

From 1D to D (and beyond !)

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P. Donohue (Orsay)

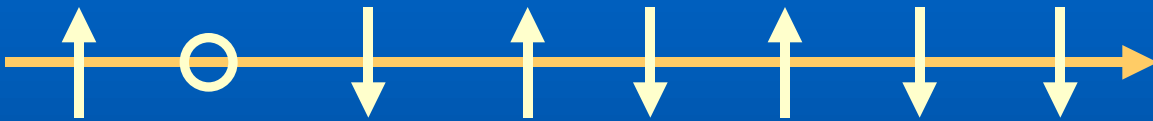
A. Georges (ENS)

A. Lichtenstein (Nijmegen)

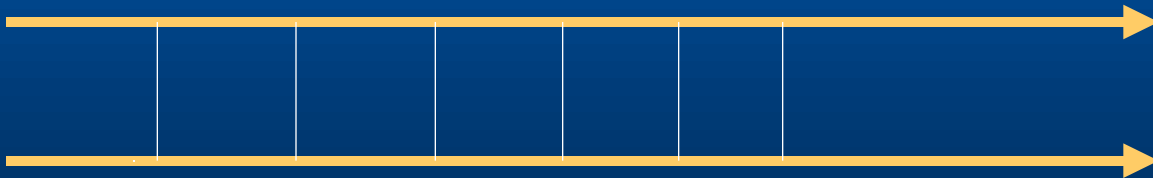
A. Lopatin (Rutgers)

E. Orignac (ENS)

Question

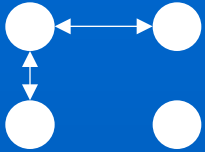


- 1 chain : Luttinger Liquid



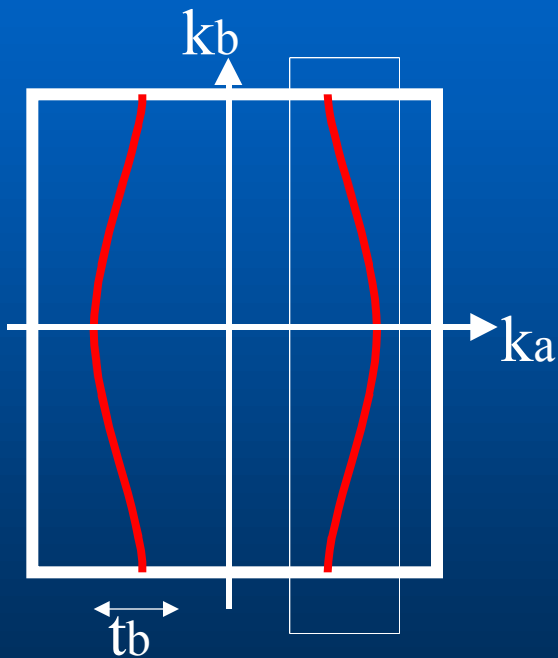
- Coupled chain : Fermi liquid ????

How to go from 1d to 3d



$$t_a > t_b > t_c$$

3000K, 300K, 20K



- High Energy (T, ω): **1D**
- Low Energy (T, ω): **2D, 3D**

Dimensional crossover

How is it modified by interactions ? T^* ?

- Renormalization arguments

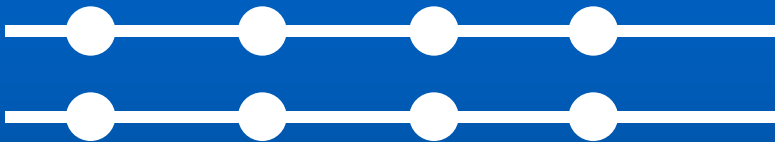
(Bourbonnais, Brazovskii+Yakovenko, Schulz)

