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international centre for theoretical physics

SMR 1302 - 4

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**WINTER SCHOOL ON LASER SPECTROSCOPY AND APPLICATIONS**

**19 February - 2 March 2001**

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***COLD COLLISIONS IN THE PRESENCE OF LIGHT***

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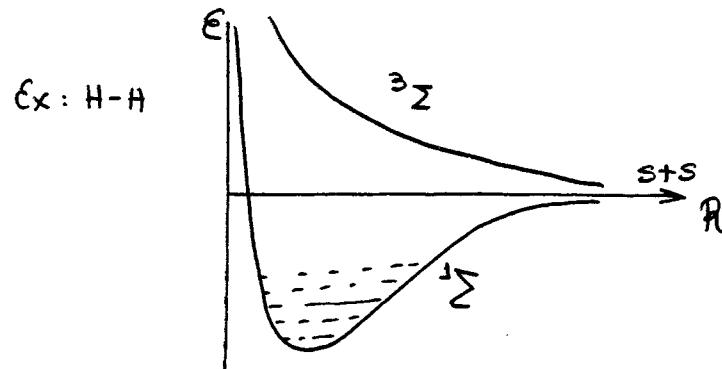
*These are preliminary lecture notes, intended only for distribution to participants.*



## II- Coed collisions in the presence of light

- \* During approximation of two atoms

electronic states  $\Rightarrow$  Molecular potentials



- \* Molecular potential has long range part, normally not important
- $\rightarrow$  coed collision  $\rightarrow$  long range is important and dominant

- { \* conventional :  $E_c \sim 1\text{eV}$   
 \* cold collision :  $E_c \sim 10^{-8} - 10^{-10}\text{eV}$

Small details are important

- \* Time of collision

$\rightarrow$  high temperature  $\Delta t \sim 10^{-12} - 10^{-14}\text{sec}$

$\rightarrow$  coed collision  $\Delta t \sim 10^{-7} - 10^{-9}\text{sec}$

- \* In the case of excited state participation

time collision  $>$  lifetime state

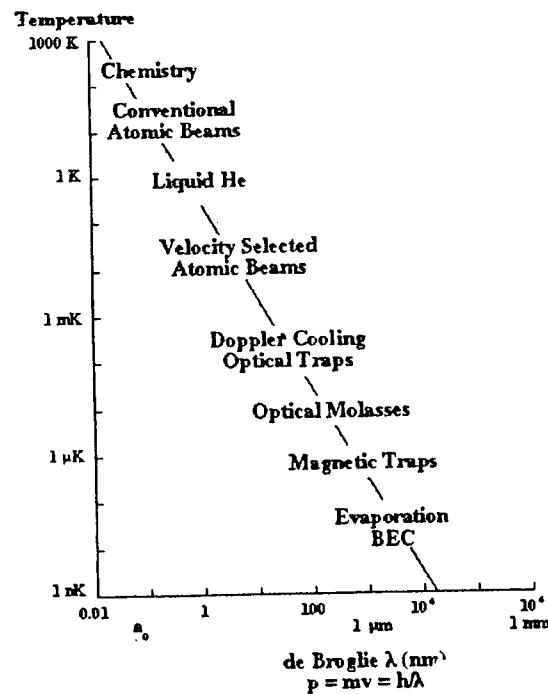
- $\Rightarrow$  spontaneous emission is important  
 (prototype for an open dissipative system)

## \* Bound States

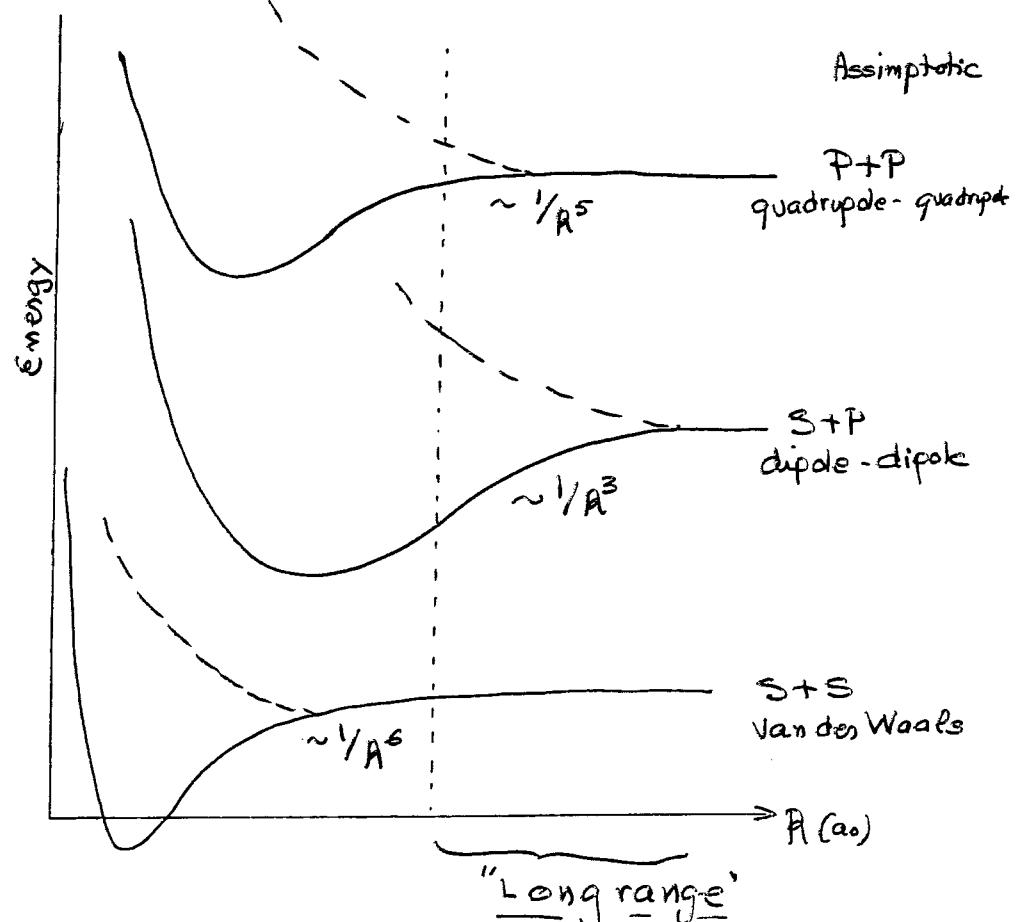
- $\rightarrow$  Low collisional energy allow to observe long range bound states

## \* Gas at quantum regime

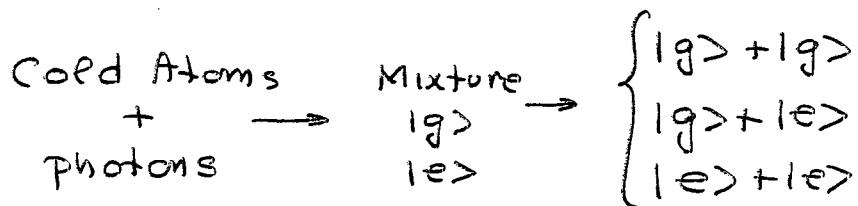
- Thermalization (evaporative cooling)
- Macroscopic properties depend on interaction



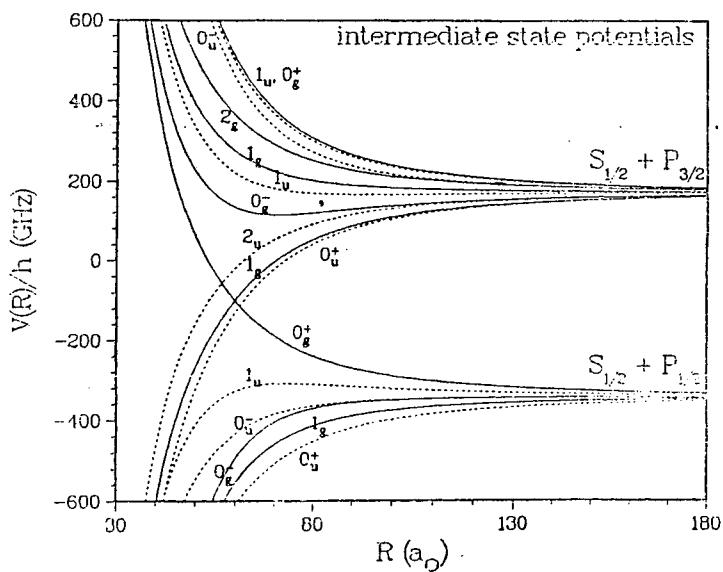
\* Depending on the internal state of the atoms → different interaction



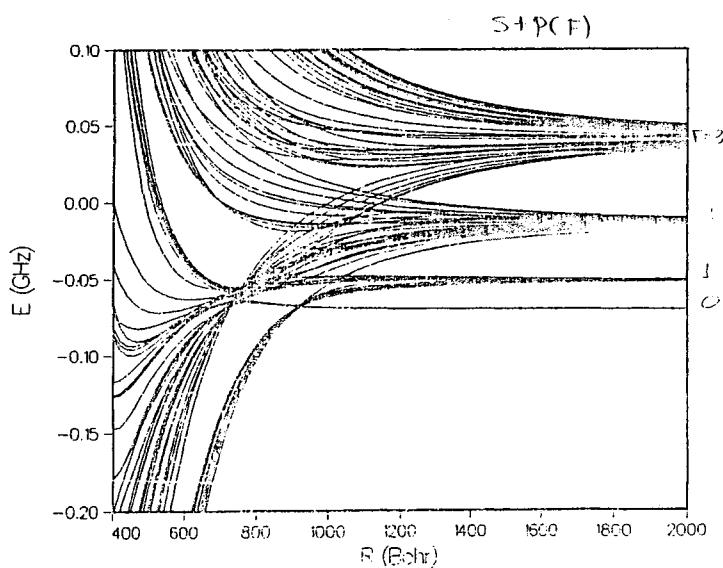
### Possibilities for collisions



- \* Fine Structure in the atoms  
➡ several symmetries



- \* Hypersfine Structure → Each level several possibilities



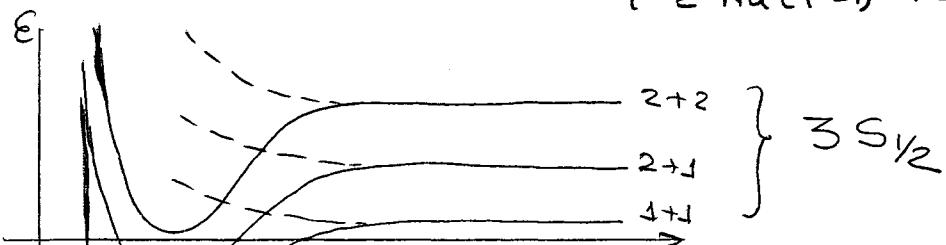
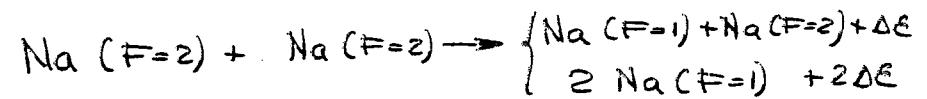
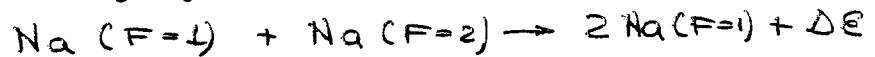
## Ground - Ground state collision

- \* Elastic process - collision is determined by the scattering length "a"  
 $\sigma \sim 4\pi a^2$   
→ determination of  $a$  is by itself interesting collision problem

Important:

- evaporative cooling
- properties of a quantum gas
- 

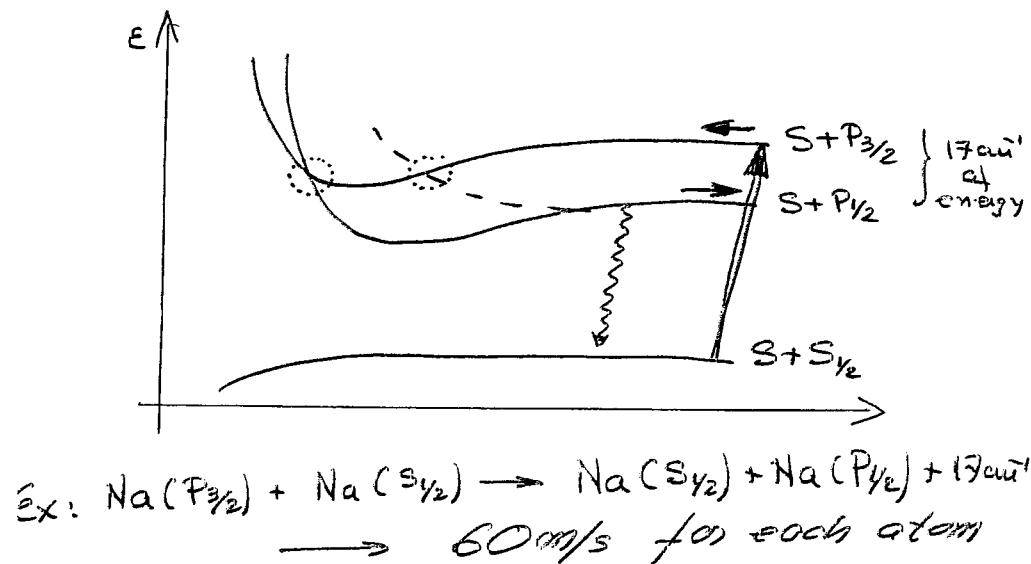
- \* Inelastic process : Hyperfine changing collision (HCC)



