

SMR/1438 - 5

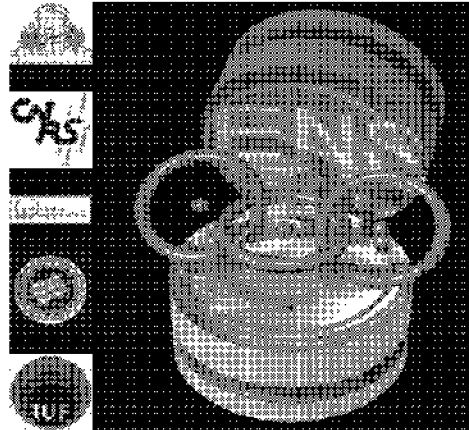
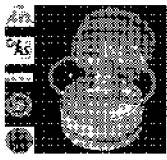
**JOINT ICTP-INFN SCHOOL/WORKSHOP ON
"ENTANGLEMENT AT THE NANOSCALE"**

(28 October – 8 November 2002)

"Controlled entanglement with Rydberg atoms and cavities"

presented by:

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Controlled entanglement with Rydberg atoms and cavities

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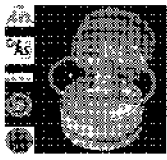
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P. Milman

L. Davidovich

Supported by CNRS, JST, EC program QUEST and QBITS



Experiments at the single quantum particle level



- Schrödinger 1952 :
« one never experiments with just one electron, one atom or one molecule. In thought experiments we sometimes assume that we do; this invariably entails **ridiculous consequences**... » (British Journal of the Philosophy of Sciences, vol 3, 1952)

- Of course we do:
 - EPR photon pairs (parametric down conversion)
 - trapped ions,
 - single electron on islands
 - Cavity QED: manipulating single atoms and a few photons

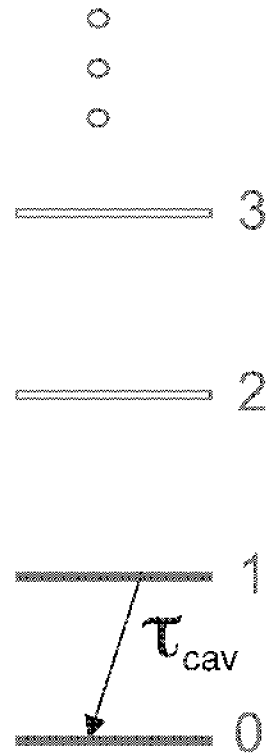
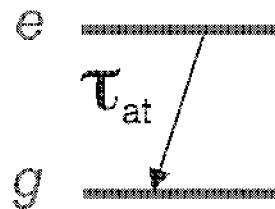
Most "*gedanken*" experiments become *real* experiments:

Quantum theory is definitely stranger than ridiculous!

Cavity QED: manipulating single atoms and photons

Two level atom
lifetime:

$$\tau_{\text{at}}$$



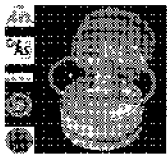
One mode of the e.m. field:
One harmonic oscillator
Photon lifetime:

$$\tau_{\text{cav}}$$

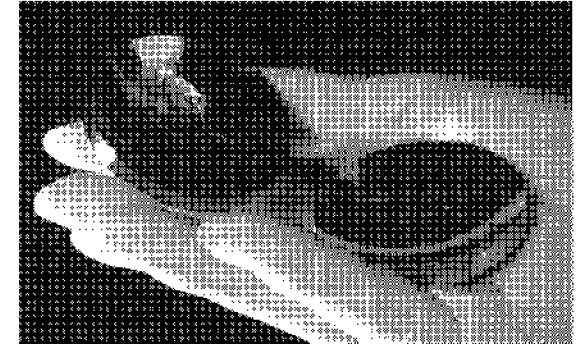
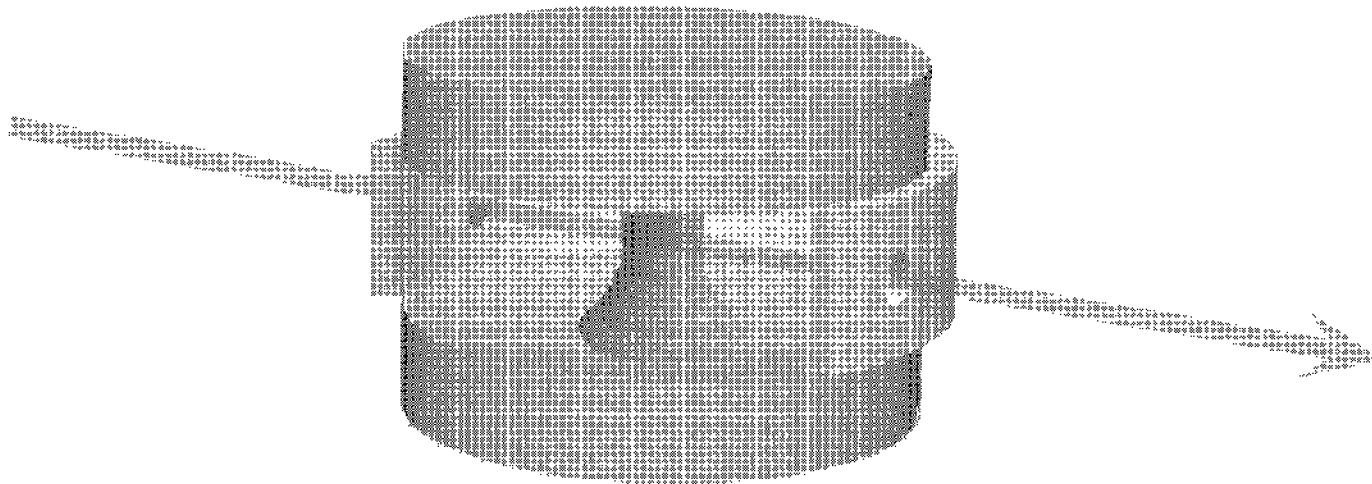
Coupling Ω

Requirement:

- $\Omega \cdot T_{\text{int}} > 2\pi$ strong atom-field coupling
- $T_{\text{int}} \ll \tau_{\text{at}}, \tau_{\text{cav}}$ weak coupling with the rest of the world



The microwave superconducting cavity



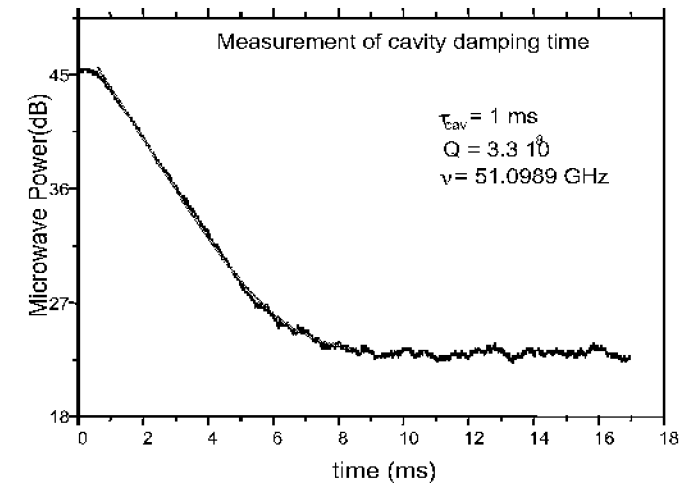
Cavity geometry:

- two spherical niobium mirrors
- a nearly closed aluminum ring
- microwave resonance:

$$\lambda=6\text{mm}, \nu_{\text{cav}}=51\text{GHz}$$

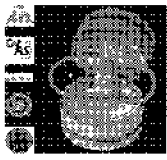
- superconducting mirrors:
- photon lifetime:

$$\tau_{\text{cav}}=1\text{ms}$$

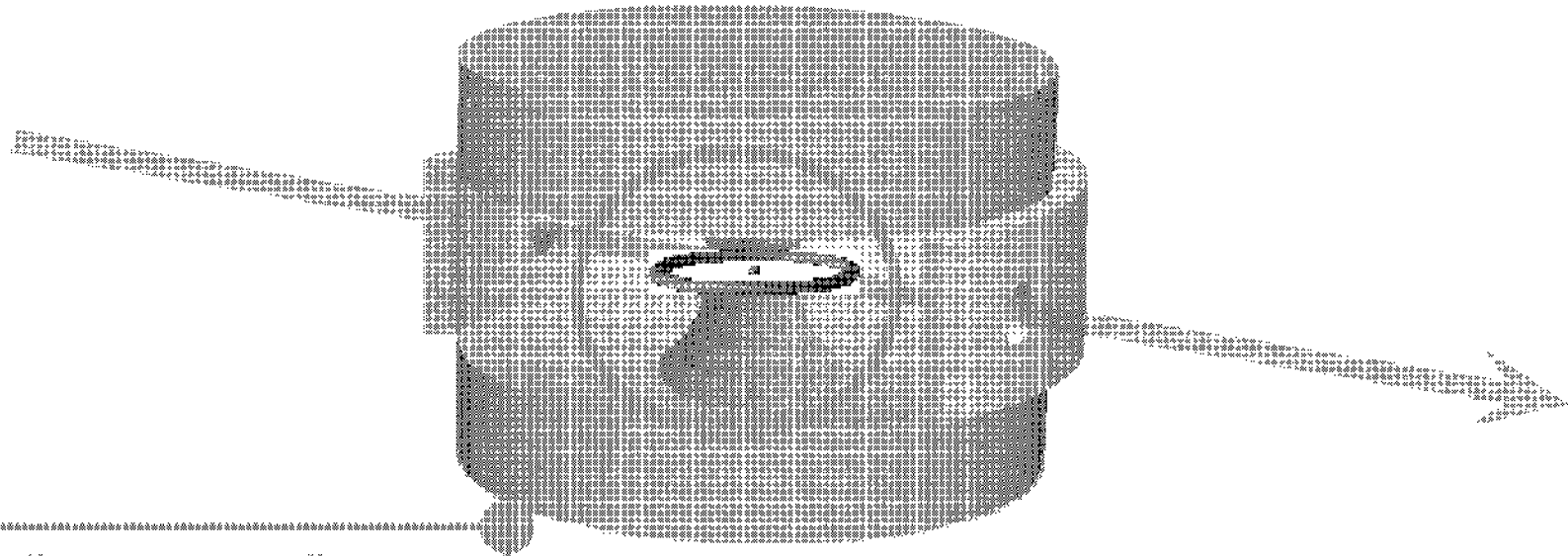


a photon "box":

