

SMR 1232 - 6

**XII WORKSHOP ON
STRONGLY CORRELATED ELECTRON SYSTEMS**

17 - 28 July 2000

***NEUTRON RESONANT MODE
IN HIGH TEMPERATURE SUPERCONDUCTORS***

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These are preliminary lecture notes, intended only for distribution to participants.

Neutron Resonant Mode in High Temperature Superconductors

B. Keimer, MPI Stuttgart

Neutron Scattering:

H.F. Fong, Princeton University

P. Bourges, Y. Sidis, CEA Saclay

L.P. Regnault, CEA Grenoble

A. Ivanov, ILL

H. He, C. Ulrich, MPI Stuttgart

Crystal Growth:

I.A. Aksay, D.L. Milius, Princeton University

G.D. Gu, UNSW, Australia

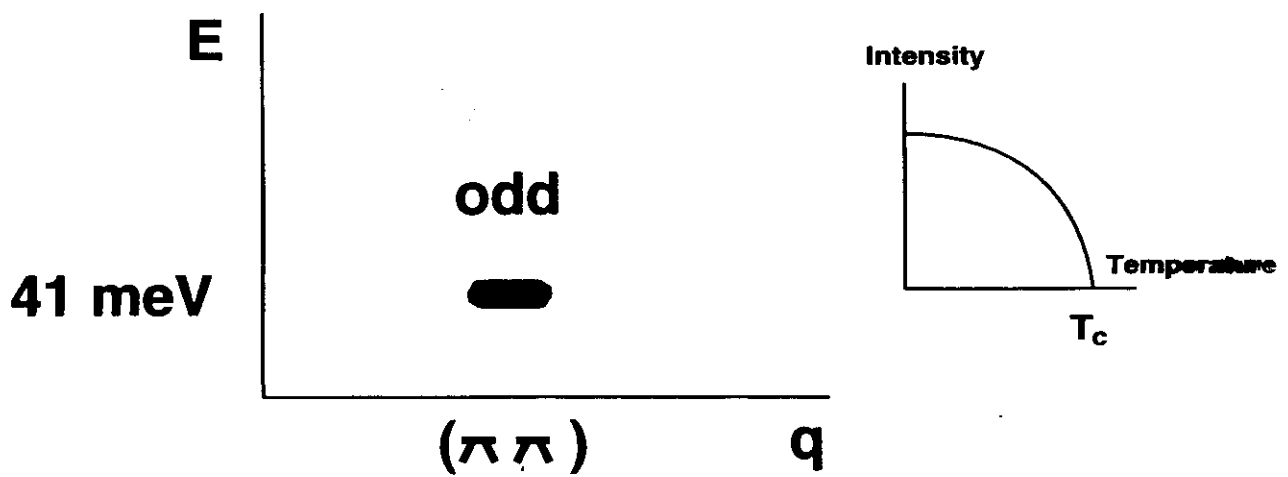
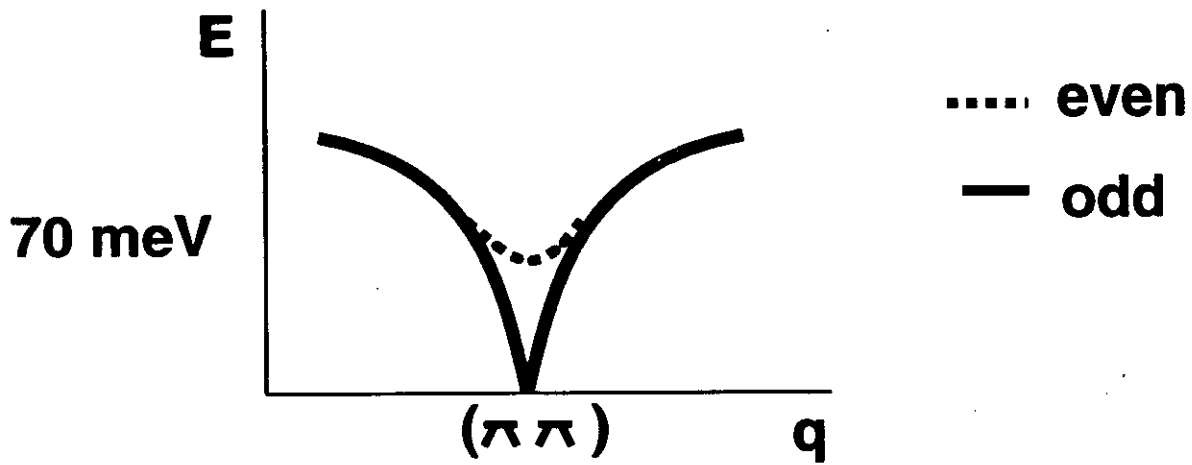
N. Koshizuka, ISTECH, Japan

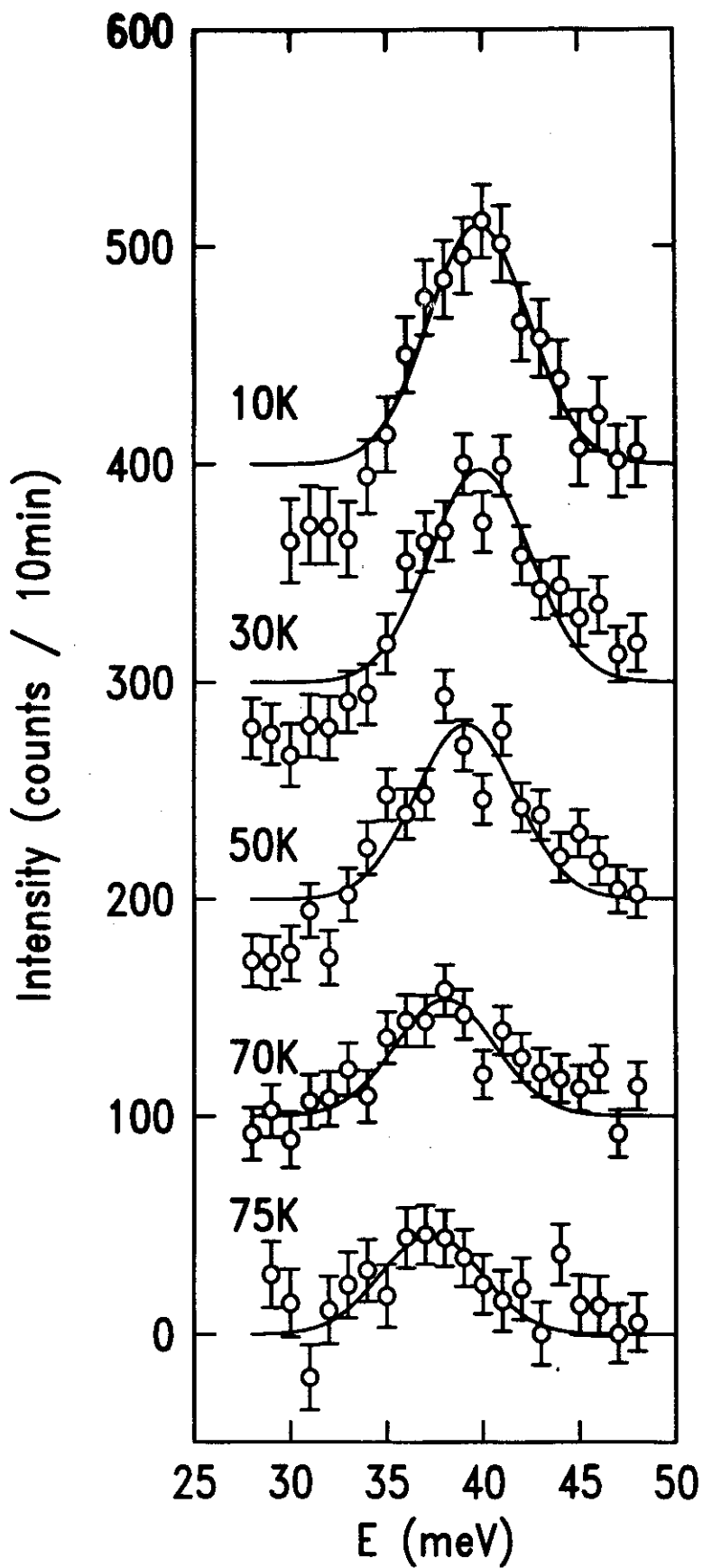
C.T. Lin, B. Liang, MPI Stuttgart

Outline:

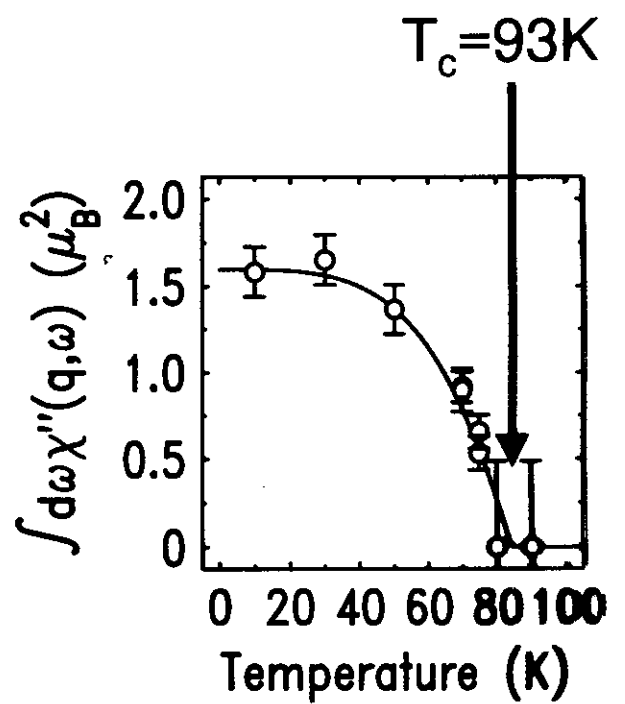
- * Introduction: Resonant Mode in optimally doped and underdoped YBCO
- * Resonant mode in optimally doped and overdoped BSCCO
- * Resonant mode and incommensurate peaks in near-optimally doped YBCO
- * Influence of $S=0$ and $S=1$ impurities on the resonant mode