\* Interaction is at short range  
Ex:  $\Delta E = 1.7 \text{ GHz} \rightarrow \Delta v = 6 \text{ m/s}$

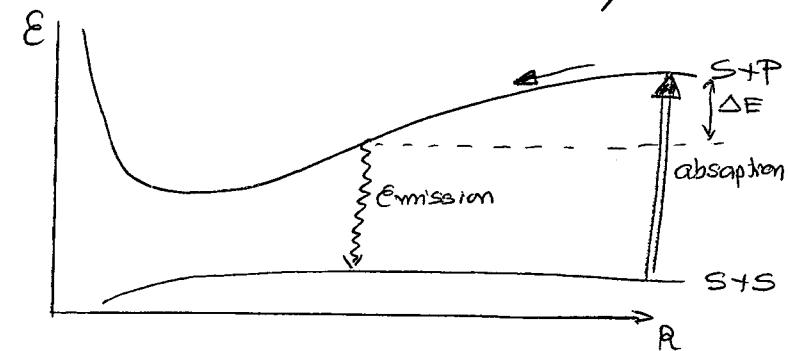
## Ground + Excited State

- \* Long range part of potentials become very important
- \* Processes where internal energy is converted in kinetic → exothermic process  
and there are two main contribution
- \* Fine Structure Change (FSC)



## \* Radiative Escape (RE)

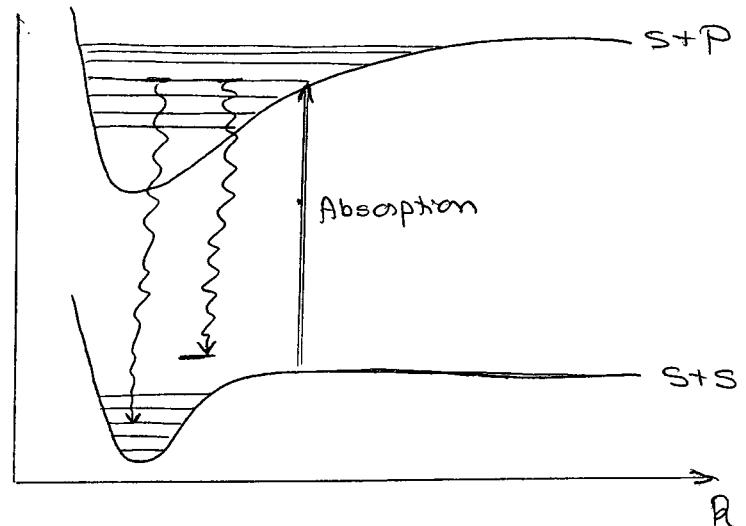
- Atoms get excited at long range → gain kinetic energy → decay



## Important points

- \* statistical nature of decay → velocity distribution
- \* Energy gained depends on excitation point, potential slope
- \* Survival is important

## Ground + Excited (Selective excitation)



\* Absorption is selective  
(Free to bound) (Franck-Condon factor)

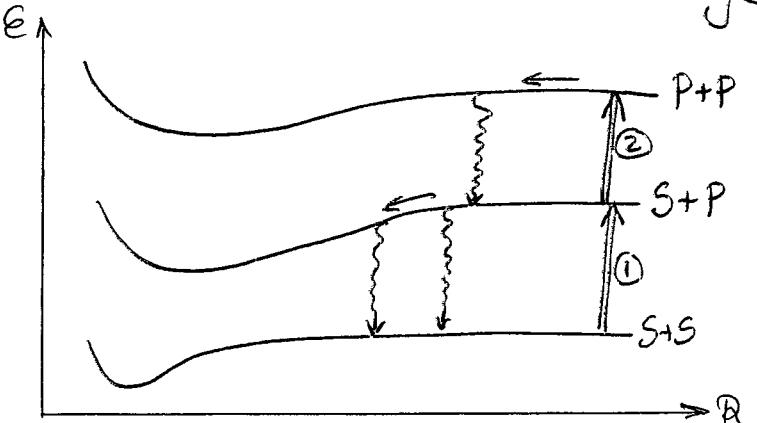
\* Emission  
 $\Rightarrow$  bound  $\rightarrow$  free (kinetic energy)  
 $\Rightarrow$  bound  $\rightarrow$  bound (Association)

Absorption reveals structure of bound states at long range (S+P)

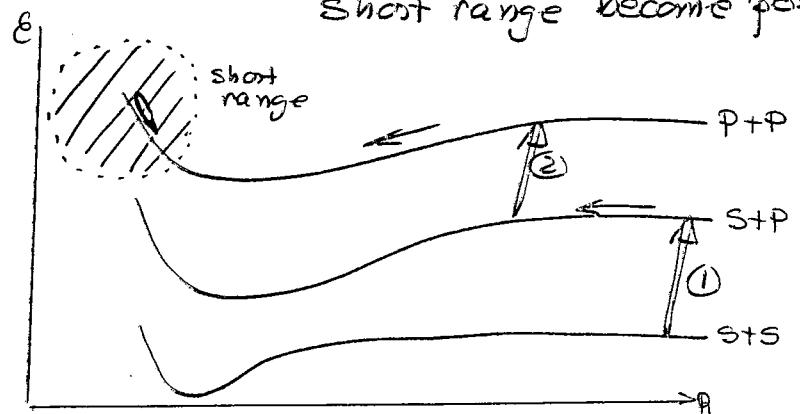
$\rightarrow$  Both processes get atoms out of the trap.

## Excite + Excite state collisions

$\rightarrow$  Starting with double excited states at long range  $\rightarrow$  hard to survive to short range

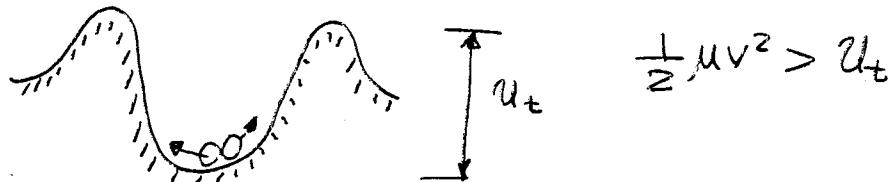


$\rightarrow$  Starting at long range but getting motion before second excitation  
 $\rightarrow$  better chances for survival  
 $\rightarrow$  Collision channels at short range become possible



II-10

Exoergic collisions  $\xrightarrow{\text{may}}$  Trap Loss  
 Kinetic energy > trap depth.



- \* Strong depending on  $U_t$   
 $\longrightarrow$  trapping parameters  
 $(\text{Int.}, \Delta, \frac{d\beta}{dt}, \text{etc})$

- \* Can be studied using variation of number of atoms in the trap

$$\frac{dN}{dt} = L - \gamma N - \beta \int_V M^3 c(r) d^3 r$$

$L$  = Loading rate

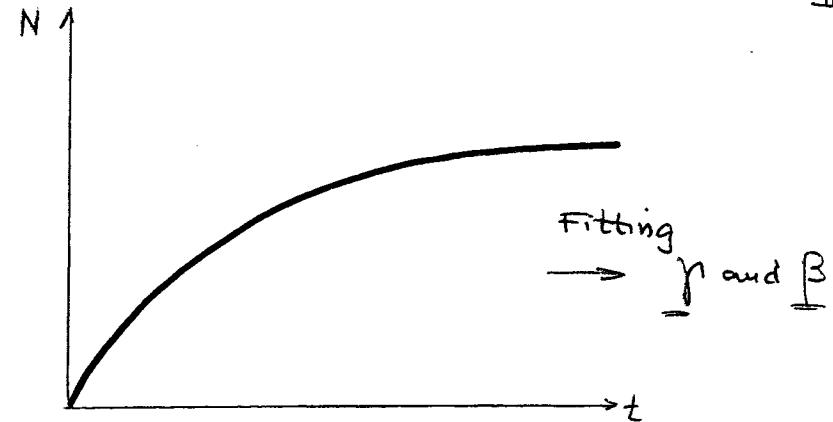
$\gamma$  = background atoms rate

$\beta$  = volumetric rate of collisions that takes to losses

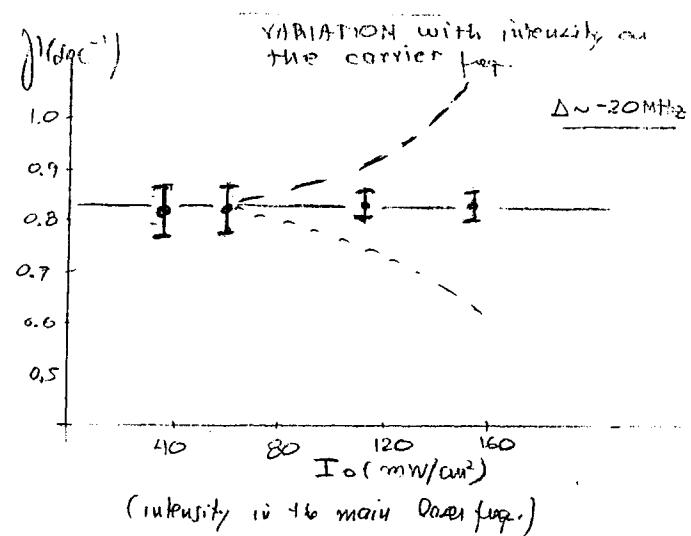
Having  $N$  atoms, each collide  $m \Gamma v$   
 $\rightarrow$  total rate  $\frac{1}{2} N m \Gamma v \rightarrow$  each collision

loss of two atoms

$$\text{rate losses} = 2 \frac{1}{2} N m \Gamma v = \underline{\beta m N} \quad (\underline{\beta = \langle \Gamma v \rangle})$$



$\gamma$  depends on back ground



(Intensity is +6 main Dose freq.)

⇒ Important facts related with the analysis of transient behavior

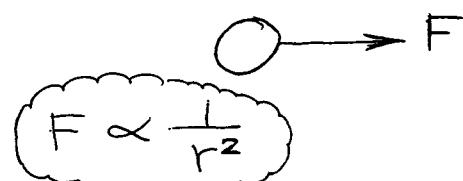
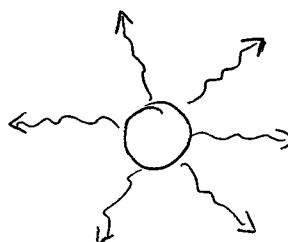
II-12

High number of atoms + High light intensity

$\Downarrow$   
radiation trapping

$\Downarrow$

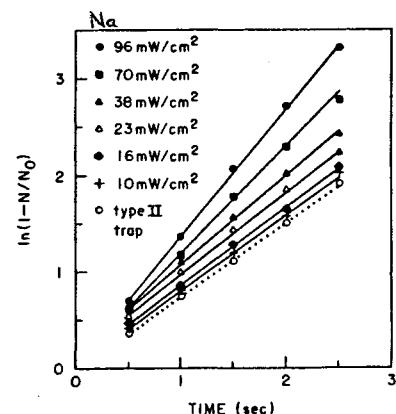
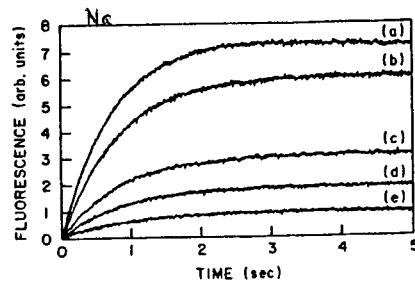
positive pressure on atoms



⇒ Due to positive pressure  $\Rightarrow$  atomic density is constant.

$$\frac{N}{V} = M_c$$

$$\rightarrow \frac{dN}{dt} = L - (\gamma + M_c \beta) N$$



$$N = N_0 [1 - \exp(-\gamma - M_c \beta t)]$$

II-13

- \* Independent measurement of  $\gamma$
- \* Determine  $\gamma + M_c \beta \rightarrow$  obtain  $\beta$

Atomic density with Gaussian Profile.

→ trapping forces more important

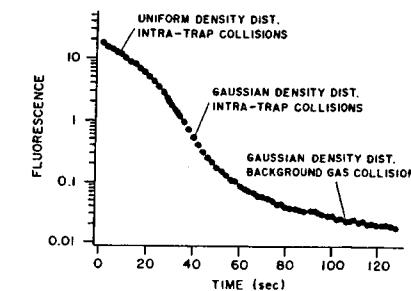
$$M_{\text{crit}} = M_c(t) e^{-r^2/w^2}$$

$$\frac{dN}{dt} = L - \gamma N - \frac{\beta N^2}{(2\pi)^{3/2} w^3}$$

$$\text{Solution: } N(t) = \frac{2L \sinh(\alpha t/2)}{\gamma \sinh(\alpha t/2) + \alpha \cosh(\alpha t/2)}$$

fitting  $\Rightarrow$  determine  $\beta$

During unloading of a MOT  $\rightarrow$  observation of several regimes



$\beta$  depends

- IT-14
- Depth of trap
  - Excited State population
  - Collision channels (FSC, RE, HCC)

$$\boxed{\beta = \beta(I, \Delta, \text{atom})}$$

Besides loading and unloading  
 → Sudden change of intensity

Load at high intensity → change to low intensity

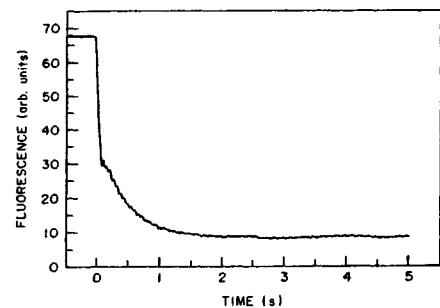


FIG. 26. Fluorescence-time spectrum. The abrupt drop at the initial moment is due to a change in light intensity without variation in the number of trapped atoms. From Santos *et al.* (1996).