$$E^* \sim t_{\perp} (t_{\perp}/t)^{\alpha/(1-\alpha)}$$

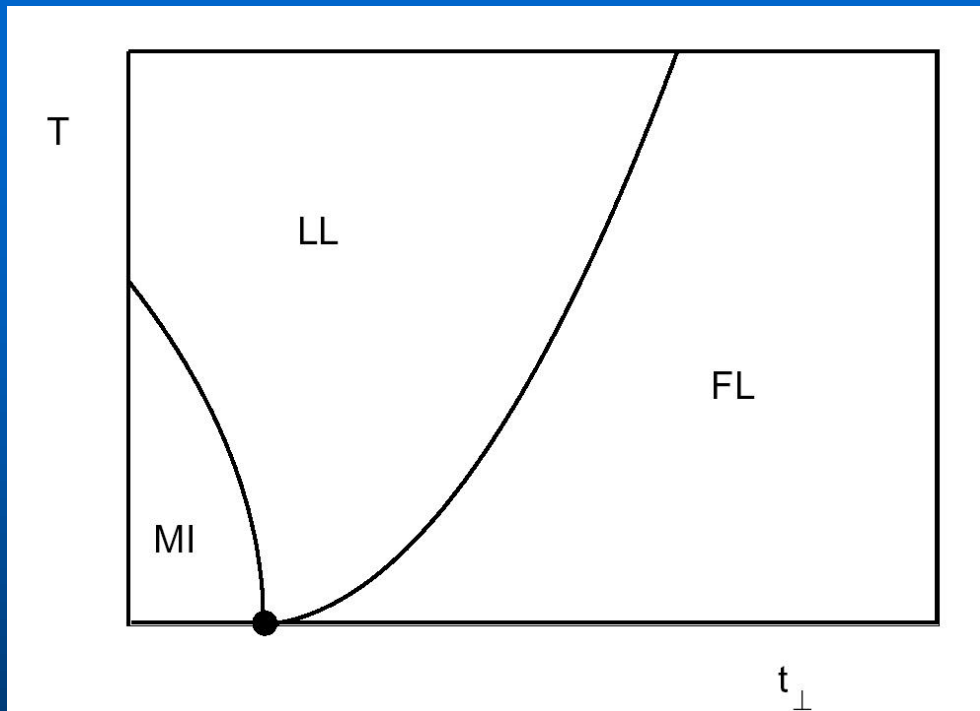
$$\alpha = \frac{1}{4} (K_{\rho} + 1/K_{\rho}) - \frac{1}{2}$$

Strong reduction of crossover temperature. But hopping still relevant !