Magnetic Excitations

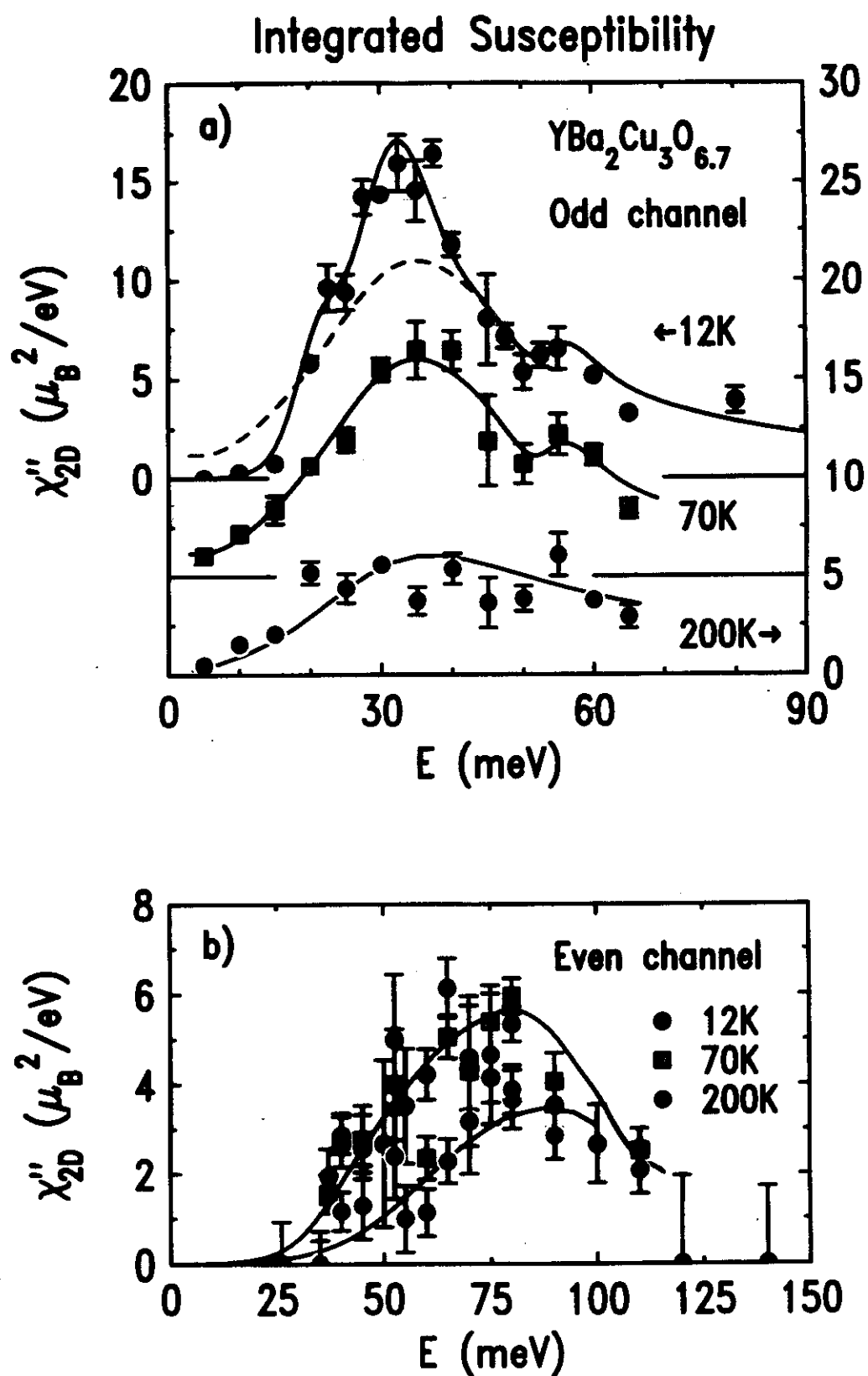




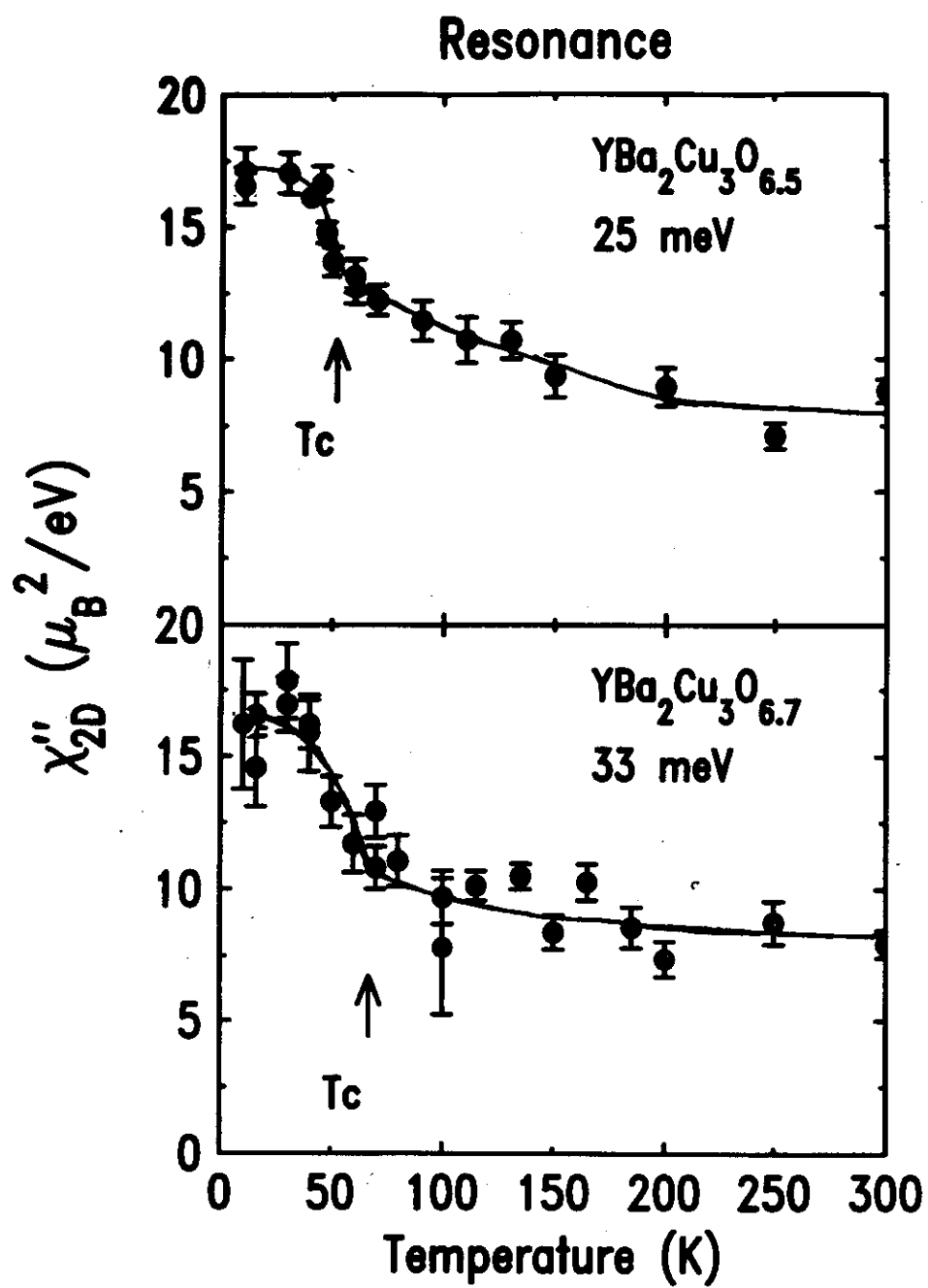
$\text{YBa}_2\text{Cu}_3\text{O}_7$

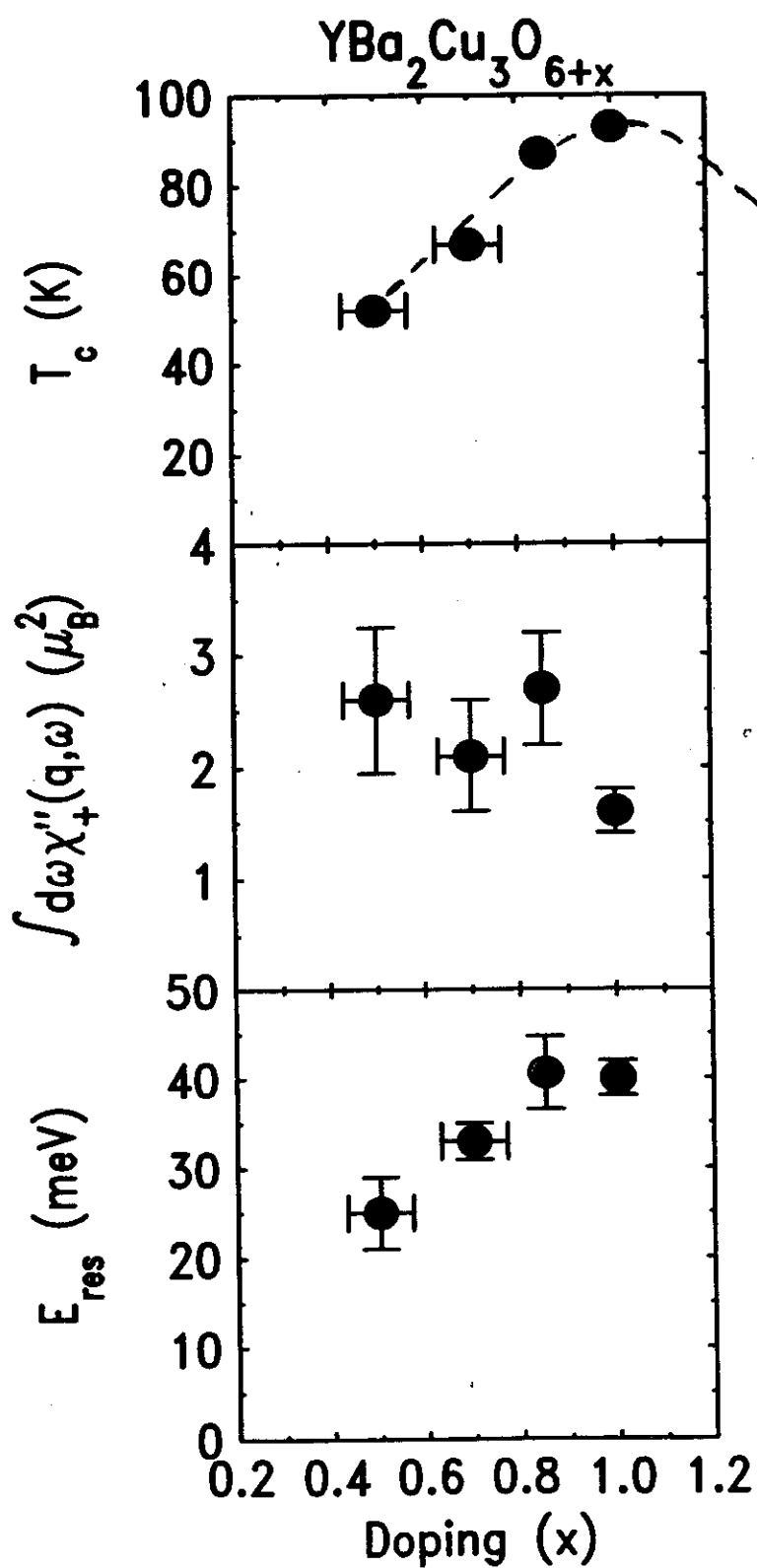


Fouq et al., PRB 1996



Fouq et al., PRB 2000





Spin Susceptibility in a BCS Superconductor

$$\chi_0(\mathbf{q}, \omega) = \sum_{\mathbf{k}} \{$$

coherence factor Fermi factor

$$\frac{1}{2} \left(1 + \frac{\epsilon_{\mathbf{k}} \epsilon_{\mathbf{k}+\mathbf{q}} + \Delta_{\mathbf{k}} \Delta_{\mathbf{k}+\mathbf{q}}}{E_{\mathbf{k}} E_{\mathbf{k}+\mathbf{q}}} \right) \frac{f(E_{\mathbf{k}+\mathbf{q}}) - f(E_{\mathbf{k}})}{\omega - (E_{\mathbf{k}+\mathbf{q}} - E_{\mathbf{k}}) + i\delta}$$

scattering of thermally excited pairs

$$+ \frac{1}{4} \left(1 - \frac{\epsilon_{\mathbf{k}} \epsilon_{\mathbf{k}+\mathbf{q}} + \Delta_{\mathbf{k}} \Delta_{\mathbf{k}+\mathbf{q}}}{E_{\mathbf{k}} E_{\mathbf{k}+\mathbf{q}}} \right) \frac{1 - f(E_{\mathbf{k}+\mathbf{q}}) - f(E_{\mathbf{k}})}{\omega + (E_{\mathbf{k}+\mathbf{q}} + E_{\mathbf{k}}) + i\delta}$$

pair annihilation

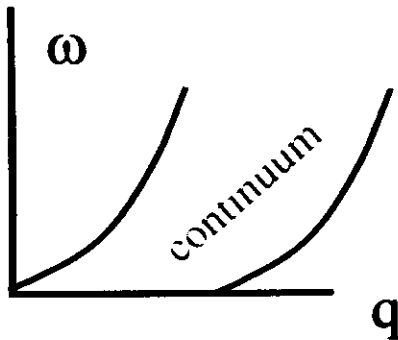