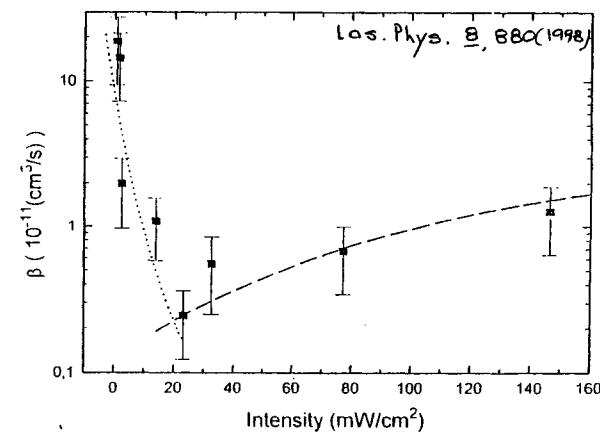
No → number at I

No' → number at I'

$$N(t) = N_0 + (N_0 - N_0') \exp \left[ -\gamma \left( 1 + \frac{M_0 \beta}{\gamma} \right) t \right]$$

→ Allow determination of  $\beta(I')$  down to very low intensity

Typical behavior of  $\beta(I)$  observed



Several regions (Interpretation)

High intensity  $\Rightarrow$  FSC + RE

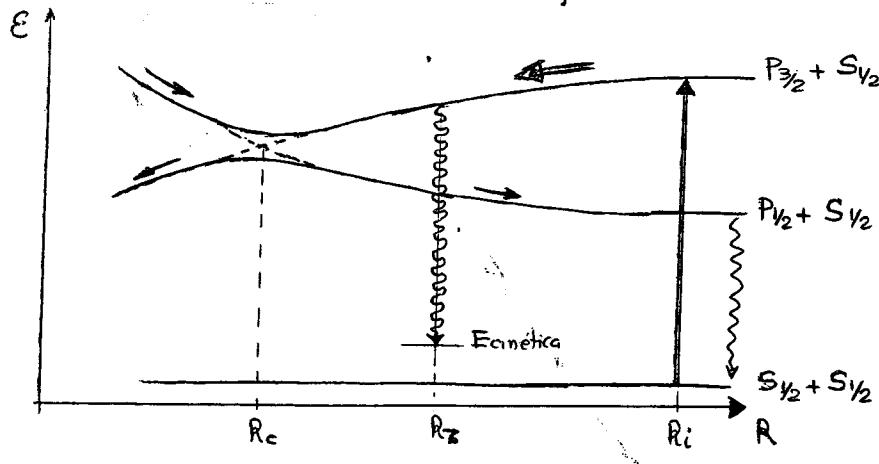
Low intensity  $\Rightarrow$  HCC

Does anybody detected FSC?  
 YES

Does anybody detected HCC?  
 NO

Main Model for  $\beta \Rightarrow$  Gallager-Pritchard  
(GP)

- \* It is a semiclassical model considering FSC + RE



- \* A pair is excited at long range ( $R_i$ )
- \*  $\Delta < 0$ , probably attractive state is excited
- \* Number of excited pairs  
(Pairs) <sub>$R_i$</sub>  =  $\frac{m^2}{2} 4\pi R_i^2 dR_i$
- \* Rate of excitation  
$$\left(\frac{I}{\hbar\omega}\right) \frac{\lambda^2}{2\pi} \frac{T}{T^2 + 4(\Delta - \frac{C_3}{\hbar R^3})^2}$$

- \* Once excited there will be decay (by spontaneous emission) at  $R_S$  ( $R_S$  is such that gained energy is larger than  $U_t$ )

$$P_{\text{rad}} \propto e^{-T t^* (R_i, R_S)} \quad (\text{AE})$$

- \* At the anti-crossing may occur FSC (position  $R_E$ )

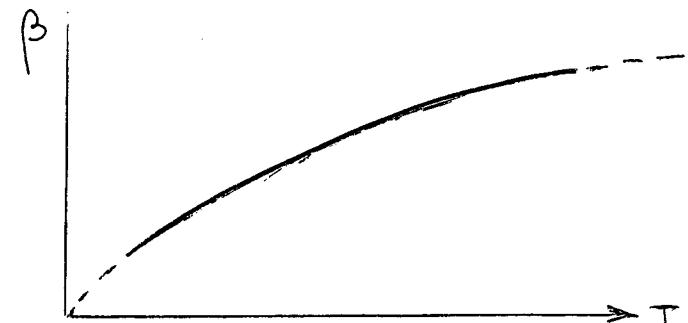
$$[P_{\text{DJ}}(R_i, R_E)] \quad (\text{FSC})$$

$\beta \propto$  Rate of Collisions with loss

$$= \int_{R_i}^{\infty} \frac{m^2 4\pi R_i^2 dR_i}{2} B(I, \Delta) [P_{\text{rad}} + P_{\text{DJ}}]$$

$B$  is a function on  $I, \Delta, R_i$

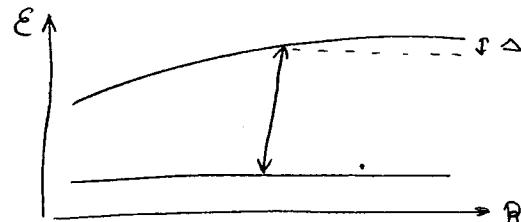
Obtained  $\beta(I)$



"decreasing excitation"

$\beta(\Gamma) \rightarrow$  behavior of trap

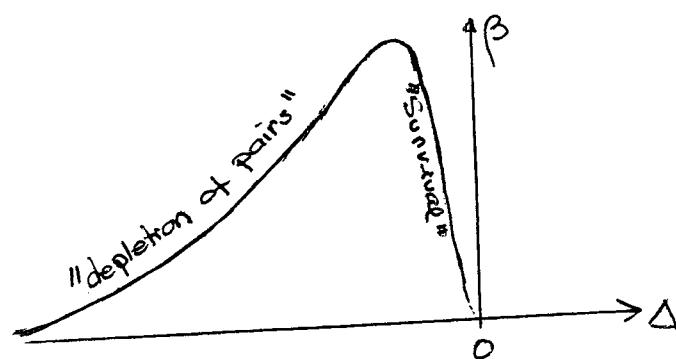
$\beta(\Delta) \rightarrow$  Potential and Survival



Case  $\Delta \gg \Gamma$

$$\beta \propto \left\{ \Delta^2 \sinh \left( \frac{\Delta z}{\Delta} \right)^{5/6} \right\}^{-1}$$

$\hbar \Delta z =$  energy gained by pair during a life-time ( $\beta$ )



Investigation of  $\beta(\Delta)$ , normally done using "catalyses"

- \* Extra laser at  $\Delta_c$ , introducing additional losses

$$\beta = \beta_t + \beta_c(\Delta_c) \frac{I_c(\Delta_c)}{I_{ref}}$$

+ trap alone      addition

- \* Measurement done by changing  $I_c(\Delta_c)$ , keeping total number constant

$$\beta_c(\Delta_c) \propto \frac{1}{I_c(\Delta_c)}$$

Typical results  $\rightarrow$  Good understand

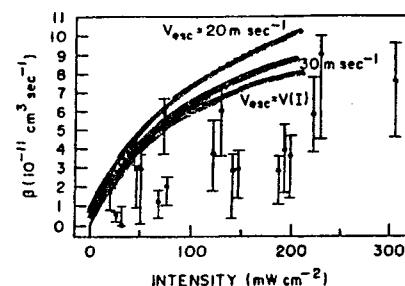


FIG. 28. Measurement of the trap loss rate constant as a function of MOT light intensity. The wider range of light intensity reveals an increase in  $\beta$  with MOT intensity. Dotted, full, and dot-dash curves are theory calculations using the optical-Bloch equation method of Band and Julienne (1992) with different assumptions about the maximum escape velocity of the MOT. The curve  $V_{esc} = V(I)$  results from a simple model in which the escape velocity is a function of the MOT intensity. Data are from Marcassa *et al.* (1993).

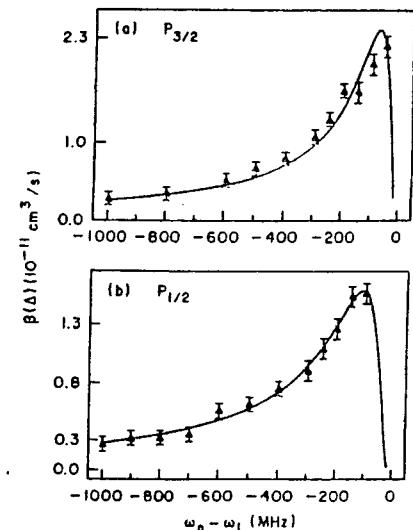


FIG. 29. Trap-loss spectrum in a Na MOT for catalysis laser detuned to the red from both atomic line-structure asymptotes. From Marcassa *et al.* (1993).

4-40

(Rb)

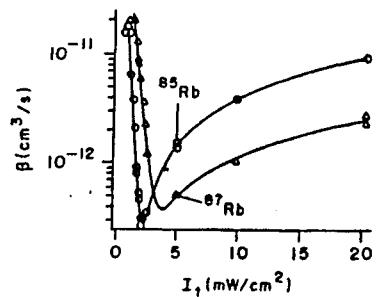


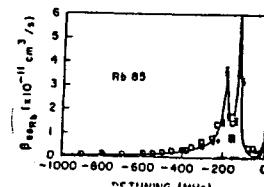
FIG. 32. Trap-loss spectrum for two isotopes of rubidium. Right-hand branch shows trap loss due to FCC and RE. Left-hand branch shows trap loss due to HCC. From Wallace *et al.* (1992).

- \* Differences between isotopes

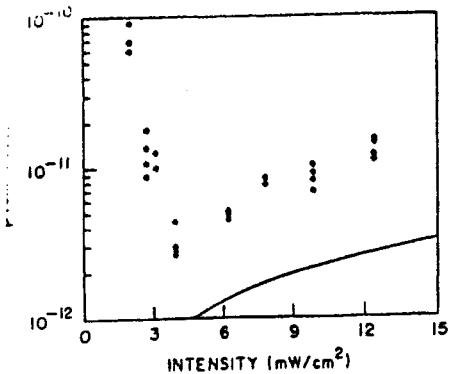
$$^{85}\text{Rb} \rightarrow \Delta E_{HF} \approx 3 \text{ GHz}$$

$$^{87}\text{Rb} \rightarrow \Delta E_{HF} \approx 6 \text{ GHz}$$

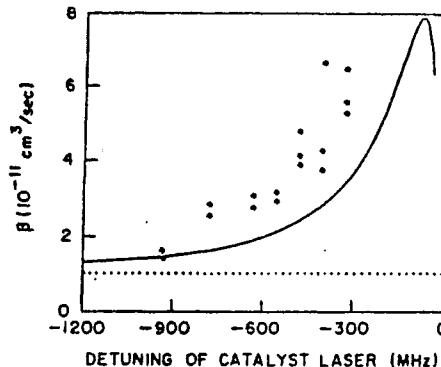
- \* Possible structure as function  $\Delta$



(Cs)



30. Trap-loss rate constant  $\beta$  as a function of MOT intensity for Cs collisions: comparison of experiment with theory, from Sesko *et al.* (1989).



(Li)

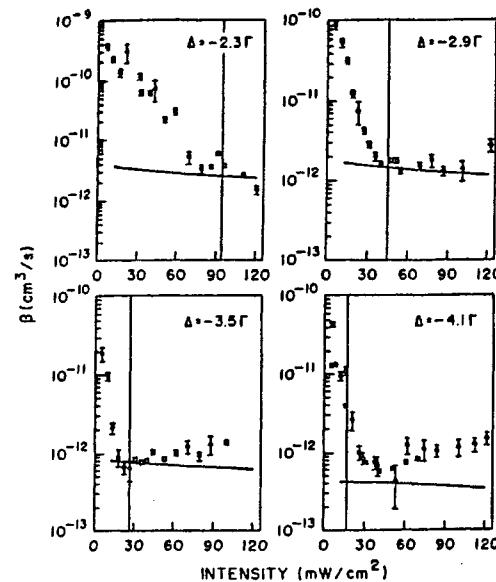
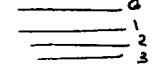


FIG. 40. Trap-loss rate constant  $\beta$  for four different detunings and a range of intensities in the Rice group's Li MOT. The vertical lines in each plot denote  $I_c$ , the calculated critical intensity required to recapture an atom released in the shallowest direction. From Ritchie *et al.* (1995).