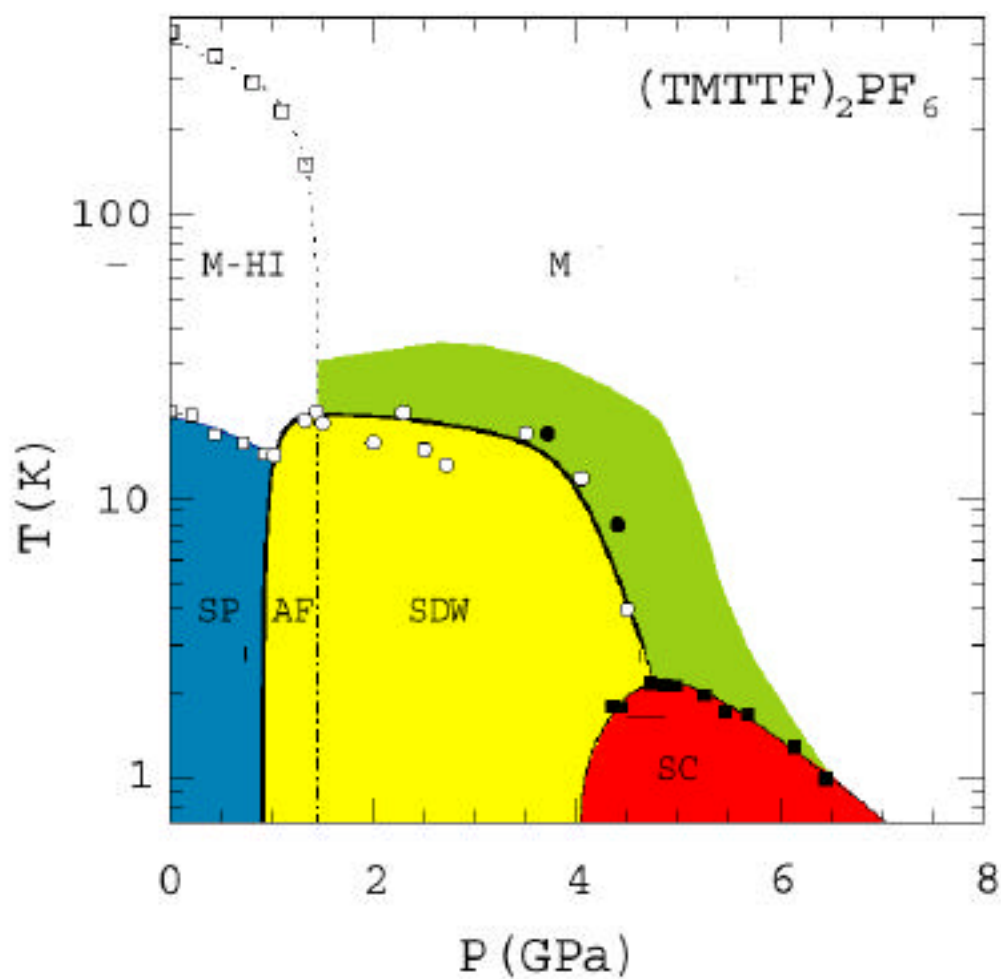
Mott insulators: confinement



Competition Mott insulator/Interchain hopping



- How to study ?
- Difficult (RG, RPA, etc.)

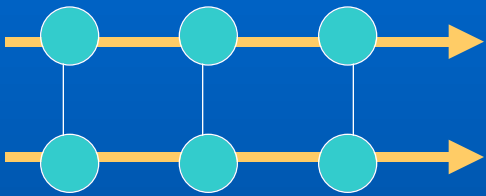


After D. Jérôme

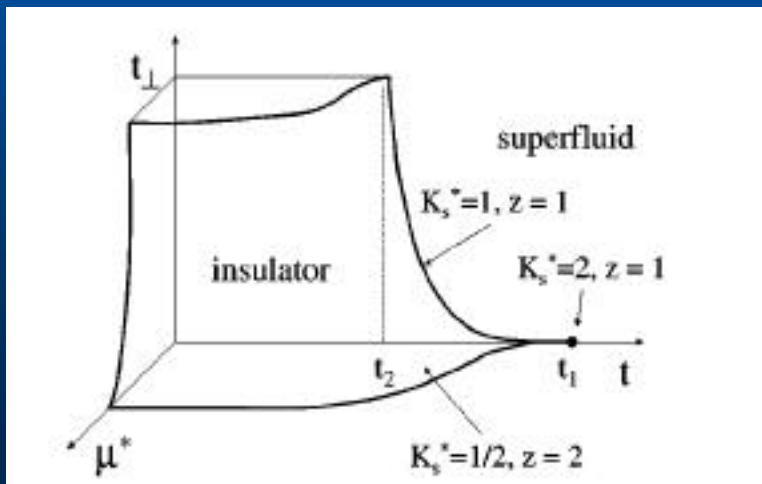
Ladders

- 2 one dimensional chains
- « Simple » analytical solution
- NO deconfinement for fermionic ladders : always insulating
- Interesting phases : Orbital antiferromagnet

