

SMR 1435/4

WORKSHOP ON THEORETICAL PLASMA ASTROPHYSICS

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Self-Organization & Singular Perturbations in Plasma Flows

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These are preliminary lecture notes, intended only for distribution to participants.

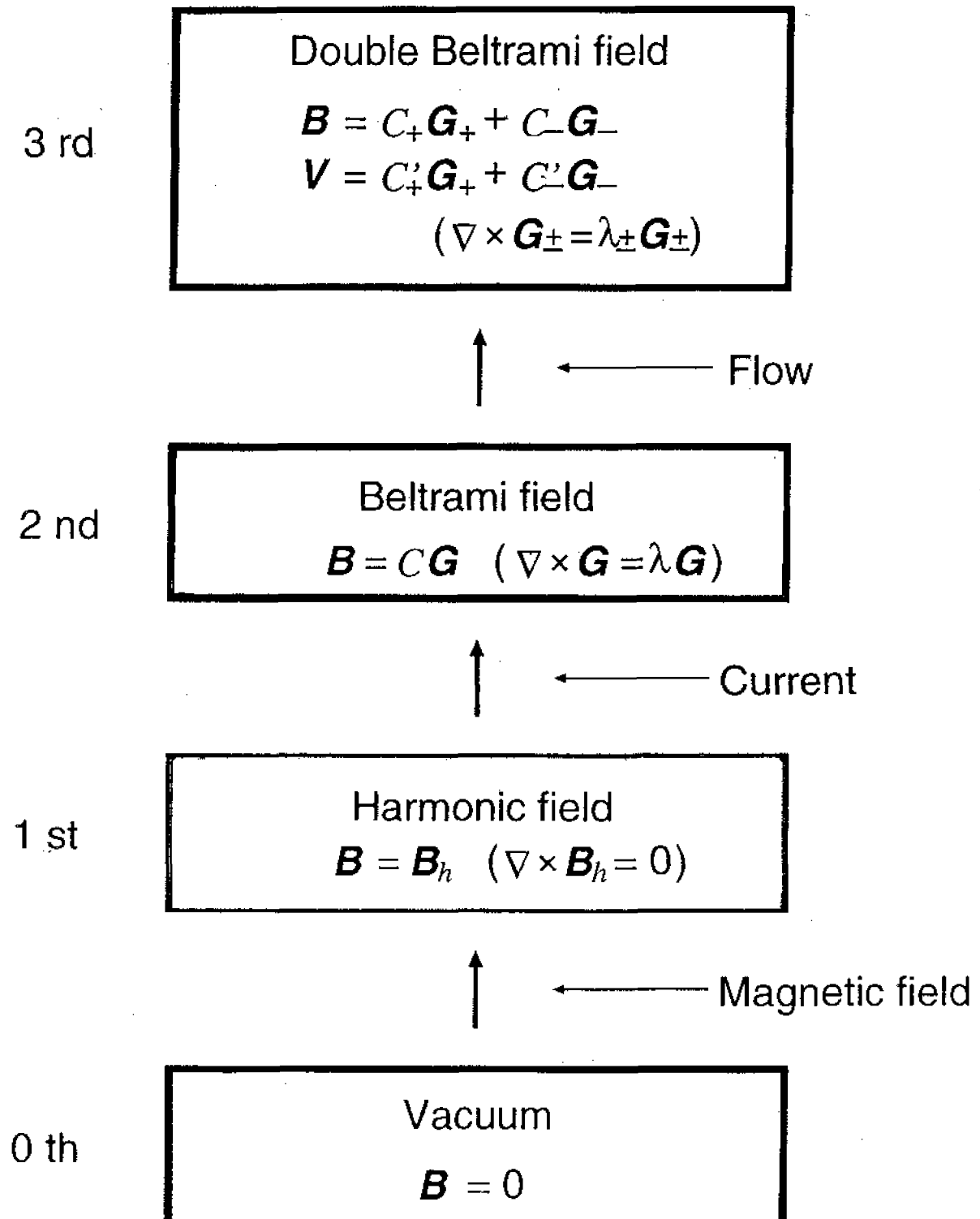
Self-organization & Singular Perturbations in Plasma Flows