- one mode
- a few photons

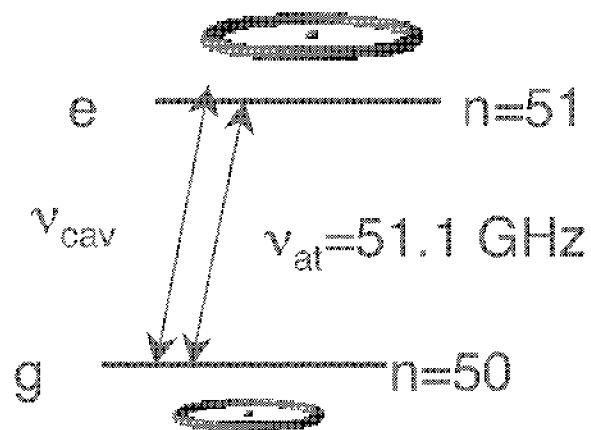


Circular Rydberg atoms

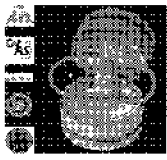


"Circular Rydberg atoms":

$$|l|=|m|=n-1$$



- long lifetime: 30 ms
- Large dipole: $a_0 \cdot n^2$
- ideal closed two level system



Experimental set-up

^{85}Rb

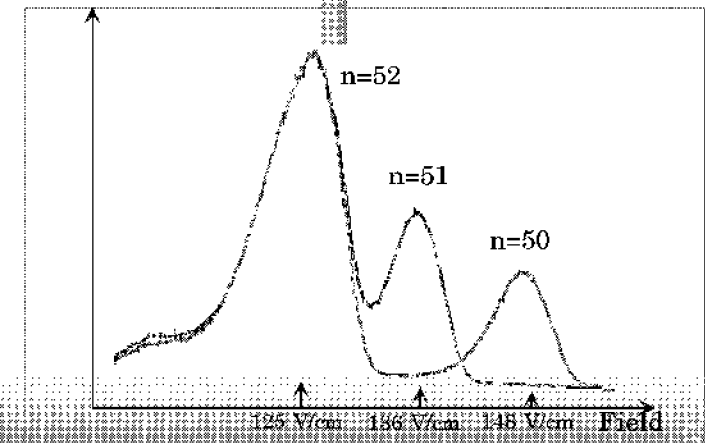
Laser velocity selection

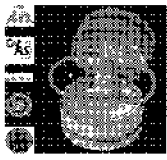
Cryogenic environment
 $T=0.6$ to 1.3 K
weak blackbody radiation

Circular atom preparation:

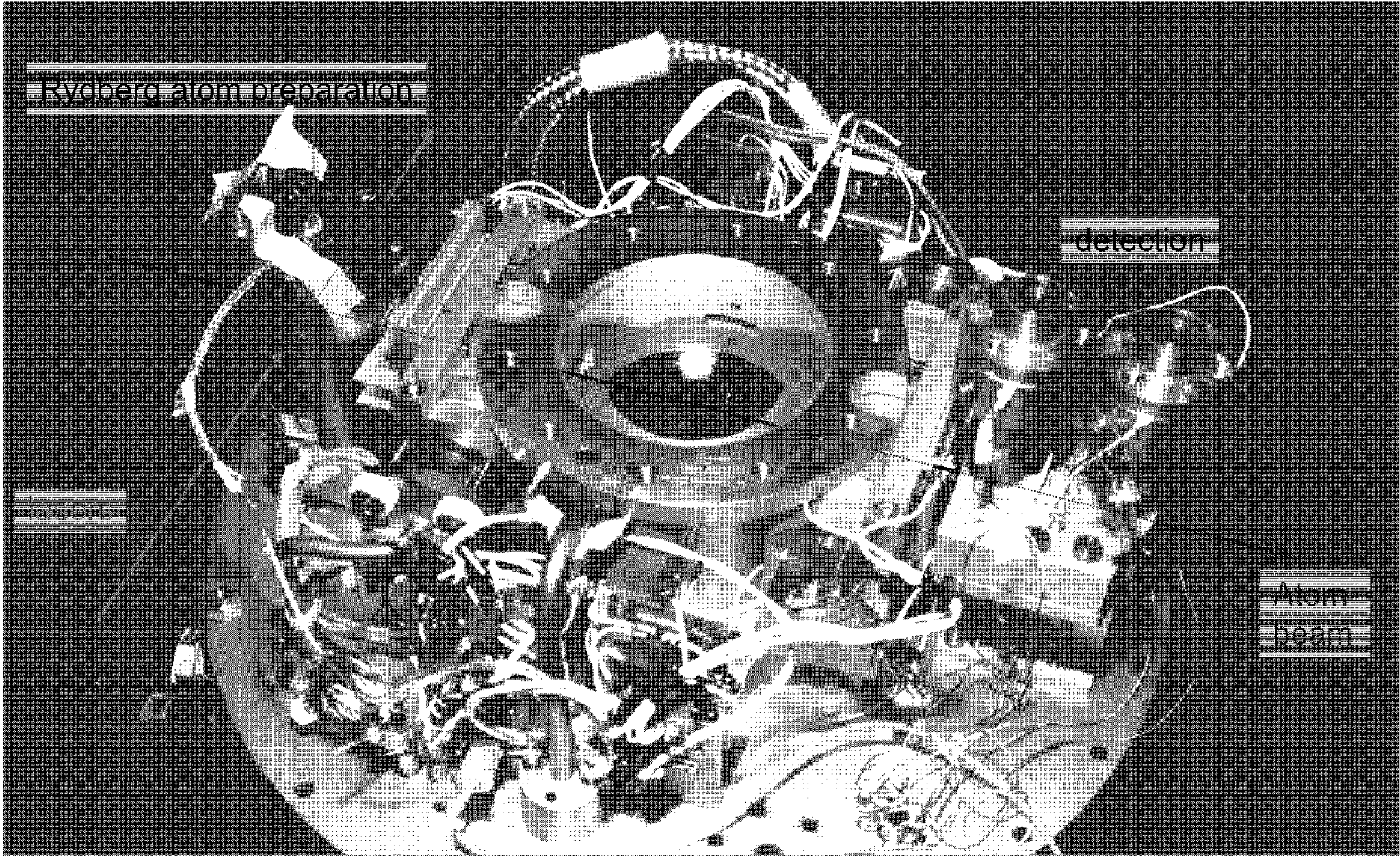
- 53 photons process
- pulsed preparation
- 0.1 to 10 atom/pulse

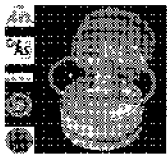
State selective
detector Atoms detected
one by one