Bosonic Ladders: deconfinement

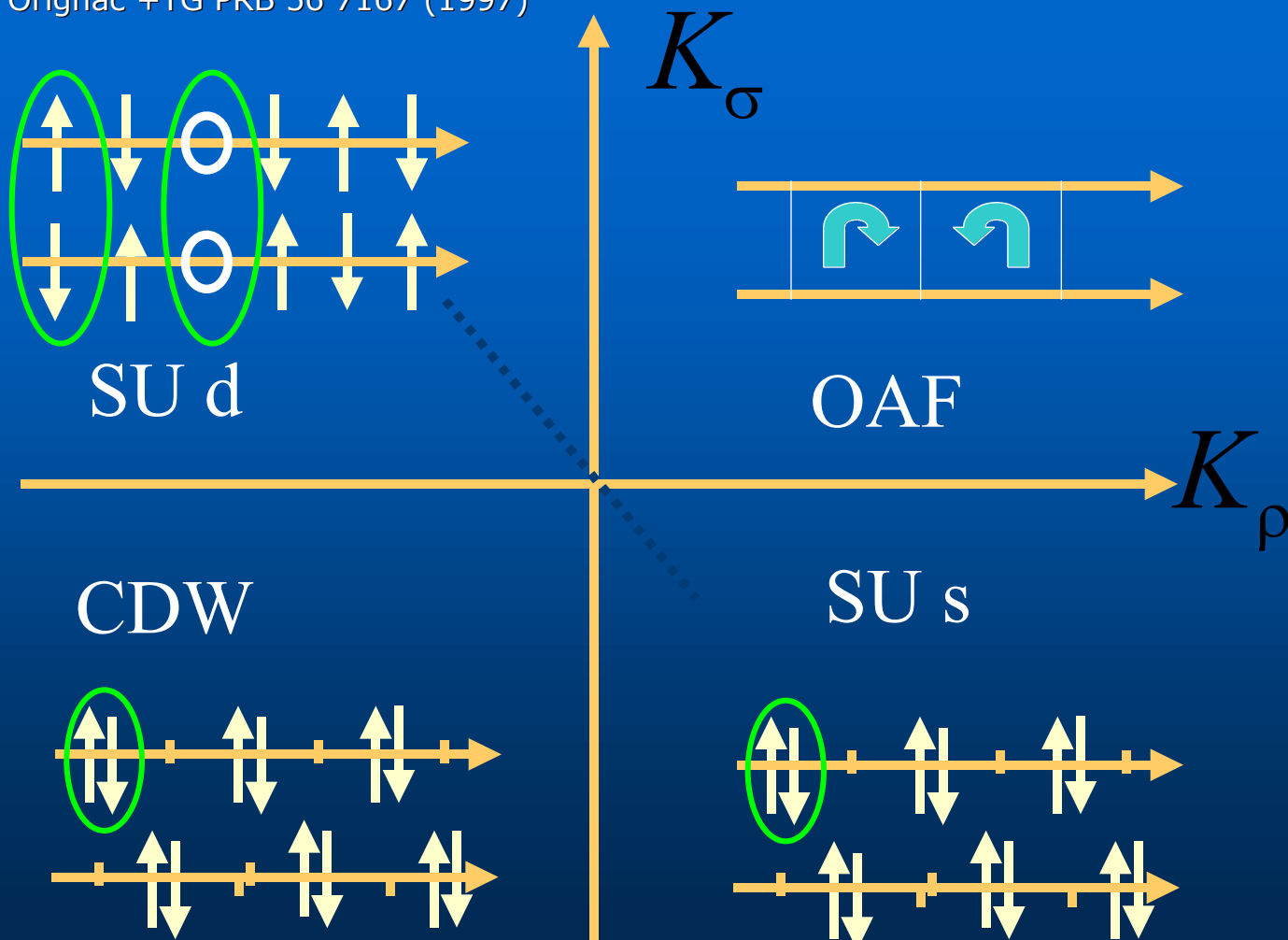


P. Donohue + TG PRB 63 180505(R) (2001)



Critical
properties of
the transition
?

E. Orignac +TG PRB 56 7167 (1997)



Infinite number of chains

S. Biermann, A. Georges, TG, A. Lichtenstein, cond-mat 0201542



$$S_{\text{eff}} = - \int \int_0^\beta d\tau d\tau' \sum_{ij,\sigma} c_{i\sigma}^+(\tau) \mathcal{G}_0^{-1}(i-j, \tau - \tau') c_{j\sigma}(\tau') \\ + \int_0^\beta d\tau H_{\text{ID}}^{\text{int}}[\{c_{i\sigma}, c_{i\sigma}^+\}], \quad (4)$$

$$G(k, i\omega_n) = \int d\epsilon_\perp \frac{D(\epsilon_\perp)}{i\omega_n + \mu - \epsilon_k - \Sigma(i\omega_n, k) - \epsilon_\perp},$$

