High-field entropy transport in cuprates

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A TALE OF TWO VELOCITIES!

(Durst, Lee '99)

An anisotropic Dirac cone



Excitation spectrum in

the vicinity of a node:

E(k) =
$$(\epsilon_k^2 + \Delta_k^2)^{1/2}$$

= $(v_F^2 k_1^2 + v_2^2 k_2^2)^{1/2}$

Fermi velocity : $v_F = d\epsilon_k/dk_1$ Gap velocity: $v_2 = d\Delta_k/dk_2$



 $\kappa_{00}/T = (nk_B^2/3) (v_F / v_2)$

Residual quasi-particle conductivity in optimallydoped cuprates





The residual quasi-particle conductivity is absent or dramatically reduced in underdoped cuprates!





The Wiedemann-Franz law

• Relates the conduction of heat to conduction of charge in an electronic system:

$$\kappa \rho / T = L_0 = \pi^2 / 3 (k_B / e)^2$$

- A fundamental property of any Fermi liquid at T=0. Theoretically robust in presence of strong correlations.
- Experimentally checked in a variety of stronglycorrelated electron systems (Heavy fermions, organic superconductors, Sr₂RuO₄...).
- Would break down in case of electron fractionalization.



Two distinct violations: i) An excess of electronic heat conductivity above 0.3K ii) sudden vanishing below 0.3K

Recent theoretical scenarios for the violation of WF law in various contexts

- Mei-Rong Li & E. Orignac, cond-mat/0201291(2002) (disordered Luttinger liquid)
- Guo-Zhu Liu, Cond-mat/0203048 (2002) (chiral symmetry breaking)
- Wonkee Kim & J. P. Carbotte, Cond-mat/0202514(2002) (ddesnsity wave)
- Houghton, Lee Marston, PRB 65, 220503 (2002) (t-j model)

Additional degrees of freedom leads to L/L_0 exceeding 1.



FIG. 1. The Lorenz ratio [Eqs. (12) and (13)] as a function of the doping for t=0.44 eV, J=0.13 eV. The ratio approaches unity in the dilute limit, $x \rightarrow 1$.





Measuring subkelvin thermal conductivity in a Bitter magnet (producing H >20T)



Normalised conductance, G/Gg, of a CBT sensor against bic voltage V. The theoretical curve of Eq. (1) is shown as a blac line.

Checking the WF law in optimally-doped Bi-2201



The amazing case of non-superconducting heavilyoverdoped LSCO





The collapse of linear term below 0.3K remains unexplained (reminiscent of PCCO!)

- Losing connection to electron thermal bath?
- A hidden, very small energy scale?



Summary

of what has been observed regarding the WF law in cuprates

- Validity of WF law in the high-field « normal » state (overdoped Tl-2201 and optimally-doped Bi-2201)
- Unexpected downward deviation from the Wiedemann-Franz law, (heavily-overdoped LSCO, and electron-doped PCCO)
- A large linear term κ_{00}/T in the superconducting state compared to the residual resitivity of the normal state (underdoped Bi-2201 and La-214)





Nernst coefficient in metals



 $\mathbf{J} = \mathbf{\sigma} \mathbf{E} + \alpha / \mathbf{I} \qquad \mathbf{I}$ $\mathbf{J}_{\mathbf{q}} = \alpha \mathbf{E} + \kappa \qquad \mathbf{T}$ In a Boltzmann picture: $\alpha \propto \partial \mathbf{\sigma} / \partial \epsilon \mid \epsilon = \epsilon_{\mathrm{F}}$ (Thus, a possible finite Nernst signal due to an energy-dependent scattering rate)

Nernst effect in optimally-doped cuprates (Ri, Heubner et al. 1994)



FIG. 3. Resistivity ρ (a) and normalized. Neural electric field $E_{\rm s}/\nabla_s T$ (b) versus temperature for an epitaxial, e-axis-oriented YBa₂Cu₂O₁₋₅ film at different magnetic fields applied parallel to the c axis of the film.

In underdoped cuprates Nernst coefficient remains finite for T>Tc



Vortex-like excitations in the Tc<T<T* window ? (First reported by Xu, Ong et al. Nature 406,486(2000))

A word of caution: the case of 2H-NbSe₂



Complicated Fermi surface may lead to a sizeable qp Nernst signal of both signs!



FIG. 2. Fermi-entries mapping of 2*H*-NbSe₂ by ARFES at $T=50\,$ K (left panel) and a theoretical Fermi surface obtained by bulk band-structure calculations (right panel). Posel (a) shows the ${\bf k}_1$ distribution of the globrosofission intensity measured at $h_2=24.5\,$ eV, and panel (b) shows the modulus of the two-data set of the two dimensional intensity gardient in comparison with the colculated Fermi-surface cut at $k_1=0$ [cf. Fig. 4(a)]. The k_1 -dependent Fermi-

Fermi surface Rossnagel et al. PRB 64,235119 (2001)

Field-dependence of the Nernst coefficient is highly non-linear!



A finite Nernst signal persists at high fields even when resistivity displays localisation!



Field-induced superconductor-insulator transition and the Nernst effect











Transport entropy of vortices



For a thermal force : $F_{th}=S_{\phi}$ $T=\eta v$ & $E_{y}=vB$ $N=E_{y}$ / $T=S_{\phi} B/ \eta$ For a Lorenz force: $F_{L}=J \Phi_{0}=\eta v \& E_{x}=vB$ $\rho=E_{x}/J=\Phi_{0}B/ \eta$ Therefore: $S_{\phi}=N\Phi_{0}/ \rho$ is the excess of entropy associated with a single vortex (Assuming that ρ is only due to flux flow!)

In the extereme underdoped case, $S_{\phi}c$ is a fraction of k_{B} . Two orders of magnitude smaller than in conventional superconductors!

 $S_{\phi} = S_n - S_s = k_B (ln\Omega_n - \Omega_s) = k_B ln(\Omega_n / \Omega_s)$ (Ω_n / Ω_s) is not much larger than unity!

Recent scenarios proposed to for a Nernst signal in an extended region of the (H,T) phase diagram

•Tesanavic, Vafek & Franz, PRB 65, 180511 (2002), Chiral symmetry breaking
•Weng & Muthukumur, cond-mat/ 0202079, Spontaneous vortices in RVB picture
•Ussishkin, Sondhi & Huse, cond-mat/0204484, Superconducting Gauusian fluctuations

•Kontani, cond-mat/0204193, AF fluctuations (No vortices)

•Ikeda, cond-mat/0203221, Quantum fluctuations due near the field-induced Superconductor-insulator Quantum Critical Point



FIG. 1. Schematic phase diagram for disordered superconducting films. *Distinct* T = 0 superconductor-insulator transitions occur at both critical disorder Δ_c and critical magnetic field R_c .

Summary

• In underdoped cuprates, a Nernst signal persists in an extended domain of the (H,T) plane

•Superconductivity seems to survive at high fields even in presence of a non-metallic resistivity

• The magnitude of extracted transport entropy in the underdoped regime is particularly small compared with conventional superconductors