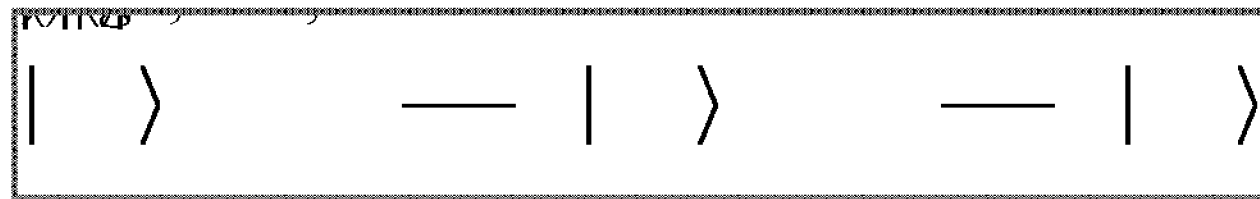
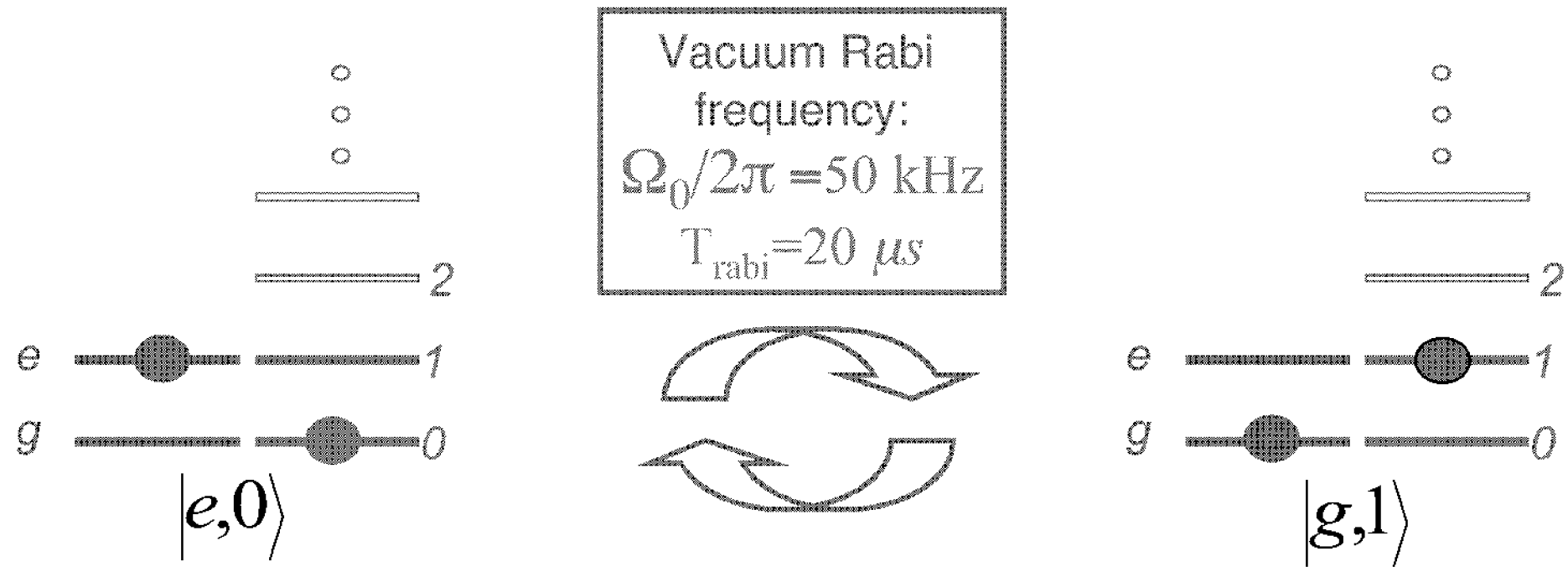


Our Cavity QED setup

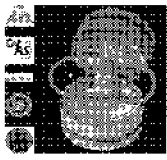




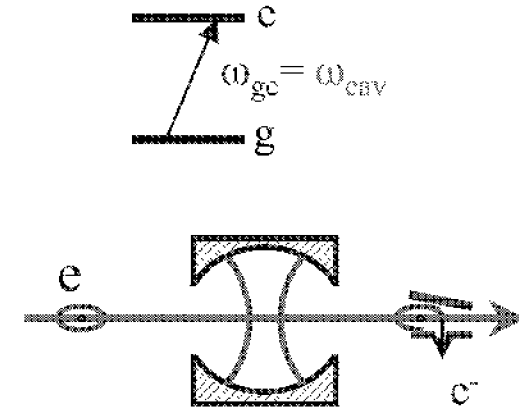
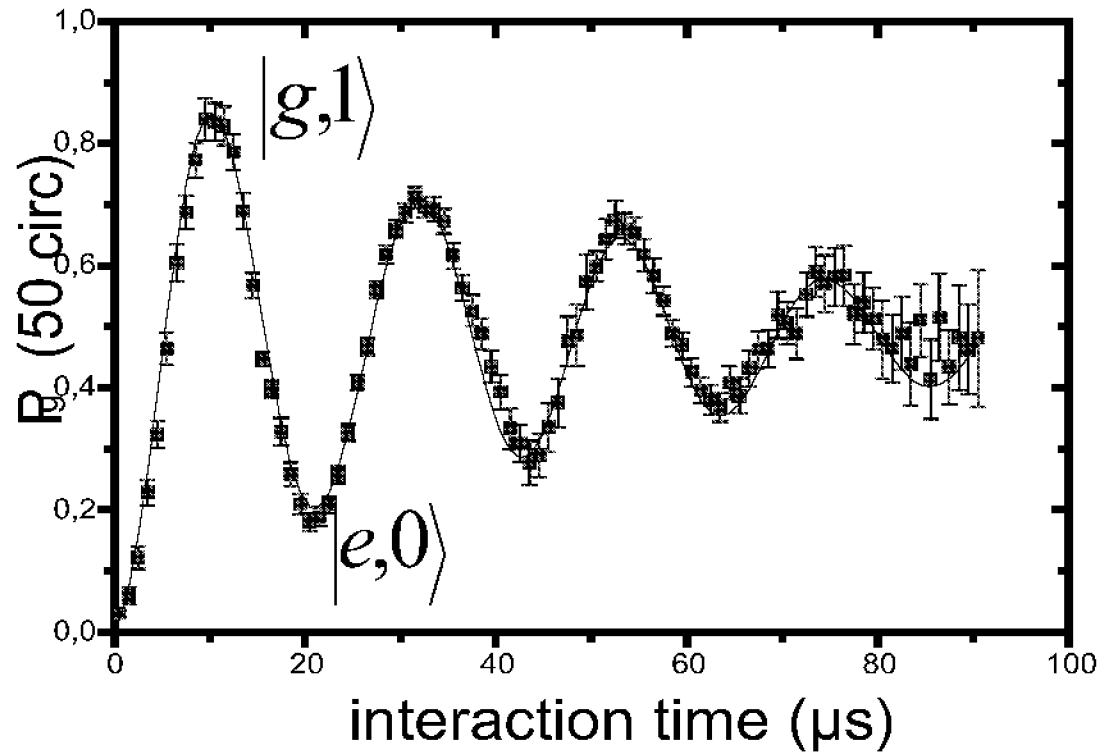
One atom-one photon: the strong coupling



- Coherent Rabi oscillation without any external field: reversible emission and absorption of one photon



Vacuum Rabi oscillation signal



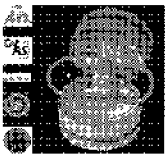
$$\Omega_0 = 47 \text{ kHz}$$

$$T_{\text{Rabi}} = 20 \mu\text{s}$$

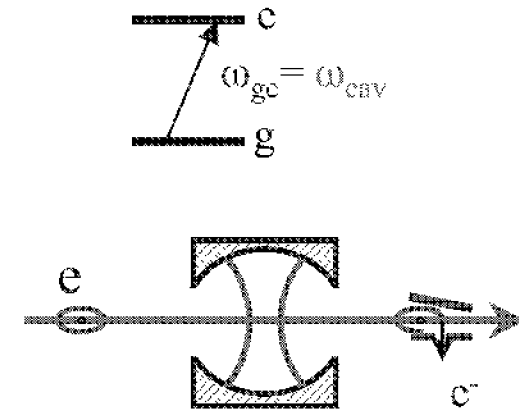
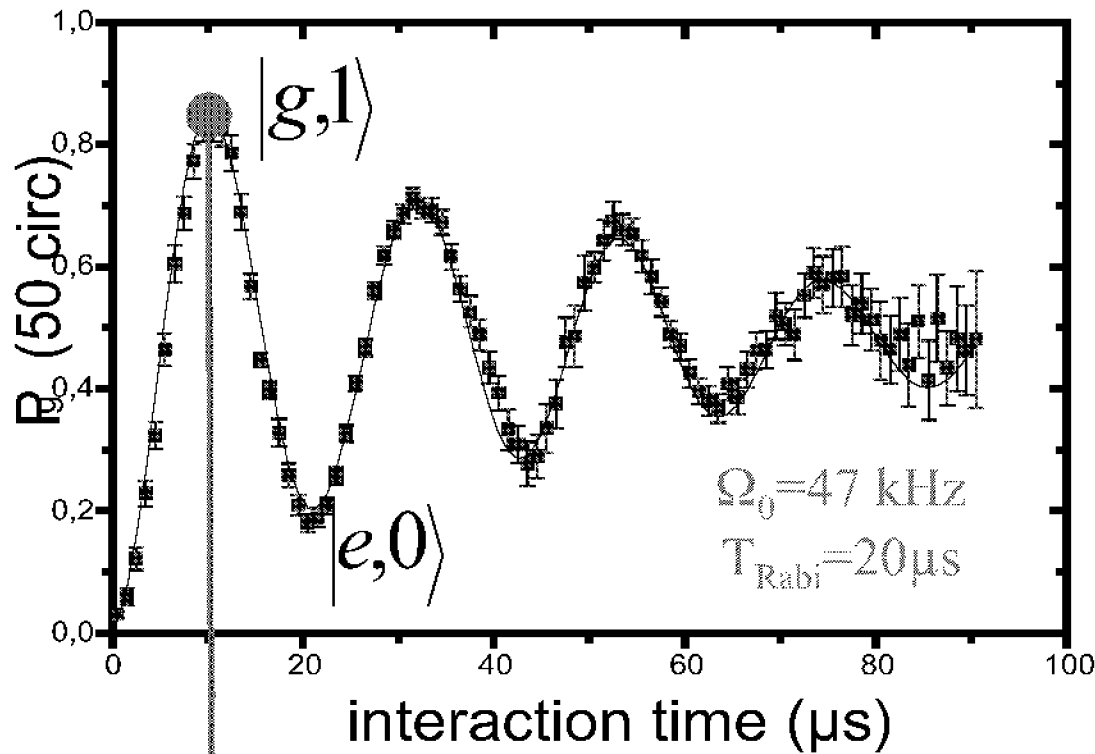
$$T_{\text{cav}} = 220 \mu\text{s}$$

$$T_{\text{at}} = 30 \text{ ms}$$

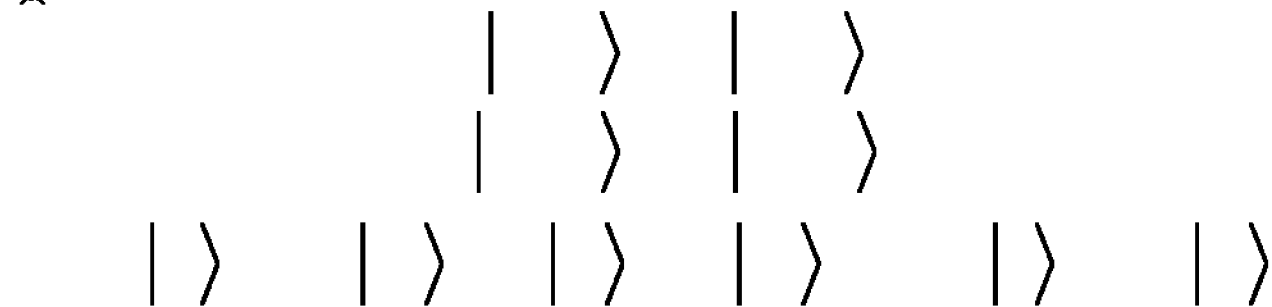
$\Rightarrow$ Strong coupling



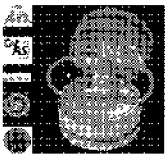
Quantum gates based on vacuum Rabi oscillation