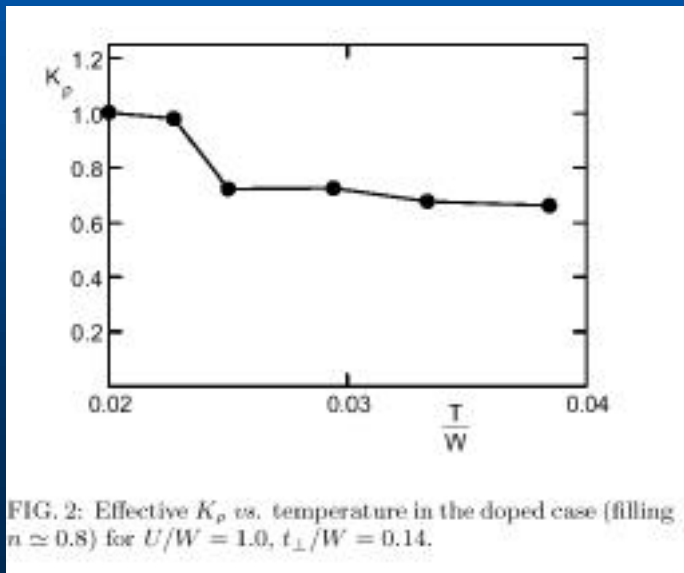
$$\text{Re } \sigma_\perp(\omega, T) \propto t_\perp^2 \int d\epsilon_\perp D(\epsilon_\perp) \int \frac{dk}{2\pi} \int d\omega' A(\epsilon_\perp, k, \omega') \\ \times A(\epsilon_\perp, k, \omega + \omega') \frac{f(\omega') - f(\omega' + \omega)}{\omega}.$$

- Self consistent theory for Σ
- Feedback of $t_{\perp}$ in Σ (a priori important for deconfinement)
- Different from RPA

$$G(k, k_{\perp}, i\omega_n) = \frac{1}{i\omega_n - \varepsilon_k - \varepsilon_{\perp}(k_{\perp}) - \Sigma_{1D}(k, i\omega_n)}$$

- Difficult to solve the equations analytically

Incommensurate case



$$T^* \propto \frac{t_\perp}{\pi} \sqrt{\frac{t_\perp}{t}} \sqrt{\frac{\theta}{1-\theta}}$$

$$T^* \propto 0.5 \frac{t_\perp}{\pi}$$

Commensurate case

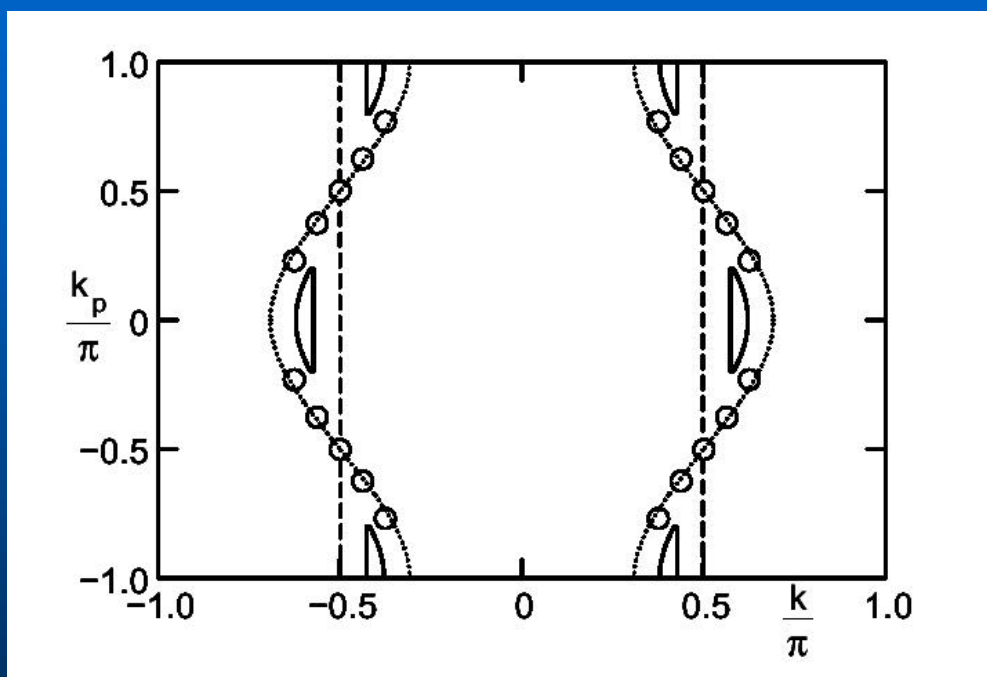


TABLE II: Effective K_{ρ} at half-filling, as a function of $t_{\perp}/W$ for $U/W = 0.65$ and $T/W = 1/40$.

