- Observations consistent with non-propagating fluctuations
- Faraday rotation “tomography” consistent with independent measures of coronal plasma state
- Observations consistent with “smooth” corona, without large waves or turbulence