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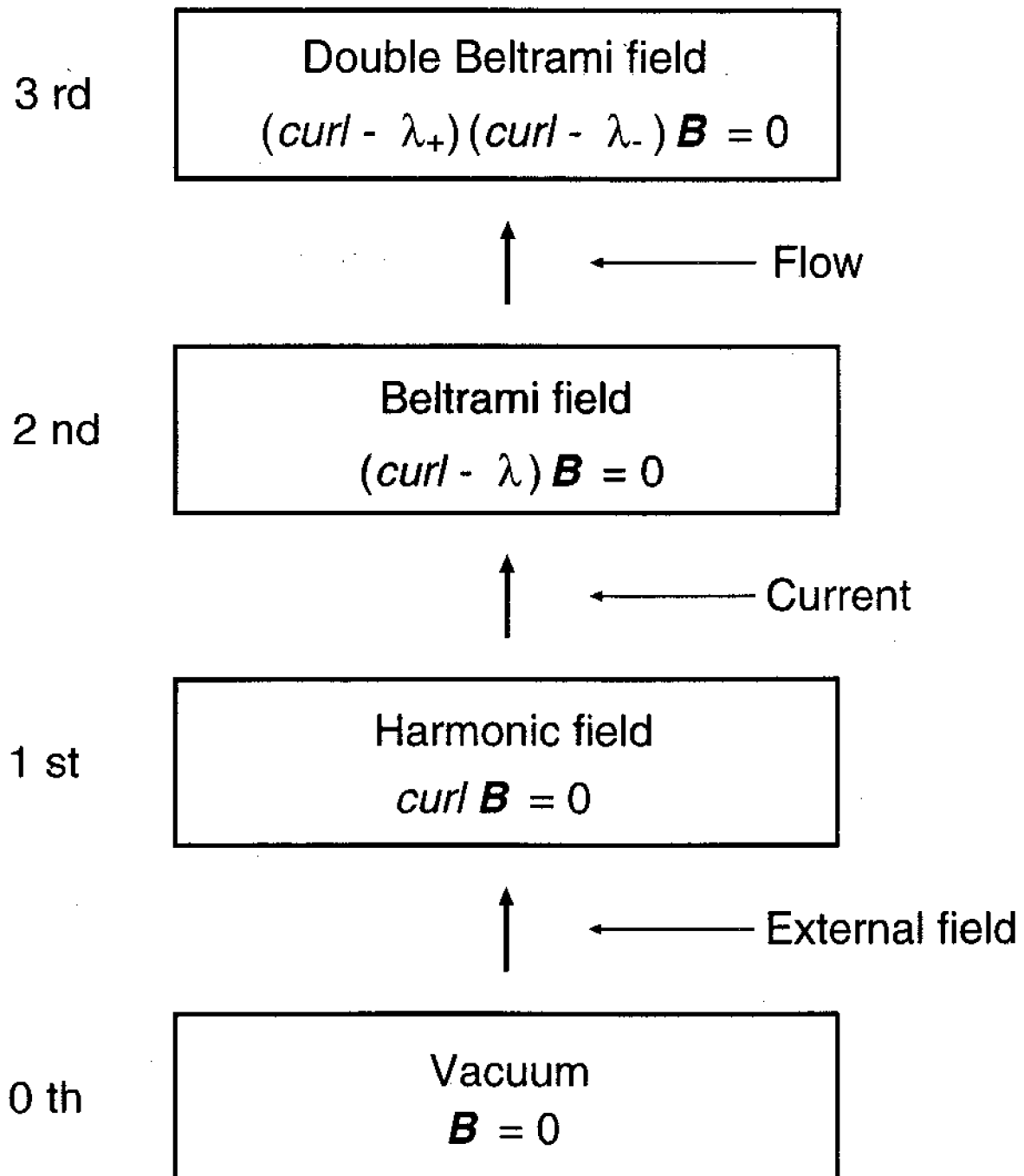
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Hierarchy of relaxed states



Hierarchy of relaxed states



Topology of Function Spaces



fine structure

Topology



coarse grain

strong topology

coercive



weak topology

fragile

dissipation



robust

Mathematical Exercise

$$\begin{cases} g(u) = \int_{\Omega} |\nabla u|^2 dx \\ H(u) = \int_{\Omega} |u|^2 dx \end{cases} \quad \left(\begin{array}{l} u|_{\partial\Omega} = 0 \\ \Omega \subset \mathbb{R}^N; \text{ bounded} \end{array} \right)$$