$t_{\perp}/W$	0.00	0.04	0.07	0.11	0.14	0.16	0.18
K_{ρ}	0.00	0.02	1.01	1.09	1.07	1.06	1.04

$$t_{\perp}^* \cup \Delta_{1D}$$

Fermi Surface



Z

$k_{\perp}/\pi$	0.23	0.38	0.50	0.62	0.77
$Z(k_{\perp})$	0.78	0.77	0.76	0.77	0.79

Hall Effect

A. Lopatin, A. Georges, TG PRB 62 (00)

If no momentum relaxation (no umklapp)
along the chains:

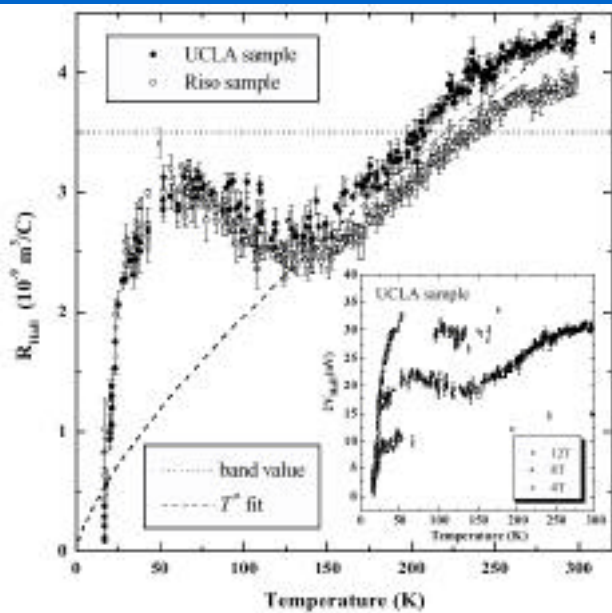
$$\rho_{yx} = \frac{H}{nec} \frac{2\alpha k_F}{v_F}.$$

With umklapp : scaling expression for Hall

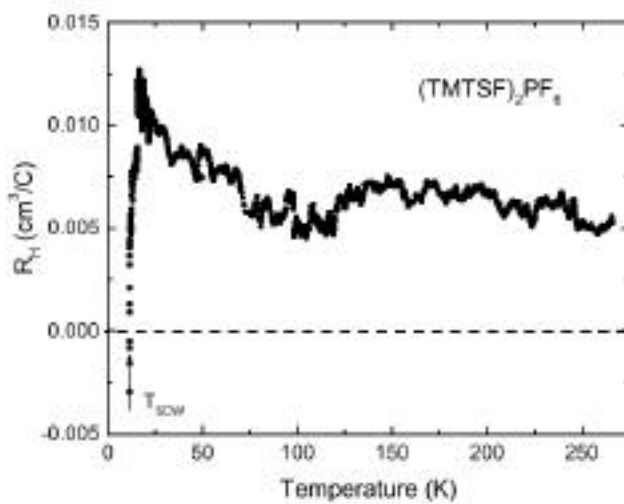
$$R_H(T) = R_H^0 \left(1 + a_1 \frac{g_3}{T^{1/x}} + a_2 \frac{g_3^2}{T^{2/x}} + \dots \right)$$

$$2x = 1/(1 - n^2 K_\rho)$$

Plot of R_h vs ρ_{xx}/T



J. Moser et al., PRL 84
2674 (00)



G. Mihaly et al., PRL 84
2670 (00)

Conclusions

- Fascinating (and poorly understood yet) crossover between 1D and higher dimensions
- Relevant for experimental systems
- Competition hopping/interactions : confinement
- chDMFT Good method to tackle the dimensional crossover