$$+ \frac{1}{4} \left(1 - \frac{\epsilon_{\mathbf{k}} \epsilon_{\mathbf{k}+\mathbf{q}} + \Delta_{\mathbf{k}} \Delta_{\mathbf{k}+\mathbf{q}}}{E_{\mathbf{k}} E_{\mathbf{k}+\mathbf{q}}} \right) \frac{f(E_{\mathbf{k}+\mathbf{q}}) + f(E_{\mathbf{k}}) - 1}{\omega - (E_{\mathbf{k}+\mathbf{q}} + E_{\mathbf{k}}) + i\delta} \}$$

pair creation

$\epsilon_{\mathbf{k}}$ band dispersion

$$E_{\mathbf{k}} = \sqrt{\epsilon_{\mathbf{k}}^2 + \Delta_{\mathbf{k}}^2}$$

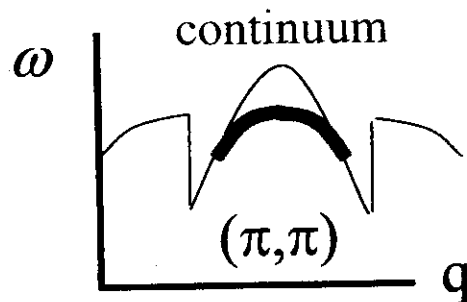
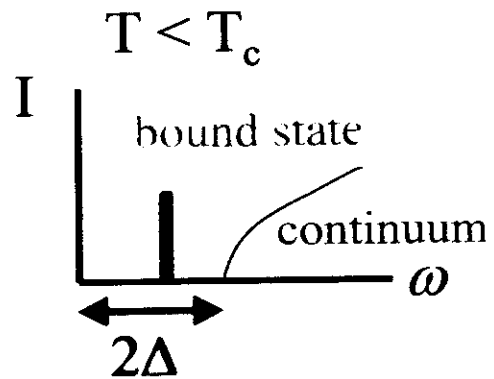
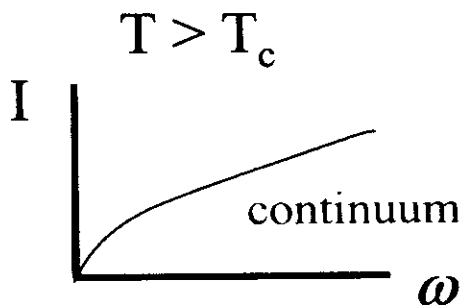
Strong-Correlation Model

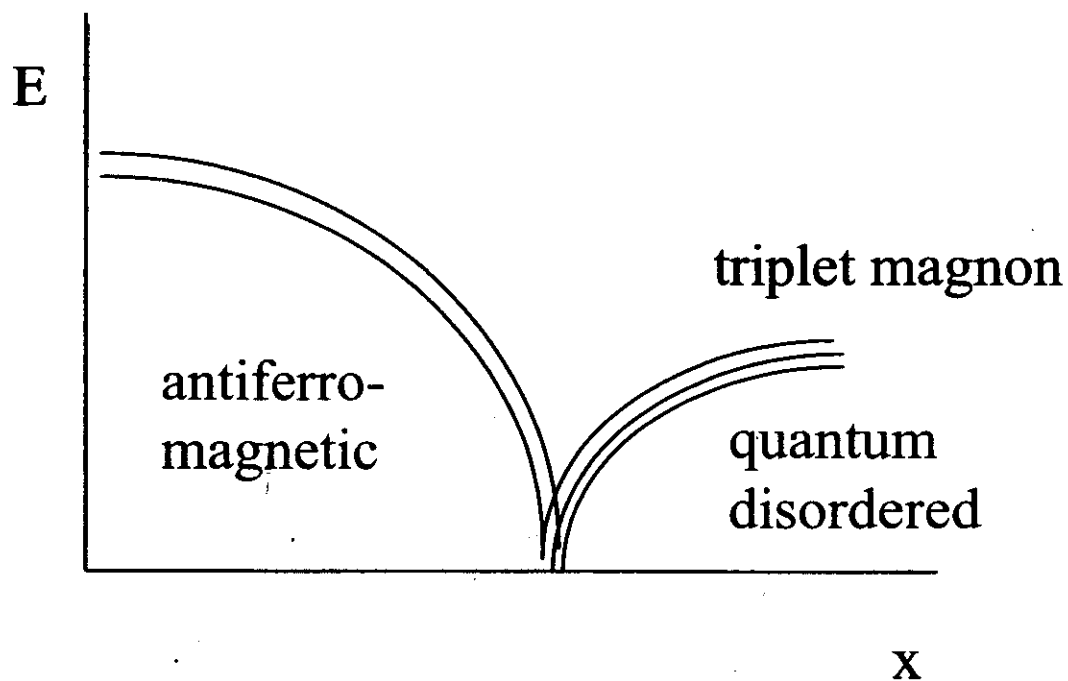


Starting point:
band susceptibility $\chi_0(q, \omega)$

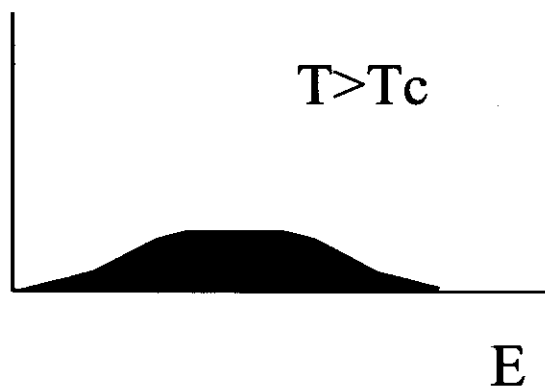
Renormalization of $\chi_0(q, \omega)$ by interactions:

Stoner model:
$$\chi(q, \omega) = \frac{\chi_0(q, \omega)}{1 - J(q) \chi_0(q, \omega)}$$



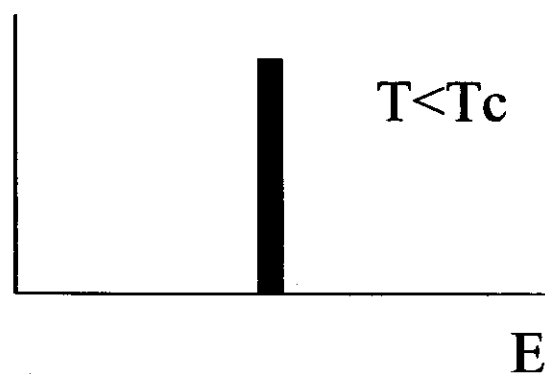


Intensity



strong quasiparticle
damping

Intensity



no quasiparticle
damping

SO(5) Model

Resonance peak is pseudo-Goldstone mode of
symmetry group linking
antiferromagnetism (3-comp. order parameter)
and d-wave superconductivity (2 comp. order parameter)

symmetry exact at quantum critical point ($E_{\text{res}}=0$)