- \* Excite state with inverted hyperfine



- ∴ Less influence of hyperfine spaghetti

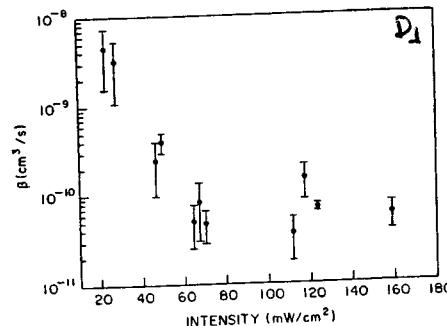
\*  $\Delta E_{FSC}$  is small only because important at low intensity  
 $\Delta E_{FSC} \sim 0.5K$

(K)

- \* It is possible to separate FSC and RE? (YES)

### First Experiment using D<sub>1</sub> line

(trap at D<sub>2</sub>, flip to D<sub>1</sub> observing decay)



- \* Absence FSC
- \* Larger values if compared with D<sub>2</sub> (D<sub>1</sub> trap much shallower)
- \* Equivalent low intensity behavior (but not value)

A  
interesting  
point

⇒ Special experiments to measure  $\beta_{FSC}$

- \* Pisa (Arimondo et al) → fluorescence in Cs
- \* Rochester (Bigelow et al) → fluorescence Cs
- \* Brasil → Fragment ionization in Rb

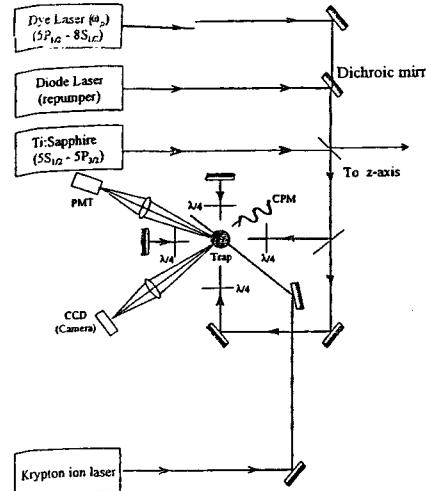


Fig. 1. Experimental setup. Rubidium is trapped in a MOT using a Ti:sapphire laser while a dye laser is tuned to the  $5P_{1/2} \rightarrow 8S_{1/2}$  transition and the krypton laser ionizes atoms out of the  $8S_{1/2}$  state. A channel electron multiplier detects the ions.

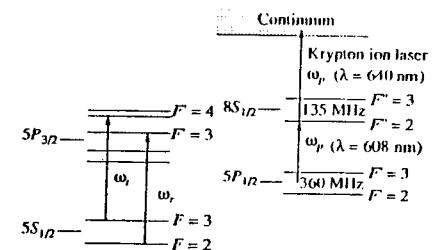
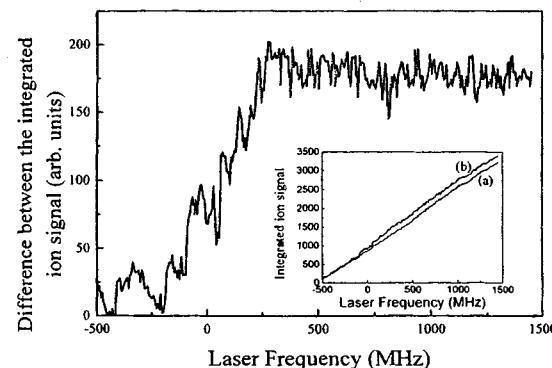


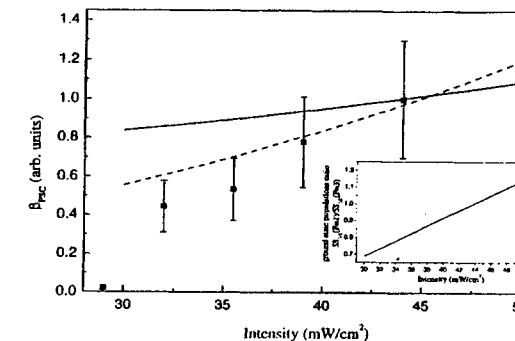
Fig. 2. Cw detection scheme. The produced  $5P_{1/2}$  atoms generated from FSC collisions are excited to  $8S_{1/2}$  state by a first photon from a dye laser ( $\lambda = 607$  nm) and then ionized by a second photon ( $\lambda = 640$  nm) provenient of a Krypton ion laser.

- Problems:
- \* Atoms  $P_{1/2}$  emerge with diff. velocities  
→ difficults for selective laser.
  - \* background ions

→ Integration.



## Investigation of $\beta_{FSC}$ with $\Delta$



From the produced ions

$$\beta_{FSC} \approx 0.15 \beta_{\text{Total}}$$

$\beta_E$  is dominant

Measurement by Rochester's group (Cs)

J.P. Shaffer et al.: Cs radiative escape and fine structure changing collision rates 329

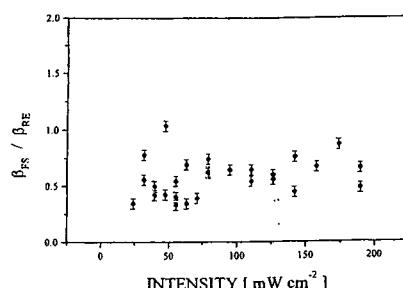


Fig. 6. The ratio  $\beta_{FSC} / \beta_{RSC}$  for a detuning of  $\Delta = -44\Gamma$ . The scatter in the plot gives a measure of the uncertainty in this value. The data points were fit to a constant to establish that  $\beta_{FSC} / \beta_{RSC} = 0.58 \pm 0.03$ .

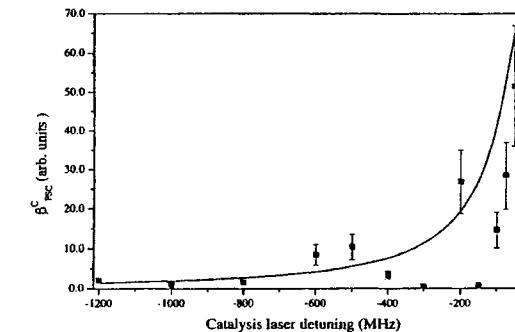


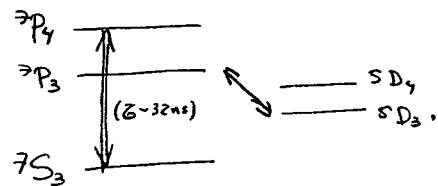
Fig. 3. Additional trap loss rate due to FSC introduced by the catalysis laser;  $\beta_{FSC}^c$ , as a function of catalysis laser detuning, operating below the transition  $5S_{1/2}(F=3) \rightarrow 5P_{3/2}(F'=4)$ ; in the frequency range of  $-1200 \text{ MHz} \leq \Delta_c \leq -50 \text{ MHz}$ . The full line represents the theoretical prediction.

About equal participation in Cs

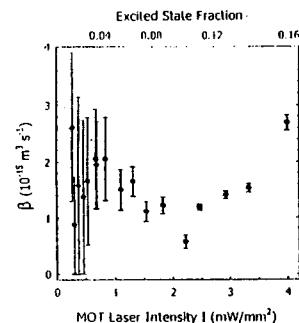
## Interesting new results in trap loss

Mot of Cr atoms

[Bradley et al PRA 61, 053407  
(2000)]



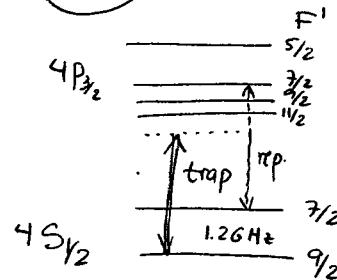
repumper  $\sim 660\text{nm}$   
trap  $\sim 426\text{nm}$   
 $I_s \sim 8.5\text{mW/cm}^2$



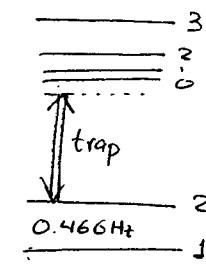
Clear minimum  
around  $2\text{mW/mm}^2$

- \* Shows increase at low intensity,  
but HCC is not present.

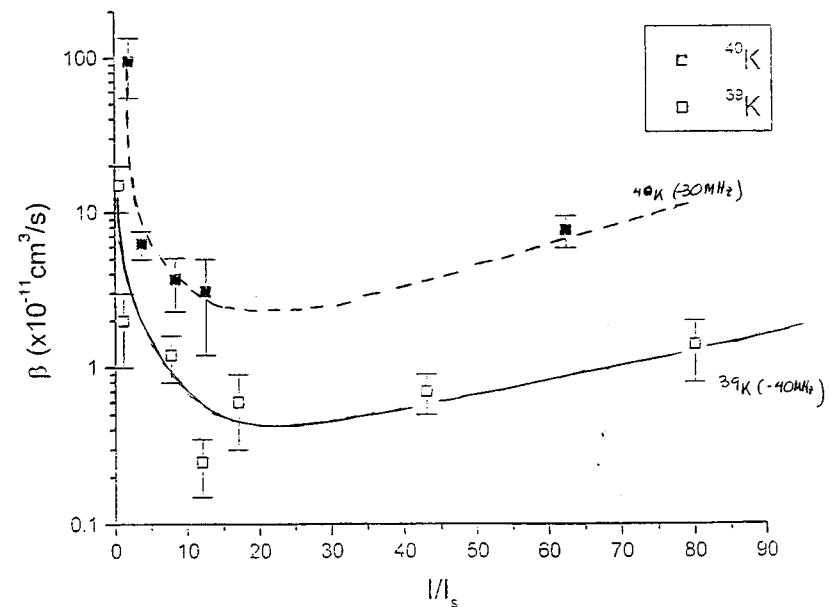
$^{40}\text{K}$  (Florence)



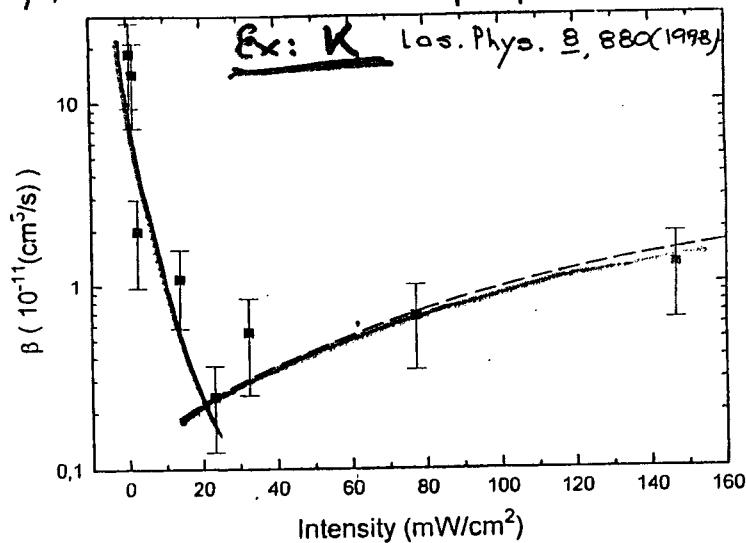
$^{40}\text{K}$   
(trapping at  
low hyperfine)



$^{39}\text{K}$   
(trapping at  
higher hyperfine)



\* Typical behavior of  $\beta$  vs I



Region I → { Radiative Escape (RE)  
{ Fine structure change (FSC)

Region II → Hyperfine change (HCC)

Recently <sup>FSC</sup>  
<sub>RE</sub>

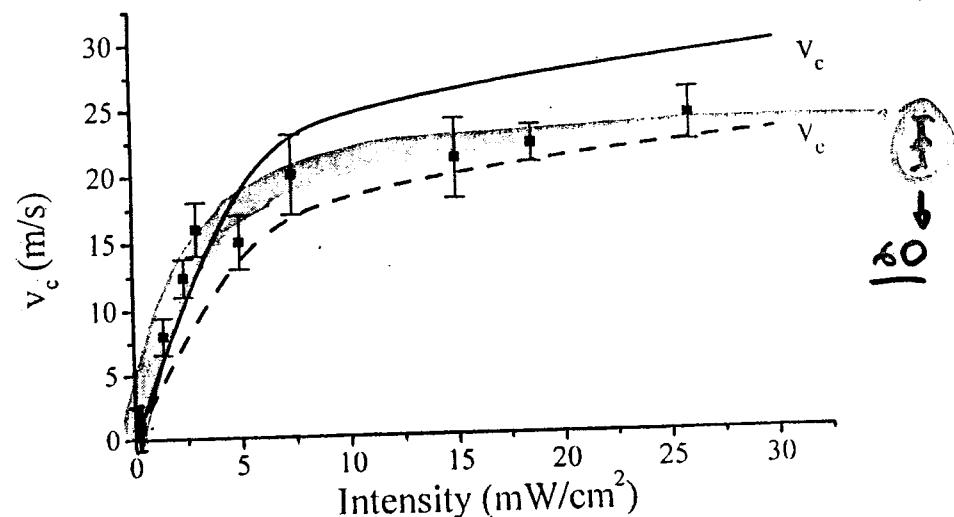
→ Trap loss is connected with  
Trap depth.