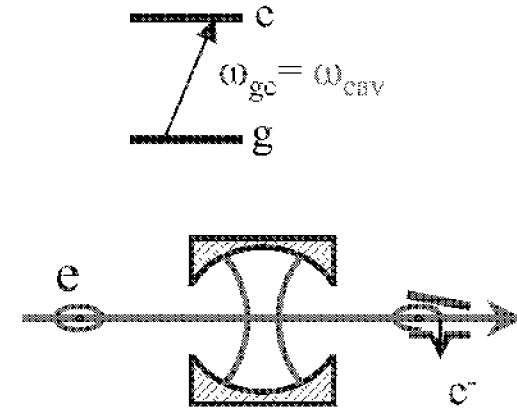
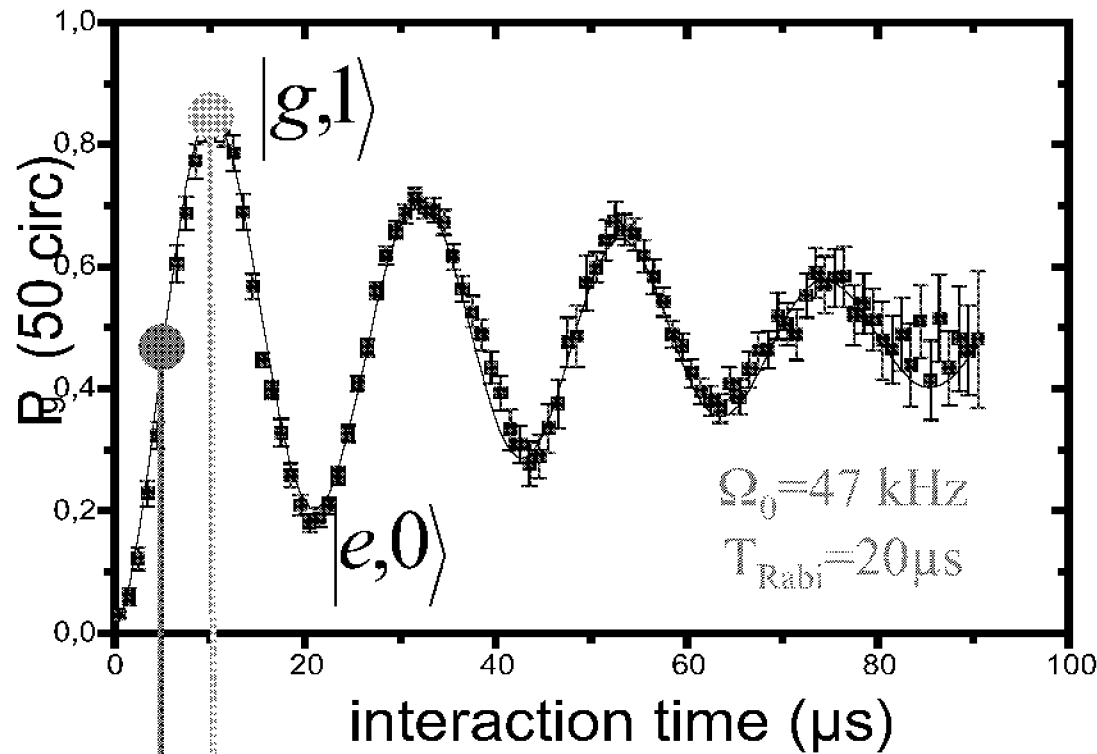
Emission/absorption of one photon



"Quantum memory", X. Maître et al. PRL 79,769 (1997)



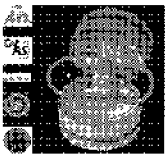
Quantum gates based on vacuum Rabi oscillation



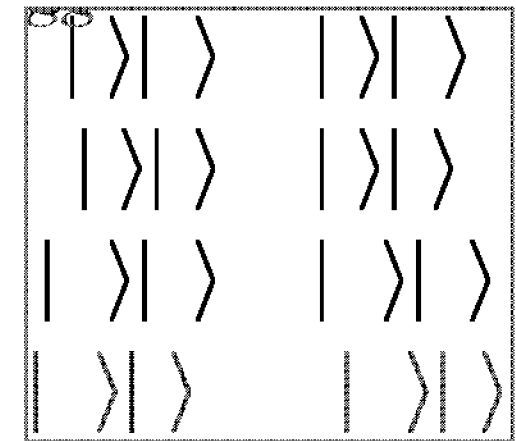
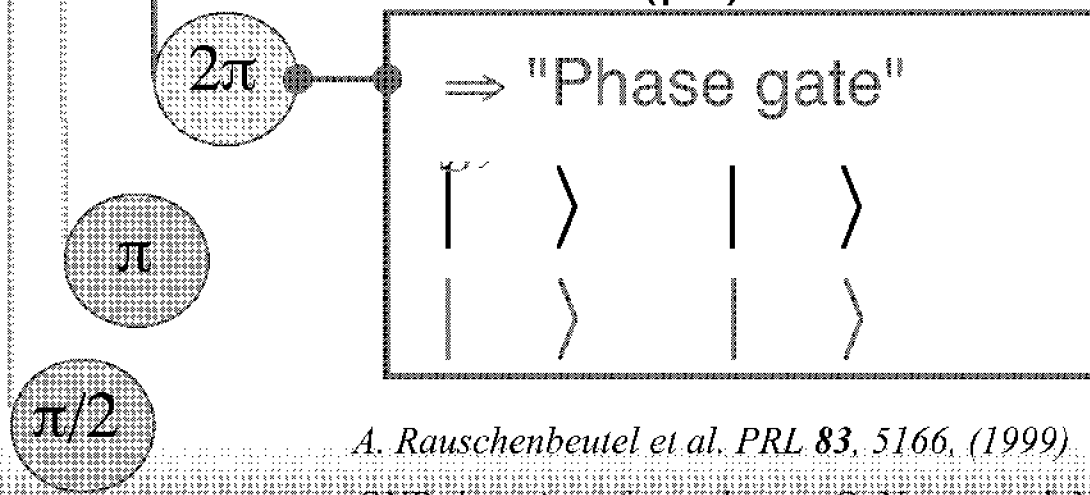
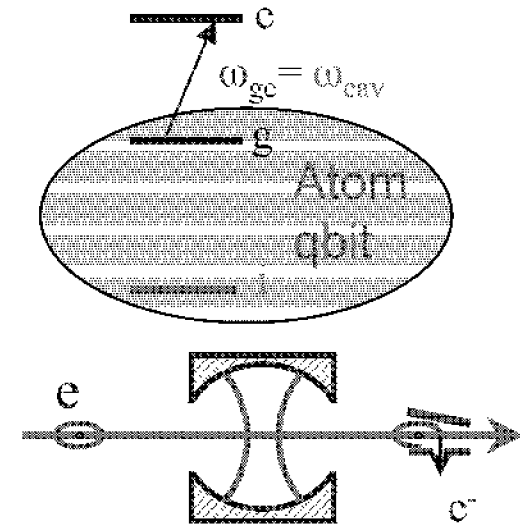
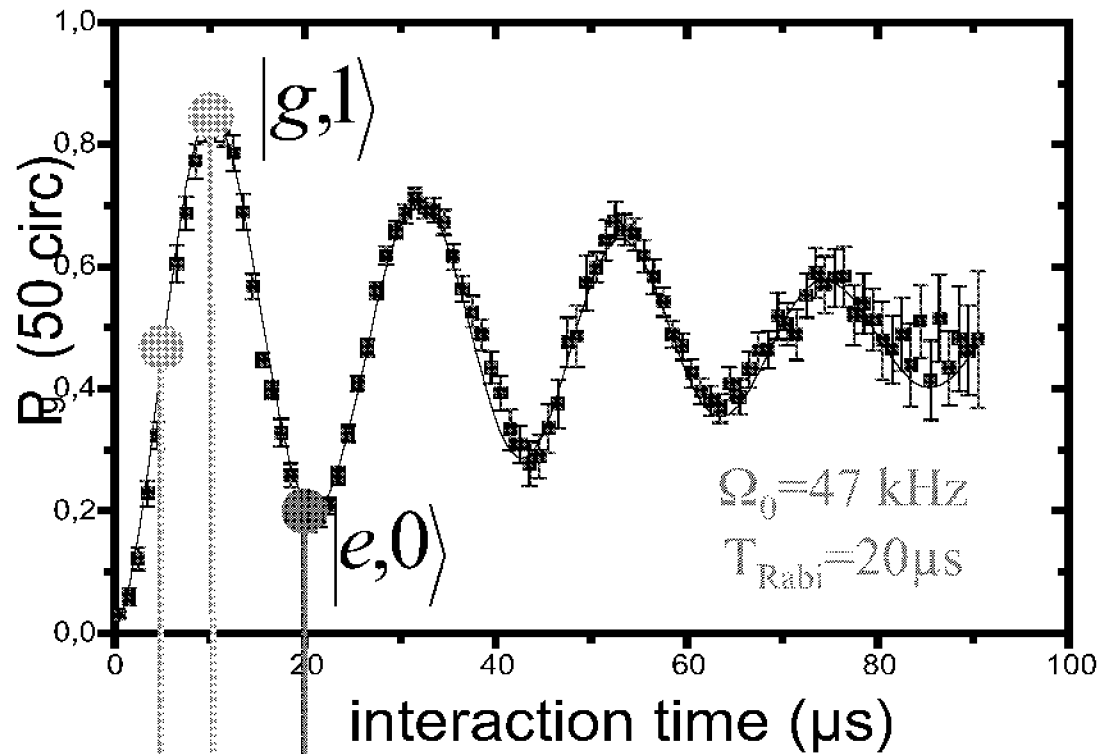
$\Rightarrow$ EPR pair preparation