(i) find a minimizer of $g(u)$ under $H(u)=1$

$$\delta [g(u) - \lambda H(u)] = 0$$

$$\rightarrow \begin{cases} -\Delta u = \lambda u \\ u|_{\partial\Omega} = 0 \end{cases}$$

$$\rightarrow u = a\varphi_j \quad (-\Delta\varphi_j = \lambda_j\varphi_j)$$

$$H(u) = |a|^2 = 1 \rightarrow a = 1$$

$$g(u) = \lambda_j \rightarrow \boxed{u = \varphi_1}$$

(ii) find a minimizer of $H(u)$ under $g(u)=1$

$$\delta [H(u) - \mu g(u)] = 0$$

$$\rightarrow \begin{cases} -\Delta u = \mu^{-1} u \\ u|_{\partial\Omega} = 0 \end{cases}$$

$$\rightarrow u = a\varphi_j$$

$$g(u) = |a|^2 \lambda_j = 1 \rightarrow a = \lambda_j^{-1/2}$$

$$H(u) = 1/\lambda_j \rightarrow \boxed{u = 0}$$

I. Self-organization

Essential characteristics

- irreversible process, but not dominated by dissipation
- generation of stable equilibrium with global simplicity (coarse grain)
- not always

conventional models

(1) Dirichlet principle (diffusion process)

$$\partial_* u = -\partial G \quad (G: \text{convex functional})$$

$$\text{ex. } \partial_* u = \nabla \cdot (D(x) \nabla u) + f(x)$$

minimizes

$$G = \frac{1}{2} \int D(x) |\nabla u|^2 dx - \int f(x) u dx$$

- Exact, if we have G .
- predicts "trivial" equilibrium.

(2) Selective dissipation

$H_0, H_1, \dots, H_N$ constants of motion (ideal)

minimize H_0 with keeping $H_1, \dots, H_N$ constant.

$$\text{ex. minimize } \int B^2 dx - \mu \int A \cdot B dx$$

→ Taylor relaxed state.

- Hypothetical model
- H_0 must be coercive
- always?
- may derive unstable solution

(3) Statistical distribution

$H_0, H_1, \dots, H_n$ constraints of motion (ideal)

maximize S' on the ensemble
(entropy)

- Ensemble must be a closed manifold
- second quantization

(4) Minimum entropy production

Minimize "entropy production" with
keeping $H_1, \dots, H_n$ const.

ex. minimize $\int \eta (\nabla \times \mathbf{B})^2 dx - \mu_1 \int B^2 dx - \mu_2 \int \mathbf{A} \cdot \mathbf{B} dx$

- may derive non-equilibrium (unstable) solutions.

A new framework

Target functional (F)

- measure of fluctuations and dissipation (turbulence)
- coercive functional

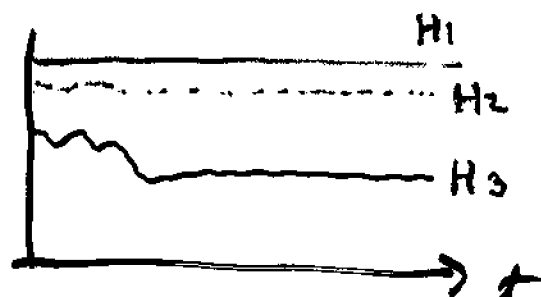
Constraints ($H_1, \dots, H_N$)

- constants of motion of ideal dynamics
- adjusted by a small dissipation
- continuous functionals in the topology of F .

Self-organization

$$\delta [F - \mu_1 H_1 - \mu_2 H_2 \dots \mu_N H_N] = 0$$

Succeeds when $H_1 \dots H_N$ are adjusted to produce a stable equilibrium so that F is minimized.