→ Trap depth is connected with  
escape velocity ( $v_e$ )

Knowledge of  $v_e$  is important  
in the interpretation of  $\beta$

II- Recently done experiments  
by our group:

\* Measurement of capture velocity ( $v_c$ )  
(Na) [Phys. Rev. A 62, 13404(2000)]



Observations

- How is  $v_c$  measured?
- Dependence with  $\frac{dB}{dt}$  and capture velocity  
(to be published)
- Nano operation  $W \sim .5 \text{ cm}$   
→ damping force dominant  
to  $v_c$

### \* Overlapping <sup>87</sup>Rb - BECs,

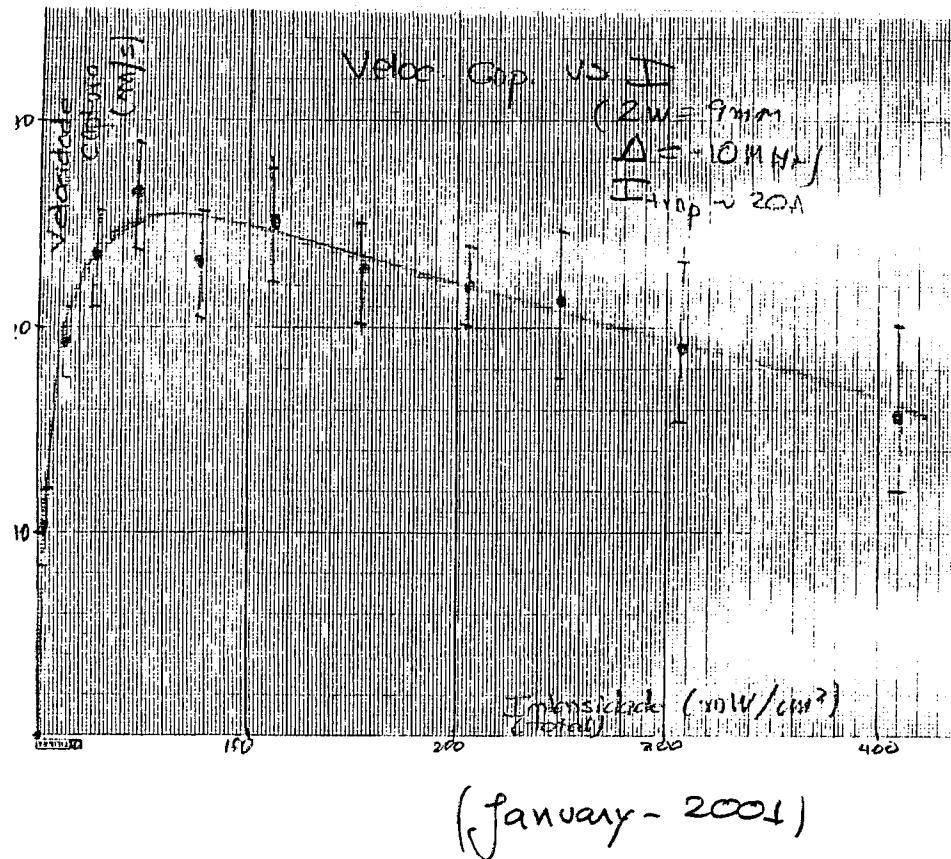
[ Myatt et al - Phys. Rev. Lett. 78, 586 (1997) ]

Overlapping BECs of mixed spin states  $|2,2\rangle$  and  $|1,-1\rangle$  shows very small spin exchange rate

$$\beta_{\text{eq}} \sim 2 \times 10^{-14} \text{ cm}^3/\text{s}$$

compared to values from MOTs

$$\beta \sim 10^{-11} \text{ cm}^3/\text{s}$$



### \* Counting coed Collisions

(Cs) [ Uebenhöhl, Kuhn, Frese, Meschede, Gomer  
 Un. Bonn - 1999 ]

- \* trapping small number
- \* directly observe two-body collisions
- \* Observation of low rate for HCC attributed to suppression due to repumper

## \* HCC in $^{87}\text{Rb}$ \*

[ Nesnidal and Walker  
Phys. Rev. A 62, 030701 (2000)]

- Detailed measurement at low intensity where HCC is dominant

$$\begin{array}{c} F=2 + F=2 \quad F=1 + F=1 \\ \swarrow \qquad \searrow \\ F=1 + F=2 \end{array}$$

RENEE C. NESNIDAL AND THAD O. WALKER

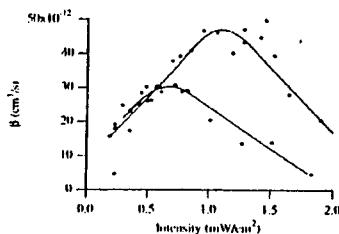


FIG. 1. Loss rates due to ultracold collisions as a function of total trap laser intensity with a detuning of  $\Delta = -11'$ . The circles and squares indicate data taken with a magnetic-field gradient of 10 and 18 G/cm, respectively. The solid lines have been included to guide the eye.

PHYSICAL REVIEW A 62 030701(R)

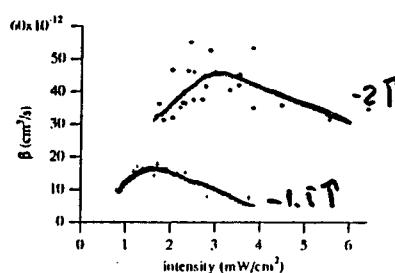
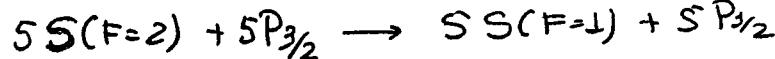


FIG. 2. Loss rates as a function of total trap laser intensity for larger laser detunings. The '+'s and circles indicate data taken at  $\Delta = -1.5T$  and  $-21T$ , respectively.

- Explanation :  $\rightarrow$  Spin-exchange in a hyperfine excited state



$\rightarrow$  Light modification of HCC (?)

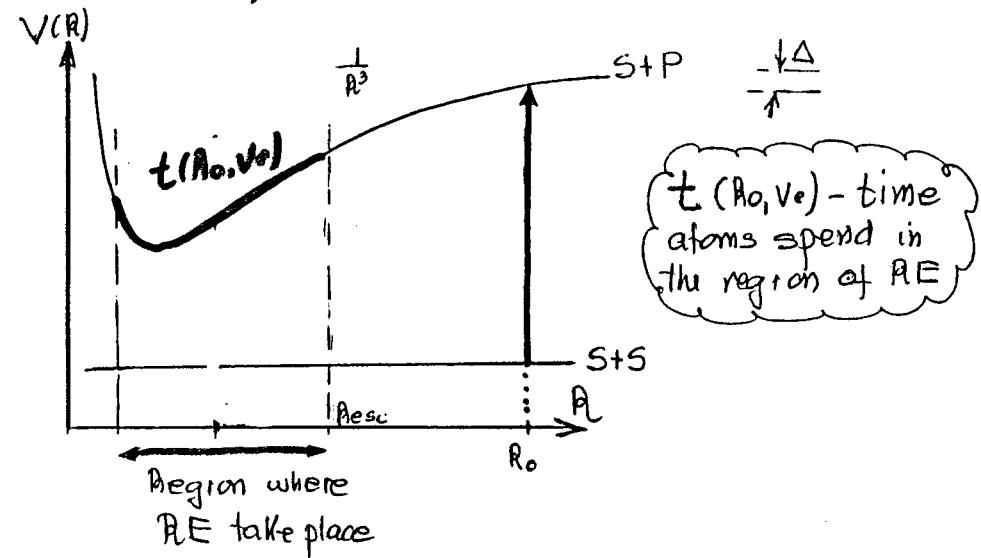
- tunning down  $\beta$  at low I.

## IV- Our new Model

- Gollagher-Pritchard (GP)
- $V_e$  (I) inferred from our previous experiment
- (We will apply to  $^{85}\text{Rb}$ )

$\Rightarrow$  Important to have the correct variation of  $V_e$  with intensity.

- Consider PE as dominant (based on previously done experiments)

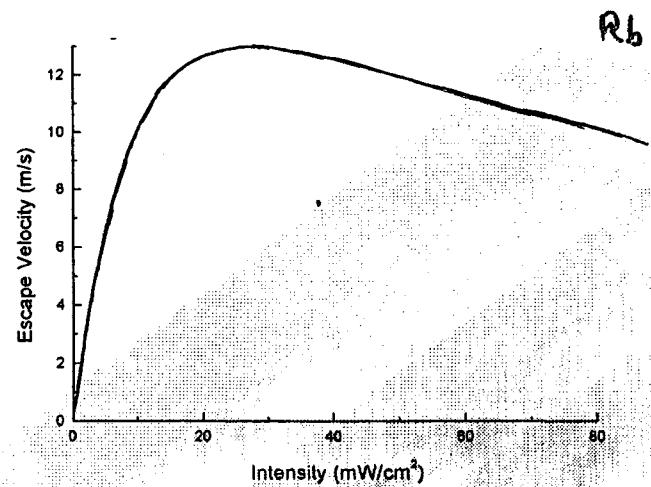


# Important terms for the model

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II-

## \* Escape Velocity ( $V_e$ )



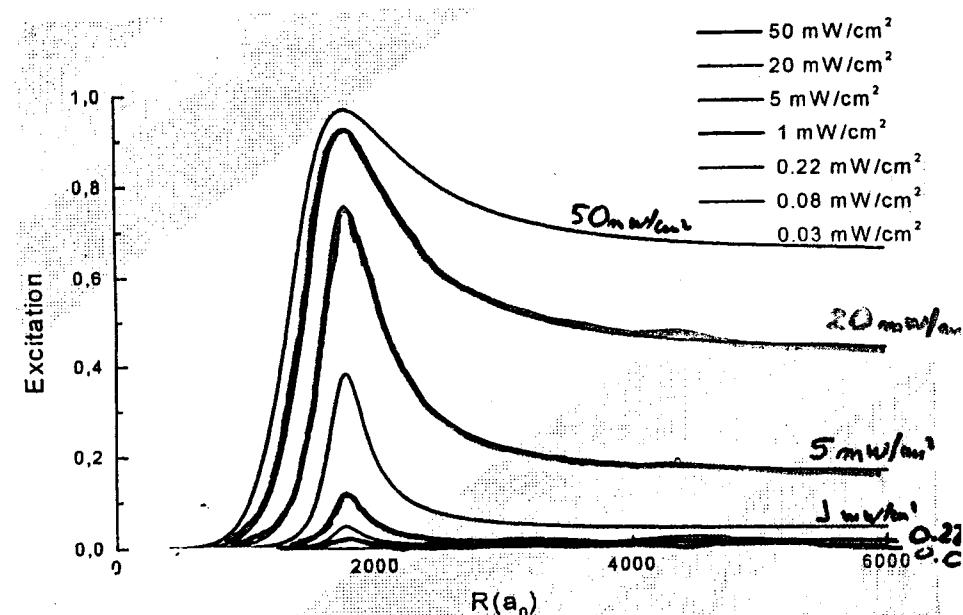
- \* Best fitting with experimental determination of ( $V_e$ )  $\Rightarrow$  damping as dominant process.

- \* Existence of maximum

$$\beta \propto \int_{\infty}^{\infty} 4\pi R_0^2 \epsilon(R_0, \omega_c, I) P_{RE}(R_0, V_e) dR_0$$

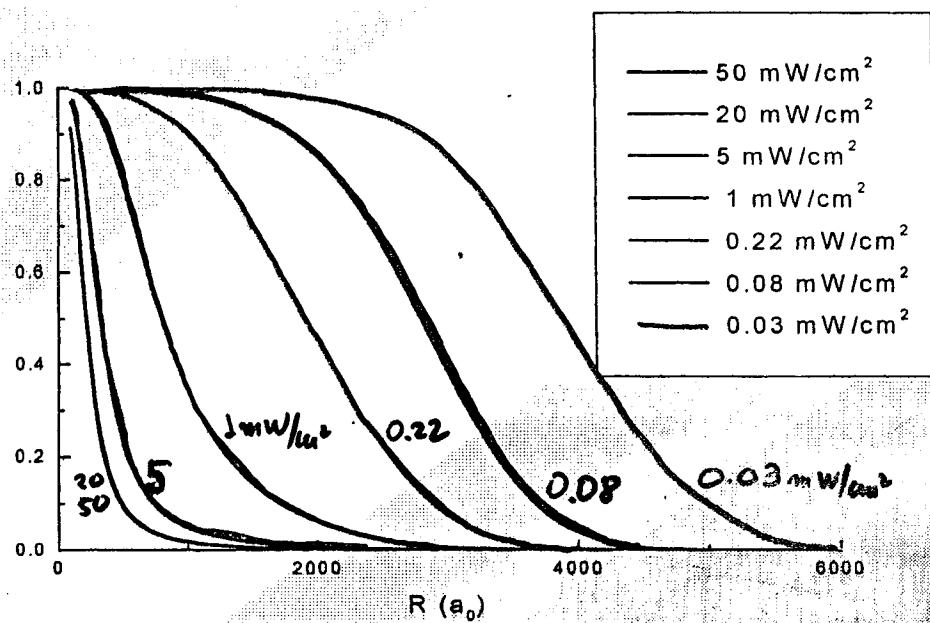
Let's observe each term

## \* excitation rate $\epsilon(R_0, \omega_c, I)$



- \*  $\epsilon$  is monotonic with  $I$
- \* asymmetry  $\rightarrow$  potential