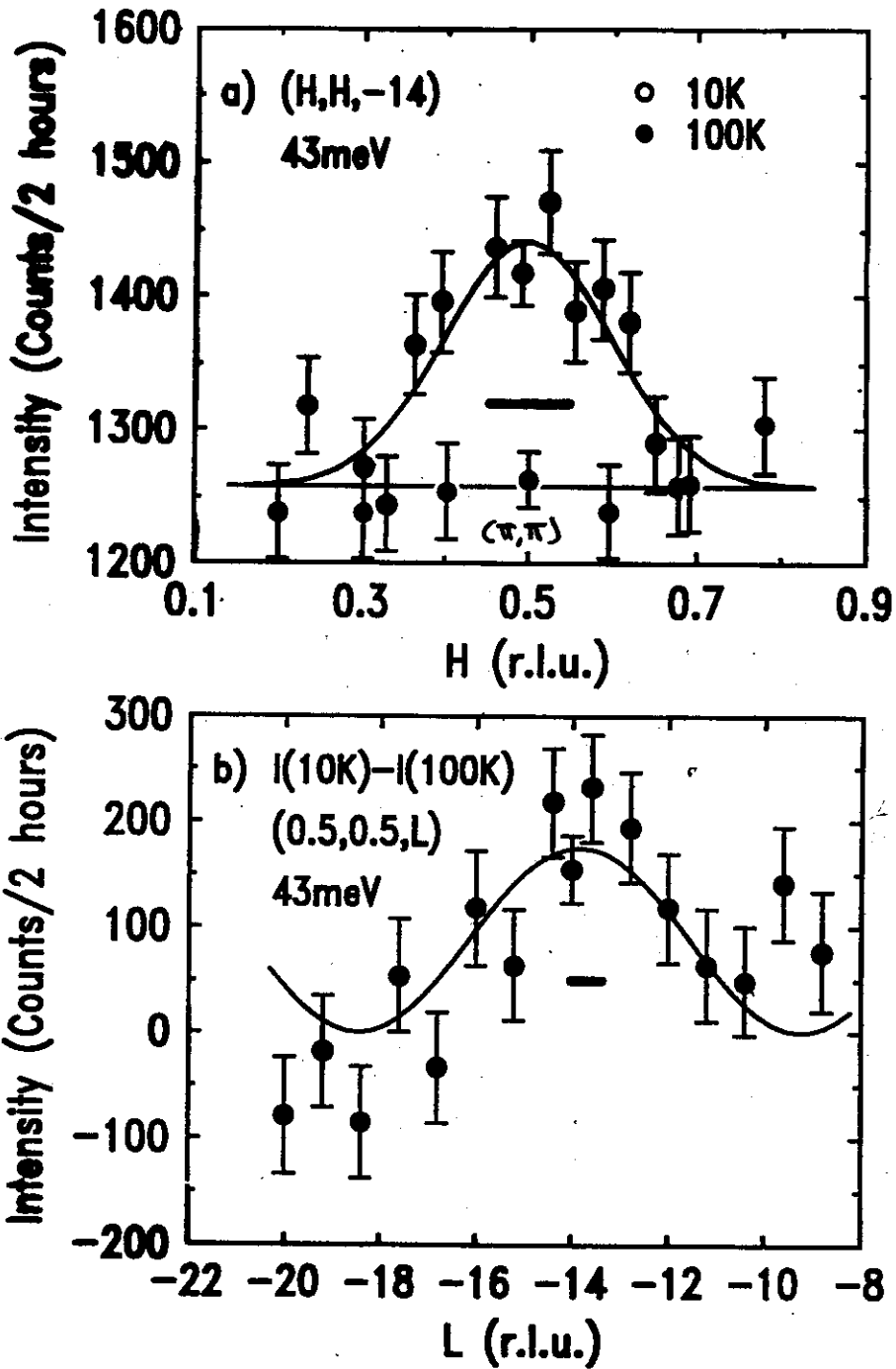
in superconducting phase:

$$I_{\text{res}} \propto \frac{\Delta^2}{n J^2}$$

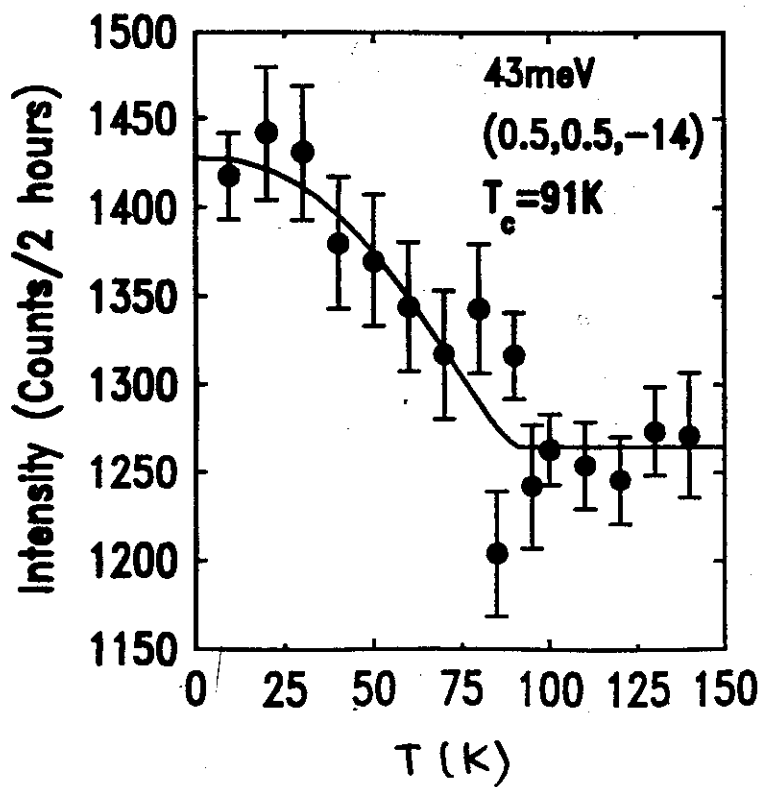
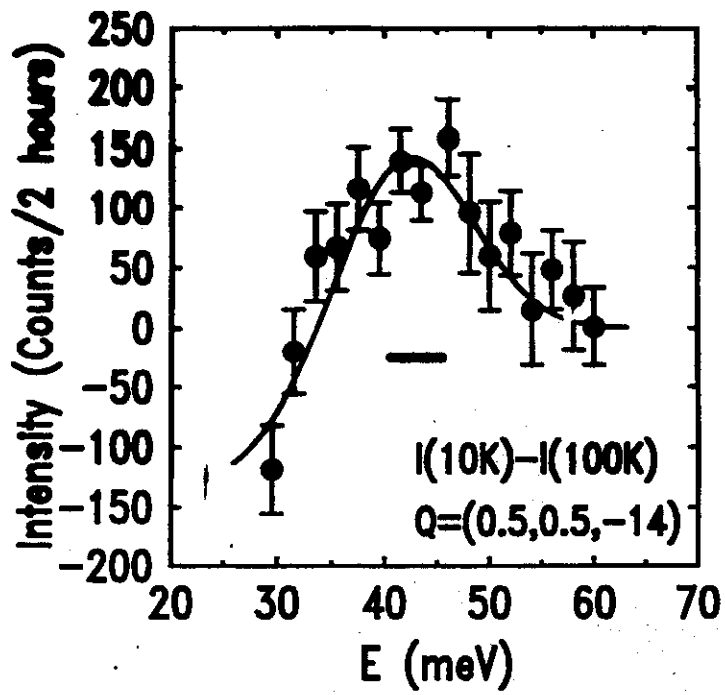
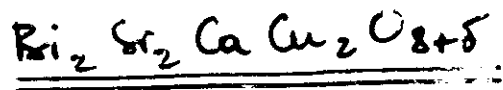
$$E_{\text{res}} \propto n J$$

Bi_2Se_3 Ca Cu₂ C_{8+δ} ($T_c = 92\text{K}$)