Self-organization in two-fluid MHD

$$\partial_t \omega_j + \nabla \times (\omega_j \times V_j) = 0$$

$$\begin{cases} \omega_1 = B, & V_1 = V_e = V - \nabla \times B \quad (\text{electrons}) \\ \omega_2 = B + \nabla \times V, & V_2 = V \quad (\text{ions}) \end{cases}$$

Constants of motion

$$\begin{cases} H_0 = \|V\|^2 + \|B\|^2 = \int (|V|^2 + |B|^2) dx \\ H_1 = (\text{alt}^T \omega_1, \omega_1) = \int A \cdot B dx \\ H_2 = (\text{alt}^T \omega_2, \omega_2) = \int (A + V) \cdot (B + \nabla \times V) dx \end{cases}$$

(Note: H_0 is not coercive to control H_2 .
 H_2 is not convex)

measure of fluctuations and dissipation

$$F = \frac{\|\nabla \times \omega_2\|^2}{\|\omega_2\|^2} \quad (\text{ion anisotropy})$$

self-organization

$$\delta [F - \mu_0 H_0 - \mu_1 H_1 - \mu_2 H_2] = 0$$

Euler-Lagrange equation

$$(\text{curl} - \lambda_1)(\text{curl} - \lambda_2)(\text{curl} - \lambda_3) u = 0 \quad (T)$$

$$(u = B \text{ or } V)$$

→ "triple Beltrami" solution

$$u = c_1 G_1 + c_2 G_2 + c_3 G_3$$

$$(\text{curl } G_j = \lambda_j G_j)$$

→ u is a stable equilibrium if u is a "double Beltrami" with small enough λ_1 and λ_2 .

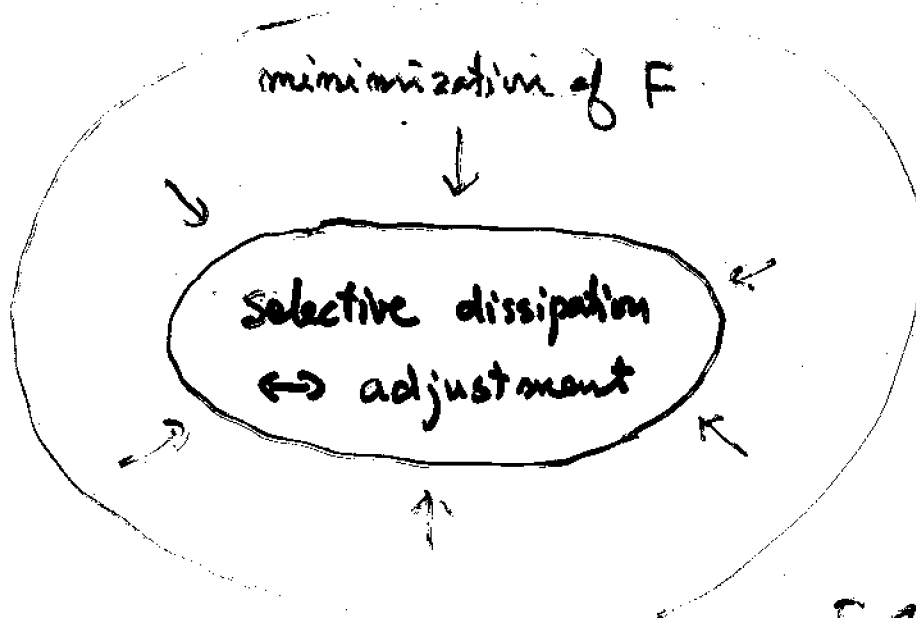
→ This occurs when H_0, H_1, H_2 satisfies a relation

$$H_0 - \mu_1' H_1 - \mu_2' H_2 = 0$$

$$(\Leftrightarrow \delta [H_0 - \mu_1' H_1 - \mu_2' H_2] = 0)$$

$$\Leftrightarrow (\text{curl} - \lambda_1)(\text{curl} - \lambda_2) u = 0 \quad (D)$$