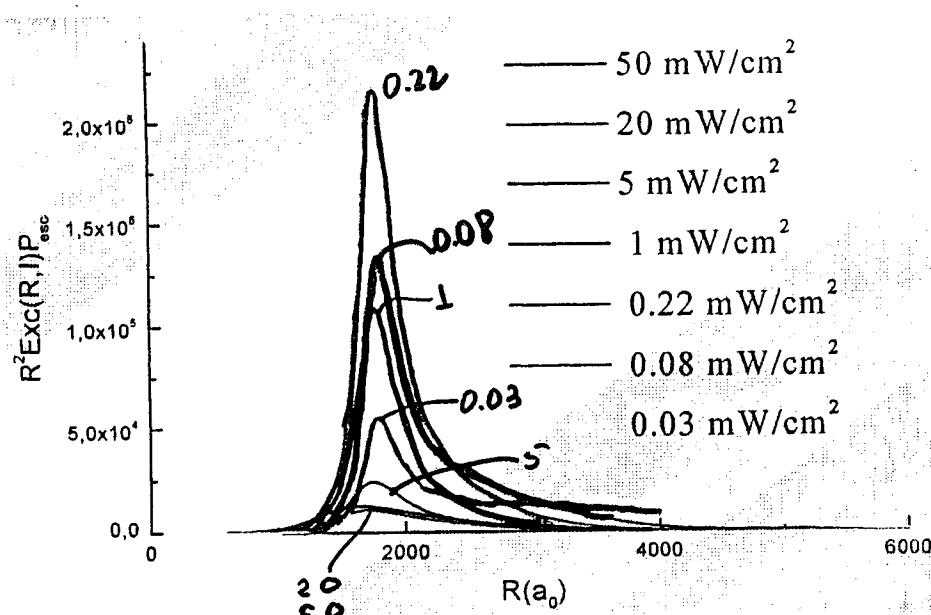
## \* Radiative Escape Probability (PRE)



II-36

\* | Overall product f (integrand)

II-37

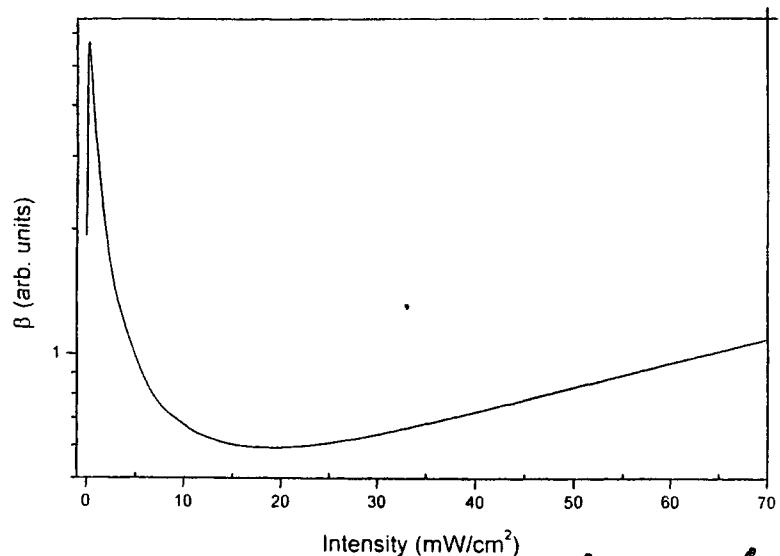


- \* PRE always decreases with  $R$   
(Large  $R$  smaller chance to survive)
- \* Higher intensity → increase  $V_e$   
atoms must survive to short range
- \* Intense variation at lower intensity  
(more pronounced than at higher intensities)

$\beta \propto \text{Integral}$   
(area)

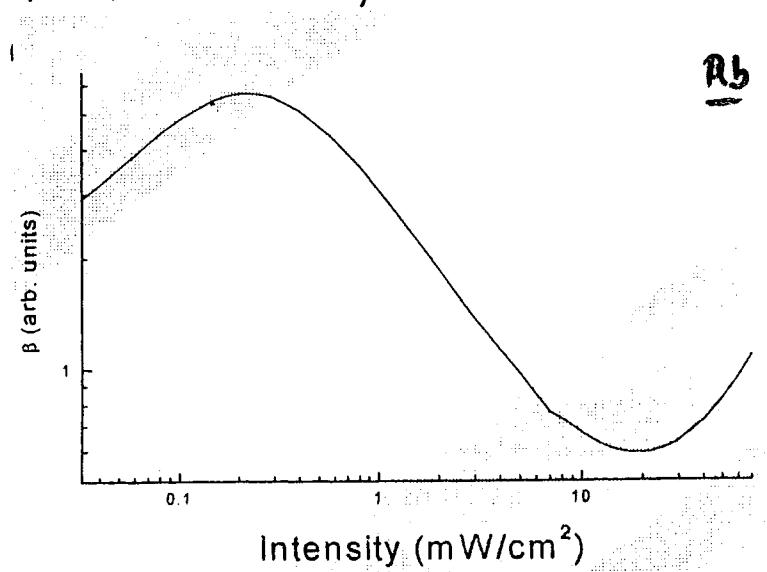
\* Overall behavior of  $\beta(I)$

II-3



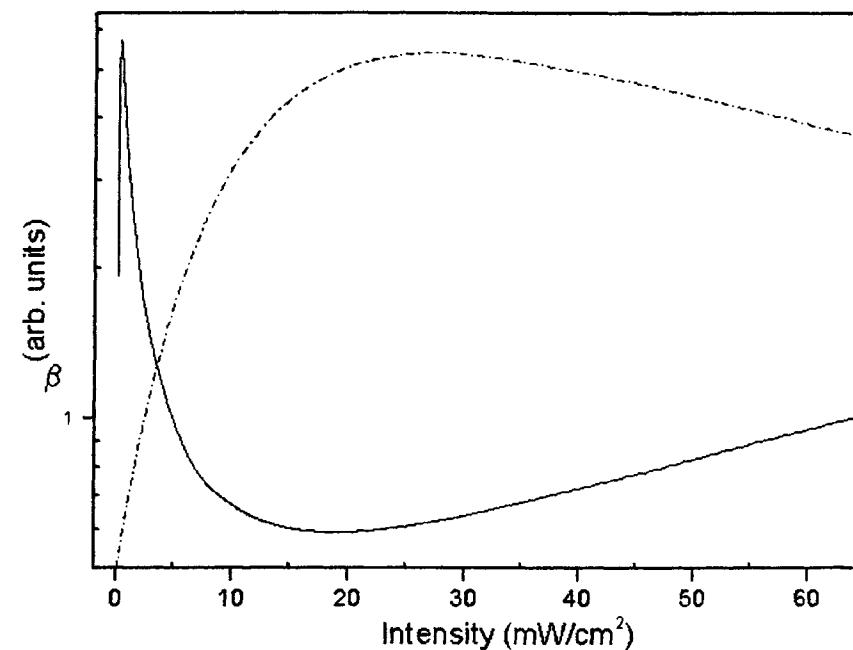
→ can fit well most part of  
existent  $\beta$  vs  $I$

\* At low intensity:



$$\frac{\partial \beta}{\partial I} \rightarrow 0$$

Fig.1: G.D. Telles et al., Alternative Interpretation for the Magneto Optical Trap Losses at Low Light Intensit



\* Is HCC necessary to explain trap loss?

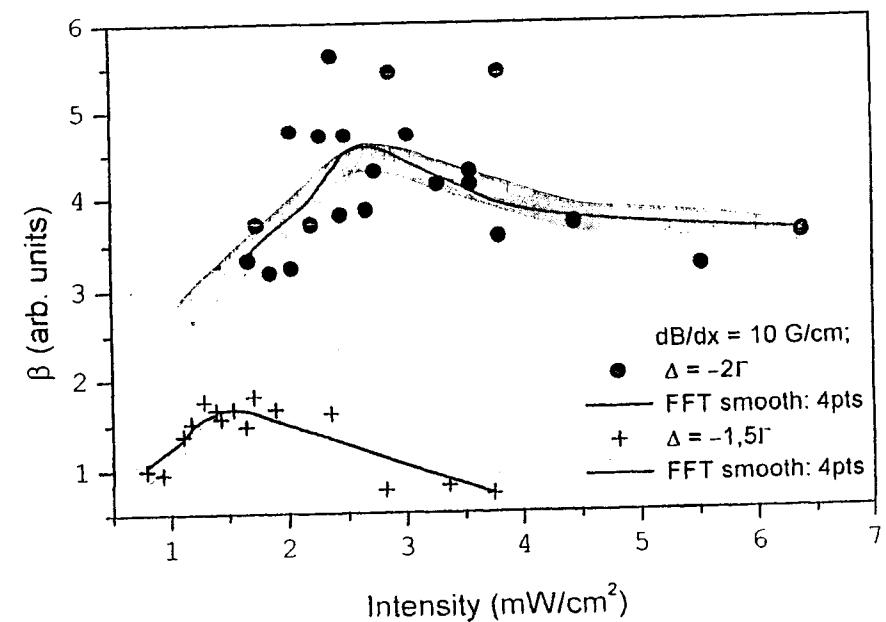
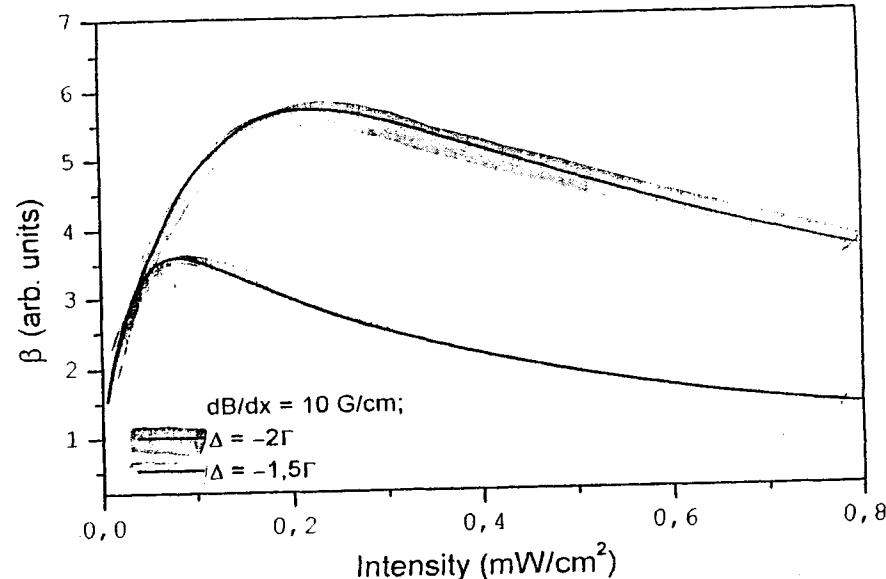
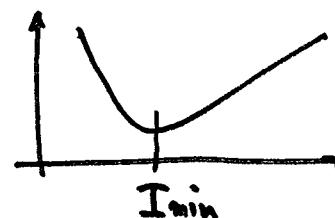
II-39

T

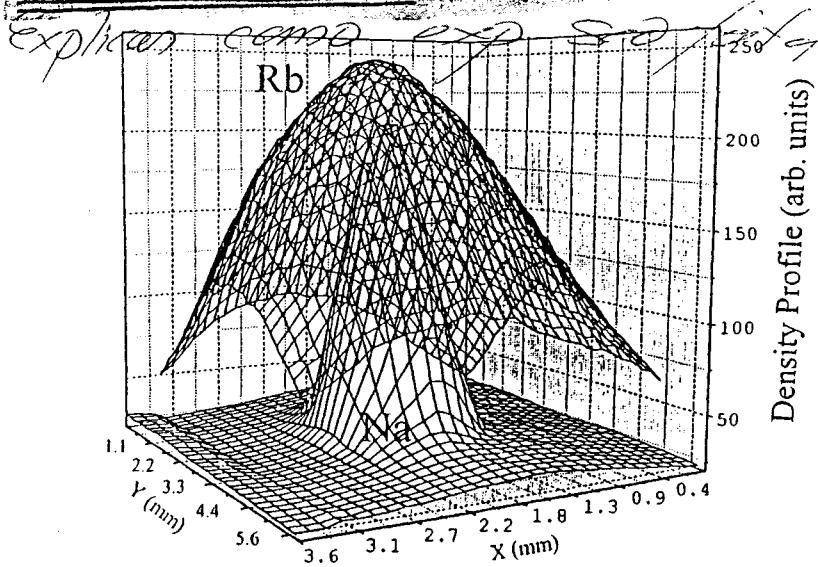
\* What is the real level of participation of HCC in trap loss at low intensity?

\* Comparison with experimental data → position of minimum

Species	$\Delta (\Gamma)$	$I_{\text{exp}} (\text{mW/cm}^2)$	$I_{\text{mod}} (\text{mW/cm}^2)$	Reference
Li	-3.5	28	25	N.W.M. Ritchie et al PRA 51, R890 (95)
Na	-1	5	7.9	S.R. Muniz et al, PRA 55, 4407 (97)
K	-7	80	81	L.G. Marcassa et al, accepted PRA (00)
Rb	-1	3	2.5	C.D. Wallace et al, PRL 69, 897 (92)
Cs	-1	4	1.4	Sesko et al, PRL 63, 961 (89)

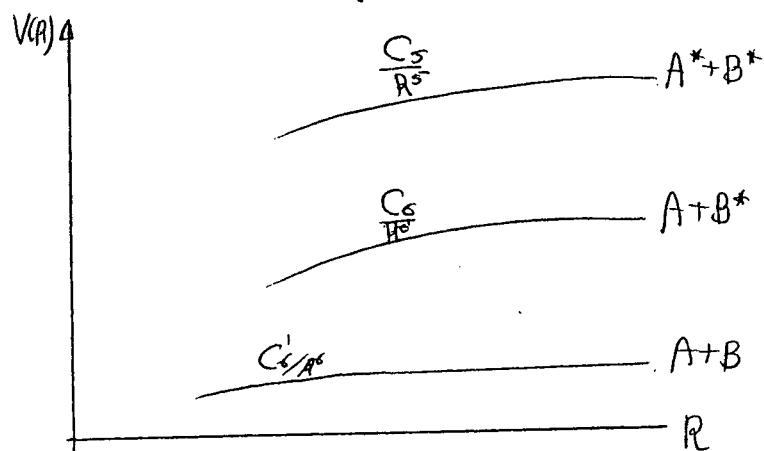


# Heteronuclear Systems



How we measure trap loss?