0.08 cm² single crystal

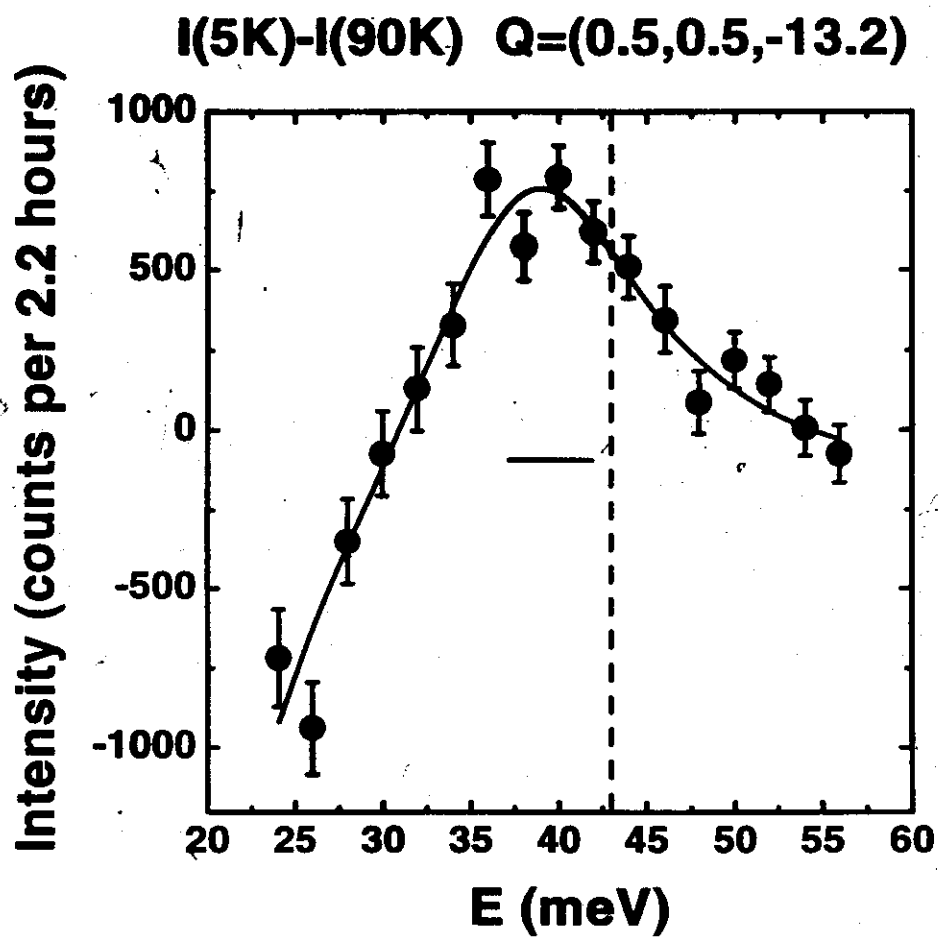


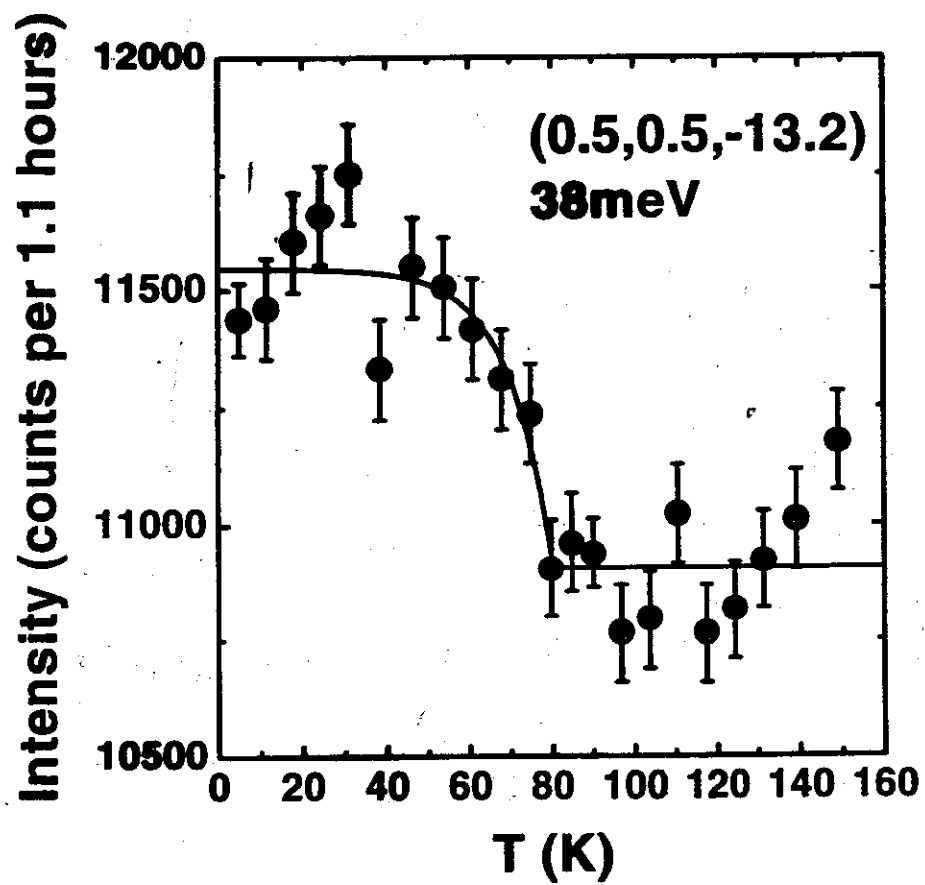
Fong et al., Nature

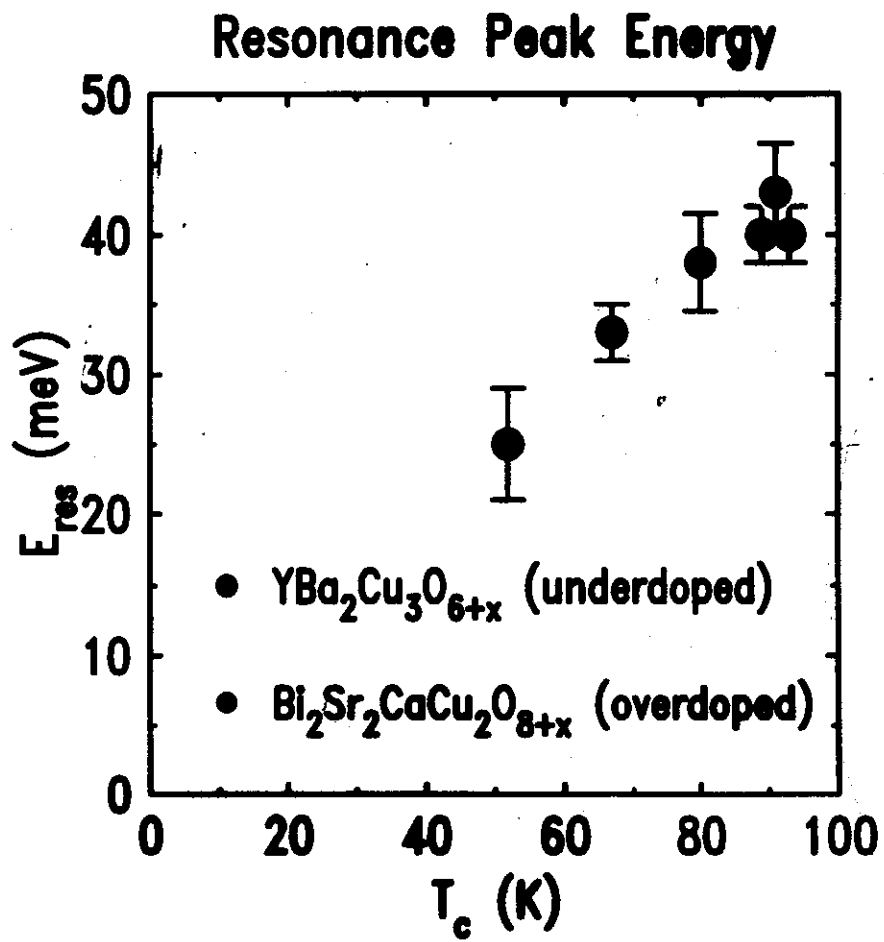


Fong et al.
 Nature (99)

Overdoped $\text{Bi}_{1-x}\text{Ca}_x\text{Cu}_2\text{O}_{8+\delta}$ ($T_c = 85\text{ K}$)







$$\frac{E_{res}}{kT_c} \sim 5.5$$

$\text{YBa}_2\text{Cu}_3\text{O}_{6.85}$ ($T_c = 87\text{K}$)

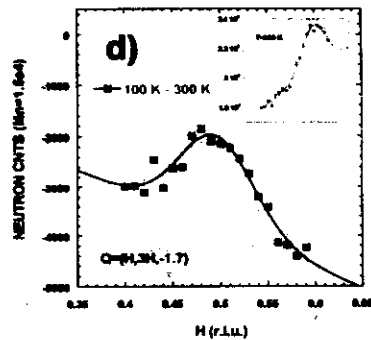
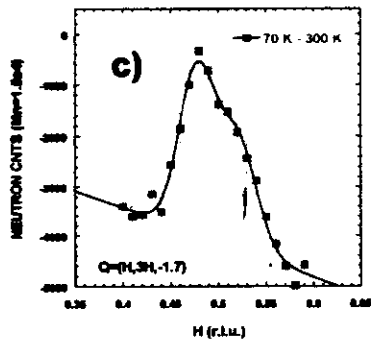
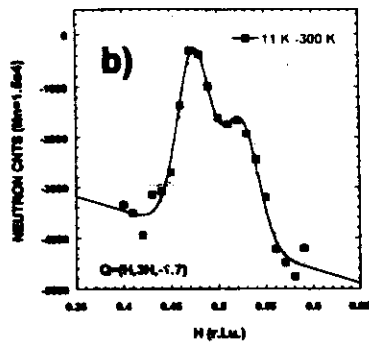
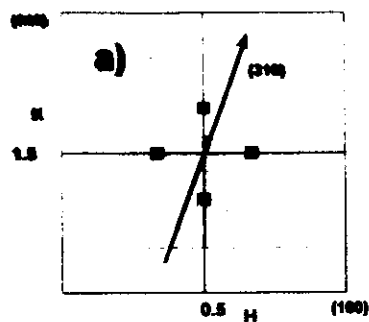
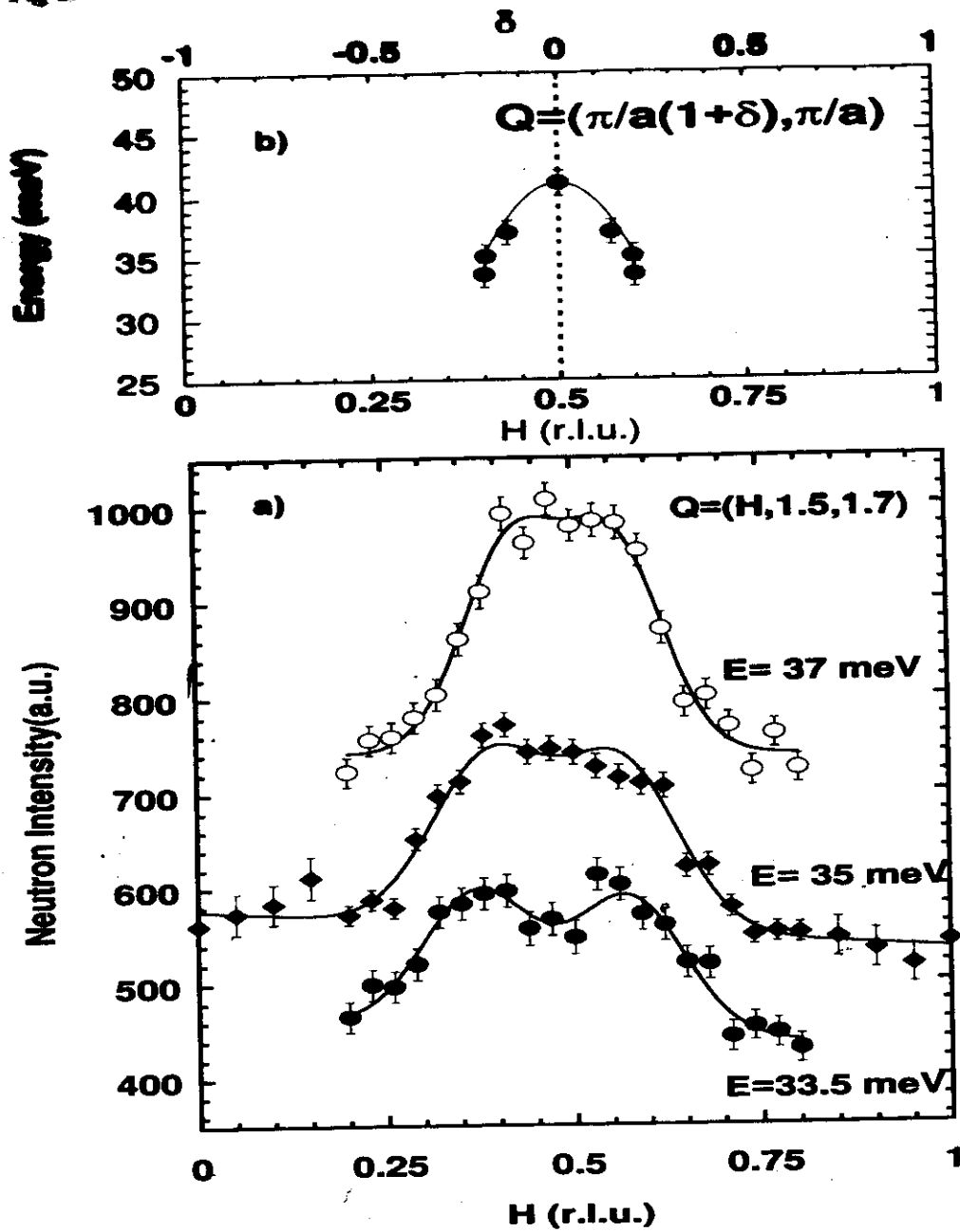
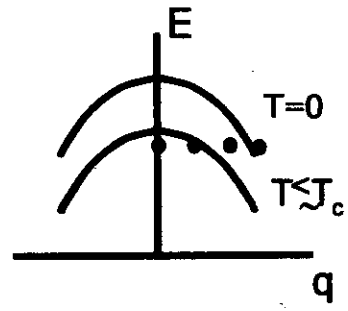


