

SMR 1232 - 14

**XII WORKSHOP ON
STRONGLY CORRELATED ELECTRON SYSTEMS**

17 - 28 July 2000

***DO WE KNOW EVERYTHING ABOUT
THE MAGNETISM OF CUPRATES?***

Gabriel AEPPLI
N.E.C. Research Institute Inc.
NJ-08540 Princeton, U.S.A.

These are preliminary lecture notes, intended only for distribution to participants.

Do we know everything about the magnetism of cuprates?

Beyond Heisenberg

R.Coldea(ORNL,RAL)

P.Dai(ORNL)

H.A.Mook(ORNL)

S.M.Hayden(Bristol)

Z.Fisk(NHMFL)

S-W.Cheong(Lucent/Rutgers)

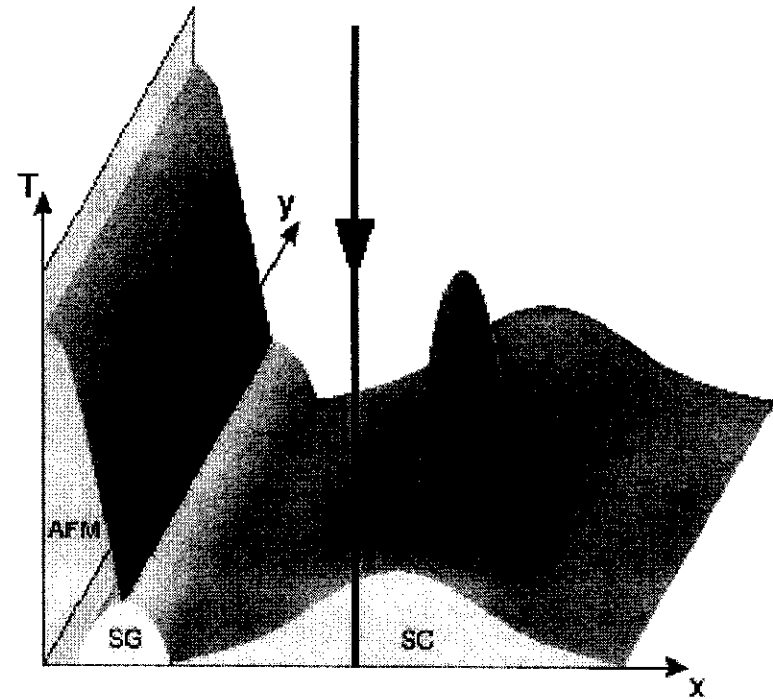
T.G. Perring(RAL)

F.Dogan(UW)

T.E.Mason(ORNL)

Things we know

- 2-dimensionality
- phase diagram
- Heisenberg model with strong nearest neighbor exchange



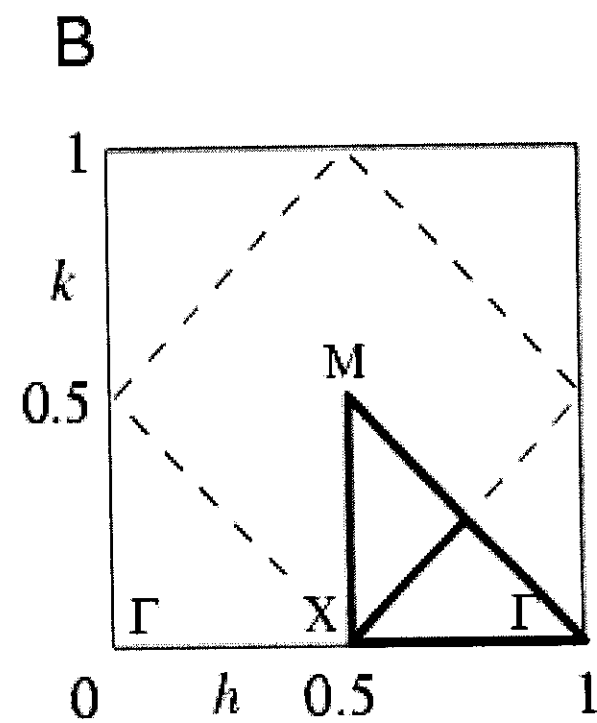
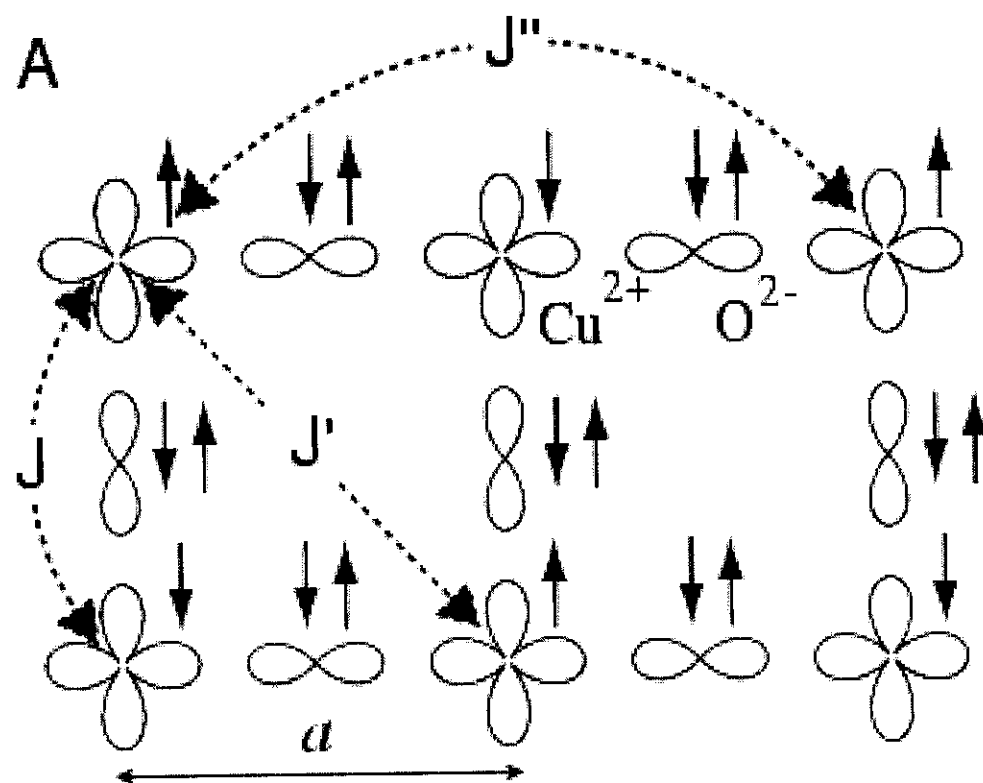
From “Nearly Singular Magnetic Fluctuations in the Normal State of a High- T_c Cuprate Superconductor”, G. Aeppli, T. E. Mason, S. M. Hayden, H. A. Mook, J. Kulda, *Science* **278**, pp. 1432-1435 (1997)

Naïve(Heisenberg)picture of spin physics in hi- T_c materials

- Starting point is AFM Heisenberg Hamiltonian
- a QCP between AFM and Spin Fluid phases with pairing built in
- damping to take account of metallicity
- appropriate embellishments for spare ground states(e.g. black mountain) found in some alloy series

Is Heisenberg picture all we need to understand spin sector of hi- T_c problem?

- Revisit insulator(Dai et al,Cond Mat 2000)
- look at excitations in the s.c.(Coldea et al,Cond Mat 2000)



Spin wave theory

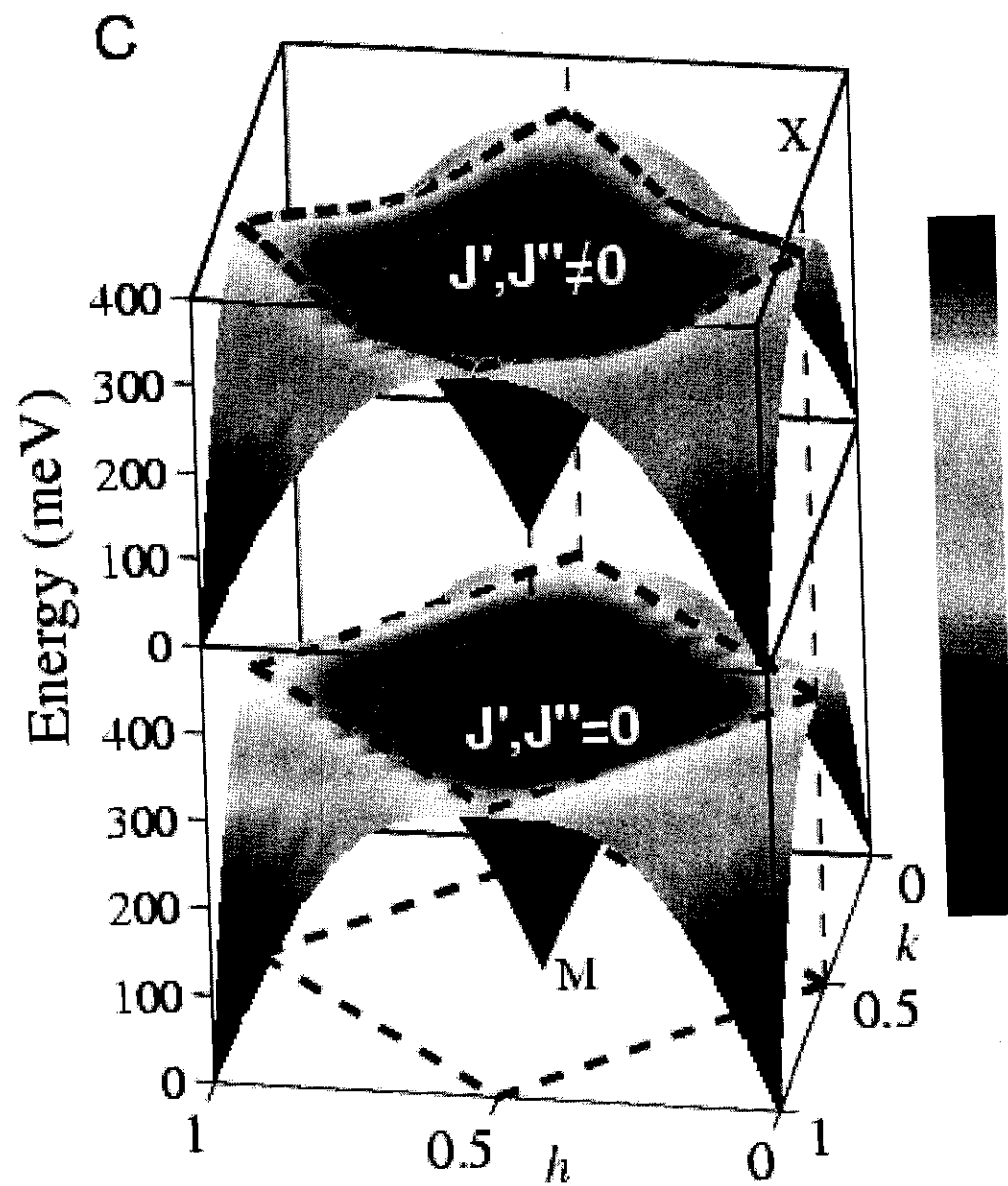
$$\mathcal{H} = \sum_{i,\delta} J \mathbf{S}_i \cdot \mathbf{S}_{i+\delta} + \sum_{i,\delta'} J' \mathbf{S}_i \cdot \mathbf{S}_{i+\delta'} + \sum_{i,\delta''} J'' \mathbf{S}_i \cdot \mathbf{S}_{i+\delta''}$$

$$\hbar\omega_Q = 2Z_c J \sqrt{A_Q^2 - B_Q^2}$$

$$A_Q = 1 - \frac{J'}{J} (1 - \cos(2\pi h) \cos(2\pi k)) - \frac{J''}{J} \left(1 - \frac{\cos(4\pi h) + \cos(4\pi k)}{2} \right)$$

$$B_Q = \frac{\cos(2\pi h) + \cos(2\pi k)}{2}$$

Z_c =quantum renormalization factor =1.18



Why neutrons?