$\beta_{A/B}$  (loss of A in the presence of B)



II-41

- \* Mechanisms of loss RE, FSC (HCC)

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\*  $\beta_{A/B}$  as a function of  $I_A$

As  $I_A$  increases:
 

- \* Ve varies (increases)
- \* Population  $A^*$  increases

\* the behavior of  $\beta_{A/B}$  vs  $I_A$  will depend on:

- \* competition between  $Ve$ ,  $A^*$
- \* what mechanism is involved:  $A + B^*$ ;  $A^* + B$ ;  $A^* + B^*$
- \* relative  $C_s$  or  $C_s$  coeff.
- \*

\* Mass ratio  $\frac{m_A}{m_B}$  is important

\* Detuning

\* Each system  $\Rightarrow$  peculiarities

$$\frac{dN_A}{dt} = -\gamma N_A - (\beta) \text{Möller} - (\beta') \text{HMB cur}$$

Na/Rb

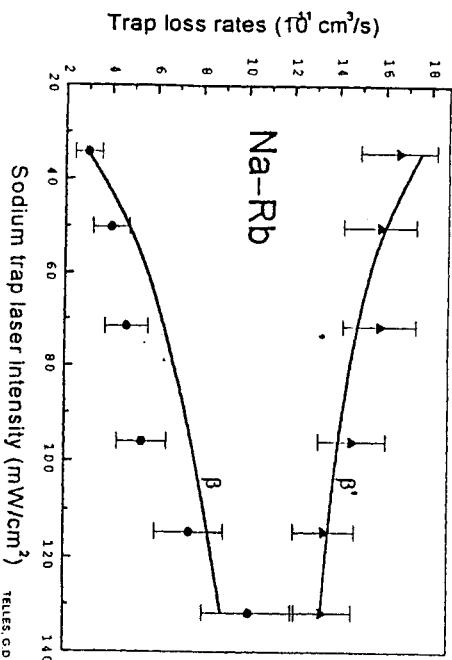
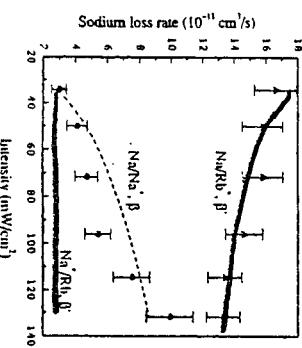
Phys. Rev. A 59, R23 (1999)

change AceB → change B'

K/Rb

Phys. Rev. A - Accepted (2000)

$\Delta \sim 40 \text{ MHz}$



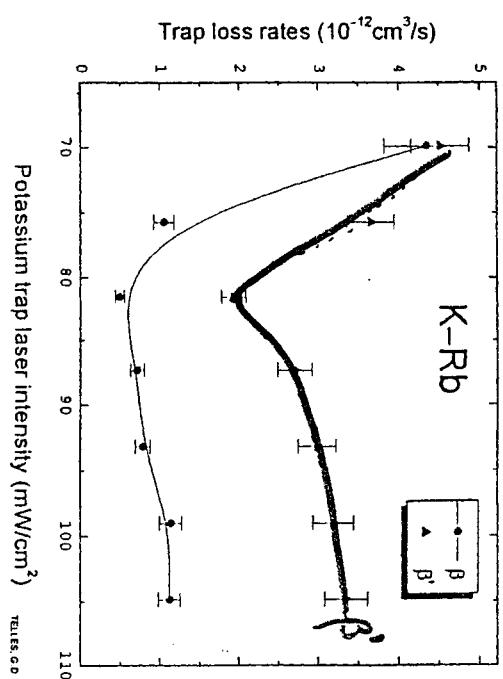
- \*  $\beta'$  and  $\beta$  → opposite behavior
- \* Model GP-like using RE
- \* Two possibilities {  $\begin{cases} \text{Na/Rb}^* \\ \text{Na}'/\text{Rb} \end{cases}$

$$C_S(\text{Na}/\text{Rb}^*) \approx 10 C_S(\text{Na}'/\text{Rb})$$

Na/Rb\* dominant

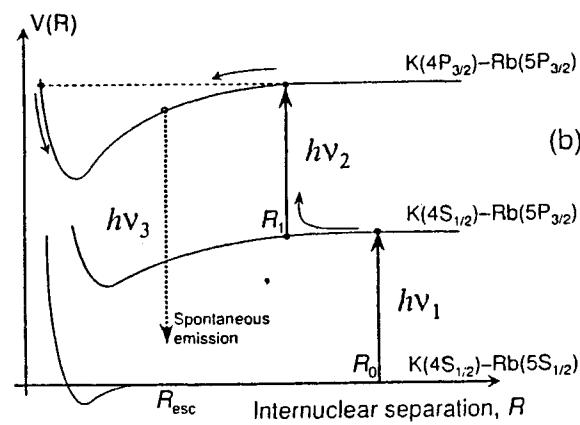
- \* K\*-Rb channel → All repulsive
- \* Only K-Rb\*  $\Rightarrow$  does not explain relative value
- \* Possibility of K\*-Rb\*

\* Low intensity we have interpreted in the paper as HCC (?)  
(New analyses)  $\equiv$



## Double excited participation:

II-45



## Reciprocity

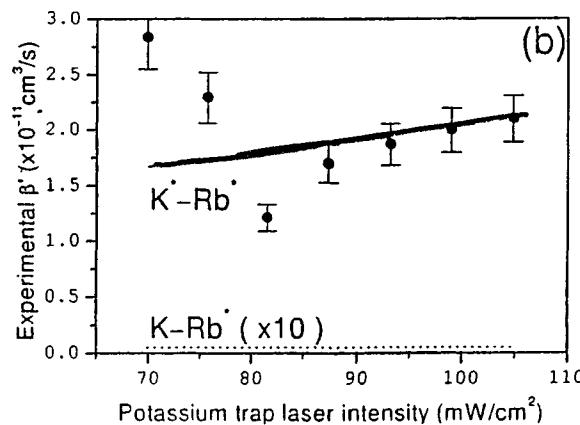
II-46

$$\beta'_{Rb/K} \ll \beta'_{K/Rb} \quad \text{at } (70-100\text{mW})$$

- \* FSC release  $\sim 100K$   $\rightarrow$  enough to eject both atoms

$\Rightarrow$  RE must be dominant

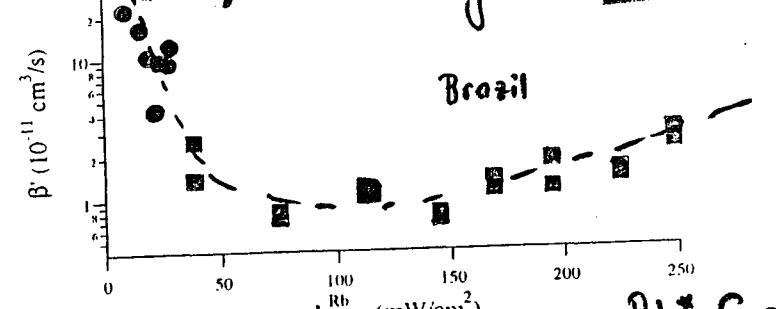
- \*  $\frac{\beta'_{K/Rb}}{\beta'_{K/K}} \sim 3$   $\left\{ \begin{array}{l} \text{- mass ratio} \\ \text{- survival probability} \\ \text{- HCC (?)} \end{array} \right.$   
(model OK)



$\Rightarrow K^* - Rb^*$  dominant

Rb/Cs Brazil + Pisa [Phys. Rev. A (2001)?]  
(process)

Italy good - behavior !!!



$$\beta'_{Rb/Cs} \sim 20 \beta'_{Cs/Rb}$$

$Rb^* - Cs \approx \text{impulse}$   
 $Rb - Cs^*$  (main channel)

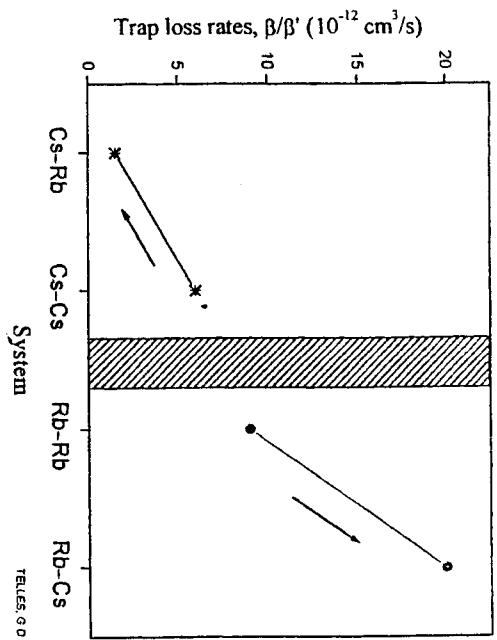
## Homo vs Hetero

J-47

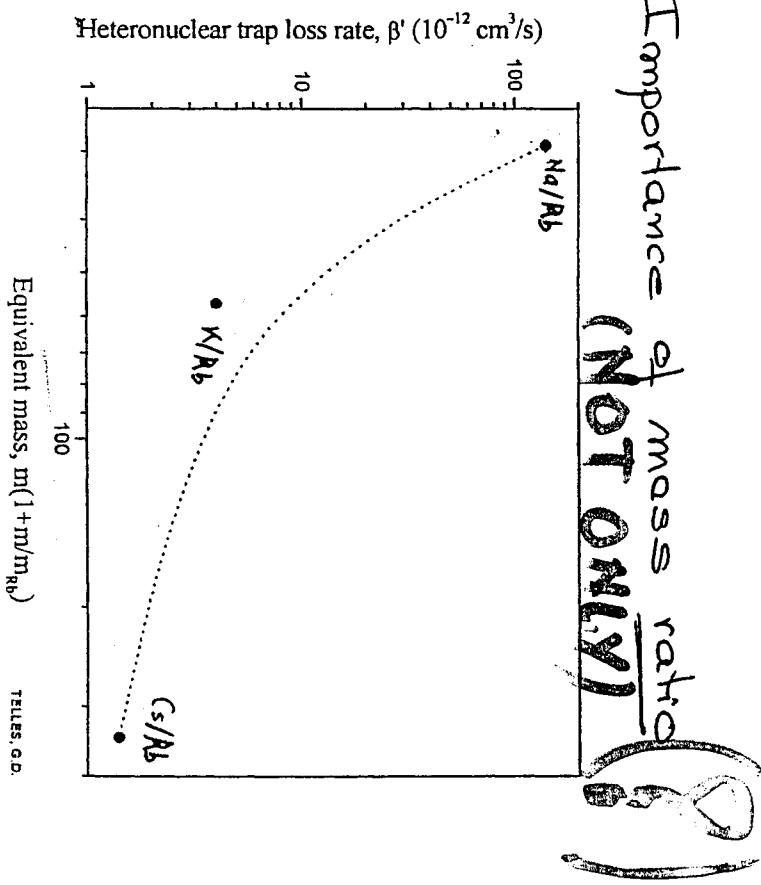
Na/K

Phys. Rev. A 52, R4340 (1995)  
Phys. Rev. A 60, 3892 (1999)

J-48



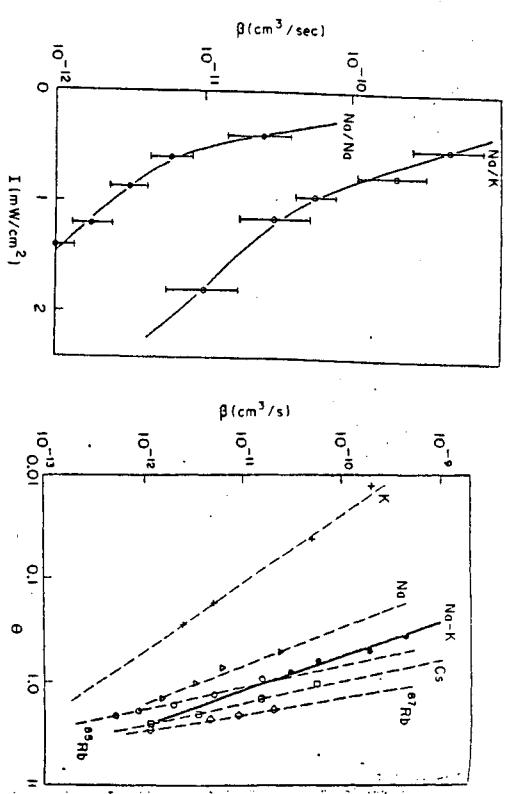
\* Importance of mass ratio (NOT ONLY)