$$| \chi \rangle = \frac{1}{\sqrt{2}} (| \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle)$$

E. Hagley et al. PRL 79,1 (1997)

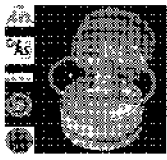


Quantum gates based on vacuum Rabi oscillation

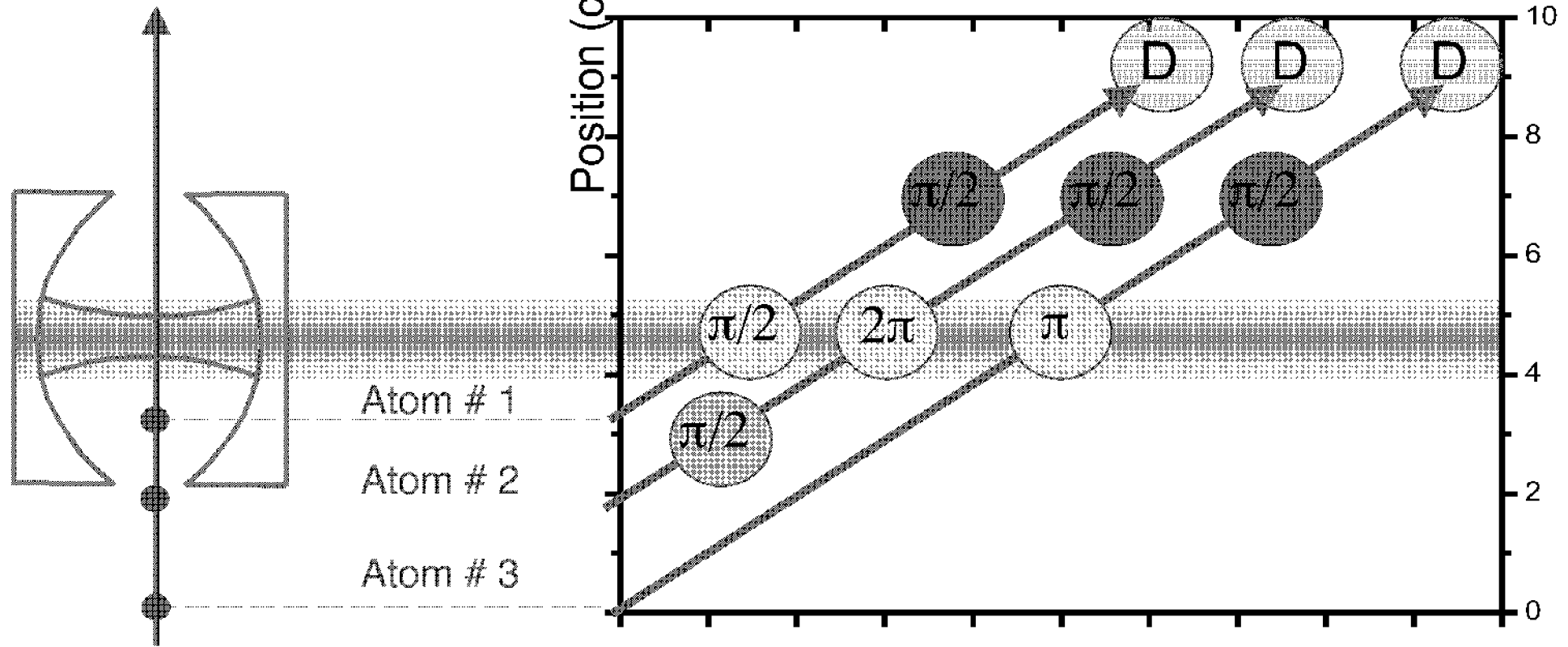


A. Rauschenbeutel et al. PRL 83, 5166, (1999)

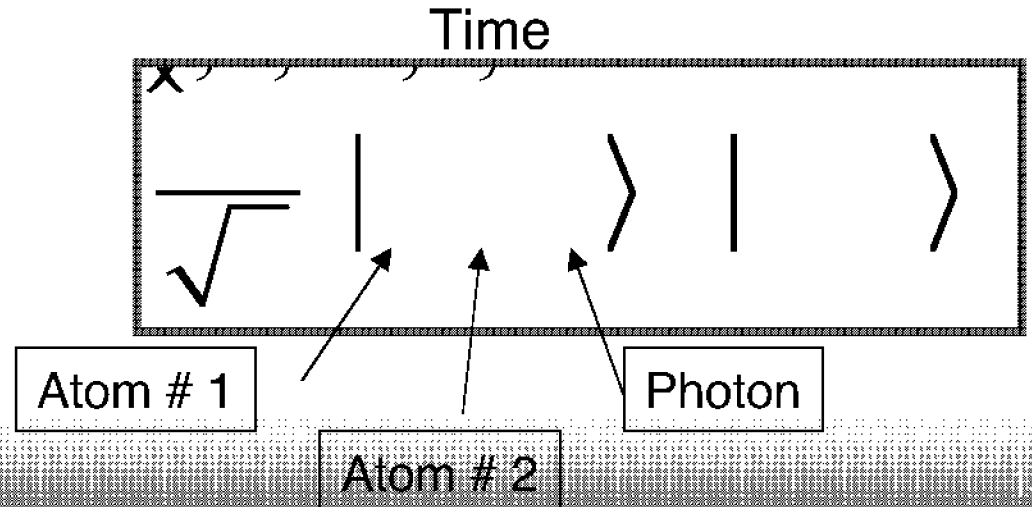
QND detection of one photon: G. Noges et al. Nature 400, 239 (1999)

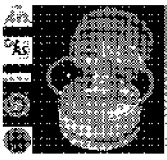


Step by step preparation of a GHZ triplet



- 0 • Resonant Rabi oscillation in C
- $\pi/2$ • Classical $\pi/2$ pulse
- D • Détection

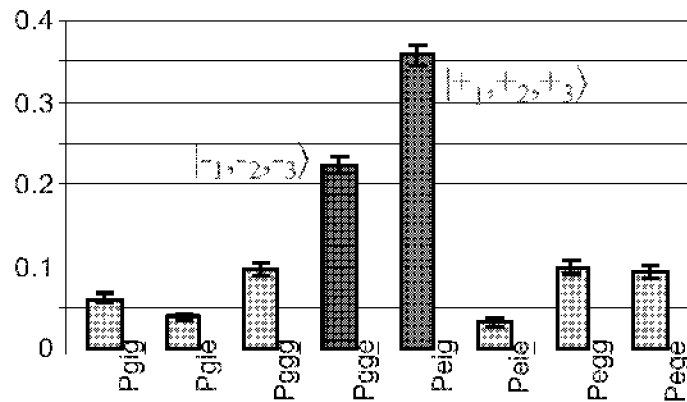




Fidelity of preparation of the GHZ state

$$\frac{1}{\sqrt{2}} (|+\rangle + |-\rangle)$$

Measurement of populations



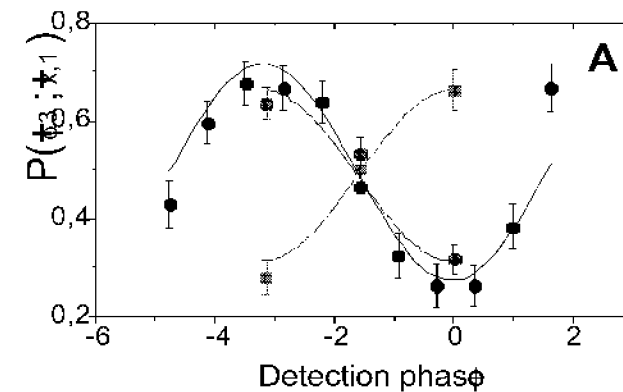
Rauschenbeutel et al., Science **288**, 2024 (2000)