- Electrons have charge and spin, many excellent probes of charge sector(e.g.electron & photon spectroscopies described elsewhere in this meeting), many interesting properties of solids in spin sector(e.g. magnetism)
- want neutral & easily interpretable probe of spin sector

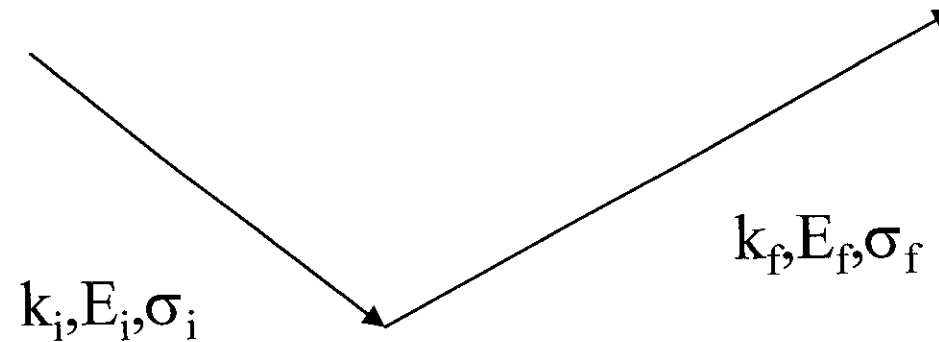
Neutrons are an ideal probe because-

- they carry spin(1/2) but not charge
- $E = \hbar^2 k^2 / 2m$ with $m = \text{mass of proton} = \text{mass of hydrogen}$ so that
 $\lambda = \text{interatomic spacings}$
- wide spectral range with modern sources
 $1 \text{ meV} < E < 1 \text{ eV}$
- weak coupling, so simple expressions for cross-sections

Some useful facts

- $\hbar^2/2m=2.07\text{meV}\cdot\text{\AA}^2$
- $1\text{meV}=11.6\text{K}$
- $4.14\text{meV}=1\text{THz}=10^{12}\text{Hz}$
- a 5 meV neutron has a wavelength of 4.05Å
& velocity of 976m/sec
- $m=1.82\times 10^3 m_e$

Scattering experiments



$$Q = k_i - k_f$$

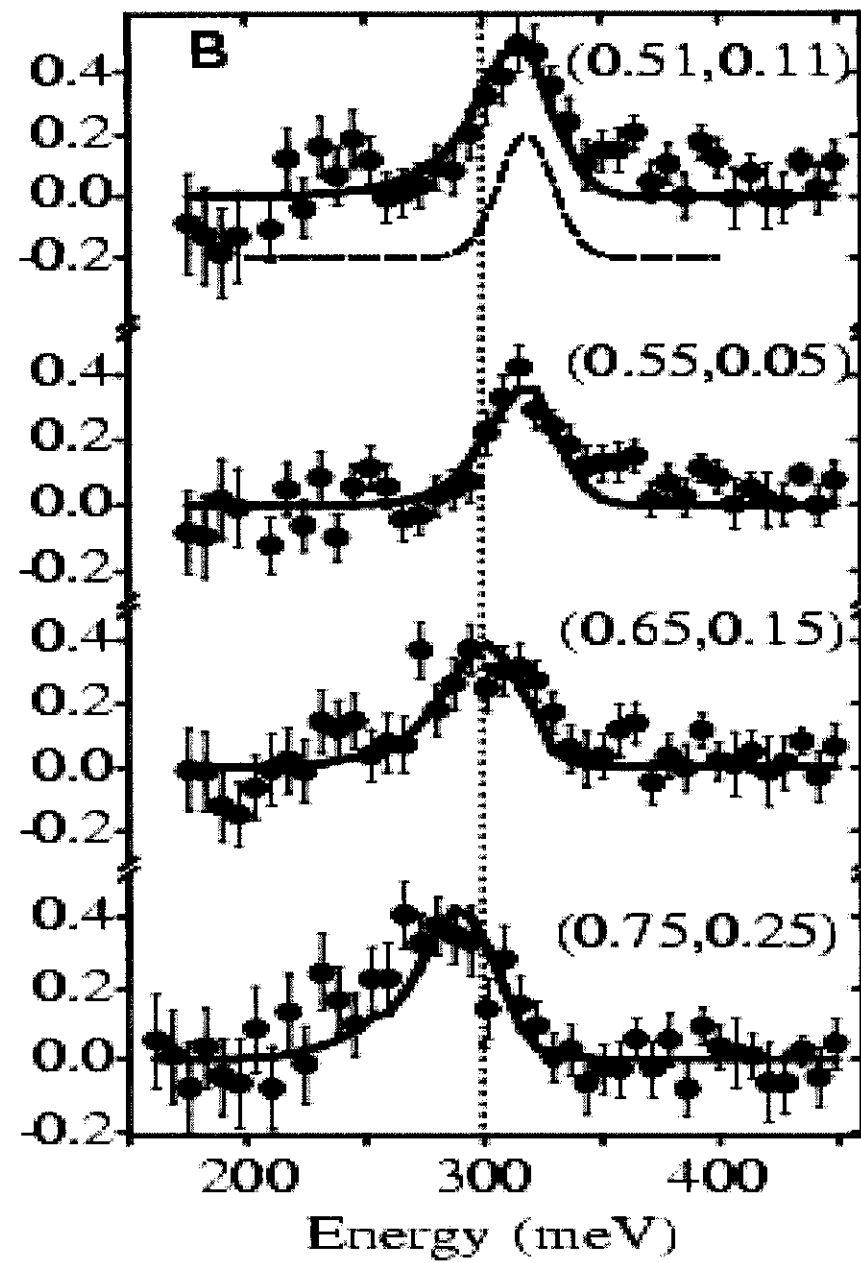
$$\hbar\omega = E_i - E_f$$

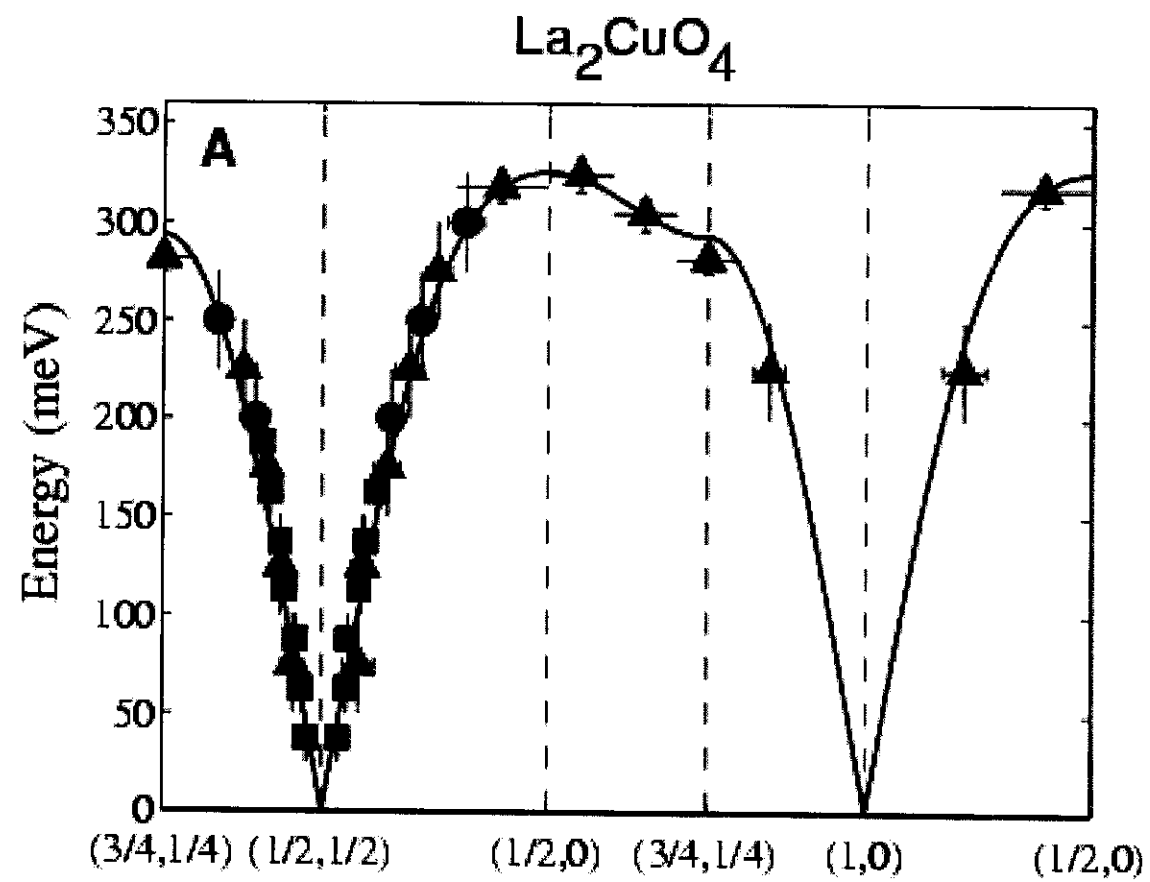
Measure differential cross-section = ratio of outgoing flux per unit solid angle and energy to ingoing flux = $\delta^2\sigma/\delta\Omega\delta\omega$

inelastic neutron scattering

Fermi's Golden Rule

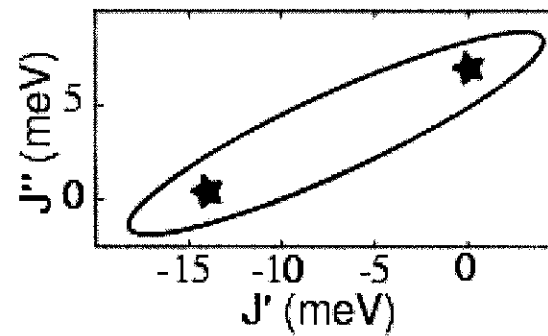
- at $T=0$,
 $\delta^2\sigma/\delta\Omega\delta\omega = \sum_f |\langle f | S(Q)^+ | 0 \rangle|^2 \delta(\omega - E_0 + E_f)$ where
 $S(Q)^+ = \sum_m S_m^+ \exp i \mathbf{q} \cdot \mathbf{r}_m$
- for finite T , $\delta^2\sigma/\delta\Omega\delta\omega = k_f/k_i S(Q, \omega)$ where
 $S(Q, \omega) = n(\omega + 1) \text{Im} \chi(Q, \omega)$





$$\mathcal{H} = \sum_{i,\delta} J \mathbf{S}_i \cdot \mathbf{S}_{i+\delta} + \sum_{i,\delta'} J' \mathbf{S}_i \cdot \mathbf{S}_{i+\delta'} + \sum_{i,\delta''} J'' \mathbf{S}_i \cdot \mathbf{S}_{i+\delta''}$$

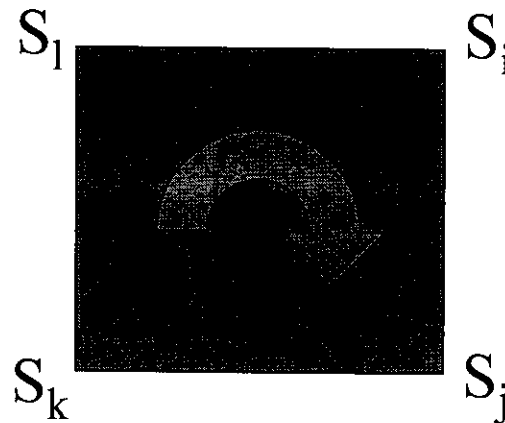
$$J=0.11\text{eV}(J''=0), 0.14\text{eV}(J'=0)$$