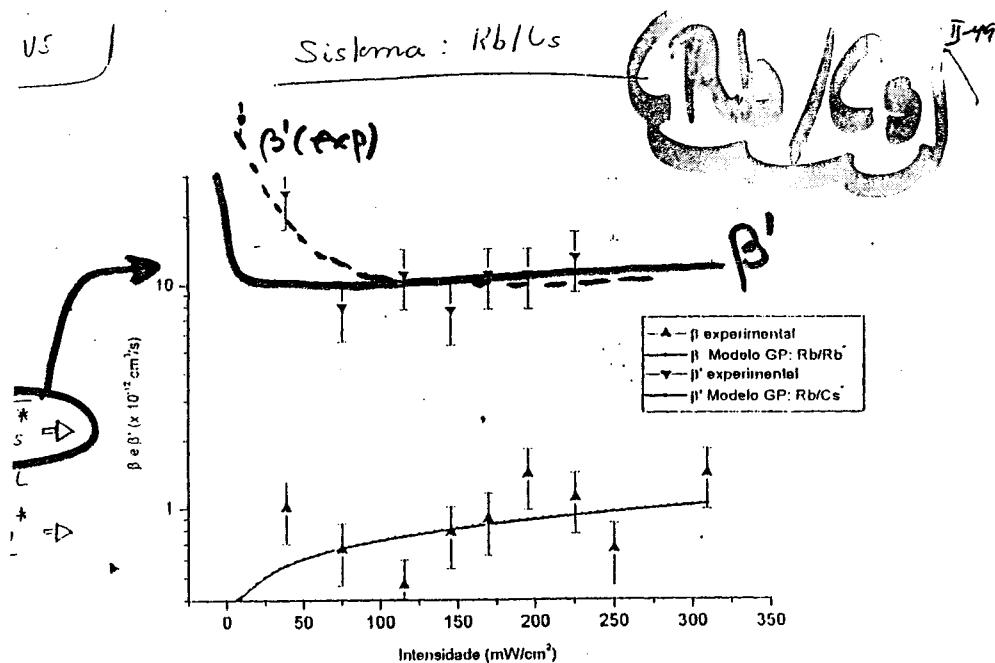
□

Rev. Mod. Phys. 71, 4 (1999)

$$\Theta \sim \frac{4}{\lambda^2} \frac{T^3}{\Delta^3} \frac{I}{I_{\text{sat}}} \frac{\omega}{m}$$



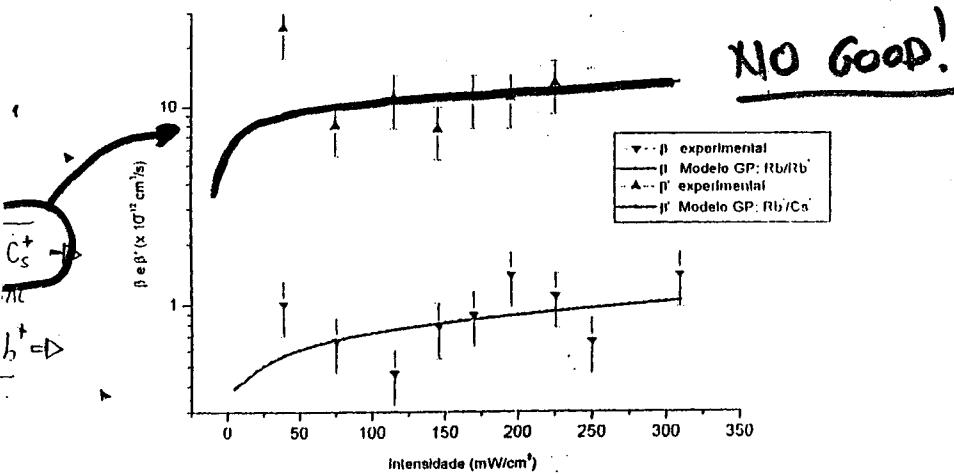
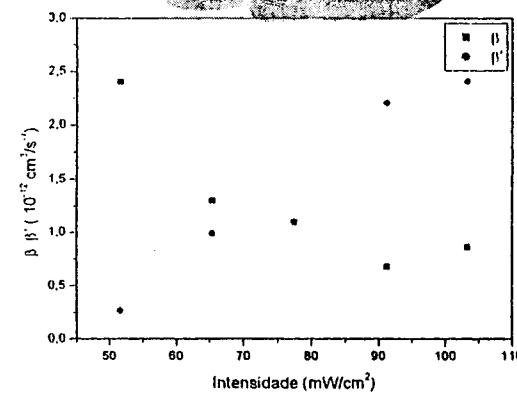
VS



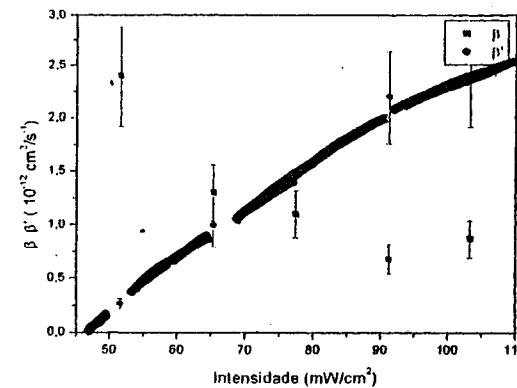
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K/K  
Cs/K



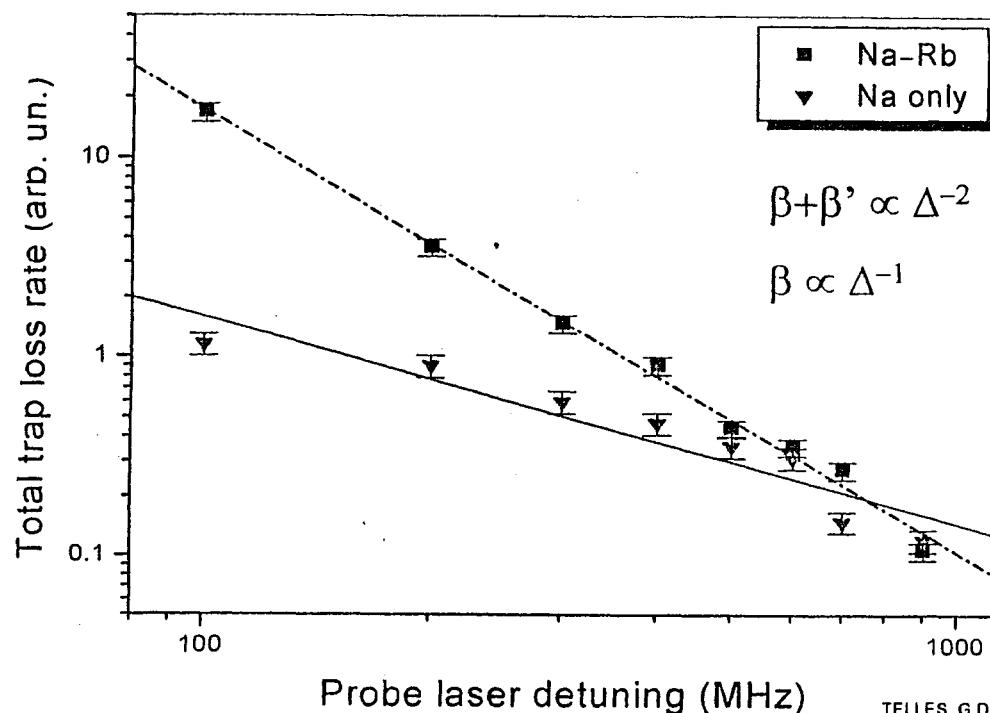
$\Delta = -\Delta \text{Hf}t$



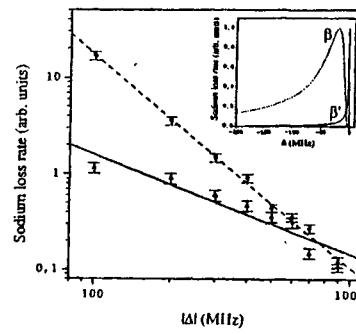
# Detuning Dependence

J-51

\* Extra laser  $\Rightarrow$



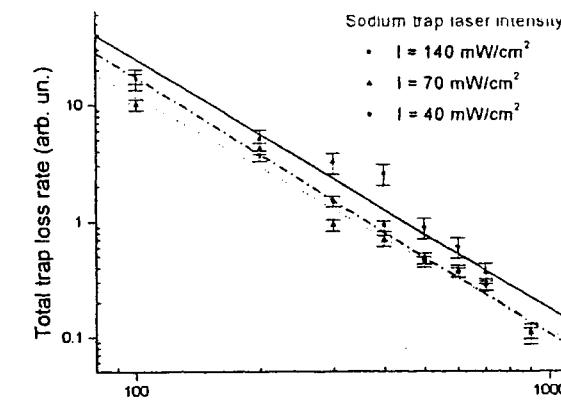
\* extra laser to excite Rb  $\rightarrow$  no variation



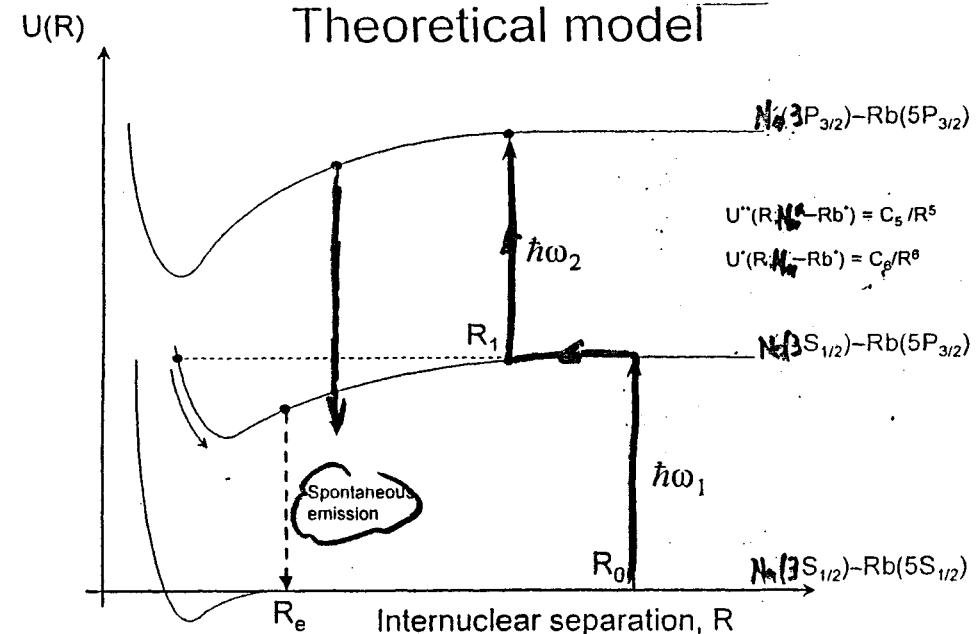
Sodium trap laser intensity:

- $I = 140 \text{ mW/cm}^2$
- $I = 70 \text{ mW/cm}^2$
- $I = 40 \text{ mW/cm}^2$

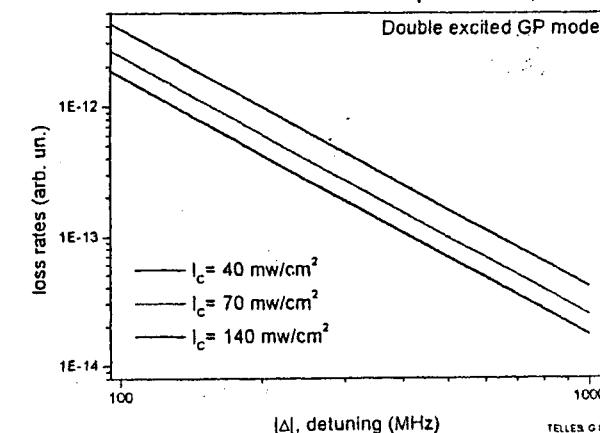
J-52



## Theoretical model



Double excited GP model



Lista publicações pelo grupo São Carlos  
em colisões frias

- Marcassa et al - Phys. Rev. A, R4583 (1993)
- Santos et al - Phys. Rev. A, 52, R4340 (1995)
- Bagnato et al - Los. Phys. 4, 1062 (1994)
- Bagnato et al - Phys. Rev. Lett. 70, 3225 (1993)
- Bagnato et al - Phys. Rev. A, R2523 (1993)
- Napolitano et al - Phys. Rev. Lett. 73, 1352 (1994)
- Marcassa et al - Phys. Rev. Lett. 73, 1911 (1994)
- Marcassa et al - Phys. Rev. A 52, R913 (1995)
- Marcassa et al - J. Phys. B 29, 3051 (1996)
- Zilio et al - Phys. Rev. Lett. 76, 2033 (1996)
- Muniz et al - Phys. Rev. A, 55, 4407 (1997)
- Marcassa et al - Braz. J. Phys. 27, 238 (1997)
- Santos et al - Los. Phys. 8, 850 (1998)
- Bagnato -Weiner - Science Sp 7, 50 (1996)
- S. Miranda et al - Phys. Rev. A 59, 882 (1999)
- G. Telles et al - Phys. Rev. A 59, R23 (1999)
- Weiner, Bagnato, Zilio, Julienne - Rev. Mod. Phys. 71, 1 (1999)
- Marcassa et al - Eur. J. Phys. D 7, 317 (1999)