— View in the topology of F —



$F \uparrow$
no-self-organization

Singular Perturbation



DB model

$$u = G_+ G_+ + G_- G_-$$

($u = B$ and V)

$$\nabla \times G_{\pm} = \lambda_{\pm} G_{\pm}$$

$$\lambda_+ \gg \lambda_-$$

microscopic model : ε -model $\rightarrow f_{\varepsilon}$

$\downarrow \lim_{\varepsilon \rightarrow 0} ?$

macroscopic model : 0-model $\rightarrow f_0$

- Σ -model = 2-fluid MHD

$$\left\{ \begin{array}{l} \partial_t A = (\mathbf{v} - \underline{\varepsilon} \nabla \times \mathbf{B}) \times \mathbf{B} - \nabla(\phi - \underline{\varepsilon} p_e) \\ \partial_\star (\underline{\varepsilon} \mathbf{v} + A) = \mathbf{v} \times (\mathbf{B} + \underline{\varepsilon} \nabla \times \mathbf{v}) - \nabla(\underline{\varepsilon} \frac{v^2}{2} + \underline{\varepsilon} p_i + \phi) \end{array} \right.$$

$$\left(\varepsilon = \frac{c/\omega_{pi}}{l} \right)$$

- $\mathcal{O}$ -model = MHD

$$\left\{ \begin{array}{l} \partial_\star \mathbf{B} + (\mathbf{v} \cdot \nabla) \mathbf{B} - (\mathbf{B} \cdot \nabla) \mathbf{v} = 0 \\ \partial_\star \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} - (\mathbf{B} \cdot \nabla) \mathbf{B} = -\nabla(p + \frac{B^2}{2}) \end{array} \right.$$

DB field (static sol. of the ε -model)

$$\partial_* \mathcal{Q}_j - \nabla \times (\psi_j \times \mathcal{Q}_j) = 0$$

$$(e) \begin{cases} \mathcal{Q}_1 = B \\ \psi_1 = \psi - \varepsilon \nabla \times B \end{cases}$$

$$(i) \begin{cases} \mathcal{Q}_2 = B + \varepsilon \nabla \times \psi \\ \psi_2 = \psi \end{cases}$$

Beltrami field

$$\begin{cases} \mathcal{Q}_1 = a \psi_1 \\ \mathcal{Q}_2 = b \psi_2 \end{cases}$$

$$\Rightarrow \begin{cases} B = C_+ G_+ + C_- G_- \\ \psi = (a' + \varepsilon \lambda_+) C_+ G_+ + (a' + \varepsilon \lambda_-) C_- G_- \end{cases}$$

$$\lambda_{\pm} = \frac{1}{2\varepsilon} \left[(b - a') \pm \sqrt{(b - a')^2 - 4(1 - b/a)} \right]$$

$$\text{writing } b/a = 1 + \delta$$

$$\begin{cases} \lambda_- \approx \frac{\delta}{\varepsilon} \left(\frac{1}{2} - a \right)^{-1} \rightarrow 0(1) \\ \lambda_+ \approx -\frac{1}{\varepsilon} \left(\frac{1}{2} - a \right) \rightarrow \infty \end{cases}$$

How is the small scale created?

Characteristics of the 0-model.

$$(\nabla\psi)^2 (\partial_t\psi + v \cdot \nabla\psi) (\partial_t\psi + (v+B) \cdot \nabla\psi)^2 (\partial_t\psi + (v-B) \cdot \nabla\psi)^2 = 0$$

$\uparrow$ elliptic $\uparrow$ flow $\nwarrow \nearrow$ Alfvén

Quasi-static: $\partial_t = 0$

$$(\nabla\psi)^2 (v \cdot \nabla\psi) (B-v) \cdot \nabla\psi)^2 (B+v) \cdot \nabla\psi)^2 = 0$$

① non-integrability (chaos) of characteristics
(3D)

② non linearity (fields are unknown)

intersecting characteristics

"structures" $\leftarrow$ $\updownarrow$ ($\rightarrow$ shocks)

dispersion = singular perturbation

Summary

- ∞ -dim vector space
 - different norms induce different topology
- coercive forms
 - measure fluctuations and dissipation
 - controle self-organization
 - const. of motion must be adjusted
($\leftarrow$ self-organisation)
 - controle stability (Lyapunov function)

DB model

Mahajan - Y., PRL 81 (1998), 4863.

Variational Principle

Y. - Mahajan, PRL. 88 (2002), 095001.

Beltrami Fields

Y. - Giga, Math. Z. 204 (1990), 235.

Y., J. Math. Phys. 33 (1992), 1252.

Y. - Mahajan, J. Math. Phys. 40 (1999), 5080.