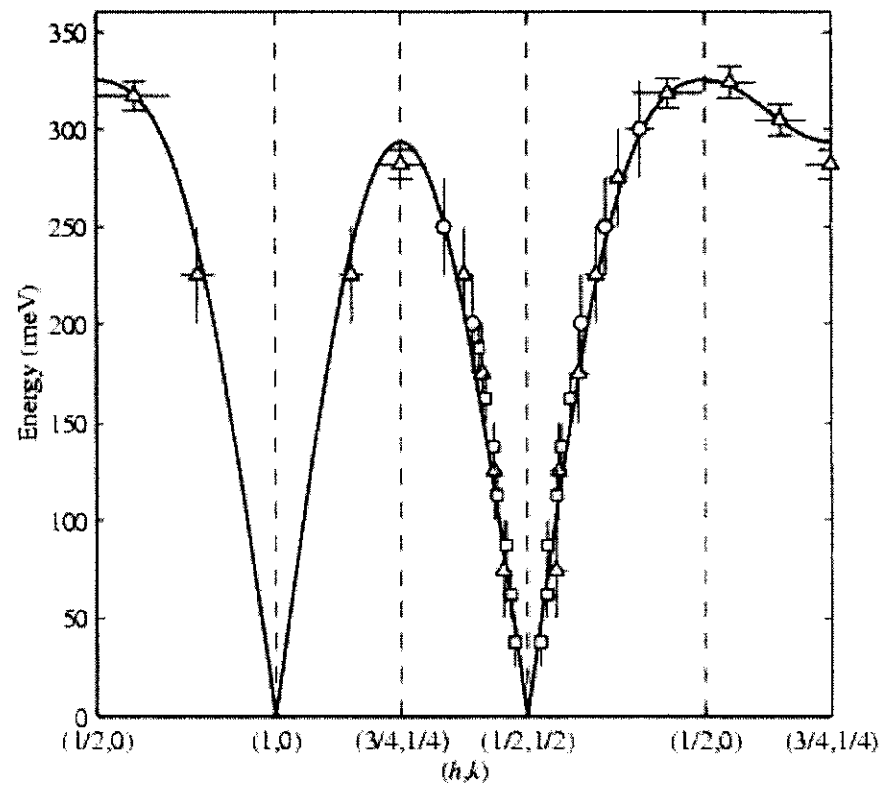
A ferromagnetic next nearest neighbor coupling

- Not expected
- are we using the wrong Hamiltonian?
- consider ring exchange terms which provide much better fit to small cluster calculations and explain light scattering anomalies , i.e.

$$H = \sum J S_i S_j + J_{\text{ring}} S_i S_j S_k S_l$$



La_2CuO_4 (un-renormalised J's) $J=162.7$, $J_c=64.32$



More detail on J_c

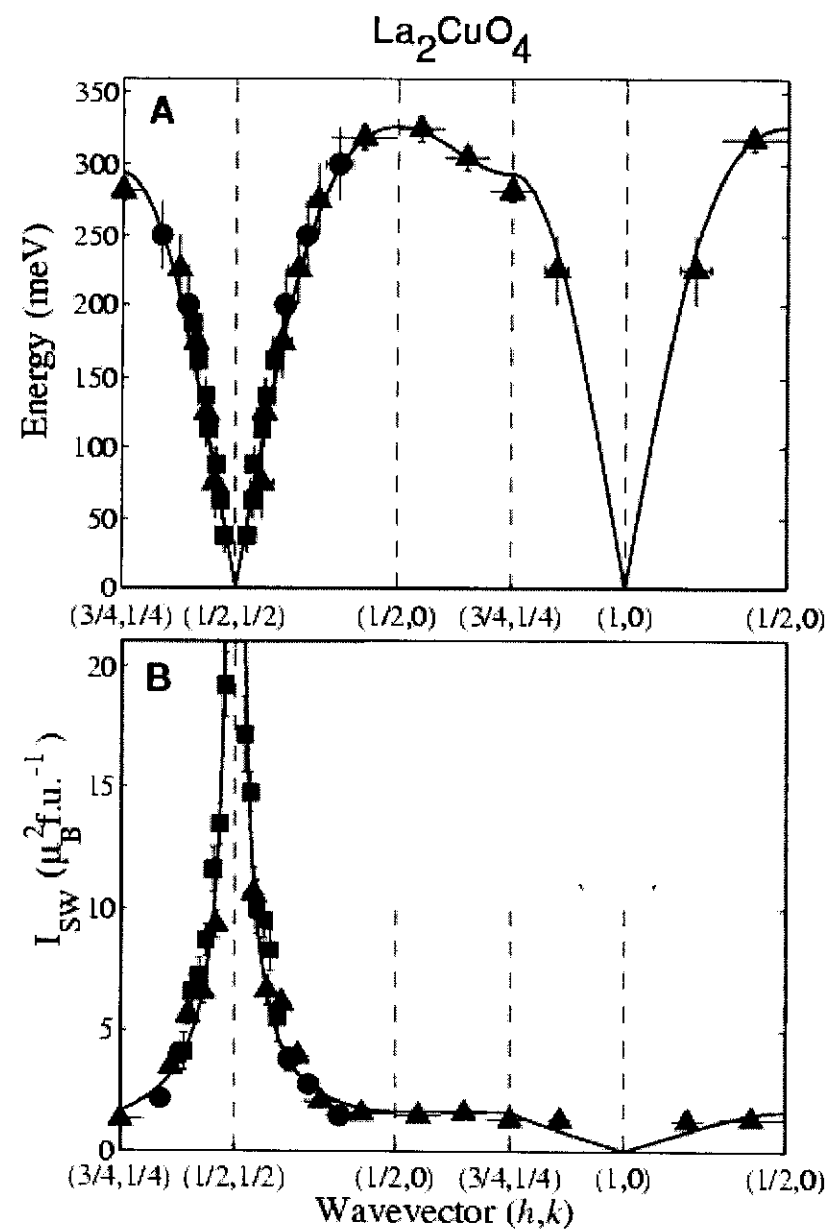
$$H = J \sum_{i,j} \mathbf{S}_i \cdot \mathbf{S}_j + J' \sum_{i,i'} \mathbf{S}_i \cdot \mathbf{S}_{i'} + J'' \sum_{i,i''} \mathbf{S}_i \cdot \mathbf{S}_{i''} \\ + J_c \sum_{i,j,k,l} [(\mathbf{S}_i \cdot \mathbf{S}_j)(\mathbf{S}_k \cdot \mathbf{S}_l) + (\mathbf{S}_i \cdot \mathbf{S}_l)(\mathbf{S}_k \cdot \mathbf{S}_j) - (\mathbf{S}_i \cdot \mathbf{S}_k)(\mathbf{S}_j \cdot \mathbf{S}_l)]$$

$$J=143.3 \pm 5 \text{ meV and } J_c=48.4 \pm 9 \text{ meV.}$$

$$H = -t \sum_{i,j,\sigma=-} (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}) + U \sum_i n_{i-} n_{i+}$$

$$J = 4t^2/U - 24t^4/U^3 \quad J_c = 80t^4/U^3 \quad \text{and} \quad J' = J'' = 4t^4/U^3$$

$$t=0.32 \pm 0.02 \text{ eV and } U=2.6 \pm 0.3 \text{ eV}$$



GA-NEC-Trieste

Conclusions- parent insulator- need to go beyond Heisenberg

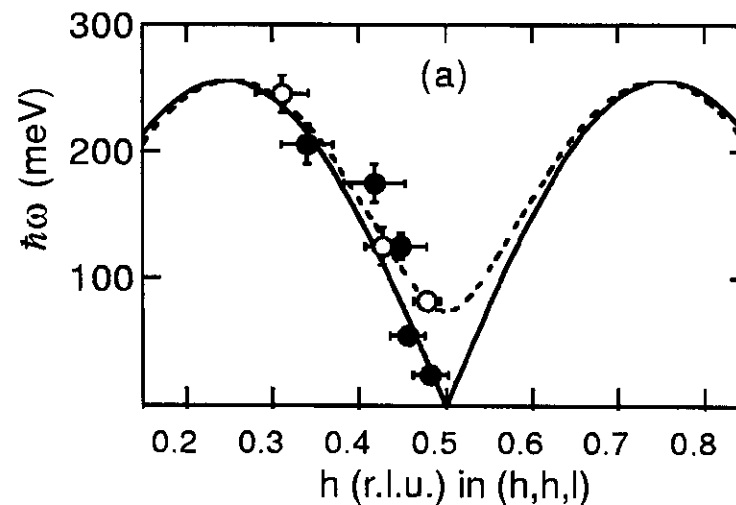
- Substantial(on scale of T_c) further neighbor interactions
- **J_{ring} a term to worry about** - data most consistent with cluster calculations with cyclic(4-spin) interactions(e.g. Schmidt '90 with $J_{\text{ring}}/J=30\%$)
- remarkable consistency also with Hubbard model parameters from electronic spectroscopies - **charge neutral low energy probe gives independent estimate of effective electronic parameters t & U**
- once further neighbor interactions are taken into account, spin wave theory gives excellent account of absolute amplitudes as well as dispersion

What happens for Y123

- Higher T_c
- ‘Cleaner’
- easier to adjust stoichiometry
- many more data of other types

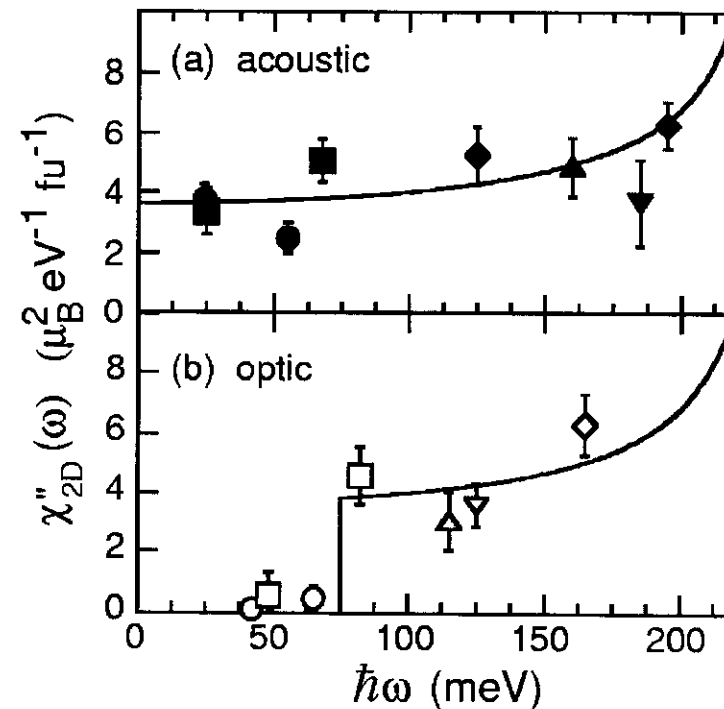
Y123 is a bilayer AFM -bilayers give rise to acoustic(in-phase) & optic(out-of-phase) spin waves

“High-Frequency Spin
Waves in $\text{YBa}_2\text{Cu}_3\text{O}_{6.15}$ ”,
S. M. Hayden, G. Aeppli,
T. G. Perring, H. A. Mook,
F. Dogan, *Phys. Rev. B*
54(10), pp. R6905-6908,
(1996)



Q-averaged spin wave spectral weights

“High-Frequency Spin Waves in $\text{YBa}_2\text{Cu}_3\text{O}_{6.15}$ ”, S. M. Hayden, G. Aeppli, T. G. Perring, H. A. Mook, F. Dogan, *Phys. Rev. B* **54**(10), pp. R6905-6908, (1996)



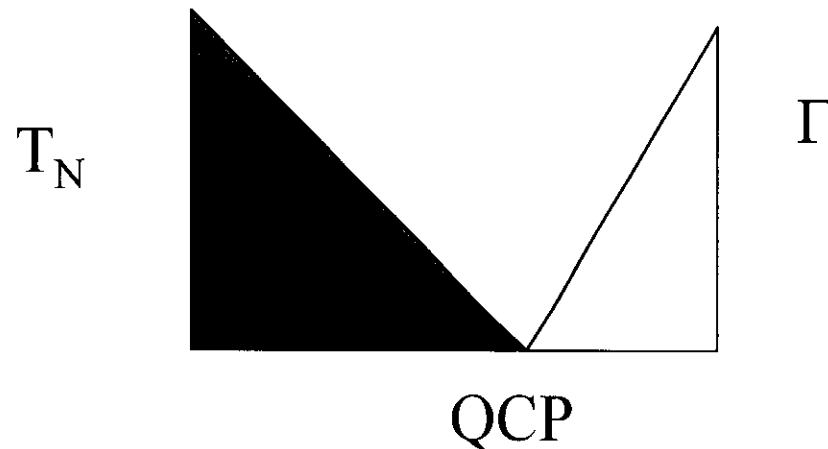
- Spin wave dispersion for $\text{Y123O}_{6.15}$ & *amplitudes* very well described by conventional theory with significant coupling between planes in bilayers and in-plane J slightly less than for La_2CuO_4
- acoustic DOS is flat for low E , gap in optic DOS
- what happens on doping?