Measurement of: " $\sigma_{z1}, \sigma_{z2}, \sigma_{z3}$ "

$$\Rightarrow P_{\text{long}} = P_{+++} + P_{---} = 0.58 (0.02)$$

$$\langle \sigma_{x1} \sigma_{x2} \sigma_{x3} \rangle$$

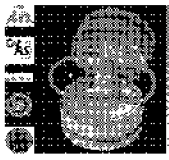
Measurement of coherences



Measurement of: $\sigma_{x1}, \sigma_{x2}, \sigma_{x3}$

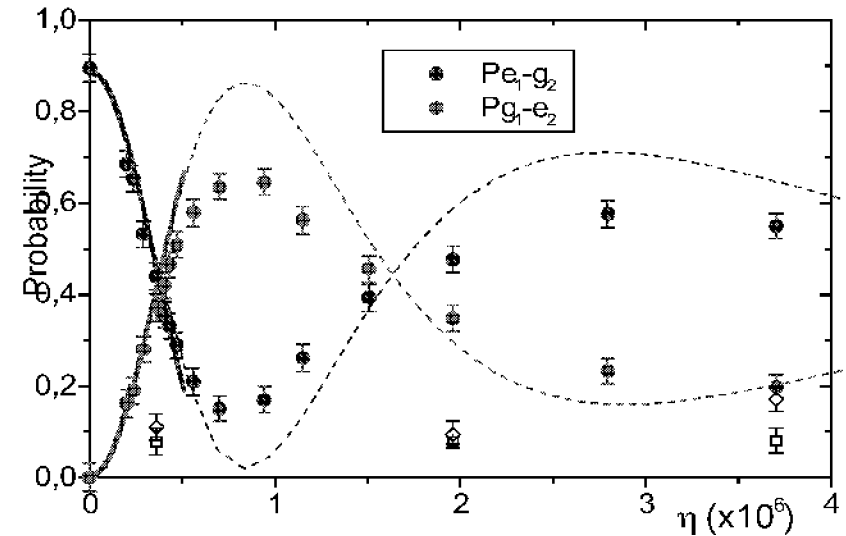
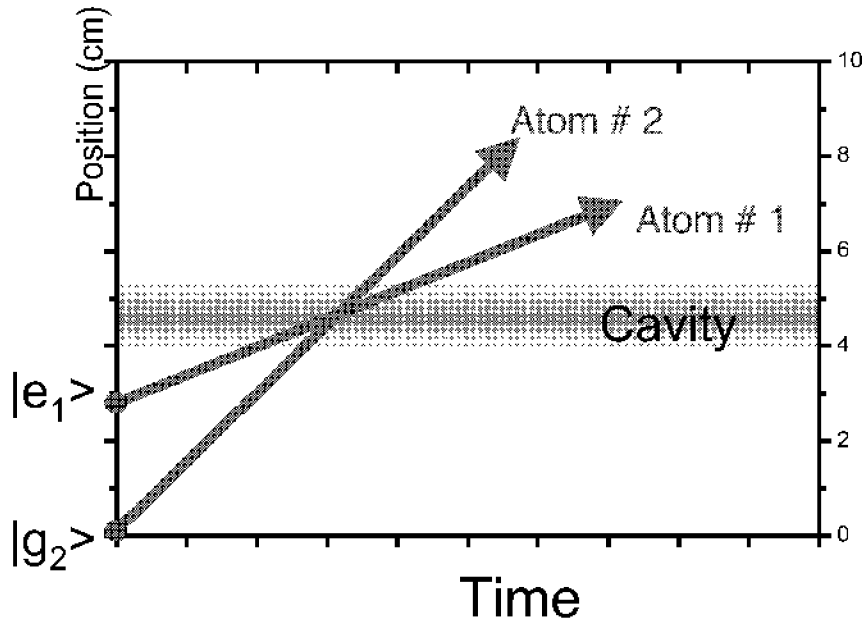
$$\Rightarrow A = \langle \sigma_{x1} \sigma_{x2} \sigma_{x3} \rangle = -0.28 (0.03)$$

$\Rightarrow$ The prepared state is not separable

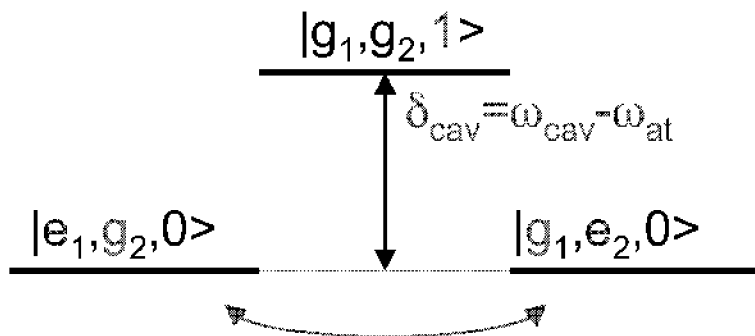


Direct atom-atom entanglement: non-resonant interaction with the cavity

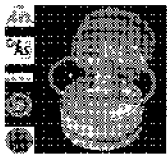
- Two atoms with different velocity cross the cavity at the same time:



Rabi oscillation
 $\Rightarrow$ Entanglement
 Unensitive
 to cavity damping



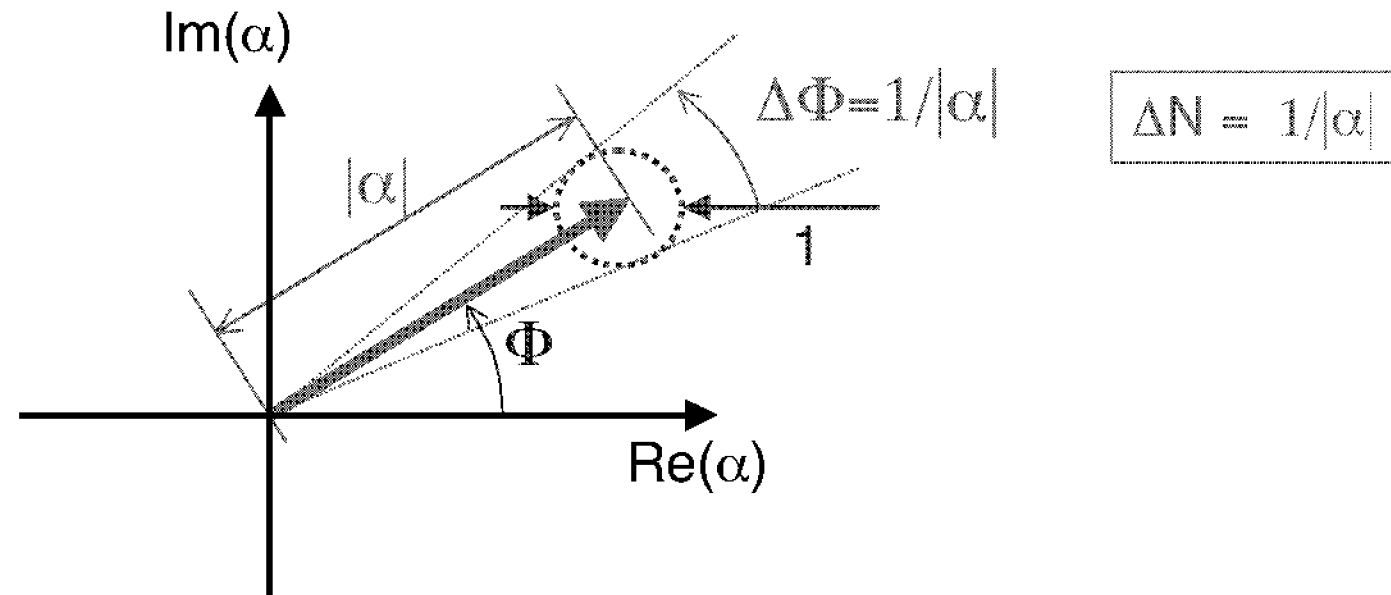
Effective coupling: —



Coherent state of the field

$$| \alpha \rangle = \frac{1}{\sqrt{n!}} e^{-|\alpha|^2/2} \frac{\alpha^n}{n!} | n \rangle$$

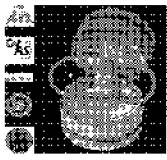
- Representation in the complex plane (Fresnel vector):