Fig. 2



near-optimally doped $\text{YBa}_2\text{Cu}_3\text{O}_{6.85}$ ($T_c = 87\text{ K}$)



$\text{YBa}_2\text{Cu}_3\text{O}_{6.85}$

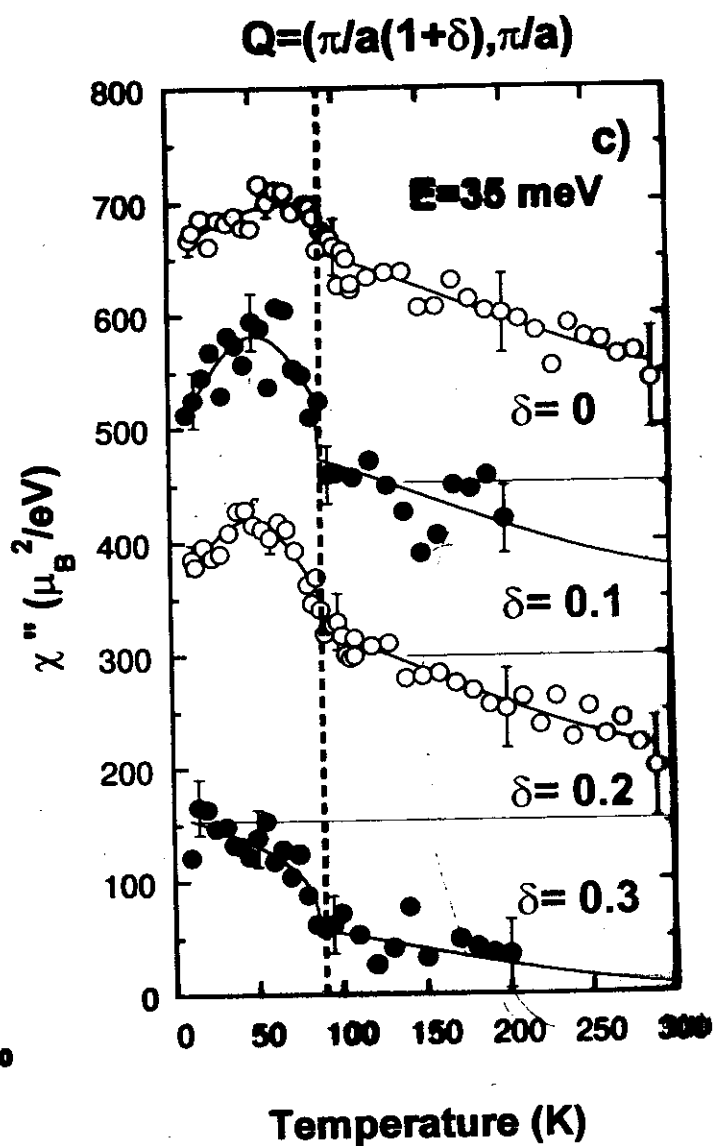
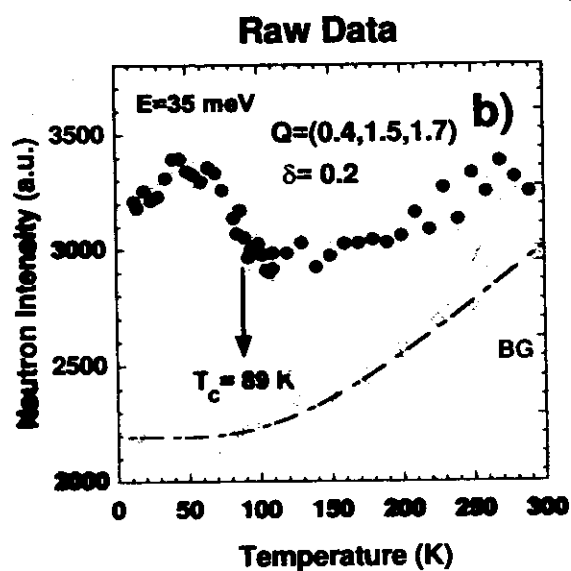
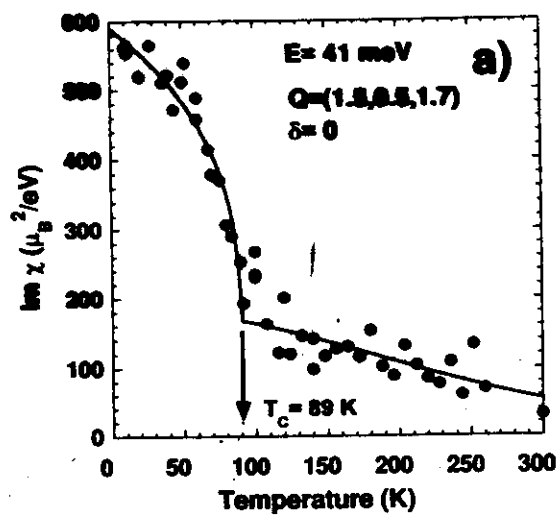
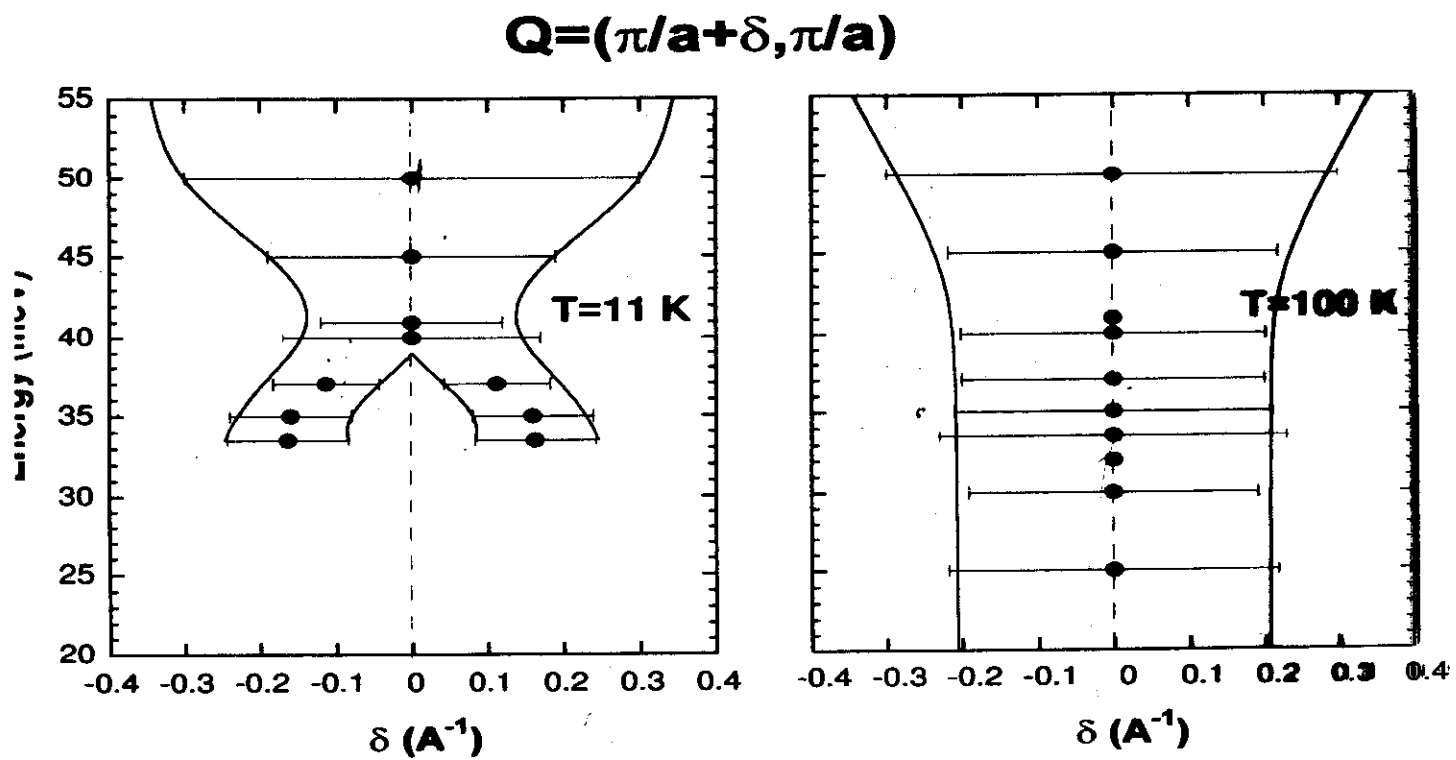


Fig 3:



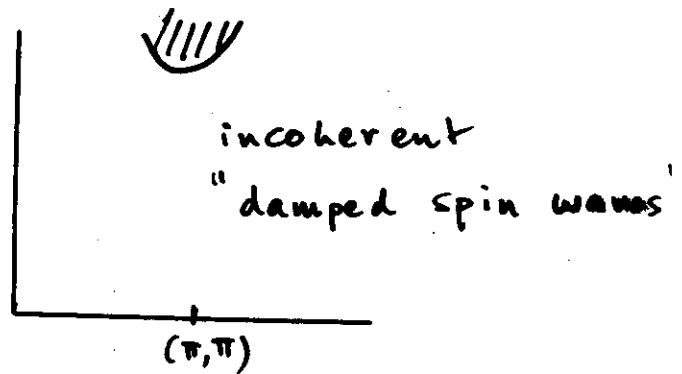
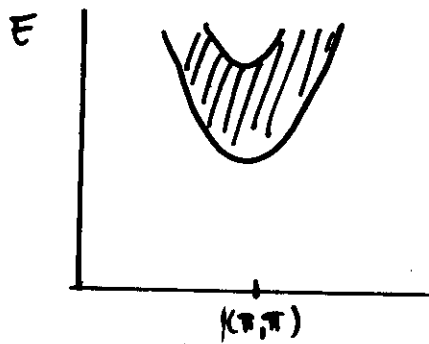
Bourges et al., Science (1990)