Might expect (simple quantum critical point for a metal emerging from an AFM)-

- Replacement of AFM Bragg peaks and SW peaks by quasielastic acoustic term and damped optical term, i.e. $\delta(\omega) + \text{SW} \rightarrow$

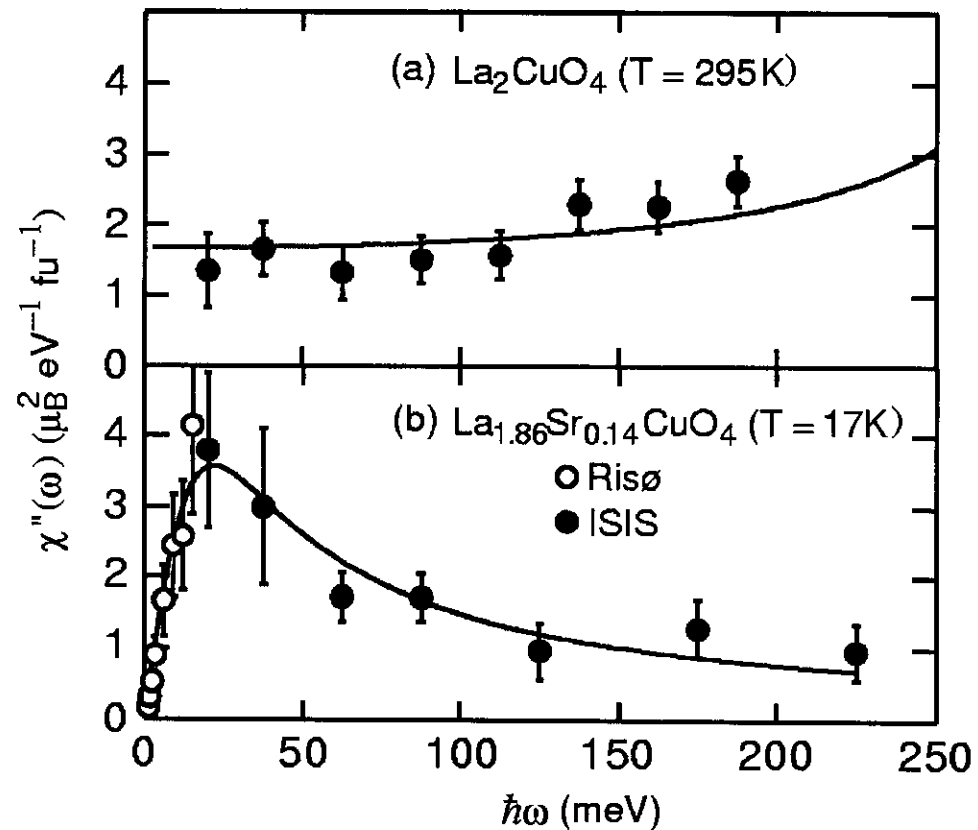
$$\text{Im}\chi_{\text{acoustic}} = \Gamma\omega/(\omega^2 + \Gamma^2) \quad \text{and}$$

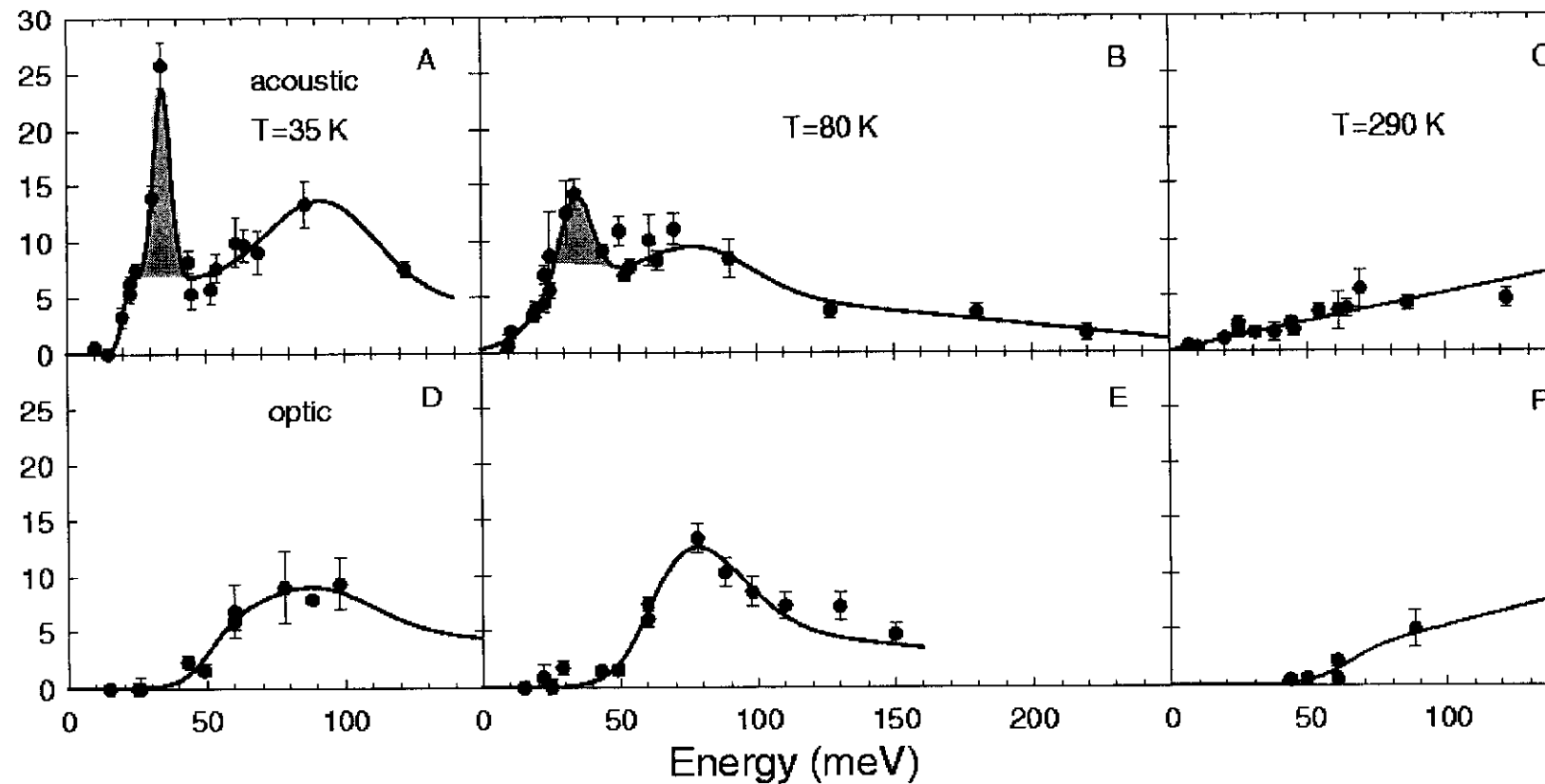
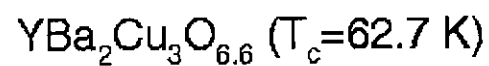
$$\text{Im}\chi_{\text{optic}} = \Gamma\omega[1/((\omega - \omega_o)^2 + \Gamma^2) + 1/((\omega + \omega_o)^2 + \Gamma^2)]$$



This is actually what seems to happen for $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$

“Comparison of the High-Frequency Magnetic Fluctuations in Insulating and Superconducting $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ”, S. M. Hayden, G. Aeppli, H. A. Mook, T. G. Perring, T. E. Mason, S-W. Cheong, Z. Fisk, *Phys. Rev. Letts.* **76(8)**, pp. 1344-1347 (1996).