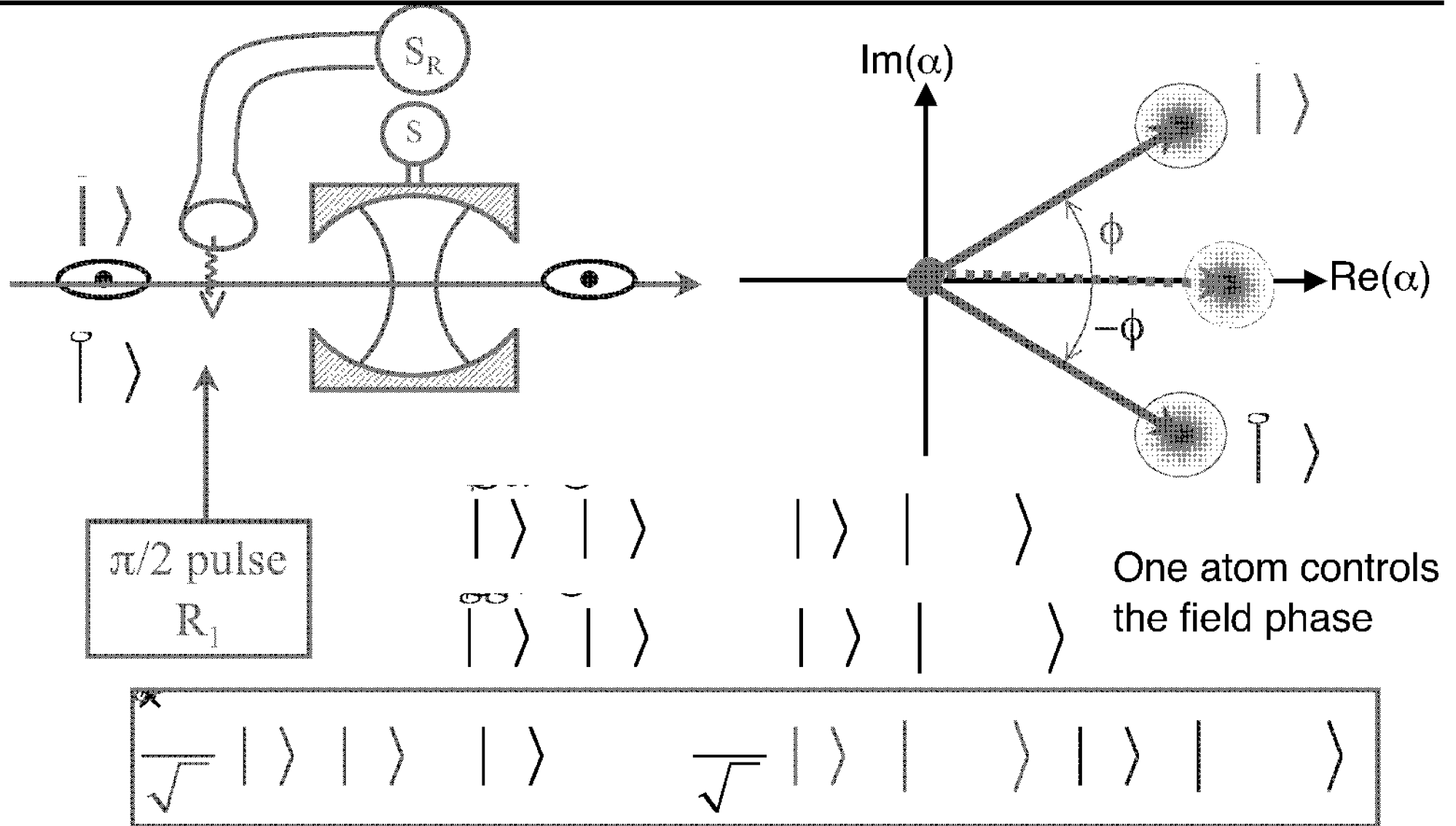
- Heisenberg uncertainty relation:

$$\Delta N \cdot \Delta \Phi > 1$$

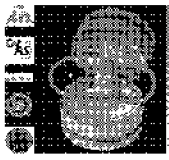
- 0 to tenth of photons easily injected in C using a classical source



Entanglement by non-resonant interaction manipulation of the classical phase

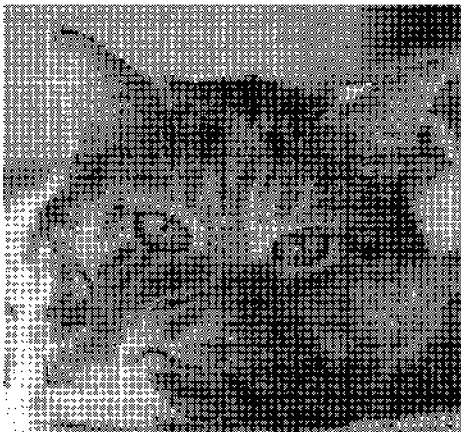


Single atom phase shift: 45° !



"Schrödinger cat" and decoherence

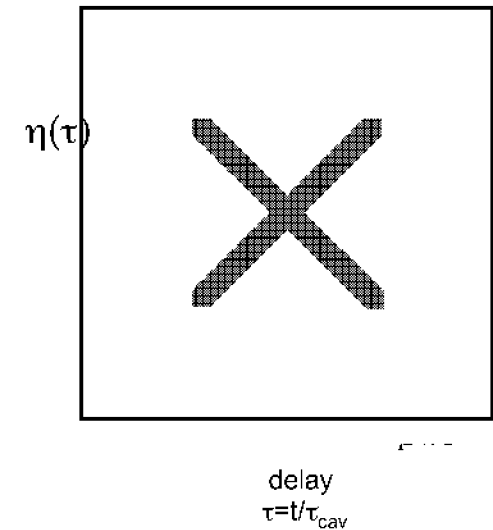
$$|\Psi_{cat}\rangle = 1/\sqrt{2} (|e, \nearrow\rangle + |g, \searrow\rangle)$$



Decoherence time

—

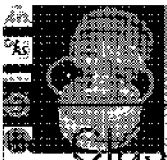
—



M. Brune et al.,
Phys. Rev. Lett. 77, 4887 (1996)

⇒ Observe decoherence for larger fields

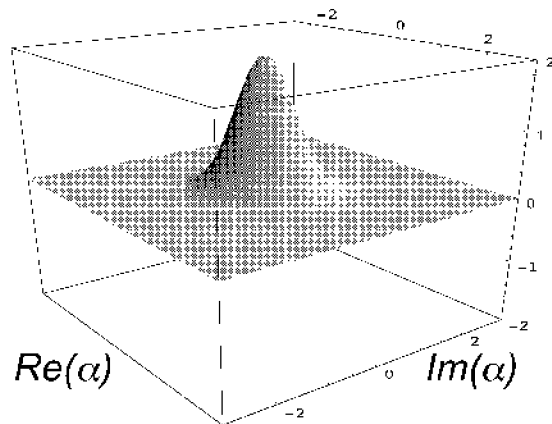
⇒ measure the complete field state: Wigner distribution



Examples of Wigner functions

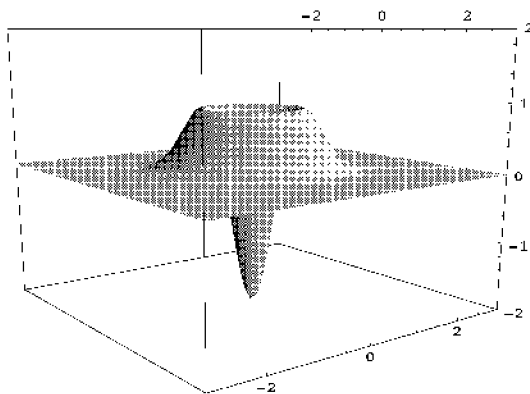
Classical states

Vacuum $|0\rangle$

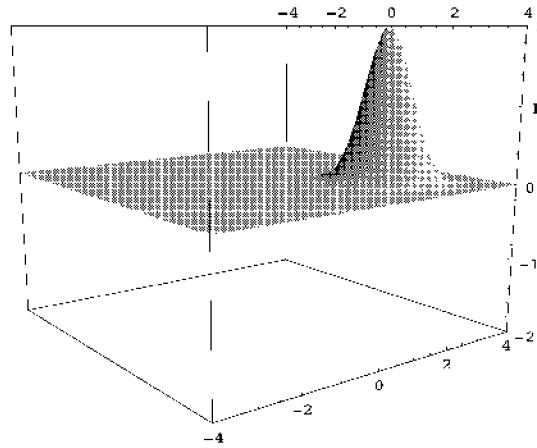


• Quantum states

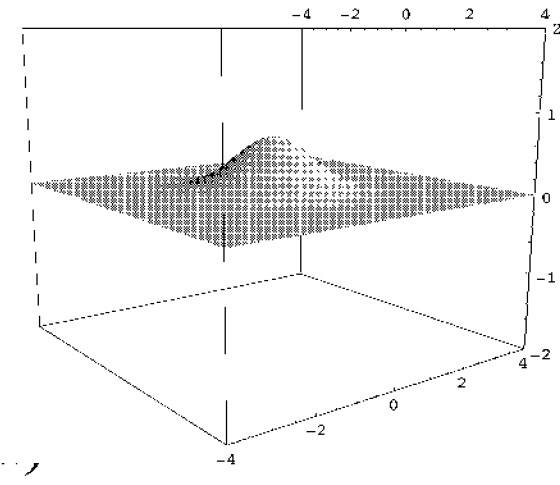
Fock state $|1\rangle$



Coherent state $|\alpha\rangle$
($\alpha=1.5+1.5i$)

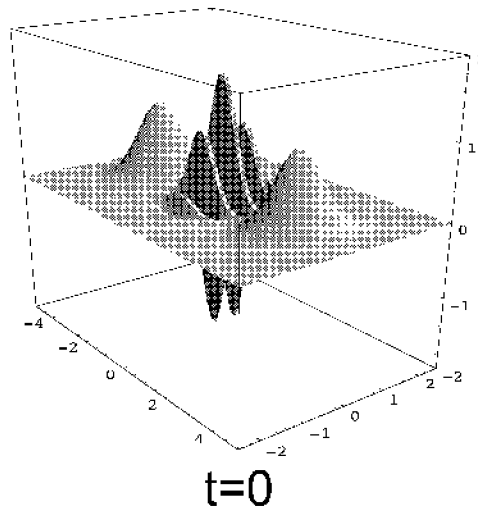


Thermal field $n_{th}=1$

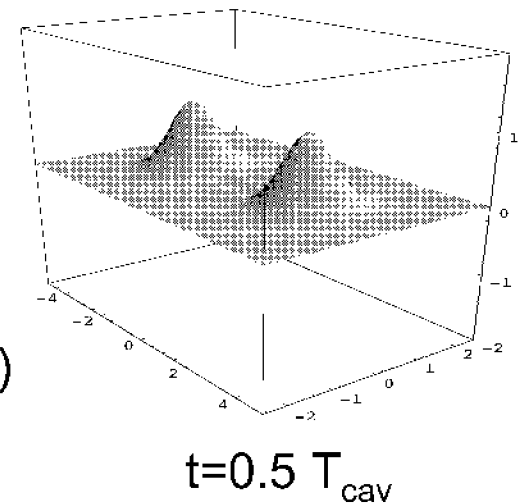


Schrödinger cat's state

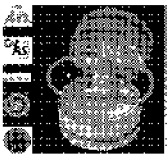
$$| \rangle - \frac{1}{\sqrt{2}} (| \rangle + | \rangle)$$