One way to look at the spectra:

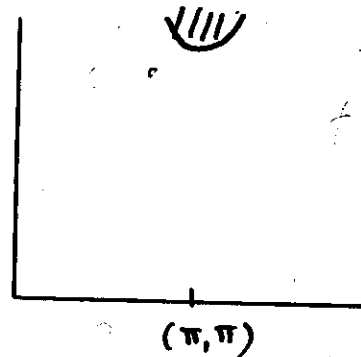
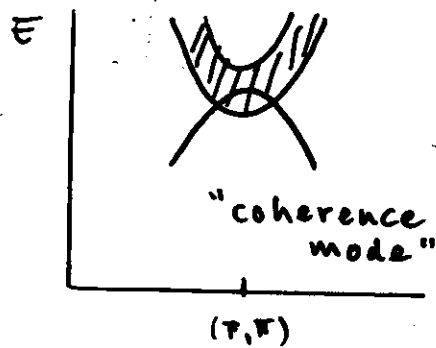
$T > T_c$

odd

even

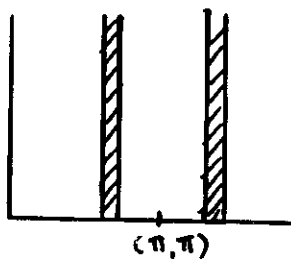


$T < T_c$

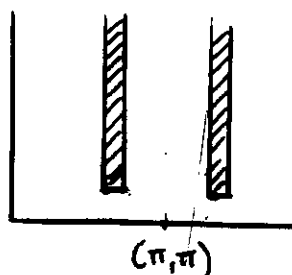


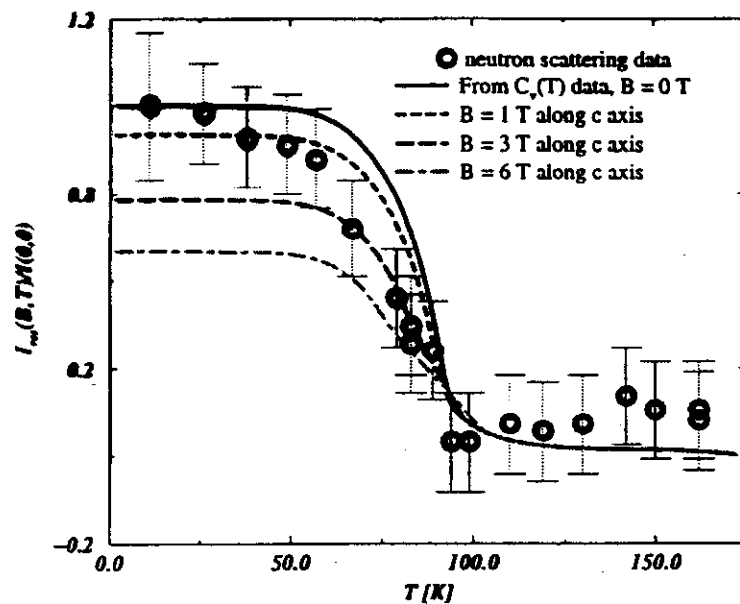
compare LSCO:

$T > T_c$

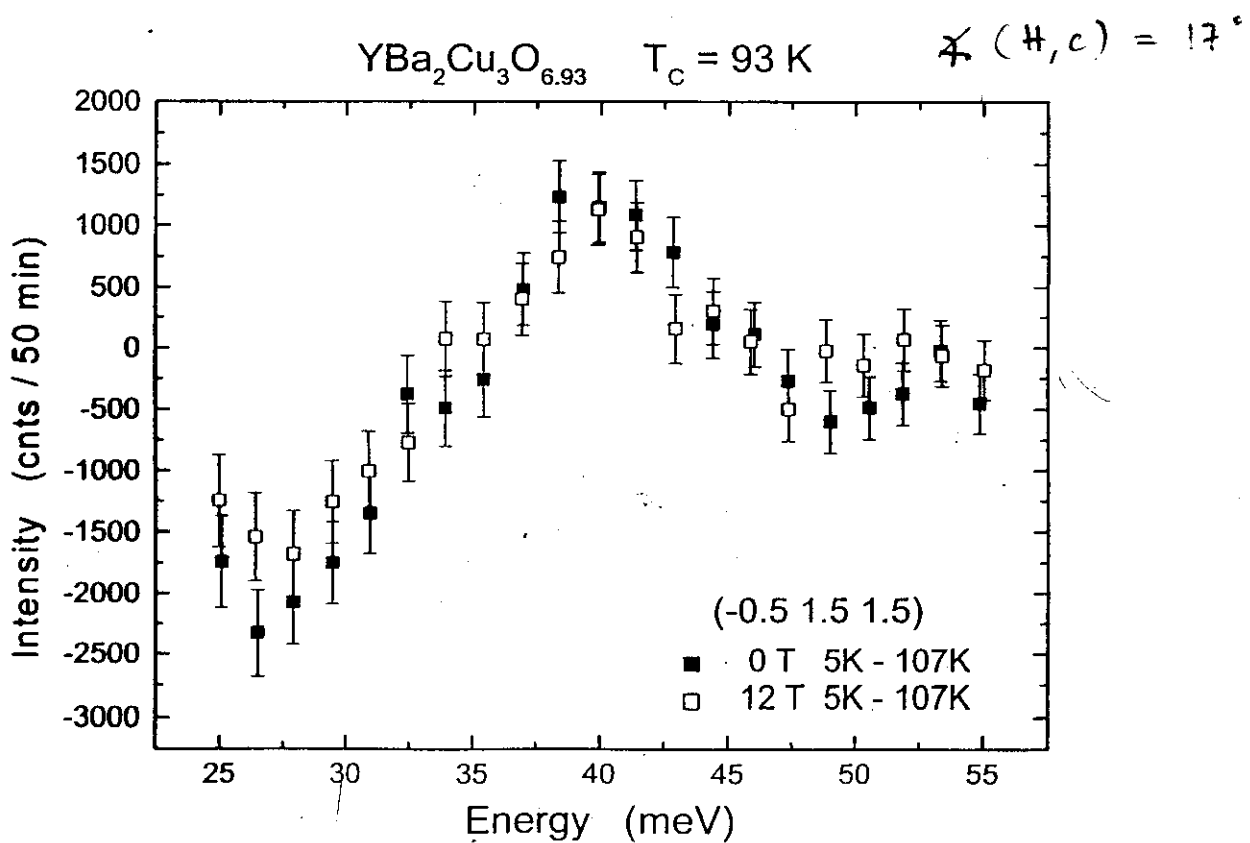


$T < T_c$





B. Janko, cond-mat/1912.0000



Impurity Effects

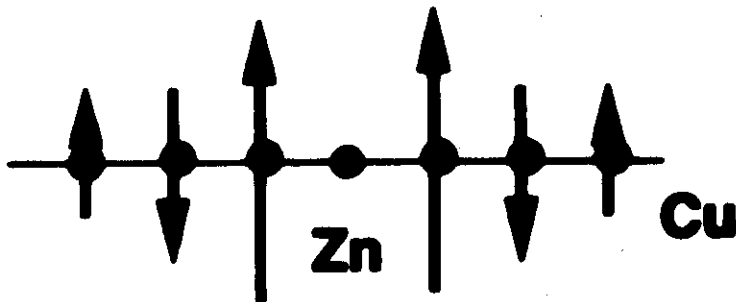
ordinary superconductors:

nonmagnetic impurities: weak effect

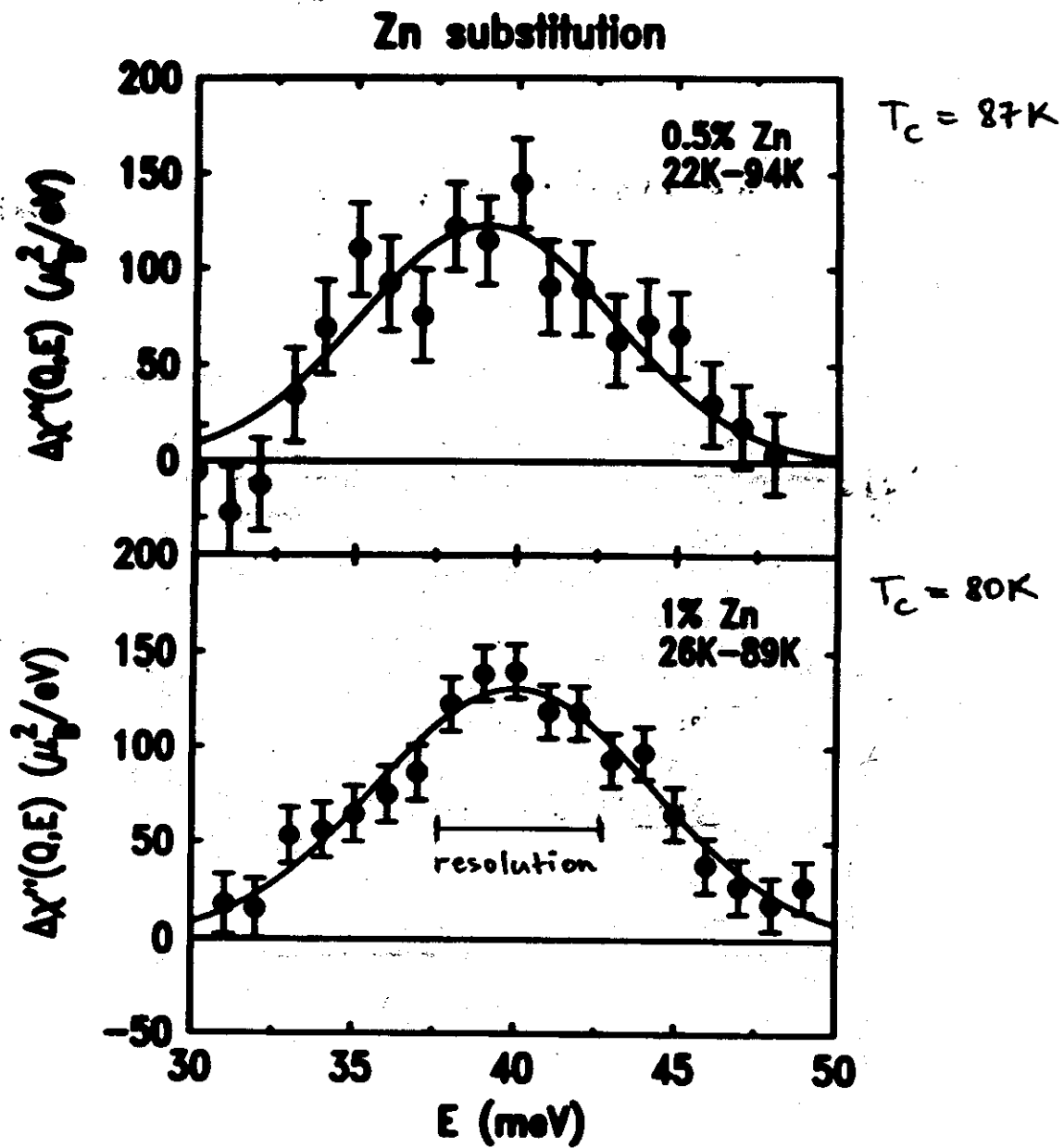
magnetic impurities: strong effect

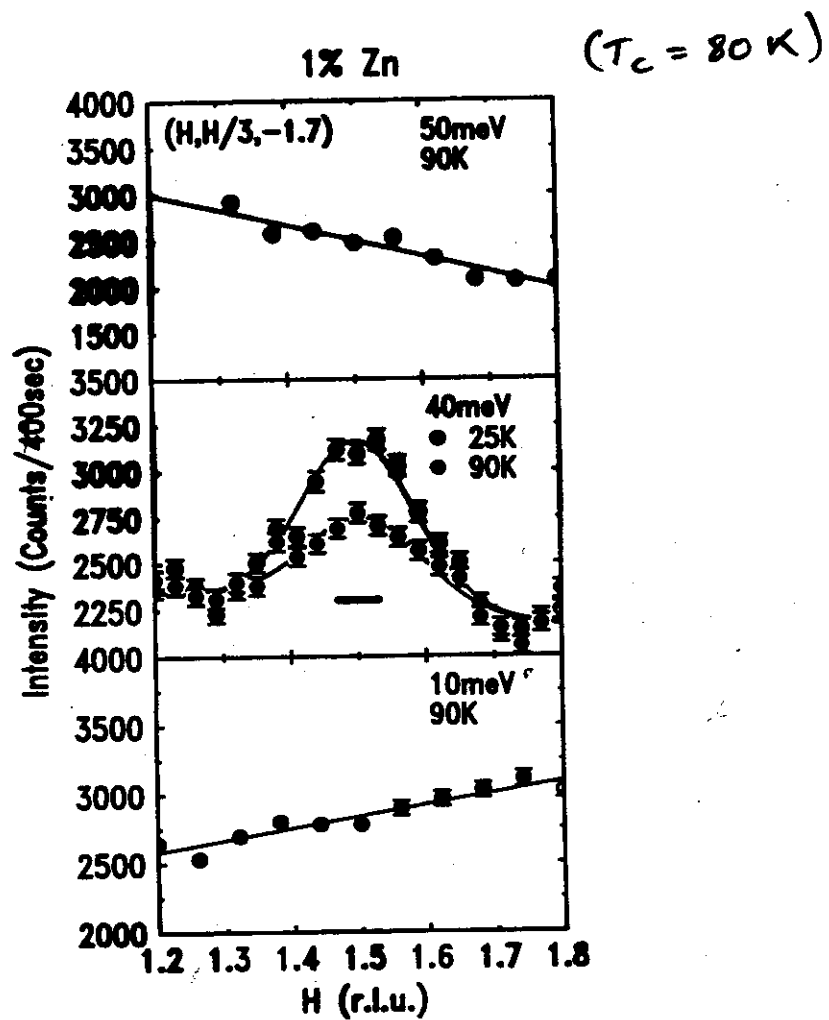
Zn in cuprates:

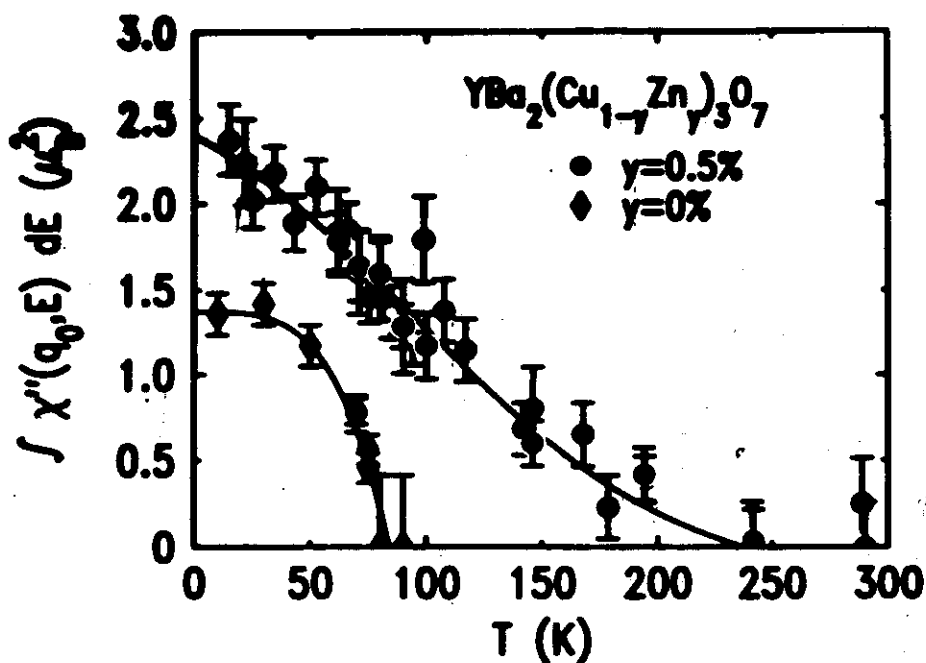
- $3d^{10}$: "nonmagnetic", but
- *transport* experiments indicate efficient scattering of conduction electrons
- *magnetization* and *NMR* experiments show local magnetization



- *neutron* experiments (at high Zn concentrations) shows low energy magnetic fluctuations







observable spectral weight at low temperature (μ_B^2)

YBCO7

YBCO7 : 0.5% Zn

$\int dE \chi''(\pi, \pi, E)$

1.6 ± 0.5

2.2 ± 0.5

$\int d^2q dE \chi''(\vec{q}, E)$

0.048

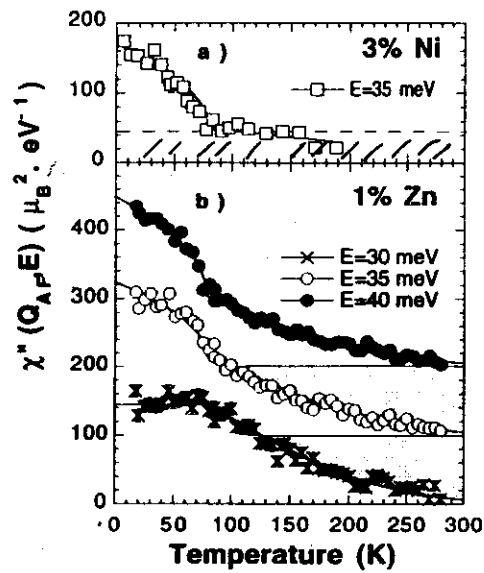
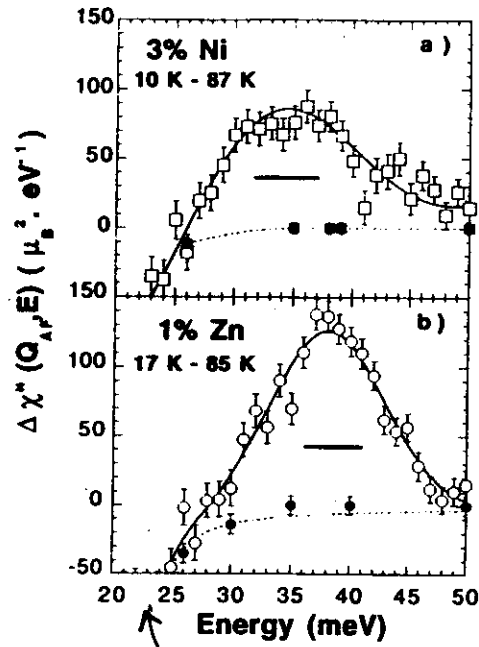
0.058

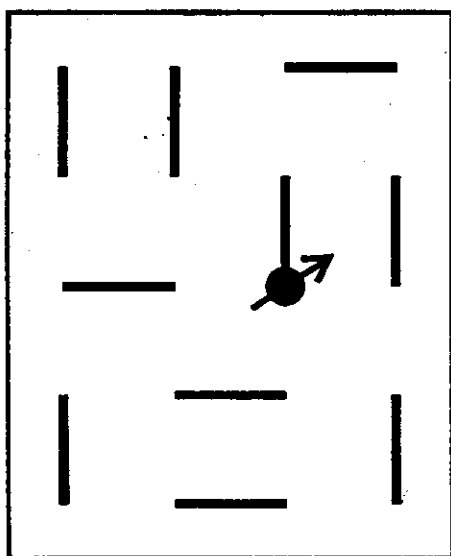
removed completely
(within detection limit)
above T_c

half removed
above T_c

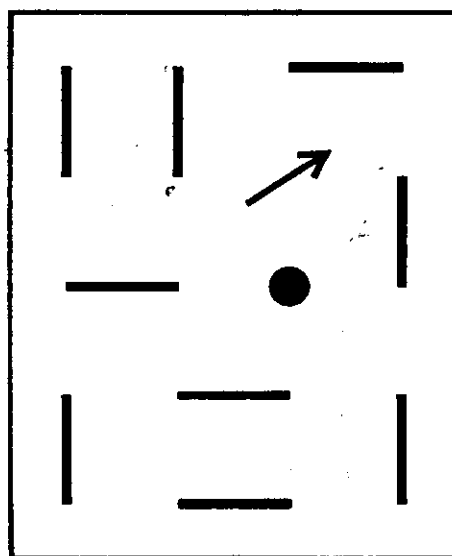
Fong et al., PRL (99)

$$\frac{E_{res}}{kT_c} \sim 5.5$$





$S=1$ impurity



$S=0$ impurity

Kilian et al., cond-mat 9904315

Conclusions:

- * Resonant mode energy determined by T_c , not doping
- * Incommensurate peaks disappear above T_c , follow same temperature dependence as resonant mode
- * $S=0$ and $S=1$ impurities affect resonant mode strongly and *differently*