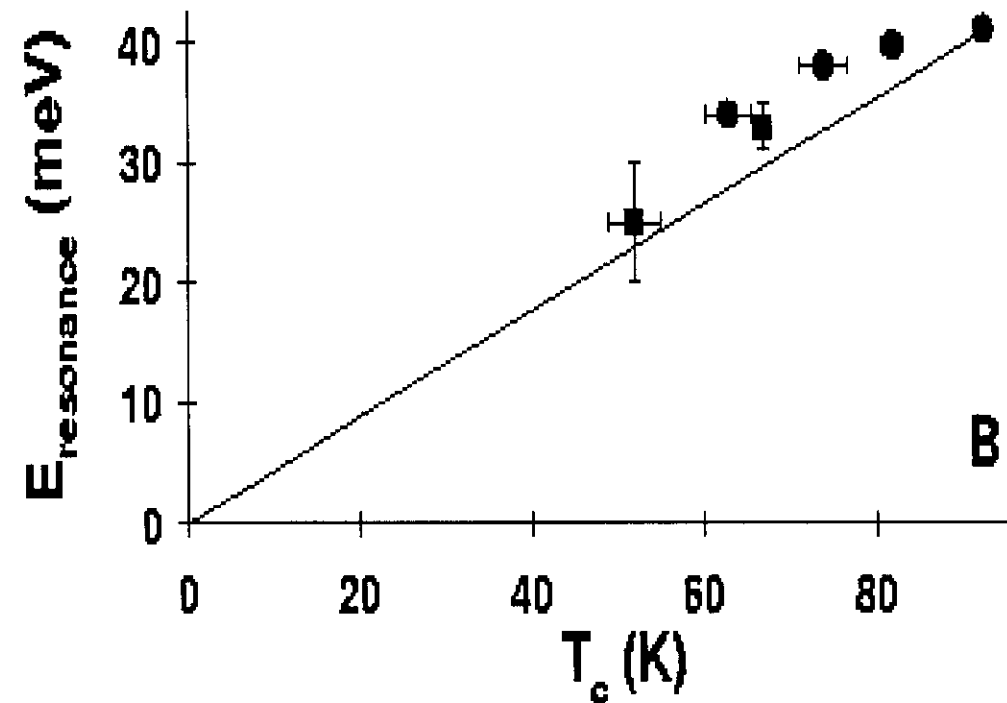


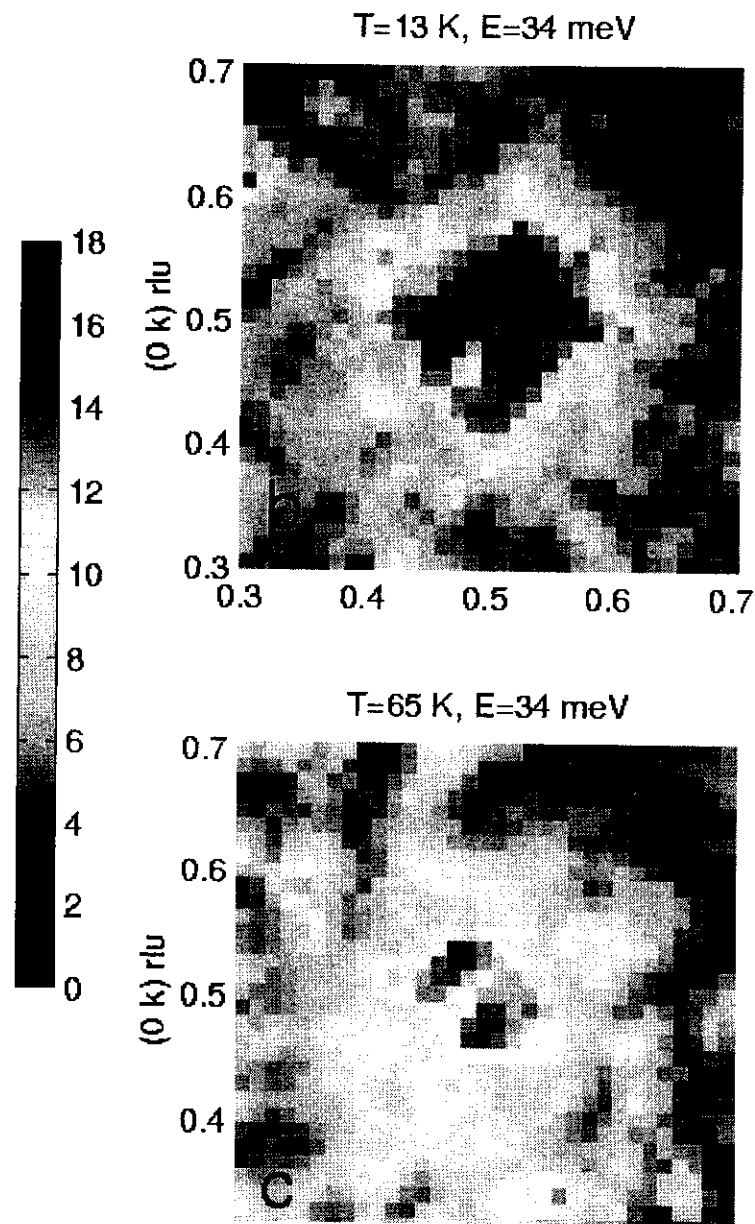
“The Magnetic Excitation Spectrum and Thermodynamics of High- T_c Superconductors”, P. Dai, H. A. Mook, S. M. Hayden, G. Aeppli, T. G. Perring, R. D. Hunt, F. Dogan, *Science* **284**, pp. 1344-1347, (1999).

Y123 seems much more complicated

- Optic mode(expected)
- spin gap in SC phase(expected because of spin singlet pairing in BCS-type states)
- pseudogap in normal phase, with onset T^* (not expected)
- ‘resonance’ (not expected)

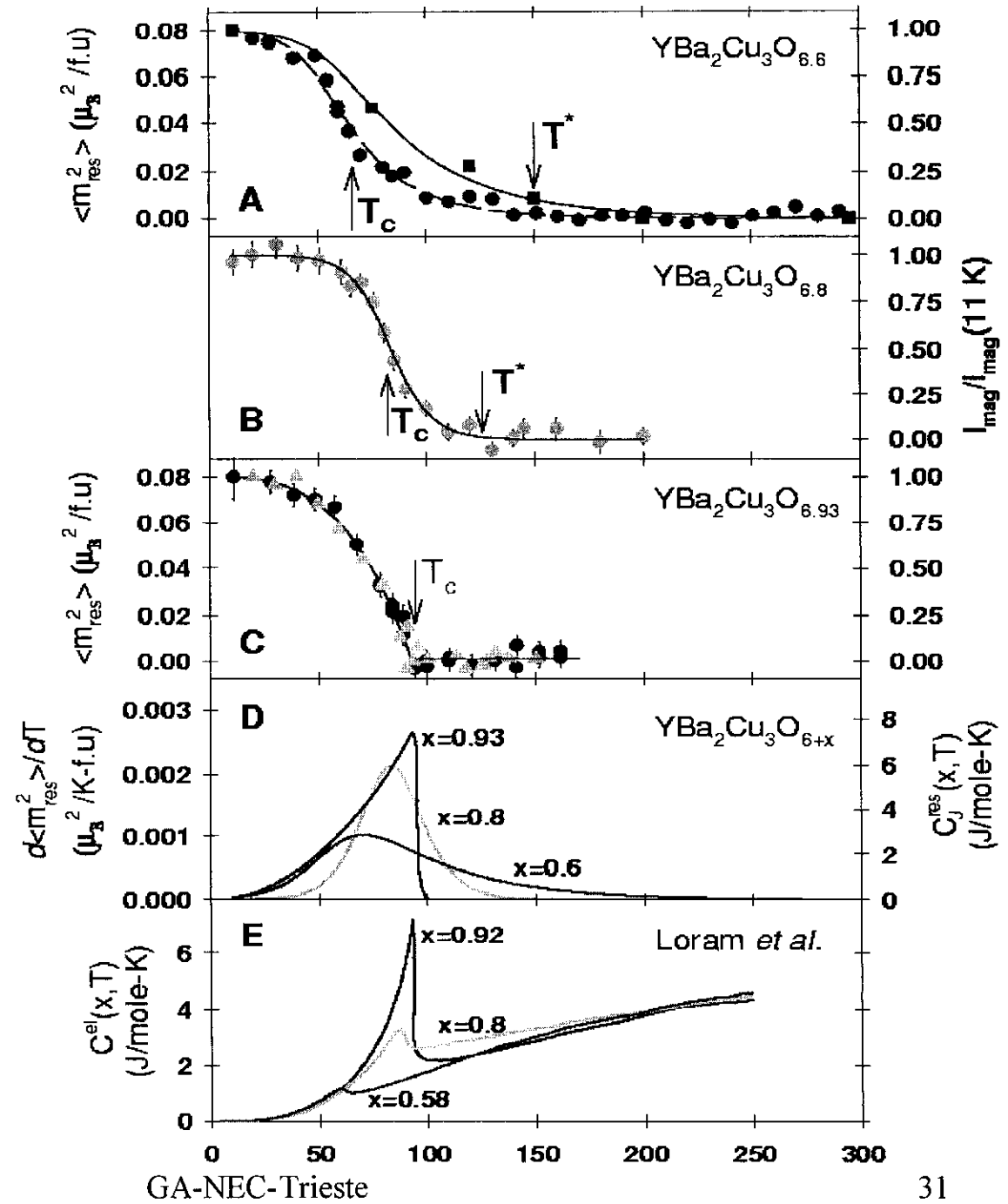
Focus on ‘resonance’ because it is easy to see (but beware that it contains only $\sim 1\%$ of $S(S+1)$)



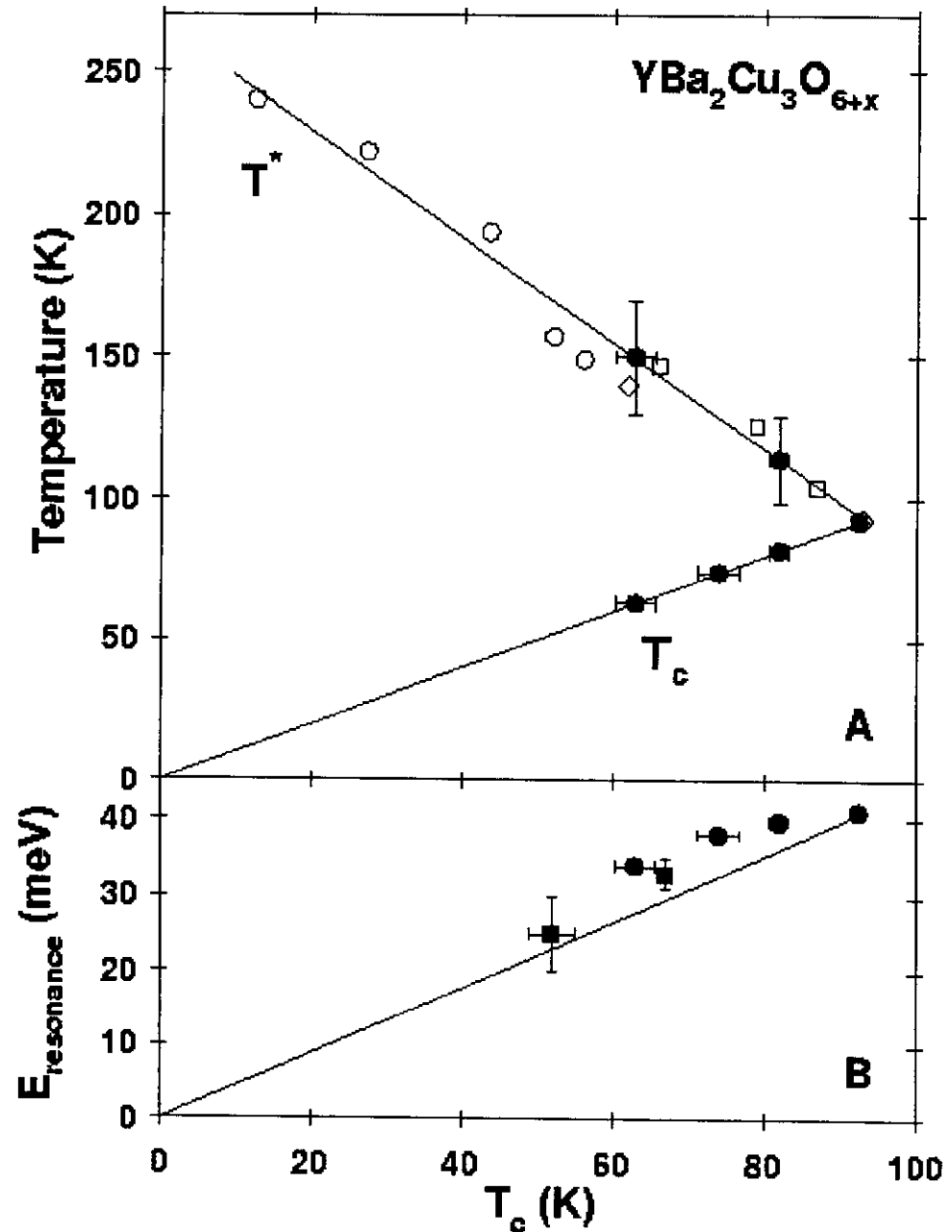


“Spin Fluctuations
in $\text{YBa}_2\text{Cu}_3\text{O}_{6.6}$ – A
New Picture”, H.
A. Mook, P. Dai, S.
M. Hayden, G.
Aeppli, T. G.
Perring, F. Dogan,
Nature **395**, pp.
580-582, (1998)

“The Magnetic
Excitation
Spectrum and
Thermodynamics
of High- T_c
Superconductors”
, P. Dai, H. A.
Mook, S. M.
Hayden, G.
Aeppli, T. G.
Perring, R. D.
Hunt, F. Dogan,
Science **284**, pp.
1344-1347,
(1999).



“The Magnetic
 Excitation
 Spectrum and
 Thermodynamics
 of High- T_c
 Superconductors”,
 P. Dai, H. A.
 Mook, S. M.
 Hayden, G. Aeppli,
 T. G. Perring, R. D.
 Hunt, F. Dogan,
Science **284**, pp.
 1344-1347, (1999).