($\alpha=3$)



Negativity of W is a good criterion of quantumness



Principle: measure $W(\alpha)$ by measuring field "parity"

- Based on : *Cahill and Glauber, PR 177, 1857 and 1882 (1969)*



$\hat{D}(\alpha)$

Displaced field density operator

$$| \rangle \begin{cases} | \rangle & \text{if } n \text{ even} \\ | \rangle & \text{if } n \text{ odd} \end{cases}$$

« parity » operator

⇒ Method of measurement of W :

A) Apply $D(-\alpha)$

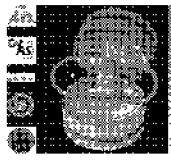


Classical source coupled to the cavity mode : easy

B) Measure parity operator

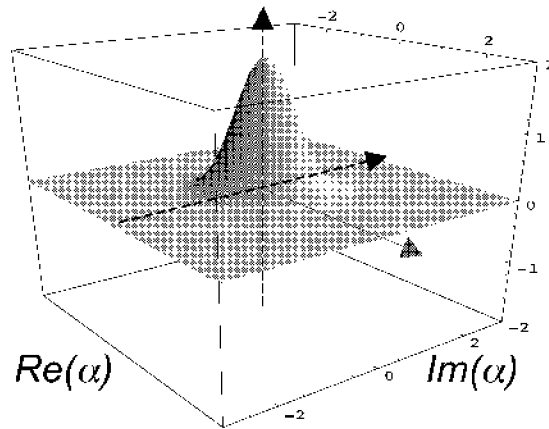
*B. Englert et al., Opt. Comm. **100**, 526 (1993),*

*Lutterbach and Davidovich, PRL **78** (1997) 2547*

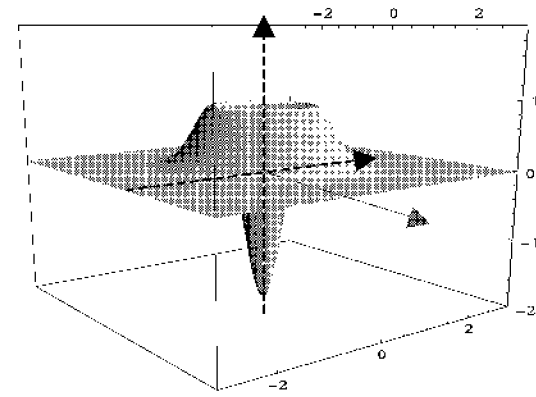


Measuring $W(a)$ for vacuum or one photon state:

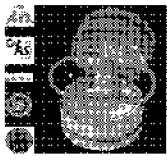
Vacuum $|0\rangle$



Fock state $|1\rangle$

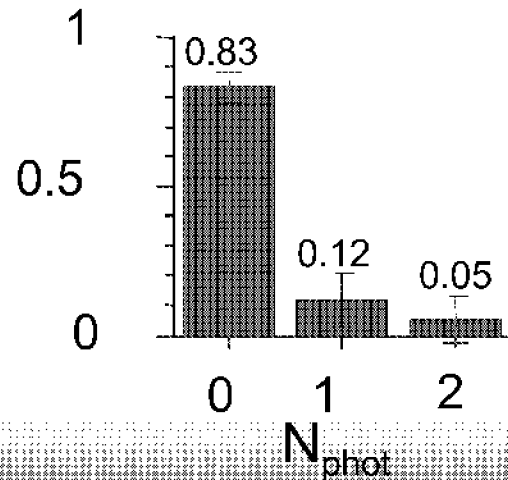
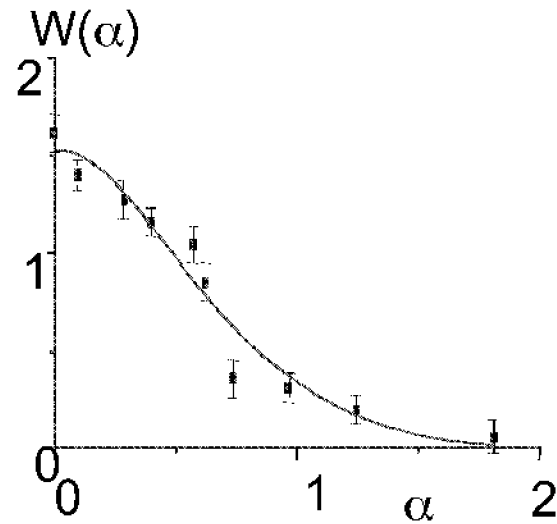


These distributions have a cylindrical symmetry
We only measure a radial cut
by varying $|a|$ at a fixed phase

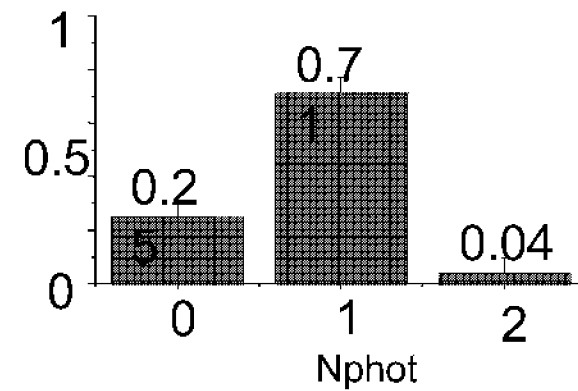
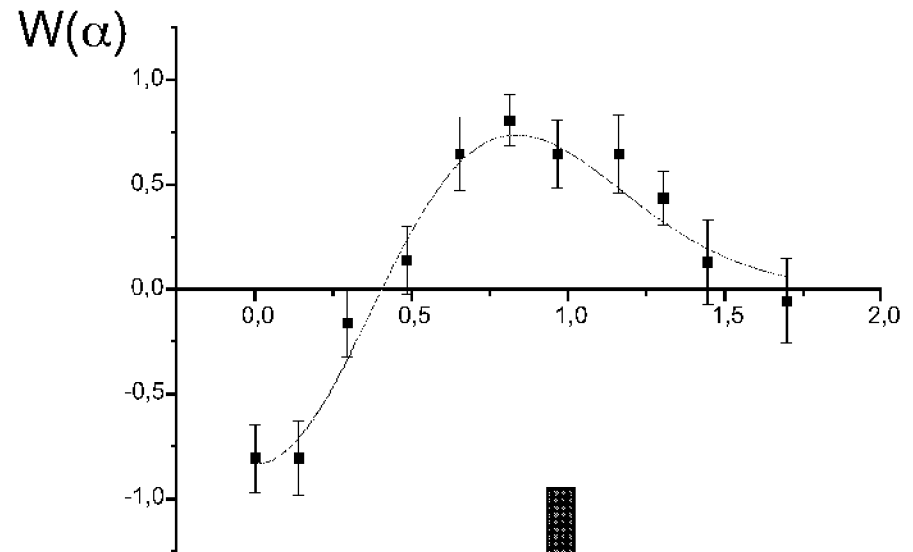


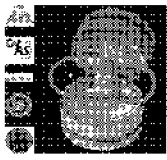
Measured Wigner functions

"Vacuum"



One photon Fock state





Conclusion and perspectives

Rydberg atoms and cavities:

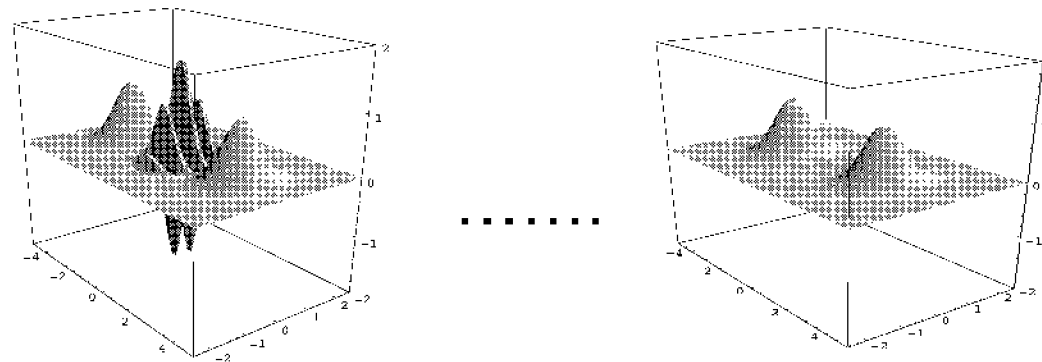
efficient tools for entanglement "engineering":

- manipulation of entangled qubits:
 - longer cavity damping time
 - better detection efficiency
 - use of non-resonant gates

⇒ Manipulation of up to 6 atoms and two cavity modes

⇒ demonstration of elementary quantum algorithm: error correction, Grover

- Preparing and characterizing mesoscopic superposition states:



⇒ Observing a movie of decoherence?