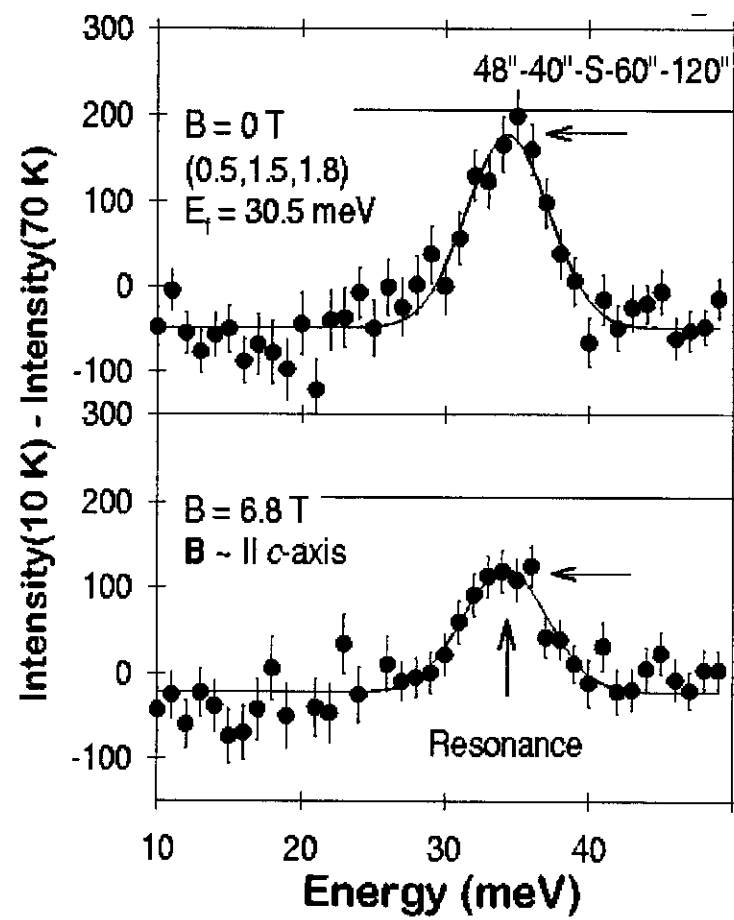
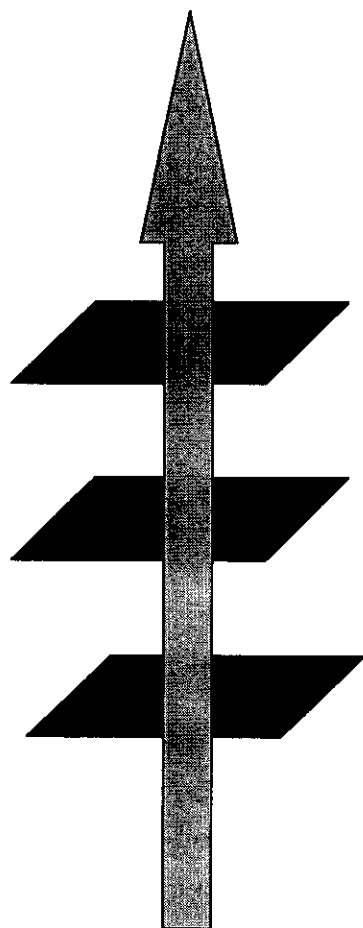


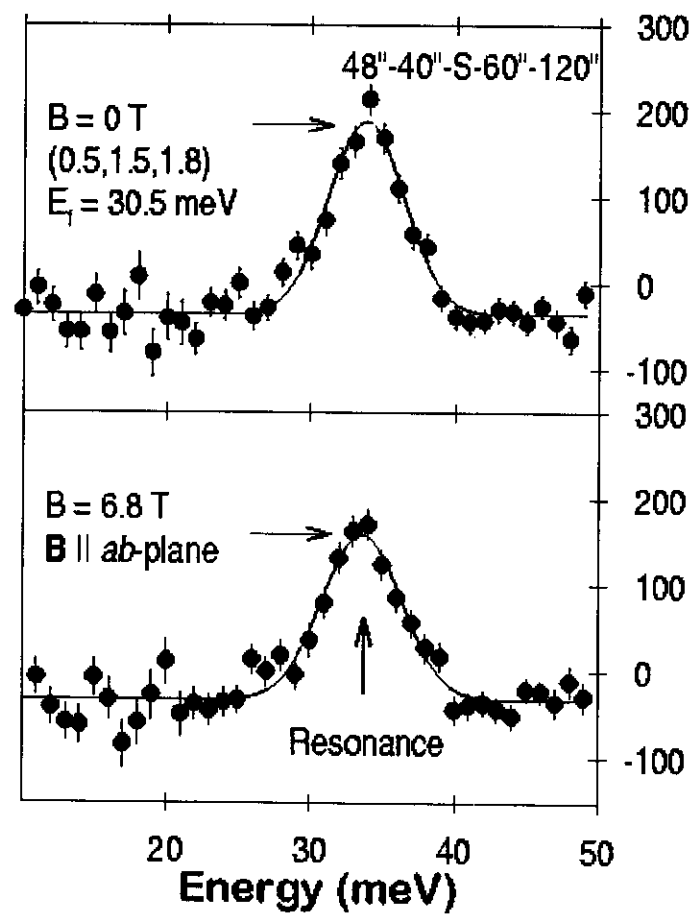
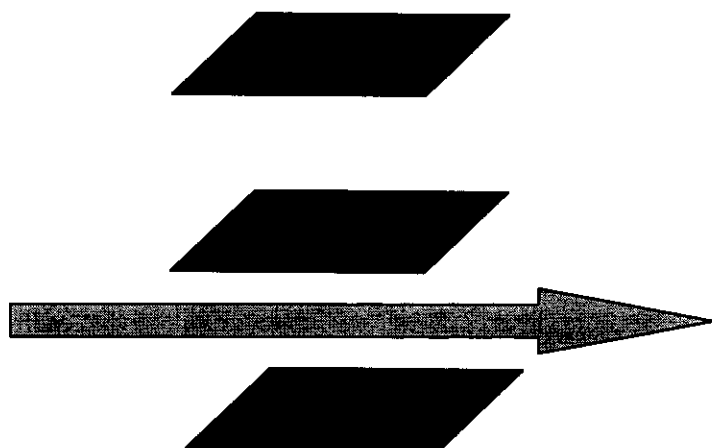
resonance

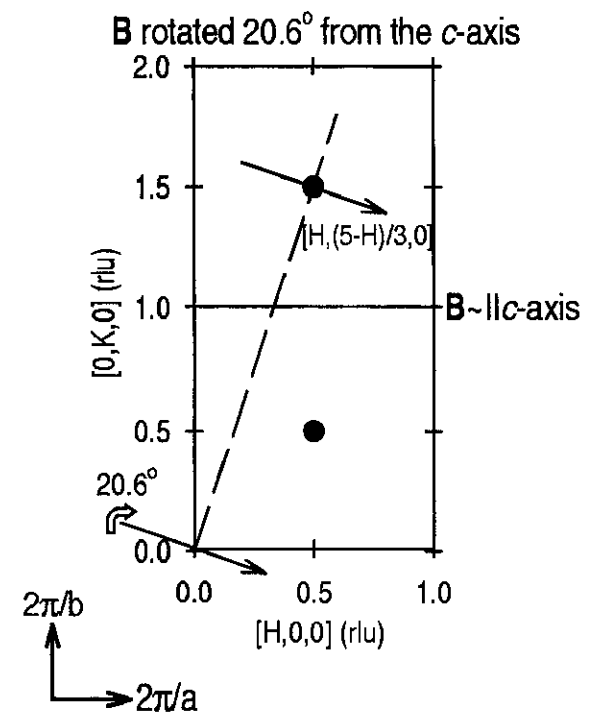
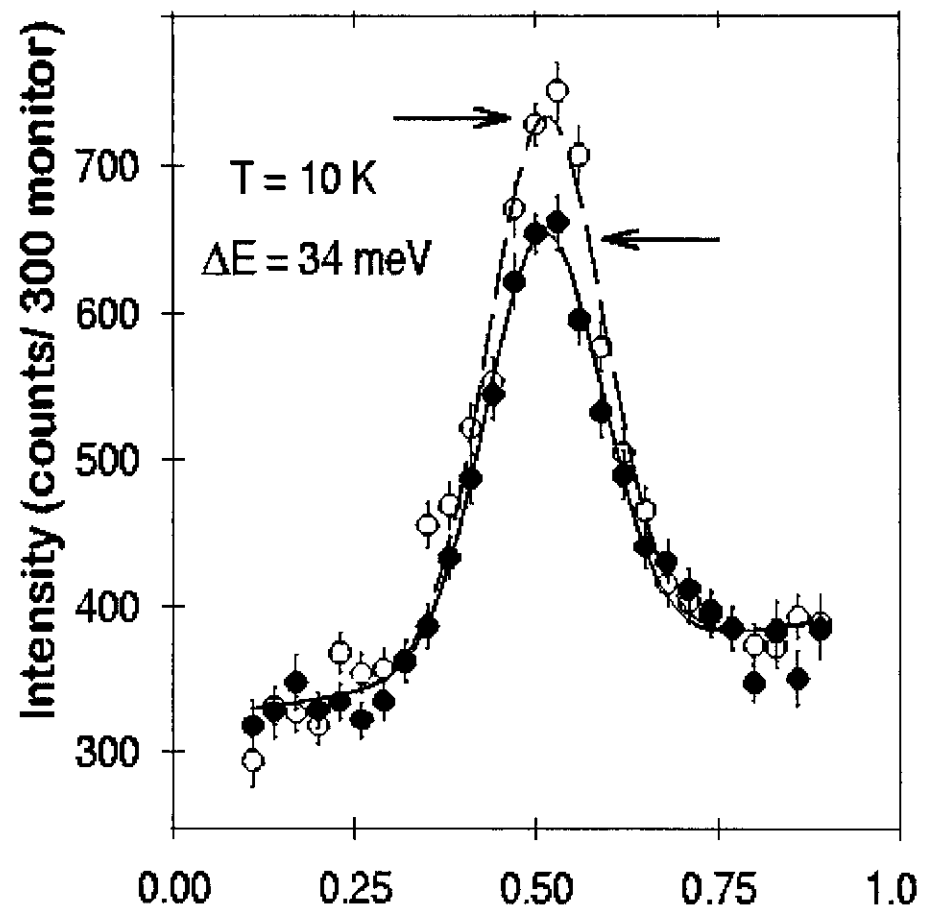
- an extra feature in the *acoustic* mode spectrum beyond spin gap
- seems connected with s.c., and pseudogap regime above T_c
- $dI(\text{resonance})/dT \sim C$ as doping, T is varied
- sharper in Q,E in s.c. phase than in normal state

Are all of these properties simply related to spin pairing(from some underlying Heisenberg picture) or is there something more profound?

- test by applying a magnetic field - if spin pairing is the only physics, expect Zeeman effect(which should be isotropic) to dominate, if orbital effects matter, will see anisotropic effects at energy scales different from $g\mu_B H$ (Dai et al., Cond Mat 2000)





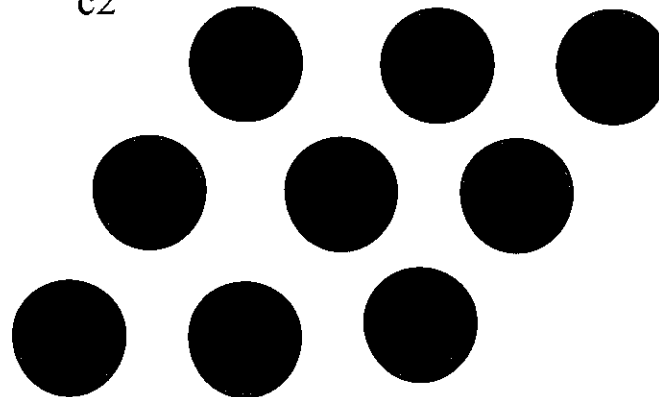


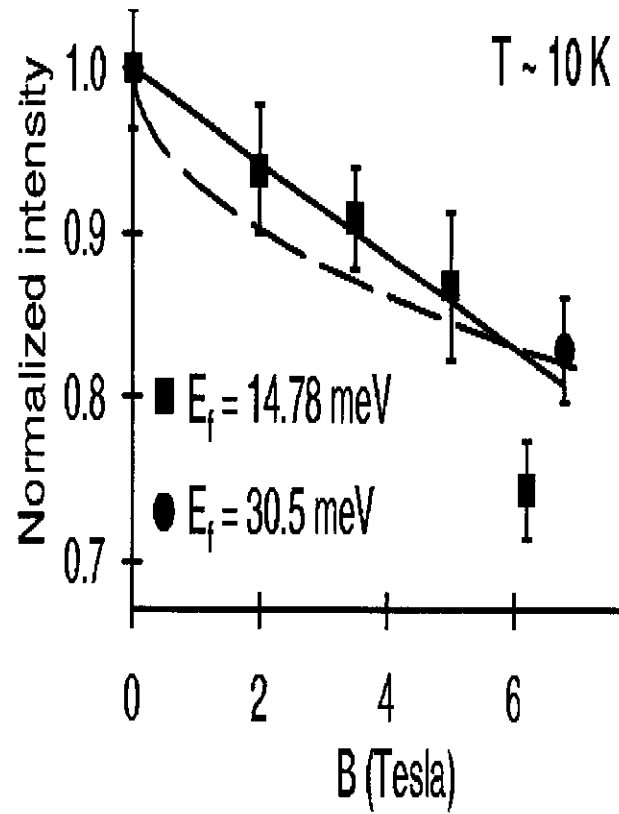
Very modest field, very strong effect

- 30% effect at 34meV for $H=6.5\text{T}$ along c
 $6.5\text{T} \sim 0.7\text{meV} \sim 7\text{K} \ll T_c \sim 60\text{K} \ll E_{\text{res}} \sim 360\text{K}$
- much smaller effect for $H=6.8\text{T}$ along CuO_2 planes
- conclude that not a spin(Zeeman) effect, but an orbital effect

In a superconductor, most natural ‘orbital’ effect is from vortex inclusions

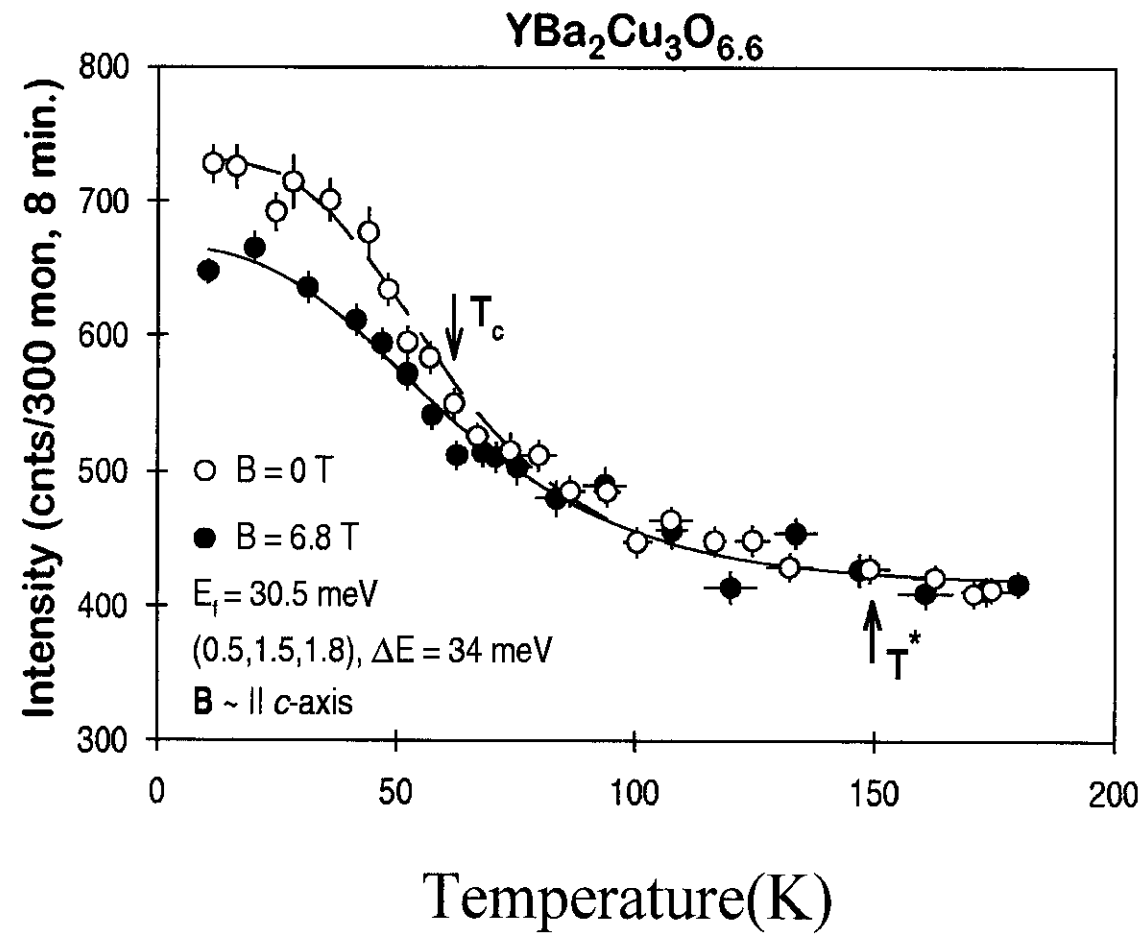
- vortices ~ normal state inclusions
- sc state has sharp resonance and normal state has no sharp resonance
- resonance intensity ~ 1 - volume fraction for normal state inclusions ~ $1 - B/B_{c2}$



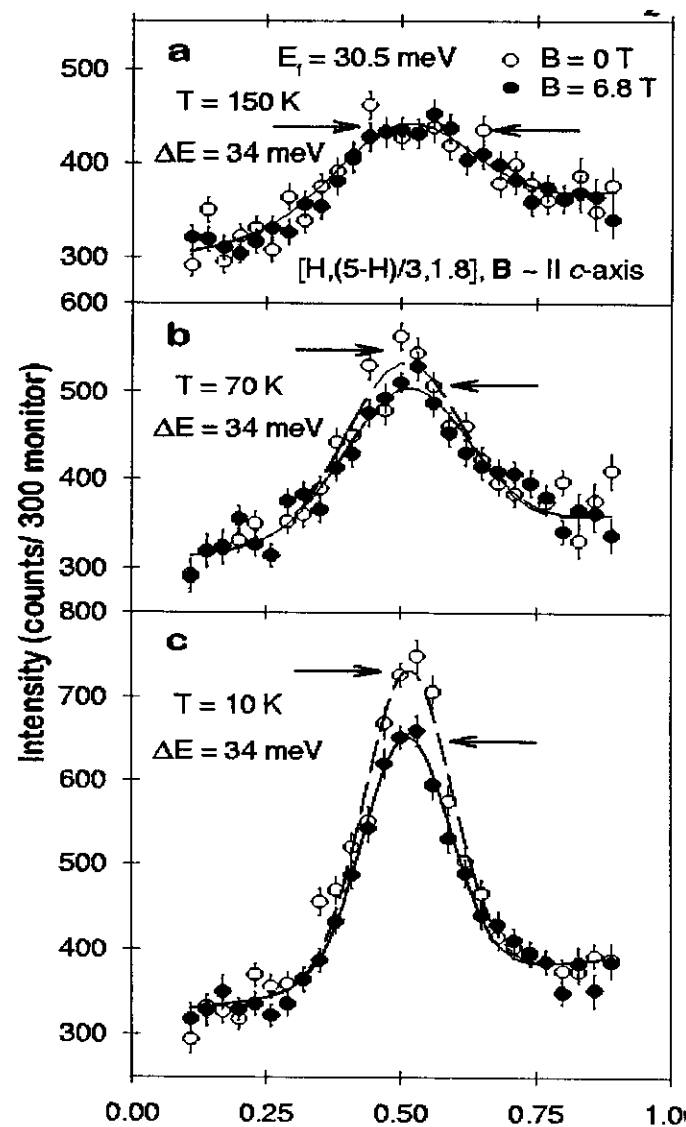


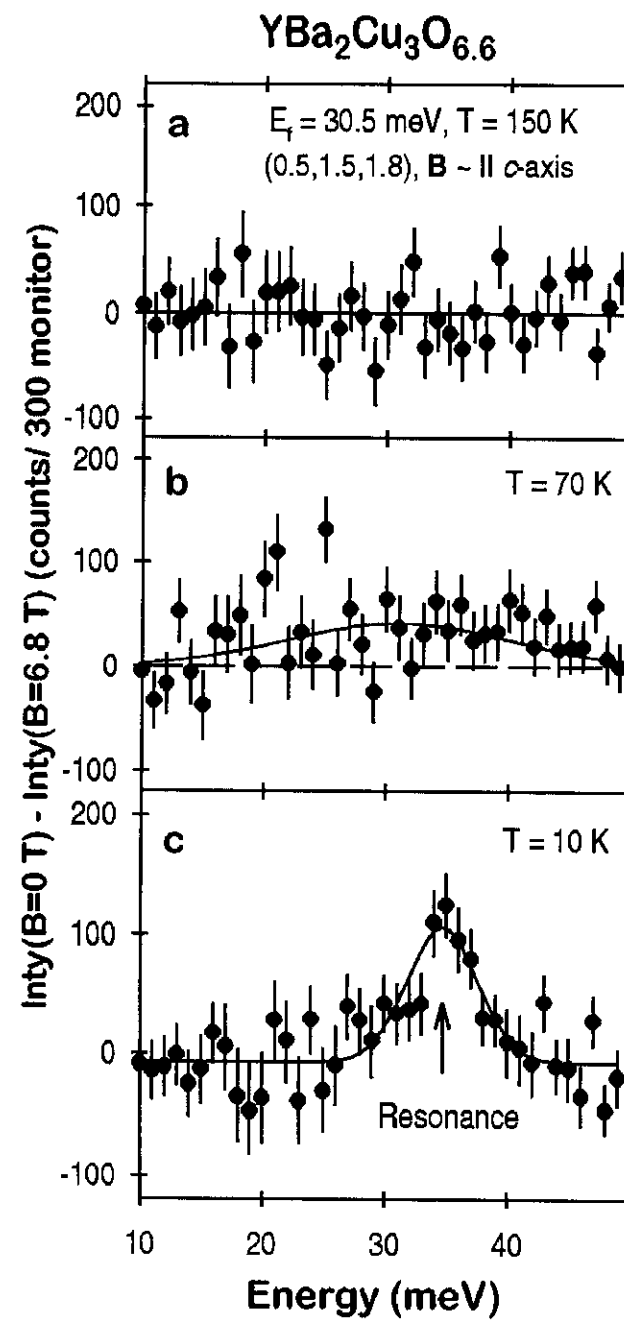
$$I/I_0 = 1 - B/B_{\text{char}} \text{ where}$$

$$B_{\text{char}} = 36\text{ T} \sim B_{c2} = 45\text{ T}$$

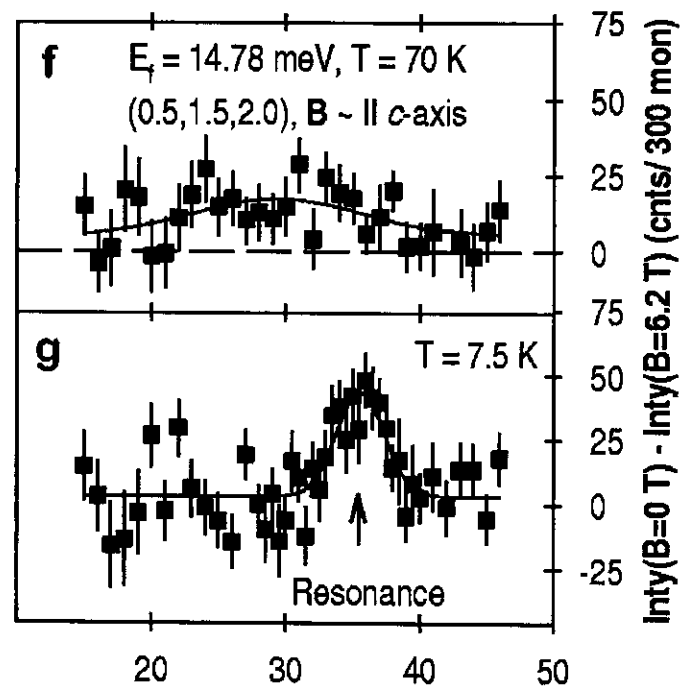


Is there really life above T_c ?

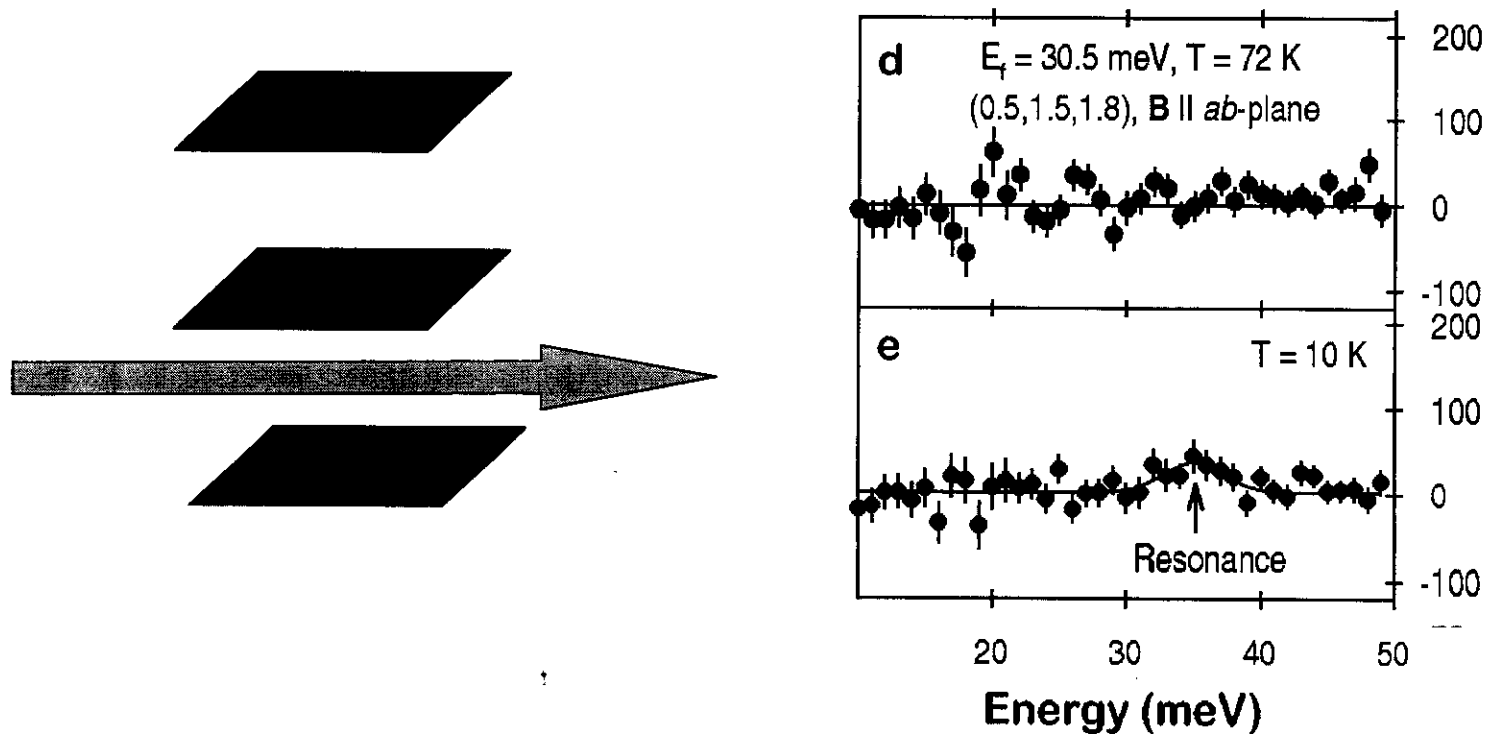




Yes, but let's change instrument configuration to make sure



What about anisotropic response to field?



Orbital effect for $T > T_c$

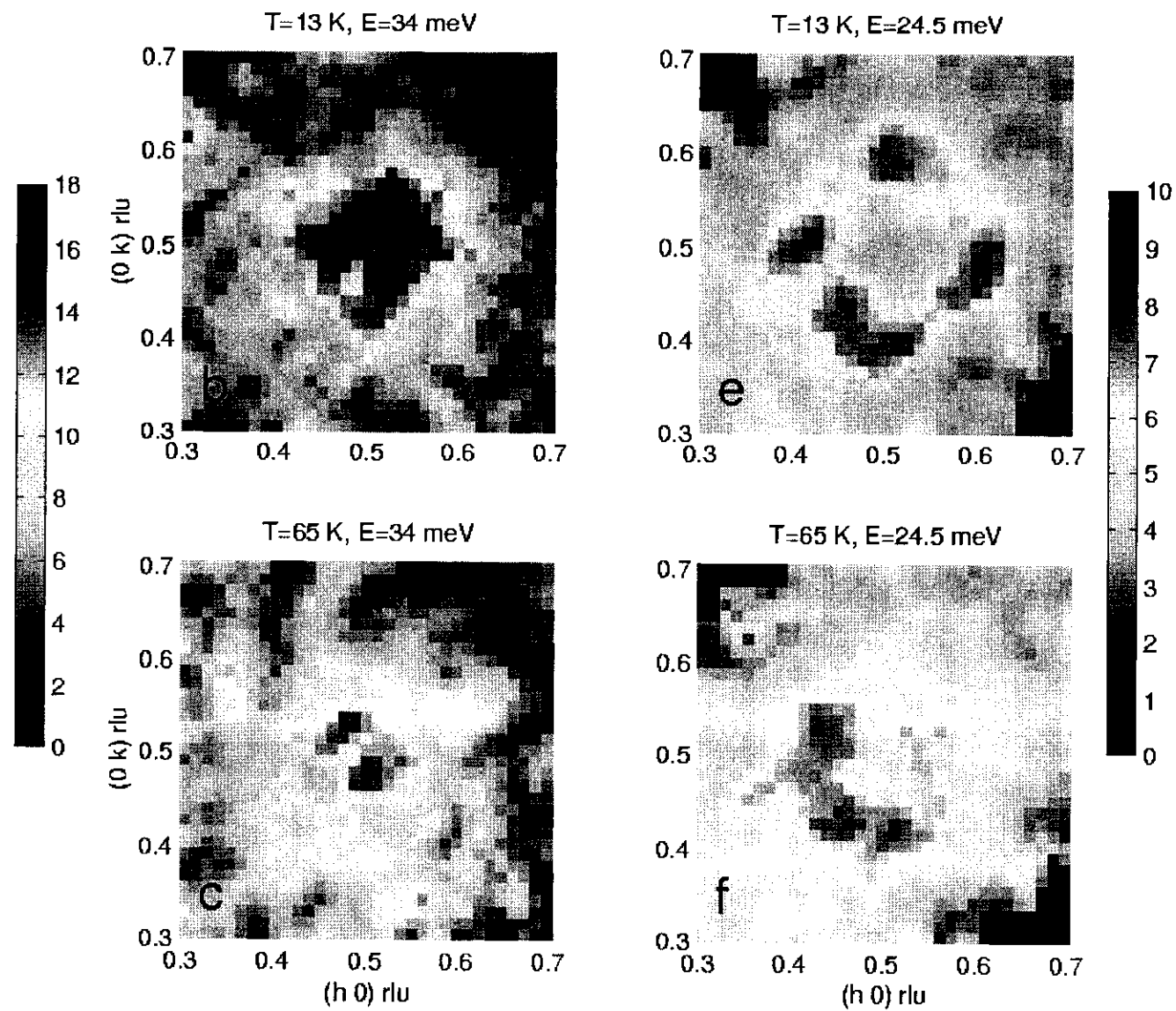
- $I \sim |\langle \psi_f | S^+ | \psi_0 \rangle|^2$
- Wavefunctions giving rise to resonance contain non-trivial orbital as well as spin terms - contrast with spin-Peierls where we also have sharp singlet-triplet excitation but only a Zeeman response to field
- one natural explanation is to invoke short range s.c. phase coherence as well as spin pairing

The ‘resonance’ as a measure for s.c. order

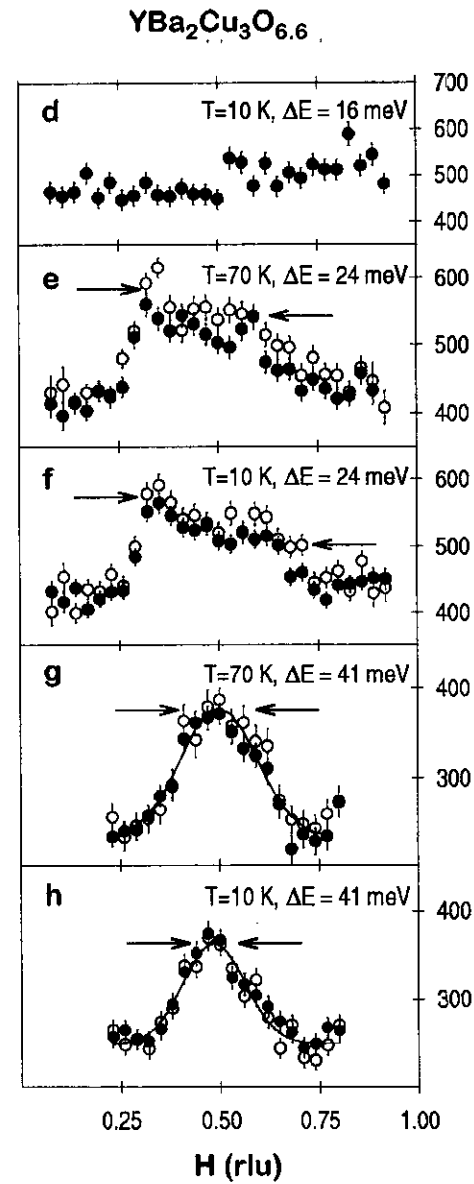
- Integrated weight $\sim$ # of pairs
- width in frequency $\sim 1/(\text{phase coherence time})$
- width in k-space $\sim 1/(\text{phase coherence length})$
- T-dependence in agreement with LBL(Orenstein)
optical data showing s.r. phase coherence showing
up for $T < T^{**} < T^*$ -----> 2-stage pre-s.c. regime
involving pseudogap formation followed by onset
of phase coherence

An oddity- seeming non-conservation of spectral weight when field is varied at fixed T

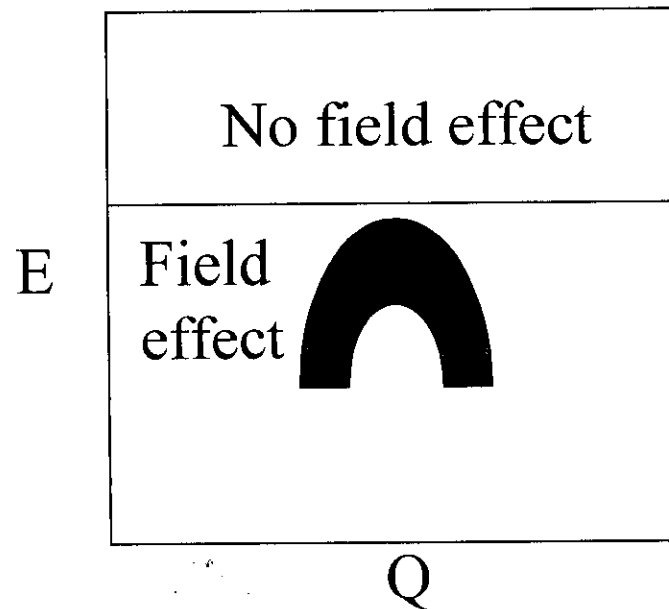
- even though at fixed B, S is conserved
- reminiscent of STM data - loss of ‘coherence peaks’ at gap edge but no obvious compensating changes
- because Q,E is a huge place, search for missing weight...



Hunting for missing spectral weight



- no luck in hunt - no increasing signals so far anywhere
- decreasing incommensurate signals for $E=24\text{meV} < E_{\text{res}}$, but no decrease in signal for $E=41\text{meV} > E_{\text{res}}$ ---> is E_{res} the top of an 'excitation band' sensitive to pairing & s.c. phase coherence



Is Heisenberg picture(perhaps with some damping and disordered ground states) all we need to understand spin sector of hi-Tc problem? -**NO**

- Revisit insulator - find ring exchange term
- look at excitations in the s.c. - discover that resonance is sensitive to orbital effects, most likely due to orbital aspect of s.c.pairing